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ISOTOPE EFFECTS AND NELIUM RETENTION BEHAVIOR IN VANADIUM TRITIDE

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ABSTRACT

The relaxation times of the M. T. and ³He nuclei have been measured in vanadium hydride and tritate assumes. Substantial isotope effects in both the phase transition temperatures and diffusion parameters have been found. When compared to hydrides, the tritide samples have lower transition temperatures and faster mobilities. The differences in the occupancies of the interstitial sites are largely responsible for these isotope effects. Nost of the helium atoms generated by tritium decay remain trapped in microscopic bubbles formed within the VT_X lattice. Evidence is presented for the gradual growth of the helium bubbles over periods of hundreds of days.

INTRODUCTION

The vanadium-hydrogen phases exhibit unusually large isotope effects with respect to phase compositions, crystal structures, and diffusion properties. Although most previous studies have been on the vanadium hydrides and deuterides, 1-3 some results for the tritide phases are now available. The present paper summarizes extensive nuclear magnetic resonance (NMR) measurements on the two tritide compositions VT 0.50 and VT 0.75 as well as the corresponding hydrides VH 0.50 and VR 0.76. Isotope effects in the phase transitions and diffusion properties were

obtained from the temperature dependences of the proton and triton relaxation times. Furthermore, the relaxation times of the tritium decay product (i.e., 3He -- the light helium (sotope) have been monitored for several years in both $v_{0.50}^{\rm T}$ and $v_{0.75}^{\rm T}$. The age dependences of the He relaration times are very similar to the behavior previously observed in several other metal critides. 8-10 The rather substantial quantities of helium that are retained in the VT samples are believed to be in the form of very high pressure microscopic gas bubbles. 8-10 This conclusion is consistent with recent transwission electron wicroscopy (TEM) studies 11 in VT_. The He relaxation times have also indicated that the mean radius of the helium bubbles in VT has increased by a factor of about three between 100 and 1600 days.

EXPERIMENTAL SECTION

All of the V(H,T) samples had been prepared from the same rod of zone-refined vanadium metal that was purchased from Materials Research Corporation, Orangebury, New York. Synthesis details as well as summaries of various properties of the VH_X and VT_X samples have been published. The powders were sealed in 7-mm and 9-mm o.d. evacuated glass tubes for the NNR measurements. The pulsed NNR techniques and spectrometer that were used to measure the

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spin-lattice $\{T_{ij}\}$ and spin-spin $\{T_{2n}\}$ relaxation times have also been previously described. 4,5,8 Although some T_1 values were determined at several resonance frequencies, 4,5 most of the proton and triton relaxation time data were obtained at 34.5 MHz. In order to minimize deleterious effects from helium generation, all of the triton data were collected within a few Veeks after Synthesis. All of the ³He NMR experiments were performed at a resonance frequency of 45.7 MHz and 3He data were obtained only at room temperature. After the completion of the NMR studies, each YT sample was analyzed for its final composition. Portions of the powders were thermally decomposed by induction heating to a nominal maximum temperature of about 1000°C in a calibrated volumetric system based upon an all metal Topler pump. The compositions of the evolved gases were determined by wass spectrographic analysis.

RESULTS AND DISCUSSIONS

Summaries of the phase transitions for $V(H,D,T)_{0.50}$ and $V(H,D,T)_{0.75}$ are presented in Table I where the phase compositions and transition temperatures are based upon previous resiscivity, thermal analysis, x-ray and neutron diffraction, microscopy, and MMR measurements. 1-4 Detailed descriptions of the various phases ere available in Refs. 1-4. The a-phases are bodycentered cubic where the hydrogen isotopes predominantly occupy the tetrahedral intersticial sites in a statistically random fashion. All of the other V(H,D,T), phases result from at least partial ordering on either octahedral sites (i.e., $\theta_{\rm H}$, $\theta_{\rm D}$, $\theta_{\rm T}$, $\delta_{\rm H}$, and $\epsilon_{\rm H}$) or tetrahadral sites (i.e., Y_D , δ_D , ξ_D , and presumably analogous phases for $YT_{0.75}$) which generate tetragonal or orthorhombic distortions of the vanadiom lattice. 1,4 The low-temperature phase compositions of VT_{0.50} and VT_{0.75} have not been definitely established but are based upon the character of the triton relexation times and expected analogies to the $\mathrm{VD}_{_{\mathbf{N}}}$ phases. Further work is needed to complete the VI_ phase diagram. Table I. Summary of phase transitions for several isotopic vanadium-hydrogen compositions. Transition temperatures correspond to a heating sequence. Descriptions of various phases are given in References 1-4.

