Materials Research Bulletin 48 (2013) 180-187

Contents lists available at SciVerse ScienceDirect



Materials Research Bulletin



journal homepage: www.elsevier.com/locate/matresbu

Structural characterization, thermoluminescence and EPR studies of Nd₂O₃:Co²⁺ nanophosphors

B. Umesh^{a,b}, B. Eraiah^{b,**}, H. Nagabhushana^{c,*}, S.C. Sharma^c, D.V. Sunitha^c, B.M. Nagabhushana^d, J.L. Rao^e, C. Shivakumara^f, R.P.S. Chakradhar^{g,***}

^a Department of Humanities, PVP Polytechnic, Dr. AIT Campus, Bangalore 560 056, India

^b Department of Physics, Bangalore University, Bangalore 560 056, India

^c Prof. C.N.R. Rao Centre for Advanced Materials, Tumkur University, Tumkur 572 103, India

^d Department of Chemistry, M.S. Ramaiah Institute of Technology, Bangalore 560 054, India

^e Department of Physics, Sri Venkateswara University, Tirupati 517 502, India

^f Solid State and Structural Chemistry Unit, Indian Institute of Science, Bangalore 560 012, India

^g CSIR – National Aerospace Laboratories, Bangalore 560 017, India

ARTICLE INFO

Article history: Received 29 March 2012 Received in revised form 28 August 2012 Accepted 2 September 2012 Available online 7 September 2012

Keywords:

A: Oxides

B: Chemical synthesis

C: X-ray diffraction

D: Electron paramagnetic resonance (EPR)

D. Luminescence

1. Introduction

Nanotechnology is a growing area of research, primarily due to its numerous applications in engineering/science. Therefore, the synthesis of new nanomaterials and improving its properties are of current research interest to many researchers. It has been found that the physical properties of individual nanoparticles can be very different from those of their bulk counterparts [1]. Rare earth compounds have been extensively used in high performance luminescent devices, magnets, catalysts, and other functional materials because of their electronic, optical and chemical characteristics resulting from the 4f shell of their ions [2-4]. These properties depend strongly on the material composition

and structure, which are sensitive to the bonding states of rare

earth ions. If the rare earth oxides were fabricated in the form of

one dimensional nanostructure, they would have new properties

Corresponding author.

ABSTRACT

 Nd_2O_3 : Co²⁺ (1–4 mol%) nanophosphors (15–25 nm) have been prepared via low temperature solution combustion method. Scanning electron micrograph (SEM) shows that the product is highly porous in nature. The stokes line in the Raman spectrum at ~ 2000 cm⁻¹ is assigned to F_{e} mode and the anti-stokes lines are assigned to a combination of $A_g + E_g$ modes. With increase of Co²⁺ concentration, the intensity of F_{g} mode decreases, whereas the combination of $A_{g} + F_{g}$ modes completely disappears. Electron paramagnetic resonance (EPR) spectrum exhibits two resonance signals with effective g values at g = 2.25 and g = 2.03. Thermoluminescence (TL) response of Nd₂O₃:Co²⁺ nanopowders with γ dose 0.23– 2.05 kGy was studied. The activation energy (E) and frequency factor (s) are estimated using Chen's glow peak shape method and obtained to be in the range 0.45–1.67 eV and 1.8×10^4 to 4.0×10^{12} s⁻¹, respectively. It is observed that the TL glow peak intensity at 430 K increases linearly with γ dose which is suitable for radiation dosimetry.

© 2012 Elsevier Ltd. All rights reserved.

as a result of both their marked shape-specific and quantum confinement effects. They would also act as electrically, magnetically, or optically functional host materials for rare earth ions in phosphor displays, waveguide devices, and for the active material in lasers [5].

Studies on radiation induced defects in insulating materials have been interesting over the last few decades. TL is one such radiation induced defect related process in crystalline materials. In TL, while heating the irradiated material, energy stored in the material is released with the emission of light and the intensity of the emitted light as a function of temperature forms a glow curve. The position, shape and intensity of the glow peaks are related to the properties of the trapping states responsible for TL [6]. The main applications of these materials are in radiation dosimetry for personnel and environmental monitoring, etc. Recent studies on oxide luminescent nanomaterials show that they have a potential application in the radiation dosimetry of high doses of ionizing radiation, where the conventional microcrystalline phosphors saturate [6–8]. This saturation occurs due to the ionized zones overlapping each other in the micro materials at higher doses. However, with the use of very tiny particles such as nanoscale thermoluminescent dosimetry (TLD) materials, this problem is overcome to a major extent [9].

^{*} Corresponding author. Tel.: +91 9945954010.

Corresponding author. Tel.: +91 9449804014.

E-mail addresses: eraiah@rediffmail.com (B. Eraiah), bhushanvlc@gmail.com (H. Nagabhushana), sreechakra72@yahoo.com (R.P.S. Chakradhar).

