



Effect of Sm^{3+} substitution on structural and magnetic investigation of nano sized Mn–Sm–Zn ferrites

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Abstract: Nano size $\text{Mn}_{0.4}\text{Zn}_{0.6}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$ ($x = 0.00, 0.01, 0.02, 0.03, 0.04, \text{ and } 0.05$) ferrites were prepared by solution combustion method. The structural and magnetic properties of samples were characterised by X-ray diffractometer, Fourier transform infrared spectrometer, transmission electron micrographs and magnetic properties at room temperature. The X-ray diffraction patterns and two prominent absorption bands in the frequency range $375\text{--}589\text{ cm}^{-1}$ confirmed the single phase with spinel cubic structure. The average nano crystallite sizes were in best agreements with transmission electron microscope images. Magnetic studies revealed the narrow hysteresis loops of ferrimagnetic nature at room temperature. The values of saturation magnetization (M_s), remanence magnetization (M_r), coercivity (H_c), remanence ratio (M_r/M_s), magneton number, anisotropy constant and Yaffet–Kittle angle decreased with the increase in Sm^{3+} ion concentration was attributed to relative number of ferric ions on the tetrahedral sites diminished and reduced the Sm–Fe interaction.

Keywords: Nanoferrites; Magnetization; Anisotropy constant; Y–K angle

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1. Introduction

In recent years, nanoferrites have attracted much attention due to their technological applications as microwave devices, high speed digital tapes and disk recording, ferrofluids, magnetic refrigeration systems, biosensors, recording colour imaging, imaging drug delivery and hyperthermia for tumor therapy [1, 2]. The high-permeability and low-power losses at megahertz frequency can be achieved using nanocrystalline ferrites in place of micro-magnetic ferrites [3]. Mn–Zn nanoferrites have excellent properties, particularly their high initial permeability, high saturation magnetization, high resistivity and low power loss [4–7]. Nanosized Mn–Zn ferrites are extensively used as inductors, transformers cores, rod antennas, loading coils, deflection yokes, choke coils, noise filters, recording

heading, magnetic amplifiers and electromagnetic interference devices (EMI) etc. Moreover, these ferrites are very important in biomedicine as magnetic carriers for bioseparation, enzymes and proteins immobilization, magnetic resonance imaging (MRI), sensing as well as therapeutic applications such as ac and ac magnetic field—assisted cancer therapy [8–10]. The most significant properties in the nano-region depend on the shape, size, crystallinity and occupation of the nano-particle [11].

The rare earth substituted ferrites are nowadays under extensive investigations and are becoming promising materials for various applications. Studies on rare earth ferrites have shown that they are more effective for the relief of acute and chronic pain or discomfort due to the various injuries and ailments [12]. It is well known that the rare earth ions have unpaired $4f$ electrons screened by $5s$ and $5p$ orbital and are almost not affected by the potential field of surroundings ions. The substitution of rare earth ions into spinel ferrites results in occurring of $4f\text{--}3d$ couplings which determine the magneto crystalline

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anisotropy in ferrites and can also improve the electrical and magnetic properties of spinel ferrites [13].

Nanosized Mn–Zn ferrites are synthesized by several methods, such as solid state reaction [14], chemical coprecipitation [15], hydrothermal synthesis [16, 17], sol–gel combustion [18, 19], spray pyrolysis methods [20], citrate precursor [21] etc. However, most of these methods are economically unfeasible for large scale production of ferrites. Therefore, in the present investigation, the combustion reaction synthesis has been used successfully to obtain nanosized particles. This method has several advantages due to its simplicity, short reaction time and inexpensive nature, allowing for the production of fine, homogeneous crystalline powders without the risk of contamination, lower consumption of energy and elimination of intermediary calcinations stages during synthesis [22]. In the present investigation, the effect of Sm^{3+} ions on structural and magnetic properties of Mn–Zn nanoferrites prepared by combustion method is reported.

