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Analysis of localized beam losses in the Booster extraction straight section and the Booster to Storage Ring transfer line

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1 Introduction

The NSLS-II accelerators are installed within radiation shielding walls that are designed to attenuate the radiation generated from an assumed beam loss power (current multiplied by the assumed energy) within the beam channel to a level of <0.5mrem/h at the outer surface of the bulk shield walls of the accelerator. This Bulk Shielding Power assumes that the beam current is intercepted by a thick target (i.e., 30 cm of copper or steel) along the Booster or transport line beam axis and the dose rate level is calculated transverse to the beam axis. Any operational losses greater than this BSP level are expected to be addressed by installing supplemental shielding near that loss point in order to attenuate this radiation level outside the shield wall to the design level of <0.5mrem/hr.

The purpose of this exercise is to determine the potential beam loss locations, for further identification of the radiation risks they impose outside the existing bulk shielding, and specification of the supplemental shielding design necessary to mitigate those risks.

We will present the beam systems description along with their planned operating ranges, which forms the basis of our analysis. Next, we will describe the method used for each system to determine potential beam loss locations. Finally, as a conclusion, we will present the results of our studies in a form of convenient to use table which includes all the information necessary for analysis of radiation risks induced by the determined beam losses.

In the following analysis we assume, unless the opposite is indicated in the clearest possible terms, that the accelerator is configured in any feasible way that its design allows.

2 Description of the beamlines

We consider three beamlines: the Booster extraction straight section (ES), the Booster to Storage Ring (BSR) transfer line phase 1 and the Booster to Storage Ring transfer line phase 2. Figure 1 (obtained from [1]) shows the beamlines under consideration.

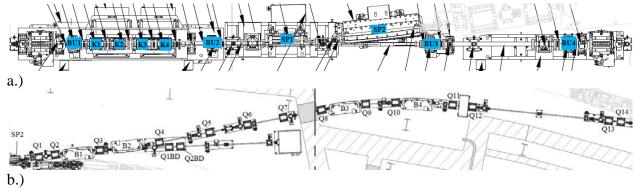


Figure 1: Schematics of ES (a) and BSR phase 1 and phase 2 (b).

The extraction section consists of four slow bumps forming local closed 4-bump BU1, BU2, BU3 and BU4 (their respective lattice file names are BRBU1SE, BRBU2SE, BRBU3SE and BRBU4SE; their respective engineering names are BR-XSBUM1, BR-XSBUM2, BR-XSBUM3 and BR-XSBUM4), four fast kickers K1, K2, K3 and K4 (their respective lattice file names are

BRK1SE, BRK2SE, BRK3SE and BRK4SE; their respective engineering names are BR-XSKIC1-1, BR- XSKIC1-2, BR- XSKIC2-1 and XSKIC2-2), pulsed septum magnet SP1 (its lattice file name is BRSP1SE and its engineering name is BR-XSSMP1), and DC septum magnet SP2 (its engineering name is BR-XSSMD1). These magnets are utilized for the extraction of the beam from the Booster into the BSR line.

The phase 1 of BSR line (a.k.a. dump line) consists of five quadrupoles Q1, Q2, Q3, Q1BD and Q2BD (their respective engineering names are BS-Q1, BS-Q2, BS-Q3, BS-Q1BD and BS-Q2BD), and one bending magnet B1 (its engineering name is BS-B1). This magnets transport beam to the beam dump when bending magnet B2 is turned off.

The BSR line phase 2 includes magnets Q1, Q2, Q3 and B1, and additionally has quadrupoles Q4, Q5, Q6, Q7, Q8, Q9, Q11, Q12, Q13, Q14 (BS-Q4, BS-Q5, BS-Q6, BS-Q7, BS-Q8, BS-Q9, BS-Q11, BS-Q12, BS-Q13, BS-Q14 in engineering notation) and bends B2, B3 and B4 (BS-B2, BS-B3, BS-B4 in engineering notation). This magnets transport beam to the Storage Ring. Injection into the Storage Ring is performed with SP3 and IS, which are the DC and pulsed Storage Ring injection septa respectively. Their engineering names are BS-SP3 and SR-IS-SP1.

