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Investigation of deep level defects in copper irradiated bipolar junction transistor

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1. Introduction

Bipolar junction transistors (BJTs) are still being extensively used in space and other radiation rich environments. These devices are sensitive to high energy particle irradiation. Considerable amount of data are available on the radiation effects of γ -rays, fast neutrons, electrons and protons on semiconductor devices [1–5]. However, there appears to be rather little work on heavy ion induced effects and consequent characterization of defects by DLTS technique. Deep level defect monitoring plays an important role in designing a semiconductor device suitably for various applications. Thus a study of radiation induced effects in semiconductor devices is important to observe changes in electrical characteristics and to get basic information regarding the generation and annihilation of defects.

Exposure of semiconductor devices to high energy particle radiation is known to generate variety of defects. The nature of these defects generated by irradiation process depends on the properties

ABSTRACT

Commercial bipolar junction transistor (2N 2219A, npn) irradiated with 150 MeV Cu¹¹⁺-ions with fluence of the order 10^{12} ions cm⁻², is studied for radiation induced gain degradation and deep level defects. *I–V* measurements are made to study the gain degradation as a function of ion fluence. The properties such as activation energy, trap concentration and capture cross-section of deep levels are studied by deep level transient spectroscopy (DLTS). Minority carrier trap levels with energies ranging from $E_C - 0.164$ eV to $E_C - 0.695$ eV are observed in the base–collector junction of the transistor. Majority carrier trap levels are also observed with energies ranging from $E_V + 0.203$ eV to $E_V + 0.526$ eV. The irradiated transistor is subjected to isothermal and isochronal annealing. The defects are seen to anneal above 350 °C. The defects generated in the base region of the transistor by displacement damage appear to be responsible for transistor gain degradation.

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of target as well as impinging high energy particle. To investigate these deep level defects several techniques are in practice. Deep level transient spectroscopy (DLTS) is now an established technique for detecting and characterizing variety of defects in semiconductor devices. DLTS is a high frequency capacitance transient thermal scanning method useful in observing a wide variety of traps in semiconductor devices [6].

The BJT used in the present study has been thoroughly studied in our earlier work for 24 MeV protons, 8 MeV electrons and ⁶⁰Co γ -rays induced effects [7–9]. A DLTS study of deep level defects in Li-ion irradiated transistor (chosen from the same batch) is also reported earlier [10].

2. Experimental details

Commercial BJT (2N 2219A, npn) manufactured in an indigenous technology from Continental Device India Ltd. (CDIL) has been selected for the present study. This device is a switching transistor with standard configuration (base thickness is 2.0 μ m and oxide thickness is 1.2 μ m) suitable for low and high frequency

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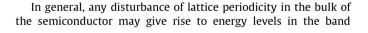
operation. Decapped transistor is exposed to 150 MeV Cu¹¹⁺-ions using 15 UD 16 MV pelletron Tandem Van de Graff accelerator facility at Inter-University Accelerator Centre, New Delhi. The transistor in the biased condition (CE mode) is irradiated by Cu¹¹⁺-ions with three different fluences 3×10^{11} ions cm⁻², 1×10^{12} ions cm⁻² and 1×10^{13} ions cm⁻². During irradiation, the target chamber is maintained at room temperature (300 K) and low pressure (7.5×10^{-9} Torr). The ion fluences are calculated by measuring the ion beam current and irradiation time. The ion beam current is fixed at ~ 1 particle nano ampere (pnA). Output characteristics of the transistor are studied at a constant base current ($I_{\rm B}$) of 50 μ A. The collector voltage ($V_{\rm CE}$) is varied from -0.1 V to 1.5 V in steps of 0.01 V. Gummel plots are acquired by varying base emitter voltage ($V_{\rm BE}$) from 0 V to 0.7 V in steps of 0.01 V at constant collector voltage (V_{CF}) of 5 V.

