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The effects of electron divergence on the point design

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FSC Meeting on Electron Divergence
5 August 2010

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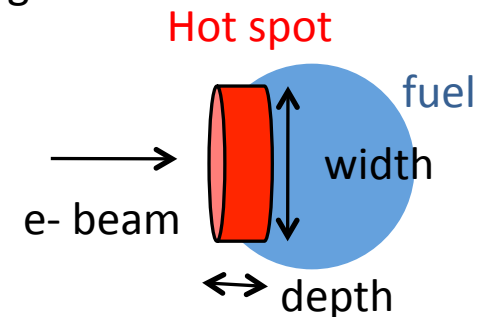
Release No. LLNL-CONF-447431.

Summary: Divergent electron source pushes us to width > depth hot spots, large beam energies; need to try focusing tricks

- **Electron source:** from full PIC LPI simulations [A. Kemp, L. Divol]:
 - Energy spectrum: modified two-temperature, ponderomotive scaling with laser intensity. Large fraction of energetic electrons don't fully stop in hot spot.
 - Angle spectrum: highly divergent, hard to hit small hot spot.

- **Width > depth ignition:** to lower laser intensity, and subtend divergent beam, use large radius beams and hot spots with width > depth.

- Requires more deposited energy than small hot spot, but we can't hit a small spot!
- Look for ignition w/ low hot-spot temperatures (4-6 keV).



- **Ignition-scale transport simulations:** with implicit PIC code LSP.

- Comparison with Atezni's ignition condition.
- Burn calculations in HYDRA on LSP final conditions.

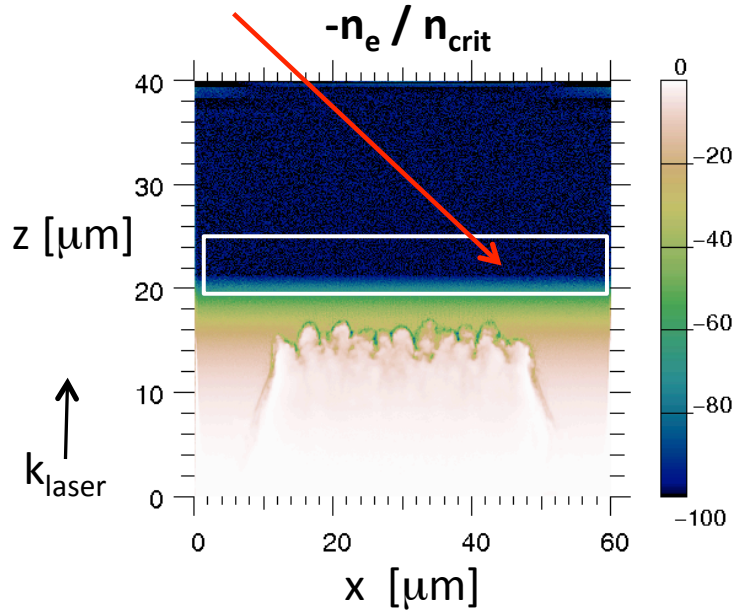
- **Magnetic focusing** shows promise – resistivity gradients, external B fields – not discussed further here.

We directly compare electrons in PSC “white box” with an LSP “white box” in a test run with an excited electron beam (forces are included)

PSC run “3Dpre”:

- 3D Cartesian, 1 μm wavelength.
- pre-plasma $n_e \sim \exp[z / 3.5 \mu\text{m}]$.
- Peak dens = 100 n_{crit} . Data at time 365 fs.
- best focus: $I_{\text{las}}(r) = I_0 \exp[-(r/18.3 \mu\text{m})^8]$
 $I_0 = 1.37 \text{ E}20 \text{ W/cm}^2$. $T_{\text{pond}} = 4.63 \text{ MeV}$.

“white box”: Forward-going e- with kinetic energy between 0.5 and 30 MeV

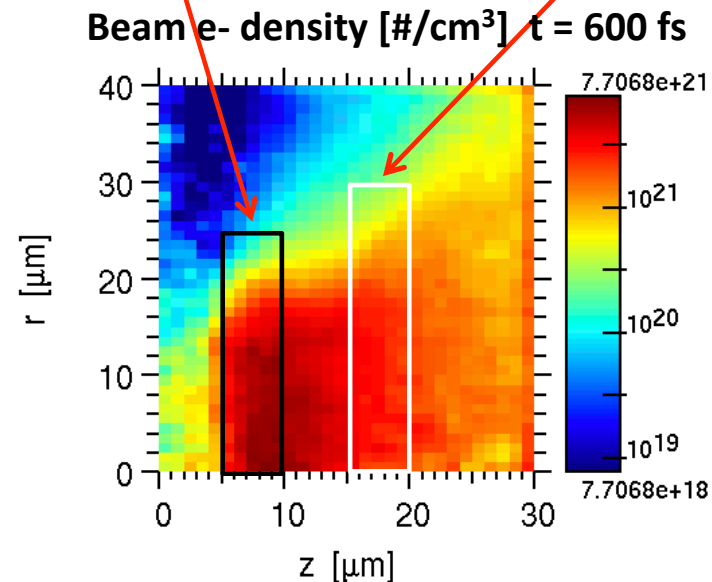


LSP run “source5”:

- 10 g/cc, 100 eV, Z=6 carbon.
- Beam radial profile same as laser intensity I_{las} .
- Forces from E/B fields included.
- No dE/dx or scattering since not in PSC run, and would be larger at LSP’s high density.

