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Materials Degradation and Detection (MD2): Deep Dive Final Report

JS McCloy
RO Montgomery
P Ramuhalli
RM Meyer
S Hu

Y Li
CH Henager, Jr.
BR Johnson

February 2013



Pacific Northwest
NATIONAL LABORATORY

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Richland, Washington 99352

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Acronyms

ABI	automatic ball indenter
ACPD	alternating current potential drop
AE	acoustic emission
AFAM	atomic force acoustic microscopy
AM	acoustic microscopy
AMR	anisotropic magneto-resistance
APT	atom probe tomography
ASME	American Society of Mechanical Engineers
B&PV	boiler and pressure vessel code
BEMM	ballistic electron magnetic microscopy
BMG	bulk metallic glass
BWR	boiling water reactor
DBTT	ductile-brittle transition temperature
DCPD	direct current potential drop
EBSD	electron back-scatter diffraction
EC	eddy current
EDS	energy dispersive spectroscopy
EMAT	electromagnetic acoustic transducer
FORC	first order reversal curve
GMR	giant magneto-resistance
HF	high frequency
IASCC	Irradiation-assisted stress corrosion cracking
IGSCC	Intergranular stress corrosion cracking
IR	infrared
ISI	in-service inspection
LVDT	linear variable differential transformer
LWR	light water reactor
MAE	magnetic acoustic emission
MBN	magnetic Barkhausen noise
MD2	Materials Degradation and Detection
MFM	magnetic force microscopy
MOKE	magneto-optic Kerr effect
MOS	magneto-optic sensor
MTR	materials test reactor
N/A	not applicable
NDE	non-destructive evaluation
NLU	nonlinear ultrasonics
ODS	oxide dispersion strengthened

OSURR	Ohio State University Research Reactor
PEEM	photoelectron emission microscopy
PFM	piezoresponse force microscopy
PIE	post irradiation evaluation
PMMD	Proactive Management of Materials Degradation
PNNL	Pacific Northwest National Laboratory
PWR	pressurized water reactor
PZT	lead zirconate titanate
RPV	reactor pressure vessel
SAM	scanning acoustic microscopy
SANS	small angle neutron scattering
SAXS	small angle x-ray scattering
SCC	stress corrosion cracking
SEM	scanning electron microscopy
SEMPA	scanning electron microscopy polarization analysis
SHPM	scanning Hall probe microscopy
SNOM	scanning near-field optical microscopy
SOTA	state-of-the-art
SP-STM	spin-polarized scattering tunneling microscopy
SQUID	superconducting quantum interference device
SSM	scanning SQUID microscopy
TEM	transmission electron microscopy
TGSCC	transgranular stress corrosion cracking
TMR	tunnel magneto-resistance
UFM	ultrasonic force microscopy
UQ	uncertainty quantification
US	ultrasonic
V&V	verification and validation
XMCD	x-ray magnetic circular dichroism

1 Executive Summary

An effort is underway at Pacific Northwest National Laboratory (PNNL) to develop a fundamental and general framework to foster the science and technology needed to support real-time monitoring of early degradation in materials used in the production of nuclear power. The development of such a capability would represent a timely solution to the mounting issues operators face with materials degradation in nuclear power plants. The envisioned framework consists of three primary and interconnected “thrust” areas including 1) microstructural science, 2) behavior assessment, and 3) monitoring and predictive capabilities. A brief state-of-the-art assessment for each of these core technology areas is discussed in the paper.

The focus of this effort is to move beyond the current approach of in-service inspection through periodic nondestructive examination (NDE) of structural materials by developing the ability to use real-time monitoring of material degradation. The intent of real-time monitoring of materials degradation is to provide a better understanding of the surface and volumetric material changes occurring during the early stages of the incubation and micro-damage accumulation. By detecting the presence of the conditions conducive to the material degradation mechanisms at an early stage in the process, better insights are gained about the state of the material that can be used to introduce actions to mitigate further progression of the degradation mechanisms or understand the margin to failure. However, important scientific and technological breakthroughs are needed in characterizing, detecting, and understanding the early phases of environmental degradation in materials used in the production of nuclear energy. A critical step in achieving this objective is to develop an appropriate means to detect minor changes in material microstructures at the onset of degradation. In this case “appropriate detection” would need to satisfy several criteria: 1) be non-destructive; 2) be able to probe large volumes; 3) have sub-millimeter spatial resolution; 4) be capable of real-time/ on-line monitoring; 5) be able to operate remotely; 6) be able to withstand radiation environments at high temperatures.

Sections 2 and 3 of this report are adapted from a paper presented at the 2012 3rd International Conference on Nuclear Power Plant Life Management (PLiM) discussing the outline of this proposed program, with additional information added. Following, we present two detailed case studies of candidate NDE technologies: magnetic methods and acoustic methods. In each category we discuss some of the theory related to sensing, some of the modeling work associated with understanding sensor materials and irradiation effects, and some of the methods for imaging microstructural properties related to the sensing modality. Finally, we discuss the research gaps and a proposed path forward, and acknowledge some complementary work being performed at other institutions.

2 Introduction and Motivation

The loss in capacity factor due to material degradation issues for PWRs and BWRs has generally decreased over the last few decades (Figure 1), as plants have replaced problematic components and improved their inspection and maintenance programs (Ford 2007). Nonetheless, an increase in capacity factor of 1% in the U.S. nuclear fleet has the potential to offset ~20 billion pounds of CO₂ (or other greenhouse gas) emissions each year (OECD 2011). However, it is anticipated that costs associated with management of materials performance will continue to increase as plants age. One method to improve capacity factors of nuclear power plants and to reduce the costs associated with managing materials performance is through the implementation of successful programs for the real-time monitoring of materials degradation.

