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Research Article

Remarks on Constitutive Modeling of Nanofluids

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We discuss briefly the constitutive modeling of the stress tensor for nanofluids. In particular, we look at the viscosity of nanofluids containing multiwalled carbon nanotubes (MWCNTs) stabilized by cationic chitosan. MWCNTs can be used either to enhance or reduce the fluid base viscosity depending on their weight fractions. By assuming that MWCNT nanofluids behave as generalized second-grade fluid where the viscosity coefficient depends upon the rate of deformation, a theoretical model is developed. A simplified version of this model, similar to the traditional power-law model, is used in this study. It is observed that the theoretical results agree well with the experimental data.

1. Introduction

Nanofluids are made by adding nanoscale particles in low volumetric fractions to a fluid in order to enhance or improve their rheological, mechanical, optical, and thermal properties. The base fluid can be any liquid such as oil, water, ethylene glycol, or conventional fluid mixtures. Limited available studies on nanofluid viscosity have been reported [1–19]. In most of these studies, the behavior of the viscosity and the shear stress of nanofluids have been interpreted using the widely used empirical model developed by Casson [20]

$$\tau^{1/2} = \tau_0^{1/2} + \mu_{\infty}^{1/2} \dot{\gamma}^{1/2}. \tag{1}$$

In this equation, τ_0 is the yield stress, μ_∞ is the suspension viscosity at infinite shear rate, and $\dot{\gamma}$ is the shear rate. One of the inherent limitations of such empirical models is that they are, in general, one-dimensional in nature and it is not that easy or straightforward to generalize and obtain the appropriate 3-dimensional form, which are often necessary to solve general 3-dimensional problems. Nevertheless, this equation has been found to be successful for a range of parameters and a class of fluids. Phuoc and Massoudi [14] used this equation and obtained the values for μ_∞ being 0.1225 cp and 0.0225 cp for Fe₂O₃—deionized water nanofluids with polyvinylpyrrolidone (PVP) or polyethylene

oxide (PEO) as a dispersant, respectively. These values are about two orders of magnitude lower than the viscosity of the base fluid (a liquid prepared with PVP as a dispersant (DW-0.2% PVP) had a viscosity similar to that of water, while the viscosity of water with PEO as a dispersant (DW-0.2% PEO) was about 12.5 cp). Choi et al. [8] used this equation and calculated the intrinsic viscosities of CrO_2 —ethylene glycol, γ -Fe₂O₃, α -Fe₂O₃—EG and Ba-ferrite-EG nanofluids at infinite shear rate and reported a decrease of the viscosity with an increase in the particle volume fraction. This could be problematic, since the intrinsic viscosity should reach the viscosity of the base fluid in case of dilute suspensions or increase as the particle volume fraction increases if the suspension is dense enough.

In general, most complex, that is, nonlinear, materials exhibit unusual and peculiar characteristics such as viscoelasticity (as, for example, identified by creep or relaxation experiments, often exhibiting memory effects), yield stress, normal stress differences. The science of studying nonlinear fluids is "Rheology" and according to Reiner [21, p. 457]: "rheology started when Bingham in 1916 investigated concentrated clay-suspensions, and Bingham and Green in 1919 investigated oilpaints." The non-linear time-dependent response of complex fluids constitutes an important area of mathematical modeling of non-Newtonian fluids. For many

practical engineering cases, where complex fluids such as paint and slurries are used, the shear viscosity can be a function of one or all of the following:

- (i) Time,
- (ii) shear rate,
- (iii) concentration,
- (iv) temperature,
- (v) pressure,
- (vi) electric field,
- (vii) magnetic field,
- (viii)

Thus, in general, $\mu = \mu(t, \pi, \theta, \phi, p, \mathbf{E}, \mathbf{B}, ...)$, where t is the time, π is some measure of the shear rate, θ the temperature, ϕ the concentration, p the pressure, E the electric field, and B the magnetic field. Of course, in certain materials or under certain conditions, the dependence of one or more of these can be dropped. A more appropriate question is not what the shear viscosity should be, but rather, what the stress tensor of a given fluid should be. Bingham [22] was one of the first scientists who proposed a constitutive relation for the stress tensor of a viscoplastic material in a simple one dimensional shear flow, where the relationship between the shear stress and the rate of shear was described in terms of a yield function $F = 1 - \tau_0/|\tau|$ where τ_0 is the yield stress and τ is the shear stress. For many fluids such as polymers, slurries, and suspensions, some generalizations have been made to model shear-dependent viscosities. These fluids are known as the power-law or the generalized Newtonian fluid models; these widely used models are deficient in many ways; for example, they cannot predict the normal stress differences or yield stresses and they cannot capture the memory or history effects [23, 24].

