

Effects of electron temperature and density on ion-dust bremsstrahlung spectrum in dusty plasmas

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The effects of electron temperature and density on the ion-dust grain bremsstrahlung process in dusty plasmas are studied. The ion-dust bremsstrahlung radiation cross section is obtained as a function of the dust charge, dust radius, Debye length, collision energy, radiation energy, electron density, and electron temperature by using the Born approximation. It is shown that the ion-dust bremsstrahlung radiation cross section decreases with an increase in the electron density in dusty plasmas. It is also shown that the electron temperature suppresses the bremsstrahlung radiation cross section. In addition, the effect of electron temperature on the ion-dust bremsstrahlung process is found to be more significant than the effect of electron density in dusty plasmas. © 2009 American Institute of Physics. [DOI: 10.1063/1.3126453]

The bremsstrahlung¹⁻⁷ in plasmas has received much attention since this process is known to be one of the most fundamental processes and the continuum radiation due to the bremsstrahlung process has been a diagnostic tool for investigating the plasma parameters. Recently, there has been a considerable interest in physical processes in plasmas encompassing highly charged dust grains.⁸⁻¹² In addition, it has been shown that the collective dust-plasma interactions are ubiquitous in various astrophysical and laboratory plasmas.¹⁰ Hence, the collision and radiation processes have been extensively investigated in order to obtain the information on the plasma parameters such as the plasma density and temperature in dusty plasmas. In most cases, the dusty plasmas are consisting of thermal electrons and ions and negatively charged dusty grains, i.e., three-component dusty plasmas. Hence, the ion-dust grain interaction in dusty plasmas would be different from that in only two-component ion-dust plasmas due to the influence of the plasma electrons on dusty plasmas. Thus, in this paper we have investigated the effects of electron density and temperature on the bremsstrahlung process due to the interaction of ion and negatively charged dust grain in three-component dusty plasmas. The ion-dust grain bremsstrahlung radiation cross section is obtained as a function of the dust charge, dust radius, Debye length, collision energy, radiation energy, electron density, and electron temperature by using the Born approximation for the initial and final states of the projectile ion.

In the Born approximation, the differential ion-dust grain bremsstrahlung cross section¹³ $d^2\sigma_b$ can be obtained by the second-order nonrelativistic perturbation method,

$$d^2\sigma_b = d\sigma_C dW_\omega, \quad (1)$$

where the differential elastic scattering cross section $d\sigma_C$ can be represented by

$$d\sigma_C = \frac{1}{2\pi\hbar v_0^2} |\bar{V}(\mathbf{q})|^2 q d\mathbf{q}, \quad (2)$$

\hbar is the rationalized Planck constant, v_0 is the initial relative collision velocity, $\bar{V}(\mathbf{q})$ is the Fourier transformation of the ion-dust grain interaction potential $V(\mathbf{r})$,

$$\bar{V}(\mathbf{q}) = \int d^3\mathbf{r} e^{-i\mathbf{q}\cdot\mathbf{r}} V(\mathbf{r}), \quad (3)$$

$\mathbf{q}(=\mathbf{k}_0-\mathbf{k}_f)$ is the momentum transfer, and \mathbf{k}_0 and \mathbf{k}_f are, respectively, the wave vectors of the initial and final states of the projectile ion. Here, dW_ω represents the differential photon emission probability within the frequency between ω and $\omega+d\omega$,

$$dW_\omega = \frac{\alpha}{4\pi^2} \Lambda_i^2 \sum_{\hat{\mathbf{e}}} |\hat{\mathbf{e}} \cdot \mathbf{q}|^2 \frac{d\omega}{\omega} d\Omega, \quad (4)$$

where $\alpha(=e^2/\hbar c \cong 1/137)$ is the fine structure constant, $\Lambda_i(=\hbar/m_i c)$ is the Compton wavelength of the ion, m_i is the mass of the ion, c is the velocity of the light, $\hat{\mathbf{e}}$ is the unit photon polarization vector, and $d\Omega$ is the differential solid angle. In dusty plasmas consisting of electrons, ions, and negatively charged dusty grains, the quasineutral condition¹⁴ at equilibrium is represented by

