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# Disappearance of plasmaspheric hiss following interplanetary shock

Zhenpeng Su

*University of Science and Technology of China*

Hui Zhu

*University of Science and Technology of China*

Fuliang Xiao

*Changsha University of Science and Technology*

Yuming Wang

*University of Science and Technology of China*

Chao Shen

*Chinese Academy of Sciences, Beijing*

*See next page for additional authors*

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**Authors**

Zhenpeng Su, Hui Zhu, Fuliang Xiao, Yuming Wang, Chao Shen, Min Zhang, Shui Wang, C A. Kletzing, W. S. Kurth, G. B. Hospodarsky, Harlan E. Spence, Geoffrey Reeves, H. O. Funsten, J. B. Blake, D. N. Baker, and J. R. Wygant

## RESEARCH LETTER

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## Key Points:

- First report on shock-induced disappearance of plasmaspheric hiss
- Interruption of chorus access into plasmasphere by enhanced Landau damping
- Removal of source electrons for chorus by the shrinking magnetopause

## Correspondence to:

Z. Su,  
szpe@mail.ustc.edu.cn

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## Disappearance of plasmaspheric hiss following interplanetary shock

Zhenpeng Su<sup>1,2,3</sup>, Hui Zhu<sup>1,2,3,4</sup>, Fuliang Xiao<sup>5</sup>, Huinan Zheng<sup>1,2,3</sup>, Yuming Wang<sup>1,3</sup>, Chao Shen<sup>6</sup>, Min Zhang<sup>1</sup>, Shui Wang<sup>1</sup>, C. A. Kletzing<sup>7</sup>, W. S. Kurth<sup>7</sup>, G. B. Hospodarsky<sup>7</sup>, H. E. Spence<sup>8</sup>, G. D. Reeves<sup>9</sup>, H. O. Funsten<sup>10</sup>, J. B. Blake<sup>11</sup>, D. N. Baker<sup>12</sup>, and J. R. Wygant<sup>13</sup>

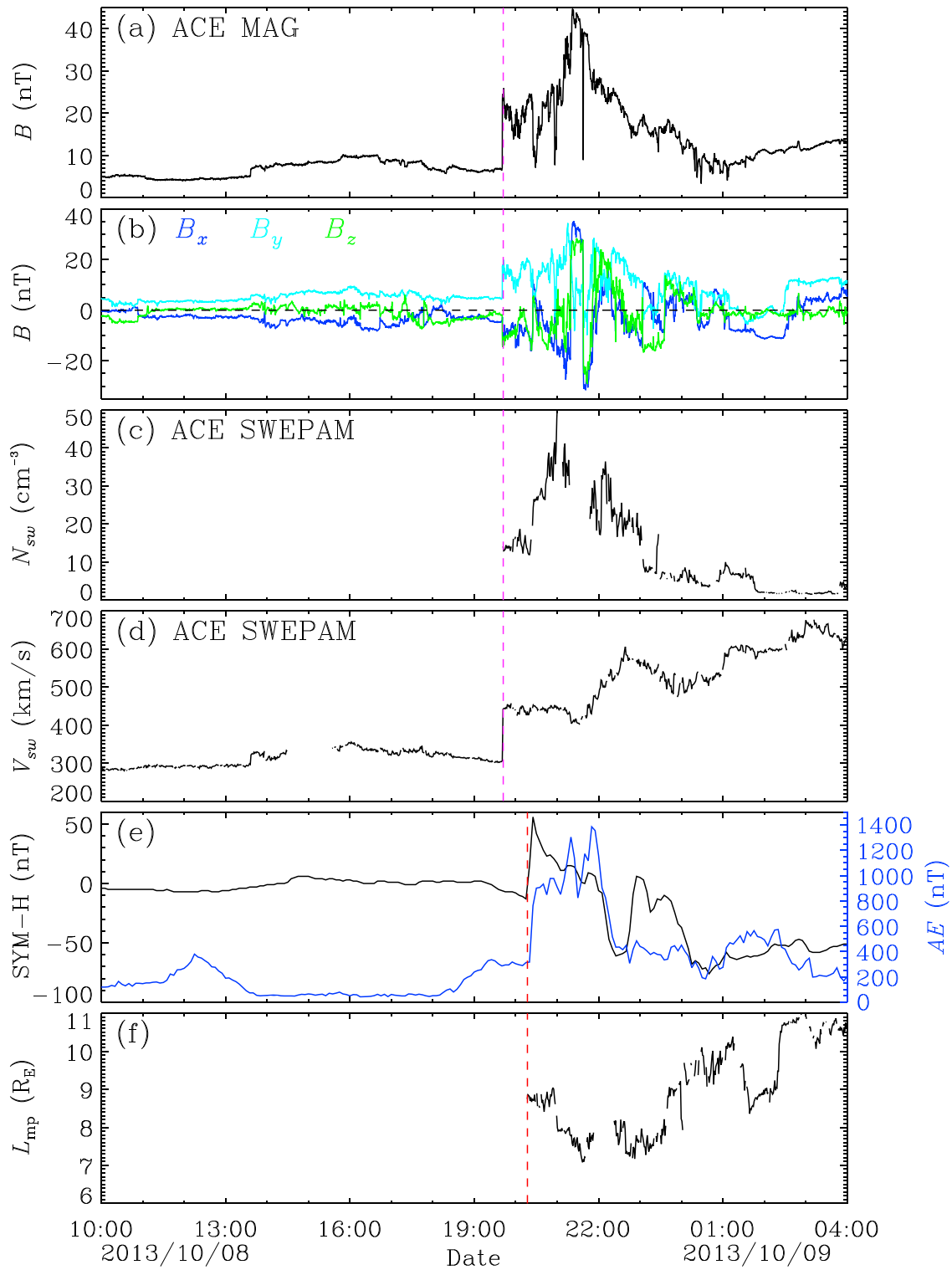
<sup>1</sup>CAS Key Laboratory of Geospace Environment, Department of Geophysics and Planetary Sciences, University of Science and Technology of China, Hefei, China, <sup>2</sup>State Key Laboratory of Space Weather, Chinese Academy of Sciences, Beijing, China, <sup>3</sup>Collaborative Innovation Center of Astronautical Science and Technology, Hefei, China, <sup>4</sup>Mengcheng National Geophysical Observatory, School of Earth and Space Sciences, University of Science and Technology of China, Hefei, China, <sup>5</sup>School of Physics and Electronic Sciences, Changsha University of Science and Technology, Changsha, China, <sup>6</sup>State Key Laboratory of Space Weather, Center for Space Science and Applied Research, Chinese Academy of Sciences, Beijing, China, <sup>7</sup>Department of Physics and Astronomy, University of Iowa, Iowa City, Iowa, USA, <sup>8</sup>Institute for the Study of Earth, Oceans, and Space, University of New Hampshire, Durham, New Hampshire, USA, <sup>9</sup>Space Science and Applications Group, Los Alamos National Laboratory, Los Alamos, New Mexico, USA, <sup>10</sup>ISR Division, Los Alamos National Laboratory, Los Alamos, New Mexico, USA, <sup>11</sup>The Aerospace Corporation, Los Angeles, California, USA, <sup>12</sup>Laboratory for Atmospheric and Space Physics, University of Colorado Boulder, Boulder, Colorado, USA, <sup>13</sup>School of Physics and Astronomy, University of Minnesota, Twin Cities, Minnesota, USA