$$\begin{split} & \beta_{\mathrm{H}} - \mathrm{vR}_{0 - 50} \frac{446 \ \mathrm{K}}{} - \epsilon_{\mathrm{H}} - \mathrm{vR}_{0 - 50} \frac{470 \ \mathrm{K}}{} - \alpha_{\mathrm{H}} - \mathrm{vR}_{0 - 50} \\ & \beta_{\mathrm{D}} - \mathrm{vD}_{0 - 50} \frac{407 \ \mathrm{K}}{} - \alpha_{\mathrm{D}} - \mathrm{vD}_{0 - 50} \\ & \beta_{\mathrm{T}} - \mathrm{vT}_{0 - 50} \frac{372 \ \mathrm{K}}{} - \alpha_{\mathrm{T}} - \mathrm{vT}_{0 - 50} \\ & \delta_{\mathrm{H}} - \mathrm{vR}_{0 - 75} \frac{210 \ \mathrm{K}}{} - \epsilon_{\mathrm{H}} - \mathrm{vR}_{0 - 75} \frac{455 \ \mathrm{K}}{} - \alpha_{\mathrm{H}} - \mathrm{vR}_{0 - 75} \\ & \alpha_{\mathrm{D}} - \mathrm{vD}_{0 - 75} \frac{151 \ \mathrm{K}}{} - \delta_{\mathrm{D}} - \mathrm{vD}_{0 - 75} \frac{211 \ \mathrm{K}}{} - \delta_{\mathrm{D}} - \mathrm{vD}_{0 - 75} \\ & \frac{217 \ \mathrm{K}}{} - \alpha_{\mathrm{D}} - \mathrm{vD}_{0 - 75} \frac{?}{} - \zeta_{\mathrm{T}} - \mathrm{vT}_{0 - 75} \frac{225 \ \mathrm{K}}{} - \alpha_{\mathrm{T}} - \mathrm{vT}_{0 - 75} \end{split}$$

Since the NMR relexation times for VH_{0.50} and $\text{VT}_{0.50}$ have been previously compared in some detail, the temperature dependent behavior of the relaxation times for VK_{0.76} and VT_{0.75} is emphasized in the present paper. Figures 1 and I present the \mathbf{T}_1 and \mathbf{T}_{2n} data for these samples. Within experimental accuracy all the T, magnetization recoveries were exponential throughout the measured temperature ranges for both samples In contrast to the regions of nonexponential T, behavior previously found for VT_{0.50}. The frequency dependent T₁ minima indicate that diffusion is the primary spin relaxation process for both systems. The temperature behavior of the T_{2m} parameters is also consistent 12 with substantial proton and triton mobilities for temperatures above 150 K. However, the temperatures of the T, minima are significantly differeat (i.e., 282 K for VH_{0.76} and 217 K for VT_{0.75} at 34.5 MHz) which implies isotopic differences in the proton and triton diffusion parameters. This view is supported by the activation energies (E_g) obtained from conventional analyses 5, 12 of the proton and triton relaxation times. The available $\mathbf{E}_{_{_{\mathbf{J}}}}$ values for several V(H,D,T), compositions are compared in Table II along with the expected phase and interstitial site occupancy. The results for V(N,D,T) are

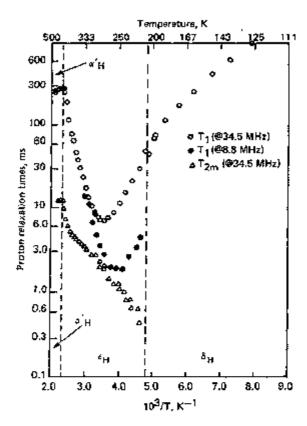


Fig. 1. Proton relexation times for VR $_{0.76}$ that were obtained during cooling sequences. The vertical dashed lines indicate the transitions between the hydride phases from References I-3.

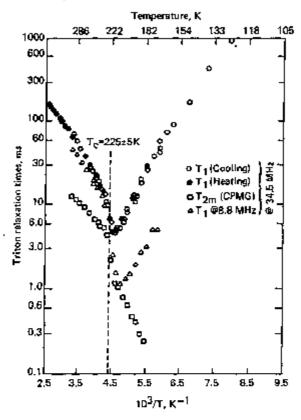


Fig. 2. Triton relaxation times for VT obtained during heating and cooling cycles. A phase transition from the high temperature ophase is indicated at 225±5 K; however, no significant changes in the T, and T data are observed at the expected $\hat{\tau}_{\rm p} + \delta_{\rm m}$ $\tau_{\rm p}$ phase transitions that would be analogous to VP 0.75 behavior.