DLTS spectra are recorded for both unirradiated transistor and three different transistors of the same batch (date code) exposed to Cu¹¹⁺-ion for different fluences. The DLTS system (IMS-2000, M/s. Lab Equip, India) employed for the present study consists of a boxcar averager, a pulse generator, a thousand point digitizer, a voltage generator and a high speed capacitance meter. The pulse generator is capable of generating pulses of widths ranging from 100 ns to 10 s. The pulse height could be programmed from – 12 V to +12 V. The boxcar averager is capable of generating seven rate windows. The time constants can be varied from 1 ms to 2 s. In the present study, DLTS spectra are recorded with a reverse bias of 5 V and pulse width of 19.2 ms applied to base–collector junction. The trap concentration, activation energy and capture crosssection of different deep levels are determined by DLTS spectra.

The transistors are subjected to isothermal and isochronal (30 min) annealing. The annealing temperature in the furnace can be maintained constant for several hours with an accuracy of 1 °C. During isochronal (30 min) annealing, the temperature is varied from 100 °C to 500 °C. During isothermal (100 °C) annealing, the annealing time is varied from 30 min to 480 min. DLTS spectra are recorded at different stages of thermal annealing and the characteristics of several deep level defects are monitored.

3. Results and discussion

3.1. I-V measurements



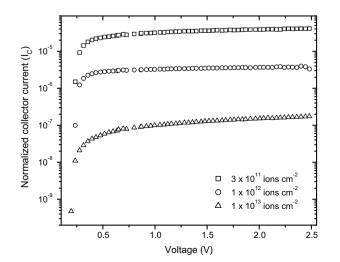


Fig. 1. Collector characteristics of Cu-ion irradiated transistor for three different ion fluences (at constant $I_B = 50 \mu$ A).

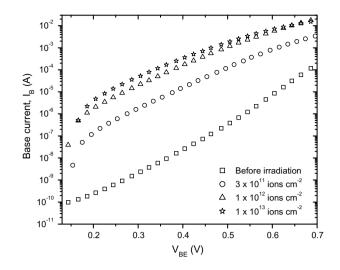


Fig. 2. Variation of $I_{\rm B}$ with $V_{\rm BE}$ for three different ion fluences (at constant $V_{\rm CE}$ = 5 V).

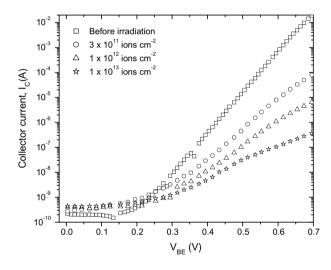


Fig. 3. Variation of $I_{\rm C}$ with $V_{\rm BE}$ for three different ion fluences (at constant $V_{\rm CE}$ = 5 V).

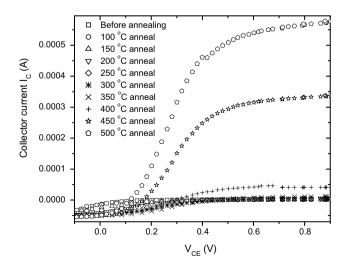


Fig. 4. Collector characteristics of the Cu-ion irradiated transistor at different isochronal annealing temperatures (at constant $I_B = 50 \mu$ A).

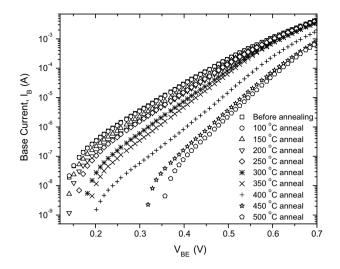


Fig. 5. Variation of $I_{\rm B}$ with $V_{\rm BE}$ at different isochronal annealing temperatures (at constant $V_{\rm CE}$ = 5 V).

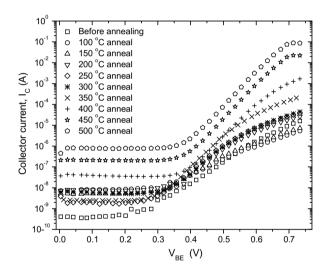


Fig. 6. Variation of $I_{\rm C}$ with $V_{\rm BE}$ at different isochronal annealing temperatures (at constant $V_{\rm CE}$ = 5 V).