**Abstract** Plasmaspheric hiss is one of the important plasma waves controlling radiation belt dynamics. Its spatiotemporal distribution and generation mechanism are presently the object of active research. We here give the first report on the shock-induced disappearance of plasmaspheric hiss observed by the Van Allen Probes on 8 October 2013. This special event exhibits the dramatic variability of plasmaspheric hiss and provides a good opportunity to test its generation mechanisms. The origination of plasmaspheric hiss from plasmatrail chorus is suggested to be an appropriate prerequisite to explain this event. The shock increased the suprathermal electron fluxes, and then the enhanced Landau damping promptly prevented chorus waves from entering the plasmasphere. Subsequently, the shrinking magnetopause removed the source electrons for chorus, contributing significantly to the several-hours-long disappearance of plasmaspheric hiss.

## 1. Introduction

Plasmaspheric hiss is a broadband whistler mode emission in the terrestrial plasmasphere and plasmaspheric plumes [Russell et al., 1969; Thorne et al., 1973; Summers et al., 2008]. It was widely considered to be structureless and incoherent, but its fine structures have been reported recently [Summers et al., 2014]. Hiss waves typically occur in the frequency range from ~0.1 kHz to several kilohertz [Hayakawa and Sazhin, 1992] and even extend to the lower frequency 20 Hz [Li et al., 2013]. The generation mechanism for plasmaspheric hiss is still under debate. Candidate mechanisms include the excitation by electron cyclotron instability in the outer plasmasphere [Thorne et al., 1979; Chen et al., 2014; Summers et al., 2014], and the origination from lightening whistlers [e.g., Sonwalkar and Inan, 1989; Green et al., 2005] or whistler mode chorus waves outside of the plasmasphere [Bortnik et al., 2008, 2009].

Through cyclotron resonance, hiss waves can cause the pitch angle scattering of radiation belt electrons over a wide energy range from ~0.1 MeV to several MeV [Horne and Thorne, 1998; Summers et al., 1998]. Such physical process contributes to the formation of the slot region separating the inner and outer radiation belts during quiet times [e.g., Lyons and Thorne, 1973; Abel and Thorne, 1998; Albert, 1994; Meredith et al., 2007] and the precipitation loss of outer radiation belt electrons during storm times [e.g., Li et al., 2007; Shprits et al., 2009; Su et al., 2011a; Mourenas and Ripoll, 2012; Thorne et al., 2013; Ni et al., 2014]. Hence, the information on the global spatiotemporal distribution of hiss is required to understand and/or predict the evolution of the electron radiation belt [Shprits et al., 2009; Subbotin et al., 2010; Su et al., 2010, 2011b; Tu et al., 2013; Glauert et al., 2014]. Based on the CRRES observations, Meredith et al., [2004, 2007] showed the statistical dependence of