^{0025-5408/\$ -} see front matter © 2012 Elsevier Ltd. All rights reserved. http://dx.doi.org/10.1016/j.materresbull.2012.09.004

Besides the synthesis of Nd₂O₃:Co²⁺ (1–4 mol%) nanophosphors, this paper also provides a detailed analysis of phosphor using X-ray diffraction (XRD), scanning electron microscopy (SEM), transmission electron microscopy (TEM), Fourier transform infrared spectroscopy, Raman spectroscopy (FTIR), electron paramagnetic resonance (EPR), UV-vis spectroscopy and thermoluminescent techniques. The TL results of the reported nanomaterials have revealed very imperative characteristics such as simple glow peak structure. linear increase in intensity with dose, simple trap distribution. This has encouraged us to study the TL of Nd₂O₃:Co²⁺ nanophosphor using γ irradiation in the dose range 0.23-2.05 kGy.

2. Experimental

2.1. Synthesis of Nd₂O₃:Co²⁺nanophosphor

An aqueous solution containing stoichiometric amounts of analar grade neodymium nitrate (Nd(NO₃)₃; Qualigens Fine chemicals; 99%), cobalt nitrate (Co (NO₃)₂; Sigma Aldrich, 99.9%) and oxalyl dihydrazide (C₂H₆N₄O₂; ODH) fuel were dissolved in a minimum quantity of double distilled water in a cylindrical Petri dish of 150 ml capacity. ODH was used as a fuel in the combustion synthesis and it is prepared in our laboratory by the reaction of diethyl oxalate $(C_6H_{10}O_4)$ and hydrazine hydrate (H_6N_2O) as described in the literature [10]. The mixture was dispersed well using a magnetic stirrer for about 5 min. A petri dish containing the heterogeneous mixture was placed in a pre-heated muffle furnace maintained at 400 \pm 10 °C. The reaction mixture initially undergoes thermal dehydration followed by ignition with liberation of large gaseous products such as oxides of nitrogen and carbon. The flame temperature during combustion reaction was measured using an optical pyrometer which is kept inside the furnace. The flame temperature during combustion was found to be 1100 °C. The combustion reaction persists for 2-3 min yielding a highly voluminous and foamy product.

The powder X-ray diffraction studies were carried out using Philips X-ray diffractometer (model PW 3710) with CuK_a radiation $(\lambda = 1.5405 \text{ Å})$. The surface morphology of the samples was examined using scanning electron microscopy (JEOL JSM 840A) by sputtering technique with gold as covering contrast material. Transmission electron microscopy (TEM) analysis was performed on a Hitachi H-8100 (accelerating voltage up to 200 kV, LaB₆ filament) equipped with EDS (Kevex Sigma TM Quasar, USA). The UV-vis spectra were recorded on a UV-3101 Shimadzu Visible spectrometer. The EPR spectrum was recorded at room temperature using a JEOL-FE-1X EPR spectrometer operating at the X-band frequency (\approx 9.205 GHz) with a field modulation frequency of 100 kHz. The magnetic field was scanned from 0 to 500 mT and the microwave power used was 10 mW. A powdered specimen of 100 mg was taken in a quartz tube for EPR measurements. Raman spectroscopic studies were performed on Renishaw In-via Raman spectrometer with 633 nm He-Cd laser and a Leica DMLM optical microscope equipped with $50 \times$ objective, thus providing a laser spot of 2 µm in diameter. TL measurements were carried out at room temperature using Nucleonix TL reader using $Co^{60} \gamma$ -source as excitation in the dose range 0.23-2.05 kGy.

3. Results and discussion

3.1. Structural characterization

Fig. 1(a–d) shows powder X-ray diffraction (PXRD) patterns of Co^{2+} (1–4 mol%) doped Nd₂O₃ phosphor. All the powder diffraction peaks correspond to hexagonal phase with A-type of Nd₂O₃ (JCPDS No. 83-1353). Further, no secondary phase or impurity peaks was

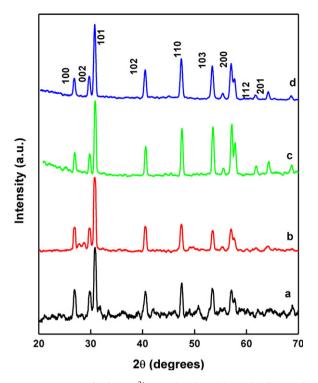


Fig. 1. PXRD patterns of Nd_2O_3 : Co^{2+} nanophosphor: (a) 1 mol%, (b) 2 mol%, (c) 3 mol%, (4) 4 mol%.

observed which confirms the phase purity of the sample. The crystallite size was estimated from the broad PXRD peaks using the Scherer's equation [11].

The crystallite size was found to be in the range 15–25 nm. A small shift in the position of main peak (1 0 1) to the lower side of 2θ values. A peak shift in XRD profiles arises due to the presence of size effect or micro strains or presence of both of them [12]. Further, the particle size of Co²⁺ (1–4 mol%) Nd₂O₃ sample was also estimated from the intense Raman line (2000 cm⁻¹) broadening using the equation [13].

$$\Gamma(\mathrm{cm}^{-1}) = 10 + \frac{124.7}{D_{\mathrm{R}}} \tag{1}$$

where Γ (cm⁻¹) is the full-width at half-maximum of the Raman active mode peak and D_R is the particle size of a sample. By substituting Γ (cm⁻¹) = 461 cm⁻¹ in the above relation, the particle size of the Nd₂O₃ sample was obtained to be 12 nm. This is in good agreement with Williamson and Hall (W–H) and Scherer's method.

