BNL 50199 (T-549)



MOGUS – A CODE FOR EVALUATING THE MOTT SCATTERING CROSS SECTION AND THE GOUDSMIT-SAUNDERSON ANGULAR MULTIPLE-SCATTERING DISTRIBUTION FOR USE IN ELECTRON TRANSPORT CALCULATIONS





July 1969

BROOKHAVEN NATIONAL LABORATORY

ASSOCIATED UNIVERSITIES, INC.

under contract with the

UNITED STATES ATOMIC ENERGY COMMISSION

BNL 50199 (T-549) (Mathematics and Computers – TID-4500)

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July 1969

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Printed in the United States of America Available from Clearinghouse for Federal Scientific and Technical Information National Bureau of Standards, U.S. Department of Commerce Springfield, Virginia 22151 Price Printed Copy \$3.00, Microfiche \$0.65

March 1970

575 copies

ABSTRACT

The formulas required for the calculation of the Mott cross section for scattering of electrons by an unscreened nuclear charge and the Goudsmit-Saunderson angular multiple-scattering distribution for electrons are summarized, and a FORTRAN-IV code which performs these calculations is given. The occurrence of the calculated quantities in electron transport theory is outlined in an introductory section. Subsequent sections contain a detailed summary of the transport model and analytical expressions for the calculated quantities, an outline of the code, a guide to the use of the code, illustrative results, and the code itself.

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I. INTRODUCTION: DEFINITIONS AND ASSUMPTIONS

Consider an energetic electron which originates at energy E_0 and slows down in an infinite homogeneous medium. A coordinate system is established such that the electron is produced at the origin and proceeds along a trajectory initially tangent to the z axis. The direction of flight at a point on the trajectory is characterized by ω and ϕ , the polar and azimuthal angles respectively (relative to the z axis) of the tangent to the trajectory, or alternately by $\mu = \cos \omega$ and ϕ .

A dimensionless energy τ is defined as the ratio of the electron's kinetic energy E to its rest mass energy, or

$$\tau = E (MeV)/0.510976$$
 (1)

Quantities needed to describe the slowing-down process are

- s (τ_0, τ) path length, or total distance traveled by the electron in slowing down from τ_0 to τ .
- $s_t(\tau_0)$ $s(\tau, 0)$, total path length for slowing down from τ to (in effect) 0.

 $t(\tau_0, \tau)$ fraction of the total path length $s_t(\tau_0)$ which an electron has yet to travel when its energy is τ . By definition

$$\mathbf{t}(\tau_{0},\tau) = [\mathbf{s}_{t}(\tau_{0}) - \mathbf{s}(\tau_{0},\tau)] / \mathbf{s}_{t}(\tau_{0}) \quad . \tag{2}$$

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 $A(\mu, t)d\mu$ probability that when the fractional residual path length $[A(\mu, s)d\mu]$ of an electron equals t (or the path length traversed equals s) the polar direction cosine lies within $d\mu$ of μ . Azimuthal symmetry is assumed.

A is normalized by the relation

$$\int_{-1}^{1} A(\mu, t) d\mu = \int_{0}^{\pi} A(\omega, t) \sin \omega d\omega = 1 , \qquad (3)$$

and hence, since the initial direction cosine is $\mu = 1$,

$$A(\mu, t=1) = \delta(\mu-1)$$
 . (4)

In general, electron slowing-down problems cannot be solved exactly; approximations must be made regarding the precise functionalities of $s(\tau_0, \tau)$ and $A(\mu, t)$, and for all but the simplest models numerical computations of one kind or another are required. Descriptions of models that have been adopted for various transport calculations are given in review articles by Berger¹ and Zerby and Keller.² A common model involves the use of the continuous-slowing-down approximation for energy loss calculations, so that a unique relationship between energy and path length is prescribed, and the choice of a multiple scattering formula for $A(\mu, t)$, of which the most accurate currently known is that of Goudsmit and Saunderson.³ Lateral spatial deflections are generally ignored in the absence of a satisfactory way to handle them, and one of various available empirical or semiempirical relations between range and energy is customarily invoked to correlate distance traveled in a specific direction with total distance traveled.

A code has been written which sets up the prerequisites for numerical transport calculations based on the model just described. For a given set of system parameters, the code generates a table of path length s and fractional residual path length t vs. dimensionless energy τ , and performs calculations which prepare for the subsequent generation of the Goudsmit-Saunderson distribution. Thereafter, if a range decrement r and initial energy E_o are specified, the program calculates the final energy E_f from an empirical relation due to Katz and Penfold, ⁴ interpolates in the previously generated t-E table to obtain the fractional residual path length t(E_f), and generates the Goudsmit-Saunderson distribution A(μ ,t) for specified values of the deflection angle $\omega = \cos^{-1}\mu$. Alternately, the fractional residual path length itself may be specified and A(μ ,t) calculated directly.

In the course of generating the Goudsmit-Saunderson distribution, it is necessary to calculate the Mott cross section for scattering of electrons by an unscreened nuclear charge.⁵ This calculation is performed by a routine which is called internally in the angular distribution calculation, and may also be called externally if values of this cross section are separately desired.

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In Section II the analytical expressions which form the basis of these calculations are presented. Section III contains an outline of the code, indicating where and how the code uses the material in Section II. Section IV contains a user's guide to the code, including required statements in the calling program, core storage requirements, and descriptions of additional routines which the user must supply. Section V contains illustrative results generated by the code and estimated running times, and Section VI contains the code itself and sample calling programs.

A. Path Length vs. Energy

0

At the point of origin of an energetic electron, the dimensionless energy τ and the path length s equal τ_0 and 0, respectively. τ and s are related in the continuous-slowing-down approximation by

$$\tau(\mathbf{s}) = \tau_0 - \int_0^{\mathbf{s}} |d\tau/d\mathbf{s}'| d\mathbf{s}' , \qquad (5)$$

and

$$s(\tau_{0},\tau) = \int_{\tau}^{\prime 0} |d\tau'/ds|^{-1} d\tau' .$$
 (6)

An expression for the stopping power $(d\tau/ds)_c$ due to collisions with atomic electrons has been derived, based on the Møller free electron scattering cross section for large energy transfers and a Born approximation formula due to Bethe for small transfers^{6,7}:

$$-(d\tau/ds)_{c}(cm^{-1}) = 0.301314 \frac{\rho Z}{W} \frac{(\tau+1)^{2}}{\tau(\tau+2)} \left[\ln \frac{0.130548 \tau^{2}(\tau+2)}{I^{2}} + f^{-}(\tau) -\delta(\tau,Z) - C_{s}(\tau,Z) \right] , \qquad (7)$$

$$f^{-}(\tau) = \frac{1 + \tau^{2}/8 - (2\tau + 1)\ln 2}{(\tau + 1)^{2}} , \qquad (8)$$

where ρ , Z, and W denote the density (g/cm³), atomic number, and atomic

weight, respectively, of the stopping medium, I (MeV) is the so-called mean ionization potential, δ is a correction for the mean energy loss reduction due to polarization of the stopping medium by the energetic electrons, and C_S accounts for effects which occur when the incident electron velocity approximately equals the electron orbital velocity in a given shell. The numerical constants of Eq. (7) represent the following combinations of physical and natural constants:

$$0.301314 \approx \frac{2\pi e^4 N_{av}}{(m_o c^2)^2}$$

$$0.13058 \approx \frac{1}{2} (m_o c^2)^2 ,$$

where e is the electronic charge, N_{av} is Avogadro's number, and m_0c^2 is the electron rest mass energy.

Caldwell⁸ has determined values of I for several materials, and Sternheimer⁹ has calculated values of δ based on these I values. Results of both sets of calculations are tabulated by Nelms.¹⁰ Huson⁷ has obtained an empirical expression for the shell correction C_s by fitting experimental data:

$$C_{s}(\tau, Z) = \frac{Z^{\frac{1}{2}}}{x+a} \left(1 - \frac{b}{x^{2}}\right) ,$$
 (9)

$$x = \frac{(137)^2 \tau(\tau+2)}{(\tau+1)^2 Z} , \qquad (9a)$$

$$a = \frac{Z^2(1 + 2/Z)}{1 + Z^2/30} , \qquad (9b)$$

$$b = \frac{36(Z^3 + 200)}{(10+Z)(Z^3 + 3000)} \quad . \tag{9c}$$

At higher energies, bremsstrahlung production makes an increasingly important contribution to the mean energy loss. The total stopping power due to both radiative and collision losses is approximately given by the formula¹¹

$$-(d\tau/ds) = -(d\tau/ds) [1 + (Z\tau/1400)], \quad Z\tau <<1400.$$
(10)

The validity of Eq. (10) is discussed in Section III. F.

B. Range vs. Energy

Katz and Penfold⁴ have proposed empirical range-energy relations based on a compilation of experimental thick absorber transmission data. These relations, in terms of r_t (cm) (defined in Section I), E (MeV), and ρ (g/cm³) are as follows:

1. $0.01 \le E \le 3 \text{ MeV}$ $r_t(E) = 0.412E^n/\rho$ (11) $n = 1.265 - 0.0954 \ln E$;

$$E(r_{t}) = \exp\left\{6.629979035639 - 5.241090146751[1.261848095652 - 0.3816 \ln(\rho r_{t})]^{\frac{1}{2}}\right\} .$$
(12)

2.
$$1 \leq E \leq 20$$
 MeV

$$r_{+}(E) = \rho^{-1}(0.530E - 0.106),$$
 (13)

$$E(r_t) = 1.886792453 \rho r_t + 0.2 .$$
 (14)

C. Mott Scattering Cross Section

The ratio of the Mott scattering cross section $\sigma_{\rm m}(\tau,\theta)$ for scattering of electrons by an unscreened nuclear charge to the Rutherford cross section $\sigma_{\rm r}(\tau,\theta)$ occurs as a parameter of the Goudsmit-Saunderson distribution. The latter cross section is (in cm²)

$$\sigma_{\rm r}(\tau,\theta) = \frac{7.961759 \times 10^{-26} \,{\rm z}^2(\tau+1)^2}{\tau^2(\tau+2)^2(1-\cos\,\theta)^2} \tag{15}$$

where

$$7.961759 \times 10^{-26} \approx e^4 / (m_o^2)^2$$

By definition, $2\pi\sigma_r \sin\theta d\theta$ is the cross section for scattering within $d\theta$ of the polar scattering angle θ .

The Mott cross section is given by 5

$$\sigma_{\rm m}(\tau,\theta) = \frac{7.944723 \times 10^{-26} \,{\rm z}^2}{\tau^2(\tau+2)^2} |F_0+F_1|^2 \,{\rm csc}^2(\theta/2)$$

$$+\frac{1.491145 \times 10^{-21}}{\tau(\tau+2)} |G_0+G_1|^2 \sec^2(\theta/2), \quad (16)$$

where

$$1.491145 \times 10^{-21} \approx (\hbar/m_{o}c)^{2} ,$$

$$7.944723 \times 10^{-26} \approx 1.491145 \times 10^{-21}/(137)^{2} .$$

The quantities F_0 , F_1 , G_0 , and G_1 in terms of

$$q = \frac{Z}{137} \frac{(\tau+1)}{\left[\tau(\tau+2)\right]^{\frac{1}{2}}}$$
(17)

 \mathbf{are}

$$F_{o} = (i/2) \exp\left\{iq \ln[\sin^{2}(\theta/2)]\right\} \Gamma(1-iq) / \Gamma(1+iq) , \qquad (18)$$

$$G_{o} = -iqF_{o} \cot^{2}(\theta/2) , \qquad (19)$$

$$F_{1} = (i/2) \sum_{k=0}^{\infty} (-)^{k} [kD_{k} + (k+1)D_{k+1}] P_{k}(\cos \theta) , \qquad (20)$$

$$G_{1} = (i/2) \sum_{k=0}^{\infty} (-)^{k} [k^{2}D_{k} - (k+1)^{2}D_{k+1}] P_{k}(\cos \theta) , \qquad (21)$$

where Γ is the gamma function and P_k the Legendre polynomial of order k. The coefficients D_k are

$$D_{k} = \frac{e^{-i\pi k}}{k + iq} \frac{\Gamma(k - iq)}{\Gamma(k + iq)} - \frac{e^{-i\pi\rho_{k}}}{\rho_{k} + iq} \frac{\Gamma(\rho_{k} - iq)}{\Gamma(\rho_{k} + iq)} , \qquad (22)$$

where

$$\rho_{\rm k} = \left[{\rm k}^2 - \left({\rm Z} / 137 \right)^2 \right]^{\frac{1}{2}} \ . \tag{23}$$

The series for F_1 and G_1 converge quite slowly, so much so that direct summation is all but impossible, even on the most advanced computers. To facilitate convergence, Feshbach¹² has applied the Euler transformation¹³ to Eqs. (20) and (21), and Sherman¹⁴ has applied two consecutive transformations, the first due to Yennie, Ravenhall, and Wilson¹⁵ and the second the Euler transformation. Calculations have shown, however, that the introduction of the second transformation does not improve convergence, and in fact often results in a considerable decrease in the convergence rate and hence an increase in computing time. The present code consequently omits the Euler transformation.

The Yennie, Ravenhall, and Wilson transformation is as follows. If a function $f(\theta)$ may be expressed as an infinite Legendre series,

$$f(\theta) = \sum_{k=0}^{\infty} a_k P_k(\cos \theta) , \qquad (24a)$$

then

$$(1-\cos \theta)^{m} f(\theta) = \sum_{k=0}^{\infty} a_{k}^{(m)} P_{k}^{(\cos \theta)} , \qquad (24b)$$

where

$$a_{k}^{(i+1)} = a_{k}^{(i)} - \frac{k+1}{2k+3} a_{k+1}^{(i)} - \frac{k}{2k-1} a_{k-1}^{(i)} .$$
(24c)

By definition, $a_k^{(0)} \equiv a_k^{(i)}$, and $a_{-1}^{(i)} \equiv 0$.

It has been found most efficient to apply this transformation to Eqs. (20) and (21) with m = 3 for $Z \leq 50$, and m = 4 for Z > 50. The transformed equations are

$$\mathbf{F}_{1} = \frac{1}{2}\mathbf{i}(1-\cos\theta)^{-m}\sum_{k=0}^{\infty}\mathbf{f}_{k}^{(m)}\mathbf{P}_{k}^{(\cos\theta)}, \quad \theta > 0$$
(25)

$$f_{k} = (-)^{k} [kD_{k} + (k+1)D_{k+1}];$$
 (25a)

$$G_{1} = \frac{1}{2}i(1-\cos \theta)^{-m} \sum_{k=0}^{\infty} g_{k}^{(m)} P_{k}(\cos \theta) , \quad \theta > 0 , \qquad (26)$$

$$g_{k} = (-)^{k} [k^{2} D_{k} - (k+1)^{2} D_{k+1}]$$
 (26a)

The coefficients $\{f_k^{(m)}\}\$ and $\{g_k^{(m)}\}\$ are generated from $\{f_k\}\$ and $\{g_k^{(m)}\}\$ respectively using (24c).

The calculational procedure is as follows. F_0 and G_0 are calculated from Eqs. (18) and (19). F_1 and G_1 are calculated from Eqs. (25) and (26), with D_k being given by (22), ρ_k by (23), and m = 3 ($Z \le 50$) or 4 (Z > 50). The cross section σ_m is calculated from (16).