black box: e- beam excited

LSP white box



- beam $f(E, \theta) = f_E(E) * f_\theta(\theta)$; could add several together.
- number (not just energy) flux important: controls currents, B

Beam energetics, and two-temperature energy spectrum

Energy spectrum: quasi two-temperature,
scaled ponderomotively

$$dN/d\varepsilon = \frac{1}{\varepsilon} \exp[-\varepsilon/\tau_1] + b_2 \exp[-\varepsilon/\tau_2]$$

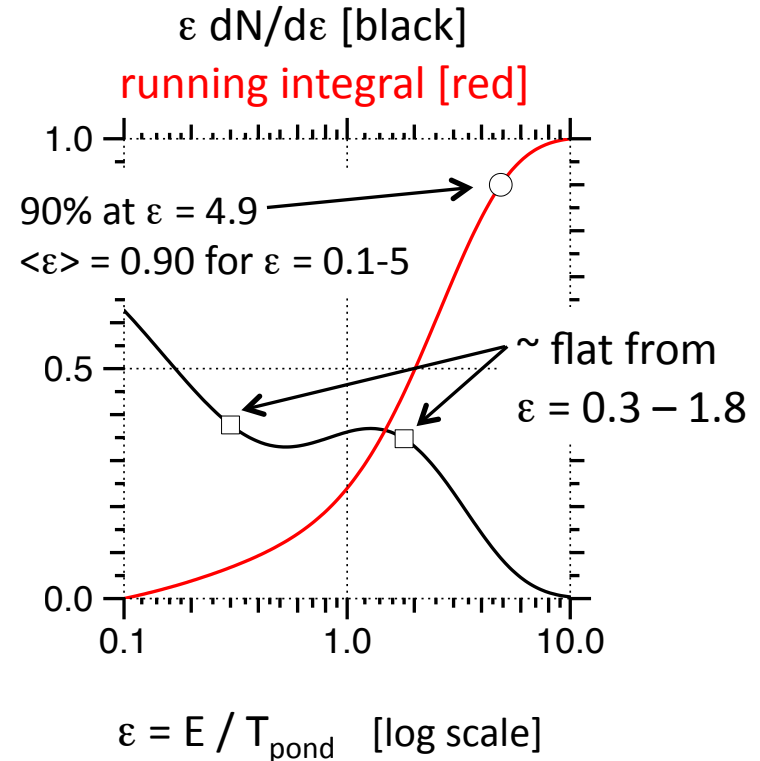
Note asymmetry in the two terms!

$$\frac{T_{\text{pond}}}{m_e c^2} := [1 + a_0^2]^{1/2} - 1$$

$$\sim a_0 := \text{sqrt} \left[\frac{I_{\text{las}} \lambda^2}{1.37 \cdot 10^{18} \text{ W cm}^{-2} \mu\text{m}^2} \right]$$

$$\varepsilon = E / T_{\text{pond}} \quad \tau_i = T_i / T_{\text{pond}}$$

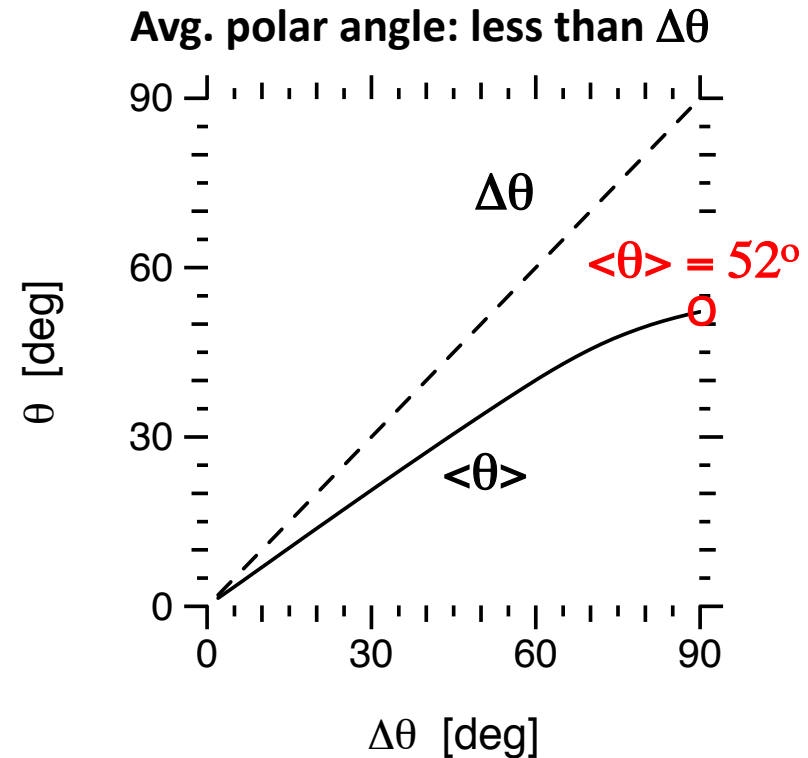
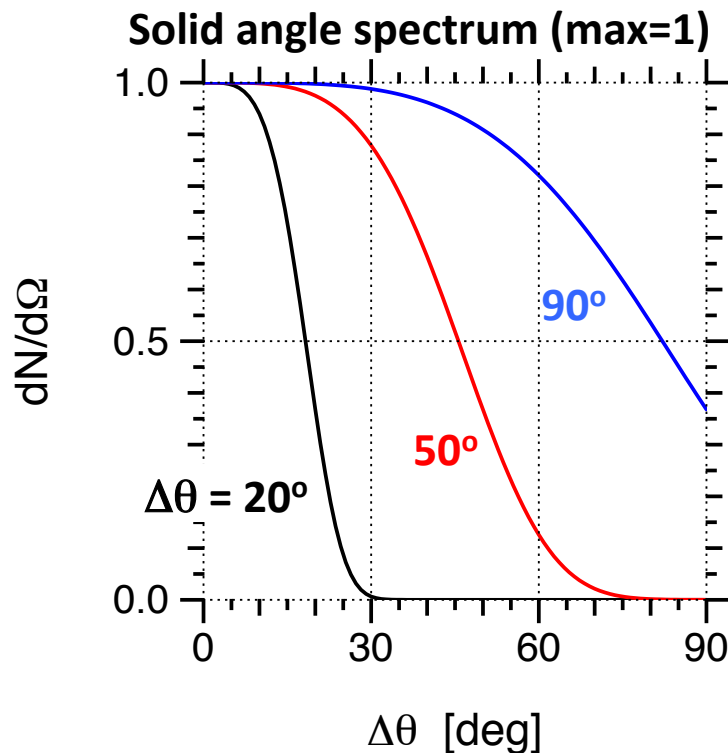
$$\tau_1 = 0.19 \quad \tau_2 = 1.3 \quad b_2 = \frac{n_2}{\tau_2} = 0.82$$