The trend in component aging issues over time has shown that although fewer failures have occurred, the remaining mechanisms have increasing significance (Staehele 2012). Cracking in smaller components characterized early failures (such as sensitized stainless steel small-bore piping and bolts in reactor internals). These types of cracking, while still prevalent, represent a smaller fraction of component or material failures at present and have been mitigated through managing the operating conditions (water chemistry, etc.) as well as by replacement of components. At present, the major issues are with respect to large components (such as the pressure vessel head, pressurizer surge lines, welds in large bore piping, etc.), which appear to be impacted by late-onset forms of material degradation. Continued reliability of these components are challenging as they represent high-value components that are safety-critical, and as such, can result in expensive shut-down of operating plants. An estimate of the costs of component failures for BWR reactors is shown in Table I.

Current in-service inspection practices are based on requirements in the ASME B&PV Code which were originally developed in the 1960's. The driving philosophy of the Code requirements is the management of fatigue degradation (Doctor 2008). The code specifies whether volumetric, surface, or visual examinations should be performed depending on the functional requirements of the component and the risks associated with its failure. In addition to the type of inspection, the code specifies the frequency of periodic exams and sampling criteria for performing exams of class 1, 2, and 3 components. The philosophy of current ISI requirements is challenged by the existence of diverse and challenging degradation mechanisms in nuclear power plants, such as stress corrosion cracking (SCC). In addition, as plants continue to age, it can be anticipated that new degradation mechanisms will continue to manifest in components (Wilkowski et al. 2002).

While current in-service inspection processes have been successful at detecting cracks in large-scale components, it is almost always after the fact (i.e., after the material degradation has reached a stage where replacement of the component is the only option). An example of this situation is the corrosion and subsequent cracking in the weld region of the control rod drive mechanism penetrations in the reactor vessel head at Davis-Besse. The environmental degradation incurred during operation resulted in the replacement of the entire reactor vessel

head (IAEA 2007). The consequences of the material wastage due to boric acid corrosion may have been mostly averted had real-time monitoring of crack formation and growth been available to provide indications of material degradation leading to leakage of primary coolant onto the component. Further reviews find that experience supports a need for the development of real-time materials degradation detection technologies that can monitor environmentally assisted material degradation of reactor structural components and provide sufficient lead time for maintenance, repair, or other mitigation activities. The program under development at PNNL addresses this need. An associated issue is the ability to determine critical locations to place monitoring capabilities to ensure that degradation that has immediate impacts on safety is detected in a timely fashion. In this, the PNNL effort will leverage previous (and ongoing) activities in Material Degradation Assessment (Proactive Materials Degradation Assessment (PMDA), Extended Materials Degradation Assessment (EMDA), and the industry-led Materials Degradation Matrix (MDM)) (Andresen et al. 2007; Lain 2008; Pathania 2008; Stark 2008). These assessments provide an indication of critical locations and materials, as well as the known (and possible future) degradation mechanisms associated with these locations and materials. Gaps in the technical knowledge required to achieve materials degradation detection are also identified in this assessment.

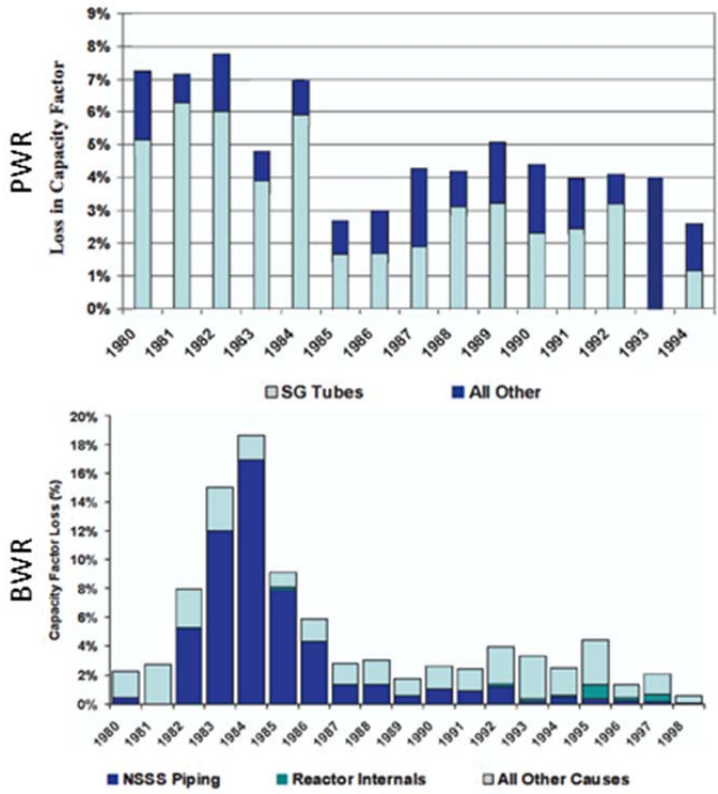


Figure 1. Plant capacity losses from materials degradation are >\$1B/ year loss

Table I. BWR component replacement cost in \$M (1999 cost basis, \$10K/man-rem, \$0.5M/day replacement power)

Component Replacement	Est. Repair Cost	Rad. Exposure Costs	Replacement Power Costs	Total Cost
Shroud Support to Vessel Weld	250	55	182	487
CRD Stub Tube to Vessel Weld	50	7	113	170
In-Core Housing	50	2.25	50	102
Recirculation Piping	20	12	63	95
Recirc. Inlet/Outlet Nozzles	25	3	22	50
Jet Pump Riser Brace	10	0.26	32	42
Core Support Plate	4	0.3	25	29
Top Guide	8	0.15	15	23
Instrumentation Penetrations	6	2	13	21

3 Overall approach to develop real-time monitoring of materials degradation

The framework that is being used to develop the science and technology needed for real-time monitoring of materials degradation in nuclear power reactors consists of three main interconnected thrust areas, described below and illustrated in Figure 2:

- (1) **Microstructural science:** This thrust area will leverage developments in computational methods at the meso-scale and analytical and nano-scale microscopy to understand the onset and progression of prototypical degradation mechanisms in structural materials, under the influence of temperature, neutron irradiation, chemical attack, and other stressors. The focus will be on understanding degradation precursors (i.e., material microstructural changes that occur prior to crack coalescence and growth). A key area of concentration is the development of meso-scale quantitative property measurement techniques (physical property or micro-property imaging) to measure magnetic, acoustic, and electrical properties on sub-micron length scales.
- (2) **Behavior assessment:** This thrust area will develop multi-scale, multi-physics computational models of material behavior changes caused by degradation mechanisms. Meso-scale models of degradation evolution and the resulting changes in physical properties of the material (such as acoustic, magnetic, and electrical) will be coupled to macro-scale models of the sensing physics to obtain insights into the sensitivity of measurement methods

to early stages of damage. Advances in quantitative micro-property imaging technologies (e.g., magnetic force microscopy, acoustic force atomic microscopy, etc.) from thrust area (1) will be leveraged to verify and validate the models.

- (3) **Monitoring and predictive capabilities:** The focus of this area is to leverage materials science, computational methods, and physical property imaging [including the developments in thrust areas (1) and (2)] to develop environmentally hardened sensors to monitor materials behavior under harsh conditions. In addition, the macroscopic measurements from these sensors will be utilized, along with the multi-scale, multi-physics models, to develop inverse analysis methods that attempt to quantify the level of damage at the microstructural level.

At the heart of each of these thrust areas are common research activities: computational modeling and model-based inverse analysis, micro-scale physical property measurements, and engineering-scale nondestructive measurements. A state-of-the-art assessment of these core technology areas is discussed next, along with a proposed plan for addressing the S&T needs for developing continuous monitoring technologies.

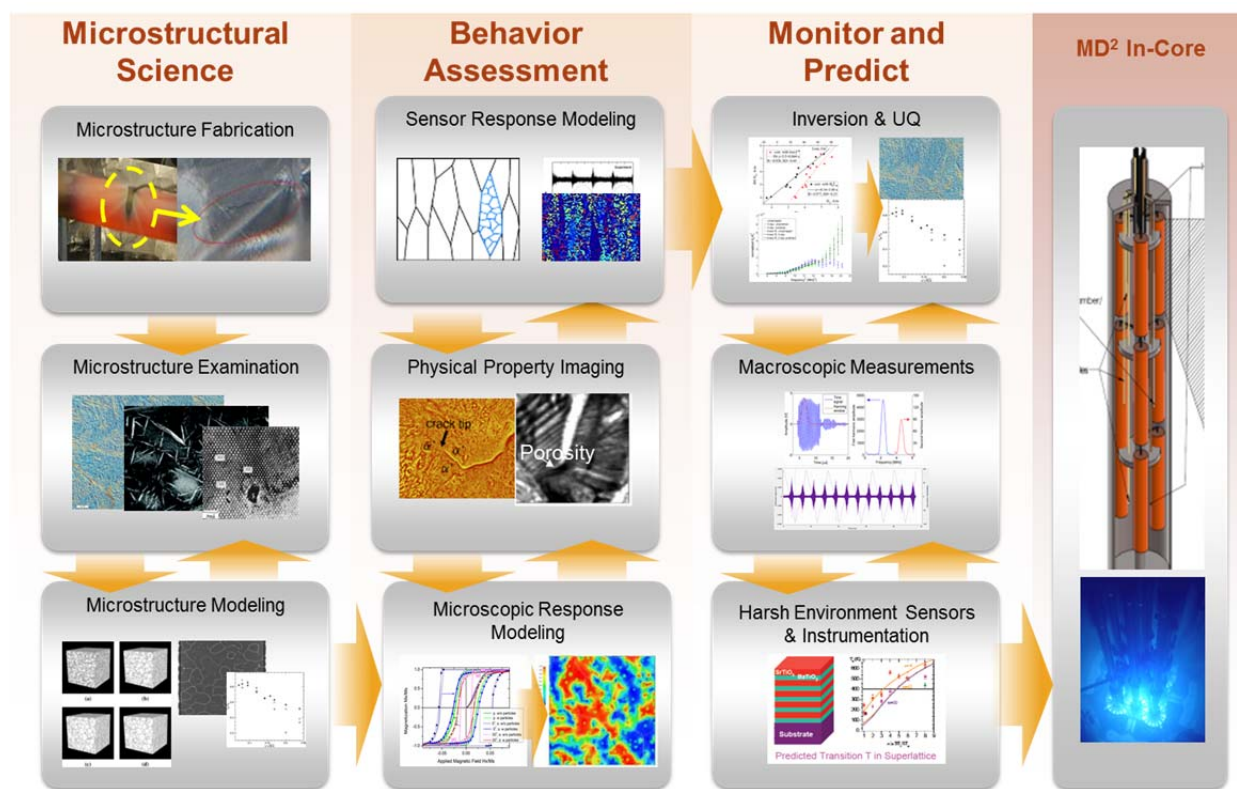


Figure 2. Conceptual illustration of the integrated parts of Material Degradation and Detection (MD2)

3.1 Gaps and Path Forward

Given the available technical literature in this area, and the gaps identified in the different thrust areas, the following interconnected approach is proposed and shown schematically in Figure 3, with an example for acoustic NDE, showing specific gaps where verification and validation (V&V) are needed. Note: reference acronym section as appropriate.