The preceding calculational method clearly fails when $\theta = 0$, since both σ_m and σ_r are infinite at this angle; in addition, convergence is poor for small but nonzero values of θ , regardless of the transformation used. Fortunately, an analytical expression exists for σ_m/σ_r at small angles, namely¹

$$\sigma_{\rm m}/\sigma_{\rm r} \approx 1 + (\pi\beta^2 q/\sqrt{2}) \cos \gamma (1 - \cos \theta)^{\frac{1}{2}}, \qquad (27)$$

where

$$\beta^2 = \tau (\tau+2)/(\tau+1)^2$$
, (27a)

$$\cos \gamma = \operatorname{Re} \left\{ \frac{\Gamma(\frac{1}{2} - iq)}{\Gamma(\frac{1}{2} + iq)} \quad \frac{\Gamma(1 + iq)}{\Gamma(1 - iq)} \right\} , \qquad (27b)$$

and q is given by Eq. (17). An immediate result is

$$\sigma_{\rm m}/\sigma_{\rm r} = 1, \ \theta = 0$$
 (28)

Eq. (16) for $\sigma_{\rm m}$ appears to blow up at $\theta = \pi$ as a result of the sec²($\theta/2$) term; however, from Eq. (19) $G_0(\pi) = 0$, and since $P_k(-1) = (-)^k$, from Eq. (21) $G_1(\pi) = 0$ as well. Eq. (16) thus becomes for $\theta = \pi$

$$\sigma_{\rm m}(\pi) = \frac{7.944723 \times 10^{-26} \,{\rm z}^2}{\tau^2(\tau+2)^2} |F_{\rm o}^* + F_{\rm 1}^*|^2 \quad , \tag{29}$$

where

$$F_{0}^{*} = \frac{1}{2}i\Gamma(1-iq)/\Gamma(1+iq)$$
, (29a)

$$F_1^* = i \sum_{k=1}^{\infty} kD_k$$
 (29b)

D. The Goudsmit-Saunderson Angular Multiple-Scattering Distribution

The evaluation of the Goudsmit-Saunderson distribution using recursion relations developed by Spencer¹⁶ is discussed by Berger.¹ The present section is an expanded (and to an extent corrected) version of Berger's discussion.

The distribution is given by

$$A[\mu, t(s)] = \sum_{\ell=0}^{\infty} (\ell + \frac{1}{2}) \exp\left\{-\int_{0}^{s} G_{\ell}(s') ds'\right\} P_{\ell}(\mu) , \qquad (30)$$

where μ is the cosine of the angle of deflection at path length s relative to

the flight direction at s = 0, $A(\mu, t)$ is the probability defined in Section I, t(s) is given by Eq. (2), $P_{\rho}(\mu)$ is the Legendre polynomial of order ℓ , and

$$G_{\underline{\ell}}(s') = 2\pi N \int_{0}^{\pi} \sigma[\tau(s'), \theta] [1 - P_{\underline{\ell}}(\cos \theta)] \sin \theta \, d\theta \quad .$$
(31)

In Eq. (31), N is the numer of atoms/cm³, $\sigma(\tau, \theta)$ is the single scattering cross section, and $\tau(s)$ is given by Eq. (5).

Berger derives an expression for σ by applying a screening correction obtained from Moliere's multiple scattering theory¹⁷ to the Rutherford cross section σ_{r} , and multiplying the result by σ_{m}/σ_{r} , where σ_{m} is the Mott cross section. The factor σ_{m}/σ_{r} is in turn expressed as the sum of the small angle analytical formula [Eq. (27) + the screening term] for this ratio and a function h(θ), which is the true ratio σ_{m}/σ_{r} minus the small angle formula; this permits the integration of (31) to be carried out with relatively high precision in the small angle region, where the cross section has both its highest values and greatest degree of variation. Finally, the factor Z² in the Rutherford formula is replaced with Z(Z+1) to take into account inelastic scattering. The resulting expression multiplied by $2\pi N$ [expressed in terms of the density ρ (g/cm³) and atomic weight (W)] is

$$2\pi N\sigma(\tau,\theta) = \left[0.301314 \rho Z(Z+1)/W\right] \left[(\tau+1)/\tau(\tau+2)\right]^2 (1-\cos\theta+2\eta)^{-2} \\ \times \left\{1 + (\pi q/\sqrt{2}) \cos\gamma (1-\cos\theta+2\eta)^{\frac{1}{2}} + h(\tau,\theta)\right\},$$
(32)

where q is given by Eq_{\circ} (17),

$$\eta = 1.7 \times 10^{-5} Z^{2/3} [\tau(\tau+2)]^{-1} (1.13 + 3.76q^2), *$$
(33)

and

$$h(\tau, \theta) = (\sigma_{\rm m}^{\prime}/\sigma_{\rm r}^{\prime}) - 1 - (\pi q/\sqrt{2}) \cos \gamma (1 - \cos \theta + 2\eta)^{\frac{1}{2}}.$$
 (34)

The ratio σ_m / σ_r is given by the expressions of Section II.C. Berger notes in his paper a further small corrective term to $\sigma(\tau, \theta)$ suggested by Spencer, but he has since found this term to be suspect, ¹⁸ and it is accordingly omitted here.

The function $h(\tau, \theta)$ defined by Eq. (34) has been found to be well represented by a polynomial of the form¹⁶

$$h(\tau,\theta) = \sum_{j=1}^{J} h_j(\tau) (1 - \cos \theta + 2\eta)^{j/2}$$
(35)

where J is approximately 6. This expression is substituted into Eq. (32), and the resulting expression for $2\pi N\sigma$ is substituted into Eq. (31) to yield

$$G_{\ell}[\tau(s)] = 0.301314 \frac{\rho Z(Z+1)}{W} \left[\frac{(\tau+1)}{\tau(\tau+2)} \right]^{2} \left\{ p(-2,\ell) + [h_{1}(\tau) + (\pi q/\sqrt{2}) \cos \gamma] \right. \\ \times \left. p(-3/2,\ell) + \sum_{j=2}^{J} h_{j}(\tau) p[\frac{1}{2}(j-4),\ell] \right\}$$
(36)

where

$$p(m, \ell) = \int_{-1}^{1} (1 - \mu' + 2\eta)^{m} \left[1 - P_{\ell}(\mu') \right] d\mu' \quad .$$
(37)

*Berger's version of this formula, ¹ which omits the factor $[\tau(\tau+2)]^{-1}$, is incorrect.

Spencer has developed recursion relations for the terms p(m, l), which are as follows¹⁷:

$$p(m, 0) = 0$$
 , (38)

$$p(-2,1) = ln(1 + \eta^{-1}) - (1+\eta)^{-1}$$
, (39)

$$\ell p(-2, \ell+1) = (2\ell+1)(1+2\eta)p(-2, \ell) - (\ell+1)p(-2, \ell-1) - (2\ell+1)(1+\eta)^{-1},$$

$$\ell \ge 1$$
(40)

$$p(-3/2,1) = 2(2\overline{\eta})^{3/2} (1+\overline{\eta})^{-1}$$
, (41)

$$p(-3/2, \ell+1) = \overline{\eta} p(-3/2, \ell) + p(-3/2, 1) , \qquad \ell \ge 1$$
 (42)

$$p(m+1, \ell) = (1+2\eta)p(m, \ell) + p(m, 1) - \frac{\ell+1}{2\ell+1} p(m, \ell+1) - \frac{\ell}{2\ell+1} p(m, \ell-1)$$
(43)

where

$$\overline{\eta} = 1 - 2\eta \left[-1 + (1 + \eta^{-1})^{1/2} \right] \quad . \tag{44}$$

As an example, consider the generation of the $p(m, \ell)$ terms needed to evaluate Eq. (36) for J = 5. The required values of m for any ℓ are -2, -1, 0, -3/2, -1/2, and +1/2. The following schematic table indicates the order in which the p's are evaluated for values of ℓ from 1 to 4 inclusive: a (1) under (m = -2, $\ell = 1$) indicates that p(-2, 1) is calculated first, etc.

$m \rightarrow$	0	-1	-2	l	-3/2	$-\frac{1}{2}$	$\frac{1}{2}$	← m
	(6)	(3)	(1)	1	(1)	(3)	(6)	
	(9)	(5)	(2)	2	(2)	(5)	(9)	
	(12)	(8)	(4)	3	(4)	(8)	(12)	
	(15)	(11)	(7)	4	(7)	(11)	(15)	
		(14)	(10)	5	(10)	(14)		
			(13)	6	(13)			

According to this table, the first $p(m, \ell)$'s to be calculated are p(-2, 1)[from Eq. (39)] and p(-3/2, 1) [Eq. (41)]; next, p(-2, 2) [Eq. (40)] and p(-3/2, 2) [Eq. (42)]; then p(-1, 1) and p(-1/2, 1) [Eq. (43)], and so on. The row corresponding to a given value of ℓ contains the six p's needed to evaluate Eq. (36) for that ℓ .

If $G_1(\tau)$ is evaluated from Eq. (36) as indicated above, and the result is expressed first in terms of s [using Eq. (5) et seq.] and then in terms of t [as defined by Eq. (2)], it is found¹⁶ that the resulting function is well approximated by

$$G_{1}(t) = c_{1}[t(t+c_{2})]^{-1} .$$
(45)

The constants c_1 and c_2 are best determined by solving Eq. (45) using the exact values of $G_1(t)$ at t = 1 ($\tau = \tau_0$) and at a second value of t near t = 1. It is further found¹⁶ that to a good approximation

$$G_{\ell}(t) = G_{1}(t)G_{\ell}(1)/G_{1}(1)$$
 (46)

From Eqs. (2), (45), and (46),

$$\int_{0}^{S} G_{\ell}(s') ds' = s_{t} \int_{t}^{1} G_{\ell}(t') dt' = \frac{s_{t}c_{1}}{c_{2}} \frac{G_{\ell}(1)}{G_{1}(1)} \ln \left[\frac{t+c_{2}}{t(1+c_{2})}\right]^{*}.$$
 (47)

It is thus seen that $G_{\ell}(t)$, given implicitly by Eq. (36), need only be calculated for two values of t when $\ell = 1$ [to determine c_1 , c_2 , and $G_1(1)$], and for one value, t = 1, when $\ell > 1$ [since only $G_{\ell}(1)$ occurs in Eq. (47)].

All relations needed for the evaluation of the distribution, Eq. (30), have now been established. The full calculational procedure is summarized in Section III.

E. Convergence of the Goudsmit-Saunderson Series

A limitation of the Goudsmit-Saunderson distribution is that the series of Eq. (30) requires a prohibitively large number of terms for convergence when the fractional residual path length t exceeds a critical value; in fact, it appears that the convergence of the series may be asymptotic, which is to say the most accurate result is obtained by summing a finite number of terms, beyond which accuracy decreases. The onset of this phenomenon is quite dramatic: in one case studied, the series converged in 75 terms for t = 0.95, while convergence was far from achieved after 500 terms for t = 0.97.

*This formula is also given incorrectly in ref. 1.

It has been found that if the Yennie-Ravenhall-Wilson transformation, (24) with m = 3, is applied to the Goudsmit-Saunderson series for $\mu < 1$, convergence is usually achieved quite rapidly in instances where direct summation fails. The transformed series is

$$A(\mu,t) = (1-\mu)^{-3} \sum_{\ell=0}^{\infty} a_{\ell}^{(3)} P_{\ell}(\mu) , \qquad \mu < 1$$
 (48)

where

$$a_{\ell} = (\ell + \frac{1}{2}) \exp[-\int_{0}^{s} G_{\ell}(s') ds'] ,$$

 G_{ℓ} is given by (31), and $\{a_{\ell}^{(3)}\}\$ are generated from $\{a_{\ell}\}\$ using (24c).

If A (μ =1) is plotted against A (μ =cos 5°) for a variety of atomic numbers and penetrations and a moderately broad range of initial energies, the resulting points very nearly fall on a single curve. (See Figure 1.) This suggests an approximate method of obtaining A (μ =1,t) when (30) does not converge for μ =1 but (30) or (48) converges for μ =cos 5°: namely, evaluate A(cos 5°,t) and interpolate on the appropriate curve of Figure 1 to obtain A(1,t). It is then in turn possible to calculate A(μ ,t) if cos 5° < μ < 1 and (48) diverges for μ but converges for cos 5°. It has been observed^{3b} that at small angles the multiple scattering distribution behaves like a modified Gaussian,

$$A(\mu,t) \approx A(1,t)e^{-K(1-\mu^2)}$$

If K is chosen so that the given function fits the known values of A(1,t) and $A(\cos 5^{\circ},t)$ (the former having been obtained from the latter by interpolation from Figure 1), the resulting expression is

$$A(\mu,t) \approx A(1,t) \exp\left[\frac{1-\mu^2}{1-\cos^2 5^\circ} \ln \frac{A(1,t)}{A(\cos 5^\circ,t)}\right], \ \mu > \cos 5^\circ.$$
 (49)

The approximation just presented is not likely to be required in most calculations, since (30) or (48) converge almost everywhere, failing only when both the energy and the penetration are quite low. For completeness, it should be noted that if the penetration is <u>exceptionally</u> small and the energy is low, even approximation (49) fails, for one of two reasons: first, Eq. (48) may not converge for $\mu = \cos 5^{\circ}$, and second, the $[A(\mu=1), A(\mu=\cos 5^{\circ})]$ pair may lie beyond the curve of Figure 1. The implication in either case is that the distribution is very nearly a delta function, and hence angular deflections may be neglected with little attendant error. It is emphasized, however, that the occurrence of such a situation in a transport calculation is most unlikely.

III. OUTLINE OF THE CODE

A. Path Length vs. Energy: TTAB, DSCAL, DMUS, DENCOR, GAUS

The five named routines perform the calculations outlined in Section II. A. DENCOR generates the polarization correction δ of Eq. (7) by applying an interpolation formula to the appropriate curve of ref. 10. DMUS calculates $(d\tau/ds)$ from (9), and DSCAL calculates $|d\tau/ds|^{-1}$, the integrand of (6). TTAB carries out the integration of (6) using GAUS, a numerical quadrature subroutine which must be written or otherwise obtained by the user (see Section IV. B.). The quantities generated by TTAB include the total path length s_t for a given initial energy E₀, a table of path length s(E₀, E) vs. energy E, and the fractional residual path length t corresponding to each tabulated s. Both energy E and dimensionless energy τ are stored as alternative independent variables.

The table of s vs. E is generated by TTAB as follows. At the initial energy E_0 the path length s by definition equals 0. An energy E_1 is calculated as E_0 -h, where h is an energy decrement chosen by the user, and 12-point Gaussian quadrature is used to calculate $s(E_0, E_2)$ from Eq. (6). A second energy $E_2 = E_1$ -h is generated, integration over the range (E_2, E_1) is carried out, and the result is added to $s(E_0, E_1)$ to obtain $s(E_0, E_2)$. A final energy E=0 is ultimately reached, and $s_t = s(E_0, 0)$ is generated. Once the energy falls below 0.1 MeV, the slope dE/ds begins to change with increasing rapidity (see the tables of ref. 10), and so 20-point quadrature is used in this range. The values of s so calculated are estimated to be no more than 0.1% in error except at the very lowest energies, say less than 0.05 MeV, where the error may rise to 1 to 2%. Naturally, the smaller the energy decrement h the less is the error.

B. Range-Energy-Path Length: RANGEN

Subroutine RANGEN performs one of four calculations, according to the value of a parameter INDEX. Let $r_t(E)$ be the range and $t(E_0, E)$ the fractional residual path length as defined in Section I, expressed in terms of E (MeV) rather than τ , and define r(E) as $r_t(E_0)-r_t(E)$, so that r(E) is the distance the electron has traveled in the initial direction of its flight when its energy is E, its initial energy having been E_0 . The energy E_0 is assumed known in the ensuing discussion.

If INDEX = 1, the total range $r_t(E_0)$ is calculated from the appropriate Katz-Penfold formula [(11) or (13)]. This calculation must be performed before RANGEN can be called with INDEX greater than 1. If INDEX = 2, a value of r is read, the residual range $r_t(E)$ corresponding to this r is calculated as $r_t(E_0)-r$, the energy $E(r_t)$ is calculated from Eq. (12) or (14), and the fractional residual path length t(E) is obtained by interpolating in the previously generated t-E table. If INDEX = 3, the last of these calculations is omitted, and only $E[r_t(r)]$ is determined. If INDEX = 4, a value of t is read, E(t) is obtained by interpolation in the t-E table, $r_t(E)$ is calculated from Eq. (11) or (13), and r(E) is obtained as $r_t(E_o)-r_t(E)$.

C. Mott and Rutherford Cross Sections: RMOTT, RM1, CGAM

RMOTT calculates the Mott cross section $\sigma_{\rm m}$ and the Rutherford cross section $\sigma_{\rm r}$ using the formulas of Section II.C. If $\theta \leq 10^{\circ}$, the ratio $\sigma_{\rm m}/\sigma_{\rm r}$ is calculated from (27), $\sigma_{\rm r}$ is calculated from (15), and $\sigma_{\rm m}$ is obtained as the product of these quantities. The term $\cos \gamma$ of (27) is generated by CGAM in the first of these calculations.