- PSC run “3Dpre”: Laser power = 1.3 PW.
- LSP beam power = 0.68 PW = 52% of laser power; sets absolute rate e-beam energy added.
- Total laser absorption is higher ~80-90%, but some is parasitic: expanding plasma,
- return current, ions, etc.

Source polar angle spectrum: super-Gaussian, large opening angle

$$\frac{dN}{d\Omega} = \exp\left[-(\theta/\Delta\theta)^4\right] \quad \Delta\theta = 90 \text{ deg for LSP source to match PSC}$$



- “The opening angle” is ill-defined: should specify $\langle\theta\rangle$, θ_{rms} , θ enclosing 90% of e-, etc.
- “The intrinsic source” is also ill-defined: my goal is LSP black-box source for transport sims that replicates e- in PSC white box. Different questions (e.g. K-alpha spot size) have different answers.
- Only fwd-going e- excited. LSP bug recently fixed, which gave bad angle spectrum when exciting > 1 ptcl / cell /step. My Anomalous 2010 talk suffered from this bug.

PSC (black) and LSP (red) e- in white boxes are similar – adequate for transport and design studies

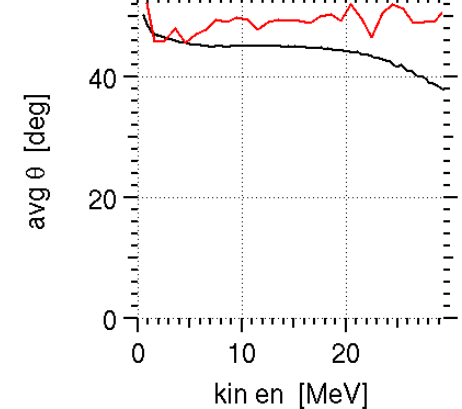
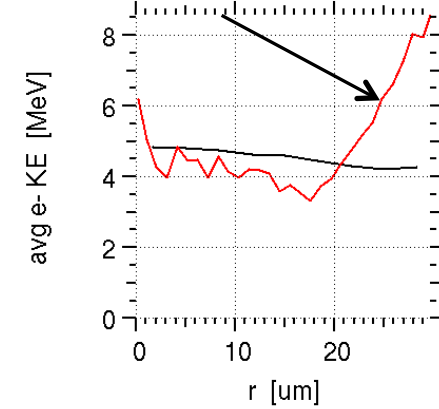
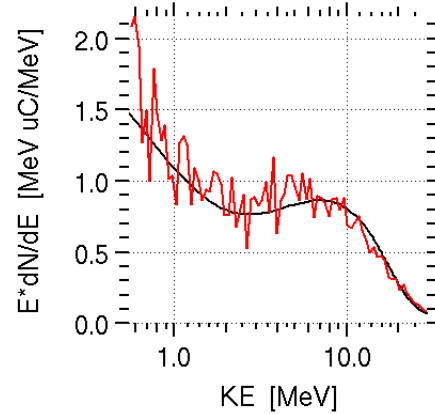
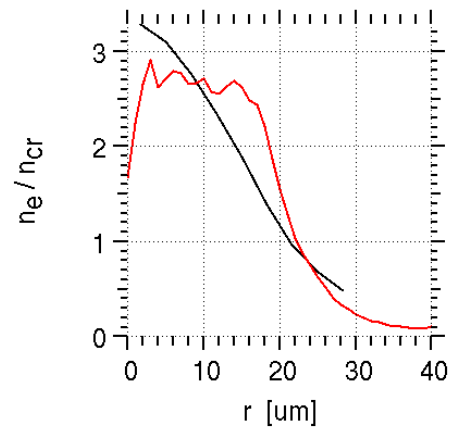
e- energy in white box [Joules]: 14.9 PSC, 15.0 LSP. 0.6% error!

Avg. polar angle theta [deg]: 46.8 PSC, 48.6 LSP, 3.7% error. Smaller than 52 deg in source region.