**Figure 1.** Interplanetary and magnetospheric parameters: (a) solar wind magnetic field magnitude  $B$ ; (b) three components ( $B_x$ ,  $B_y$ , and  $B_z$ ) of solar wind magnetic field in the GSM coordinate system; (c) solar wind proton number density  $N_{sw}$ ; (d) solar wind bulk speed  $V_{sw}$ ; (e) geomagnetic activity indices  $SYM-H$  and  $AE$ ; (f) magnetopause nose location  $L_{mp}$ .

hiss waves on substorm activity. Based on the THEMIS observations, *Golden et al.* [2012] gave a statistical hiss distribution model driven by solar wind parameters and geomagnetic activity indices. Based on the Cluster observations, *Agapitov et al.* [2013] presented the statistical characteristics of hiss wave normal angles. Based on the Polar observations, *Tsurutani et al.* [2015] emphasized the influence of solar wind ram pressure on hiss waves. However, it remains to be determined to what extent these statistical models can reproduce the realistic variability of hiss waves.

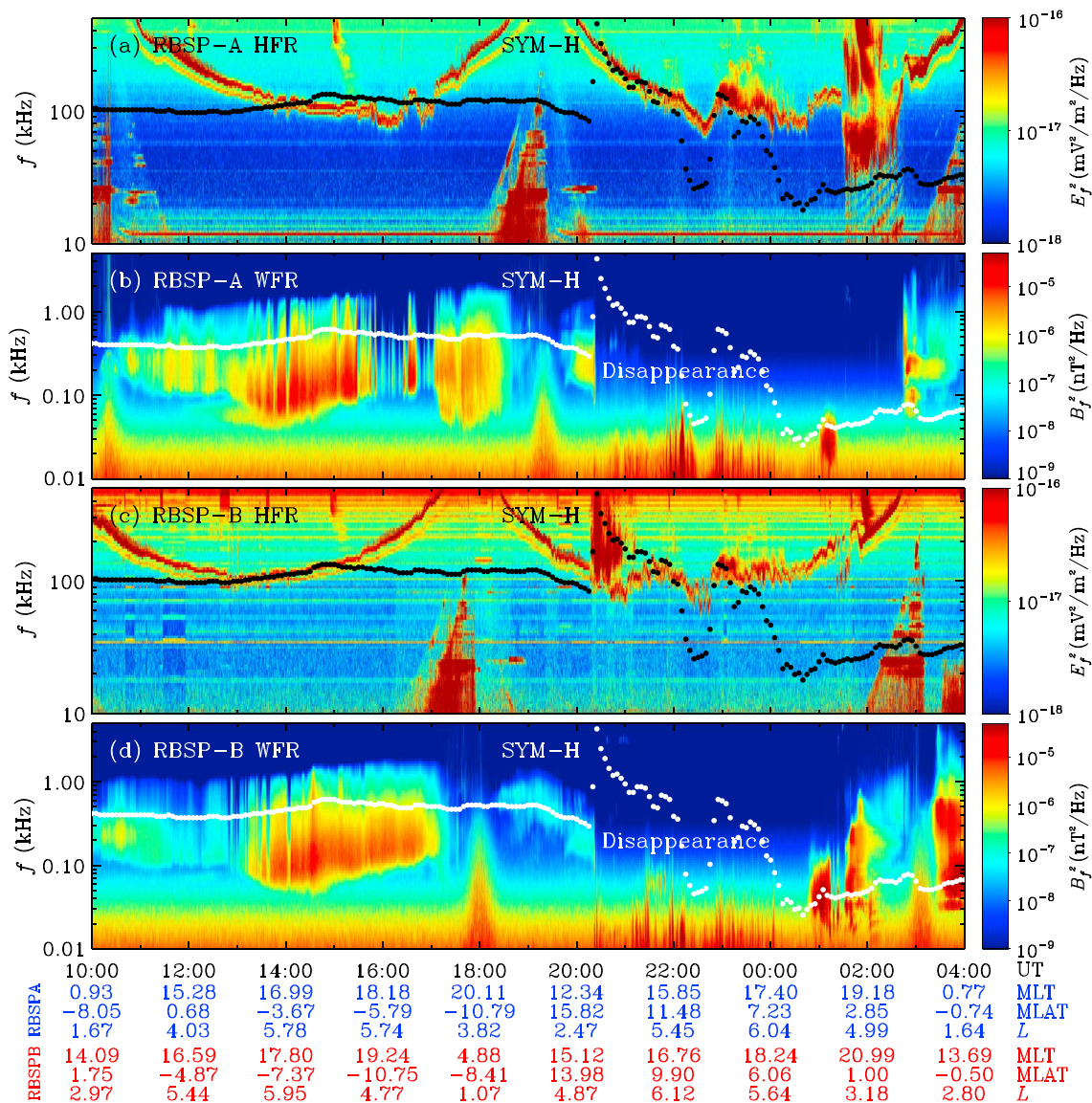
In this letter, we report a plasmaspheric hiss event observed by the Van Allen Radiation Belt Storm Probes (RBSP) [*Mauk et al.*, 2013] during 8–9 October 2013. An interplanetary shock triggered a strong substorm (with maximum  $AE \approx 1400$  nT) but simultaneously quenched the plasmaspheric hiss for about 5 h, contrary to the statistical picture that strong substorm [*Meredith et al.*, 2004] or solar wind with high ram pressure [*Tsurutani et al.*, 2015] enhance hiss activity. Note that *Thorne et al.* [1974] had reported two substorm events with the reduction of duskside hiss. To our best knowledge, this is the first report on the shock-induced disappearance of plasmaspheric hiss. This special event exhibits the significant and complex variability of plasmaspheric hiss and provides a good opportunity to test the generation mechanisms for plasmaspheric hiss.