In an effort to obtain a model that does exhibit both normal stress effects and shear-thinning/thickening, Man [25] modified the constitutive equation developed by Rivlin and Ericksen [26] (see also [27, 28]) for a second-grade fluid by allowing the viscosity coefficient to depend upon the rate of deformation; that is, $\mu_{\text{eff}} = \mu \Pi^{m/2}$, where Π is the second invariant of the symmetric part of the velocity gradient, and m is a material parameter. When m < 0, the fluid is shearthinning, and if m > 0, the fluid is shear-thickening. In this paper, the viscosity of nanofluids containing multiwalled carbon nanotubes (MWCNTs) stabilized by cationic chitosan is studied. MWCNTs can be used either to enhance or reduce the fluid base viscosity depending on their weight fractions. By assuming that MWCNT nanofluids behave as a generalized second-grade fluid, where the viscosity coefficient depends upon the rate of deformation, a theoretical model is developed, and comparisons are made with the experimental data.

2. Constitutive Modeling

Mathematically, the purpose of constitutive relations in mechanics is to supply connections between kinematic, mechanical, and thermal fields providing a suitable formulation of a problem which can be solved for properly posed problems. Just as different figures in geometry are defined as idealizations of natural objects, continuum mechanics seeks to establish particular relations between the stress tensor and the motion of the body for "ideal materials" [27]. In some instances, it may be necessary to represent the same real material by different ideal materials in different circumstances. A classic example is that of the theory of incompressible viscous fluids, which gives an excellent description of the behavior of water flowing through pipes but is useless for the study of the propagation of sound waves through water. While a constitutive equation is a postulate or a definition from the mathematical standpoint, physical experience remains the first guide, perhaps reinforced by experimental data. Constitutive relations are required to satisfy some general principles. Wang and Truesdell [29, page 135] list six general principles: (1) determinism, (2) local action, (3) equipresence, (4) universal dissipation, (5) material frame indifference, and (6) material symmetry. Constitutive relations should hold equally in all inertial coordinate systems at any given time (often referred to as coordinate invariance requirement). This would guard against proposing a relation in which a mere change of coordinate description would imply a different response in the material. Many of the so-called "power-law" models used in describing non-Newtonian fluids are not coordinate invariant. In general, this difficulty can easily be overcome by stating the equations either in tensorial form or by using direct notations not employing coordinates at all. The principle of material frame-indifference (sometimes referred to as objectivity), which requires that the constitutive equations be invariant under changes of frame, is perhaps the most important of all. It is a consequence of a fundamental principle of classical physics that material properties are indifferent, that is, independent of the frame of reference of the observer. This principle requires that constitutive relations depend only on frame-indifferent forms (or combinations thereof) of the variables pertaining to the given problem (see Massoudi [30] for further details). Among other approaches to model complex materials, one can list (i) using physical and experimental models, (ii) doing numerical simulations, (iii) using statistical mechanics approaches, and (iv) ad hoc approaches.

In general, based on available experimental observations, it can be said that many nanofluids exhibit characteristics similar to those of non-linear materials such as colloidal suspensions, polymers, rubber, and granular materials. The main points of departure from linear behavior are the following:

- (1) the ability to shear-thin or shear-thicken,
- (2) the ability to creep,
- (3) the ability to relax stresses,
- (4) the presence of normal stress differences in simple shear flows,
- (5) the presence of yield stress.

To the best of our knowledge, it has not been reported whether nanofluids exhibit normal stress effects, the non-linear phenomena related to the stresses that are developed orthogonal to the planes of shear. Therefore, we propose to use a general model that can exhibit both the normal stress effects and the shear-thinning/thickening effects. To do so, we assume that nanofluids, such as the one studied in the present work, behave as generalized second-grade fluids. For a second grade fluid, the Cauchy stress tensor is given by [26–28]

$$\mathbf{T} = -p\mathbf{1} + \mu \mathbf{A}_1 + \alpha_1 \mathbf{A}_2 + \alpha_2 \mathbf{A}_1^2, \tag{2}$$

where p is the indeterminate part of the stress due to the constraint of incompressibility, 1 is the identity tensor, μ is the coefficient of viscosity, α_1 and α_2 are material moduli which are commonly referred to as the normal stress coefficients. The kinematical tensors A_1 and A_2 are defined through

$$\mathbf{A}_{1} = \mathbf{L} + \mathbf{L}^{T},$$

$$\mathbf{A}_{2} = \frac{d\mathbf{A}_{1}}{dt} + \mathbf{A}_{1}\mathbf{L} + \mathbf{L}^{T}\mathbf{A}_{1},$$