2. Experimental details

In the present study the polycrystalline $\text{Mn}_{0.4}\text{Zn}_{0.6}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$ were prepared by solution combustion route, using stoichiometric molar amounts of manganese nitrate [$\text{Mn}(\text{NO}_3)_2 \cdot 4\text{H}_2\text{O}$], zinc nitrate [$\text{Zn}(\text{NO}_3)_2 \cdot 6\text{H}_2\text{O}$], and iron nitrate [$\text{Fe}(\text{NO}_3)_3 \cdot 9\text{H}_2\text{O}$], samarium nitrate [$\text{Sm}(\text{NO}_3)_2$] as oxidizer and mixture of urea [NH_2CONH_2] and glucose [$\text{C}_6\text{H}_{12}\text{O}_6$] as fuels. The stoichiometric compositions of metal nitrates and fuels were evaluated based on the total oxidizing and reducing valences of the component. The oxidizer to fuel ratio was taken as 1: (60:40). All the metal nitrates and fuels were diluted with 30 ml double distilled water and thoroughly mixed using magnetic stirrer with the speed of 800 rpm for 1 h until the reactants were dissolved completely to get homogenous solution. This homogeneous solution containing redox mixture was taken in a Pyrex dish and kept in a pre heated muffle furnace maintained at 450 ± 10 °C. Initially, the solution boiled, then frothed and ignited to yield fine powder of Sm^{3+} co-substituted Mn–Zn rapidly because these were exothermic. The whole combustion process was complete in <20 min, whereas the reaction time of the actual ignition was <5 s. The samples were pressed into pellets by applying the pressure of 5-ton/ cm^2 for 5 min. The pellets were sintered at 1000 °C for 8 h and furnace cooled to room temperature. The X-ray diffraction (XRD) patterns were obtained at room temperature on Philips make PW-3710 X-ray powder diffractometer in the range of 20–80 with CuK_α radiation ($\lambda = 0.154056$ nm). The crystallite size was calculated by Debye–Scherrer method, given by Eq. (1) below. The IR absorption spectra

of powdered samples were recorded in the range of 350–800 cm^{-1} on Perkin-Elmer FT-IR spectrum one spectrometer using KBr pellet technique. The distance between magnetic ions and hopping lengths in octahedral sites and tetrahedral sites were calculated by using Eq. (2) [23]

$$D_P = \frac{0.94\lambda}{\beta \times \cos \theta} \quad (1)$$

$$L_A = a \frac{\sqrt{3}}{4} \quad \text{and} \quad L_B = a \frac{\sqrt{3}}{2} \quad (2)$$

The magnetization measurement was carried out in vibrating sample magnetometer (VSM) mounted on an electromagnet with a bipolar source of maximum applied field of 1.5 T at room temperature. The Remanence ratio (S), Magneton number (η_B), Anisotropy constant (K), and Y–K angle of the samples with different Sm^{3+} concentration is calculated using Eqs. (3)–(6) [24, 25].

$$R = \frac{M_r}{M_s} \quad (3)$$

$$(\eta_B) = \frac{MW \times M_s}{5585} \quad (4)$$

$$K = \frac{Hc \times M_s}{0.96} \quad (5)$$

$$\cos \alpha_{Y-k} = \frac{\eta_B + 5(1-X)}{5(1+X)} \quad (6)$$

3. Results and discussion

Figure 1 shows the XRD pattern of as-prepared $\text{Mn}_{0.4}\text{Zn}_{0.6}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$ ($x = 0.00, 0.01, 0.03, 0.05$)