Williamson and Hall suggested a method [14] combining the domain size and lattice micro strain effects on line broadening, when both are operative. The W–H approach considers the case when the domain effect and lattice deformation are both simultaneously operative and their combined effects give the final line broadening FWHM (β), which is the sum of β (grain size) and β (lattice distortion). This relation assumes a negligibly small instrumental contribution compared to the sample-dependent broadening. The strain (ε) and crystallite sizes evaluated from W–H method are listed in Table 1. It is observed that, the strain is more in 3 mol% Co doped sample and it decreases with increase of Co²⁺ (4 mol%) concentration.

Fig. 2 shows the Rietveld refinement performed on the Nd_2O_3 : Co^{2+} (3 mol%) nanopowder, which shows highest intensity. The Rietveld refinement is a method in which the profile intensities obtained from step-scanning measurements of the powders allow to estimate an approximate structural model for the real structure. In our work, the Rietveld refinement was performed through the

Table 1

Crystallite size, band gap and strain values in $Nd_2O_3{:}Co^{2\ast}\;(1{-}4\,mol\%)$ nanophosphor.

Nd ₂ O ₂ :Co ²⁺ (mol%)	Band gap (eV)	Particle size (nm)		Strain (×10 ⁻³)
	Eg	Debye–Scherrer's method (<i>d</i>)	W–H method (D)	
1	5.34	15	15	9.83
2	5.38	17	16	14.5
3	5.40	20	18	20.7
4	5.46	18	18	7.81

FULLPROF program [15]. We utilize the psedo-voigt function in order to fit the several parameters to the data point: one scale factor, one zero shifting, four back ground, three cell parameters, five shape and width of the peaks, one global thermal factors and two asymmetric factors. A typical analysis of the Nd₂O₃:Co²⁺ (3 mol%) nanopowder in Fig. 2 presents the experimental and calculated PXRD patterns obtained by the Rietveld refinement. The packing diagram of corresponding Nd₂O₃:Co²⁺ nanopowder after Rietveld refinement was shown in Fig. 3. The refined parameters such as occupancy and atomic functional positions of the Nd₂O₃ nanopowder are summarized in Table 2. The fitting parameters (R_p , R_{wp} and χ^2) indicate a good agreement between the refined and observed PXRD patterns for the Nd₂O₃:Co²⁺ nanopowder.

The surface morphology of Co^{2+} (1-4 mol%) doped Nd₂O₃ nanopowders were studied and SEM images are shown in Fig. 4. It is observed from SEM micrographs that, the crystallites have more porous, agglomeration with irregular morphology. This type of porous network with lot of voids is typical features of combustion synthesis in oxide powders due to escaping gases [16]. The porous powders are highly friable which facilitates easy grinding to obtain finer particles. It is observed from the SEM micrographs that the crystallites have no uniform shape. This is believed to be related to the non-uniform distribution of temperature and mass flow in the combustion flame [17,18]. Further the porosity increases with increase of Co²⁺ concentration. Fig. 5 shows the transmission electron micrograph of Co²⁺ (3 mol%) doped Nd₂O₃ which shows that the crystallites are fused with each other. The average particle size was observed in the range of 15-25 nm. This was also confirmed by Debye-Scherrer's equation and W-H plots. The grain size determined from W-H plots is slightly higher than those calculated using Scherrer's formula (Table 1). The small variation in the values is due to the fact that in Scherrer's formula, strain component is assumed to be zero and observed broadening of

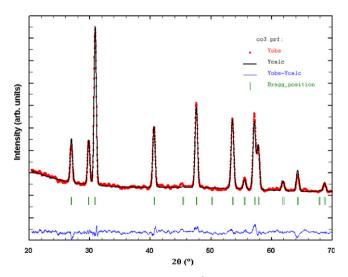


Fig. 2. Rietveld refinement of Nd₂O₃:Co²⁺ (3 mol%) nanophosphor.

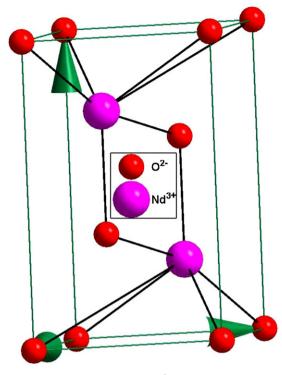


Fig. 3. Packing diagram of Nd₂O₃:Co²⁺ (3 mol%) nanophosphor.

diffraction peak is considered as a result of reducing grain size only. Further, it is observed that the strain for 3 mol% Co^{2+} doped sample is greater than that of 1, 2 and 4 mol% Co^{2+} . This increase in strain causes the broadening and shift in XRD peaks in Co doped sample.