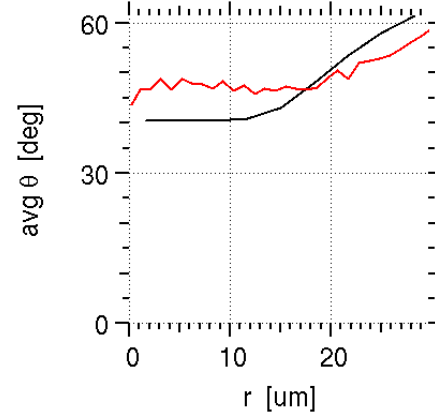
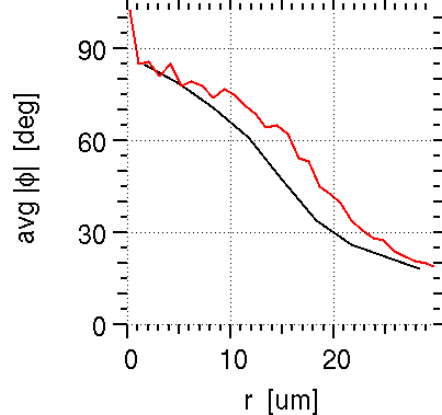
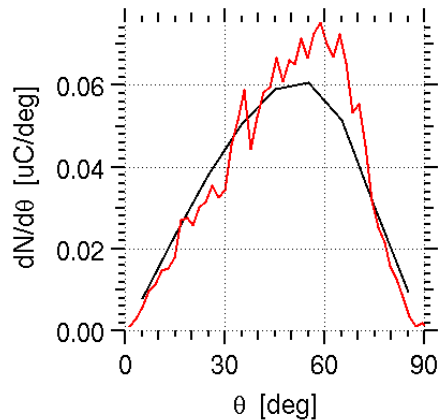
resistive B in LSP may deflect

low-KE e-; density low at large r

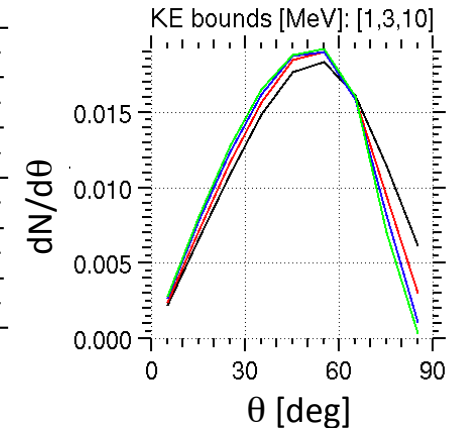
All LSP e-, not just some KE range



ϕ = angle b/t r and ρ , in xy plane¹
small $|\phi|$ -> radial outward drift



PSC dN/dtheta in en. bins:
higher en.bit more fwd-going

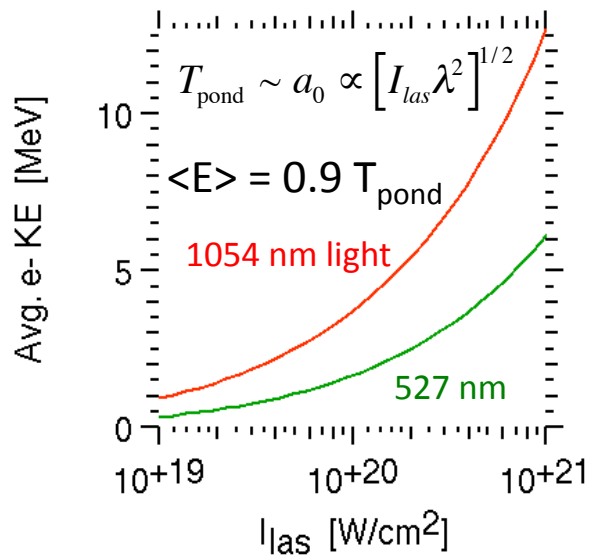


¹A. Debayle et al, IFSA 2009 Proceedings: radial outward drift w/ Gaussian laser spot; azimuthally uniform LSP source develops outward drift as it propagates to white box.

- **Electron source**
- **Width > depth ignition**
- **Ignition-scale transport simulations**

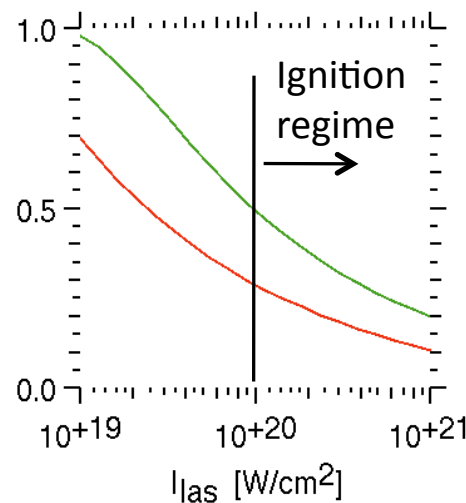
Usefulness of energy spectrum decreases with laser intensity, but energy deposited still increases; green light better than red