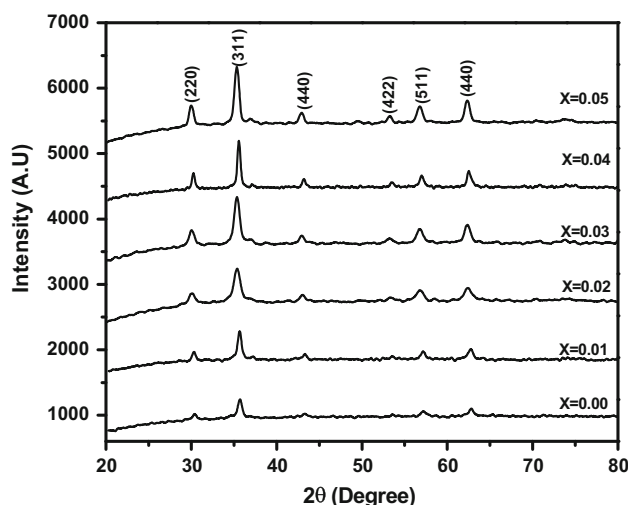


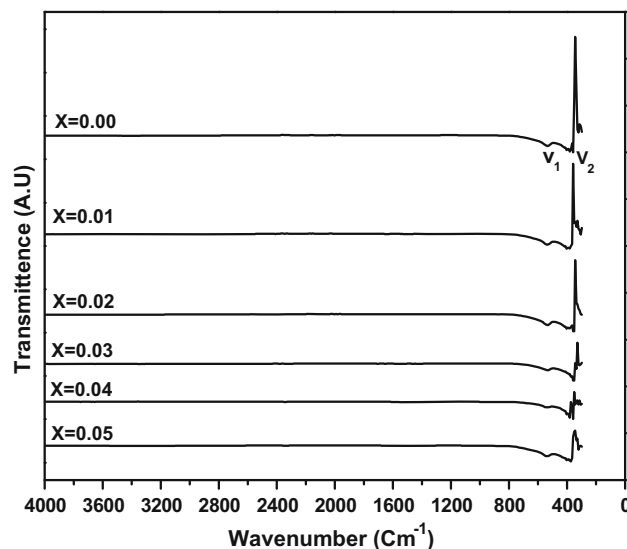
Fig. 1 XRD patterns of as-synthesized $\text{Mn}_{0.4}\text{Zn}_{0.6}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$ ($x = 0.00, 0.01, 0.02, 0.03, 0.04$ and 0.05)

Table 1 Data on lattice parameter (A), crystallite size (D_p) by XRD, particle size (T) by TEM, hopping length (L_A and L_B) and position of absorption bands of Mn_{0.4}Zn_{0.6}Sm_xFe_{2-x}O₄ (x = 0.00, 0.01, 0.02, 0.03, 0.04, and 0.05)

Composition (X)	Lattice parameter (a) Å	Particle size (XRD) nm	Particle size (TEM) nm	Hopping lengths		Absorption bands	
				A-site (L _A)	B-site (L _B)	v ₁ (cm ⁻¹)	v ₂ (cm ⁻¹)
0.00	8.344	27	30	3.613	7.226	529	343
0.01	8.410	16	18	3.641	7.283	537	380
0.03	8.412	13	14	3.642	7.285	537	358
0.05	8.417	10	12	3.644	7.289	537	372

nanostructured particles from which proper phase identification cannot be observed. Analysis of XRD pattern confirms the presence of strong diffraction peaks corresponding to the planes (1 1 1), (2 2 0), (3 1 1), (4 0 0), (4 2 2), (5 1 1/3 3 3), (4 4 0), indicating the formation of spinel cubic structure. After certain limit, Sm accumulates at the grain boundaries of Mn–Zn Ferrite and forms orthorhombic phase. This may be due to the presence of SmFe₂O₃ phase formed due to high reactivity of iron with samarium [26]. The lattice constant (Table 1) is found to increase with increase in Sm³⁺ ion concentration which is attributed to the fact that the ionic radii of Fe³⁺ (0.67 Å) is smaller than that of Sm³⁺ (0.99 Å) ions substituted for Sm³⁺ ions. This confirms the occupancy of samarium on octahedral site [26]. The particles size of the powder samples estimated from the full width at half maxima of pattern (FWHM) of their (311) XRD planes by employing Debye–Scherrer relation [27]. The average particle size of the samples is in the range of 10–27 nm.