Fig. 6 shows the FTIR spectra of Nd_2O_3 :Co²⁺ (1–4 mol%) phosphor. The peaks at 407 and 670 cm⁻¹ are the characteristic metal-oxygen (Nd-O) vibrations. A sharp peak at 3615 cm⁻¹ represents M-OH peak. The band between 3500 and 3800 cm⁻¹ is due to the (O-H) vibration of H₂O absorbed by Nd₂O₃ phosphor. The band between 1350 and 1500 cm^{-1} is assigned to the carbonate peaks and these bands are observed in all the samples. Fig. 7 shows the Raman spectra of $Nd_2O_3:Co^{2+}$ (1–4 mol%) nanophosphors recorded at room temperature with an excitation wavelength of 633 nm He-Cd laser. It is a nondestructive and potential probe for investigating the structural properties of nanomaterials [19]. Different features were observed with increase of Co²⁺ dopant namely (i) broadening of the intense Raman line, (ii) shift of the Raman peak positions and (iii) change in intensity of the Raman peaks. A strong Raman peak (Stokes line) at 2000 cm⁻¹a-1along with weak peaks (anti stokes) at 1442, 1629, 1840, 1942, 2061, 2109 cm⁻¹ was observed up to 3 mol% Co²⁺ concentration. The strong and intense peak at 2000 cm⁻¹ has been assigned to F_g mode and other weak peaks at 1442 cm⁻¹, 1629 cm⁻¹, 1840 cm⁻¹,

Table 2	2
---------	---

Rietveld refinement parameter of Nd₂O₃:Co²⁺ doped (3 mol%) nanophosphor.

Atoms	Oxidation state	Wyckoff notation	Positional parameters			Occupancy
			x	у	Ζ	
Nd_1	3	2d	0.3333	0.6667	0.2482 (2)	0.97
Со	2	2d	0.3333	0.6667	0.2482 (2)	0.03
01	2	1a	0.0000	0.0000	0.0000	1.00
02	2	2d	0.3333	0.6667	0.6474 (5)	1.00

Crystal system = Hexagonal.

Lattice parameter, a = 3.821 (2)Å, C = 5.996 (3)Å, volume = 75.82 (6)Å³.

Space group = P-3m1 (164).

 R_{factors} : R_{p} = 7.35, R_{wp} = 10.2, R_{Bragg} = 5.62, R_{F} = 4.57.

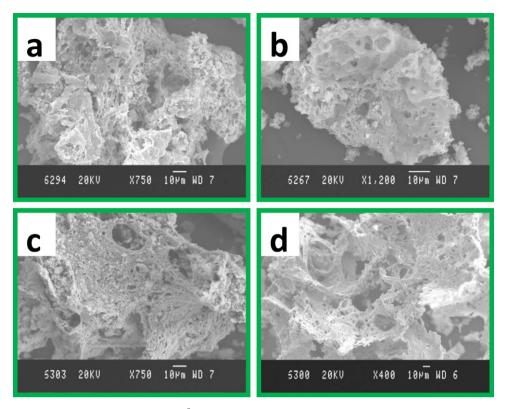


Fig. 4. SEM images of Nd₂O₃:Co²⁺ nanophosphor: (a) 1 mol%, (b) 2 mol%, (c) 3 mol%, (4) 4 mol%.

1942 cm⁻¹, 2061 cm⁻¹, 2109 cm⁻¹ are assigned to combination of F_g and E_g modes. The strong Raman intensity for 2000 cm⁻¹ band indicates a large polarizability change during the vibration. Therefore this band is expected to be more sensitive to changes in chemical bonding. It is observed that the intensity of the Stokes and anti-stoke lines decreases with increase of Co²⁺ concentration (inset of Fig. 7). Further, it is observed that, with increase of Co²⁺ concentration the intense peak (2000 cm⁻¹) shifted towards lower wavenumber and for higher concentration of Co²⁺ (4 mol%), the peaks at ~1838 cm⁻¹, 2055 cm⁻¹ and 2105 cm⁻¹ completely vanishes.

Many authors [20,21] reported that broadening of the Raman lines in rare earth oxides is due to thermal agitation. It is well known that in the case of nano systems, the Raman frequency undergoes a red shift with a decrease in the radius of the nano solid. Pan et al. [22] suggested that this red shift arises from the

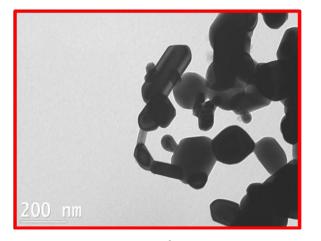


Fig. 5. TEM picture of Nd₂O₃:Co²⁺ (3 mol%) nanophosphor.

cohesive bond weakening of the lower co-ordinated atoms near the surface region of the nanograin. Another factor responsible for the shift may be the short range interactions between a phonon and photon, which describe the covalent bonding and thus are correlated to bond energy and bond length. Another factor governing this shift may be that when the size is decreased, the

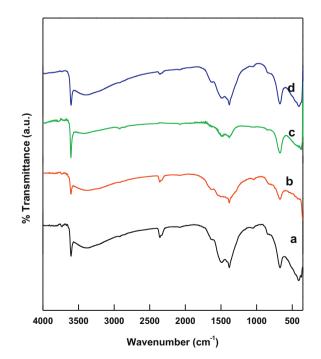


Fig. 6. FTIR spectra of $Nd_2O_3:Co^{2\ast}$ nanophosphor (a) 1 mol%, (b) 2 mol%, (c) 3 mol%, (4) 4 mol%.