Avg. e- energy in our 2-T spectrum

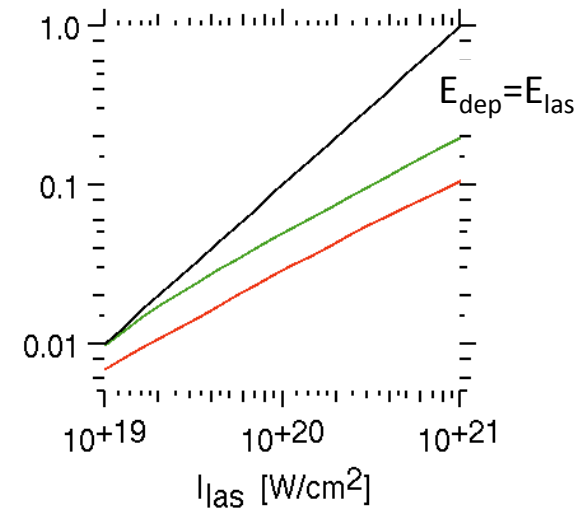


Energy deposited by 2-T spectrum in 1.2 g/cm² of 300 g/cm³ DT

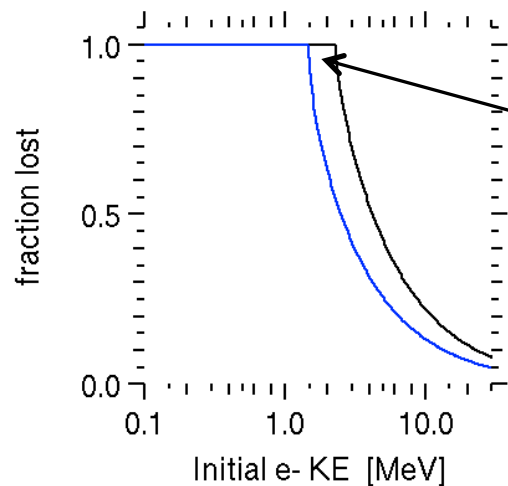
Fraction of energy deposited



Total deposited energy if $I_{\text{las}} \sim E_{\text{las}}$



Fraction of energy deposited DT, 300 g/cm³ (for one e-)



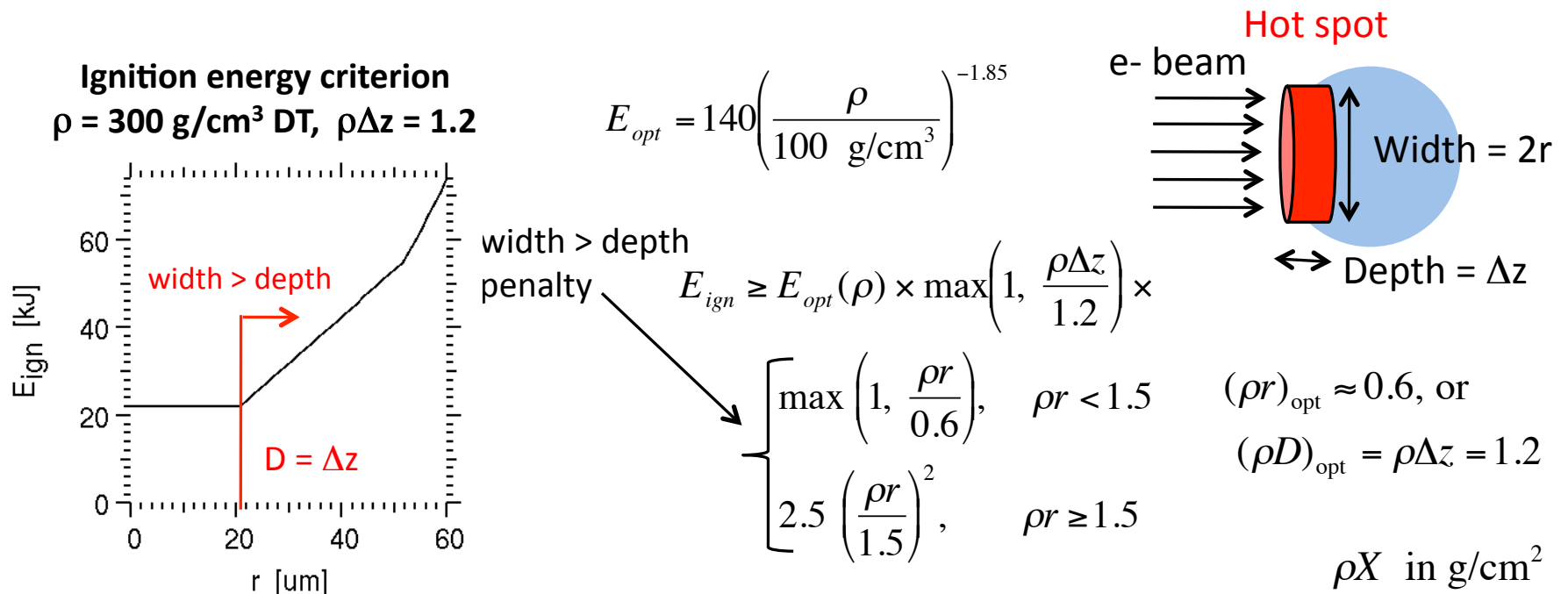
$E_{\text{crit}} [\text{MeV}] \sim 1.2 \rho \Delta x [\text{g/cm}^2]$

$\rho \Delta x = 2 \text{ g/cm}^2$
1.2

You still win as you increase laser energy (for fixed spot and pulse), but slower than linearly.

Ignition requirements for fast ignition: width > depth mode

- **Fuel:** $\rho \sim 300 \text{ g/cm}^3$, $\rho r > 2 \text{ g/cm}^2$ – should give energy gain ~ 100 w/ $\sim 1 \text{ MJ}$ indirect-drive compression laser.
- **Ignition energy:** TN burn not yet in our transport simulations. We rely on 2D rad-hydro studies by Atzeni et al.¹: Collimated, mono-energetic beam into spherical fuel.



One pays an energy penalty to ignite a width>depth hot spot, but it's the better choice if one's beam is too energetic or divergent – which our e- beams are.