Infrared spectra of the samples for different compositions are shown in Fig. 2. Infrared spectra of the samples show two prominent absorption bands, v₁ near 600 cm⁻¹ attributed to tetrahedral complexes and v₂ around 400 cm⁻¹ assigned to octahedral complexes. The positions of the absorption bands are listed in Table 1. The difference in frequencies between v₁ and v₂ is due to changes in bond length (Fe³⁺–O²⁻) at tetrahedral and octahedral sites [28]. The slight change in band position, particularly in v₁ and v₂ suggests that the method of preparation, grain size and porosity can influence in locating the band position [29]. It can be seen that the absorption band v₁ does not show any splitting or shoulders and hence the possibility of Fe²⁺ ions at A-sites is ruled out [30]. On substitution of Sm³⁺ ion content, v₁ continues to be widen is attributed to atomic mass and volume of the dopant, which affects Fe–O distances on octahedral sites, suggesting that Sm³⁺ ions occupy octahedral lattice sites [31]. The shift occurs in absorption bands v₁ and v₂ for each octahedral and tetrahedral site due to the perturbation occurring in the Fe³⁺–O²⁻ bond by introducing Sm³⁺ ions [31].

**Fig. 2** FTIR spectra of as-synthesized Mn_{0.4}Zn_{0.6}Sm_xFe_{2-x}O₄ (x = 0.00, 0.01, 0.02, 0.03, 0.04, and 0.05)

The TEM micrographs of samples with Sm³⁺ content (x = 0.03 and 0.05) annealed at 350 °C are shown in Fig. 3. It is clear that the particles are spherical in shape and are agglomerated. This agglomeration can be attributed to magnetic interaction arising among ferrite nano-particles. The particle size (D_T) calculated from TEM, is larger than crystal size estimated by XRD, indicating that majority of particles in samples are multi-crystal particles. The addition of Sm³⁺ ions reduces the grain growth probably due to deposition of Sm³⁺ on the grain boundaries which in turn hampers its motion [32].

Figure 4 shows the magnetic hysteresis curves for different compositions of the samples at room temperature very narrow hysteresis loops, reveal the behaviour of soft magnetic material. From the Table 2, it is clear that the values of saturation magnetization (M_s), remanence magnetization (M_r), remanence ratio (M_r/M_s), and coercivity (H_c) and magneton number (n_B) decrease with increase in Sm³⁺ ion concentration. According to Neel's two sub-lattice model of ferrimagnetism, the net moment is given by the formula μ_B = M_B(x) – M_A(x), where M_A and M_B are