## 2. Event Overview

Figure 1 plots the interplanetary and magnetospheric parameters from 10:00 UT on 8 October 2013 to 04:00 UT on 9 October 2013. The solar wind magnetic field  $\mathbf{B}$ , ion number density  $N_{sw}$  and bulk speed  $V_{sw}$  were observed by the Magnetic Fields Experiment [*Smith et al.*, 1998] and the Solar Wind Electron, Proton, and Alpha Monitor [*McComas et al.*, 1998] on board the Advanced Composition Explorer [*Stone et al.*, 1998]. The time-shifted interplanetary parameters can be used to determine the location of the magnetopause [*Shue et al.*, 1998]. The geomagnetic indices  $SYM-H$  and  $AE$  were provided by the World Data Center for Geomagnetism, Kyoto. The interplanetary shock driven by a coronal mass ejection can be clearly identified at 19:42 UT on 8 October 2013. After about 35 min, the shock compressed the magnetosphere (with the magnetopause nose shrinking to  $L_{mp} \sim 8.5$ ) and caused the geomagnetic storm sudden commencement (abrupt increase of  $SYM-H$  index from  $-10$  nT to 50 nT). Slightly later, the shock further triggered a strong substorm with the maximum  $AE \approx 1400$  nT.

Figure 2 shows the waves observed by the Electric and Magnetic Field Instrument Suite and Integrated Science instrument suite (EMFISIS) [*Kletzing et al.*, 2013] on board the twin RBSP satellites. The time range from 10:00 UT on 8 October 2013 to 04:00 UT on 9 October 2013 covered approximately two orbital periods of the RBSP satellites. At the shock arrival time (corresponding to the peak of  $SYM-H$ ), the RBSP-A satellite was in the slot region ( $L \approx 3.3$ ) and the RBSP-B satellite was approximately at the center of outer radiation belt ( $L \approx 5.3$ ). The upper hybrid resonance bands (bright lines) were clearly visible in the electric power spectral densities (Figures 2a and 2c) of the High-Frequency Receiver (HFR). These upper hybrid frequencies (positively correlated with the background electron density [*Kurth et al.*, 2014]) were above 60 kHz most of the time, indicating the locations of RBSP satellites in the plasmasphere. In the HFR channel, the shock-induced disturbances were quite evident for RBSP-B but invisible for RBSP-A. The plasmaspheric hiss waves with frequencies from 50 Hz to 1 kHz (Figures 2b and 2d) were detected by the Waveform Receiver (WFR). In the first orbital period, both RBSP satellites almost continuously received hiss waves. Around 13:00 UT on 8 October 2013, there was an intensification in the hiss spectrums of both RBSP satellites, probably caused by a weak substorm (Figure 1e). Around 17:00 UT on 8 October 2013, the hiss intensity observed by RBSP-A was modulated by the background electron density [*Chen et al.*, 2012d]. In the second orbital period, the hiss observed by both RBSP satellites abruptly ceased at the peak of the  $SYM-H$  index. After about 5 h, the hiss waves recovered with much stronger intensities (Figure 2d).

Figures 3 and 4 present the observations of background electron distributions and electromagnetic fields by the twin RBSP satellites. The omnidirectional/differential electron fluxes were collected by the Helium Oxygen Proton Electron (HOPE) Mass Spectrometer [*Funsten et al.*, 2013], Magnetic Electron Ion Spectrometer (MagEIS) [*Blake et al.*, 2013], and Relativistic Electron-Proton Telescope (REPT) [*Baker et al.*, 2013] of the Energetic Particle, Composition, and Thermal Plasma (ECT) suite [*Spence et al.*, 2013]. The magnetic field and electric field were provided by the EMFISIS Magnetometer and Electric Fields and Waves (EFW) instruments [*Wygant et al.*, 2013]. For the electron fluxes over a wide energy range, the RBSP-A observations responded weakly to the shock arrival, but the RBSP-B observations exhibited a sudden enhancement up to 2 orders of magnitude with the shock arrival. When the RBSP-A satellite went into the outer region ( $L > 3.6$ ), it observed the drift echoes of



**Figure 2.** (a–d) Electric/magnetic field spectral density in the HFR/WFR channels observed by the twin RBSP satellites. Note that the superposed solid dots represent the *SYM-H* index.

relativistic electrons (Figure 3d). Both spacecraft detected an increase in the magnetic field strength after the shock hit (most easily observed in Figure 4e due to the lower preshock field). The shock-induced variation in electric fields detected by both RBSP satellites had comparable peak amplitudes (10 mV/m) but different fluctuating periods (~3 min for RBSP-A and ~6 min for RBSP-B).

