

Curvature of the Lanthanide Contraction: An Explanation

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Abstract

A number of studies have shown that for isostructural series of the lanthanides (elements La through Lu) a plot of equivalent metal-ligand bond lengths versus atomic number differs significantly from linearity and can be better fit as a quadratic equation. However, for hydrogen type wave functions it is the inverse of the average distance of the electron from the nucleus (an estimate of size) that varies linearly with effective nuclear charge. This generates an apparent quadratic dependence of radius with atomic number. Plotting the inverse of lanthanide ion radii (the observed distance minus the ligand size) as a function of effective nuclear charge gives very good linear fits for a variety of lanthanide complexes and materials. Parameters obtained from this fit are in excellent agreement with the calculated Slater shielding constant, k .

Keywords: Lanthanides, X-ray diffraction, Structure elucidation, Electrostatic interactions, Bond theory.

1. Introduction

The lanthanide contraction, the change in ion size with atomic number, was one of the first structural features of the lanthanides and papers referring to it date back to at least 1926 [1]. Because early diffraction studies were time-consuming the entire lanthanide series (here defined as La through Lu, excepting the radioactive Pm) was typically not examined and the decrease in ionic radius was usually shown using representative members of the series [2] and the contraction was assumed to be linear with atomic number [3]. Modern advances in structure analysis have yielded a number of whole or near-whole series of structures for isostructural compounds [4-8], generating much more data for a plot of the lanthanide contraction and showing that the contraction is significantly curved, a feature that we and others in the field have generally accepted.

The early x-ray and neutron diffraction characterization of the series $[\text{Ln}(\text{H}_2\text{O})_9][\text{CF}_3\text{SO}_3]_3$ was especially significant [6, 9], because of its relevance to the solution chemistry of the lanthanide aqua ions. This series is isostructural and the cations have C_{3h} crystallographic site symmetry. Thus each vertex and each cap H_2O is equivalent within its class. In an ionic model the capping ligands have more ligand-ligand repulsion and so have longer Ln-O distances. This result is mirrored in modern calculations [10]. The authors of the $[\text{Ln}(\text{H}_2\text{O})_9][\text{CF}_3\text{SO}_3]_3$ papers struggled to explain both the curvature in the Ln-O distances across the series as well as the different curvature seen for the two different water molecules [6]. Subsequently, Quadrelli examined this and the isostructural series $[\text{Ln}(\text{H}_2\text{O})_9][\text{EtOSO}_3]_3$ [8], and found both to be more closely fit by a quadratic expression [11].

More recently, Ibers *et al.* determined the structure of a series of compounds with the general formula RbLnSe_2 [4]. The lanthanide ions lie in octahedral coordination environments, and each coordinating selenium atom is identical. The Ln-Se internuclear distances of the complexes were found to be more closely fit by a quadratic expression, as were the Ln-Te distances in the series CsLnZnTe_3 [5].

We earlier presented an explanation of the apparent quadratic nature of the contraction based on the nuclear shielding of the f orbitals [12] and Slater's model for atomic shielding [13-15]. That paper described a structural series for a multidentate ligand. Here we expand upon that model and show its effective explanation of Ln^{3+} ionic radii for a number of published series of isostructural complexes.

2. Modeling and plotting ionic radius from crystal-structure data

For hydrogen atom wave functions, it is not the average distance of the electron from the nucleus, $\langle r \rangle$, that varies linearly with effective nuclear charge but rather the average inverse, $\langle 1/r \rangle$. This suggests that $1/r(x)$, the inverse of the lanthanide ionic radius (where x is the number of f electrons in the lanthanide ion of interest, treated as a continuous variable), can be

represented by a linear equation:
$$\frac{1}{r(x)} = a + bx \quad (1).$$

As will be shown, this is mathematically equivalent to the Slater shielding formula.

Here the intercept $a = \frac{1}{r_0}$ (2)

is the inverse of the ionic radius of lanthanum and the slope

$$b = \frac{\partial \left(\frac{1}{r(x)} \right)}{\partial x} \quad (3)$$

is the incremental change in the inverse of the ionic radius from one lanthanide to the next. Thus, the internuclear distance between the lanthanide and the coordinating atom of the ligand, $d(\text{Ln-L})$, is given by

$$d(\text{Ln-L}) = \frac{1}{a+bx} + r_{(L)} \quad (4)$$

where $r_{(L)}$ is the radius of the coordinating ion on the ligand.

Eq. (4) can be expanded through a Maclaurin series [16] to give an equation of the form

$$f(x) = f(0) + f'(0)x + f''(0)x^2/2 + \dots \quad (5)$$

which, if terminated at the third term, gives the polynomial described by Quadrelli

$$d(\text{Ln-L}) = A_0 - A_1x + A_2x^2 \quad (6)$$

where

$$A_0 = 1/a + r_{(L)}$$

$$A_1 = b/a^2$$

$$A_2 = b^2/a^3$$

Equations 1-6 suggest that a plot of $1/r(x)$ as a function of the number of f electrons will be linear. The ionic packing model, with derived ionic radii, was pioneered by Goldschmidt [17] and Pauling [18], culminating in tabulations such as the widely used values of Shannon [19]. The relevant Ln-L internuclear distances are shown in Table I. By subtracting the known ionic radii of the ligating atoms from the Ln-L internuclear distances in simple structures, the ionic radii of the Ln ions can be determined for a series of experimentally determined distances.