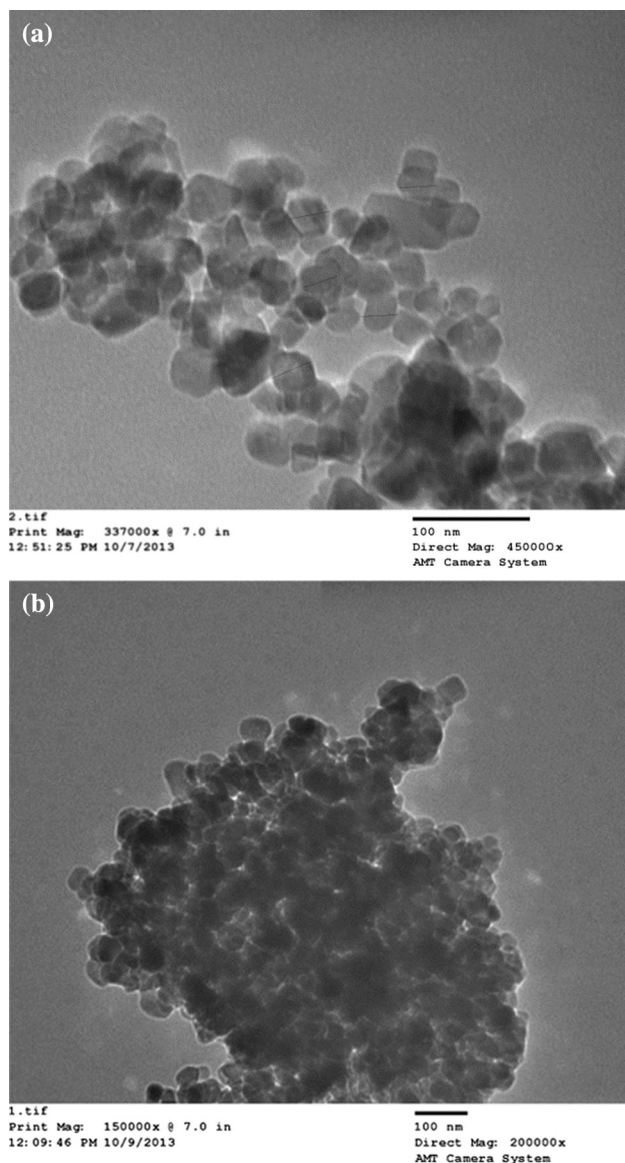


Fig. 3 TEM images of as-synthesized (a) $\text{Mn}_{0.4}\text{Zn}_{0.6}\text{Sm}_{0.03}\text{Fe}_{1.97}\text{O}_4$ and (b) $\text{Mn}_{0.4}\text{Zn}_{0.6}\text{Sm}_{0.05}\text{Fe}_{1.95}\text{O}_4$

the A and B sub-lattice magnetic moments in μ_B [33]. The existence of random canting of particle surface spins, surface effects and the occurrence of a glassy state have been reported to play an active role in the decline of magnetization values [34]. The magnetic moment decreases significantly in Sm^{3+} substituted samples which is attributed to small free ion of magnetic moment of Sm^{3+} , $n_B = 1.58\mu_B$ [35] as against for Fe^{3+} ; $n_B = 5\mu_B$ at room temperature [36, 37]. The substituted Sm^{3+} ion with $4f$ and $5p$ electrons, occupies on octahedral site and hence Sm-Fe interaction is weaker than Fe-Fe interaction. This dilutes the interaction in octahedral site and decreases the Curie temperature.

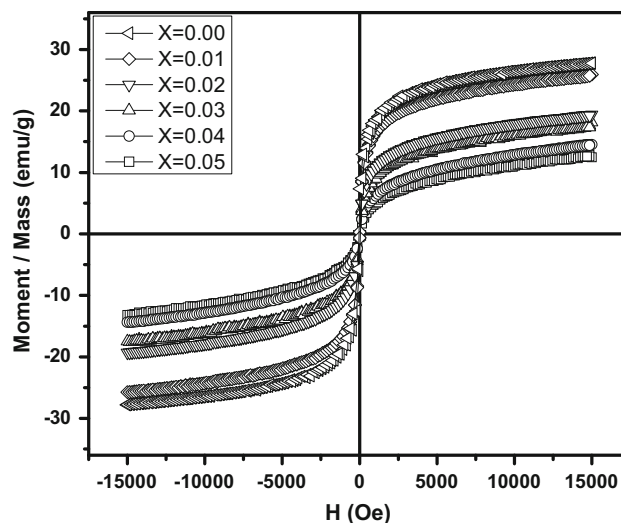


Fig. 4 Hysteresis loop of as-synthesized $\text{Mn}_{0.4}\text{Zn}_{0.6}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$ ($x = 0.00, 0.01, 0.02, 0.03, 0.04, \text{ and } 0.05$)

In ferrites, the coercive force is obtained by reversal of the directions of the wall movement and that of domain rotation on reversing the direction of the applied magnetic field. Generally, the effective pinning for domain wall causes the coercivity; it is known that the larger grain size decreases H_C [38]. In the present investigation, the coercive values are low; hence probability of domain rotation is also lower. The materials with larger grain size have been used to achieve lower core loss [38].