$$n_{e0} + Z_d n_{d0} = n_{i0} = n_0, \quad (5)$$

where n_{j0} ($j=e, i, d$) are, respectively, the equilibrium densities for electrons (e), ions (i), and dust grains (d), $-Z_d$ is the charge number of the dusty grains, and n_0 is the total plasma density. In the case of $|e\phi/k_B T_j| \ll 1$ ($j=e, i$), the Poisson's equation¹² in three-component dusty plasmas can be written as

$$\nabla^2 \phi = 4\pi e (n_e + Z_d n_d - n_i) \approx 4\pi e n_{e0} \frac{e\phi}{k_B T_e} \left(1 + \frac{n_0 T_e}{n_{e0} T_i} \right), \quad (6)$$

where ϕ is the electrostatic potential in dusty plasmas, k_B is the Boltzmann constant, and T_j is the temperature of species

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j. The effective Debye length λ_{eff} in three-component dusty plasmas is then given by

$$\lambda_{\text{eff}} = \lambda_D \left(\frac{n_{e0}}{n_0} + \frac{T_e}{T_i} \right)^{-1}, \quad (7)$$

where $\lambda_D (\equiv k_B T_e / 4\pi n_0 e^2)^{1/2}$ is the standard Debye length in electron-ion plasmas and the correction factor $(n_{e0}/n_0 + T_e/T_i)$ stands for the influence of the dusty plasma on the Debye length. Since the correction factor is usually greater than unity, the effective Debye length in dusty plasmas is found to be greater than that in customary electron-ion plasmas. Hence, in spherical polar coordinates with their origin at the center of the dust grain, the interaction potential $V_{i-d}(r)$ between the ion with charge ze and negatively charged dust grain with charge $-Z_d e$ in dusty plasmas is represented by the Yukawa form with the effective Debye length λ_{eff} ,

$$V_{i-d}(r) = -\frac{Z_d z e^2}{r} \exp(-r/\lambda_{\text{eff}}). \quad (8)$$

For the sake of simplicity, the dust grains are assumed to be spherical shapes throughout in this work. After some algebra, the Fourier transformation $\bar{V}_{i-d}(q)$ of the ion-dust grain interaction potential is then given by

$$\begin{aligned} \bar{V}_{i-d}(q) &= -Z_d z e^2 \int \frac{d^3 \mathbf{r}}{r} \exp(-r/\lambda_{\text{eff}} - i\mathbf{q} \cdot \mathbf{r}) \\ &= -\frac{4\pi Z_d z e^2 a}{q} \frac{\exp(-a/\lambda_{\text{eff}})}{(a/\lambda_{\text{eff}})^2 + (qa)^2} \\ &\quad \times [(a/\lambda_{\text{eff}}) \sin(qa) + (qa) \cos(qa)], \end{aligned} \quad (9)$$

where a is the radius of the spherical dust grain. Hence, the ion-dust bremsstrahlung cross section in dusty plasmas is then given by

$$\begin{aligned} d^2 \sigma_b &= \frac{16}{3} \frac{\alpha^3 a_0^2 m_e}{\bar{E}_0} Z_d^2 z^2 \exp(-2\bar{\lambda}_{\text{eff}}^{-1}) \\ &\quad \times \left(\frac{\bar{\lambda}_{\text{eff}}^{-1} \sin \bar{q} + \bar{q} \cos \bar{q}}{\bar{\lambda}_{\text{eff}}^{-2} + \bar{q}^2} \right)^2 \bar{q} d\bar{q} \frac{d\omega}{\omega}, \end{aligned} \quad (10)$$