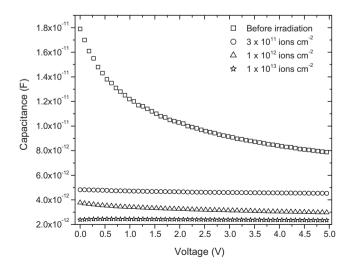


Fig. 7. C-V characteristics of the Cu-ion irradiated transistor for three different fluences.

gap. Radiation induced defects may have such energy levels with them and these defects can have a major impact on the electrical characteristics of the transistor [11]. Fig. 1 exhibits the collector characteristics of the Cu-ion irradiated transistor for three different fluences. This plot indicates that the collector current decreases with an increase in the ion fluence. Figs. 2 and 3 exhibit Gummel plots (variation of base current $I_{\rm B}$, collector current $I_{\rm C}$ with base

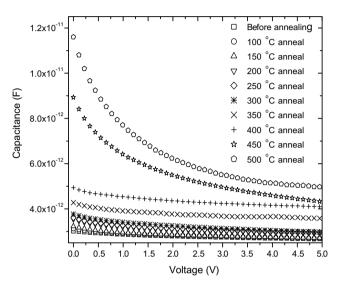


Fig. 8. *C*–*V* characteristics of the Cu-ion irradiated transistor at different isochronal annealing temperatures.

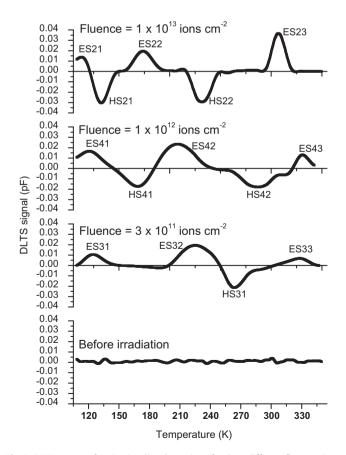


Fig. 9. DLTS spectra of Cu-ion irradiated transistor for three different fluences. Rate window is fixed at 127.9 s^{-1} .

voltage V_{BE}) before and after irradiation. It is observed that the base current increases with ion fluence. However, the collector current decreases as the ion fluence increases. This increase in base current and decrease in collector current upon irradiation could be attributed to the increased recombination of charge carriers in the base– collector junction of the transistor.

One important aspect of characterization of BJTs for radiation induced effects is radiation induced gain degradation. The BJTs are found to be particularly vulnerable to ionizing radiation and the transistor gain degradation is the primary cause for parametric shifts and functional failures. The degradation of forward current gain of BJT when exposed to radiation is dependent largely on nature and energy of radiation and to some extent on the dose rate. It is well known that the degradation of the transistor can occur due to increased recombination in the base region due to displacement damage caused upon irradiation. The increased recombination in the neutral base leads to an increase in the base current which in turn results in a decrease in the collector current. When recombination centers are generated in the base region of the transistor, it leads to an increase in the base current by decreasing the minority carrier lifetime [12]. A decrease in the minority carrier lifetime will be reflected in the degradation of forward current gain of the transistor.

Fig. 4 shows the collector characteristics of the transistor irradiated with 1×10^{12} ions cm⁻² at different isochronal (30 min) annealing temperatures. The isochronal annealing clearly shows that there is a major recovery in the collector current above 400 °C. Figs. 5 and 6 exhibit the variation of I_B and I_C as a function of V_{BE} at different isochronal (30 min) annealing temperatures. It is found that while collector current increases, the base current decreases with an increase in annealing temperature.

3.2. C-V measurements

Fig. 7 shows the capacitance–voltage (*C–V*) characteristics of the base–collector junction of the transistor before and after irradiation. The plot shows that there is a considerable degradation in the *C–V* characteristics of the transistor after irradiation. This would indicate that there is a partial loss of charge carriers in the base–collector junction of the transistor upon irradiation. This loss could be attributed to an increased recombination in the base–collector junction of the irradiated transistor. Fig. 8 shows the *C–V* characteristics of the base–collector junction of the transistor. Fig. 8 shows the *C–V* characteristics of the base–collector junction of the transistor irradiated with 1×10^{12} ions cm⁻² at different isochronal (30 min) annealing temperatures. The *C–V* characteristics appear to improve after 350 °C temperature.