### 3. Discussion

There are generally three candidate generation mechanisms for plasmaspheric hiss (see section 1). Considering the poor connection between lightning and shock, we exclude the origination from lightening whistlers [Sonwalkar and Inan, 1989; Green et al., 2005]. If plasmaspheric hiss originates from plasmatrough chorus [e.g., Bortnik et al., 2008, 2009; Chen et al., 2012c], the following four stages would be involved: (1) excitation of chorus outside the plasmasphere, (2) propagation and Landau damping or cyclotron amplification of chorus, (3) refraction of chorus into the plasmasphere, and (4) amplification of hiss within the plasmasphere. In fact, the isolated occurrence of the fourth stage corresponds to the other generation mechanism, excitation by electron cyclotron instability [e.g., Thorne et al., 1979; Chen et al., 2014; Summers et al., 2014]. Obviously, the first





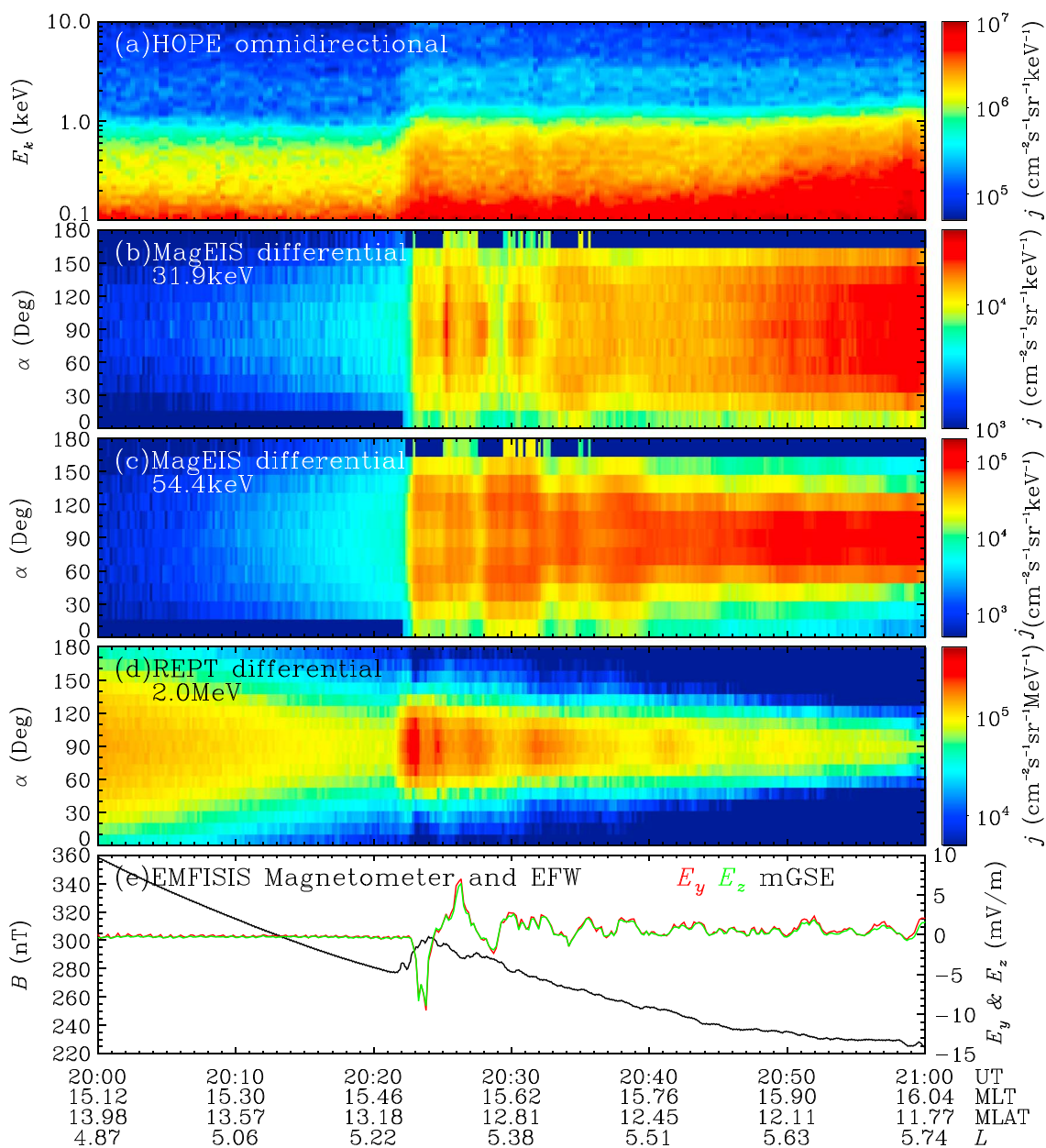


Figure 4. RBSP-B observations on the (a–d) omnidirectional/differential electron fluxes and (e) electromagnetic fields around the shock arrival time.