$$\mathbf{L} = \text{grad}\mathbf{u}.$$
(3)

where d/dt is the total time derivative, given by $d(\cdot)/dt = \partial(\cdot)/\partial t + [\operatorname{grad}(\cdot)]\mathbf{u}$, where \mathbf{u} is the velocity vector. The thermodynamics and stability of fluids of second grade have been studied in detail [28], where it is shown that if the fluid is to be thermodynamically consistent in the sense that all motions of the fluid meet the Clausius-Duhem inequality and that the specific Helmholtz free energy of the fluid be a minimum in equilibrium, then

$$\mu \ge 0,$$
 $\alpha_1 \ge 0,$
 $\alpha_1 + \alpha_2 = 0.$
(4)

By allowing the viscosity coefficient to depend on the rate of deformation, Man [25] modified the constitutive equation, (2) and proposed the following:

$$\mathbf{T} = -p\mathbf{1} + \mu \Pi^{m/2} \mathbf{A}_1 + \alpha_1 \mathbf{A}_2 + \alpha_2 \mathbf{A}_1^2, \tag{5}$$

where

$$\Pi = \frac{1}{2} \operatorname{tr} \mathbf{A}_1^2 \tag{6}$$

is the second invariant of the symmetric part of the velocity gradient, and m is a material parameter. When m < 0, the fluid is shear-thinning, and if m > 0, the fluid is shear-thickening. A subclass of models given by (7) is the generalized power-law model, which can be obtained by setting $\alpha_1 = \alpha_2 = 0$ in (7) (see [31–34] for further discussions of this model). Notice that if the normal stress parameters α_1 and α_2 are zero, then

$$\mathbf{T} = -p\mathbf{1} + \mu \Pi^{m/2} \mathbf{A_1}. \tag{7}$$

In this paper, we will use this simplified form, which can also be considered as a generalized power-law fluid model. Using the cylindrical coordinate system for our present measurements and assuming $\mathbf{u} = w(r)\mathbf{e}_{\mathbf{z}}$ where e_z denotes a unit vector along the z direction yields the following calculations

$$\mathbf{A}_{1} = \begin{bmatrix} 0 & 0 & \frac{\partial w}{\partial r} \\ 0 & 0 & 0 \\ \frac{\partial w}{\partial r} & 0 & 0 \end{bmatrix},$$

$$\mathbf{A}_{1}^{2} = \begin{bmatrix} \left(\frac{\partial w}{\partial r}\right)^{2} & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & \left(\frac{\partial w}{\partial r}\right)^{2} \end{bmatrix},$$

$$\Pi = \frac{1}{2} \left[2 \left(\frac{\partial w}{\partial r}\right)^{2} \right].$$
(8)

The *z*-component of the stress tensor becomes

$$T_{rz} = \mu \left[\left(\frac{\partial w}{\partial r} \right)^2 \right]^{m/2} \left(\frac{\partial w}{\partial r} \right) = (\mu \dot{\gamma}^m) \dot{\gamma}, \tag{9}$$

where the shear-dependent viscosity is defined as

$$\mu_{\text{eff}} = \mu \left[\left(\frac{\partial w}{\partial r} \right)^2 \right]^{m/2} = \mu \dot{\gamma}^m. \tag{10}$$

In the next section, we will briefly discuss the results of our experimental investigation and show how this model can be used to describe the observed behavior of the fluid.

3. Experimental Evaluation

In Figures 1 and 2, we present the results on the calculated and measured viscosity and shear stress for waterbased nanofluids containing Multiwalled carbon nanotubes (MWCNT) stabilized by low molecular weight chitosan (>75% deacetylation). The measured data were reported by Phuoc et al. [15]. The calculated results are carried out using (9) for the shear stress and (10) for the viscosity with m =-0.547, -0.65, and -0.647 and $\mu = 0.134$, 0.331, and 0.523 when CNTs weight percent increased from 1 to 3 and the weight percent of the chitosan was 0.1. While using 0.2 wt% chitosan, it was found that m = -0.584 and -0.678 and $\mu =$ 0.354 and 0.641 for 2 wt% and 3 wt% CNTs, respectively. It is seen that using the generalized power-law model, with (7) as a subclass of the generalized second-grade fluid models, the measured experimental values compare well with the theoretical model. For a given weight percent of the stabilizer, increasing the CNTs weight percent has a strong effect on the value of μ . For a given value of CNTs weight percent, increasing the weight percent of the stabilizer increases both m and μ .

