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1                                   **The TRIUMF Nuclear Structure Program and TIGRESS**

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## 8 **ABSTRACT**

9 The Isotope Separator and Accelerator (ISAC) facility located at the TRIUMF laboratory in  
10 Vancouver, Canada, is one of the world's most advanced ISOL-type radioactive ion beam  
11 facilities. An extensive  $\gamma$ -ray spectroscopy programme at ISAC is centred around two major  
12 research facilities: i) the  $8\pi$   $\gamma$ -ray spectrometer for  $\beta$ -delayed  $\gamma$ -ray spectroscopy experiments  
13 with the low-energy beams from ISAC-I, and ii) the next-generation TRIUMF-ISAC Gamma-  
14 Ray Escape Suppressed Spectrometer (TIGRESS) for in-beam experiments with the accelerated  
15 radioactive ion beams. An overview of these facilities and recent results from the diverse  
16 programme of nuclear structure and fundamental interaction studies they support is presented.

## 17 **I. INTRODUCTION**

18 Nuclear physics has undergone a renaissance due to the development of new radioactive-ion  
19 beam facilities at several laboratories around the world. Of these, the Isotope Separator and  
20 ACcelerator facility (ISAC), which is of the ISOL type, located at TRIUMF in Vancouver,  
21 Canada, currently has the highest beam power available worldwide at 50 kW. ISAC employs a  
22 500 MeV proton beam, delivered by a cyclotron, with currents up to 100  $\mu$ A on targets ranging  
23 from SiC to Ta. Some of the most notable achievements are beams of  $^{11}\text{Li}$  up to  $5 \times 10^4 \text{ s}^{-1}$ ,  $^{26}\text{Al}$

1 at  $3 \times 10^{10} \text{ s}^{-1}$ , the separation of isomeric from ground state Ag beams with the laser ion source,  
2 and the observation of  $t_{1/2}=3.1 \text{ ms}$   $^{179\text{m}}\text{Lu}$  beam released from the Ta target. The ISAC facility  
3 (hereafter referred to as ISAC-I) is able to provide mass-separated beams at low-energy  
4 (typically 30 keV) as well as accelerated up to approximately 1.7 MeV/nucleon and mass 30. An  
5 extension of the laboratory, known as ISAC-II, will be able to accelerate beams up to mass 150  
6 and approximately 6.5 MeV/nucleon when completed. The wide variety of intense radioactive  
7 beams provides for a full programme of nuclear astrophysics, nuclear structure, weak interaction  
8 tests, and materials science. In the present work, the focus is on the nuclear structure programme  
9 using the  $8\pi$  spectrometer at ISAC-I and the TIGRESS spectrometer that will be located at  
10 ISAC-II.

## 11 **II. NUCLEAR STRUCTURE AT ISAC-I**

12 The nuclear structure programme at ISAC-I is centred (although not exclusively) on the  $8\pi$   
13 spectrometer and its auxiliary detectors. The  $8\pi$  spectrometer consists of twenty HPGe detectors  
14 with a relative efficiency of 20–25% surrounded by annular BGO shields for Compton  
15 suppression and a BGO “back-plug” detector. The detectors are positioned at the hexagon  
16 positions of a truncated icosahedron. The radioactive beam typically at 30-keV energy is  
17 deposited in the centre of the array, on either a metallic target or on a moving tape collector  
18 (MTC). The target-to-detector distance is approximately 14 cm, and the array efficiency is  
19  $\sim 1.5\%$  at 1.33 MeV. The tape of the MTC forms a continuous loop, typically 120 m in length,  
20 and has its movement governed by the programming of a controller unit.