3 Description of beamline elements

Each of the four bumps BU1-BU4 is an "H"-shape magnet consisting of a yoke and two coils [2]. The bumps nominal bending angle at 3 GeV is 7.5 mrad. All four bumps are powered in series. The peak current fed to the magnet is 1500 A (the power supply can provide about 10% higher current). The pulse length is 3 ms for one sinusoid. The bump's horizontal gap (horizontal distance between the yokes material) is about 152 mm. Figure 2 shows the general view of the bump magnet.

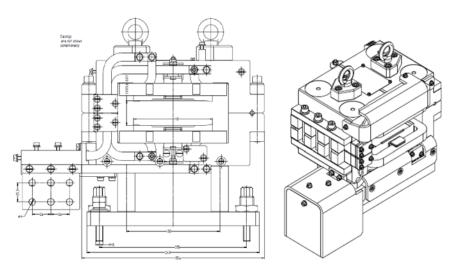


Figure 2: The central cross-section (left) and a general view of the bump magnet (right).

The fast kicker [2] is shown in Figure 3. The kickers provide nominal bending angle at 3 GeV of 1.45 mrad each. The kicker's pulse has 215 nsec rise time and 310 nsec flat-top. The flat top current is 2.3 kA. The power supply (PS) at maximum can provide about 10% higher current. Kickers K1 and K2 are powered in series. Kickers K3 and K4 are powered in series. Each kicker's horizontal gap is 60 mm.

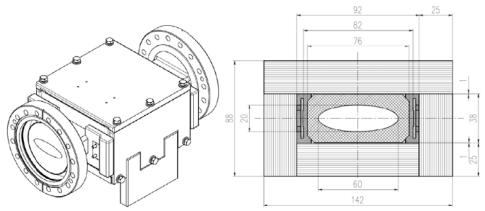


Figure 3: A general view (left) and the central cross-section of the kicker magnet.

The pulsed septum has a nominal bending angle of 48mrad [2]. The nominal current at the peak field from pulsed septum PS is 10.22kA (it can be about 10% higher). The extraction channel is 18 mm x 16 mm (horizontal x vertical) and accommodates the thin-walled (0.3 mm) stainless steel vacuum chamber of 15mm diameter. Figure 4 shows the cross-section and the top view of the pulsed septum.

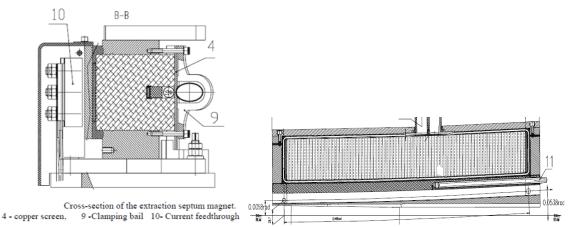


Figure 4: The cross-section (left) and the top view (right) of the pulsed extraction septum.

The DC septum is a "C"-shape bending magnet. The nominal bending angle of the DC septum is 104 mrad [3]. A nominal power supply current is 440 A and the maximum current is 500 A [4]. The DC septum accommodates vacuum chamber of 30mm diameter. The distance from the reference trajectory to the septum yoke is about 140mm. The view of the DC septum is given in Figure 5.

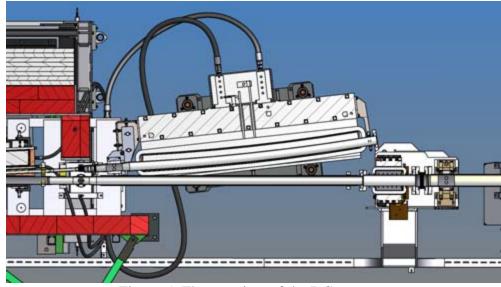


Figure 5: The top view of the DC septum.

The DC injection septum (SP3) is identical to SP2. Its nominal bending angle is 114 mrad.