3.3. DLTS measurements

In principle, irradiation by any ion can damage transistors through both ionization and displacement. Displacement damage is a bulk phenomenon which results in the generation of several

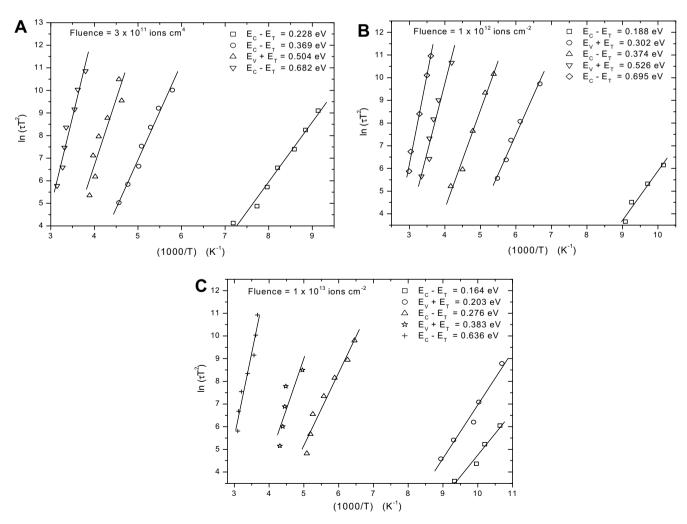


Fig. 10. Arrhenius plots for Cu-ion fluences: (A) 3×10^{11} ions cm⁻²; (B) 1×10^{12} ions cm⁻²; (C) 1×10^{13} ions cm⁻².

types of defects such as vacancy, interstitial, di-vacancy, Frenkel pair, vacancy–impurity complexes namely A-center (V–O), E-center (V–P), interstitial clusters and di-interstitial or higher order complexes called D-center [11].

The deep level defects generated by irradiation of transistors by Cu-ions are characterized using DLTS technique. Fig. 9 exhibits the DLTS spectra of Cu-ion irradiated transistor for three different ion fluences. Fig. 10 exhibits the Arrhenius plots of deep level defects for three different ion fluences. Three minority carrier and one majority carrier deep level defects are observed in the DLTS spectra of the Cu-ion irradiated transistor with fluence of 3×10^{11} ions cm⁻². For the ion fluences 1×10^{12} and 1×10^{13} ions cm⁻², three minority carrier and two majority carrier defect levels are observed. The trap concentration, capture cross-section and introduction rate of all the deep level defects are calculated from the DLTS spectra and presented in Table 1.

Recombination of electron-hole pairs at the defect levels generated upon displacement damage is the most important physical phenomenon responsible for the gain degradation. Mainly four kinds of recombination processes are observed in semiconductor devices: (i) Shockley Read Hall (SRH) or multi-phonon recombination; (ii) radiative recombination; (iii) Auger recombination and (iv) non-radiative recombination. Radiative recombination is important in direct band gap semiconductors like GaAs. Auger recombination is observed in either direct or indirect band gap semiconductors when the carrier concentration is high. Further, the radiative, non-radiative and Auger recombination lifetimes are independent of trap concentration [13]. SRH recombination is particularly important in indirect band gap semiconductors such as Si [13–17]. The values of SRH recombination lifetimes calculated using DLTS data are tabulated in Table 1.

The transistor irradiated with fluence 1×10^{12} ions cm⁻² is subjected to both isothermal and isochronal annealing and the characteristics of each defect are monitored by recording DLTS spectra. However, no significant changes are observed in the DLTS spectra as a function of annealing time at a given constant temperature (isothermal annealing). After isochronal annealing, significant changes are observed in the DLTS spectra. The total defect concentration decreases and the effective recombination life time increases with an increase in isochronal (30 min) annealing temperature as shown in Fig. 11. Fig. 12 exhibits the variation of concentration of different deep level defects as a function of isochronal (30 min) annealing temperature. The annealing behavior of different deep level defects in silicon has been studied by several

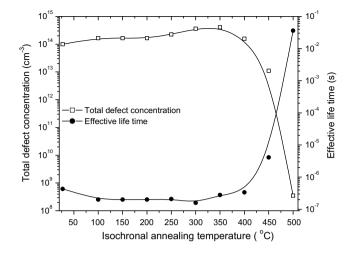


Fig. 11. Variation of total defect concentration and effective recombination life time as a function of isochronal annealing temperature.