If $\theta > 10^{\circ}$, the full calculational procedure of (16) <u>et seq.</u> is carried out to obtain $\sigma_{\rm m}$. The generation of the coefficients $\{D_k\}$ [Eq. (22)] as they are needed is carried out in RM1. The series of (25) and (26) are summed until twice in succession the addition of another term changes the value of F (or G) by less than five parts in 10^6 .

If RMOTT is called for $\theta = 180^{\circ}$, the value 179° is substituted and a message to this effect is printed out. The difference between $\sigma_{\rm m}(179^{\circ})$ and $\sigma_{\rm m}(180^{\circ})$ is insufficient to warrant the inclusion of Eq. (29) in the program. D. Goudsmit-Saunderson Distribution: GOUD, HFIT, CGAM, G1G, GLG

To generate $A(\mu,t)$ for a given initial energy, penetration, and angular deflection, GOUD must be called three times, with a parameter INDEX successively taking on values of 1, 2, and 3. When GOUD is called with INDEX = 1

for a given initial energy τ_0 , it in turn calls HFIT, which calculates $h(\theta)$ from Eq. (34) at IK equally spaced angles from 0 to π , and fits $h(\theta)$ exactly at these points with a polynomial of the form of Eq. (35). IK equals 7 in the program as currently written; it may be changed to any integer from 3 to 10 inclusive by redefining it in subroutine DEFINE. CGAM calculates $\cos \gamma$ from Eq. (27b) in the course of the evaluation of $h(\theta)$. GOUD then calls G1G, which uses the calculated coefficients $\{h_j\}$ to evaluate $G_1(t)$ from Eqs. (36)-(44) at two values of t: t = 1 ($\tau = \tau_0$) and the value once removed from t = 1 in the table generated as indicated in Section III.A. The resulting values of G_1 are used to determine c_1 and c_2 from Eq. (45), and control is then returned to the calling program.

The first value of the deflection angle cosine μ for which the distribution is calculated for a given fractional residual path length t must be 1, a calculation performed when GOUD is called with INDEX = 2. There are several reasons for this procedure. It is useful to distinguish the first calculation for a given t from subsequent calculations, since the values of certain quantities which do not depend on μ must be calculated only once, and thereafter may be read as stored variables. The value $\mu = 1$ is a convenient choice for isolation, since the Legendre polynomials $\{P_{\underline{\ell}}(1)\}$ of (30) are all equal to 1, and hence the recursive calculations used to generate the polynomials for other values of μ are in this case superfluous. In addition, at small penetrations (t \approx 1) convergence problems occur, and the subsequent method of calculation depends on the behavior of the Goudsmit-Saunderson series for $\mu = 1$ (see Section II.E).

When GOUD is called with INDEX = 2 (for $\mu = 1$) or INDEX = 3 (for $\mu \neq 1$), terms of the series (30) are evaluated from (47) and summed. The terms $\{G_{\underline{\ell}}(1)\}$ of (47) are calculated from (36)-(44) by subroutine GLG. The summation is terminated when twice in succession either the addition of a term changes the value of $A(\mu, t)$ by less than one part in 10⁵ or the absolute value of the term added is less than 10⁻⁷. The factors

$$(\ell + \frac{1}{2}) \exp\left[-\int_0^s G_{\ell}(s') ds'\right]$$

are stored, so that they need only be calculated once for a given t, regardless of the number of values of μ for which (30) is evaluated.

If one of the given convergence criteria is not satisfied after 150 terms have been summed and $\mu \leq \cos 5^{\circ}$, the transformed series for $A(\mu, t)$, Eq. (48), is summed, while if (30) fails to converge for $\mu > \cos 5^{\circ}$, the approximate method described in the paragraph following (48a) is used. The returned value of one of the parameters of GOUD named LT indicates the procedure followed. If $3 \leq LT \leq 150$, the series (30) converged, with LT being the number of terms required for convergence. If LT > 150, then (30) did not converge but (48) did, in LT-150 terms. If LT = 1, (30) failed to converge and the approximation following (48a) was used. Finally, if LT = 0, even the approximation failed, and the returned value of $A(\mu, t)$ is meaningless. E. Miscellaneous Routines: COMPLX, CONJGX, DEFINE, GAMCAR, GAUS, LEQ

COMPLX generates a complex number from two real arguments CONJGX calculates the conjugate of a complex number DEFINE performs preliminary bookkeeping chores GAMCAR calculates the gamma function of a complex argument GAUS performs numerical quadrature LEQ solves simultaneous linear algebraic equations

F. Restrictions and Limitations of the Code

The code may be used to evaluate the Goudsmit-Saunderson distribution for a single-element stopping medium having one of the following atomic numbers: Z = 3, 4, 6, 12, 13, 26, 29, 47, 50, 79, 82, 92. To perform the calculation for an atomic number not included in this list, a new function DENCOR must be written to generate the polarization correction $\delta(E)$ of Eq. (7). The details of the calculation of δ are given in refs. 7 and 9. To calculate the distribution for a stopping medium which is a compound or mixture of elements, TTAB must be rewritten so that a weighted average of singleelement stopping powers is calculated and integrated to obtain the path length vs. energy table. This calculation is also outlined in ref. 7.

Three factors place an upper bound on the initial energy E_0 . First, Function DENCOR generates $\delta(E)$ by interpolation on the curves of ref. 10, which apply to $E \leq 10$ MeV. For initial energies $E_0 > 10$ MeV, DENCOR must be rewritten as indicated in the previous paragraph. Second, the range-energy relations of Katz and Penfold [Eqs. (11)-(14)] apply only to $E \leq 20$ MeV, so that for $E_0 > 20$ MeV another range-energy relation must be found and written into the code. The third limitation results from the use of the approximate factor $1 + (Z\tau/1400)$ of Eq. (10) to estimate bremsstrahlung losses. An approximate upper limit on E_0 may be obtained by limiting the bremsstrahlung loss as calculated by the given expression to at most 10% of the collisional loss, which implies $ZE_0 < 72$, where E_0 is in MeV. For initial energies too high to satisfy the criterion, a more precise expression for the radiative energy loss should be used in subroutine DMUS (see, e.g., ref. 2).

It has been noted that the Goudsmit-Saunderson series fails to converge at low initial energies and small penetrations; in general, the lower the energy, the broader the range of penetrations in which this failure occurs. To insure that the series or the transformed series [Eq. (48)] converges down to penetrations small enough for angular deflections to be negligible, a lower bound should be imposed on the initial energy. $E_0 = 0.05$ MeV has been found adequate for this purpose.

Most of the program limitations apart from those inherent in the theoretical model are imposed by the dimensioning of certain arrays. The maximum number of intervals permitted in the generated t vs. E table is 150; in practice, no more than 100 are ever likely to be required for precision. If for some reason it is desired to increase this number, the dimensions of arrays E (or EG), TAU (or TAUG), and T (or TG) must be increased to the new maximum plus two in every routine in which Common BLOCKG occurs, the dimension of S must be similarly increased in Subroutine TTAB, and the number 150.0 in the second and third executable statements of TTAB should be changed to the new maximum.

Subroutine RMOTT, which generates the Mott cross section, provides for the evaluation of up to 90 terms of the series of Eqs. (25) and (26); if convergence to within five parts in 10^6 is not achieved in this many terms, a message to this effect is printed and the program is terminated. If it should ever prove necessary to increase the maximum number of terms, and this is most unlikely, the dimensions of complex arrays UF and UG in RMOTT must be changed to the desired maximum, and the terminal value of the DO 200 statement must be similarly altered.

The maximum number of terms provided by Subroutine GOUD for the evaluation of (30) and (48) for the Goudsmit-Saunderson distribution is 150. It is not recommended that this limit be increased, but if it is desired to do so, the dimension of GX (or GXG) in Common BLOCKG and the dimensions of GL and YX in Subroutine GOUD should be increased to 10 more than the new maximum, and the terminal indices of the DO 115, DO 150, and DO 490 statements in GOUD should be increased to the new maximum.

The parameter restrictions of the code are summarized below.

Atomic number: Z = 3, 4, 6, 12, 13, 26, 29, 47, 50, 79, 82, 92;Initial energy: $0.05 \le E_0 \le 10$ MeV; Combination of Z and E_0 : $ZE_0 < 72$.
IV. GUIDE TO THE USE OF THE CODE

A. Calling Program for Energy Loss and Angular Distribution Calculations

The user must supply FORTRAN statements and define variables in the order given. \emptyset refers to the letter, 0 to the number.

(1) Statement.

CØMMØN/BLØCKG/RHØG, ZG, AG, RHØZAG, ALPHG, ALPHSG, AIMG, NANGG, ANGG(37), ANGRG(37), CØSG(37), NSTEPG, EG(152), TAUG(152), TG(152), RTG, STG, NTG, N1G, N2G, F1G, F2G, F3G, F4G, F5G, F6G, F7G, IKG, C1G, C1XG, BMG(10,10), PG(10), GXG(160), JXG, JHG, JPG, JH1G, JH2G, IØPTG, ASCØRG, BSCØRG, DSCØRG, NCRVG

The user should take care not to define in his program any of the variable names in this common block except as subsequently indicated. The easiest way to do this is to avoid ending variable names with the letter G.

(2) Define or read in as data: NCRVG

The value of this variable, which specifies the stopping element, must be an integer from 1 to 11 inclusive. The possible values of NCRVG and the corresponding elements and atomic number are as follows: 1(Li,3), 2(Be,4), 3(Graphite, 6), 4(Mg,12), 5(Al,13), 6(Fe,26), 7(Cu,29), 8(Ag,47), 9(Sn,50), 10(Pb,82), 11(U,92).

(3) Statement: CALL DEFINE

DEFINE defines a number of constants for subsequent energy loss and

angular distribution calculations including RHOG, ZG, and AG, the density (g/cm^3) , atomic number and atomic weight respectively of the stopping medium. In addition, the routine generates an array ANGG, which contains angles 0°, 5°, 10°,...,175°,179°, another array ANGRG, which contains the same angles in radians, and a third array CØSG, which contains the cosines of the angles. The latter two arrays are needed in subsequent calculations, but any or all of them are also available to the user for other purposes, e.g., printed tabulation of data.

(4) Define:

EI: initial electron energy, MeV. See Section III.F. for bounds on this parameter.

EMIN: energy below which the electron is essentially at the end of its range. A suggested value is 0.01.

NSTEPG: number of intervals between EI and EMIN. A suggested value is 50, the maximum is 150.

(5) Statement: CALL TTAB (EI, EMIN)

TTAB generates a table of TG vs. TAUG (and EG), with EG varying from EI to EMIN in NSTEPG equal decrements, as well as STG, where

TG = t, defined by Eq. (2), dimensionless

EG = electron energy, MeV

TAUG = τ , defined by Eq. (1), dimensionless

 $STG = s_{+}$, the total path length corresponding to initial energy EI, cm.

The path length s for each tabulated value EG is not stored in common, but if desired may be calculated directly from STG (s_t) and TG (t) using Eq. (2), or printed out directly in TTAB. If only path length vs. energy data are required, the program terminates at this point.

(6) Statement: CALL GOUD (1, DUMMY, DUMMY, DUMMY, DUMMY)

This call of GOUD sets up the Goudsmit-Saunderson distribution calculation for initial energy EI, as discussed in Section III.D. The last four arguments of the calling statement are dummy variables, as the observant reader may have guessed.

(7) Statement: CALL RANGEN (1, RTG, EI, 1.0)

RANGEN here calculates RTG, the extrapolated range in cm corresponding to initial energy EI, using (11) or (13).

If the distribution is to be calculated for a given distance traveled in the initial direction of the electron's flight, a conversion must be made to obtain t, the fractional residual path length at the end of this distance. If t is known to begin with, the next two steps are omitted.

(8) <u>Define</u>: R, the straight-line penetration, or the distance in cm traveled in the initial direction of the electron's flight.

(9) Statement: CALL RANGEN (2, R, EN, TN)

RANGEN first calculates the energy EN possessed by the electron after it has traversed the straight-line distance R, and then interpolates in the t-E table generated in Step (5) to obtain the fractional residual path length TN = t(EN). The details of the calculation are described in Section III. B.

If Steps (8) and (9) are executed, the next two steps are omitted.

(10) <u>Define</u>: TN, the fractional residual path length at the point for which the distribution is to be calculated.

(11) Statement (optional): CALL RANGEN (4, R, EF, TN)

For the given fractional residual path length TN, RANGEN calculates EF, the energy at the end of the path, and then R, the penetration in the initial flight direction. The details of the calculation are given in Section III. B.

(12) Statement: CALL GOUD (2, TN, 1.0, AGS, LT)

The Goudsmit-Saunderson distribution calculation for the fractional residual path length TN is set up, and the value of the distribution at $\mu = 1$ ($\omega = 0$) is determined. Quantities calculated include:

AGS = $A(\mu = 1, t = TN)$, Eq. (30)

LT = number of terms of the series (30) needed for convergence to within one part in 10⁵. If LT = 1, the value of AGS is an approximation, good to within approximately 10%. If LT = 0, the value of AGS cannot be determined, but to a good approximation the distribution may be taken to be a delta function.

In addition the computer stores in common for each term of the series from L = 1 to L = LT

GXG(L) = expression (51).

These terms are functions of t but not of μ . In calculations for a given t, each of them is determined once and stored, and thereafter the stored values are used directly. The individual terms GXG(L), which as the Legendre expansion coefficients of the multiple scattering law are sometimes separately needed in the course of a calculation, are thus available to the user.

(13) <u>Statement:</u> IF(LT-1) n_1, n_2, n_3

Statement n_1 should begin a segment of the program based on the fact of the nonconvergence of the Goudsmit-Saunderson series for $\mu = 1$ and the failure of the approximation formulas. It may be assumed that the distribution is nearly a delta function peaking at $\mu = 1$. Statement n_2 should begin a segment which either uses or rejects the returned value of AGS, which has been generated by interpolation on the curves of Figure 1. If it is desired to use the value, n_2 should equal n_3 , the statement number of the next step, to which control proceeds if the Goudsmit-Saunderson series converges.

(14) <u>Define</u>: AMU = μ , a direction cosine for which A(μ ,t) is to be calculated.

(15) Statement: CALL GOUD (3, TN, AMU, AGS, LT)

This statement causes to be calculated AGS = A (μ = AMU, t = TN). The meaning of LT is as given following Step (12), with the addition that if LT > 150 the transformed series (48) was used to calculate AGS, and converged in LT-150 terms.

(16) Statement: IF (LT-1) n_1, n_2, n_3

 n_1 , n_2 , and n_3 are as defined following Step (13). Normally n_2 should equal n_3 , the continuation of the program.

Once Steps (1)-(3) have been executed, each calculation that involves a new initial energy must begin with Step (4); when a new t (or r) is introduced at the same initial energy, the calculation recommences with Step (10) [or Step (8)], and for different values of μ at a given t only Steps (14)-(16) are required.

B. User-Written Subroutines for Energy Loss and Angular Distribution Calculations

The user must supply the following subroutines in addition to those given in Section VI.

1. SUBROUTINE GAUS (FUNC, A, B, NORD, INT)

GAUS carries out the integration

$$\int_{A}^{B} FUNC(x)dx$$

by dividing the interval $A \le x \le B$ into INT equal subintervals, and performing

NORD-point quadrature on each subinterval. A SHARE subroutine of this type which uses Gaussian quadrature is available to institutions participating in this system; otherwise, the routine may be written using standard tables. ¹⁹ The code calls for NORD = 12 and 20 and INT = 1, so that a user-written routine need only provide for these values. If the user's version of GAUS does not provide for subdivision of the initial range (INT = 1 only), the defining statement should still contain five arguments, the fifth being a dummy integer. Any routine that calls GAUS should declare the name FUNC in an EXTERNAL statement.

2. SUBROUTINE LEQ (A, B, N, M, IA, IB, DET)

LEQ, which is part of many computer libraries, solves the matrix equation AX = B and evaluates the determinant of the matrix A. The subroutine parameters are defined as follows:

A is a two-dimensional array with dimensions (IA, IX), where both IA and IX are greater than or equal to N;

B is a one- or two-dimensional array with dimensions (IB, IY), where $IB \ge N$ and $IY \ge M$. If B is one-dimensional, IY may be omitted;

> N is the number of equations and unknowns; M is the number of vector solutions desired (=1 in the code); IA, IB are the first dimensions of arrays A and B respectively; DET contains the determinant of A.