The pulsed injection septum (IS) is a C-shaped magnet. It accommodates the vacuum chamber with 21 mm horizontal aperture. The septum yoke is touching the vacuum chamber. The nominal bending angle of pulsed septum is 100 mrad. The nominal current at the peak field from pulsed septum PS is 7.9kA (it can be about 10% higher).

The four BSR bends (B1-B4) are identical "C"-shape magnets. Their nominal bending angles are given in Table 5. Bends power supplies can provide 375 A at maximum [4]. The measurement of B1 gives the conversion coefficient of 0.00502 Tm/A [5]. Bends accommodate vacuum chamber of 120 mm width [6]. The distance from the reference trajectory to the yoke of the bend is about 220mm. The top view of the DC septum is shown in Figure 6.

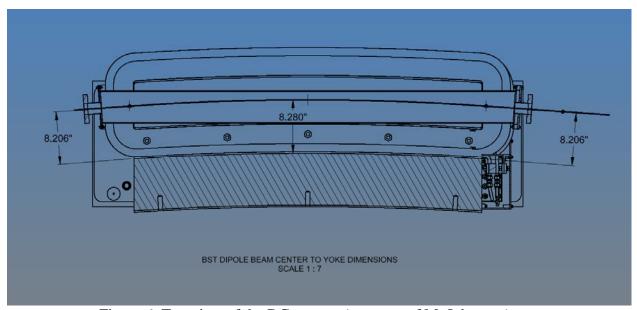


Figure 6: Top view of the DC septum (courtesy of M. Johansen)

The sixteen BSR quadrupoles (Q1-Q14, Q1BD and Q2BD) are identical. Their nominal gradients are given in Table 5. The maximum currents for quads power supplies are given in Table 1 [4]. The measured conversion coefficients for the quads are shown in Table 2 [5] (we use median value of 0.171 T/m/A for quads Q12-Q14 that were not measured). The quadrupoles accommodate vacuum chamber of 38mm diameter [6].

Magnet Name	BSR-Q1	BSR-Q2	BSR-Q3	BSR-Q4	BSR-Q5	BSR-Q6	BSR-Q7	BSR-Q8
PS Max Current	175	200	175	175	200	200	175	200
Magnet Name	BSR-Q9	BSR-Q10	BSR-Q11	BSR-Q12	BSR-Q13	BSR-Q14	BSR-Q1BD	BSR-Q2BD
PS Max Current	175	175	200	175	200	175	175	175

Table 1: The maximum currents of BSR quads power supplies (courtesy of G. Ganetis).

	Magnet Name	BSR-Q1	BSR-Q2	BSR-Q3	BSR-Q4	BSR-Q5	BSR-Q6	BSR-Q7	BSR-Q8
	Q coeff [T/m/A]	0.171	0.170	0.171	0.171	0.171	0.158	0.163	0.171
	Magnet Name	BSR-Q9	BSR-Q10	BSR-Q11	BSR-Q1BD	BSR-Q2BD			
ĺ	Q coeff [T/m/A]	0.171	0.170	0.173	0.171	0.170			

Table 2: Conversion coefficients for BSR quads (courtesy of R. Fliller).

4 Analysis of beam angles in the Booster extraction straight section

Although the booster ES magnets are supposed to be used only at extraction energy of 3 GeV (2 GeV extraction might be intentionally used as well), technically they can fire at any point of Booster ramp. Since these magnets are powered in series there are three scenarios to consider: either kickers K1 and K2 fire, or K3 and K4 fire, or four bumps BU1-4 fire. If the phase space available to the beam centroids at the entrance to the first of the chained magnets is defined by geometric acceptance of upstream drift (see below), and if BU2-4 are treated independently from BU1, then the three described scenarios cover all possible local beam losses that can happen under any feasible ES settings.