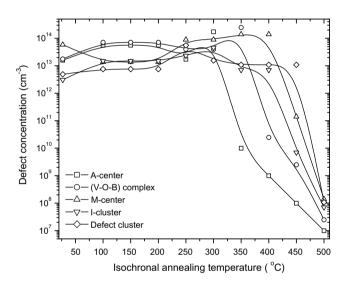


Fig. 12. Variation of defect concentration as a function of isochronal annealing temperature.

Table 1
Data obtained from DLTS analysis of the Cu-ion irradiated transistor

Ion fluence (ions cm ⁻²)	Defect label	Activation energy (eV)	Trap concentration (cm ⁻³)	Capture cross-section $(7 \times 10^{-18} \text{ cm}^2)$	Introduction rate η (cm ⁻¹)	Recombination life time (s)	Identified defect type	Literature comparison	
								Defect position	Reference
3×10^{11}	ES31 ES32 HS31 ES33	$E_{\rm C} - 0.228 \pm 0.018$ $E_{\rm C} - 0.369 \pm 0.019$ $E_{\rm V} + 0.504 \pm 0.044$ $E_{\rm C} - 0.682 \pm 0.030$	$\begin{array}{c} 4.5\times10^{12}\\ 2.0\times10^{13}\\ 9.3\times10^{12}\\ 2.2\times10^{12} \end{array}$	$\begin{array}{c} 2.3\times10^{-15}\\ 1.1\times10^{-15}\\ 9.8\times10^{-15}\\ 8.8\times10^{-14} \end{array}$	15.1 65.8 31.0 7.3	$\begin{array}{l} 1.3\times 10^{-05}\\ 4.8\times 10^{-06}\\ 1.1\times 10^{-06}\\ 4.4\times 10^{-07}\end{array}$	Di-vacancy (C _i –O _i) I-cluster Cluster	$E_{\rm C} - 0.23$ $E_{\rm C} - 0.36$ $E_{\rm V} + 0.50$	[21] [21] [25]
1 × 10 ¹²	ES33 ES41 HS41 ES42 HS42 ES43	$E_{\rm C} = 0.082 \pm 0.030$ $E_{\rm C} = 0.188 \pm 0.031$ $E_{\rm V} + 0.302 \pm 0.018$ $E_{\rm C} = 0.340 \pm 0.020$ $E_{\rm V} + 0.526 \pm 0.029$ $E_{\rm C} = 0.695 \pm 0.017$	$\begin{array}{c} 2.2 \times 10 \\ 1.6 \times 10^{13} \\ 1.8 \times 10^{13} \\ 5.8 \times 10^{13} \\ 3.1 \times 10^{12} \\ 4.9 \times 10^{12} \end{array}$	$\begin{array}{l} 4.6 \times 10^{-15} \\ 4.2 \times 10^{-16} \\ 2.8 \times 10^{-16} \\ 1.3 \times 10^{-15} \\ 2.7 \times 10^{-14} \end{array}$	7.3 15.9 17.6 57.8 3.1 4.9	$\begin{array}{c} 4.4 \times 10 \\ 2.0 \times 10^{-06} \\ 1.6 \times 10^{-05} \\ 6.5 \times 10^{-06} \\ 2.2 \times 10^{-05} \\ 6.4 \times 10^{-07} \end{array}$	A-center (V–O–B) M-center I-cluster Cluster	$ E_{C} - 0.18 E_{V} + 0.30 E_{C} - 0.33 E_{V} + 0.50 - - $	- [26] [27] [28] [25] -
1×10^{13}	ES21 HS21 ES22 HS22 ES23	$\begin{split} E_{\rm C} &- 0.164 \pm 0.030 \\ E_{\rm V} + 0.203 \pm 0.026 \\ E_{\rm C} &- 0.276 \pm 0.023 \\ E_{\rm V} + 0.383 \pm 0.071 \\ E_{\rm C} &- 0.636 \pm 0.030 \end{split}$	$\begin{array}{c} 1.6\times10^{13}\\ 1.0\times10^{14}\\ 2.4\times10^{12}\\ 1.7\times10^{13}\\ 3.0\times10^{14} \end{array}$	$\begin{array}{l} 8.5 \times 10^{-16} \\ 9.0 \times 10^{-15} \\ 3.6 \times 10^{-17} \\ 3.0 \times 10^{-16} \\ 9.0 \times 10^{-15} \end{array}$	1.6 10.1 0.2 1.7 30.4	$\begin{array}{c} 1.2\times10^{-05}\\ 1.6\times10^{-07}\\ 1.6\times10^{-03}\\ 2.0\times10^{-05}\\ 3.2\times10^{-08} \end{array}$	A-center Di-vacancy (B _i -O _i) K-center Cluster	$E_{\rm C} - 0.17$ $E_{\rm V} + 0.21$ $E_{\rm C} - 0.27$ $E_{\rm V} + 0.38$	[25] [25] [27] [27]