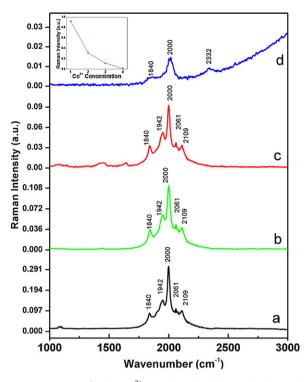


Fig. 7. Raman spectra of Nd_2O_3 : Co^{2+} nanophosphor: (a) 1 mol%, (b) 2 mol%, (c) 3 mol%, (4) 4 mol% (inset: variation of Raman peak intensity at 2000 cm⁻¹ vs Co²⁺ concentration).

momentum conservation will be relaxed and the Raman active modes will not be limited at the center of the Brillounin zone.

The UV-visible absorption spectra of different Co^{2+} (1–4 mol%) doped Nd₂O₃ phosphor are shown in Fig. 8. The spectra exhibit a broad and prominent absorption band with a maximum at ~226–240 nm along with weak absorption bands at ~292 nm. The maximum absorption arises due to transition between valence band to conduction band. The weak absorption in the UV-visible region is expected to arise from transitions involving extrinsic states such as surface traps or defect states or impurities [23]. Smaller size particles are found to have high surface to volume

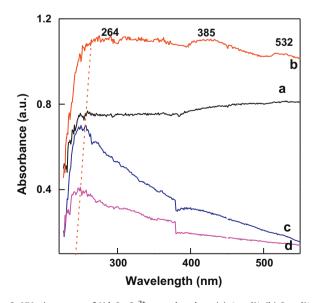


Fig. 8. UV-vis spectra of Nd_2O_3 : Co^{2+} nanophosphor: (a) 1 mol%, (b) 2 mol%, (c) 3 mol%, (4) 4 mol%.

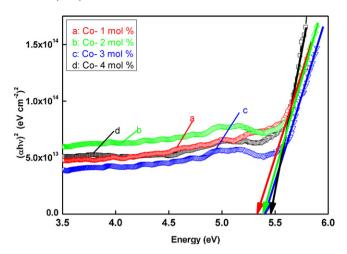


Fig. 9. Band gap of Nd_2O_3 : Co^{2+} nanophosphor: (a) 1 mol%, (b) 2 mol%, (c) 3 mol%, (4) 4 mol%.

ratio. This results in increase of defects distribution on the surface of nanomaterials. Thus the lower is the particle size, the nanomaterials exhibit strong and broad absorption bands [24].

The optical band gap energy (E_g) of Co²⁺ (1–4 mol%) doped Nd₂O₃ phosphor was estimated by Wood and Tauc relation [25] and shown in Fig. 9. The E_g values for different Co²⁺ (1–4 mol%) in Nd₂O₃ phosphor is found to be in the range 5.34–5.46 eV. The value obtained for 3 mol% Co²⁺ is well agreement to those reported in the literature [26]. Further, the E_g values are lesser for 1, 2 mol% Co²⁺ whereas for 4 mol% Co^{2+} , the E_g value is enhanced. The variation in band gap values in 1, 2, and 4 mol% Co²⁺ can be related to the degree of structural order-disorder in the lattice which is able to change the intermediately energy level distribution within the band gap. This is attributed to particle size effect. The nanomaterials have large surface to volume ratio as a result the formation of voids on the surface as well as inside the agglomerated particles. Such voids can cause fundamental absorption in the UV wavelength range [27]. When 3 mol% Co²⁺ doped into Nd₂O₃ matrix it becomes more ordered structure with less defects. Consequently intermediate energy levels (deep and shallow holes) are minimized within the optical band gap. Further, the E_{g} values depends on various factor the preparation methods and different experimental conditions. In particular these key factors can favor or inhibit the formation of structural defects, which are able to control the degree of structural order-disorder of the material and consequently the number of intermediately energy levels within the band gap.

The EPR spectrum of Nd_2O_3 :Co²⁺ (4 mol%) nanopowder at room temperature is shown in Fig. 10. The EPR spectrum of Co²⁺ ions in general is observed only at low temperature because the spin lattice relaxation time is extremely short for Co²⁺ ions. At higher temperature the spectrum become broader probably due to short relaxation time characteristic of high spin state of the Co²⁺ ion [28]. The Co²⁺ with d⁷ electron configuration and ⁴F ground state has seven fold orbital degeneracy. The ⁴F state of d⁷ ions in octahedral field splits to ⁴T triplet ground state. Spin-orbital interaction leads to splitting of this triplet so that the lowest state is Kramers doublet with isotropic g = 4.3. The fields of lower symmetry mix this state with nearly lying excited states, and g-value becomes anisotropic [29]. Therefore the g values become anisotropic and are sensitive to variations in crystal field [30]. While Co²⁺ ions at the tetragonal sites in the lattice should produce a single EPR line with the Lorentzian line shape with a g-factor of ~ 2.25 [31]. In the present work, the EPR spectrum of Co²⁺ ions in Nd₂O₃ nanopowder exhibit resonance signals with anisotropic g values at g = 2.25 and

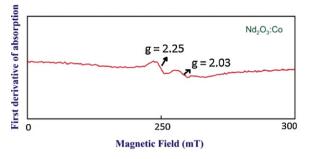


Fig. 10. EPR spectrum of Nd₂O₃:Co²⁺ (4 mol%) nanopowder at room temperature.

g = 2.03. These two resonance signals are attributed to Co²⁺ ions at the cubic sites of Nd₂O₃ lattice.