The elements of the solution matrix X are stored in B(I,K), where I = 1, 2, ..., N and K = 1, 2, ..., M. The original A and B matrices are destroyed by the routine.

C. Calling Program for the Mott Cross Section

The values of the Mott to Rutherford cross section ratio $\sigma_m^{\prime}/\sigma_r^{\prime}$ needed for the Goudsmit-Saunderson distribution calculation are generated within the code. If it is desired to obtain σ_m^{\prime} in a separate calculation, the following procedure must be followed:

(1) Define:

Z = atomic number of the scattering medium

E = initial electron energy, MeV

(2) Define:

THE TA = center-of-mass scattering angle, radians

 $AMU = \cos(\theta)$

(3) Statement: CALL RMOTT (Z, E, AMU, THE TA, R, SIGM, SIGR, 1)

The calculated quantities are

SIGM =
$$\sigma_{m}(E, \theta)$$
, cm²
SIGR = $\sigma_{r}(E, \theta)$, cm²
R = σ_{m}/σ_{r}

The details of the calculation are described in Section III. C.

If cross sections are to be calculated for additional angles at the same Z and E, the next two steps are successively executed for each calculation. If Z or E changes, the calculation must begin with Step (1).

- (4) Define: THE TA, AMU [as in Step (2)]
- (5) Statement: CALL RMOTT (Z, E, AMU, THETA, R, SIGM, SIGR, 2)

If RMOTT is called for $\theta = 0$, it returns the correct value of R (=1), and does nothing about SIGM and SIGR, both of which are infinite. It should be clear that the numbers returned for these variables are in this case meaningless, and should not be used.

D. Storage Requirements

The code for the generation of the Goudsmit-Saunderson distribution, including the calling program GOUDS given in Section VI, occupies about 25,000 (octal) or 11,000 (decimal) locations.

V. ILLUSTRATIVE RESULTS

Typical results generated by the code are displayed in Tables 1-8, along with the times in seconds required to generate them. The latter figures represent central processor time on a CDC 6600 computer; a rough estimate of the equivalent calculating time on an IBM 7094 may be obtained by multiplying the given figures by 5.

The results of path length vs. energy calculations for two atomic numbers and a single initial energy are shown in Table 1. The time required for calculations of this sort is approximately $0.0012n_i$ seconds, where n_i is the number of intervals in the table.

Tables 2-4 list values of the Mott to Rutherford cross section ratio for several energies, angles, and atomic numbers. The time required for the calculation of a single value is approximately 0.05 seconds.

Tables 5-8 give representative values of the Goudsmit-Saunderson distribution $A(\omega, t)$ for two atomic numbers and two initial energies. The final energies E_f corresponding to the tabulated fractional residual path lengths t are shown, as are the values of LT for each calculation. To repeat from the discussion of Section III.D., if $3 \leq LT \leq 150$, the Goudsmit-Saunderson series (30) converged to within one part in 10^5 in LT terms; if LT > 150, the transformed series (48) converged in LT-150 terms; if LT = 1, the approximation given in the paragraph following (48a) was used, and if LT = 0, the value of A could not be determined.

A number of characteristic features of the distribution may be observed in these tables. At small values of t, corresponding to large penetrations, the distribution approaches isotropy: in the limit as t approaches 0, $A(\omega,t) \equiv 0.5$. As t increases, the distribution becomes increasingly biased in the initial direction of the electron's flight ($\omega = 0$), and the number of terms required to evaluate the series expression (30) for $A(\omega,t)$ also increases. As t approaches 1, (30) fails to converge and the transformed series (48) is summed instead: this occurs first at large angles and low initial energies, and eventually over the entire range of angles $\omega > 0$ and at all initial energies for t sufficiently close to 1.

The computing times given in Tables 5-8 do not include the times required to generate the required t-E tables (≈ 0.1 seconds per table), and to call GOUD with INDEX = 1 (≈ 0.2 seconds).

Figure 1 displays a collection of pairs $[A(\omega=0^{\circ}), A(\omega=5^{\circ})]$. The points shown represent data for Z = 13, 50, and 82, and a variety of fractional residual path lengths t. The fact that the curves shown fit the data as well as they do despite the range of energies and atomic numbers forms the basis for the approximation given after (48a).

VI. FORTRAN IV CODE

Listings follow of Programs GCAL and SIGMOT, which are sample calling programs for Goudsmit-Saunderson distribution and Mott cross section calculations respectively, and subroutines GOUD, DEFINE, TTAB, DENCOR, DSCAL, DMUS, HFIT, RANGEN, CGAM, G1G, GLG, RMOTT, RM1, GAMCAR, CMPLX, and CONJGX. The functions of each of these routines are described in Section III.

GCAL is set up so that a single set of parameters is read, the distribution is calculated, and the program terminated. If it is desired to run the program for several sets of data, the CALL EXIT statement (GCAL 051) should be changed to GO TO 25 (for a new stopping element), GO TO 30 (for the same stopping element and a new initial energy), or GO TO 40 (for the same stopping element and initial energy and a new penetration). SIGMOT is set up to read an atomic number Z, calculate the Mott to Rutherford cross section ratio for several energies and angles, read another Z, etc.

	PROGRAM GCAL(INPUT,OUTPUT) GCA	L 001
	COMMON/BLOCKG/RHOG,ZG,AG,RHOZAG,ALPHG,ALPHSG,AIMG,NANGG,ANGG(37),AGCA	L 002
	1NGRG(37),COSG(37),NSTEPG,EG(152),TAUG(152),TG(152),RTG,STG,NTG,N1GGCA	L 003
	2,N2G,F1G,F2G,F3G,F4G,F5G,F6G,F7G,IKG,C1G,C1XG,BMG(10,10),PG(10),GXGCA	004
:	3G(160),JXG,JHG,JPG,JH1G,JH2G,IOPTG,ASCORG,BSCORG,DSCORG,NCRVG GCA	_ 005
4	FORMAT(1H 10X3HE = F6.3,6X3HS = E10.3,6X3HT = E10.3) GCA	_ 006
25	READ 1,NCRVG GCA	_ 007
1	FORMAT(I10) GCA	_ 008
	CALL DEFINE GCA	009
2	FORMAT(2F10.0,I10) GCA	_ 010
30	READ 2,EI,EMIN,NSTEPG GCA	_ 011
	CALL TTAB(EI, EMIN) GCA	012
	CALL RANGEN(1,RTG,EI,1.0) GCA	_ 013
	PRINT 3,ZG,EI,STG,RTG GCA	_ 014
3	FORMAT(1H120X*ENERGY VS. FRACTIONAL RESIDUAL PATH LENGTH*//5X3HZ =GCA	015
	1F5.1,6X4HEI =F5.2,6X*TOTAL PATH LENGTH ST =*E10.3,6X*RANGE RT =*E1GCA	_ 016
	GCA GCA	017
	DO 35J=1,NTG GCA	_ 018
	S=STG*(1.0-TG(J)) GCA	_ 019
35	PRINT 4,EG(J),S,TG(J) GCA	_ 020
	PRINT 5,ZG,EI GCA	_ 021
5	FORMAT(1H110X*GOUDSMIT-SAUNDERSON DISTRIBUTION Z =*F5.1,5X4HEI =FGCA	022
	15•2/) GCA	023
	CALL GOUD (1, DUM, DUM, DUM) GCA	L 024
40	READ 8,TN GCA	_ 025
8	FORMAT(F10+0) GCA	026
	PRINT 17 GCA	_ 027
17	FORMAT(1H08X1HT9X2HEF8X5HOMEGA7X3HAGS7X5HTERMS) GCA	_ 028
	CALL RANGEN(4,R,EF,TN) GCA	_ 029
	CALL GOUD (2,TN,1.0,AGS,LT) GCA	_ 030
	PRINT 6,TN,EF,ANGG(1),AGS,LT GCA	_ 031
6	FORMAT(1H04XF7•4,3XF7•4,4XF7•1,3XE11•4,4XI3) GCA	_ 032
	DO 90J=2,3 GCA	_ 033
	CALL GOUD(3,TN,COSG(J),AGS,LT) GCA	_ 034
	IF(LT-150) 6:,62,63 GCA	_ 035
62	PRINT 7, ANGG (J), AGS, LT GCA	_ 036
7	FORMAT(1H 25XF7.1,3XE11.4,4XI3) GCA	_ 037
	GO TO 90 GCA	_ 038
63	LT=LT-150 GCA	_ 039
	PRINT 20,ANGG(J),AGS,LT GCA	_ 040
2Ŭ	FORMAT(1H 25XF7•1,3XE11•4,4XI3,*(YRW)*) GCA	_ 041
90	CONTINUE	_ 042
	DO 100J=4,NANGG,3 GCA	_ 043
	CALL GOUD(3,TN,COSG(J),AGS,LT) GCA	_ 044
93	IF(LT-150) 94,94,95 GCA	_ 045
94	PRINT 7, ANGG(J), AGS, LT GCA	_ 046
	GO TO 100 GCA	_ 047
95	LT=LT-150 GCA	_ 048
	PRINT 20, ANGG(J), AGS, LT GCA	- 049
100	GCA	- 050
	GCA GCA	. 051
	END GCA	. 052

SAMPLE CALLING PROGRAM FOR MOTT CROSS SECTION CALCULATION



	PROGRAM SIGMOT(INPUT,OUTPUT)	SIGM	001
	DIMENSION ANG(20), ANGR(20), COSA(20), E(10), BETAS(10), Q(10), SLAMB(10	SIGM	00
	1),R(10,20)	SIGM	005
	DATA NE, $(E(1), 1=1, 4)/4, 10, 00, 1, 0, 01/2$	SIGM	004
	DATA NANG, (ANG(J), J=2, 13)/13,5.0,10.0,15.0,30.0,60.0,90.0,120.0,12	SIGM	005
•		SIGM	006
2	FORMAT(FIU.0)	SIGM	007
3	FORMAT(INIZOX42HRATIO OF MOTI TO RUTHERFORD CROSS SECTIONS//4H Z =	SIGM	008
		SIGM	009
4	FORMAT(IH 5XIHEF7.2.9(4XF5.2))	SIGM	010
5	FORMAT(1H 5HTHETA/)	SIGM	011
8	FORMAT(1H010XF6.3,* SECONDS CP TIME FOR PRECEDING CALCULATION.*/)	SIGM	012
9	FORMAT(1H + 5, 1, 10(2x+7.4))	SIGM	013
	FACT=3.14159; //180.0	SIGM	014
	$ANG(1) = 0 \cdot 0$	SIGM	015
	$ANGR(1) = 0 \cdot 0$	SIGM	015
	$\cos(1) = 1 \cdot 0$	SIGM	017
	DO 21J=2, NANG	SIGM	018
	ANGR(J)=ANG(J)*FACT	SIGM	019
21	COSA(J)=COS(ANGR(J))	SIGM	020
30	READ 2,Z	SIGM	021
	IF(Z.EQ.0.0) CALL EXIT	SIGM	022
	PRINT 3,Z	SIGM	023
	$PRINT 4_{9}(E(I)_{9}I=1_{9}NE)$	SIGM	024
	PRINT 5	SIGM	025
	DO 60I=1,NE	SIGM	026
	J=1	SIGM	027
	CALL RMOTT(Z,E(I),COSA(J),ANGR(J),R(I,J),SIGM,SIGR,1)	SIGM	028
	DO 60J=2,NANG	SIGM	029
60	CALL RMOTT(Z,E(I),COSA(J),ANGR(J),R(I,J),SIGM,SIGR,2)	SIGM	030
	DO 65J=1,NANG	SIGM	031
	GO TO 30	SIGM	033
65	PRINT 9 , ANG(J), (R(I,J), I=1,NE)	SIGM	032
	END	SIGM	034

SUBROUTINES

	SUBROUTINE GOUD(INDEX,TX,COSX,AGS,LT)	GOUD	001
С	INDEX=1 FOR NEW EI, =2 FOR NEW TX,OLD EI, =3 FOR AGS(COSW,TX,EI)		
	COMMON/BLOCKG/RHO,Z,A,RHOZA,ALPHA,ALPHAS,AIMEAN,NANG,ANG(37),ANGR	(GOUD	002
	137), COSA(37), NSTEP, E(152), TAU(152), T(152), RT, ST, NT, N1, N2, F1, F2, F3	,GOUD	003
	2F4,F5,F6,F7,IK,C1,C1X,BM(10,10),P(10),GX(160),JX,JH,JP,JH1,JH2,IO	PGOUD	004
	3T,ASCOR,BSCOR,DSCOR,NCRV	GOUD	005
	DIMENSION G1(2),B(10),BI(10),PAI(5,7),PBI(5,7),PA(5,7),PB(5,7),GL	(GOUD	006
	1160),AF(7,3),YX(150),ATOL(15),AT5L(6,2)	GOUD	007
	DATA NT5,MR,MRP/6,3,4/	GOUD	800
	DATA(ATOL(J), J=1,6)/693147,2.30259,2.99573,3.91202,4.31749,4.60	5GOUD	009
	117/	GOUD	010
С	ATOL CONTAINS LOGS OF .5,10,20,50,75,100	GOUD	011
	DATA (AT5L(J), J=1,6)/693147, 2.27213, 2.91777, 3.71357, 4.00733, 4.1	5GOUD	012
	1888/	GOUD	013
С	AT5L(COL 1) CONTAINS LOGS OF •5,9•7,18•5,41,55,64	GOUD	014
	DATA (AT5L(J), J=7, 12)/693147, 2.25129, 2.86220, 3.59731, 3.87120, 4.	OGOUD	015
	12535/	GOUD	016
С	AT5L(COL 2) CONTAINS LOGS OF •5•9•5•17•5•36•5•48•56	GOUD	017
	GO TO (20,85,120), INDEX	GOUD	018
	20 LMAXT=0	GOUD	019
	NC5=1	GOUD	020
	$IF(E(1) \bullet L1 \bullet 0 \bullet 25) NC5=2$	GOUD	021
	CALL HFII(I)BEIAI)BEIASI)E(AI)QI)FP)	GOUD	022
		GOUD	025
		GOUD	024
	F D A = 1 + 0 + C F A I + C	COUD	025
	$E TADI = 1 \bullet 0 = (ETAI + ETAI) \bullet (SQRT(TSA) = 1 \bullet 0)$	GOUD	020
	D1(1)=C1A/(DETAS1*TAU(1)*(TAU(1)+2.0)) D1(2)=D1(1)*(D(1)+ED)	GOUD	021
	D(2) = D(1) = (1	GOUD	020
	40 BI(+1)=BI(1)*P()	GOUD	030
	$G1(1)=G1G(FTAI) \bullet FTABI \bullet PAI \bullet PBI \bullet BI \bullet F3X \bullet F4X \bullet F5X)$	GOUD	031
	I=N2	GOUD	032
	CALL HFIT(I)BETA, BETAS, ETA, Q, FQ)	GOUD	033
	F3Y=1.0+1.0/ETA	GOUD	034
	F4Y=1.0/(1.0+ETA)	GOUD	035
	$F5Y=1 \cdot 0 + ETA + FTA$	GOUD	036
	ETAB=1.0-(ET/+ETA)*(SQRT(F3Y)-1.0)	GOUD	037
	B(1) = C1X/(BETAS*TAU(I)*(TAU(I)+2.0))	GOUD	038
	B(2)=B(1)*(P(1)+FQ)	GOUD	039
	DO 55J=2,IK	GOUD	040
	55 B(J+1)=B(1)*P(J)	GOUD	041
	G1(2) = G1G(ETA)ETABPAPBPBF3YF4YF5Y)	GOUD	042
	H=JH	GOUD	043
	PAI(1,JH1)=((AJH+AJH+1.0)*(F5X*PAI(1,JH)-F4X)-(AJH+1.0)*PAI(1,JX)) GOUD	044
	1/AJH	GOUD	045
	PBI(1,JH1)=ETABI*PBI(1,JH)+PBI(1,1)	GOUD	046
	DO 75J=2,JH	GOUD	047
		GOUD	048
	ALF	GOUD	049
	JM1=J-1	GOUD	050
	PAI(J,LF)=F5X*PAI(JM1,LF)+PAI(JM1,1)-((ALF+1.0)*PAI(JM1,LF+1)+ALF	*GOUD	051
	1PAI(JMI,LF-1))/(ALF+ALF+1.0)	GOUD	052
	PBI(J)LF)=F5X*PBI(JM1)LF)+PBI(JM1)])-((ALF+1.0)*PBI(JM1)LF+1)+ALF		053
	IPBI(JM19L+-1))/(AL++AL++1+0)	GOUD	054
			055
		GOUD	057
		0000	1001