21 A number of auxiliary detectors systems have been implemented in the  $8\pi$  spectrometer;  
22 SCEPTAR – the SCintillating Electron Positron Tagging ARray, PACES – the Pentagonal Array  
23 for Conversion Electron Spectroscopy, and DANTE – the Dipentagonal Array for Nuclear

1 Timing Experiments. SCEPTAR consists of twenty pieces of 1.6 mm thick plastic scintillator  
2 surrounding the target position, arranged so that there is a one-to-one match with the Ge  
3 detectors. Being relatively thin, SCEPTAR acts as a  $\Delta E$  detector for electrons and positrons, and  
4 is useful to indicate that a  $\beta$  decay has occurred. PACES incorporates five Si(Li) detectors 5 mm  
5 thick that replace the upstream half of SCEPTAR, and is most useful for conversion electron  
6 spectroscopy. DANTE consists of ten BaF<sub>2</sub> detectors positioned in the open pentagonal  
7 positions of the truncated icosahedron and is used for fast-timing measurements. This  
8 combination of detector systems is unique in the world, and provides for comprehensive  
9 spectroscopy of nuclei following  $\beta$  decay.

10 The  $8\pi$  nuclear structure program has examined nuclei from  $^{11}\text{Li}$  to  $^{179}\text{Lu}$ . An initial study  
11 [1] with the  $8\pi$  array to observe the  $\beta$  decay of a beam of  $10^3 \text{ }^{11}\text{Li s}^{-1}$  suggested that the  
12 theoretical prediction of a  $\beta$  decay occurring in the core and leaving intact the halo neutrons  
13 could be realized. Further studies [2] with greatly increased sensitivity using  $8\pi$ +SCEPTAR and  
14 a beam of  $>10^4 \text{ }^{11}\text{Li s}^{-1}$  have been performed that will yield a much more detailed picture. Figure  
15 1 shows portions of spectra obtained from the two  $^{11}\text{Li}$  experiments. The improvement in quality  
16 and statistics of the data is obvious, and the detailed  $\gamma$ -ray lines-shapes now available allow for  
17 precise determination of the level life-times and neutron decay branches. Future studies in this  
18 direction are planned, such as the heavier halo system  $^{14}\text{Be}$ .

19 A number of theoretical calculations predict that the familiar shell gaps that give rise to the  
20 “magic” numbers, or closed major shells, may change drastically in neutron-rich nuclei, and thus  
21 a major line of research is the mapping of the nuclear shells as one proceeds away from stability.  
22 The investigation of shell structure in the neutron-rich Mg and O isotopes is particularly  
23 interesting due to the “island of inversion” where the changes in the shell structure cause the

1 disappearance of the  $N=20$  neutron shell with the consequence that nuclei become deformed  
2 rather than retain a spherical shape, and other nuclei that would normally be predicted to be  
3 bound are unbound. An initial experiment [3] examining the decay of  $^{32}\text{Na}$  demonstrated that  
4 spectroscopy is possible with beam intensities as low as  $1\text{ s}^{-1}$ . More detailed studies must await  
5 the development of the actinide target for an expected 2–3 orders-of-magnitude boost in beam  
6 intensity. The very-neutron rich Cd and Pd isotopes are of interest in order to probe the  
7 quenching of the  $N=82$  shell below  $^{132}\text{Sn}$ . This has an important consequence on observed  
8 elemental abundances, particularly on the abundance of nuclei in the mass 120 region [4-6].  
9 This programme will soon commence now that Ag beams have been developed. A further  
10 region of interest, due to predicted shell closures at  $N=32$  and  $N=34$ , is the neutron-rich Ca  
11 isotopes and an experiment was recently performed to observe the  $\beta$  decay of the  $^{52}\text{K}$ . The  $^{52}\text{K}$   
12 beam intensity was  $\sim 10\text{ s}^{-1}$ , and sufficient data were collected to permit  $\gamma\gamma$  coincidences, analysis  
13 of which is ongoing [7].

14 For many years it has been known that nuclei near  $N=90$  exhibit a rapid change in the shape  
15 of the ground state as the neutron number increases from  $N=88$  to  $N=92$ . The excited-state  
16 spectrum also undergoes a rapid evolution from that resembling a vibrational system to a  
17 rotational one. Very recently [8,9], based on improved spectroscopy through  $^{152}\text{Eu}$  decay and  
18 IBM calculations, it was suggested that nuclei at  $N=90$  are at a critical point of a shape phase  
19 transition. A new model, based on a solution of the Bohr Hamiltonian with a square-well  
20 potential in the  $\beta$  shape degree of freedom, has been developed centred on this idea [10]. It has  
21 been claimed [11] that the new model successfully describes the level structures of a series of  
22  $N=90$  isotones; however, the lack of precision data prevents a definitive conclusion.