4.1 Beam angles analysis at low energy

Since at low energy (150 MeV) electron beam might be a freshly injected one rather than the stored one, the proper method of defining the beam centroids phase space at the entrance of the magnets is to define it by the acceptance of the respective upstream drift. This technique, although it results in overestimates of the lost beam angles provides a bullet-proof guarantee that all possible beams entering the magnet are considered. The described method (below we will refer to it as a "transfer line approach") is illustrated in Figure 7.

Starting with the beam centroids phase space described above we track each beam centroid through the first chained magnet. We stop tracking the beam if it hits the magnet yoke and record the angle at which such beam was "lost". Here, by the lost beam we mean that it is not a bright electron beam traveling towards the concrete wall anymore.

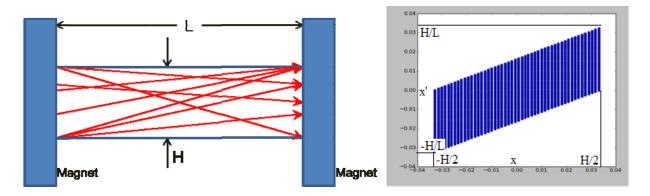


Figure 7: The picture on the left schematically shows the beam centroids (red) originating at the previous magnet and filling the drift upstream of the magnet under consideration. The plot on the right shows the phase space available to the beam centroids at the entrance of the considered magnet.

Next, we record the beam angle at the exit of the magnet and track the beam through the drift to the next chained magnet. If the beam makes it to the entrance of the next magnet without hitting the vacuum chamber, then the described tracking procedure is repeated for this magnet. This routine is continued until the exit of the last chained magnet is reached. Figure 8 illustrates this procedure.

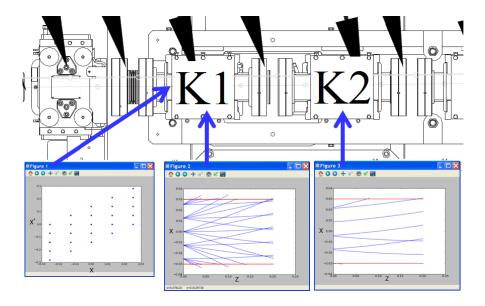


Figure 8: The left plot shows the phase space at the entrance of K1, which is defined by the geometry of upstream drift, filled with 25 beam centroids. The central plot shows beams tracked through K1, and the right plot shows beams (which survived the K1 and K1-K2 drift) tracked through K2. Red lines represent magnet yoke. The tracking is shown for maximum K1-2 current only.

					θmax R	θmax L
element	θnom	bending	θmax R	θmax L	inside	inside
name	[rad]	direction	[rad]	[rad]	[rad]	[rad]
BRBU1SE	0.0075	L	0.18311	0.33311	no hits	no hits
BRK1SE	0.00145	L	0.144824	0.167069	0.279992	0.283182
BRK2SE	0.00145	L	0.067584	0.112075	0.096549	0.115519
BRK3SE	0.00145	L	0.15959	0.179214	0.356008	0.356008
BRK4SE	0.00145	L	0.073657	0.112905	0.110485	0.127403
BRBU2SE	0.0075	R	0.278052	0.113052	no hits	no hits
BRBU3SE	0.0075	R	0.150948	0.011695	no hits	no hits
BRBU4SE	0.0075	L	0.007797	0.009882	no hits	no hits

Table 3: Maximum angles of lost beams and beams exiting the ES magnets at 150 MeV. The 4th and the 5th columns give angles at magnet exit. The 6th and the 7th columns give angles of the beams lost inside the magnet. "L" and "R" stand for left and right bending direction.

The described tracking is performed for each beam centroid (we fill the initial phase space with 2500 centroids). The currents of the chained magnets are scanned in the full range of respective PS.

Finally the maximum angles of the beams lost inside and exiting each magnet are determined. We are interested in both the horizontal angle to the left and angle to the right of beam direction of motion. Results of this analysis are presented in the Table 3.