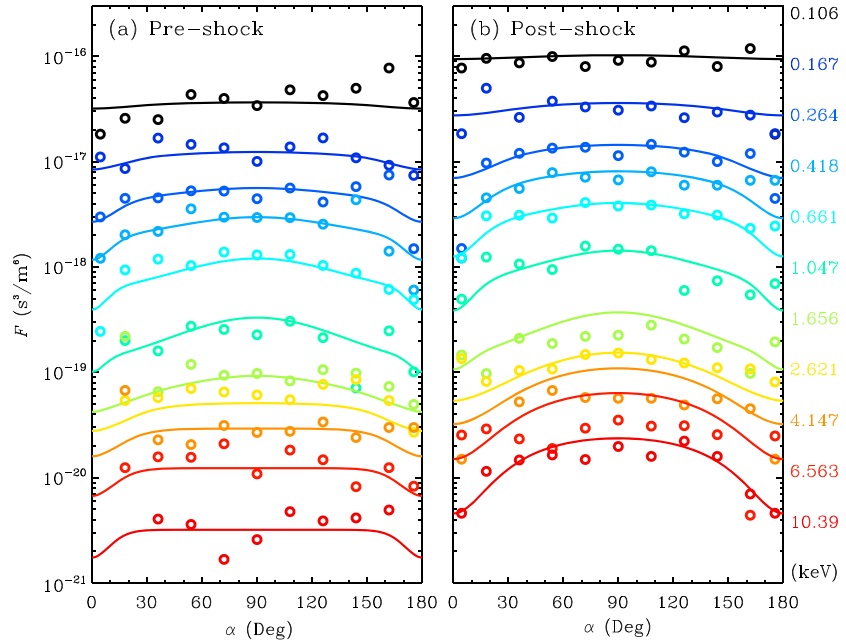
A detailed and accurate analysis of the electron energization process is beyond the scope of this letter and is left for future investigation.

In addition, the shock and the subsequent coronal mass ejection could highly compress the magnetosphere and remove the outer magnetospheric plasma through the “magnetopause shadowing” process [e.g., Turner *et al.*, 2012]. The shock-induced compression of the plasmasphere could also occur, as shown in previous simulations [e.g., Samsonov *et al.*, 2007] and observations [e.g., Zhang *et al.*, 2012]. During the geomagnetic storm driven by the shock and the subsequent coronal mass ejection, the plasmaspheric plume might gradually form due to the enhanced convection electric field [e.g., Goldstein *et al.*, 2005].

### 3.2. Excitation of Chorus Outside the Plasmasphere

Generally speaking, the chorus waves are excited through the cyclotron resonance with anisotropic suprathermal electrons [e.g., Kennel and Petschek, 1966; Nunn *et al.*, 1997; Omura *et al.*, 2008; Li *et al.*, 2009]. As discussed in section 3.1, we expect a prompt increase of suprathermal electron fluxes after the shock, which could





**Figure 5.** Suprathermal electron PSDs (a) before and (b) after shock at  $L = 8.0$  with circles for the hypothetical observations and lines for the simulations. Note that the suprathermal electron PSDs in the plasmatrough ( $L = 8.0$ ) have been assumed to be the tenfold of those observed by RBSP-B in the plasmasphere ( $L = 5.3$ ).

favor the excitation of chorus waves. Without the subsequent chorus damping during propagation (see section 3.3), such prompt energization process would intensify (rather than reduce) the plasmaspheric hiss. During 21:00–24:00 UT on 8 October 2013, the magnetopause nose shrank to about  $L_{mp} = 7–8$  (see Figure 1f), quite close to the preshock magnetopause  $L \gtrsim 6.2$  (Figure 2). The drastic removal of source electrons for chorus by the magnetopause shadowing could contribute largely to the several-hours-long disappearance of plasmaspheric hiss waves.

### 3.3. Propagation and Landau Damping or Cyclotron Amplification of Chorus

The chorus waves experience Landau damping/cyclotron amplification by the suprathermal electrons in the course of propagation toward higher latitudes [e.g., Bortnik et al., 2011; Chen et al., 2012a]. The intuitive idea is that the enhanced Landau damping would interrupt the propagation of chorus. We attempt here to estimate the shock-induced variation in the Landau damping rate.