1 configurations for the clover detectors in TIGRESS; a “fully-forward” high efficiency mode, in  
2 which the front BGO shields are pulled back to allow the Ge detectors to be moved forward, and  
3 “fully-suppressed” optimum peak-to-total mode. Detailed GEANT4 simulations [14] indicate  
4 that the full array will achieve a photo-peak efficiency of 17% for a single  $\gamma$ -ray at 1 MeV.

5 A number of auxiliary detectors are planned for TIGRESS. The first to be commissioned is  
6 BAMBINO, which consists of 150- $\mu\text{m}$  thick Si CD-S2 detectors from Micron Technology Inc.  
7 that will be used for light-ion Coulomb excitation. Two such detectors can be mounted 3 cm  
8 from the target for both forward and backward hemispheres, covering a solid angle of  $\sim 1.15\pi$  sr,  
9 and each have 24 rings in  $\theta$  for angles between  $20^\circ$  and  $49^\circ$  and between  $131^\circ$  and  $160^\circ$ , and 16  
10 sectors in  $\phi$  for  $360^\circ$  coverage. Figure 3 shows a photograph of BAMBINO in its target chamber  
11 during a commissioning experiment. Also planned is a Si  $\Delta E$ - $E$  detector in a box-like geometry  
12 to be used for Coulomb excitation studies, single-nucleon transfer, and inelastic scattering, a  
13 segmented Bragg detector to be able to identify ions on an event-by-event basis, a CsI detector  
14 for light-ion identification for Coulomb excitation and fusion-evaporation studies, and a neutron  
15 detector employing deuterated liquid scintillator. A very powerful complement is the  
16 ElectroMagnetic Mass Analyzer (EMMA), a recoil spectrometer of the electric-dipole –  
17 magnetic-dipole – electric-dipole type [15], that is expected to be available in 2009.

18 The first experiment, performed at ISAC-I, conducted using TIGRESS and BAMBINO was  
19 the Coulomb excitation of  $^{20,21}\text{Na}$ . The system was commissioned with a beam of stable  $^{21}\text{Ne}$ .  
20 The beams had energies of 1.7 MeV/nucleon and intensities of  $2\text{--}6 \times 10^6 \text{ s}^{-1}$ , and impinged on a  
21 target of  $^{\text{nat}}\text{Ti}$  that was  $450 \mu\text{g}/\text{cm}^2$  thick. The resulting  $\gamma$ -ray spectra for  $^{21}\text{Ne}$  and  $^{21}\text{Na}$  beams  
22 (with a preliminary Doppler correction and without background subtraction) obtained in  
23 coincidence with particles detected in BAMBINO are shown in Fig. 4. Of particular interest is

1 the appearance of the 511-keV annihilation line that arises from the  $\beta^+$  decay of  $^{21}\text{Na}$  beam  
2 particles; the weakness of the peak indicates that the amount of scattered beam in the target  
3 chamber is well under control and will not interfere with the extraction of the matrix elements.  
4 By analyzing the particle- $\gamma$  coincidence yield as a function of scattering angle, the matrix  
5 element for Coulomb excitation can be uniquely determined.

#### 6 **IV. CONCLUSIONS**

7 Highlights of the nuclear structure studies at the TRIUMF ISAC facility were outlined. The  
8  $8\pi$  spectrometer and its associated auxiliary detectors is a world unique facility for nuclear  
9 physics utilizing  $\beta$  decay, in that  $\gamma$ -ray singles, coincidences, conversion electrons,  $\beta$ -particle  
10 tagging, and lifetime measurements can be performed simultaneously. A broad programme of  
11 nuclear structure, nuclear astrophysics, and Standard Model tests is underway. Nuclear structure  
12 investigations with TIGRESS have just commenced with the Coulomb excitation of radioactive  
13  $^{21}\text{Na}$ , and preliminary spectra from the experiment have been presented.

