

### §13. Island Dynamics in LHD Plasmas (I)

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We have found that a significant enhancement of the externally imposed island with  $n/m=1/1$  occurs in the LHD discharges with collisional, finite beta plasma. For the collisionless plasmas, on the other hand, suppression of the island has been observed.

Figure 1(a) shows that when the density is ramped up slowly, a transition takes place from a normal state to equilibrium with a large island. The stored energy ( $W_p$ ) initially increases and nearly saturates and at  $t = 1.6$  s, the transition takes place. The electron temperature profiles before and after the transition are shown in Fig. 1(b). After the transition, local flattening of the temperature profile at  $\nu/2\pi = 1$  ( $R = 4.20$  m) is seen and leads to a nearly uniform drop ( $\sim 150$  eV) in the temperature in the core, resulting in a significant reduction in  $W_p$ . The flattening region of  $\sim 120$  mm in  $R$ , corresponding to  $w = 0.12$  is much wider than that in the vacuum configuration (40 mm in  $R$ ). The edge profile beyond the island ( $R > 4.25$  m), however, remains almost unchanged. From the transition, the amplitude of the magnetic fluctuation increases, indicating that some MHD activities are involved.

In Fig. 2, the stored energy ( $W_p$ ) is plotted as a function of density for various  $w_{ex}$ . For  $w_{ex} = +0.045$ , no clear island is detected at any density. The  $W_p$  scales as  $(n)^{1/2}$ , but saturates and drops at high density, very typical density dependence in LHD. For  $w_{ex} = -0.07$  (the minus sign means that the spatial phase of the island is different by  $\pi$  compared with that with plus sign), the island starts to appear when the density exceeds a threshold value and at the same time  $W_p$  starts to become lower than that with  $w_{ex} = +0.045$ . This threshold value decreases with increasing  $|w_{ex}|$ . For a large vacuum island ( $w_{ex} = -0.11$ ), however, an island is seen at any density, as expected. For high-density regime,  $W_p$  and hence  $\tau_E$  depend on  $w_{ex}$  strongly. When  $n$  is below  $1.5 \times 10^{19} \text{ m}^{-3}$ , the stored plasma energy ( $W_p$ ) is insensitive to  $w_{ex}$ .

Data in Fig. 1, 2 clearly shows that under certain plasma conditions (i.e., higher density and larger vacuum island), plasma can generate a large island, larger than that in the vacuum configuration, which in turn deteriorate the energy confinement severely. Such plasma effects are the strongest for the most inward shifted

configuration ( $R_{ax} = 3.5$  m). For  $R_{ax} = 3.6$  m, this effect is much weaker and the threshold density is much higher.

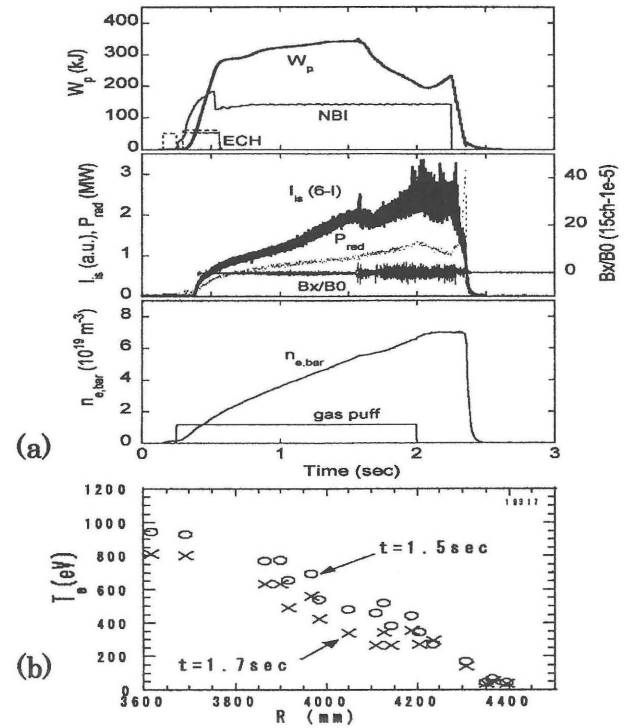


FIG. 1. (a) Time evolution of the stored energy in an LHD discharge ( $R_{ax}=3.5$ m,  $B=2.8$ T) shows a transition from a normal state to an equilibrium with a large island. (b) The electron temperature profiles before ( $t=1.5$  sec) and after ( $t=1.7$  sec) the transition.

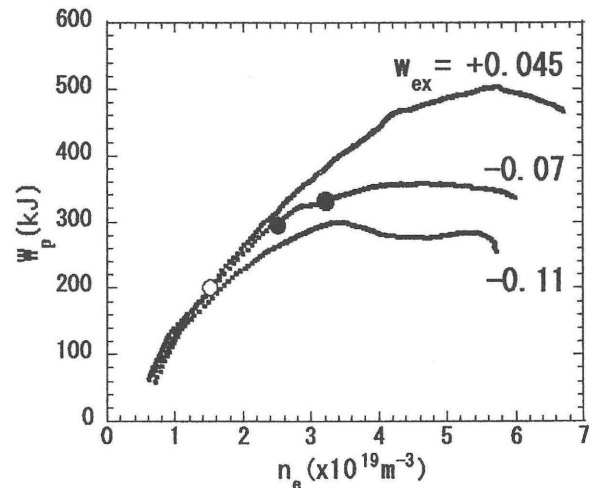


FIG. 2. The stored plasma energy is plotted as a function of the density for various vacuum island sizes ( $B = 2.83$  T,  $R_{ax} = 3.5$  m). For  $w_{ex} = -0.07$  case, island appears at higher density, (o : no island), (• : island)