

## §9. Inward Turbulent Particle Transport in NBI-Heated Plasmas

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Edge plasma turbulence study are carried out using a Langmuir probe array in the CHS heliotron/torsatron. The Langmuir probe array consists of 4 sets of a triple probe, which are separated radially[1]. Each triple probe set has 4 molybdenum electrodes. Radial profiles of time-averaged plasma parameters such as electron temperature  $T_e$ , electron density  $n_e$  and plasma potential  $V_s$  are measured by 4 set probes for one discharge. Floating potentials are measured by two electrodes in each probe set, which are separated by 4.5 mm in the poloidal direction. Fluctuations of poloidal electric field  $E_\theta$  are derived from these floating potentials. Here, fluctuations of electron temperature  $T_e$  are ignored. The turbulent particle flux is expressed as

$$\Gamma_{turb} = (2/B_t) \int_0^\infty \gamma_{n_e E_\theta} \cos \alpha_{n_e E_\theta} [P_{n_e} P_{E_\theta}]^{1/2} df,$$

where  $B_t$ ,  $\gamma$ ,  $\alpha$ ,  $P$  and  $f$  are the toroidal magnetic field, coherence, phase, power spectrum density, and frequency, respectively. Therefore, the flux is reduced by reduction of fluctuation amplitude, decorrelation of  $n_e$  and  $E_\theta$  or change of relative phase between them, but inward flux is only realized by the change in the relative phase between both fluctuations.

In some shots, the averaged inward flux is seen around  $r/\langle a \rangle \sim 0.94$  during the NBI heating phase. We investigated the parameter space where the inward flux is realized, changing fuel gas,  $B_t$ , NBI heating power  $P_{NBI}$  and  $n_e$  near the edge, in the inward-shifted configuration of  $R_{ax} \sim 0.92$  m. As seen from Fig.1, the inward flux is observed in the space of  $T_e > 15$  eV and  $n_e < 3 \times 10^{18} \text{ m}^{-3}$  at  $r/\langle a \rangle \sim 0.94$ . The inward particle flux appeared in several fuel gases (H and He) and at  $B_t = 0.9 - 1.4$  T. This parameter space approximately corresponds to the low collisionality regime (the effective collision frequency normalized by the transit frequency of a circulating particle  $\nu^* < 1$ ) less than lower bound of Pfirsch-Schlüter regime, as seen from Fig.1.

The radial electric field  $E_r$ , its shear  $E_r'$  and curvature  $E_r''$  are thought to be the possible candidates for the reversal of fluctuation-induced particle flow. The inward flux obviously correlates with  $E_r'$ , that is, the large positive  $E_r'$  is generated near the relevant region where the inward flux takes place. Moreover,  $T_e$  profile in the low density discharge clearly has a plateau or slightly hollow structure, of which location approximately corresponds to that of the  $1/q = 1$  rational surface  $r/\langle a \rangle|_{1/q=1} = 0.94 - 0.97$ . The sign of the phase  $\cos(\alpha_{n_e E_\theta})$  between  $n_e$  and  $E_\theta$  in an inward flow discharge is in  $\cos(\alpha_{n_e E_\theta}) \sim -1$  up to 50 kHz, in contrast to that in an outward flow discharge ( $\cos(\alpha_{n_e E_\theta}) > 0$ ), while the coherence  $\gamma_{n_e E_\theta}$  remains unchanged in both discharges. The correlation between the

flow reversal and  $E_r$ ,  $E_r'$  and  $E_r''$  is further investigated for many shots in CHS. The dependence of the turbulent particle flux  $\Gamma_{turb}$  on  $E_r'$  is shown in Fig.2, where  $\Gamma_{turb}$  and  $E_r'$  are evaluated at  $r/\langle a \rangle \sim 0.94$ . From Fig.2, the critical value that induces the large inward turbulent flux ( $> 1 \times 10^{20} \text{ m}^{-2} \text{ s}^{-1}$ ) seems to be  $E_{rc}' \sim 1 \times 10^6 \text{ Vm}^{-2}$ , although the inward flux is observed even in the range of  $0 < E_r' < E_{rc}'$ . This threshold in  $E_r'$  is about twice of that in the edge ECH experiment[2].

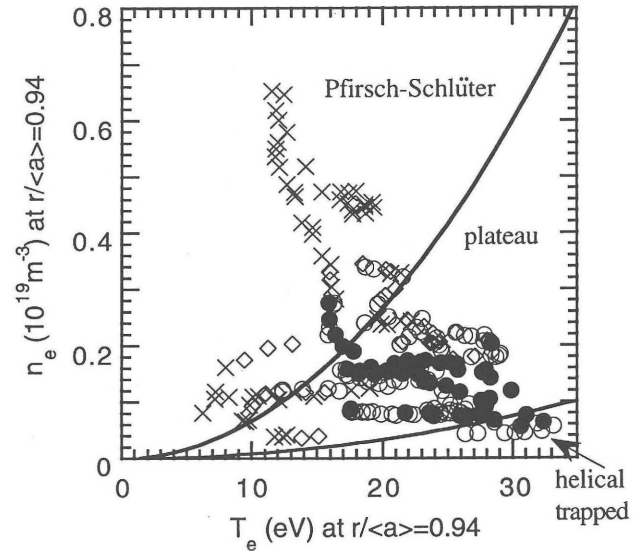


Fig. 1 Data points of  $\Gamma_{turb}$  at  $r/\langle a \rangle \sim 0.94$  on the parameter space defined by  $T_e$  and  $n_e$  at  $r/\langle a \rangle \sim 0.94$ . The solid circles, open circles, open squares and crosses respectively indicate the data with  $\Gamma_{turb} < -3 \times 10^{19}$ ,  $-1 \times 10^{19} > \Gamma_{turb} > -3 \times 10^{19}$ ,  $0 > \Gamma_{turb} > -1 \times 10^{19}$  and  $\Gamma_{turb} > 0 \text{ m}^{-2} \text{ s}^{-1}$ . Inward particle fluxes are clearly observed in the range of  $T_e > 15$  eV and  $n_e < 3 \times 10^{18} \text{ m}^{-3}$ .

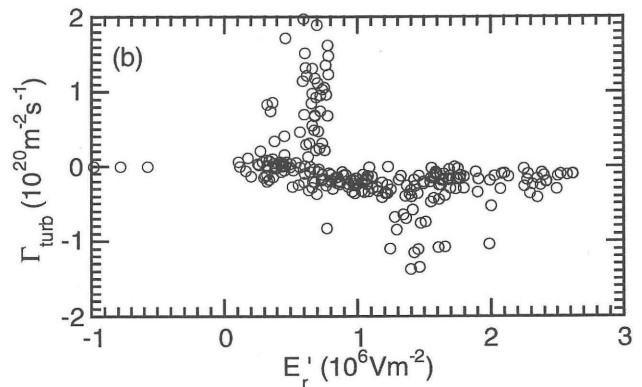


Fig. 2 Dependence of  $\Gamma_{turb}$  at  $r/\langle a \rangle \sim 0.94$  on the radial electric field shear  $E_r'$  at  $r/\langle a \rangle \sim 0.94$ . When  $E_r'$  is more than  $1 \times 10^6 \text{ Vm}^{-2}$ , large inward  $\Gamma_{turb}$  is observed.

### Reference

- [1] K. Ohkuni, *et al.*, Rev. Sci. Instrum. **72**, 446 (2001)
- [2] M.G.Shats, K.Toi, K. Ohkuni, *et al.*, Phys. Rev. Lett. **84**, 6042 (2000)