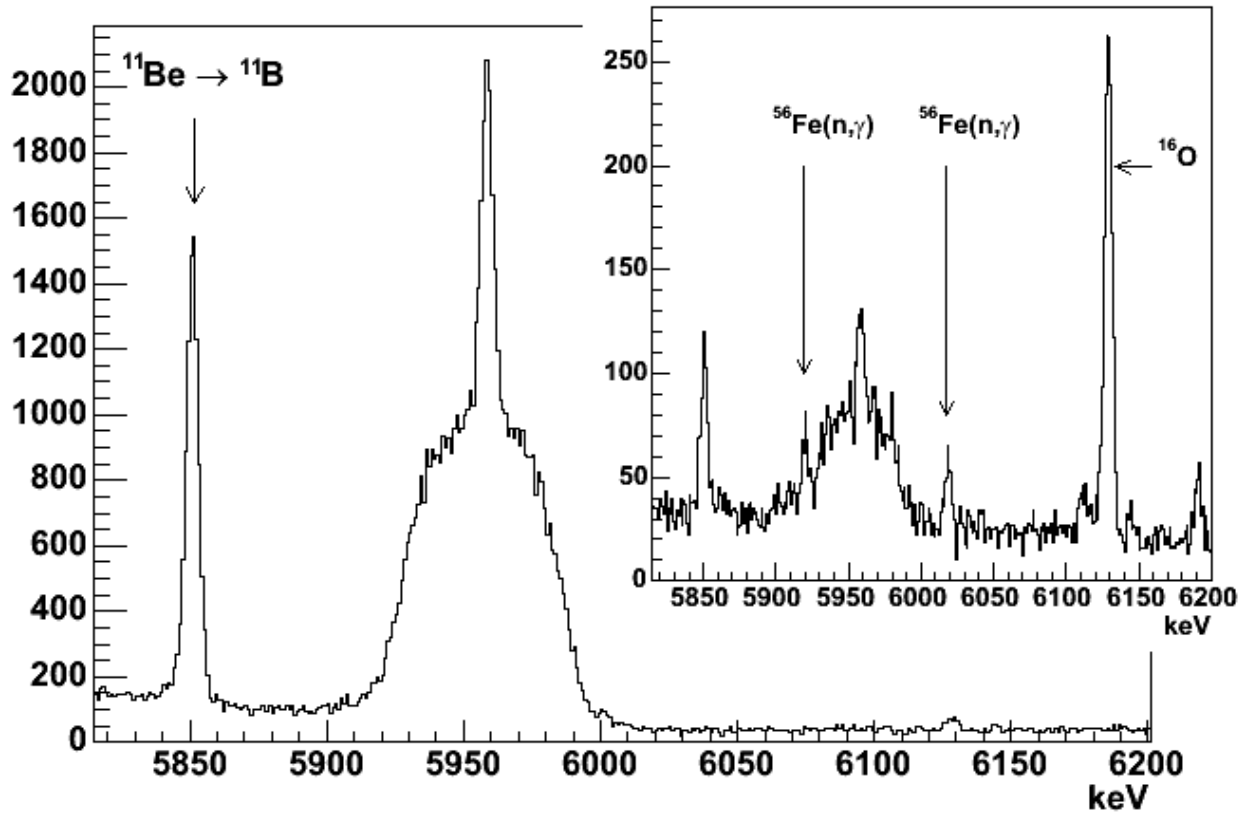
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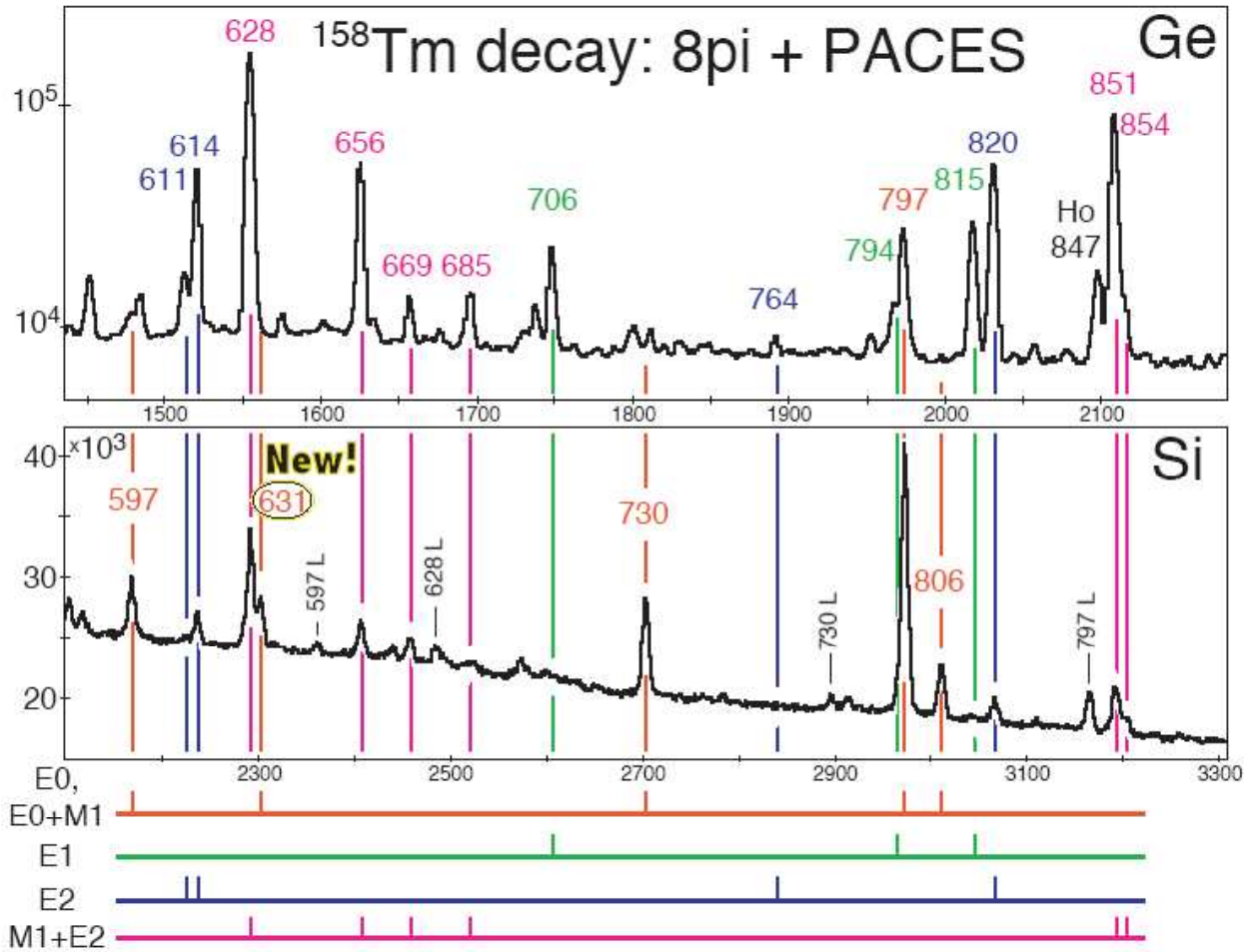
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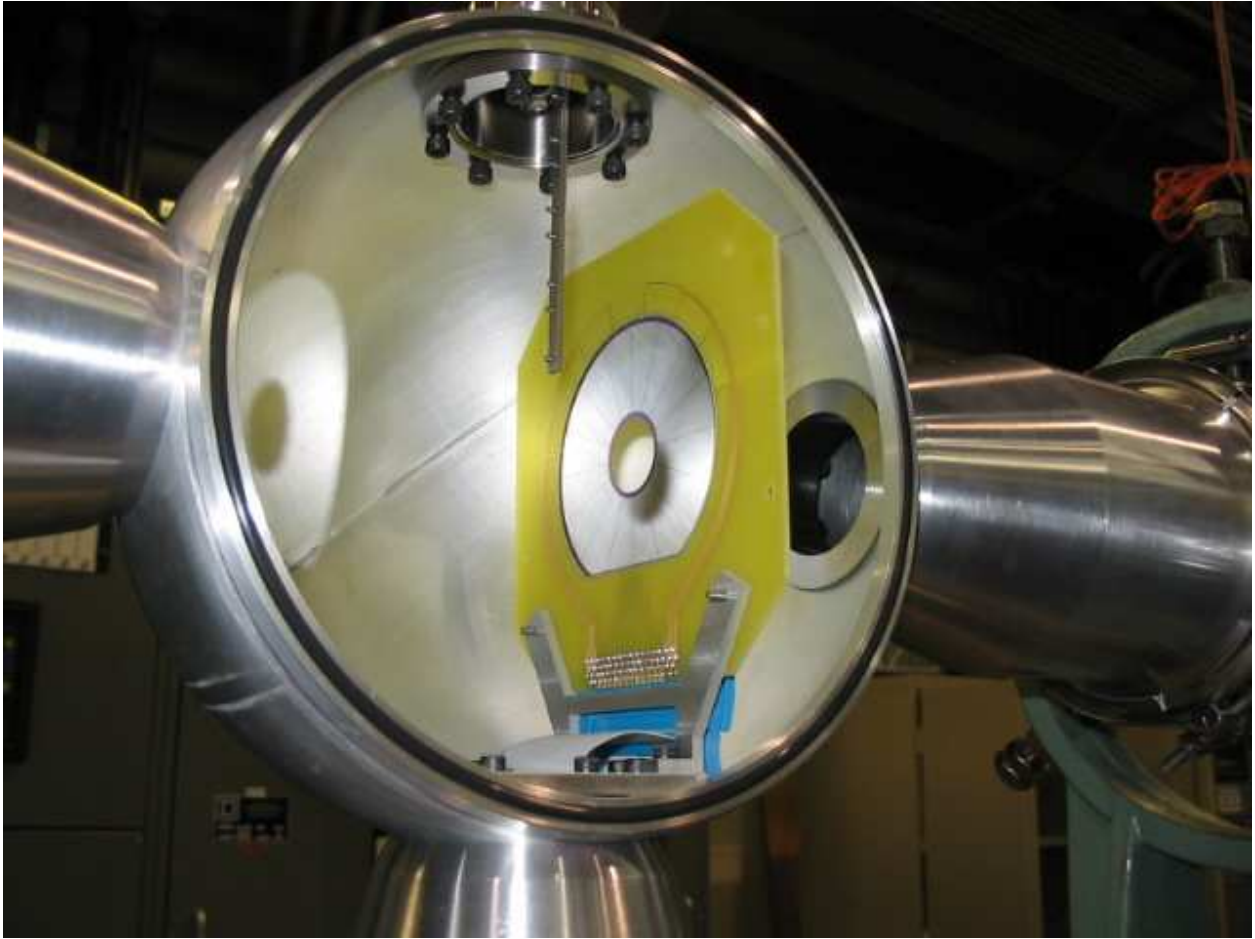
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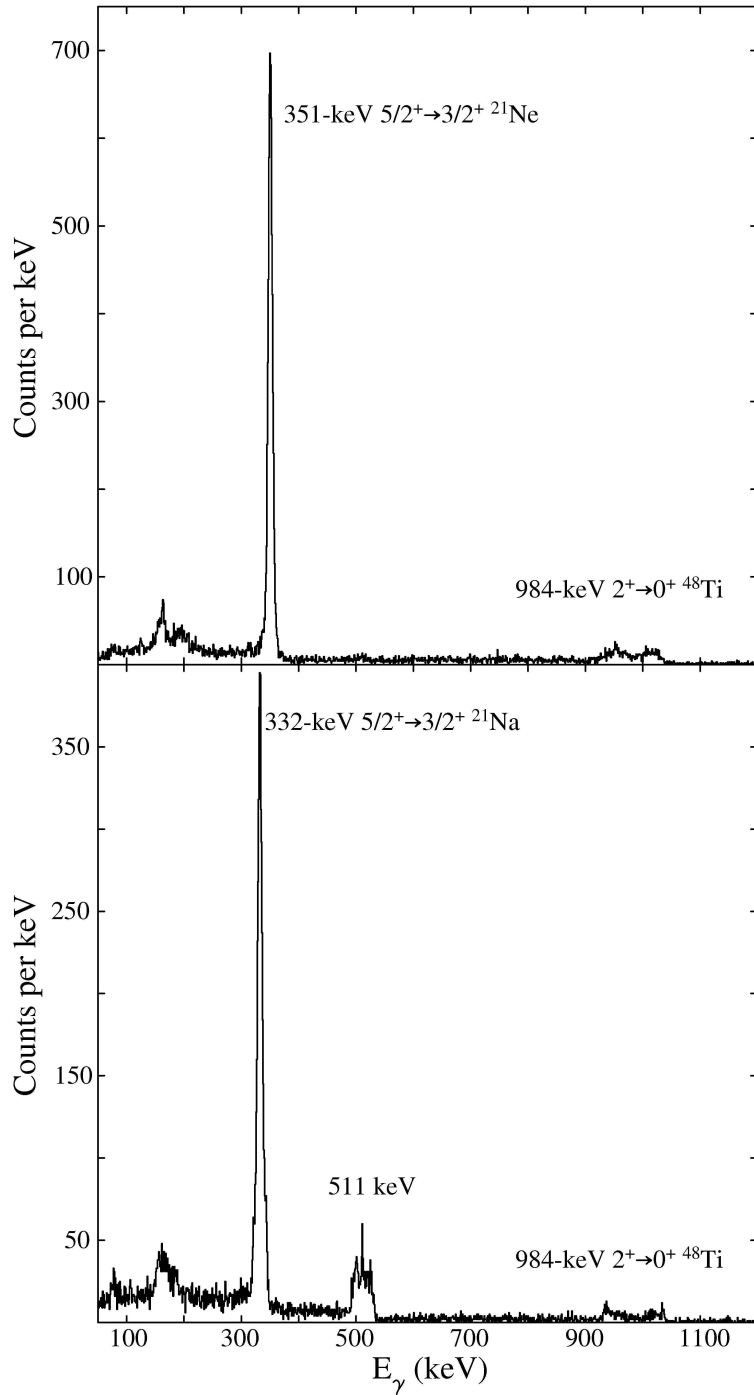
1 Figure 1. Comparison of data obtained from two  $^{11}\text{Li}$  decay experiments. In the first experiment  