where $a_0 (\equiv \hbar^2 / m_e e^2)$ is the first Bohr radius of the hydrogen atom, m_e is the mass of the electron, $\bar{E}_0 (\equiv m_e v_0^2 / 2Ry)$ is the scaled initial collision energy, $Ry (\equiv m_e e^4 / 2\hbar^2 \approx 13.6 \text{ eV})$ is the Rydberg constant, $\bar{q} (\equiv qa)$ is the scaled momentum transfer, and $\bar{\lambda}_{\text{eff}} (\equiv \lambda_{\text{eff}} / a)$ is the scaled effective Debye length. This expression of the bremsstrahlung cross section is reliable when the kinetic energy of the projectile ion is greater than the interaction energy between the ion and dust grain due to the prerequisite of the Born approximation. It has been also shown that the continuum spectrum due to the bremsstrahlung process would be investigated through the bremsstrahlung radiation cross section¹⁵ defined as $d^2 \chi_b / d\bar{\varepsilon} d\bar{q} \equiv \hbar \omega (d\sigma_b / \hbar d\omega d\bar{q})$, where $\bar{\varepsilon} (\equiv \varepsilon / Ry)$ is the scaled photon energy and $\varepsilon (\equiv \hbar \omega)$ is the photon energy. After some mathematical manipulations, the ion-dust grain bremsstrahlung radiation cross section $d^2 \chi_b / d\bar{\varepsilon}$ in three-component dusty plasmas is then obtained by the following analytic form:

$$\begin{aligned} \frac{d^2 \chi_b}{d\bar{\varepsilon}} &= \frac{4}{3} \frac{\alpha^3 a_0^2}{\bar{E}_0} (\eta Z_d z)^2 \exp(-2\xi^{-1} \bar{\lambda}_D^{-1}) \\ &\quad \times \left| \ln \left\{ \frac{1 + [\bar{a} \eta \xi \bar{\lambda}_D (\sqrt{\bar{E}_0} + \sqrt{\bar{E}_f})]^2}{1 + [\bar{a} \eta \xi \bar{\lambda}_D (\sqrt{\bar{E}_0} - \sqrt{\bar{E}_f})]^2} \right\} \right. \\ &\quad + (1 + 2\xi^{-1} \bar{\lambda}_D^{-1}) [\cosh(2\xi^{-1} \bar{\lambda}_D^{-1}) + \sinh(2\xi^{-1} \bar{\lambda}_D^{-1})] \\ &\quad \times \left\{ \sum_{l=1}^2 \{ \text{Ci}[(-1)^l 2i \xi^{-1} \bar{\lambda}_D^{-1} + 2\bar{a} \eta (\sqrt{\bar{E}_0} + \sqrt{\bar{E}_f})] \right. \\ &\quad \left. - \text{Ci}[(-1)^l 2i \xi^{-1} \bar{\lambda}_D^{-1} + 2\bar{a} \eta (\sqrt{\bar{E}_0} - \sqrt{\bar{E}_f})] \right\} \\ &\quad + i \sum_{l=1}^2 \{ \text{Si}[2i \xi^{-1} \bar{\lambda}_D^{-1} + (-1)^l 2\bar{a} \eta (\sqrt{\bar{E}_0} + \sqrt{\bar{E}_f})] \\ &\quad \left. - \text{Si}[2i \xi^{-1} \bar{\lambda}_D^{-1} + (-1)^l 2\bar{a} \eta (\sqrt{\bar{E}_0} - \sqrt{\bar{E}_f})] \right\} \\ &\quad + \frac{2 \cos[\bar{a} \eta (\sqrt{\bar{E}_0} + \sqrt{\bar{E}_f})]}{1 + [\bar{a} \eta \xi \bar{\lambda}_D (\sqrt{\bar{E}_0} + \sqrt{\bar{E}_f})]^2} \\ &\quad - \frac{2 \cos[\bar{a} \eta (\sqrt{\bar{E}_0} - \sqrt{\bar{E}_f})]}{1 + [\bar{a} \eta \xi \bar{\lambda}_D (\sqrt{\bar{E}_0} - \sqrt{\bar{E}_f})]^2} \\ &\quad - \frac{2\bar{a} \eta \xi \bar{\lambda}_D (\sqrt{\bar{E}_0} + \sqrt{\bar{E}_f}) \sin[\bar{a} \eta (\sqrt{\bar{E}_0} + \sqrt{\bar{E}_f})]}{1 + [\bar{a} \eta \xi \bar{\lambda}_D (\sqrt{\bar{E}_0} + \sqrt{\bar{E}_f})]^2} \\ &\quad \left. + \frac{2\bar{a} \eta \xi \bar{\lambda}_D (\sqrt{\bar{E}_0} - \sqrt{\bar{E}_f}) \sin[\bar{a} \eta (\sqrt{\bar{E}_0} - \sqrt{\bar{E}_f})]}{1 + [\bar{a} \eta \xi \bar{\lambda}_D (\sqrt{\bar{E}_0} - \sqrt{\bar{E}_f})]^2} \right|^2, \end{aligned} \quad (11)$$