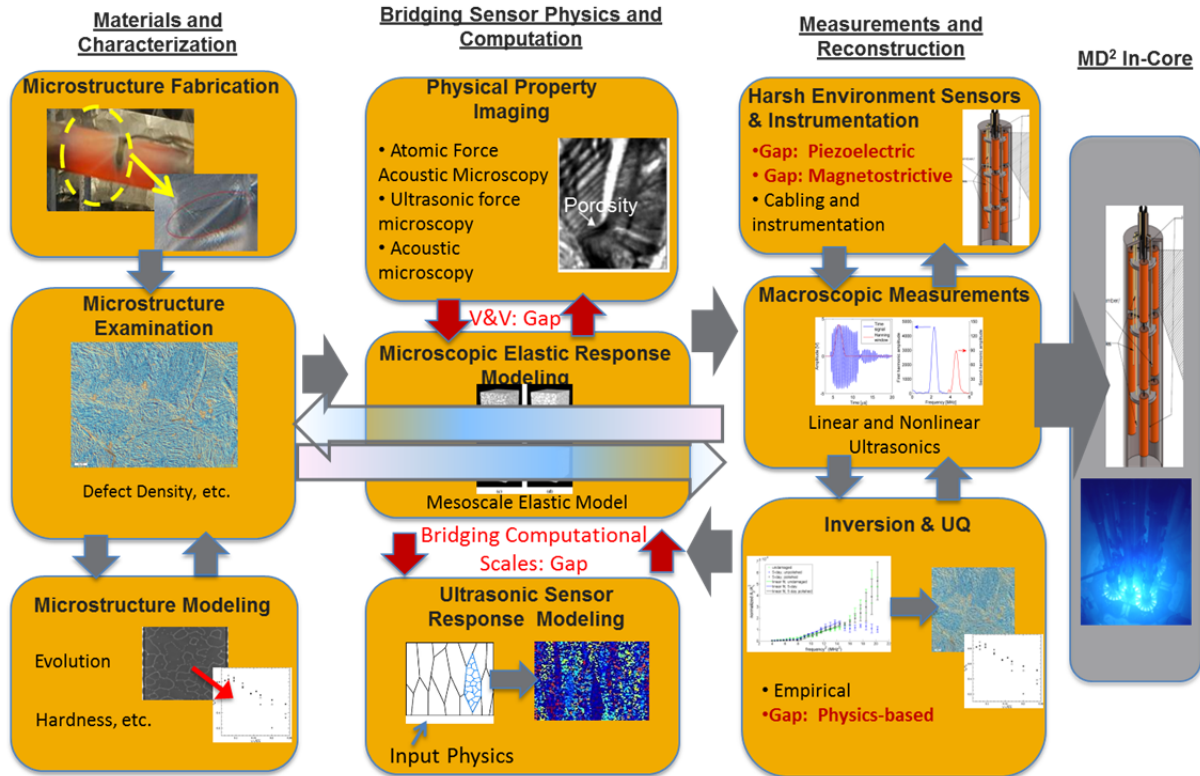


Figure 3. Schematic of approach for acoustic NDE sensor evaluation

1. Microstructure fabrication – Specimens are fabricated under known environmental conditions (known stress, neutron flux, temperature, etc.). Outcome: A set of specimens with different forms of microstructure, with and without degradation (used in steps 2 and 8).
2. Microstructure examination – Microscopy (imaging) used to characterize the microstructures. Tools to include APT, TEM, Environmental TEM, etc. Outcome: set of parameters defining the microstructure. Should also include thermodynamic and kinetic data (of use in steps 3 and 4).
3. Microstructure Modeling: Atomistic/phase field models of microstructure evolution. Outcome: Basic quantity such as hardness.

4. Microscopic response modeling: Phase field model providing the response to applied stress at meso-scale. Outcome: Elastic modulus tensor C_{ijkl} or magnetic hysteresis loop (or other relevant quantity).
5. Physical Property Imaging: Microscale imaging of elastic moduli variation (or magnetic permeability variation). Tools include MFM, AFAM, UFM, AM, etc. Data used to validate phase field models. **Gap:** V&V at these length scales. The tools identified above are likely to provide the necessary validation data for mesoscale models.
6. Sensor Response modeling: Computational Models of applied energy interaction with material at bulk scales. For instance, modeling of ultrasonic wave propagation in the bulk material. Models need to be coupled to mesoscale models (Step 4) to integrate results from mesoscale models for bulk-scale predictive modeling. Tools for macroscale modeling include FEM, BEM, etc. to solve wave equation in anisotropic media. Outcome: Bulk-scale computational model of sensor physics. **Gap:** Bridging computational scales (from meso to macro). Multiscale modeling approaches may offer a route to bridging this gap.
7. **Gap:** Harsh environment sensors and instrumentation: Combination of modeling and experiment to characterize and develop sensor materials, sensors and instrumentation (piezoelectric, magnetostrictive, etc.) for harsh environment operation (high temp/rad). Tools include computational models and materials characterization systems being investigated currently. AFAM, etc. are useful tools to examine piezoelectric materials at microscopic scales during sensor operation. Outcome: If successful, a method to characterize sensor materials for performance in different environments, and design robust sensors. Used in step 8.
8. Macroscopic measurements: Macroscopic measurements on different microstructures (step 1) using sensors (step 7). Outcome: Data sets for use in bulk-scale model (step 6) V&V as well as microstructure reconstruction (step 9).
9. Inversion and uncertainty quantification (UQ): Using sensor data at bulk scale, determine state of material and associated uncertainty. Tools include computational inverse models and characterization tools from previous steps, signal processing tools. Outcome: Robust approaches to fieldable inversion and UQ tools. **Gap:** Inversion to determine material state with minimal historical (trending) information. **Gap:** Uncertainty quantification.

3.2 Material degradation modes in LWR primary coolant systems

The majority of the structural components used in the primary system of LWRs are fabricated from the family of Fe-Cr-Ni metals often employed in many high-temperature

corrosive environments. Several varieties of these materials are implemented including, carbon steels (e.g., A553B), stainless steels (e.g., 304 and 308), and nickel-based alloys (e.g., Alloy 600 and X750) and the corresponding weld materials used to join them. In general, these materials have exhibited excellent performance in high temperature and pressure water and steam systems. However, the unique combination of high temperature, corrosive water chemistry, and neutron flux leads to a harsh environment and activates materials degradation processes that reduce the capacity of these components to perform their desired operating and safety functions. The material degradation modes active in LWR primary coolant systems can be grouped into the following categories: generalized corrosion, localized corrosion, and mechanical. Table II summarizes the major types of material degradation mechanisms and the materials most susceptible in LWR systems.