Considering that the plasmasphere had extended to at least  $L = 6.2$  (Figure 2), we specifically calculate the Landau damping rates at a selected plasmatrough location  $L = 8$ . Based on the statistical results of Li et al. [2010], we assume that the suprathermal electron fluxes at  $L = 8$  equaled tenfold of those observed by RBSP-B in the outer plasmasphere at  $L = 5.3$ . Figure 5 shows the “hypothetically observed” (circles) and modeled (lines) suprathermal electron phase space density (PSD) before and after the shock arrival. The modeled PSD  $F$  contains  $N = 4$  components

$$F(v_{\perp}, v_{\parallel}) = \sum_{i=1}^N F_i, \tag{1}$$

$$F_i = \frac{n_i}{(\sqrt{\pi} V_{th_i})^3} \exp \left[ - \left( \frac{v_{\parallel}}{V_{th_i}} - V_{dr_i} \right)^2 \right] \times \left\{ \frac{\Delta_i}{\alpha_{1_i}} \exp \left( - \frac{v_{\perp}^2}{\alpha_{1_i} V_{th_i}^2} \right) + \frac{1 - \Delta_i}{\alpha_{1_i} - \alpha_{2_i}} \times \left[ \exp \left( - \frac{v_{\perp}^2}{\alpha_{1_i} V_{th_i}^2} \right) - \exp \left( - \frac{v_{\perp}^2}{\alpha_{2_i} V_{th_i}^2} \right) \right] \right\}. \tag{2}$$

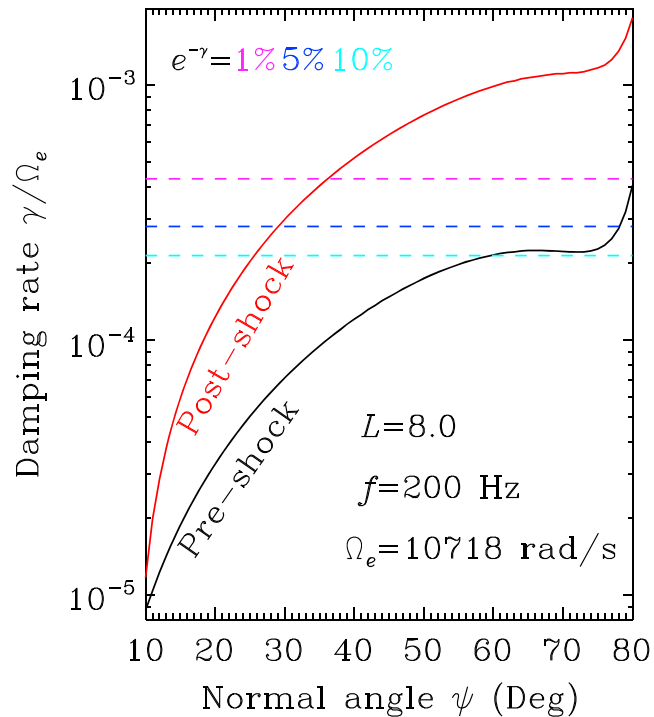
For the  $i$ th plasma component,  $n_i$  is the density,  $V_{th_i} = \sqrt{2T_i/m_e}$  and  $V_{dr_i}$  are the field-aligned thermal velocity and the normalized drift velocity,  $\alpha_{1_i}$  and  $\alpha_{2_i}$  characterize the temperature anisotropy and the size of loss cone, and  $\Delta_i$  controls the depth of loss cone. The total electron number density is taken to be constant

**Table 1.** Fitting Parameters for Electron PSDs Before and After the Shock

Component	$n_i(\text{m}^{-3})$	$T_i$ (keV)	$\Delta_i$	$\alpha_{1_i}$	$\alpha_{2_i}$	$V_{\text{dr}_i}$
<i>Pre</i>						
1	$2.2900 \times 10^6$	0.0028	1.0000	1.0000	0.5000	0
2	$1.3000 \times 10^5$	0.0348	1.0000	1.0000	0.5000	0
3	$6.0000 \times 10^4$	0.2054	0.6000	1.2491	0.1249	0
4	$2.0000 \times 10^4$	2.7321	0.6000	1.0000	0.2000	0
<i>Post</i>						
1	$1.9550 \times 10^6$	0.0028	1.0000	1.0000	0.5000	0
2	$3.2000 \times 10^5$	0.0432	1.0000	1.0000	0.5000	0
3	$1.6000 \times 10^5$	0.2843	0.5000	1.2100	0.2420	0
4	$6.5000 \times 10^4$	3.0960	0.4000	1.0615	0.6369	0

$N_e = 2.5 \times 10^6 \text{ m}^{-3}$  [Sheeley et al., 2001]. Adopting those parameters listed in Table 1 leads to reasonable agreements between the modeled and hypothetically observed distributions. Clearly, the suprathermal electron PSDs increased by a factor of 2–6 with the larger anisotropy after the shock.

These energy and pitch angle-dependent PSDs can be used to analyze the amplification/damping of plasma waves [e.g., Kennel and Engelmann, 1966; Shklyar, 2011; Su et al., 2014]. Figure 6 presents the calculated damping rates using the Waves in Homogeneous Anisotropic Magnetized Plasma (WHAMP) code [Ronmark, 1982]. The background magnetic field is specified as  $B = 61 \text{ nT}$  (i.e., the equatorial magnetic strength at  $L = 8$  in the typical dipole field), and the wave frequency is selected as  $f = 200 \text{ Hz}$  (i.e., the peak frequency of hiss spectrum in Figure 2). It is found that the damping rate increased by a factor of ~3–5 for waves with normal angles  $\psi > 20^\circ$ . As illustrated in the previous ray-tracing simulations [Bortnik et al., 2011; Chen et al., 2012a], the chorus rays that can access the plasmasphere usually have wave normal angles  $\psi \approx 40^\circ - 70^\circ$ . The intensities of these chorus rays with  $\psi \approx 40^\circ - 70^\circ$  could be reduced to much less than 1% within 1 s after the shock. Such enhanced Landau damping might be responsible for quenching plasmaspheric hiss [Bortnik et al., 2007].