where $\eta \equiv \sqrt{m_e / m_i}$, $\bar{E}_f (\equiv \bar{E}_0 - \bar{\varepsilon})$ is the scaled final projectile energy, $\bar{a} (\equiv a / a_0)$ is the scaled radius of the dust grain, $\bar{\lambda}_D \equiv \lambda_D / a$, $\xi \equiv (n_{e0}/n_0 + T_e/T_i)^{-1}$, and $\text{Ci}(x) (= -\int_x^\infty dt \cos t/t)$ and $\text{Si}(x) (= \int_0^x dt \sin t/t)$ are the cosine and sine integrals,¹⁶ respectively. Very recently, an excellent discussion¹⁷ on the additional part of the electrostatic potential in dusty plasmas is given due to the ion absorption by the dust grain and ion-neutral collisions. It has been found that the electrostatic potential would be represented only by the standard Debye–Hückel form if there is no ion flux on the surface of the dust grain, i.e., nonabsorbing dust grains. In this work, we just retain only the Debye–Hückel form of the interaction potential since we consider the bremsstrahlung process due to the scattering of the ion by the nonabsorbing dust grain. However, if the ion flux exists on the dust grain, the attractive part of the potential has to be included in the ion-dust bremsstrahlung process. Hence, it would be expected that the additional part of the interaction potential suppresses the ion-dust grain bremsstrahlung spectrum obtained by the Yukawa-type Debye–Hückel potential. Since the Born parameter $(Z_d z e^2 / r_i E) e^{-r_i / \lambda_{\text{eff}}}$, where r_i is the interparticle distance between the ion and dust grain, for the ion-dust interaction in typical circumstances of dusty plasmas is usually less than unity, the Born approximation would be quite reliable to investigate the high energy ion-dust bremsstrahlung process in dusty plasmas. Very recently, the accurate form of the effec-

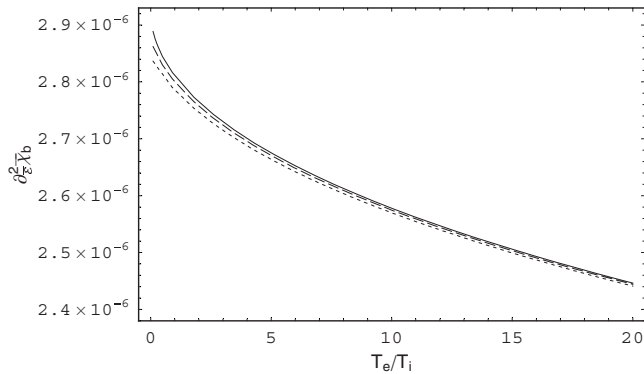


FIG. 1. The scaled ion-dust grain bremsstrahlung radiation cross section $\partial_{\bar{\epsilon}}^2 \bar{\chi}_b$ as a function of the temperature ratio T_e/T_i for $\bar{E}_0=5$, $\bar{\epsilon}=2$, $T_e=5 \times 10^3$ K, $n_0=10^{10}$ cm $^{-3}$, and $a=9.7 \times 10^{-5}$ cm. The solid line represents the case of $n_{e0}/n_0=0$. The dashed line represents the case of $n_{e0}/n_0=0.2$. The dotted line represents the case of $n_{e0}/n_0=0.5$.