4.2 Beam angles analysis at high energy

At high energy (here we consider 2 GeV and 3 GeV beams) the electron beam is stored in the Booster. This shrinks the phase space available to the beam centroids at the entrance to ES magnets significantly. Indeed, prior to the moment the fast magnets are turned on the stored beam in the ES shall be within the acceptance of 8.15 m long drift. Therefore, when the magnet is switched on the borders of the phase space at its entrance are defined by the following equation:

$$\theta_{max1}(h) = \frac{\frac{H}{2} + h}{L_{1}}, \quad -H/2 \le h \le -H(1/2 - L_{1}/L)$$

$$\theta_{max2}(h) = \frac{\frac{H}{2} - h}{L - L_{1}}, \quad -H(1/2 - L_{1}/L) \le h \le H/2$$

$$\theta_{min1}(h) = \frac{\frac{H}{2} + h}{L - L_{1}}, \quad -H/2 \le h \le H(1/2 - L_{1}/L)$$

$$\theta_{min2}(h) = \frac{\frac{H}{2} - h}{L_{1}}, \quad H(1/2 - L_{1}/L) \le h \le H/2$$
(1)

The meaning of H, h, L and L_1 is illustrated in Figure 9.

Although ES bumps are slow in comparison to beam revolution time, the phase space defined by eq. 1 can be immediately used at BU1 entrance since the four bumps are designed to create only the localized distortion of the stored beam trajectory.

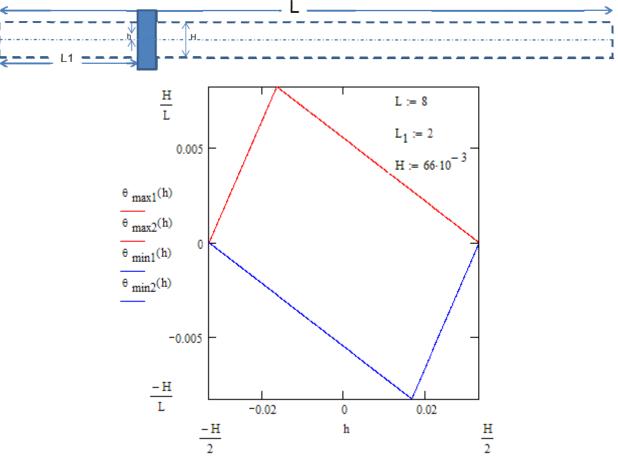


Figure 9: The schematic of ES drift and the phase space in front of the ES magnet for the stored beam.

For an entrance to K1 we use the phase space described above increased in positive angles by the maximum value available from BU1 kick. Similarly, for K3 the respective phase space shall be increased by maximum summed kicks of BU1, K1 and K2, and for BU2 entrance, the phase space must be increased by the sum of maximum kicks of BU1 and K1-K4.

Such approach significantly reduces the initial phase space available to beam centroids as compared to transfer line approach, while guaranteeing that all possible beams were taken into account.

The rest of the tracking procedure for high energy beams is identical to the one described for 150 MeV beams. The results of beam angles analysis are given in Table 4.

				2 GeV	beam			3 GeV	beam	
					θmax R	θmax L			θmax R	θmax L
element	θ nom	bending	θmax R	θmax L	inside	inside	θmax R	θmax L	inside	inside
name	[rad]	direction	[rad]							
BRBU1SE	0.0075	L	0.008098	0.020473	no hits	no hits	0.008098	0.016348	no hits	no hits
BRK1SE	0.00145	L	0.008098	0.022866	0.008098	0.022866	0.008098	0.017943	0.008098	0.017305
BRK2SE	0.00145	L	0.008098	0.025258	0.008098	0.024301	0.008098	0.019538	0.008098	0.0189
BRK3SE	0.00145	L	0.008098	0.027651	0.008098	0.027172	0.008098	0.021133	0.008098	0.021133
BRK4SE	0.00145	L	0.008098	0.030043	0.008098	0.030043	0.008098	0.022728	0.008098	0.022409
BRBU2SE	0.0075	R	0.020473	0.032518	no hits	no hits	0.016348	0.022728	no hits	no hits
BRBU3SE	0.0075	R	0.030218	0.020526	no hits	no hits	0.024598	0.017413	no hits	no hits
BRBU4SE	0.0075	L	0.010573	0.015038	no hits	no hits	0.010335	0.014478	no hits	no hits

Table 4: Maximum angles of lost beams and beams exiting the ES magnets at 2 GeV and 3 GeV. The 4th, 5th, 8th and 9th columns give angles at magnet exit. The 6th, 7th, 10th and 11th columns give angles of the beams lost inside the magnet. "L" and "R" stand for left and right bending direction.