	JJ=J+J-1	GOUD	058
77	GL(2)=GL(2)+BI(JJ)*PAI(J,2)+BI(JJ+1)*PBI(J,2)	GOUD	05
	GO TO (79,81),IOPT	GOUD	060
79	GL(2)=GL(2)+FI(JP)*PAI(JH,2)	GOUD	061
	GO TO 83	GOUD	062
81	GL(2)=GL(2)+BI(JP)*PAI(JH,2)+BI(JP+1)*PBI(JH,2)	GOUD	063
83	FACT=1.0/(G1(1)-G1(2)*T(N2))	GOUD	064
	P(1)=G1(1)*G1(2)*T(N2)*(T(N2)-1•0)*FACT	GOUD	065
	P(2)=(G1(2)*T(N2)**2-G1(1))*FACT	GOUD	066
	RETURN	GOUD	067
85	FACT=(TX+P(2))/(TX*(1.0+P(2)))	GOUD	068
	COSMAX=-1.1	GOUD	069
	LMAX=0	GOUD	070
	NYX=0	GOUD	0/1
	AG5=0.0	GOUD	072
		GOUD	013
		COUD	074
00	IF(FACT) 90,90,95	GOUD	076
90		GOUD	077
		GOUD	078
		6000	079
		GOUD	080
	RETION	GOUD	081
95		GOUD	082
	FACT = -ST * P(1) * AIOG(FACT) / (G1(1) * P(2))	GOUD	083
	GX(1) = 1.5 * FXP(FACT*G1(1))	GOUD	084
	$GX(2)=2.5 \times EXP(FACT \times GL(2))$	GOUD	085
97	AGP=0.5+GX(1)+GX(2)	GOUD	086
	NTR=1	GOUD	087
	DO 115L=3,150	GOUD	088
	AL=L	GOUD	089
	IF(L-LMAXT) 100,100,98	GOUD	090
98	LMAXT=L	GOUD	091
	GL(L)=GLG(L,ETAI,ETABI,PAI,PBI,BI,F4X,F5X)	GOUD	092
100	IF(L-LMAX) 105,105,104	GOUD	093
104	GX(L)=(AL+0•5)*EXP(FACT*GL(L))	GOUD	094
	IF(GX(L).GT.1.UE6) GO TO 1150	GOUD	095
	LMAX=L	GOUD	096
105	AGS=AGP+GX(L)	GOUD	097
	IF(GX(L)/AGS.LE.1.0E-5) 109,110	GOUD	098
109	IF(NTR.EQ.2) GO TO 117	GOUD	099
	NTR=2	GOUD	100
	AGL=AGS	GOUD	101
		GOUD	102
110		GOUD	104
110		GOUD	104
1150	NCON=2	GOUD	106
11)0	NCONX = 1	GOUD	107
	COSW = COSA(3)	GQUD	108
	GO TO 400	GOUD	109
1151	IF(AGS.LT.0.!) GO TO 116	GOUD	110
	AG10=AGS	GOUD	111
	COSW=COSA(2)	GOUD	112
	NCON=3	GOUD	113
	GO TO 400	GOUD	114
1152	IF(AG10.LT.10.0.AND.AG10/AGS.LT.0.5) GO TO 116	GOUD	115
	AG5L=ALOG(AGS)	GOUD	116
	DO 1153J=2,NT5	GOUD	
	IF(AG5L+LE+AT5L(J+NC5)) 1155+1153	GOOD	118-

-	1153	CONTINUE	GOUD	119
		GO TO 116	GOUD	120
	1155	AG5=AGS	GOUD	121
		JM1=J-1	GOUD	122
		AGOL=ATOL(JM1)+(AG5L-AT5L(JM1,NC5))*(ATOL(J)-ATOL(JM1))/(AT5L(J,NC5))*	CGOUD	123
	1	15)-AT5L(JM1,NC5))	GOUD	124
		AGU=EXP(AGOL)	GOUD	125
		AGS=AG0	GOUD	126
		$ECON=(AGOL-AG5L)/(1 \cdot 0 - COSA(2) * * 2)$	GOUD	127
		LT=1	GOUD	128
		RETURN	GOUD	129
	116	NCONV=1	GOUD	130
	1161	LT=0	GOUD	131
		RETURN	GOUD	132
	117	LT=L	GOUD	133
		ADNM=AGS+AGL]-AGL-AGL	GOUD	134
		IF(ABS(ADNM).LE.1.0E-12) GO TO 118	GOUD	135
		AGS=(AGS*AGL1-AGL**2)/(AGS+AGL1-AGL-AGL)	GOUD	136
	118	AGO=AGS	GOUD	137
		RETURN	GOUD	138
	120	NCON=1	GOUD	139
		COSW=COSX	GOUD	140
		GO 10 (122,121), INDI	GOUD	141
	121	$IF(COSW \in Q \circ 1 \circ 0) = 97 \circ 123$	COUD	142
	122		GOUD	145
			GOUD	144
	122		GOUD	145
	125	IE(NCONV, E0, 1) GO TO 1161	GOUD	147
		$IF(NCONV \cdot EQ \cdot 1) GO TO 1101$	GOUD	148
	125	$IF(COSW_{\bullet} E_{\bullet}COSMAX) = 60 TO 400$	GOUD	149
	117	PL2=COSW	GOUD	150
		PL1=0.5*(3.0*PL2**2-1.0)	GOUD	151
		AGP=0.5+PL2*GX(1)+PL1*GX(2)	GOUD	152
		NTR=1	GOUD	153
		DO 150L=3,100	GOUD	154
		AL=L	GOUD	155
		PL=((AL+AL-1.0)*COSW*PL1-(AL-1.0)*PL2)/AL	GOUD	156
		IF(L-LMAXT) 132,132,130	GOUD	157
	130	LMAXT=L	GOUD	158
		GL(L)=GLG(L,ETAI,ETABI,PAI,PBI,BI,F4X,F5X)	GOUD	159
	132	IF(L-LMAX) 135,135,134	GOUD	160
	134	GX(L) = (AL+0.5) * EXP(FACT*GL(L))	GOUD	161
		IF(GX(L).GT.1.0E6) GO TO 491	GOUD	162
			GOUD	163
	135		GOUD	164
			GOUD	105
	100	1F(ABS(PL)-1.0E-6) 140,140,138	COUD	166
	130	IF (ABS(DAGS/AGS)+LE+I+UE=3+UR+ABS(DAGS)+LE+I+UE=7) 1399140	GOUD	167
	139	IF (NIR+EQ+2) GU 10 170	GOUD	160
			GOUD	170
			GOUD	171
		60 TO 145	GOHD	172
	140	NTR=1	GUID	173
	145	PL 2=PL 1	GOUD	174
		PL1=PL	GOUD	175
		AGP=AGS	GOUD	176
	150	CONTINUE	GOUD	177
	-	COSMAX=COSW	GOUD	178
-		GO TO 400	GOUD	179

170	LT=L	GOUD	180
171	ADNM=AGS+AGL1-AGL-AGL	GOUD	18
	IF(ABS(ADNM).LE.1.0E-12) GO TO 175	GOUD	182
	AGS=(AGS*AGL1-AGL**2)/ADNM	GOUD	183
175	GO TO(180,1151,1152,176),NCON	GOUD	184
176	COSW=COSX	GOUD	185
	AG5=AGS	GOUD	186
	ECON=ALOG(AG0/AG5)/(1.0-COSA(2)**2)	GOUD	187
_	GO TO 493	GOUD	188
180	RETURN	GOUD	189
400	AMCON=1+0/(1+0-COSW)**MR	GOUD	190
	AMTST=1.0E-7/AMCON	GOUD	191
	IF(NYX•GE•1) GO TO 441	GOUD	192
	AF (1,1)=0.5	GOUD	193
	AF(1,2) = GX(1)	GOUD	194
	AF(2,1)=0.5-0.33333333333333*GX(1)	GOUD	195
	AF(1,3)=GX(2)	GOUD	196
	AF(2,2) = AF(1,2) - 0.4 + AF(1,3) - AF(1,1)	GOUD	197
	K=2	GOUD	198
	AF(3,1) = AF(2,1) - 0.3333333333333338 + AF(2,2)	GOUD	199
	DO 440J=2,MR	GOUD	200
	JM1=J-1	GOUD	201
	K=K+1	GOUD	202
	AK=K	GOUD	203
	DO 4301=1,JM1	GOUD	204
	AF(1,1) = AF(1,2)	GOUD	205
	AF(1,2) = AF(1,3)	GOUD	206
	IF(1-EQ-1) 427,428	GOUD	207
427	AF(1,3) = GX(K)	GOUD	208
	GO TO 430	GOUD	209
428		GOUD	210
	AK=AK-1.0	GOUD	211
	AF(1,3)=AF(1M1,2)=(AK+1+0)*AF(1M1,3)/(AK+AK+3+0)=AK*AF(1M1,1)/(AK+	GOUD	212
4.00		COUD	213
430		GOUD	214
	$AF(J,3) = AF(JMI,2) = \bullet 4285/14285/1448F(JMI,3) = \bullet 6666666666666666666666666666666666$	COUD	215
	Ar (J+1, 2) = Ar (J, 2) + 0.4 + Ar (J, 3) - Ar (J, 1)	GOUD	210
	$IF(J \in Q \cap R)$ (J I) 440	GOUD	217
	AF (J+291) ≈AF (J+191) - U • 333333333333333*AF (J+192)	COUD	218
440		GOUD	219
	$1 \times 1 = AF(MRP) = 1$	6000	220
	YX(1) = AF(MRP)(2)	COUD	221
441		GOUD	222
		GOUD	223
		GOUD	225
	N R = 1	6000	225
		GOUD	220
		GOUD	228
	TE(1=NYX) 465-465-445	6000	220
445	$\frac{1}{1} \left(\frac{1}{1} + 1$	GOUD	230
449		GOUD	221
		GOUD	222
	$DU = 4001 \pm 19MR$	GOUD	232
	AF(1,2) - AF(1,2)	GOUD	234
	NI (192)-NI (193) TELT.EO.1 (452-454	GOUD	225
452	II VIELWEI / 77677474 TE/K_IMAYT) /577./577./571	GUUD	236
472	エレントーレー キャント・キャント キャント キャント キャント キャント キャント・キャント オート・キャント オート・キャント オート・キャント オート・キャント オート・キャント オート・オート オート オート・オート オート オート オート オート オート オート オート オート オート	6000	237
+921	$G[(K)=G[G(K)=TAI_FTARI_PAI_PRI_RI_FAY_F5Y)$	6000	238
4522	ΟΕικι-ΟΕΟικλεικελεικοιλεκαιλεστιλοτικά καλισκη ΤΕΙΚ-ΕΜΔΥΝ Δ525-Δ525-Δ524	GOUD	23
4526	$GX(K) = (\Delta K + \Omega_{\bullet}S) * FXP(FACT*CL(K))$	GOUD	240
サノムサ	UNINI-INNIUTIINNIUTINULINII	5500	

IF(GX(K).GT.1.0E6) GO TO 491 GOUD 241 GOUD 242 LMAX=K 4525 AF(1,3)=GX(K) GOUD 243 GOUD 244 GO TO 460 GOUD 245 454 IM1=I-1 $AK = AK - 1 \cdot 0$ GOUD 246 AF(I,3)=AF(IM1,2)-(AK+1.0)*AF(IM1,3)/(AK+AK+3.0)-AK*AF(IM1,1)/(AK+GOUD 247 GOUD 248 $1AK - 1 \cdot 0$ GOUD 249 460 CONTINUE YX(L)=AF(MR,;)-(AL+1.0)*AF(MR,3)/(AL+AL+3.0)-AL*AF(MR,1)/(AL+AL-1.GOUD 250 10) GOUD 251 GOUD 252 NYX=L 465 DAGS=YX(L)*PL GOUD 253 GOUD 254 AGS=AGP+DAGS IF(ABS(PL)-1.0E-6) 475,475,468 GOUD 255 468 IF(ABS(DAGS/AGS).LE.1.0E-5.0R.ABS(DAGS).LE.AMTST) 469,470 GOUD 256 469 IF(NTR.EQ.2) GO TO 495 GOUD 257 GOUD 258 NTR=2 GOUD 259 AGL=AGS AGL1=AGP GOUD 260 GOUD 261 GO TO 475 470 NTR=1 GOUD 262 GOUD 263 475 PL2=PL1 GOUD 264 PL1=PL AGP=AGS GOUD 265 490 CONTINUE GOUD 266 491 IF (NCON.GT.1.OR.NCONV.GT.0.OR.COSW.LE.COSA(2)) GO TO 116 GOUD 267 GOUD 268 IF(AG5.GT.0.0) GO TO 493 GOUD 269 NCON=4 GOUD 270 COSW = COSA(2)GO TO 400 GOUD 271 493 AGS=AG0*EXP(-ECON*(1.0-COSW**2)) GOUD 272 GOUD 273 LT=1GOUD 274 RETURN 495 LT=L+150 GOUD 275 GOUD 276 AGS=AGS*AMCON GOUD 277 AGL=AGL*AMCON AGL1=AGL1*AMCON GOUD 278 GO TO 171 GOUD 279 END GOUD 280 SUBROUTINE DEFINE **DEFI 001** COMMON/BLOCKG/RHO,Z,A,RHOZA,ALPHA,ALPHAS,AIMEAN,NANG,ANG(37),ANGR(DEFI 002 137), COSA(37), NSTEP, E(152), TAU(152), T(152), RT, ST, NT, N1, N2, F1, F2, F3, DEFI 003 2F4,F5,F6,F7,IK,C1,C1X,BM(10,10),P(10),GX(160),JX,JH,JP,JH1,JH2,IOPDEFI 004 3T, ASCOR, BSCOR, DSCOR, NCRV DEEI 005 DIMENSION ZT(12), AT(12), RHOT(12), AIMT(12) **DEFI 006** DATA(ZT(J), J: 1, 11)/3.0, 4.0, 6.0, 12.0, 13.0, 26.0, 29.0, 47.0, 50.0, 82.0, DEFI 007 192.0/ DEFI 008 DATA(RHOT(J), J=1,11)/0.53, 1.816, 2.26, 1.74, 2.7, 7.86, 8.92, 10.5, 7.31, DEFI 009 111.337,18.485/ DEFI 010 DATA(AT(J),J=1,11)/6.94,9.02,12.01,24.32,26.97,55.85,63.57,107.88,DEFI 011 1118.7,207.21,238.07/ DEFI 012 DATA(AIMT(J), J=1,11)/3.9E-5,6.41E-5,7.81E-5,1.56E-4,1.63E-4,3.37E-DEFI 013 **DEFI 014** 14,3.77E-4,6.60E-4,7.09E-4,1.18E-3,1.325E-3/ DATA SQRT2, PI/1.414214, 3.1415927/ DEFI 015 DATA NANG, NANG1, ANG(1), ANGR(1), COSA(1)/37, 36, 0.0, 0.0, 1.0/ DEFI 016 DATA IKL, NML, N1, N2/0,0,1,3/ **DEFI 017** N1 MUST BE 1, N2 MAY BE BETWEEN (AND INCLUDING) 2 AND 10 **DEFI 018 DEFI 019** Z=ZT(NCRV)