time polarization interaction potential¹⁸ has been given in partially ionized dense plasmas, taking into account the quantum mechanical effects of diffraction and symmetry of particles and plasma screening effects. However, the ion-atom polarization bremsstrahlung process is neglected in this work since the dust plasma is assumed to be completely ionized. For typical circumstances of dusty plasmas, it has been shown that $Z_d \approx 100-1000$, $a \approx 0.01-1$ μm , and $\lambda_D/a \approx 5-100$.⁸

In order to specifically investigate the effects of electron temperature and density on the ion-dust bremsstrahlung process in a dusty plasma, we set $Z_d=200$, $a=0.1$ μm , $\lambda_D/a=50$, $m_i/m_e=1840$, and $z=1$. Figure 1 shows the scaled bremsstrahlung radiation cross section $\partial_{\bar{\epsilon}}^2 \bar{\chi}_b [\equiv (d^2 \chi_b / d\bar{\epsilon}) / \pi a_0^2]$ in units of πa_0^2 for the interaction of the ion with the negatively charged dust grain in dusty plasmas as a function of the temperature ratio T_e/T_i for various values of the density ratio n_{e0}/n_0 . In addition, Fig. 2 represents the scaled bremsstrahlung radiation cross section $\partial_{\bar{\epsilon}}^2 \bar{\chi}_b$ as a function of the density ratio n_{e0}/n_0 for various values of the

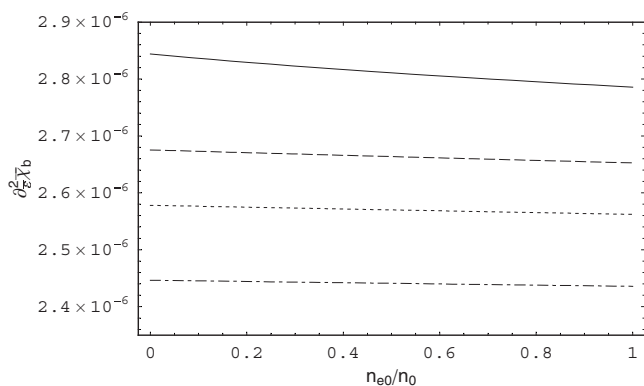


FIG. 2. The scaled ion-dust grain bremsstrahlung radiation cross section $\partial_{\bar{\epsilon}}^2 \bar{\chi}_b$ as a function of the temperature ratio n_{e0}/n_0 for $\bar{E}_0=5$, $\bar{\epsilon}=2$, $T_e=5 \times 10^3$ K, $n_0=10^{10}$ cm $^{-3}$, and $a=9.7 \times 10^{-5}$ cm. The solid line represents the case of $T_e/T_i=0.5$. The dashed line represents the case of $T_e/T_i=5$. The dotted line represents the case of $T_e/T_i=10$. The dot-dashed line represents the case of $T_e/T_i=20$.

temperature ratio T_e/T_i . As shown in these figures, it is found that the ion-dust bremsstrahlung radiation cross section decreases with an increase in the ratio of the electron temperature to the ion temperature in a dusty plasma. It is also shown that the electron density suppresses the ion-dust bremsstrahlung radiation cross section. It is interesting to note that the effect of electron density on the bremsstrahlung radiation cross section diminishes with an increase in the electron temperature since the high energy electrons can be readily repelled by the negatively charged dusty grains in dusty plasmas.¹² The effect of electron density is found to be important for low-temperature ratios. It is also found that the effect of electron temperature on the bremsstrahlung radiation cross section is found to be more significant than the effect of electron density in dusty plasmas. Hence, we have found that the effects of electron temperature and density play important roles in the ion-dust grain bremsstrahlung process in dusty plasmas containing electrons, ions, and negatively charged dust grains. These results would provide useful information on the ion-dust bremsstrahlung emission spectrum and also the radiation due to the interaction between dust particle chains and streaming ions in the plasma sheath in dusty plasmas.

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