5 Analysis of beam angles in the Booster to Storage Ring transfer line

To determine the beam centroids phase space at the entrances of elements of both phases of BSR line we adapt the transfer line approach described in section 4.1. Next, we track all the beam centroids through respective element. Each beam centroid is tracked through the magnet at its nominal, maximum and minimum possible settings. Finally, we determine the maximum angles of the beams lost inside the magnet as well as the maximum horizontal angles of the beams exiting the magnet both to the left and to the right direction with respect to the direction of beam motion.

For the dipole magnet, the uniform field is used in tracking, which gives a conservative estimate of angles in the bending direction. In the direction opposite to the bending the maximum angle occurs for the switched off magnet.

For the quadrupole the beam hitting the vacuum chamber is considered to be lost, since the space around the chamber is closely packed with material of quad poles and windings.

In the following calculations we assume that the quadrupoles can switch polarity and that the bends always have designed polarity.

The described routine is performed at two energies - at 2 GeV and at 3 GeV. The lowest energy at which one might intentionally extract the beam into BSR dump line is 2 GeV. Apparently, the maximum angles of beams are obtained at that energy. We also perform 3 GeV analysis, since it is the nominal Booster extraction energy.

The described procedure was also utilized for analysis of both pulsed and DC septa (both in Booster extraction and Storage Ring injection), since all these magnets are stand-alone bends.

Figure 10 illustrates the tracking procedure on example of bend B1. The results of performed analysis are summarized in Table 5.

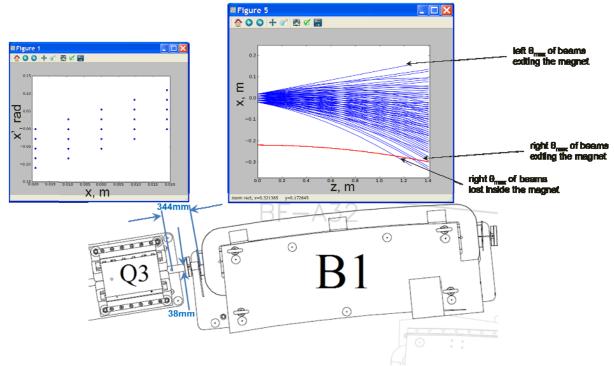


Figure 10: The phase space available to the beam centroids at the entrance of B1 (left plot) is defined by the geometric acceptance of Q3-B1 drift (its horizontal dimensions are given on the lower schematics). The phase space is filled with 25 beam centroids (for the clarity of the picture we reduced the number of beam centroid usually used in tracking by a factor of 100). Each of these beams is tracked through B1 (right plot) at its minimum, nominal and maximum settings. The red line on the right plot represents B1 yoke.

References

- [1] BR-PH-LAT-0003 (REV B)
- [2] NSLS-II Booster Final Design Report
- [3] G. Wang, BSR lattice file
- [4] Private communications with G. Ganetis
- [5] Private communications with R. Fliller
- [6] Private communications with B. Wahl
- [7] Private communications with M. Johansen