		A=AT(NCRV) RHO=RHOT(NCRV) AIMEAN=AIMT(NCRV) RHOZA=0.301314*RHO*Z/A C1=0.130548/AIMEAN**2 C1X=RHOZA*(Z+1.0) FACT=PI/180.0	DEFI DEFI DEFI DEFI DEFI DEFI	020 02 022 023 024 025 026
		DO 21J=2,NANG1 ANG(J)=ANG(J-1)+5.0 ANGR(J)=ANG(J)*FACT COSA(J)=COS(ANGR(J)) IF(ABS(COSA(J))-1.0E-5) 20,20,21	DEFI DEFI DEFI DEFI	027 028 029 030 031
	20 0	COSA(J)=0.0 CONTINUE ANG(NANG)=179.0 ANGR(NANG)=ANG(NANG)*FACT COSA(NANG)=COS(ANGR(NANG))	DEFI DEFI DEFI DEFI DEFI	032 033 034 035 036
с		IK=7 IK MUST NEVER BE LESS THAN 3 OR MORE THAN 9 JX=IK/2 JH=JX+1 JP=JX+JH	DEFI DEFI DEFI DEFI DEFI	037 038 039 040 041
	31	JH1=JH+1 JH2=JH+2 IF(JX-(IK+1)/2) 31,32,32 IOPT=2 GO TO 33 IOPT=1	DEFI DEFI DEFI DEFI	042 043 044 045 046 047
	33 /	ALPHA=Z/137•0 ALPHAS=ALPHA**2 Z2=Z*Z Z3=Z*Z2 ASCOR=SORT(Z)	DEFI DEFI DEFI DEFI	048 049 050 051 052
	É [F F	BSCOR=Z2*(1•0+2•0/Z)/(1•0+Z2/30•0) DSCOR=36•0*(Z3+200•0)/((Z3+3000•0)*(10•0+Z)) F6=Z**(2•0/3•0) F1=1•7E-5*F6 E2=3•76*ALPHAS	DEFI DEFI DEFI DEFI	053 054 055 056 057
	F F F	F3=PI*ALPHA/SQRT2 F4=0.160*F6 F7=3.33*ALPHA RETURN	DEFI DEFI DEFI DEFI	058 059 060 061
	3 (13	SUBROUTINE TTAB(EI,EMIN) COMMON/BLOCKG/RHO,Z,A,RHOZA,ALPHA,ALPHAS,AIMEAN,NANG,ANG(37),ANGR(37),COSA(37),NSTEP,E(152),TAU(152),T(152),RT,ST,NT,N1,N2,F1,F2,F3,	TTAB TTAB	001 002 003
	2 F 3 T E E	F4,F5,F6,F7,IK,C1,C1X,BM(10,10),P(10),GX(160),JX,JH,JP,JH1,JH2,IOF F,ASCOR,BSCOR,DSCOR,NCRV DIMENSION S(152) EXTERNAL DSCAL DATA FX/1.95703907/	TTAB TTAB TTAB TTAB TTAB	004 005 006 007 008
	10 N F 1 F 11 F	IF(NSTEP-150) 11,11,10 NSTEP=150 PRINT 1 FORMAT(1H0*TTAB HAS RESET NSTEP TO 150*) HSTEP=(EI-EMIN)/FLOAT(NSTEP)	TTAB TTAB TTAB TTAB TTAB	009 010 011 012 013
	N N N	NT=NSTEP+2 NORD=12	TTAB TTAB	015

	IND=1	TTAB	017
	E(1)=EI	TTAB	018
	TAU(1)=E(1)*FX	TTAB	019
	DO 20J=2,NST1	TTAB	020
	E(J)=E(J-1)-HSTEP	TTAB	021
	IF(E(J).LE.0.1.AND.IND.EQ.1) 15,16	TTAB	022
15	NC=J	TTAB	023
	IND=2	ТТАВ	024
16	TAU(J) = E(J) * FX	TTAB	025
20	CONTINUE	TTAB	026
30	E(NT)=0.0	TTAB	027
	TAU(NT)=0.0	TTAB	028
	S(1)=0.0	ΤΤΑΒ	029
	DO 50J=2,NT	ттав	030
	IF(J.EQ.NC) NORD=20	TTAB	031
	BMIN=TAU(J)	ТТАВ	032
	BMAX=TAU(J-1)	TTAB	033
50	S(J)=S(J-1)+GAUS(DSGAL,BMIN,BMAX,NORD,1)	TTAB	034
	ST=S(NT)	TTAB	035
	$T(1) = 1 \cdot 0$	ΤΤΑΒ	036
	FY=1.0/ST	ΤΤΑΒ	037
	DO 80J=2,NST1	TTAB	038
80	T(J)=FY*(ST-S(J))	TTAB	039
	T (NT) = O • O	ТТАВ	040
	RETURN	TTAB	041
	END	TTAB	042

FUNCTION DENCOR(TAUX, NCRV)	ENC	001
DENSITY CORRECTION FROM NELMS TABLES		
NCRV=1(LITHIUM-Z=3),2(BE-4),3(C-6),4(MG-12),5(AL-13),6(FE-26),7(CU-	29),	•
8(AG-47),9(SN-50 OR AU-79),10(PB-82),11(U-92)		
DIMENSION ET(15), DENLIM(11), DELC(15,11) D	ENC	002
DATA(ET(J), J: 1, 15)/0.2, 0.3, 0.4, 0.5, 0.6, 0.7, 0.8, 0.9, 1.0, 2.0, 3.0, 5.0D	ENC	003
1,7.0,8.5,10.(/ D	ENC	004
DATA(DENLIM(J), J=1,11)/0.0,0.05,0.04,0.03,0.05,0.05,0.035,0.015,0.D	ENC	005
103,0.0,0.0/ D	ENC	006
DATA (DELC(J), J=1,15) /.01,.10,.16,.22,.28,.348,.41,.47,.52,1.02,1D	ENC	007
1.45,2.03,2.51,2.83,3.1/ D	ENC	800
DATA (DELC(J), J=16,30)/.08,.145,.221,.3,.375,.45,.52,.59,.65,1.2,1D	ENC	009
1.64,2.3,2.79,3.1,3.35/ D	ENC	010
DATA (DELC(J), J=31,45)/.058,.1,.16,.22,.28,.345,.41,.47,.52,1.02,1D	ENC	011
1•43•2•03•2•5•2•8•3•05/ D	ENC	012
DATA (DELC(J), J=46,60)/.04,.042,.044,.048,.058,.075,.1,.126,.152,.D	ENC	013
1452,.77,1.25,1.64,1.88,2.1/ D	ENC	014
DATA(DELC(J), J=61,75)/.05,.051,.061,.09,.12,.158,.196,.236,.275,.6D	ENC	015
15,.99,1.52,1.92,2.17,2.35/ Di	ENC	016
DATA(DELC(J),J=76,90)/.05,.05,.055,.088,.088,.115,.145,.178,.210,.D	LNC	017
154,.83,1.31,1.68,1.93,2.15/ D	ENC	018
DATA(DELC(J), J=91, 105)/.035, 036, 038, 04, 05, 061, 08, 105, 13, 4D	ENC	019
13,.71,1.18,1.56,1.79,2.0/ D	ENC	020
DATA(DELC(J), J=106,120)/6*.015,.02,.029,.04,.21,.42,.76,1.05,1.25,D	ENC	021
11•43/ DI	ENC	022
DATA(DELC(J), J=121, 135)/3*•03, •033, •037, •04, •045, •052, •061, •15, •35D	ENC	023
1,.70,.95,1.15,1.25/ DI	ENC	024
DATA(DELC(J), J=136, 150)/9*0.0, 11, 23, 48, 71, 87, 1.01/	ENC	025
DATA(DELC(J), J=151, 165)/9*0.0, .16, .32, .60, .85, 1.02, 1.18/	ENC	026
11 E=•510976*TAUX DI	ENC	027
IF(E-ET(1)) 15,20 D	ENC	028
15 DENCOR=DENLIM(NCRV) DI	ENC	029
RETURN	ENC	030

	<pre>20 IF(E-10.0) 30,30,25 25 PRINT 1 1 FORMAT(1H010X*E IS TOO LARGE FOR DENCOR TO HANDLE*) STOP 30 D0 50J=3,15,2 IF(E-ET(J)) 55,55,50 50 CONTINUE 55 P=(E-ET(J-1))/(ET(J)-ET(J-1)) DENCOR=0.5*P*(P-1.0)*DELC(J-2,NCRV)+(1.0-P*P)*DELC(J-1,NCRV)+0.5*F 1*(P+1.0)*DELC(J,NCRV) RETURN END</pre>	DENC DENC DENC DENC DENC DENC DENC DENC	031 03 033 034 035 036 037 038 039 040 041 042
	FUNCTION DSCAL(TAU) CALL DMUS(TAU,DEDS) DSCAL=1.0/ABS(DEDS) RETURN END	DSCA DSCA DSCA DSCA DSCA	001 002 003 004 005
c	SUBROUTINE DMUS(TAUX, DEDS)	DMUS	001
C	CALCULATES DEDS=DE7DS FOR TAUX=E7.510976 COMMON/BLOCKG/RHO,Z,A,RHOZA,ALPHA,ALPHAS,AIMEAN,NANG,ANG(37),ANGR(137),COSA(37),NSTEP,E(152),TAU(152),T(152),RT,ST,NT,N1,N2,F1,F2,F3, 2F4,F5,F6,F7,IK,C1,C1X,BM(10,10),P(10),GX(160),JX,JH,JP,JH1,JH2,IOF 3T,ASCOR,BSCOR,DSCOR,NCRV CA=1.0/(TAUX+1.0)**2 CB=TAUX**2 BETAS=CA*TAUX*(TAUX+2.0) FMINUS=CA+CA*CB/8.0-(TAUX+TAUX+1.0)*CA*0.693147 X=BETAS*Z/ALPHAS	DMUS DMUS DMUS DMUS DMUS DMUS DMUS DMUS	002 003 004 005 006 007 008 009 010
C C	SCOR=ASCOR*(1.0-DSCOR/X**2)/(X+BSCOR) SCOR=SHELL CORRECTION DENCOR=DENSITY CORRECTION DEDS=-RHOZA*(ALOG(C1*CB*(TAUX+2.0))+FMINUS-SCOR-SCOR-DENCOR(TAUX,N ICRV))/BETAS DEDS=DEDS*(1.0+Z*TAUX/1400.0) RETURN END	DMUS DMUS DMUS IDMUS DMUS DMUS DMUS DMUS	011 012 013 014 015 016 017 018
	<pre>SUBROUTINE HFIT(N,BETA,BETAS,ETA,Q,FP) COMMON/BLOCKG/RHO,Z,A,RHOZA,ALPHA,ALPHAS,AIMEAN,NANG,ANG(37),ANGR(137),COSA(37),NSTEP,E(152),TAU(152),T(152),RT,ST,NT,N1,N2,F1,F2,F3, 2F4,F5,F6,F7,'(<c1,c1x,bm(10,10),p(10),gx(160),jx,jh,jp,jh1,jh2,iof 3T,ASCOR,BSCOI,DSCOR,NCRV BETAS=TAU(N)*(TAU(N)+2.0)/(TAU(N)+1.0)**2 BETA=SQRT(BETAS) ETA=F1*(1.13+F2/BETAS)/(TAU(N)*(TAU(N)+2.0)) Q=ALPHA/BETA FP=F3*BETA*CGAM(Q) IKM=IK-1 IKS=NANG/IKM JN=1 BM(1,1)=SQRT(1.0-COSA(JN)+ETA+ETA) CALL RMOTT(Z,E(N),COSA(JN),ANGR(JN),R,SIGM,SIGR,1) P(1)=R-1.0-FP*BM(1,1) D0 35J=2,IKM JN=JN+IKS BM(J,1)=SQRT(1.0-COSA(JN)+ETA+ETA)</c1,c1x,bm(10,10),p(10),gx(160),jx,jh,jp,jh1,jh2,iof </pre>	HFIT HFIT HFIT HFIT HFIT HFIT HFIT HFIT	001 002 003 004 005 006 007 008 009 010 011 012 013 014 015 016 017 019



35 50	CALL RMOTT(Z,E(N),COSA(JN),ANGR(JN),R,SIGM,SIGR,2) P(J)=R-1.0-FP*BM(J,1) BM(IK,1)=SQRT(1.0-COSA(NANG)+ETA+ETA) CALL RMOTT(Z,E(N),COSA(NANG),ANGR(NANG),R,SIGM,SIGR,2) P(IK)=R-1.0-FP*BM(IK,1) DO 50J=1,IK DO 50J=2,IK BM(J,I)=BM(J,I-1)*BM(J,1) CALL LEQ(BM,P,IK,1,10,10,DET) RETURN END	HFIT HFIT HFIT HFIT HFIT HFIT HFIT HFIT	020 021 022 023 024 025 026 027 028 029 030
	SUBROUTINE RANGEN(INDEX, RN, EN, TN)	RANG	001
	INDEX=1 FOR TOTAL RANGE RN FOR INITIAL ENERGY EN. (TN=1.0) INDEX=2,3,OR 4 RN IS THE PORTION OF THE RANGE TRAVERSED BETW AND FINAL ENERGY EN. INDEX=2 FOR EN(RN),TN(EN). INDEX=3 FOR EN(PN) INDEX=(FOR EN(TN), PN(EN).	EEN EI JUST	
	COMMON/BLOCK(/RHO,Z,A,RHOZA,ALPHA,ALPHAS,AIMEAN,NANG,ANG(37),AN	GRIRANG	002
	137),COSA(37);NSTEP,E(152),TAU(152),T(152),RT,ST,NT,N1,N2,F1,F2,	F3,RANG	003
	2F4,F5,F6,F7,1K,C1,C1X,BM(10,10),P(10),GX(160),JX,JH,JP,JH1,JH2,	IOPRANG	004
	3T, ASCOR, BSCOR, DSCOR, NCRV	RANG	005
10	GU TU (190,10,10,100),100/,100EX RX=RHO*(RT_RN)	RANG	008
10	IF(RX=0.95) 12.13.13	RANG	008
12	EN=EXP(6.629979035639-5.241090146751*SQRT(1.261848095652-0.3816	*ALRANG	009
	10G(RX)))	RANG	010
	IF(INDEX.EQ.2) 20,100	RANG	011
13	EN=1.6886/92453*RX+U.2	RANG	012
20	$DO_{30} = 2 \cdot NT$	RANG	014
20	IF(EN-E(J)) = 30,35,35	RANG	015
30	CONTINUE	RANG	016
35	TN=T(J)+(T(J-1)-T(J))*(EN-E(J))/HSTEP	RANG	017
100	RETURN	RANG	018
150	DO 160J=2,NT	RANG	019
140	IF(IN-I(J)) = 160,160,160	PANG	020
165		RANG	021
107	EN=E(J)+(E(JM1)-E(J))*(TN-T(J))/(T(JM1)-T(J))	RANG	023
	IF(EN-2.0) 170,175,175	RANG	024
170	RN=RT-•412*EN**(1•265-•0954*ALOG(EN))/RHO	RANG	025
	RETURN	RANG	026
175	RN=RT-(0.53*EN-0.106)/RHO	RANG	027
100	RETURN	RANG	028
105	$\frac{1}{2} \left(\frac{1}{2} + 1$	PANG	029
197	RETURN	RANG	031
200	$RN = (0.53 \times EN - (.106) / RHO$	RANG	032
	RETURN	RANG	033
	END	RANG	034
	FUNCTION CGAM(Q)	CGAM	001
	COMPLEX C1,C2,COMPLX,CONJGX,GAMCAR	CGAM	002
	C1 = GAMCAR(COMPLX(0.5,0))	CGAM	003
	CZ = GAMCAR (CUMPLR(1+0)(0)) $CGAM = PEAL (C2) CON (GY(C1))(C1+CON (C2)))$	CGAM	004
		СБАМ	005
	END	CGAM	007