The most impactful of the degradation modes shown in Table II is stress corrosion cracking (SCC), which has several different variants, including intergranular (IGSCC), transgranular (TGSCC), and irradiation-assisted (IASCC). Substantial research has been performed to identify the underlying causal factors that result in SCC failure in the Fe-Cr-Ni system. The research has found that SCC can become active once the conditions of environment (temperature, chemistry, neutron radiation, etc.), material susceptibility (microstructure, impurities, surface conditioning), and stress (residual, applied, etc.) align into a unique combination. Because of the large number of factors that influence SCC, quantitative predictions of the progression of the degradation process have been difficult. However, a qualitative understanding of the process has been developed and has been described by Staehle (Staehle 2012). While the progression given below is for the SCC mechanisms, the overview is general and applies to most material degradation processes.

The process of material degradation by SCC can be separated into 5 distinct stages (Staehle 2012). These stages are: (1) **Stage 1: Initial condition** – This stage represents the as-manufactured state of the material before exposure to the environment. The as-manufactured state includes surface characteristics such as finish (polishing, etching, oxidation, etc.), grain orientation/texture, second phase particles, and residual stresses. (2) **Stage 2: Precursors** – The precursor stage characterizes the exposure of the material to the operating environment. Upon exposure, certain processes may occur that promote the occurrence of SCC, such as surface pitting, depassivation of surface oxides, and buildup of irradiation damage (dislocations, amorphization of second phase particles, etc.). (3) **Stage 3: Incubation** – Progression from the precursor to the incubation stage represents the start of the active SCC process. During this stage the surface and microstructural changes occur that lead to crack initiation at the surface of the material. These changes include dislocation buildup, entry of oxygen or hydrogen, grain boundary migration of impurities, void formation, etc. (4) **Stage 4: Proto-cracks** – The formation of proto-cracks (cracks less than a few microns) results from vacancy coalescence, intergranular corrosion, etc. Proto-cracks start out as non-interacting with the applied or residual stress field, unlike the propagating cracks in Stage 5. (5) **Stage 5: Propagation** – This stage

represents the process in which the crack interacts both with the local stress field and the local chemical/irradiation conditions. Crack advance is influenced by the electrochemical potential within the crack environment, the metallurgical chemistry at the crack tip, the stress state, etc. Stable crack growth rates are observed in experiments designed to simulate the coolant chemistry and strain rates.

Table II. Degradation mechanisms and susceptible materials in LWR primary coolant systems

Mode	Degradation mechanism	Susceptible material/component
General corrosion	Boric acid corrosion	Carbon steel
	Erosion corrosion	Carbon steel, Alloy 600
	Flow-accelerated corrosion	Carbon steel, Alloy 600
Localized corrosion	General corrosion	
	Crevice corrosion	Carbon steel, Alloy 600
	Galvanic corrosion	
	Pitting corrosion	Carbon steel
	Stress corrosion cracking (IGSCC, TGSCC, IASCC, etc.)	Carbon steel, SS 304/316/308/309 CF8, Alloy 600, Alloy 82/182
Mechanical	Thermal/mechanical fatigue	Carbon steel, SS 304/316/308/309 CF8, Alloy 600, Alloy 82/182
	Thermal embrittlement	SS 308/309 CF8, Alloy 82/182
	Irradiation Embrittlement	Carbon steel, SS 304/316/308/309 CF8, Alloy 600, Alloy 82/182
	Irradiation/thermal creep	Carbon steel, SS 304/316/308/309 CF8, Alloy 600, Alloy 82/182

Current practices to manage the material degradation of primary system components in LWRs use periodic in-service inspection techniques to identify components that are in Stage 5 of the SCC process (Meyer et al. 2012). These practices characterize the presence of a crack and monitor the crack growth rate to ensure that actions are introduced prior to component failure. Because of the accelerating nature of crack growth, conservative actions are usually implemented (component repair or replacement) to ensure that rapid crack growth is prevented between inspection periods.

The goal of real-time monitoring of materials degradation is to provide a better understanding of the surface and volumetric material changes during Stage 3 (Incubation) and Stage 4 (Proto-cracks). By detecting the presence of the conditions conducive to SCC at an early stage in the process, better insights are gained about the state of the material that can be used to introduce actions to mitigate further progression of the degradation mechanisms or understand the margin to failure.

A critical step in achieving this objective is identifying a methodology that relates the changes in underlying microstructural characteristics caused by the degradation process to detectable variations in the response of the material to the external fields applied using NDE techniques such as acoustic, magnetic, and electrical methods designed to probe material structure. One approach to develop this capability is to couple the advances in meso-scale modeling of material microstructure behavior with recent advances in micro-property measurement techniques to correlate local microstructure changes to external field responses. The study underway at PNNL is centered on identifying the science and technologies required in this approach and the challenges to this development are outlined below.

3.3 Multi-scale computational modeling

Atomistic, meso-, and macro-scale modeling methods have been developed during the past few decades to study microstructure and property evolution in irradiated materials. The challenge in multi-scale modeling of material behavior is the development of physics-based modeling approaches that enable prediction of the three-dimensional microstructure evolution at the meso-scale, where time and length scales are larger than that of atomistic simulations, but smaller than that of macro-scale methods. The temporal range for typical meso-scale methods is seconds to a few hours, and the typical length scale is 10^2 's of nanometers to tens of micrometers. The phase-field approach (Figure 4) is one meso-scale method that has been successfully used in calculating the 3D microstructure evolution such as solidification, grain growth, martensitic transition, precipitation, ferroelectric/ferromagnetic transition, dislocation dynamics, deformation twin, and sintering (Chen 2002; Wang et al. 2001; Karma and Rappel 1998; Artemev et al. 2001; Hu et al. 2010). Atomistic simulations and experiments provide the thermodynamic and kinetic properties of the system that the phase-field approach requires for evaluating the microstructure behavior. The advantages of the phase-field approach are 1) the limited assumptions on microstructure morphology, 2) the ability to use a diffuse interface and topological changes instead of the difficulties of representing sharp interfaces, and 3) the continuous description of the energy landscape to take into account both short-range and long-range interactions.