Table II - Comparison of isotope effects in diffusion activation energy ($E_{\hat{a}}$) for vanadium-hydrogen phases.

Sample Composition	Phase	Hydrogen Site Occupancy	E (meV)	Temperature Range	Method
VIL.vo	αH	Tetrahedral	45±3	140 K-570 K	Corsky ^a
VD_vO	αŋ	Tetrahedral	73±3	175 K-570 K	Gorsky ⁸
vT _{vo}	ΦŢ	Tetrabedral	94±7	133 K-373 K	Gorsky ^a
VR _{0.50}	вн	Octahedral	410±20	320 K-400 K	MMR(T,)
VT _{0.50}	βT	Octahadral	310±30	290 K-345 K	MMR(T,)
VH _{0.76}	εĦ	Octahedra1	260±15	280 K-450 K	MMR(T,)
VH _{0.76}	ÉĦ	Octabedra1	180±10	210 K-280 K	NHR(T,)
vt _{0.75}	ΦT	Tetrehedral	160±10	230 K-380 K	NHR (T,)
νT _{0.75}	6T(?)	Tet rehedral	185±15	150 K-210 K	NHR (T)

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from Grosky effect measurements while the B values for ${
m VH}_{0.50}$ and ${
m VT}_{0.50}$ are from the previous NNR studies.

Several interesting isotope effects for the diffusion activation energies are apparent in Table II. First, $\mathbf{E}_{\mathbf{g}}$ increases rapidly with isotope mass in the low limit concentrations, whereas the opposite trend is observed for x = 0.50 and x = 0.75. Neither semi-classical nor current quantum theories can provide a satisfaccory and quantitative description of the hydrogen diffusion processes in the bcc metals V. Nb, or Ta although polaron models 13 do give qualitative understanding for much of the experimental results. For x = 0.50, the smaller E for the tritide has been attributed to increased disorder of the tritons among the interstitial sites of the β_p -phase which is reflected in more rapid mobilities and lower effective diffusion activation energies then are observed in the more highly ordered β_H -phase of $VH_{0.50}$. The smaller E_a for $\epsilon_B^{-\Psi h}$ 0.76 is also consistent with a direct correlation between more tapid motion (i.e., smaller E values) and increased disorder since the protons in the ϵ_{μ} -phase are randomly distributed over all the Oz octahedral sites but the protons are ordered on a subset of octabedral sites (i.e., $0z_1$) in $\beta_y = VR_{0.50}$. From the phase boundaties given in Figure 1, it is clear that $\mathbf{E}_{\mathbf{a}}$ values for the $\mathrm{VE}_{0.76}$ sample correspond to the diffusion behavior in the $\epsilon_{\mathbf{p}}$ -phase. The different E values obtained above and below the T₁ minima for ε_H-VH_{0.76} are believed to reflect complications that arise from the simultaneous presence of short-range and long-range jump processes with differing activation energies. 14 However, the \mathbf{E}_a values for $\mathbf{VT}_{0.75}$ correspond to diffusion in the disordered a-phase above 225 K and one or more of the ordered phases for tenperstures below 210 K. The relatively small E values and rapid triton mobilities observed for the "To.75 phases are probably consequences of preferred triton occupancies on tetrahedral sites. Similar diffusion parameters have been

reported 6,15 for hydrogen isotopes in NbH and TaH, where tetrehedral interstitial site occupancy is also dominant. From the E_ values in Table II, we conclude that disorder among the tetrahedral sites in VT_{0.75} has only minor influence on diffusion activation energies in contrast to the large effects found for diffusion among the octahedral interstitial sites of the β and ϵ_{H} phases. The effect of 1sotope concentration ascertsined only for α -phase VT $_{_{\mathbf{U}}}$ where $\mathbf{E}_{_{\mathbf{u}}}$ increases by a factor of 1.7 between x = 0 and x = 0.75. Similar comparisons of the published $\det a^{6,15}$ for α -NbH, and α -TaH, yield increases of 1.6 and 1.3, respectively. These larger values of $\mathbf{E}_{\mathbf{a}}$ are probably related to the contributions of hydrogen-hydrogen repulsive interactions to the diffusion process. 15 Consequently, the systematic reduction of this effect with the increased mass of the host metal is particularly intriguing and warrants further detailed study for all the hydrogen isotopes in these three metals.