	FUNCTION GIG(ETA, ETAB, PA, PB, B, FP, FQ, FR)	G1G	001
	COMMON/BLOCKG/RHU,Z,A,RHOZA,ALPHA,ALPHAS,AIMEAN,NANG,ANG,3//,ANGR 127).COSA(37).NSTEP.E(152).TAU(152).T(152).RT.ST.NT.N1.N2.E1.E2.E3	•616	003
	2F4, F5, F6, F7, IK, C1, C1X, BM(10, 10), P(10), GX(160), JX, JH, JP, JH1, JH2, IO	PGIG	004
	3T,ASCOR,BSCOR,DSCOR,NCRV	G1G	005
	DIMENSION PA(5,7),PB(5,7),B(10)	G1G	006
	PA(1,1) = ALOG(FP) - FQ	GIG	007
	PB(1,1)=2.0*(ETAB+ETAB)**1.5/(1.0+ETAB)	GIG	008
	$PA(1,2)=3 \cdot 0^*(FR*PA(1,1)-FQ)$		009
	PD(1)27 = PD(1)17 = 0 = 0 = 0	G1G	011
	AL=L-1	GIG	012
	PA(1,L)=((AL+AL+1.0)*(FR*PA(1.L-1)-FQ)-(AL+1.0)*PA(1.L-2))/AL	GIG	013
50	PB(1,L)=ETAB*PB(1,L-1)+PB(1,1)	G1G	014
	DO 60J=2,JX	GIG	015
		G1G	016
	PA(J,1)=(1+0+FK)*PA(JM1+1)=0+6666667*PA(JM1+2)	616	018
	F=JH=J+]	G1G	019
	DO = 60L=2.0 LF	GIG	020
	AL=L	GIG	021
	DFACT=1.0/(AL+AL+1.0)	GlG	022
	$PA(J_{l}) = FR*PA(JM1,L) + PA(JM1,1) - ((AL+1,0)*PA(JM1,L+1) + AL*PA(JM1,L-1)) + AL*PA(JM1,L-1) + AL*PA(JM$	1G1G	023
		GIG	024
60	PB(J9L)=FR*PB(JM19L)+PD(JM191)-((AL+100)*PD(JM19L+1)+AL*PB(JM19L* 3))*DFACT	G1G	025
	$PA(JH_{1}) = (1 \cdot 0 + FR) * PA(JX_{1}) - 0 \cdot 66666666666666666666667 * PA(JX_{2})$	G1G	027
	PB(JH,1)=(1.0+FR)*PB(JX,1)-0.666666666666666667*PB(JX,2)	GIG	028
	G1G=0.0	616	029
	DO 75J=1,JX	G16	030
		G1G	031
75	GIG=GIG+B(JJ)*PA(J,I)+B(JJ+I)*PB(J,I)	GIG	032
80	GU (U (BU,BD))10P) G1G±G1G+B(IP)+PA(IH,1)	G16	034
00	RETURN	G1G	035
85	G1G=G1G+B(JP)*PA(JH,1)+B(JP+1)*PB(JH,1)	GIG	036
	RETURN	G1G	037
	END	G1G	038
	FUNCTION GLG(L,ETA,ETAB,PA,PB,B,FP,FQ)	GLG	001
	COMMON/BLOCKG/RHO,Z,A,RHOZA,ALPHA,ALPHAS,AIMEAN,NANG,ANG(37),ANGR	GLG	002
	137) • COSA(37) • NSTEP • E(152) • TAU(152) • T(152) • KT • ST • NT • NT • NZ • F • F 2 • F 3	JGLG DGLG	003
	2F49F59F69F79TK9C19C1X96M1091039F110796X1160795X95H95F95H195H2910 3T+ASCOR+BSCOR+DSCOR+NCRV		004
•	DIMENSION $PA(5,7) \cdot PB(5,7) \cdot B(10)$	GLG	006
	AL2=L+JH-2	GLG	007
	PA(1,JH2)=((AL2+AL2+1.0)*(FQ*PA(1,JH1)-FP)-(AL2+1.0)*PA(1,JH))/AL	2 G LG	008
	PB(1,JH2)=ETAB*PB(1,JH1)+PB(1,1)	GLG	009
	DO 50J=2,JH	GLG	010
			012
	ALF=LF+L=3	GLG	013
	PA(J.LF) = FQ*PA(JM1.LF) + PA(JM1.1) ~ ((ALF+1.0) * PA(JM1.LF+1) + ALF*PA(J	MGLG	014
]	l1,LF-1))/(ALF+ALF+1.0)	GLG	015
50	PB(J,LF)=FQ*PB(JM1,LF)+PB(JM1,1)-((ALF+1.0)*PB(JM1,LF+1)+ALF*PB(JM1,LF+1)+ALF*PB(JM1,LF+1)+ALF*PB(JM1,LF)+ALF*P	MGLG	016
1	11,LF-1))/(ALF+ALF+1.0)	GLG	017
	GLG=0.0	GLG	018
			019

```
GLG
                                                                              021
 75 GLG=GLG+B(JJ)*PA(J,3)+B(JJ+1)*PB(J,3)
    GO TO (80,85), IOPT
                                                                         GLG
                                                                              022
 80 GLG=GLG+B(JP)*PA(JH,3)
                                                                         GLG
                                                                              023
                                                                              024
    GO TO 100
                                                                         GLG
 85 GLG=GLG+B(JP)*PA(JH,3) +B(JP+1)*PB(JH,3)
                                                                              025
                                                                         GLG
100 DO 110J=2,JH1
                                                                         GLG
                                                                              026
    LF = JH1 - J + 1
                                                                         GLG
                                                                              027
    DO 110I=1,LF
                                                                              028
                                                                         GLG
    PA(I,J) = PA(I,J+1)
                                                                              029
                                                                         GLG
110 PB(I,J) = PB(I,J+1)
                                                                         GLG
                                                                              030
    RETURN
                                                                         GLG
                                                                              031
    END
                                                                         GLG
                                                                              032
    SUBROUTINE RMOTT(Z,E,COSA,THETA,R,SIGM,SIGR,INDEX)
                                                                         RMOT 001
    Z=ATOMIC NUMBER, E=ENERGY(MEV), THETA=ANGLE(RADIANS), COSA=COS(THETA)
    INDEX=1 FOR NEW Z OR E, OTHERWISE=2
    SIGM=MOTT CROSS SECTION, SIGR=RUTHERFORD CROSS SECTION
   R IS THE CALCULATED VALUE OF SIGMA(MOTT)/SIGMA(RUTHERFORD)
    CODE USES ASYMPTOTIC FORMULA FOR THETA LE TEN DEGREES
   DATA PI, SQRT2, TOL/3.1415927, 1.414213562, 5.0E-6/
                                                                         RMOT 002
   COMPLEX C1,C2,C3,C4,C5,C6,F,G,DKF,DKG,F0,G0,UF0,UG0,AF(7,3),AG(7,3RMOT 003
   1),FCON,UF(100),UG(100)
                                                                         RMOT 004
   COMPLEX COMPLX, CONJGX, CEXP, GAMCAR
                                                                         RMOT 005
  1 FORMAT(1H010X*RMOTT HAS REPLACED 180 DEG. WITH 179 DEG. TO AVOID BRMOT 006
  1LOWUP .*/)
                                                                         RMOT 007
 3 FORMAT(1H0*SERIES IN RMOTT FOR THETA =*E11.4,* RADIANS DOES NOT CORMOT 008
   INVERGE IN 90 TERMS.*)
                                                                         RMOT 009
    GO TO (15,100), INDEX
                                                                         RMOT 010
 15 ALPHA=Z/137.0
                                                                         RMOT 011
    ALPHAS=ALPHA**2
                                                                         RMOT 012
                                                                         RMOT 013
    IF(Z-50) 16,16,17
16 MR=3
                                                                         RMOT 014
   MRP=4
                                                                         RMOT 015
   GO TO 19
                                                                         RMOT 016
17 MR=4
                                                                         RMOT 017
   MRP = 5
                                                                         RMOT 018
19 TAU=E/0.510976
                                                                         RMOT 019
   BETAS=TAU*(TAU+2.0)/(TAU+1.0)**2
                                                                         RMOT 020
   BETA=SQRT(BETAS)
                                                                         RMOT 021
   Q=ALPHA/BETA
                                                                         RMOT 022
   SLAMB=1.494E-21/(TAU*(TAU+2.0))
                                                                         RMOT 023
                                                                         RMOT 024
   C1=COMPLX(1 \cdot ( , Q))
   C2=GAMCAR(C1)
                                                                         RMOT 025
   FCON=CONJGX(C2)/C2
                                                                         RMOT 026
                                                                         RMOT 027
   RHOK = SQRT(1 \cdot 0 - ALPHAS)
   C4=COMPLX(RHOK,Q)
                                                                         RMOT 028
   C5=GAMCAR(C4)
                                                                         RMOT 029
   C6=C1*C2
                                                                         RMOT 030
   DKF=CEXP(COMPLX(0.0,-PI))*CONJGX(C2)/C6 - CEXP(COMPLX(0.0,-PI*RHOKRMOT 031
   1))*CONJGX(C5)/(C4*C5)
                                                                         RMOT 032
   DKG=DKF
                                                                         RMOT
                                                                              033
   CGQ=PI*ALPHA*BETA*CGAM(Q)/SQRT2
                                                                         RMOT 034
                                                                         RMOT 035
   AF(1,1)=DKF
                                                                         RMOT 036
   AG(1,1) = -DKF
                                                                         RMOT 037
   CALL RM1(1,ALPHAS,Q,DKF,DKG,C2,C4,C5,C6,AF(1,2),AG(1,2))
   AF(1,2) = -AF(1,2)
                                                                         RMOT 038
   AG(1,2) = -AG(1,2)
                                                                         RMOT 039
                                                                         RMOT 040
   AF(2,1)=AF(1,1)-0.33333333333333*AF(1,2)
   RMOT 041
   CALL RM1(2,ALPHAS,Q,DKF,DKG,C2,C4,C5,C6,AF(1,3),AG(1,3))
                                                                        RMOT 042
```

C C C C C

	$\Delta F(2 \cdot 2) = \Delta F(1 \cdot 2) - 0 \cdot 4 + \Delta F(1 \cdot 3) - A F(1 \cdot 1)$	RMOT	043
	$AG(2 \cdot 2) = AG(1 \cdot 2) - 0 \cdot 4 + AG(1 \cdot 3) - AG(1 \cdot 1)$	RMOT	04
	AF(3,1) = AF(2,1) - 0.3333333333333 + AF(2,2)	RMOT	04
	AG(3,1) = AG(2,1) - 0.3333333333333*AG(2,2)	RMOT	046
	SN=1.0	RMOT	047
	K=2	RMOT	048
		RMOT	049
		RMOT	050
		RMOT	051
		RMOT	052
		RMOT	053
	$\Delta F(1,1) = F(1,2)$	RMOT	054
	AG(T, T) = AG(T, T)	RMOT	055
	AG(1,1,1) - AG(1,1,2)	RMOT	056
	AG(1,2) = AG(1,2)	RMOT	057
	$F(1, F_0, 1) = 28$	RMOT	058
27	$\begin{bmatrix} c_{1} & c_{2} & c_{3} & c_{3} \\ c_{1} & c_{2} & c_{3} & c_$	RMOT	059
21		RMOT	060
		RMOT	061
		RMOT	062
	AG(1,5)=AG(1,5)*SN	DMOT	062
~ ^ ^		DMOT	005
28		RMOT	064
		- DMOT	005
	AF(1,3)=AF(1M1,2)=(AL+1+0)=AF(1M1,3)/(AL+AL+3+0)=AL*AF(1M1,1)/(AL*	DMOT	060
	$1 \mathbf{A} = 1 \cdot 0$	RMOT	067
	AG(1,3) = AG(1M1,2) - (AL+1,0) * AG(1M1,3) / (AL+AL+3,0) - AL*AG(1M1,1) / (AL*AL+3,0) - AL*AG(1M1,1)		060
•		RMOT	070
30	CONTINUE	RMOT	070
	AF (J,3)=AF (JM1,2)4285/14285/14*AF (JM1,3)666666666666666667*AF (JM1,1)	RMOT	071
	AG(J,3)=AG(JM1,2)4285/14285/14*AG(JM1,3)6666666666666666/*AG(JM1,1)	RMOT	072
	AF(J+1,2) = AF(J,2) - 0.4 + AF(J,3) - AF(J,1)	RMOT	073
	AG(J+1,2) = AG(J,2) - 0.4 + AG(J,3) - AG(J,1)	RMOI	074
	IF(J.EQ.MR) GO TO 40	RMOT	075
	AF(J+2,1)=AF(J+1,1)-0.33333333333*AF(J+1,2)	RMOT	076
	AG(J+2,1)=AG(J+1,1)-0.333333333333*AG(J+1,2)	RMOT	077
40	CONTINUE	RMOT	078
	JMAX=1	RMOT	079
	UFO=AF(MRP,1)	RMOT	080
	UG0=AG(MRP+1)	RMOT	081
	UF(1)=AF(MRP,2)	RMOT	082
	UG(1) = AG(MRP, 2)	RMOT	083
100	IF(COSA-1.0) 110,105,105	RMOT	084
105	R=1•0	RMOT	085
	RETURN	RMOT	086
110	SIGR=1.4943E-21*ALPHAS*(1.0-BETAS)/(BETAS**2*(1.0-COSA)**2)	RMOT	087
	IF(THETA-0.175) 115,115,120	RMOT	088
115	R=1.0+CGQ*SQRT(1.0-COSA)	RMOT	089
	SIGM=R*SIGR	RMOT	090
	RETURN	RMOT	091
120	IF(COSA+0.99999) 125,125,130	RMOT	092
125	PRINT 1	RMOT	093
	THETA = 179.0*PI/180.0	RMOT	094
	COSA = COS(THETA)	RMOT	095
120	ANG=0.5*THETA	RMOT	096
<u> </u>	SINS=SIN(ANG)**2	RMOT	097
	SFCS=1.0/(1.(-SINS))	RMOT	098
	E0====================================	RMOT	099
	GO=COMPLY(0_0_=O/(SINS*SECS))*E0	RMOT	100
		RMOT	101_
	G = G G + U G (1) * C O S A	RMOT	10
	FLAST=CABS(F)	RMOT	103

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	GLAST=CABS(G)	RMOT	104
	РК2=1•0	RMOT	105
	PK1=COSA	RMOT	106
		PMOT	107
		DMOT	107
	00 2000=2,90	RMOT	108
		RMOT	109
	JM1=J-1	RMOT	110
	PK = ((A + A - 1 - 0)) * COSA * PK 1 - (A - 1 - 0) * PK 2) / A - 1	RMOT	111
		PMOT	112
• • •	11 (J-JMAX) 100,100,140	RMOT	112
140	JMAX=J	RMOT	113
	K=K+1	RMOT	114
	AL=K	RMOT	115
	DO 155I=1,MR	RMOT	116
	$AF(I \bullet I) = AF(I \bullet 2)$	RMOT	117
	AG(1,1) = AG(1,2)	RMOT	118
		DMOT	110
	AF(1,2) = AF(1,2)	RMOT	119
	AG(1,2) = AG(1,3)	RMOT	120
	IF(I•EQ•1) 152,154	RMOT	121
152	CALL RM1(K,ALPHAS,Q,DKF,DKG,C2,C4,C5,C6,AF(1,3),AG(1,3))	RMOT	122
	SN=-SN	RMOT	123
	AF(1,3) = AF(1,3) * SN	RMOT	124
	AG(1, 3) - AG(1, 3) + SN	PMOT	125
		DMOT	122
	60 10 155	RMOT	126
154	1M1=1-1	RMOT	127
	AL=AL-1.0	RMOT	128
	AF(I,3)=AF(IM1,2)-(AL+1.0)*AF(IM1,3)/(AL+AL+3.0)-AL*AF(IM1,1)/(AL+RMOT	129
	1AL-1.0)	RMOT	130
	AG(I,3) = AG(II,1,2) - (AI,+1,0) * AG(IM1,3) / (AI,+AI,+3,0) - AI * AG(IM1,1) / (AI,+AI,+3,0) - AI * AG(IM1,0) / (AI,+3,0) / (AI,+3,0	AL +RMOT	131
		RMOT	122
155		DMOT	122
100	CONTINUE	RMOT	200
	UF(J)=AF(MR,2)-(AJ+1•0)*AF(MR,3)/(AJ+AJ+3•0)-AJ*AF(MR,1)/(AJ+AJ	-1.RMOT	134
	10)	RMOT	135
	UG(J)=AG(MR+2)-(AJ+1+0)*AG(MR+3)/(AJ+AJ+3+0)-AJ*AG(MR+1)/(AJ+AJ	-1.RMOT	136
	10)	RMOT	137
160	F=F+UF(J)*PK	RMOT	138
	G=G+UG(-1)*PK	RMOT	139
		PMOT	140
		RMOT	140
	GABS=CABS(G)	KWO I	141
	IF(ABS(PK)-1.0E-6) 176,176,165	RMOT	142
165	IF(ABS(FABS-FLAST)/FABS-TOL) 170,170,175	RMOT	143
170	IF(ABS(GABS-GLAST)/GABS-TOL) 171,171,175	RMOT	144
171	IF (NTRY-EQ.2) GO TO 210	RMOT	145
	NTRY=2	RMOT	146
	60 176	PMOT	147
176		RMOT	141
1/2		RMUT	148
1/6	FLAST=FABS	RMOT	149
	GLAST=GABS	RMOT	150
	PK2=PK1	RMOT	151
200	РКІ=РК	RMOT	152
-	DRINT 3. THETA	PMOT	153
		DMOT	156
0.1/		KMU1	104
210		RMOT	155
	F = C3*F	RMOT	156
	G=C3*G	RMOT	157
	FABS=CABS(F)	RMOT	158
	GABS=CABS(G)	RMOT	150
	CICM -CIANR*(N**2*(1_A_RFTAC)*FADC**2)(INC + CARC**2*CC)	DMOT	160
	DECTAGE CONCERNENCE CONCERNENCERNENCERNENCERNENCERNENCERNENCERNENCERNENCERNENCERNENCERNENC		100
		KMOT	101
	RETURN	RMOT	162
	END	RMOT	163