In the present work, the Y-K angles decrease with the increase in Sm^{3+} ion concentration. The decrease in the Y-K angles suggests a decrease of triangular spin arrangements on B sites to produce the increment in A-B interaction. The effect of B-B interaction can be masked by strong A-B interaction to cause the spin on B sites to be aligned parallel to each other. However, the substitution of Sm^{3+} seems to lead to canted type of arrangements on B sites to enhance the A-B interactions. Therefore, Sm^{3+} substitution brings changes in magnetization, which can be attributed to the presence of Y-K angles in the spin system on B sites. The condition for Y-K angles has been discussed from the consideration of the molecular field approximation using a noncolinear three sub-lattice model [39, 40].

4. Conclusions

Samarium substituted MnZn nanoferrites have been prepared by combustion method. XRD analysis reveals that all the prepared samples are single phase cubic spinel without any secondary phase. Increase in lattice constant value with samarium doping indicates the expansion of unit cell. The

Table 2 Data on saturation magnetization (M_s), remanence magnetization (M_r), remanence ratio (M_r/M_s), and coercivity (H_c), magneton number (n_B), anisotropy constant (k) and Y–K angle of Mn_{0.4}Zn_{0.6}Sm_xFe_{2-x}O₄ ($x = 0.00, 0.01, 0.02, 0.03, 0.04, \text{ and } 0.05$)

Samarium concentration (X)	Saturation Magnetization (M_s) in emu/gm	Remenant magnetization (M_r) in emu/gm	Coercivity (H_c) in Oe	Remanance ratio M_r/M_s	Magneton number (n_B) μ_B	Anisotropy constant (K) in erg/Oe	Yafet–Kittel angle α_{Y-K}
0.00	27.77	16.68	50.14	0.7148	1.703	1450	1.340
0.01	25.92	18.53	38.93	0.6491	1.589	1051	1.320
0.02	19.40	12.60	16.47	0.6006	1.189	333	1.242
0.03	18.13	10.82	11.01	0.5968	1.111	208	1.227
0.04	14.53	07.23	09.31	0.4975	0.891	141	1.183
0.05	12.85	05.81	04.10	0.4521	0.788	69	1.162

absorption bands in FTIR are found in the expected range of spinel ferrites. TEM images show the samples are spherical in shape and agglomerated. The magnetic behavior of pure and samarium doped magnesium ferrite particles shows soft ferrimagnetic nature. The decrease in saturation magnetization and remnant magnetization is attributed to the weak magnetic interactions.