A least-squares fit of the crystallographic data to obtain the parameters a , b , and $r_{(L)}$ of Eq. (4) for each ligand was performed by minimizing:

$$\sum_{i=1}^n (d_i(x)^{obs} - d_i(x)^{calc})^2 \quad (7)$$

where $d_i(x)^{obs}$ is the observed internuclear distance and $d_i(x)^{calc}$ is the corresponding predicted values $d(Ln-L)$ for the i^{th} crystallographic data point. In order to account for differences in lanthanide ionic radius that were due to the different coordination environments, Eq. (8) [20] was used to normalize the values of both a and b for coordination number 6 and 9

$$\frac{ir_{CN(q)}}{ir_{CN(p)}} = \left(\frac{CN(q)}{CN(p)} \right)^{\frac{1}{10}} \quad (8)$$

where $ir_{CN(q)}$ and $ir_{CN(p)}$ are the ionic radii of the lanthanide ions having coordination number (CN) q and p , respectively. Values for the radius of the vertex water oxygen atoms of $[Ln(H_2O)_9][CF_3SO_3]_3$ and $[Ln(H_2O)_9][EtOSO_3]_3$, the selenium atoms of $Rb[Ln]Se_2$, and equatorial tellurium atoms of $Cs[Ln]ZnTe_3$ were fixed according to literature values [19] in order to decrease the number of parameters to be refined. With the refined values of $r_{(L)}$ (see supporting information), the inverse of the ionic radius is graphed as a function of the number of f electrons for the four chosen series.

The series $RbLnSe_2$ ($Ln = La, Ce, Pr, Nd, Sm, Gd, Tb, Ho, Er, Lu$) [4] crystallizes with lanthanide ion site symmetry D_{3d} ; the lanthanide ions are nearly octahedral and each coordinating selenium atom is identical. A coordination number of 3 for the selenium atoms was used to fix the ionic radii of the selenium atoms to previously published values [19] for the purposes of the data refinement. A plot of $1/r(x)$ as a function of the number of f electrons is shown in Fig. 1.

The series CsLnZnTe_3 ($\text{Ln} = \text{La, Pr, Nd, Sm, Gd, Tb, Dy, Ho, Er, Tm, Yb}$) [5] crystallizes in the layered KZrCuS_3 structure type in the orthorhombic space group Cmcm . The structure is composed of LnTe_6 octahedra, with axial and equatorial tellurium atoms in the unit cell. The equatorial atom ionic radii were fixed to the 3-coordinate literature value [19], and the axial tellurium ionic radii, due to their different coordination environment, were refined. Fig. 2 shows our results for these complexes.

In the lanthanum nonaqua triflate and ethyl sulfate salts ($[\text{Ln}(\text{H}_2\text{O})_9][\text{CF}_3\text{SO}_3]_3$, $\text{Ln} = \text{La, Ce, Pr, Nd, Sm, Eu, Gd, Tb, Dy, Yb, Lu}$ and $[\text{Ln}(\text{H}_2\text{O})_9][\text{EtOSO}_3]_3$, $\text{Ln} = \text{La, Ce, Pr, Nd, Sm, Eu, Gd, Tb, Dy, Ho, Er, Tm, Yb, Lu}$) [6, 8] the cations are also tricapped-trigonal prismatic in geometry. The vertex ionic radii were fixed to a literature value for water [21] and the capping radii were refined, constrained to a single value (the values for Tm, Yb, and Lu are anomalous, which we attribute to steric crowding, and were omitted). The linear plots are shown in Fig. 3 and 4.

4. Calculating the screening constant for 4f electrons

Using Slater's revised model [14], the ionic radius of a complexed lanthanide ion can be written in the form

$$r(x) = \frac{r_0 Z_0^*}{Z_0^* + x(1-k)} \quad (10)$$

where $r(x)$ is the ionic radius of the lanthanide, Z_0^* is the effective nuclear charge of lanthanum, r_0 is the ionic radius of lanthanum, x is the number of f electrons and k is the atomic screening constant [12]. Therefore

$$\frac{1}{r(x)} = \frac{Z_0^* + x(1-k)}{r_0 Z_0^*} \quad (11)$$

and, after taking the first derivative of $1/r(x)$ with respect to x ,

$$\frac{\partial \left(\frac{1}{r(x)} \right)}{\partial x} = \frac{1-k}{r_0 Z_0^*} \quad (12)$$

we can combine Eq. (12) with Eq. (3) to get

$$k = 1 - br_0 Z_0^* \quad (13)$$

Substituting equation (2) gives:

$$k = 1 - \frac{b}{a} Z_0^* \quad (14)$$

that is, the apparent curvature of the ionic radii can be used to calculate the Slater screening constant. Calculating the screening constant using our refined values for a , b and an extrapolated value for the effective nuclear charge of lanthanum, Z_0^* in Eq. (14) [14] results in a value for k of 0.74, in good agreement with the previously reported value of 0.69 obtained from the ionization energies of the lanthanides [22].

5. Conclusions

The lanthanide contraction is a topic of considerable history and practical as well as theoretical significance. The increasing number of structures and the more complete sets of data for the entire series of lanthanides in a given structural type show that the change of lanthanide to ligand bond distance does not change linearly with atomic number. The apparent curvature of the ionic radii with atomic number is because it is $1/r$ that varies linearly with charge. This simple model gives good fits to the data and a useful value of the Slater screening constant.

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