<pre>SUBROUTINE RM1(K,ALPHAS,Q,DKF,DKG,C2,C4,C5,C6,AF,AG) COMPLEX DKF,DKG,DK1F,DK1G,C2,C4,C5,C6,AF,AG COMPLEX COMPLX,CONJGX,GAMCAR,CEXP AK1=K+1 RHOK=SQRT(AK1**2-ALPHAS) C2=C6 C4=COMPLX(RHOK,Q) C5=GAMCAR(C4) C6=COMPLX(AK1,Q)*C2 DK1F=AK1*(CEXP(COMPLX(0.0,-3.1415927*AK1))*CONJGX(C2)/C6 - CEXP(C 1MPLX(0.0,-3.1415927*RHOK))*CONJGX(C5)/(C4*C5)) DK1G=AK1*DK1F AF=DKF+DK1F AG=DKG-DK1G DKF=DK1F DKG=DK1G RETURN END</pre>	RM1 RM1 RM1 RM1 RM1 RM1 RM1 RM1 RM1 RM1	001 003 004 005 006 007 008 009 010 011 012 013 014 015 016 017 018
COMPLEX FUNCTION GAMCAR(CMX) COMPLEX CMX,COMX,GAMX,E1,E11,TX,GAM1,CLOG,CEXP COMX=CMX-1. E1=COMX+05 TX=COMX+5.5 E11=E1*CLOG(TX) GAM1=CEXP(E11-TX)*2.50662827465 GAMX=GAM1*(1.+76.18009173/(COMX+1.)-86.50532033/(COMX+2.) G+24.01409822/(COMX+3.)-1.231739516/(COMX+4.)+.120858003E-2/(COMX+ G5.)536382E-5/(COMX+6.)) GAMCAR=GAMX RETURN END	GAMC GAMC GAMC GAMC GAMC GAMC GAMC GAMC	0001 0002 0003 0004 0005 0006 0007 0008 0009 0010 0011 0012 0013
COMPLEX FUNCTION COMPLX(A,B) COMPLX=A+(0.0,1.0)*B RETURN END	COMP COMP COMP COMP	001 002 003 004
COMPLEX FUNCTION CONJGX(C) COMPLEX C XR=REAL(C) XIM=AIMAG(C) CONJGX=XR-(0.0,1.0)*XIM RETURN END	CONJ CONJ CONJ CONJ CONJ CONJ	001 002 003 004 005 006 007

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	Z =	13	$\mathbf{Z} = 82$				
E (MeV)	s (cm)	t	s (cm)	t			
0。50	0.000	1.000	0。000	1,000			
0.48	4.539×10 ⁻³	9.454×10 ⁻¹	1。672×10 ⁻³	9.486×10 ⁻¹			
0.46	9。040 ''	8.912 ''	3,337 ''	8.974 ''			
0.44	$1,350 \times 10^{-2}$	8.376 "	4。992 ''	8。465 ''			
0.42	1.791 "	7.845 "	6.637 ''	7.959 "			
0.40	2.228 "	7。320 ''	8 .271 ''	7。456 ''			
0.38	2.658 "	6。801 ''	9.890 "	6.958 ''			
0.36	3.083 ''	6.291 "	1。149×10−2	6.465 "			
0.34	3。501 ''	5.788 "	1.308 "	5.977 ''			
0.32	3.911 "	5.294 "	1.464 "	5.496 ''			
0.30	4.313 ''	4.810 "	1.619 ''	5.022 "			
0.28	4. 706 ''	4.337 "	1.770 "	4。556 ''			
0.26	5.089 "	3.877 "	1.918 "	4。100 ''			
0.24	5。461 ''	3。429 ''	2.063 "	3。654 ''			
0.22	5.820 "	2。997 ''	2.204 ''	3.220 "			
0.20	6.165 "	2.582 "	2.341 "	2。800 ''			
0.18	6。495 ''	2.185 "	2.472 ''	2.396 ''			
0.16	6.808 ''	1.809 "	2.598 ''	2.010 ''			
0.14	7.100 "	1.456 "	2.717 ''	1.644 "			
0.12	7.371 "	1.131 "	2.828 "	1.302 "			
0.10	7.617 "	8.348 $\times 10^{-2}$	2.931 "	9.867 $\times 10^{-2}$			
0.08	7.834 "	5.733 "	3。023 ''	7.032 "			
0。06	8.019 ''	3.511 "	3.103 "	4.566 ''			
0.04	8 .166 ''	1.748 "	3.169 "	2.539 "			
0.02	8.267 "	5。292×10 ⁻³	3.218 "	1.043 ''			
	$s_t(E_0) = 0.6$	08311 cm	$s_t(E_0) = 0.03251 \text{ cm}$				
	$r_t(E_0) = 0.0$	06065 cm	$r_t(E_0) = 0.$	01444 cm			

Table 1. Path Length $s(E_0, E)$, Fractional Residual Path Length t(E), Total Path Length $s(E_0, 0)$, and Range $r_t(E_0)$ for Al (Z=13) and Pb (Z=82), $E_0 = 0.5$ MeV

0.12 seconds computing time

θ (degrees)	10.0 MeV	1.0 MeV	0.1 MeV	<u>0.01 MeV</u>
0	1.000	1.000	1.000	1.000
5	1.0126	1.0118	1.0064	1.0014
10	1.0251	1.0235	1.0128	1.0027
15	1.0192	1.0191	1.0151	1.0053
30	0.9969	1.0009	1.0166	1.0109
60	0.8375	0.8611	0.9771	1.0141
90	0.5752	0.6278	0.8975	1.0072
120	0.2939	0.3769	0.8074	0.9963
150	0.0813	0.1868	0.7381	0.9873
165	0.0225	0.1343	0.7188	0.9848
170	0.0114	0。1243	0.7151	0.9843
175	0.0047	0.1183	0.7129	0.9840
179	0,0025	0.1164	0.7122	0.9839

Table 2. Ratio of Mott to Rutherford Cross Sections: Z = 13

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1.6 seconds computing time

θ (degrees)	10.0 MeV	1.0 MeV	0.1 MeV	0.01 MeV
0	1.0000	1.0000	1.0000	1.0000
5	1.0334	1.0303	1.0111	1.0013
10	1.0668	1.0606	1.0221	1.0026
15	1.1214	1.1123	1.0408	1.0057
30	1.2345	1.2226	1.1127	1.0006
60	1.2736	1.2764	1.2280	1.0162
90	1.0126	1.0548	1.2261	1.0969
120	0.5693	0.6639	1.1421	1.2207
150	0.1651	0.3048	1.0527	1.3450
165	0.0451	0.1980	1.0255	1.3840
170	0.0220	0.1774	1.0203	1.3917
175	0.0080	0.1649	1.0171	1.3963
179	0.0035	0.1609	1.0161	1.3978

Table 3. Ratio of Mott to Rutherford Cross Sections: Z = 50

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2.2 seconds computing time

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10.0 MeV	1.0 MeV	0.1 MeV	0.01 MeV
1.0000	1.0000	1.0000	1.0000
1.0366	1.0326	1.0107	1.0013
1,0732	1.0651	1.0214	1.0026
1,1250	1.1066	1.0186	1.0050
1.3583	1.3181	1.0486	1.0098
1.8737	1.8199	1.3394	0.9825
1.8930	1.8927	1.6466	0,9529
1.2465	1.3660	1.8073	1.3127
0.3940	0。6 42 4	1.8703	1.9177
0.1087	0.3984	1.8832	2.1335
0,0523	0.3501	1.8855	2.1766
0.0180	0.3206	1.8869	2.2030
0.0069	0.3111	1.8873	2.2115
	10.0 MeV 1.0000 1.0366 1.0732 1.1250 1.3583 1.8737 1.8930 1.2465 0.3940 0.1087 0.0523 0.0180 0.0069	10.0 MeV1.0 MeV1.00001.00001.03661.03261.07321.06511.12501.10661.35831.31811.87371.81991.89301.89271.24651.36600.39400.64240.10870.39840.05230.35010.01800.32060.00690.3111	10.0 MeV1.0 MeV0.1 MeV1.00001.00001.00001.03661.03261.01071.07321.06511.02141.12501.10661.01861.35831.31811.04861.87371.81991.33941.89301.89271.64661.24651.36601.80730.39400.64241.87030.10870.39841.88320.05230.35011.88550.01800.32061.88690.00690.31111.8873

Table 4. Ratio of Mott to Rutherford Cross Sections: Z = 80

2.5 seconds computing time

	Table	<u>), G</u>	Judshin	-Saulderson Distribution: $\Sigma = 13$, $E_0 = 1.0$ MeV					U Mev		
	t = 0.1 $E_{f} = 0.1928$		t = 0.5 $E_f = 0.5786$		t = 0.9 $E_f = 0.9167$			t = 0.99 $E_f = 0.9917$			
ω (degrees)	A	<u></u>	LT	A		LT	A		LT	A	$\underline{\mathbf{LT}}$
0	5.207×	10 ⁻¹	4	$1.073 \times$	10 ⁰	6	7.656×	10 ⁰	17	1.242×10^{2}	73
5	5.206	**	4	1.070	**	6	7.415	11	17	6.929×10^{1}	6 2
10	5.204	**	4	1.061	**	6	6.742	**	14	1.697 "	67
15	5.200	* *	4	1.046	**	6	5.763	**	17	3.447×10^{0}	80
30	5.179	**	4	9.691×	10-1	5	2.584	ŦT	16	1.281×10^{-1}	94
60	5.103	**	4	7.194	**	6	2.185×	10-1	18	6.284 $\times 10^{-3}$	191*
90	5.000	* *	6	4.522	**	8	2. 448×	10 ⁻²	22	1.096 "	186*
120	4.896	11	4	2.565	**	6	5.476×	10 ⁻³	21	2.896×10^{-4}	168*
150	4.821	**	4	1.508	**	6	1.85 2	11	21	9.367×10 ⁻⁵	166*
165	4.800	**	4	1.268	**	7	1.301	**	21	6.128 "	168*
180	4.793	11	4	1.190	**	7	1.140	TT	22	5.151 "	169*
Computing tin seconds	^{ne} , 0.	04			0.05			0,06	\$	0.11	

Table 5 Coudemit-Saunderson Distribution 1 0 1 4017 19 1.1

*Transformed series converged in LT-150 terms.

	Table	<u>6.</u> Go	oudsmit	-Saunder	son Di	stribu	tion; Z	<u>= 13, E</u>	$E_0 = 0.1$	1 MeV	
	t = 0 E _f =	t = 0.1 $E_f = 0.0261$		t = 0.5 $E_{f} = 0.0664$		t = 0.9 $E_{f} = 0.0939$		t = 0.99 $E_f = 0.0994$			
ω (degrees)	A		LT	A		$\underline{\mathbf{LT}}$	A		LT	A	\underline{LT}
0	5.114×	10 ⁻¹	4	9.077×3	10-1	6	6.706×	10 ⁰	19	**	0**
5	5.114	**	4	9.057	**	6	6.499	**	18	$6.169{ imes}10^{1}$	24 8*
10	5.113	**	4	8.996	**	6	5.920	11	17	1.380 "	217*
15	5.110	**	4	8.896	**	6	5.090	**	18	3,650×10 ⁰	193*
30	5.099	**	4	8.378	**	5	2.432	**	17	2.210×10^{-1}	190*
60	5.058	* *	4	6.662	11	6	$3.142 \times$	10 ⁻¹	20	1.292×10^{-2}	186*
90	5.000	**	6	4.728	11	8	5.449×	10 ⁻²	24	2.823×10 ^{−3}	180*
120	4,943	**	4	3,200	**	6	1.751	**	24	1,106 "	162*
150	4.901	**	4	2.303	11	6	9.267×	10 ⁻³	23	6.474×10 ⁻⁴	161*
165	4.890	* *	4	2.088	**	6	7.928	**	25	5.668 "	160*
180	4.886	**	4	2.017	**	7	7.529	**	27	5.424 "	163*
Computing times seconds	me, 0.	04			0。05			0,06		0.11	

*Transformed series converged in LT-150 terms.

ries failed to converge. *1

	t = 0.1 $E_{f} = 0.1928$		t = 0.5 $E_{f} = 0.5786$		t = 0.9 $E_{f} = 0.91$	67	t = 0.99 $E_{f} = 0.9917$	
ω (degrees)	A	\mathbf{LT}	A	$\underline{\mathbf{LT}}$	A	LT	A	\mathbf{LT}
0	5.000×10 ⁻¹	4	5.001×10 ⁻¹	4	9.917×10 ⁻¹	7	1.529×10^{1}	29
5	**	4	**	4	9.890 "	7	1.413 "	25
10	**	4	**	4	9.809 "	7	1.125 "	27
15	**	4	**	4	9.675 "	7	7。891×10 ⁰	2 8
30	**	4	**	4	8,995 ''	6	1.781 "	29
60	**	4	**	4	6.833 "	6	1.326×10^{-1}	35
90	**	6	5.000×10 ⁻¹	6	4.580 ''	8	3.078×10^{-2}	38
120	**	4	11	4	2.952 "	6	1.039 "	39
150	**	4	4。999×10 ⁻¹	4	2.063 "	7	3.945×10^{-3}	40
165	**	4	11	4	1.856 "	7	2,712 "	41
180	**	4	11	4	1.789 "	7	2,329 "	43
Computing time seconds	e, 0.04		0.04		0.0	5	0.09	

Table 7. Goudsmit-Saunderson Distribution: Z = 82, $E_0 = 1.0$ MeV

Table 8. Goudsmit-Saunderson Distribution: Z = 82, $E_0 = 0.1$ MeV

	t = 0.1 $E_{f} = 0.0190$		t = 0.5 $E_{f} = 0.0632$		t = 0.9 $E_{f} = 0.0933$			t = 0.99 $E_{f} = 0.0993$	
ω (degrees)	A	LT	A	\underline{LT}	A		$\underline{\mathbf{LT}}$	A	LT
0	5.000×10^{-1}	4	5.000×10 ⁻¹	4	8.778×3	10-1	8	2.460×10^{1}	1**
5	11	4	**	4	8.755	11	8	2.066 "	245*
10	**	4	**	4	8.687	**	8	1.115 "	200*
15	**	4	**	4	8,574	11	8	6.269×10 ⁰	194*
30	**	4	**	4	8.011	**	8	1.392 "	166*
60	**	4	**	4	6.308	11	8	1 _℃ 733×10 ⁻¹	165*
90	**	6	**	6	4.652	**	10	5 _° 285×10-2	166*
120	**	4	**	4	3.503	**	8	2、597 ''	165*
150	**	4	**	4	2.877	17	9	1.749 "	165*
165	**	4	**	4	2.731	**	8	1,589 "	165*
180	**	4	* *	4	2.683	**	9	1.540 "	172*
Computing tin seconds	ne, 0.04	:	0.04			0.06	;	0.12	

*Transformed series converged in LT-150 terms.

**A obtained by interpolation from Figure 1.


Figure 1

Goudsmit-Saunderson Distribution: A (ω =5°) vs. A (ω =0°) for Several Values of Z, E₀, and t