¹S. Atzeni, A. Schiavi, C. Bellei, Phys. Plasmas 14, 052702 (2007)

Energy in hot spot vs. ignition energy is our figure of merit for our no-burn simulations

- **Ignition energy:** Atzeni used collimated, mono-energetic beams; we take his ignition energy as the energy one must deposit in the hot spot.
- **Hot-spot energy:** some deposited energy is lost, so we are generous in finding hot-spot energy.
- **Hot spot construction:** $\rho > 200 \text{ g/cm}^3$ and in depth $\rho^* \Delta z = 1.2$ from cone side.
 - $E_{\text{hot-spot}} = \text{thermal} + \text{flow energy in all species.}$

$E_{\text{hot-spot}} / E_{\text{ignition}} = \text{figure of merit to compare runs;}$
accurate ignition assessment requires simulations with burn.

- **Electron source**
- **Width > depth ignition**
- **Ignition-scale transport simulations**

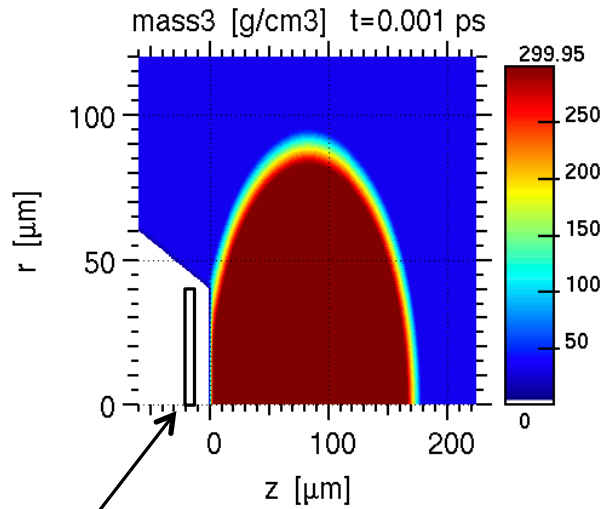
LSP¹ direct-implicit PIC code for transport modeling

- Excited electron beam (no laser) propagating to ideal targets.
- RZ geometry.
- Eulerian multi-species fluid background.
- Ideal gas EOS, fixed ionization.
- Unmagnetized Lee-More-Desjarlais background transport (e.g. electrical and thermal conductivities).
- Solodov/Davies stopping and scattering of fast e⁻ by Lemons-like algorithm.
 - Fast electrons collide with both free and bound background electrons! Only differ in “log lambda”.
- Energy conservation is excellent (usually better than 1% of beam energy).
- Problems do occur in B (and E) fields near R=0. Thanks to A. Solodov for coding for B-field smoother, which helps.

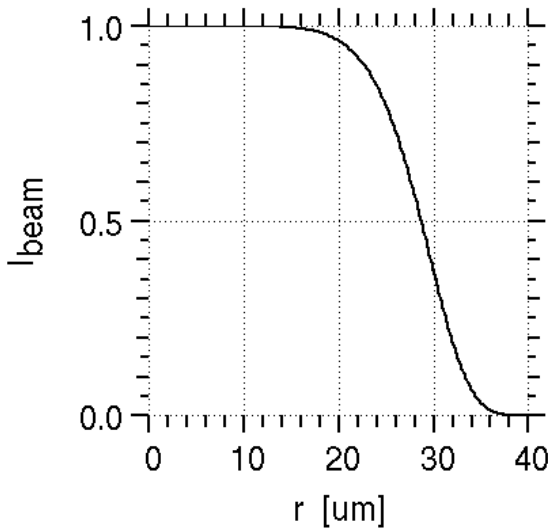
¹D. R. Welch, D. V. Rose, M. E. Cuneo, R. B. Campbell, T. A. Mehlhorn, Phys. Plasmas 13, 063105 (2006)

LSP run ign02: realistic $\Delta\theta = 90$ deg

Cone: 20 g/cc, Z=6 carbon; fuel = DT; total e- beam energy = 127 kJ; laser energy = 127/0.52 = 244 kJ



beam excited from z=-15 to -10 μm

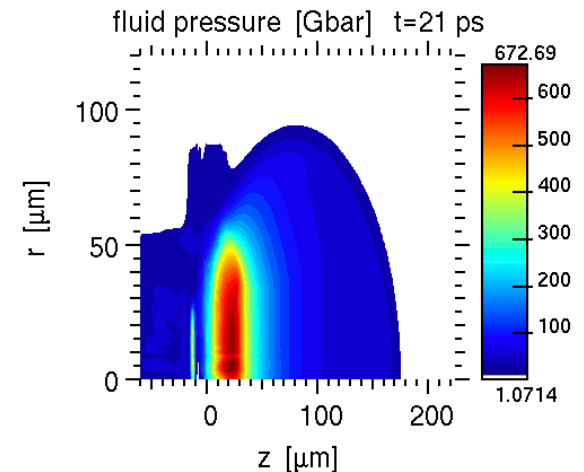
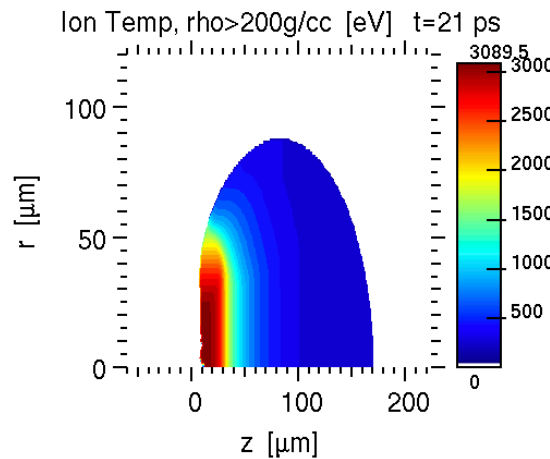
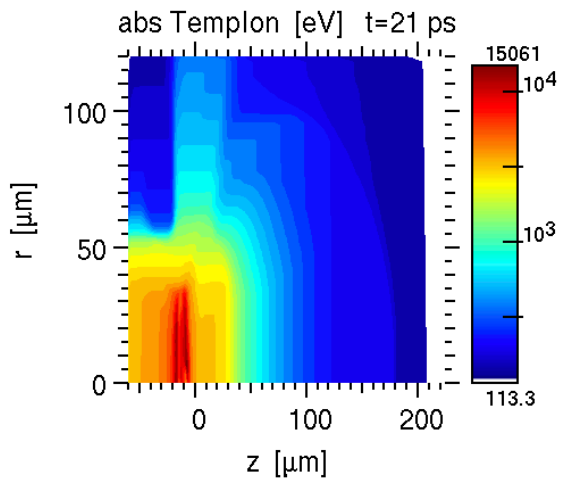