groups and documented well in the literature [18–24]. The identification of the defect type is made on the basis of their finger prints such as activation energy, annealing temperature and capture cross section. Table 1 shows the assignment of defect types with activation energy and the comparison with the reported values in the literature. The error margins in the activation energy and capture cross-section are also shown.

The defect levels $E_{\rm C} - 0.682$, $E_{\rm C} - 0.695$ and $E_{\rm C} - 0.636$ anneal at 500 °C. By comparing the measured properties of these levels with those reported in the literature, these levels could be attributed to new energy levels of defect clusters in silicon.

A comparison of *I*–*V* characteristics (Figs. 4–6), *C*–*V* characteristics (Fig. 8) and DLTS results (Figs. 11 and 12) before and after annealing suggests that the defects which annihilate in the temperature range 350 °C to 500 °C have a major contribution to the change in the base and collector currents and hence degradation of the forward current gain of the transistor.

One other important observation in the present investigation is that the nature of the defects varies with increasing radiation dose. The different type of defects that appear after irradiation with different fluence is not only interesting but is rather unusual. Similar variations have been observed in our earlier studies and are also reported earlier by others [29]. On the contrary, it is well reported in the literature that A-center (O–V) is the first defect that appears in most Si structures after irradiation [30,31]. In the present study, we have not observed A-center for a fluence of 3×10^{11} ions cm⁻². However, for higher fluences, this defect type appears in the DLTS spectra. One possible reason for this observation is that irradiation at higher fluence may result in saturation of defects and defect reordering which occurs due to the interaction between the existing defects and newly formed defects. The nature of defect could also change due to concentration of species. For example, either vacancy combination or di-vacancy combination might be favored depending upon the concentration of combining species. A gradual variation in the particle fluence and subsequent monitoring of defect levels would be useful in the identification and assignment of defect types.

DLTS monitors the thermally activated emissions, unlike in optical spectroscopy. Hence, the assignment of DLTS peaks to exact nature and type of defects could be difficult. To identify the exact microscopic structure of these defect levels, investigation by other structure sensitive techniques such as electron paramagnetic resonance (EPR) and infrared photoconductivity measurements would be required.

4. Conclusions

Commercial bipolar junction transistor (2N 2219A, npn) undergoes gain degradation upon irradiation by Cu¹¹⁺-ions. The gain degradation can be attributed to displacement damage in the base region of the transistor. Deep level defects generated in the collector-base junction of the transistor are identified as electron traps with activation energies ranging from $E_{\rm C} - 0.164$ eV to $E_{\rm C} - 0.695$ eV and hole traps with activation energies from $E_{\rm V} + 0.203$ eV to $E_{\rm V} + 0.526$ eV. DLTS studies made after annealing show that most of the defects anneal out above a temperature of 350 °C. The defects generated in the base region of the transistor by displacement damage appear to be responsible for an increase in base current through SRH or multi-phonon recombination and consequent transistor gain degradation.

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