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References

- [1] P A Shaikha, R C Kambalea, A V Rao and Y D Kolekar *J. Alloys Compd.* **492** 590 (2010)
- [2] X Cao, G Liu, Y Wang, J Li and R Hong *J. Alloys Compd.* **497** L9 (2010)
- [3] K Praveena, K Sadhana, S Bharadwaj and S R Murthy *J. Magn. Magn. Mater.* **321** 2433 (2009)
- [4] H B Yang et al. *J. Magn. Magn. Mater.* **322** 173 (2010)
- [5] A C F M Costa, V J Silva, C C Xin, D A Vieira, D R Cornejo and R H G A Kiminami *J. Alloys Compd.* **495** 503 (2010)
- [6] M M Hessien, M M Rashad, K El-Barawy, I A Ibrahim *J. Magn. Magn. Mater.* **320** 1615 (2008)
- [7] A Goldman *Modern Ferrite Technology*, 2nd edn. (Berlin: Springer) (2006)
- [8] P Papazoglou, F Eleftheriou and V T Zaspalis *J. Magn. Magn. Mater.* **296** 25 (2006)
- [9] J Wang, T Deng and Y Dai *J. Alloys Compd.* **419** 155 (2006)
- [10] S F Mansour and M A Elkestawy *Ceram. Int.* **37** 1175 (2011)
- [11] S D Shenoy, P A Joy and M R Anantharaman *J. Magn. Magn. Mater.* **269** 217 (2004)
- [12] E Ateia, *Egypt J. Sol.* **29**(2) 317 (2006)
- [13] S Jie, W Lixi, X Naicen and Z Qitu *J. Rare Earth* **28** 451 (2010)
- [14] M R Meshram, N K Agrawal, B Sinha and P S Misra *J. Magn. Magn. Mater.* **271** 214 (2004)
- [15] R Arulmurugan, G Vaidyanathan, S Sendhilnathan and B Jeyadevan *J. Magn. Magn. Mater.* **298** 83 (2006)
- [16] J F Wang, C B Ponton and I R Harris, *J. Magn. Magn. Mater.* **298** 122 (2006)
- [17] J F Wang, C B Ponton and I R Harris *J. Alloys Compd.* **403** 104 (2005)
- [18] A Mali and A Ataie *J. Alloys Compd.* **399** 245 (2005)
- [19] Q Fang, H Cheng, K Huang, J Wang, R Li and Y Jiao *J. Magn. Magn. Mater.* **294** 281 (2005)
- [20] C H Chen, M H J Emond, E M Kelder, B Meester and J Schoonman *J. Aerosol Sci.* **30**(7) 959 (1999)
- [21] A Thakur and M Singh *Ceram. Inter.* **29** 505 (2003)
- [22] A C F M Costa, V J Silva, D R Cornejo and L G F Vieira *J. Magn. Magn. Mater.* **320** e370 (2008)
- [23] E Ateia *Egypt. J. Solids* **29**(2) 329 (2006)
- [24] A A Pandit, A R Shitre, D R Shengule and K M Jadhav, *J. Mater. Sci.* **40** 423 (2005)
- [25] M L Mane et al, *J. Mol. Struct.* **1035** 27 (2013)
- [26] E Melagiriyyappa, M Veena, A Somashekarappa, G J Shankaramurthy and H S Jayanna *Indian J. Phys.* **88** 795 (2014)
- [27] H P Klug and L E Alexander, *X-ray Diffraction Procedure*. (Wiley–Inter Science, New York) Ch. 9 (1984)
- [28] O M Hameda *J. Magn. Magn. Mater.* **281** 36 (2004)
- [29] M A El Hiti, A I El Shora, A S Seoud and S M Hammed *Phase Trans.* **56** 35 (1995)
- [30] V A Halfner, N D Zverev, V P Romanov *Phys. Status Solidi A* **12** 623 (1975)
- [31] E Melagiriyyappa, H S Jayanna *J. Alloys Compd.* **482** 147 (2009)
- [32] A B Gadkari, T J Shinde and P N Vasambekar *Mater. Charact.* **60**(11) 1328 (2009)
- [33] L Neel *Proc. Phys. Soc. Lond.* **A-65** 869 (1952)
- [34] J Chand, S Verma, M Singh *J. Alloys Compd.* **552** 264 (2013)
- [35] B N Figgis, in: G Wilkinson, Gillard and J A Mc Cleverty (eds.), *Comprehensive Coordination Chemistry*, (New York: Pergamon) p 261 (1987)
- [36] K J Standley *Oxide Magnetic Materials* (Oxford: Clarendon Press) p 52 (1972)
- [37] V Naidu, S Vijayaragavan, R Legadevi and A Santhil Kumar *Int. J. Comp. Appl.* **30** 7 (2011)
- [38] P Yaseneva, M Bowker and G Hutchings *Phys. Chem. Chem. Phys.* **13** 18609 (2011)
- [39] M U Rana, M Ul-Islam, I Ahmad and T Abbas *J. Magn. Magn. Mater.* **187** 242 (1999)
- [40] N S S Murthy, M G Natera, S I Youssef, R J Begum *Phys. Rev.* **181** 969 (1969)