TL measurements probe a class of defects with energy levels in the band gap of insulating/semiconducting materials. Fig. 11 shows the TL glow curves of different concentration of Co^{2+} (1-4 mol%) in Nd₂O₃ host irradiated with 1.14 kGy. Here all the measurements were carried out at room temperature at a constant heating rate 5 °C s⁻¹. The TL intensity is found to be highest for 3 mol% of Co^{2+} . Fig. 12 shows the TL spectra for optimized concentration of Co^{2+} ions (3 mol%) in Nd₂O₃ nanophosphor irradiated with γ -rays in the dose range 0.23–2.05 kGy. Well resolved glow peaks at 430 K along with shoulder peaks in the range 150–250 K were observed.

The variation of TL glow peak intensity (430 K) with accumulated dose was studied and is shown in Fig. 13. It is observed that the TL intensity increases linearly with increasing γ -dose. This linear behavior of the nanophosphor is highly useful for dosimetric application. Further, it is noticed that there is slight shifting in the glow peak positions to lower temperature side with increase of γ dose. The appearance of more than one TL peak indicates that there are possibly different kinds of trapping centers generated due to γ irradiation. The shallow trapping center leads to the resolved peak at lower temperature and the other deeper center gives rise to shouldered peak at higher temperature. The intensity of the higher temperature glow peak increases linearly with dose. This might be due to high surface to volume ratio, which results in a higher surface barrier energy for the nanoparticles. On increasing the dose, the energy density crosses the barrier and a large number of defects are produced in the nanoparticles which ultimately keep on increasing with the dose till saturation is achieved [32].

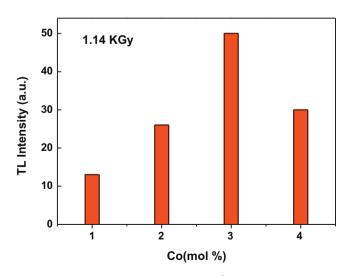


Fig. 11. Variation of TL intensity with of $Nd_2O_3{:}Co^{2+}{:}$ (a) 1 mol%, (b) 2 mol%, (c) 3 mol%, (d) 4 mol%.

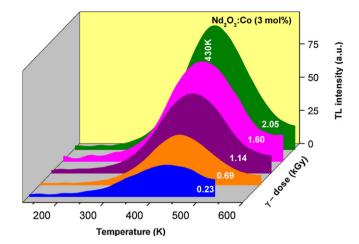


Fig. 12. TL glow curves of Nd_2O_3 : Co^{2+} (3 mol%) γ -irradiated for: (a) 0.23 Gy, (b) 0.69 Gy, (c) 1.14 kGy, (d) 1.6 kGy, (e) 2.05 kGy.

It is also observed that for lower dosed samples the traps/ defects generated appear to be very less due to less surface barrier energy. Hence the glow curve has minimum TL intensity. As the temperature increases, the surface barrier energy also increases. The increase in TL intensity with dose can be explained on the basis of track interaction model [33,34]. According to this model, the number of traps generated by the high energy radiation in a track depends upon the cross section and the length of the track inside the matrix. In the case of nanomaterials, the length of the track generated by high energy radiation is of few tenths of nanometers. At low doses, there exist a few trap centers or luminescent centers owing to the small size of the particles. As the dose increases, the TL intensity increases as the cross section would increase with increase in γ -dose.

In our earlier work [35], Nd₂O₃:Eu³⁺ nanophosphor γ irradiated for 100–400 Gy show well resolved glow peak at ~426 K at a warming rate of 5 °C s^{-1.} The TL intensity of the sample follows a linear relation with increasing dose. Strong fading was recorded initially after four days; the decay is quite slow and finally stabilizes after 21 days. Soliman [36] observed three glow peaks in the range 310–385 °C in as received Nd₂O₃ powder procured from sigma Aldrich company. However, in heat treated samples show four TL glow curves in the range 170–386 °C at a dose rate of

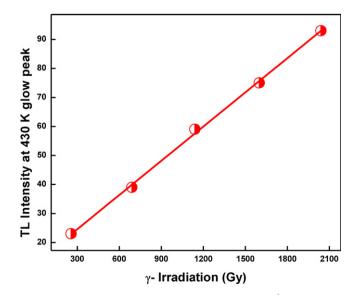


Fig. 13. Variation of TL intensity with γ dose in Nd₂O₃:Co²⁺ (3 mol%).

