

## §16. Populating Mechanism in Dense Plasma for Doubly Excited States

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We study effective rate coefficients using a collisional-radiative model (CRM) including doubly excited states. The effective rate coefficients in dense plasma were studied by Fujimoto and Kato [1]. They examined 1s-2s excitation / de-excitation using a CRM including doubly excited states  $2snl$  as well as 1s-2p excitation / de-excitation including  $2pnl$  states. They found an enhancement of the excitation and de-excitation rate coefficients with increasing electron density. We also found an enhancement due to different mechanism compared with their found mechanism [2]. In this paper, we present the populating mechanism of doubly excited states  $2l'nl$  for effective 1s-2l excitation rate coefficients in dense plasma [3].

We constructed a new CRM including singly excited  $1snl$  states and doubly excited  $2l'nl$  and  $3nl$  states. Our model contains a total of 255 states;  $1s^2$ ,  $1s$ ,  $2l$ , and  $3l$  states, 60 singly excited  $1snl$  states, 118 doubly excited  $2l'nl$ , and 70 doubly excited  $3l'nl$  states, where  $n$  is up to 20. The atomic processes considered in our model are excitation / de-excitation by electron impact, ionization / three-body recombination, radiative transitions, radiative recombination and auto-ionization / dielectronic capture. Atomic data used in our CRM for  $1snl$  and  $2l'nl$  states are the same as those described in ref.[1]. For  $3l'nl$  states we calculated the data using HULLAC code [4].

Indirect contributions to effective 1s-2l excitation rate coefficients exceed direct excitation rate coefficient at  $N_e \sim 10^{19} \text{ cm}^{-3}$  for carbon ions ( $Z = 6$ ) as shown in Fig.1. Indirect contributions through autoionization after dielectronic capture,  $1s \rightarrow 3l'nl \rightarrow 2l$ , which is a resonance contribution, are important at low density. This resonance contribution decreases with increasing electron density at  $N_e > 10^{21} \text{ cm}^{-3}$ , as predicted in Ref.[5] by dielectronic capture ladder-like (DL) excitation-ionization process  $1s \rightarrow 3l'nl \rightarrow 3l'n'l'' \rightarrow \dots \rightarrow 3l'$ . With increasing density, indirect contributions through DL excitation-ionization,  $1s \rightarrow 2l'nl \rightarrow 2l'n'l'' \rightarrow \dots \rightarrow 2l'$  increase. At very high density ( $> 10^{21} \text{ cm}^{-3}$ ), indirect processes  $1s \rightarrow 1snl \rightarrow 2l'nl \rightarrow 2l'$  increase the effective rate coefficients proportionally to  $N_e$  [2]. We also found that for 2s-1s de-excitation,  $l$ -changing transitions,  $2snl$ - $2pnl$ , are important at intermediate densities.

Two population mechanisms are considered for doubly excited states at high density in ionizing phase ( $N_{I,s} = 1$ ). The first one is DL excitation-ionization,  $1s \rightarrow 2l'nl \rightarrow 2l'$ , where contributes to the increase in the effective excitation / de-excitation. The second one is indirect process  $1s \rightarrow 1snl \rightarrow 2l'nl \rightarrow 2l'$  where the excitation  $1snl \rightarrow 2l'nl$  process is important to produce the population densities of  $2l'nl$  states. At high density, a population density for a high- $n$  doubly excited state  $q(2l'nl)$  can be written approximately using the method described in Ref.[6],

$$N_q \sim \frac{C_{q-1,q}}{C_{q,q+1}} N_{q-1} + \frac{C_{i,q}}{C_{q,q+1}} N_i. \quad (1)$$

The excitation rate coefficients are scaled as  $C_{q-1,q} \propto (n-1)^7/n^3$ ,  $C_{q,q+1} \propto n^7/(n+1)^3$ . The excitation rate coefficients  $C_{i,q}$  for  $1snl$ - $2l'nl$  do not depend on a principal quantum number  $n$ . We can derive  $N_q/g_q \propto n^{-6}$  from the first term in eq.(1), as is shown in Ref.[6], where  $g_i$  is a statistical weight. The second term in eq.(1) corresponds to the second population mechanism, and  $N_q/g_q \propto n^{-4}$  [3] is derived since  $N_i \propto n^2$  in LTE. At very high density, the second process has a stronger effect than dose the first. So  $n$ -dependence of the population density per statistical weight changes from  $n^{-6}$  to  $n^{-4}$  at very high density as shown in Fig.2.

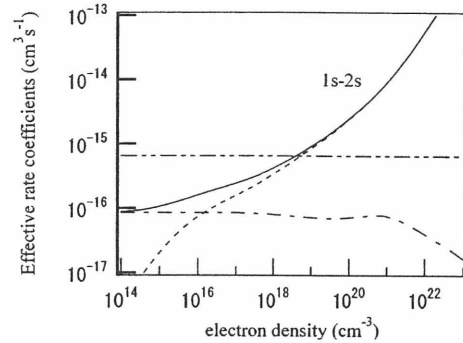


Fig.1. The effective 1s-2s excitation rate coefficients for  $Z = 6$  at  $T_e = 3.5 \times 10^5 \text{ K}$ . Solid lines: indirect contribution including  $2l'nl$  and  $3l'nl$  states, dotted lines: without  $3l'nl$  states, dot-dashed line: dielectronic capture resonances through  $3l'nl$  states, dot-dot-dashed lines: direct 1s-2s.

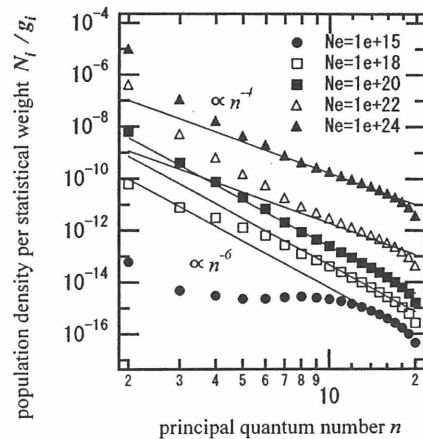


Fig.2 population density per statistical weight for doubly excited  $2snl$  states for  $Z = 6$  at  $T_e = 3.5 \times 10^5 \text{ K}$ .

### Reference

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