At the macro-scale, an important application of simulation modeling is to understand the response of measurement sensor fields (e.g., acoustic, magnetic, and electric) to material damage. Most nondestructive evaluation measurements use the physics of stress wave propagation in materials, or the interaction of electric and magnetic fields with the material. Methods to model these interactions at the macro-scale include finite element models (Ludwig and Lord 1988; Chen and Konrad 1997), finite difference methods (Kunz and Luebbers 1993), and the method of moments (Jin 2002). In addition, semi-empirical models are also often used. Challenges in applying these methods include material inhomogeneity, anisotropy, and the stochastic nature of microstructural variation in materials (e.g., thermodynamic processes which are by nature statistical, such as Boltzmann's distributions). Most of the macro-scale sensor

models are based on material property variations that are on much larger length scales than typical microstructural changes that control the onset of degradation. To gain a better understanding of the early stages of material degradation, additional and more refined information about the microstructural characteristics must be obtained from macro-scale methods. The current macro-scale/empirical methods lack the capabilities to provide this information alone and must be adapted to address the accumulation of microstructural damage at the meso-scale. A potential approach to addressing this issue is to couple meso-scale and macro-scale modeling methods to achieve a multi-scale physical model of the interaction of the sensor field with the microstructure.

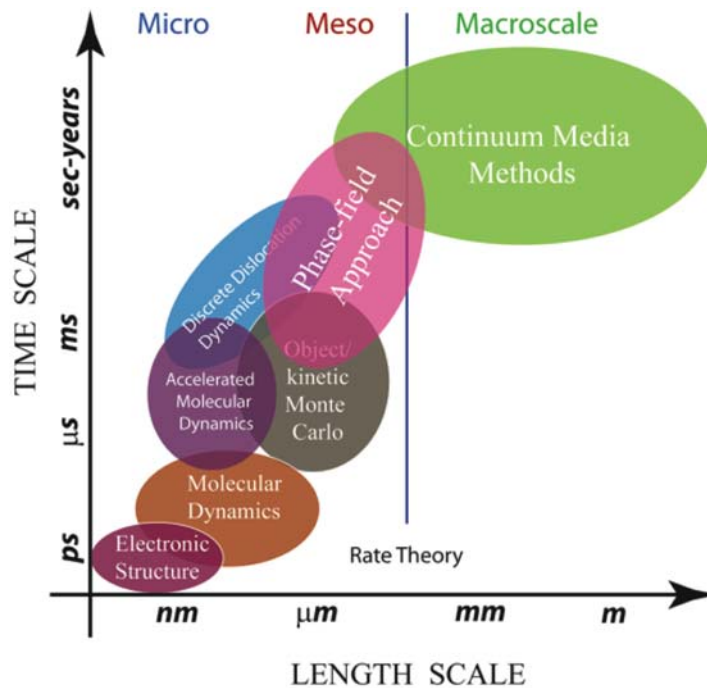


Figure 4. Rigorous meso-scale modeling of materials, phase field modeling (Hu et al. 2010)

A major difficulty in the area of computational modeling of material behavior is the ability to link simulations across length scales from a few microns to thousands of microns and higher. A number of approaches have been investigated to bridge the temporal and spatial scales from atomistic to macro-scale modeling methods. At one end of the bridging spectrum are indirect coupling methods that separately model at the meso- and macro-scales, and transfer information from one scale to the other with a passive process, using the results from lower-length scale models to inform macro-scale models. Alternatively, fully coupled simulations have been developed that use such tactics as algebraic multi-grid methodologies (Brandt 2001). The challenge in either of these cases is the ability to scale the physical processes in an accurate enough fashion to retain the appropriate level of information for diagnostic evaluations. The initial effort to enhance the modeling of sensor physics to gather higher fidelity information will

use the passive coupling approach to simplify the process while ensuring that only relevant information is passed from one scale to the next.

3.4 Microscopic physical property imaging

The processes that cause material degradation begin to occur in the microstructure at the atomic level as the atoms, ions, point defects, dislocations, grain boundaries, and second-phase particles respond and interact with the external chemical, thermal, mechanical, and irradiation environment. The ability to detect the microstructural evolution occurring in the early phases of material degradation requires identifying the impact of lower-length scale changes on key material properties at the sub-micron range. The capacity to characterize the microscopic physical properties (e.g., coercivity, elastic constants, and electrical conductivity) at the sub-micron/micron scale allows for 1) identifying the impact of microstructural changes during early stage degradation processes and 2) providing validation data to meso-scale modeling methods used to simulate the microstructure evolution during the degradation mechanisms.

Several microscopic physical property imaging techniques are derived from scanning probe microscopy capabilities, including magnetic force microscopy (MFM), atom force acoustic microscopy (AFAM), and piezoresponse force microscopy (PFM). These methods are able to resolve the impact of microstructure characteristics on nano-scale magnetization behavior (i.e., domain wall motion) and the nano-scale elastic material response. Other similar imaging methods which have somewhat lesser resolution include scanning Hall probe microscopy (SHPM), scanning SQUID microscopy (SSM), and scanning acoustic microscopy (SAM). By coupling the capabilities of these measurements with classical analytical microscopy methods, for example, transmission electron microscopy (TEM) and integrated, focused ion beam milling, scanning electron microscopy (SEM), electron backscatter diffraction (EBSD), energy dispersive spectroscopy (EDS), and atom probe tomography (APT), it now becomes possible to construct three-dimensional micro-property, micro-chemistry, and micro-crystallographic information that can be used to identify important relationships between microstructure evolution and material response. Figure 5 illustrates how the combination of structural imaging (optical and electron microscopy), chemical imaging (with microscope attachment such as EDS), and property imaging (magnetic, acoustic, and eddy current/ electrical conductivity) can combine to provide powerful diagnostic information about material degradation processes.