Table 3	
Vin atia manana atana	of Ma

Kinetic parameters of Nd ₂ O ₃ :Co ²⁺ ((1–4 mol%) estimated using Chen's glow peak shape method.
--	---

γ-Dose (kGy)	Peak	<i>T</i> _m (K)	Order of kinetics $b(\mu_g)$	Activation energy (eV)				Frequency factor (s ⁻¹)
				E_{τ}	E_{δ}	E_{ω}	Eave	
2.05	1	334	2 (0.51)	0.738	0.829	0.787	0.785	2.1E+07
	2	400	2 (0.50)	1.010	1.101	1.060	1.057	5.8E+08
	3	444	2 (0.50)	1.121	1.210	1.170	1.167	1.1E+09
	4	501	2 (0.49)	1.207	1.309	1.263	1.259	1.0E+09
1.6	1	320	2 (0.52)	0.637	0.732	0.688	0.686	3.8E+06
	2	369	2 (0.51)	1.353	1.388	1.377	1.372	6.2E+11
	3	433	2 (0.50)	0.997	1.096	1.051	1.048	1.9E+08
	4	507	2 (0.51)	1.651	1.679	1.673	1.667	5.1E+11
1.14	1	337	2 (0.52)	0.388	0.520	0.457	0.455	1.8E+04
	2	417	2 (0.51)	0.852	0.958	0.908	0.906	2.3E+07
	3	482	2 (0.49)	1.712	1.761	1.746	1.740	4.0E+12
	4	558	2 (0.49)	1.088	1.221	1.157	1.155	5.0E+07
0.69	1	325	2 (0.49)	0.649	0.757	0.704	0.703	4.8E+06
	2	395	2 (0.49)	0.944	0.928	0.882	0.918	5.2E+07
	3	439	2 (0.51)	1.322	1.126	1.108	1.185	1.7E+09
	4	496	2 (0.50)	1.492	1.159	1.134	1.262	1.2E+09
0.23	1	331	2 (0.48)	0.920	0.590	0.550	0.687	2.0E+06
	2	408	2 (0.52)	1.140	0.601	0.582	0.774	2.2E+06
	3	414	2 (0.49)	0.984	0.375	0.329	0.563	4.3E+04
	4	482	2 (0.49)	1.835	0.969	0.958	1.254	1.6E+09

600 Gy.The variation in TL glow peak temperature and the number of glow peaks in present study might be due to particle size effect, method of preparation of the phosphor, heating rate, type of ionizing radiation [37].

The dosimetric properties of TL materials depend on the kinetic parameters of its glow peak. Kinetic parameters give valuable information about mechanism responsible for the emission in material. Reliable dosimetric studies of any TL material include a good knowledge of its kinetic parameters well. The important kinetic parameters are trap depth or activation energy (E) which is the thermal energy required to liberate the trapped electrons and holes, frequency factor (s) and the order of kinetics (b). These parameters can be estimated using Chen's set of empirical equations [38].

Further, activation energy (*E*) and frequency factor can also be estimated using Chen's peak shape method [38]. The trapping parameters (*E*, *b*, *s*) estimated for different γ -doses are given in Table 3.

4. Conclusions

Different concentrations of Co²⁺ doped Nd₂O₃ nanophosphor have been synthesized by solution combustion method for the first time using oxalyl dihydrazide (C₂H₆N₄O₂; ODH) as a fuel. XRD patterns confirm that hexagonal (A-type) phase without impurity peaks. TEM and XRD studies showed that crystallite size were nanocrystalline in nature. Rietveld refinement confirmed the space group of the structure to be P-3m1 (164) with lattice parameter a = 3.821 (2) Å, c = 5.996 (3) Å and cell volume (V) = 75.82 (6) Å³. The electronic band gap of Nd₂O₃:Co²⁺ (1–4 mol%) estimated from UV–Vis spectroscopy is in the range 5.34–5.46 eV. Two resonance signals with g-values at 2.25 and 2.03 and are attributed to Co²⁺ ions. A well resolved TL glow peak at 430 K along with shouldered peaks in the range 150-250 K was observed at a warming rate of 5 °C s $^{-1}$. TL intensity of the glow peak at 430 K increases linearly with the increase of γ dose. Simple glow peak shape, simple trap distribution, linear increase in TL intensity with increase of γ dose which make Nd₂O₃:Co²⁺ (3 mol%) nanophosphor might be used in radiation dosimetry.

Acknowledgments

HN thanks to DST Nano mission for sanction of the project No. SR/NM/NS-48/2010. Prof. J.L. Rao thanks University Grants Commission (New Delhi) for the award of Emeritus Fellowship. B. Umesh thanks to the management of P.V.P.WT for encouraging higher studies and support.