		Geo	metry		Magnetic Element Data					2 GeV beam				3 GeV beam				
			Vacuum	Ref. traj. to														
element	S @ magn.	Magnetic	chamber H	I yoke dist.	PS max	B or B' max	B or B' nom		bending	K1nom	θmax R	$\theta \text{max L}$	θmax R	θmax L	θmax R	θmax L	θmax R	θmax L
name	entr. [m]	[m]	app [mm]	[mm]	current [A]	[T] or [T/m]	[T] or [T/m]	θnom [rad]	direction	[1/m^2]	[rad]	[rad]	inside [rad]	inside [rad]	[rad]	[rad]	inside [rad]	inside [rad]
BRSP1SE	0.00	0.	5 1	5 9	11242***	0.88	0.80	0.048	L		0.051	0.042	0.125	0.088	0.053	0.030	0.067	0.088
SP2	1.23	1.	2 3	140	500	0.98	0.87	0.104	L		0.164	0.134	N/A **	no hits *	0.164	0.074	N/A	no hits
Q1	3.34	0.3	5 3	3 20	175	29.92	-5.69			-0.853	0.082	0.082	0.281	0.281	0.080	0.080	0.278	0.278
Q2	4.22	0.3	5 3	3 20	200	34.03	11.69			1.754	0.058	0.058	0.078	0.078	0.053	0.053	0.075	0.075
B1	4.92	1	1 4	5 221	375	1.34	0.55	0.115	R		0.237	0.225	0.255	N/A	0.179	0.223	0.184	N/A
Q3	6.63	0.3	5 3	3 20	175	29.97	0.93			0.140	0.066	0.066	0.124	0.124	0.063	0.063	0.122	0.122
B2	7.59	1.	1 4	5 221	375	1.34	-0.73	0.154	L		0.249	0.210	N/A	0.223	0.245	0.129	N/A	no hits
Q4	9.87	0.3	5 3	3 20	175	29.99	-6.39			-0.958	0.047	0.047	0.054	0.054	0.042	0.042	0.048	0.048
Q5	12.61	0.3	5 3	3 20	200	34.21	10.46			1.570	0.049	0.049	0.057	0.057	0.044	0.044	0.049	0.049
Q6	14.48	0.3	5 3	3 20	200	31.66	-10.35			-1.552	0.040	0.040	0.045	0.045	0.034	0.034	0.037	0.037
Q7	16.64	0.3	5 3	3 20	175	28.57	5.34			0.800	0.038	0.038	0.043	0.043	0.032	0.032	0.035	0.035
Q8	18.73	0.3	5 3	3 20	200	34.11	-14.49			-2.174	0.040	0.040	0.045	0.045	0.033	0.033	0.037	0.037
В3	19.44	1.	1 4	5 221	375	1.34	0.68	0.143	R		0.218		0.237	N/A	0.149	0.241	no hits	N/A
Q9	21.15						4.72			0.708					0.065			
Q10	22.72						1.47			0.221					0.049			
B4	23.35				. 375		0.73		R		0.215			,	0.162			′
Q11	25.79						10.69			1.604				-	0.078			
Q12	26.86						-9.86			-1.479					0.047			
Q13	34.21						11.26			1.688					0.042			
Q14	34.86						-8.76			-1.314					0.067			
SP3	36.78						0.95				0.115			,	0.056			,
IS	39.60	1.			8690****	0.94	0.85		R		0.023				0.015			′
Q1BD	10.36						0.00			0.000					0.026			
Q2BD	10.96		5 3				0			0.000	0.070	0.070	0.156	0.156	0.069	0.069	0.154	0.154

^{* &}quot;no hits" - None of the possible beam trajectories cross the yoke of the magnet

Table 5: Maximum angles of lost beams and beams exiting the BSR magnets at 2 GeV and 3 GeV. The 12th, 13th, 16th and 17th columns give angles at magnet exit. The 14th, 15th, 18th and 19th columns give angles of the beams lost inside the magnet.

^{** &}quot;N/A" - No magnet yoke in this direction

^{***} Nominal 10220 [A]+10%

^{****} Nominal 7900[A] +10%

^{1.) &}quot;R" stands for deflection to the right w.r.t. beam direction, "L" stands for deflection to the left w.r.t. beam direction

^{2.)} All deflection angles are given w.r.t. beam direction downstream of the magnet