 2 (inset), the  $^{11}\text{Li}$  beam intensity was  $\sim 10^3 \text{ s}^{-1}$ , and only the  $8\pi$  Ge detectors were used. In the  
 3 second experiment, the  $^{11}\text{Li}$  beam intensity was  $\sim 1.5 \times 10^4 \text{ s}^{-1}$ , and coincidences with  $\beta$  particles  
 4 detected with SCEPTAR was employed.



1 Figure 2. Portions of the  $\gamma$ -ray (top) and electron (bottom) spectra obtained with one of the  $8\pi$  Ge  
 2 detectors and one of the PACES Si detectors in the decay of  $^{158}\text{Tm}$ . Peaks that are intense on the  
 3 electron spectrum but moderate or missing in the  $\gamma$ -ray spectrum are identified as  $E0$  transitions.



1 Figure 3. Photograph of one BAMBINO Si CD-S2 detector mounted in its target chamber. The  
2 beam enters from the left, impinges on the Ti target (not shown) and passes through the hole in  
3 the centre of the CD detector.



1  
 2 Figure 4. Portions of the  $\gamma$ -ray spectrum obtained from the Coulomb excitation of  $^{21}\text{Ne}$   
 3 (top) and  $^{21}\text{Na}$  (bottom) on a target of  $^{\text{nat}}\text{Ti}$  using BAMBINO and TIGRESS. A coincidence  
 4 condition with particles detected in BAMBINO has been applied, without background  
 5 subtraction and a preliminary Doppler correction for the scattered beam particles has been used.