References

- L.V. Interrante, M.J. Hampden-Smith, Chemistry of Advanced Materials: An Overview, Wiley VCH, New York, 1998.
- [2] X. Qu, J. Dai, J. Tian, X. Huang, Z. Liu, Z. Shen, P. Wang, J. Alloys Compd. 469 (2009) 332–335.
- [3] L. Qian, Y. Gui, S. Guo, Q. Gong, X. Qian, J. Phys. Chem. Solids 70 (2009) 688–693.
 [4] R. Bazzi, M.A. Flores-Gonzalez, C. Louis, K. Lebbou, C. Dujardin, A. Brenier, W.
- Zhang, O. Tillement, E. Bernstein, J. Lumin. 445 (2003) 102–103.
- [5] A.G. Dedov, A.S. Loktev, I.I. Moiseev, A. Aboukais, J.F. Lamonier, I.N. Filimonov, Appl. Catal. A 245 (2003) 209–220.
- [6] J. Azorin, Nucl. Tracks Rad. Meas. 11 (1986) 159-166.
- [7] G. Kitis, C. Furetta, M. Prokic, V. Prokic, J. Phys. D: Appl. Phys. 33 (2000) 1252– 1262.
- [8] P.D. Sahare, R. Ranjan, N. Salah, S.P. Lochab, J. Phys. D 40 (2007) 759-764.
- [9] S.P. Lochab, P.D. Sahare, R.S. Chauhan, N. Salah, A. Pandey, J. Phys. D 40 (2007) 1343–1350.
- [10] B.M. Nagabhushana, Ph.D. Thesis, Bangalore University, 2008.
- [11] P. Klug, L.E. Alexander, X-ray Diffraction Procedure, Wiley, New York, 1954.
- [12] S. Ekambaram, Y. Iikubo, A. Kudo, J. Alloys Compd. 433 (2006) 237–240.
- [13] Z. Wang, Z. Quan, J. Lin, Inorg. Chem. 46 (13) (2007) 5237–5242.
- [14] G.K. William, W.H. Hall, Acta Metall. 1 (1953) 22-31.
- [15] R.W.G. Wyckoff, Crystal Structures, vol. 2, Interscience, New York, 1964, pp. 4–5.
- [16] A. Jagannatha Reddy, M.K. Kokila, H. Nagabhushana, R.P.S. Chakradhar, C. Shiva-
- kumara, J.L. Rao, B.M. Nagabhushana, J. Alloys Compd. 509 (2011) 5349-5355. [17] B. Nagappa, G.T. Chandrappa, Microporous Mesoporous Mater. 106 (2007) 212-
- 218.[18] K.C. Patil, M.S. Hegde, S.T. Aruna, T. Rattan, Chemistry of Nanocrystalline Oxide Materials, Combustion Synthesis, Properties and Applications, World Scientific
- Publishing Co. Pvt. Ltd., UK, 2008. [19] A. Jagannatha Reddy, M.K. Kokila, H. Nagabhushana, J.L. Rao, C. Shivakumara, B.M.
- Nagabhushana, R.P.S. Chakradhar, Spectrochim. Acta A 81 (2011) 53–58.
- [20] W.B. White, V.G. Keramidas, Spectrochim. Acta A 28 (1972) 501–509.
 [21] W. Urben Marek, C. Cornilsen Bahn, J. Phys. Chem. Solids 48 (1984) 475–480.
- [22] L.K. Pan, C.Q. Sun, C.M. Li, J. Phys. Chem. B 108 (2004) 3404–3406.
- [23] H.Q. Cao, X.Q. Qiu, B. Luo, Y. Liang, Y.H. Zhang, R.Q. Tan, M.J. Zhao, Q.M. Zhu, Adv. Funct. Mater. 14 (2004) 243–246.
- [24] A. Emeline, G.V. Kataeva, A.S. Litke, A.V. Rudakova, V.K. Ryabchuk, N. Serpone, Langmuir 14 (1998) 5011–5022.
- [25] J. Tauc, in: F. Abeles (Ed.), Optical Properties of Solids, North-Holland, Amsterdam, 1970.

- [26] B. Umesh, B. Eraiah, H. Nagabhushanac, B.M. Nagabhushana, G. Nagaraja, C. Shivakumara, R.P.S. Chakradhar, J. Alloys Compd. 509 (2011) 1146-1151.
- [27] N. Salah, S.S. Habib, Z.H. Khan, S.A. Hamedi, S.P. Lochab, J. Lumin. 129 (2009) 192– 196.
- [28] F. Gan, H. Deng, H. Liu, J. Non-Cryst. Solids 52 (1982) 143-149.
- [29] S. Ruck, D. Stachel, Phosphorus Res. Bull. 13 (2002) 201-204.
- [30] A. Abragam, B. Bleaney, Electron paramagnetic resonance of transition ions, Oxford University Press, 1970.
- [31] S. Angelov, E. Zhecheva, R. Stoyanova, M. Atanasov, J. Phys. Chem. Solids 51 (1990) 1157-2116.
- [32] M. Chandrasekhar, D.V. Sunitha, N. Dhananjaya, H. Nagabhushana, S.C. Sharma, B.M. Nagabhushana, C. Shivakumara, R.P.S. Chakradhar, J. Lumin. 132 (2012) 1798–1806.
- [33] S. Mahajna, Y.S. Horowitz, J. Phys. D: Appl. Phys. 30 (1997) 2603–2619.
 [34] Y.S. Horowitz, O. Avila, M. Rodrigues-Villafuerte, Nucl. Instrum. Methods Phys.
- Res. B 184 (2001) 85–112.
 [35] B. Umesh, B. Eraiah, B.M. Nagabhushana, H. Nagabhushana, H.B. Premkumar, C.
- [35] B. Ohresh, B. Elalah, Shu. Nagabhushana, H. Nagabhushana, H. Nagabhushana, N. Shiwakumara, C. Shivakumara, C.H. Rayappa, R.P.S. Chakradhar, Phil. Mag. Lett. 89 (2009) 589–597.
 [36] C. Soliman, Nucl. Instrum. Methods B 251 (2006) 441–444.
- [37] H. Nagabhushana, Ph.D. Thesis, Bangalore University, 2003.
- [38] R. Chen, S.W.S. Mckeever, Theory of Thermoluminescence and Related Phenomena, World Scientific Press, Singapore, 1997.