The room temperature 3He relaxation times ${f T}_1$ and ${f T}_{2n}$ for the two ${f VT}_n$ samples are summarized in Figure 3. Although both parameters increase with age (1.e., helium content), T_{2m} increases roughly twice as fast as T, while both parameters are essentially constant after about 1800 days. behavior for the $^3\mathrm{He}$ relaxation times has been previously observed 8,10 in the tritides of Ti, Li, and U and has been related to the formation and growth of microscopic (i.e., mean diameters <10 nm) bubbles of very high pressure helium gas. The present He relaxation times in VT_ are completely consistent with this model. In fact, recent high-resolution TEM measurements 11 on S-VT_ have detected large concentrations of 1-2 nm diameter helium bubbles after only 100 days of tritium decay. From the increases in the He relaxation times, it is estimated that the mean bubble diameters in both VT samples increase by approximately a factor of three

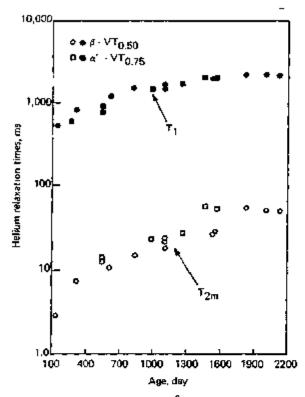


Fig. 3. Room temperature 3 Me relaxation times T_{\perp} (closed symbols) and T_{2m} (open symbols) for the VT asmples.

between 100 and 1500 days but remain essentially constant at about 5-6 nm after 1800 days for $VT_{0.50}$. This latter behavior probably reflects the irreversible rupture of the largest bubbles to release the helium from the solid as was originally proposed 8 for helium retention in VT_3 .

The ³He contents in VT_x, obtained from non-destructive NMR spin counts as well as thermal description experiments, are compared in Figure 4 to the predicted amount of ³He generated by tritium decay. Both techniques yielded ³He contents below the predictions although the uncertainty of the spin counts is about 10 to 20%. Because the ³He relaxation times had been constant for the previous several hundred days (i.e., corresponding to the continual rupture of the larges* helium bubbles), the substantial release of ²He yas observed after 2000 days at

room temperature is not particularly surprising. However, the absence of significant He contents in the gases evolved from 600-day old $VI_{0...50}$, and the very low quantity of He obtained from the $\nabla T_{0.75}$ after 1320 days is inconsistent with the essentially normal NNR signal amplitudes observed during the 3Re relaxation time experiments performed over this time period. Consequently, it is strongly suspected that not all of the helium is being released at the temperatures (i.e., approximately 1000°C) of the thermal description experiments on the younger samples. Because helium should be much more deeply trapped in isolated defects (e.g., vacancies) or very small clusters (i.e., bubble precursors), the presence of larger bubbles would facilitate helium release at lower temperatures for the older samples. explanation qualitatively reconciles the NMR spin count and thermal description results in Figure 4, and it is consistent with the incressing bubble size implied by the relaxation times in Figure 3. However, it is

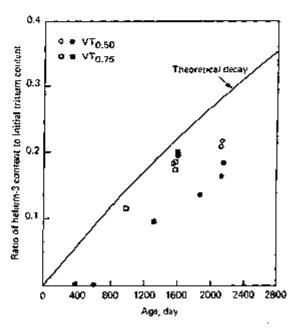


Fig. 4. Normalized belium-3 contents in VT samples from "He NMR spin-counts (open symbols) and thermal description measurements (closed symbols). The solid curve represents belium generated in VT from tritium radioactive decay.

presently impossible to estimate quantitatively how much of the helium not detected by the thermal desorption experiments was actually deeply trapped. From comparisons of the spin count and thermal description data in Figure 4, it appears that the $VT_{0.75}$ sample had released some (i.e., about 10 to 15%) of its helium content after 1600 days at room temperature, and subsequently released essentially all of its helium when heated to 1000°C. On the other hand, 2100-day old VI_{0.50} has apparently released about 25% of its helium at room temperature, but it had also retained more helium (i.e., roughly IO2) during the 1000°C desorption. Because VT_{0.75} generates helium at a rate about 1.5 times faster than ${
m VT}_{0.50}$ (where the small differences in lattice volumes are ignored), nominally identical quantities of belium had been formed and the above differences cannot be readily attributed to any major variation in total helium concentration. Furthermore, the Re relaxation times from Figure 3 indicate the mean bubble radii are also compara-Kence, other factors, which currently remain unidentified, can also influence the specific mechanisms for helium release. More information on the effects of stoichiomstry, phase composition, etc., on the distribution of helium bubbles and other defects will be zequired to resolve this issue.

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