As mentioned, a key element in understanding the response of the microstructural characteristics (e.g., defect structure, second-phase particles, grain boundary evolution) to applied fields is the meso-scale phase-field modeling. The construction of micro-scale information that can be used to establish initial conditions for the analysis models is the first step in identifying the precursor signatures of material degradation. The second step uses microscopic physical property imaging to obtain the detailed response spectrum of characterized microstructures to applied fields.

Validation of the meso-scale phase field models is necessary to ensure that the physics of microstructural behavior is being modeled accurately. One example in the literature of validating phase-field models using micro-property measurements has been performed with ferroelectric materials. A phase field model of $\text{PbZr}_{0.2}\text{Ti}_{0.8}\text{O}_3$ films (Li et al. 2005) has been compared to images of the same material taken using PFM (Ganpule et al. 2002), showing agreement between experimental and calculated material response.

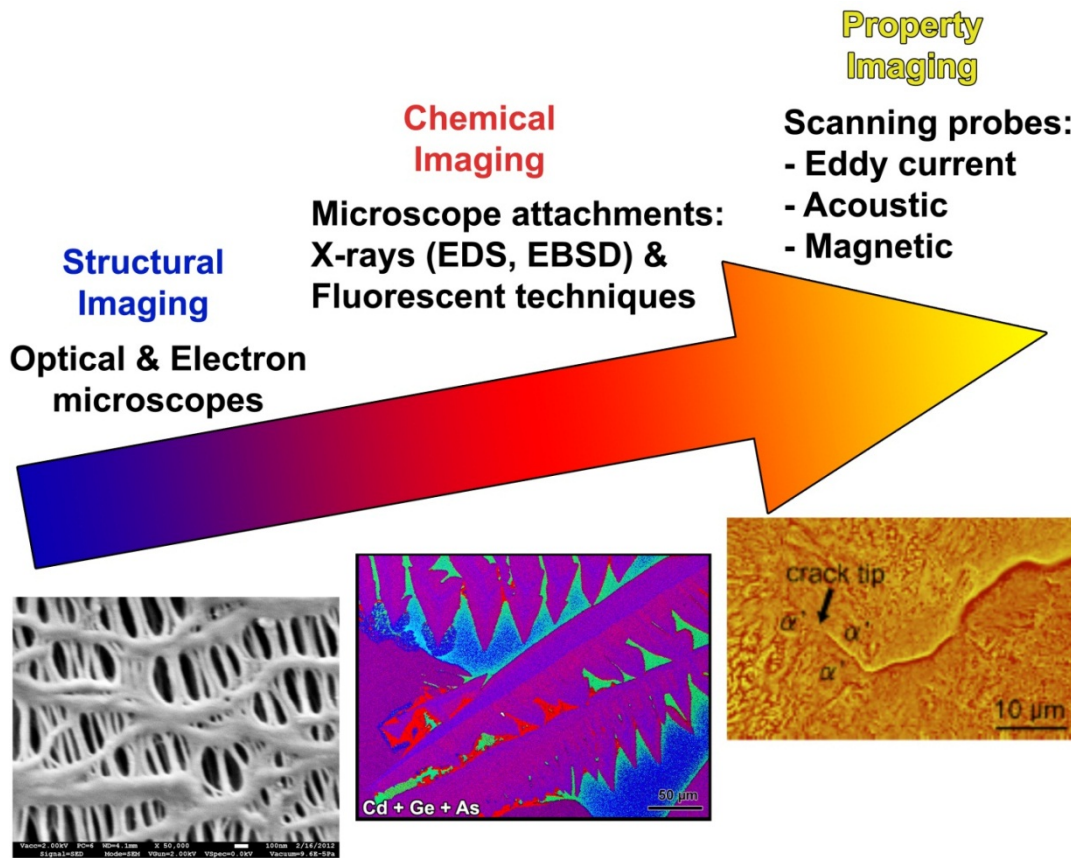


Figure 5. Illustration of the combined power of structural, chemical, and physical property imaging techniques

3.5 Engineering-scale degradation measurements - NDE

There are many possible NDE techniques which could be used to detect degradation and changes in material properties (see Figure 6). The main classes (and subclasses) of detection methods of interest are acoustic (linear and harmonic/nonlinear, with measurements of velocity/attenuation/scattering/nonlinear parameters), magnetic (Barkhausen, DC hysteresis, leakage flux), and electrical (eddy current, AC/DC potential drop, impedance tomography, impedance spectroscopy, etc.). Each detection method can be realized by multiple sensor types (e.g., magnetic: coils, Hall sensors, magneto-optic, giant magnetoresistance (GMR), anisotropic magnetoresistance (AMR), superconducting quantum interference device (SQUID); acoustic:

piezoelectric, magnetostrictive, EMAT=electromagnetic/acoustic). These NDE methods have been reviewed by many authors regarding their applicability for use in nuclear reactor plant degradation management. PNNL researchers have been engaged in the Nuclear Regulatory Commission's Proactive Management of Materials Degradation (PMMD) program and have assessed various technologies for NDE of different reactor components (Bond et al. 2009, 2011). NDE techniques for assessing creep damage in power plant steels have also been recently reviewed (Sposito et al. 2010), where the various techniques are compared for their utility in detecting flaws as a function of crack progression, building on earlier reviews (Dobmann 2006; Dobmann et al. 1992).