**Figure 6.** Normal-angle  $\psi$ -dependent damping rate  $\gamma$  for the  $f = 200 \text{ Hz}$  chorus waves at  $L = 8.0$  during the preshock/postshock times. Note that the dashed horizontal lines denote the damping rates satisfying the relations  $e^{-\gamma} = 1\%$  (magenta), 5% (blue), and 10% (cyan).

It should be mentioned that the WHAMP code fully includes the Landau damping and cyclotron amplification processes. For the present simulations, the preshock PSD with the weak anisotropy did not allow the effective amplification of waves at any normal angles, while the postshock PSD with the moderate anisotropy could amplify waves at the normal angles  $\psi < 10^\circ$  (not shown).

### 3.4. Refraction of Chorus Into the Plasmasphere

The refraction is considered to be controlled by the cold electron distribution. *Chen et al.* [2012b] have demonstrated that the location or width of the plasmopause do not affect the peak hiss intensity significantly. Hence, the potential change of plasmopause by the shock was difficult to explain the prompt disappearance of plasmaspheric hiss. However, *Chen et al.* [2009] have suggested that the location and width of the plasmaspheric plume control the access of chorus waves into the duskside plasmasphere. Within several hours after the shock, the contribution of the evolving plasmaspheric plume to the hiss disappearance may need further investigation.

### 3.5. Amplification of Hiss Within the Plasmasphere

The amplification rate of hiss waves depends on the suprathermal/energetic electron distribution [*Chen et al.*, 2012c]. As observed by the RBSP satellites (Figures 3a and 4a), the suprathermal/energetic electron fluxes exhibited significant enhancement in the outer plasmasphere, promoting the amplification of injected chorus [*Chen et al.*, 2012c] or thermal noise [*Thorne et al.*, 1979]. Hence, the shock-induced change in this stage was unable to quench plasmaspheric hiss.

## 4. Conclusions

Plasmaspheric hiss plays an important role in radiation belt dynamics during both quiet and storm times. Its global spatiotemporal distributions have been statistically investigated [*Meredith et al.*, 2004, 2007; *Golden et al.*, 2012; *Agapitov et al.*, 2013; *Tsurutani et al.*, 2015]. In a statistical sense, strong substorms or solar wind with high ram pressure can intensify the plasmaspheric hiss waves. Here we report an interesting counterexample provided by the RBSP satellites. Following the interplanetary shock on 8 October 2013, a strong substorm (with maximum  $AE \approx 1400$  nT) occurred, but simultaneously the plasmaspheric hiss waves disappeared for about 5 h. These observations clearly illustrate the significant and complex variability of plasmaspheric hiss waves.

The generation mechanism for plasmaspheric hiss is still under debate. The origination of plasmaspheric hiss from plasmatrrough chorus is suggested to be an appropriate prerequisite to explain this special event. The interplanetary shock produced the prompt acceleration of electrons over a wide range of energies and pitch angles. Landau damping in the plasmatrrough, controlled by the suprathermal (0.1–10 keV) electron fluxes, might have become too strong to allow the access of chorus into plasmasphere. As suggested in previous works [*Bortnik et al.*, 2008, 2009], the interruption of energy injection by chorus could cause the disappearance of plasmaspheric hiss. Moreover, the shock and the subsequent coronal mass ejection continuously compressed the magnetosphere and largely removed the source electrons for chorus through the magnetopause shadowing process, which could contribute significantly to the several-hours-long disappearance of plasmaspheric hiss. We reiterate that the current explanations are obtained based on the inferred variations in the plasmatrrough and magnetopause. The RBSP satellites, as well as the THEMIS or the Cluster satellites, were not in the appropriate positions to measure the dayside plasmatrrough/magnetopause properties during the time period of interest. Future global magnetospheric simulations are required to examine the proposed explanations.

The influence of the solar wind on the inner magnetosphere has long been investigated in the space physics community [e.g., *Baker et al.*, 1983, 1998; *Gonzalez et al.*, 1989; *Reeves et al.*, 1998; *Wang et al.*, 2010]. Interplanetary shocks have been reported to trigger the substorms [*Heppner*, 1955], storms [*Gonzalez et al.*, 1994], auroral brightening [*Zhou and Tsurutani*, 1999; *Meurant et al.*, 2004; *Su et al.*, 2011c], magnetospheric particle energization [e.g., *Blake et al.*, 1992; *Hudson et al.*, 1997; *Zong et al.*, 2009], and chorus intensification [*Fu et al.*, 2012]. The current observations exhibit a potentially new consequence of interplanetary shock on the inner magnetosphere, the disappearance of plasmaspheric hiss.

### Acknowledgments

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