4. Concluding Remarks

The two important constitutive relations needed for the study of flow and heat transfer in complex fluid-like

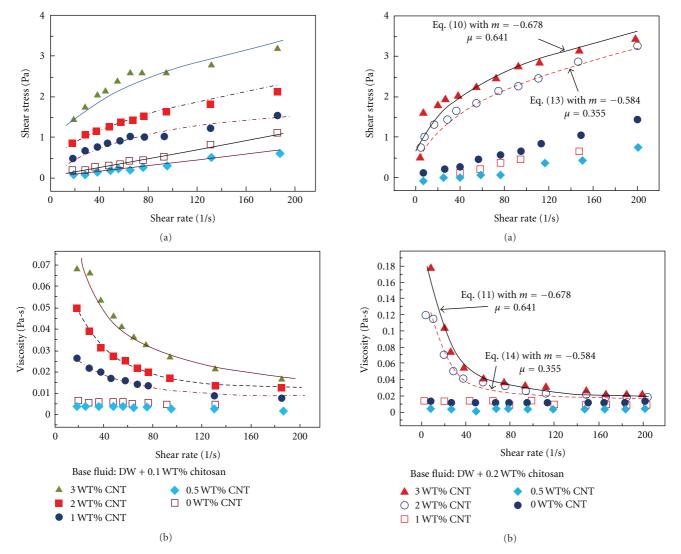


FIGURE 1: Viscosity and shear stress as a function of shear rate showing the effect of the MWCNT weight percent. The base fluid is DW + 0.1 wt% chitosan. The measured values are shown by symbols, while the calculated values are shown by the solid and dotted lines. These calculated viscosity values were obtained using (10) and the shear stresses were calculated using (9) with m = -0.547, -0.65, and -0.647 and $\mu = 0.134; 0.331$; and 0.523 for MWCNT weight percent increased from 1%, 2%, and to 3%, respectively.

FIGURE 2: Viscosity and shear stress as a function of shear rate showing the effect of the MWCNT weight percent. The base fluid is DW + 0.2 wt% chitosan. The measured values are shown by symbols, while the calculated values are shown by the solid and dotted lines. These calculated viscosity values were obtained using (10) and the shear stresses were calculated using (9) with m=-0.584 and -0.678 and $\mu=0.354$ and 0.641 for 2 wt% and 3 wt% CNTs, respectively.

materials, where the effects of radiation are ignored, are the stress tensor and the heat flux vector. From an engineering perspective, this oftentimes translates into measuring viscosity and thermal conductivity. As a result, most researchers have attempted to generalize Newton's law of viscosity and Fourier's law of heat conduction to various and more complicated cases by assuming that the shear viscosity and/or thermal conductivity could depend on a host of parameters such as shear rate, temperature, and porosity. This unfortunately does not fully take into account the non-linear and time-dependent nature (such as normal stress effects or visco-elastic response) of these complex

materials, as these generalizations cannot give rise to implicit constitutive relations.

Nanofluids represent one of the newest complex materials of the modern era. In many ways, constitutive modeling of these fluids, from a macroscopic point of view, is still at its infancy, perhaps similar to the early days of polymer rheology, rubber viscoelasticity or composite materials. With intense interest and research in the past two decades, great strides have been made, and nanofluids, due to their peculiar heat transfer and rheological properties, have been shown to contribute in many diverse ways to many industrial processes and to our lives [35]. An often neglected, yet extremely important, conceptual question, and perhaps still

an open question, is whether the same governing macroscopic balance equations can be used for nanofluids. For a recent discussion of this issue, see [36]. Since nanofluids form suspensions, theoretically from the point of view of mechanics, their thermorheological responses can be modeled either using the non-Newtonian approach or the multicomponent approach [37].

To better understand the various mechanisms in the heat transfer processes involving nanofluids, in addition to studying the thermal conductivity or radiation effects, an understanding of the mechanism for viscous dissipation is also important. A proper constitutive model for the stress tensor T, represents the first step in this direction, since the term $\mathbf{T} \cdot \mathbf{D}$ (where $2\mathbf{D} = \mathbf{A}_1$) appears in the energy equation. By assuming that nanofluids in general can behave as generalized second-grade fluids whose viscosity coefficient depends on the rate of deformation, a theoretical model has been developed. The experimental results indicate that the two important parameters in this study are related to the effects of the solid concentration on the viscosity of the base fluid and the degree of the nonlinearity of the fluid (measured through m and μ). By comparing with the measured data, the present model was found to be suitable for describing the fluid behavior. To test to see whether a particular nanofluid is capable of displaying normal stress differences, an orthogonal rheometer is needed.

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