$$P = \int da I = I_0 A_{\text{eff}}$$

$$A_{\text{eff}} = \int dr 2\pi r (I/I_0) = \pi r_{\text{eff}}^2$$

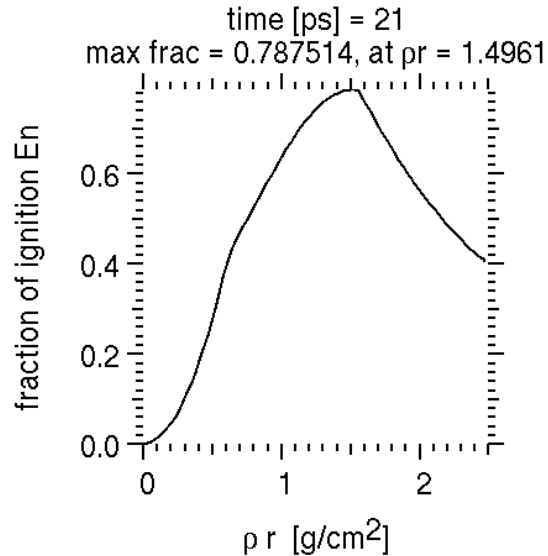
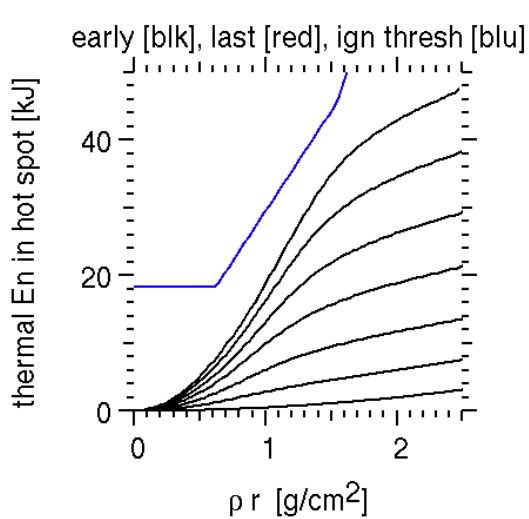
$$I_{\text{beam}}(r) = I_0 \exp[-(r/30\mu\text{m})^8]$$

$$\rightarrow r_{\text{eff}} = 28.6 \mu\text{m}$$



ign02: reached 79% of ignition condition w/ 127 kJ of beam energy (244 kJ laser); ign03 met condition w/ 185 kJ beam

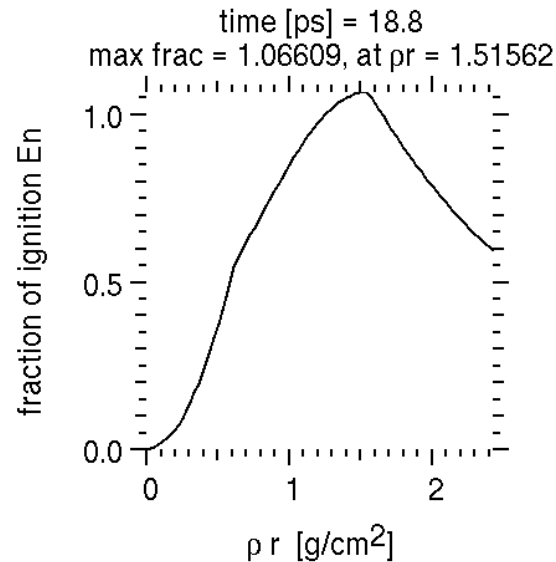
ign02



ign03: more beam en.: 185 kJ beam (356 kJ laser) added at 18.8ps, when we meet ignition condition

Have yet to:

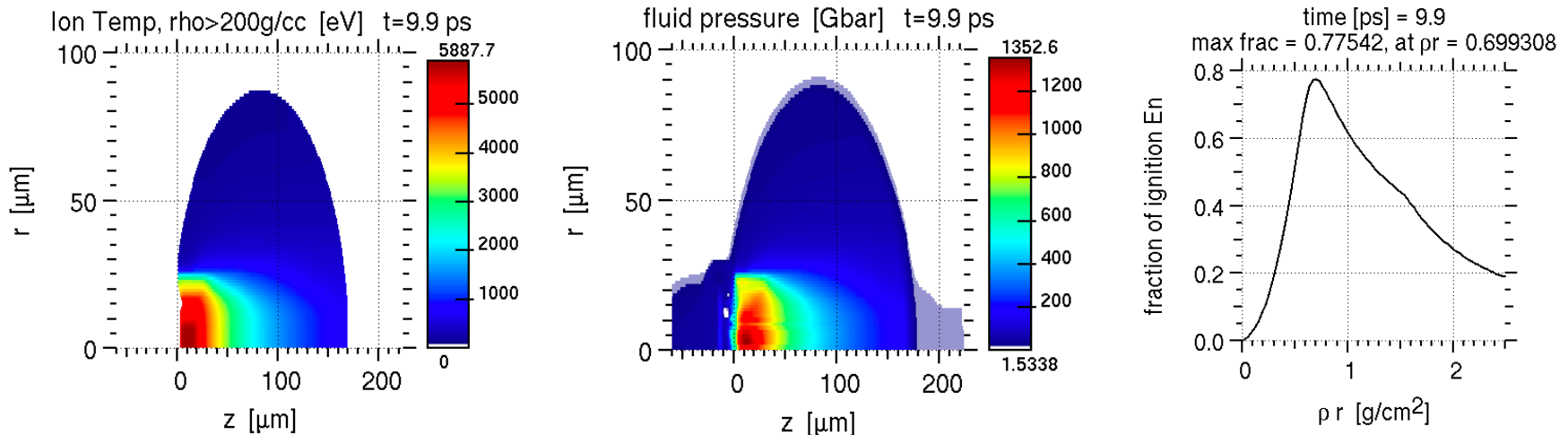
- vary spot size, standoff distance
- compare with Zuma
- try resistivity gradients



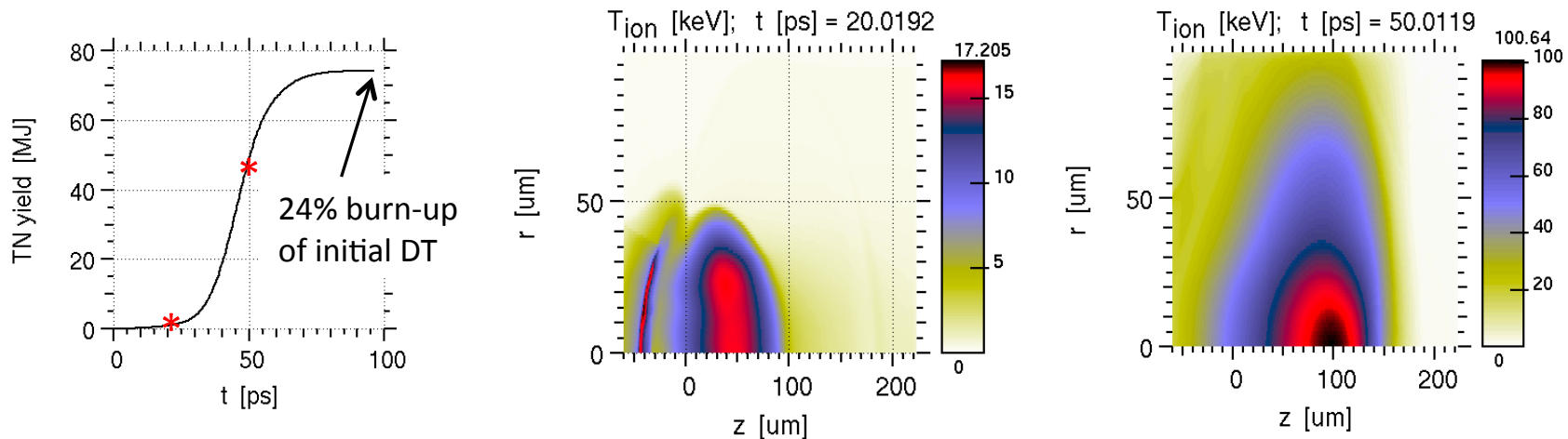
Results of HYDRA burn calculation on LSP output

- LSP run with artificially collimated source: $\Delta\theta = 20$ deg (90 deg would match PSC).
- Beam intensity profile $\exp[-(r/21 \text{ } \mu\text{m})^8]$ – smaller spot than previously.

LSP conditions at $t = 9.9$ ps (73 kJ of e-beam energy added):



HYDRA run on these conditions didn't ignite, but doubling the temperature (overkill) does:



Summary: toward electron-driven fast ignition

- **Electron source:** full-PIC simulations with pre-plasma show highly divergent angular spectra, and a modified two-temperature energy spectrum. Both facts make it hard to deposit energy in a small hot spot.
 - Can no pre-plasma help? Other laser-plasma interaction tricks?
- **Width > depth ignition:** we are pushed towards larger hot spots with $\rho*r = 1-1.5 \text{ g/cm}^2$ and temperatures $\sim 4-5 \text{ keV}$.
- **LSP Transport modeling:** Ignition condition can be achieved, with no cone-fuel standoff, with green (527 nm) laser:
 - $E_{\text{beam}} = 185 \text{ kJ}$ for realistic $\Delta\theta = 90 \text{ deg}$. Does this large, $\sim 3 \text{ keV}$ hot spot burn?
 - $E_{\text{beam}} \sim 73 \text{ kJ}$ in 10 ps and $\Delta\theta = 20 \text{ deg}$. obtains 80% of ignition condition.
 - Doesn't ignite in Hydra, but doubling the temperature (probably overkill) does: burn-up fraction of 24%.
- **Magnetic collimation:** may reduce beam divergence, improve coupling.

Collimation should be explored more, but for ignition relevant targets: scales, energies, pulse durations, geometries, etc.

Backup slides after here

LSP run source6: $\Delta\theta = 80$ deg instead of 90 deg

e- energy in white box [Joules]: 14.9 PSC, 15.0 LSP. 0.5% error!
Avg. polar angle theta [deg]: 46.8 PSC, 47.2 LSP, 0.88% error.