The relationship between conventional materials science investigation tools such as SEM/TEM, optical microscopy, and field NDE is shown in Figure 7. This figure highlights the trade-off that exists between volumetric sampling capability and scale of phenomena that the technologies may be sensitive to. The development of non-destructive material examination methods sensitive to the phenomena of degradation at early stages requires the identification of suitable observable signatures that correlate with changes in material condition indicative of damage (precursors). Such signatures could be local changes in electrical, mechanical, or thermal properties that "localize" before initiation of a macro-defect such as metal loss or crack formation. Material degradation manifests at the microstructural level in several ways and the influence of such microstructural phenomena on measurable bulk magnetic and elastic properties is described in the extensive review by Raj et al. (Raj et al. 2003). This review also highlights the important length scales for measuring some forms of degradation, and applicable NDE techniques to investigate these length scales (see Figure 8).

For a given technique, determining resolution and minimum detectable flaw size is a complex process. In general, higher frequency measurements will give better feature resolution, lower penetration (higher attenuation), be more surface sensitive, sample a smaller volume, and have a smaller signal-to-noise ratio. Physical constraints, such as skin depth for electromagnetic waves and impedance matching for acoustic waves, will also play a part in determining the best frequencies for measurement. In the reality of field measurements, as opposed to laboratory measurements, it is quite difficult to control environmental noise, so electronics noise must be decreased to the lowest level possible, and higher powers are required to increase signal levels. Ultimately, therefore, the resolution of a particular technique should not be confounded with its sensitivity (or detectability) to a particular length scale defect or feature. Additionally, having increased sensitivity does not necessarily imply increased selectivity for detecting features of interest. That being said, it is desirable to detect flaws in the pre-crack stage (Bond et al. 2011), before they have become large enough for detection using conventional NDE techniques (see Figure 7).

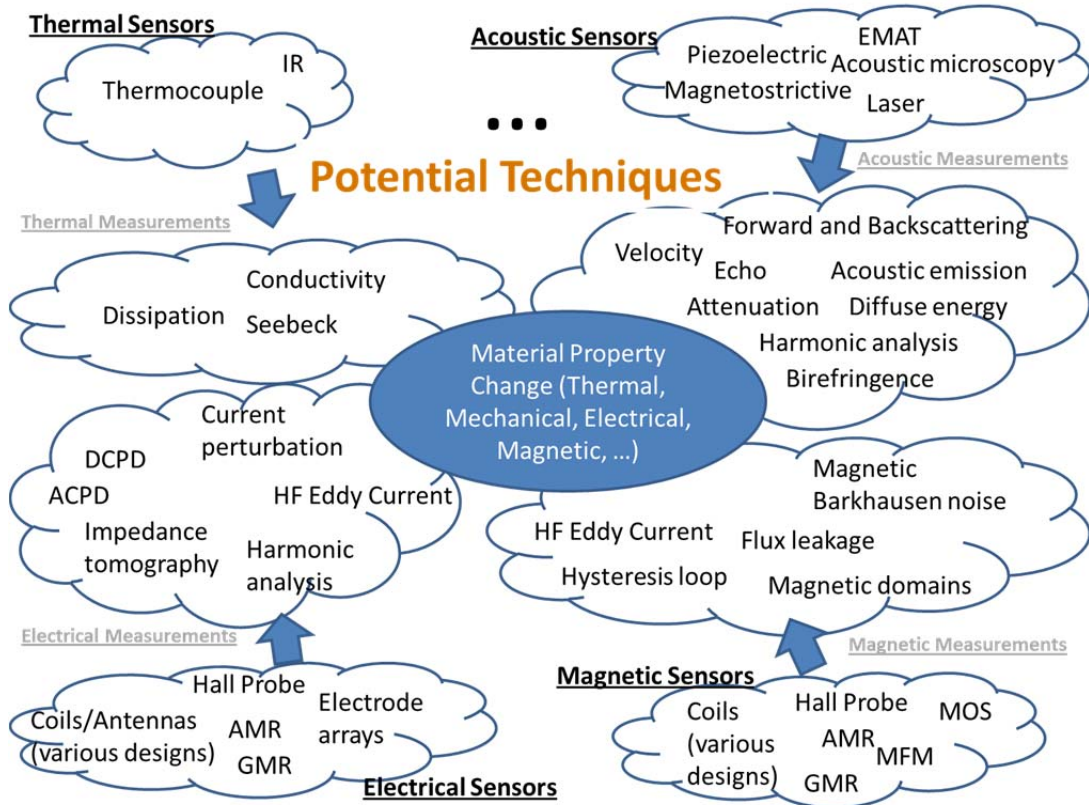


Figure 6. Potential NDE technologies and techniques (see acronym list)

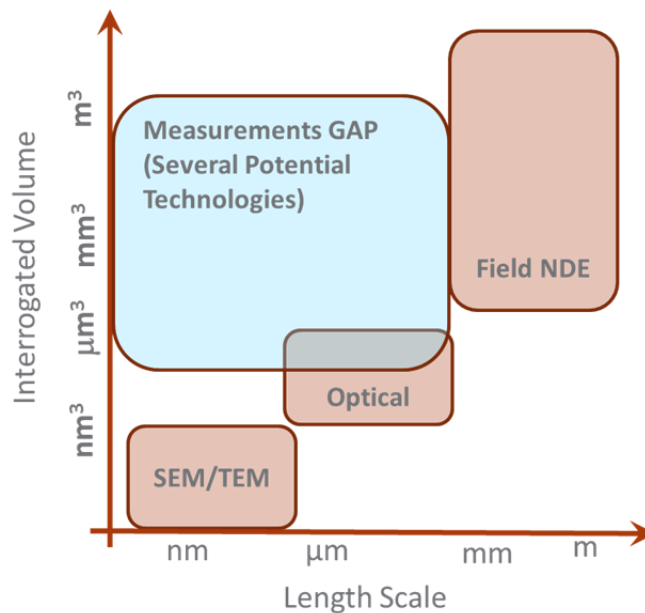


Figure 7. Conceptual description of current state of art in field NDE and the gap towards which the present program is geared

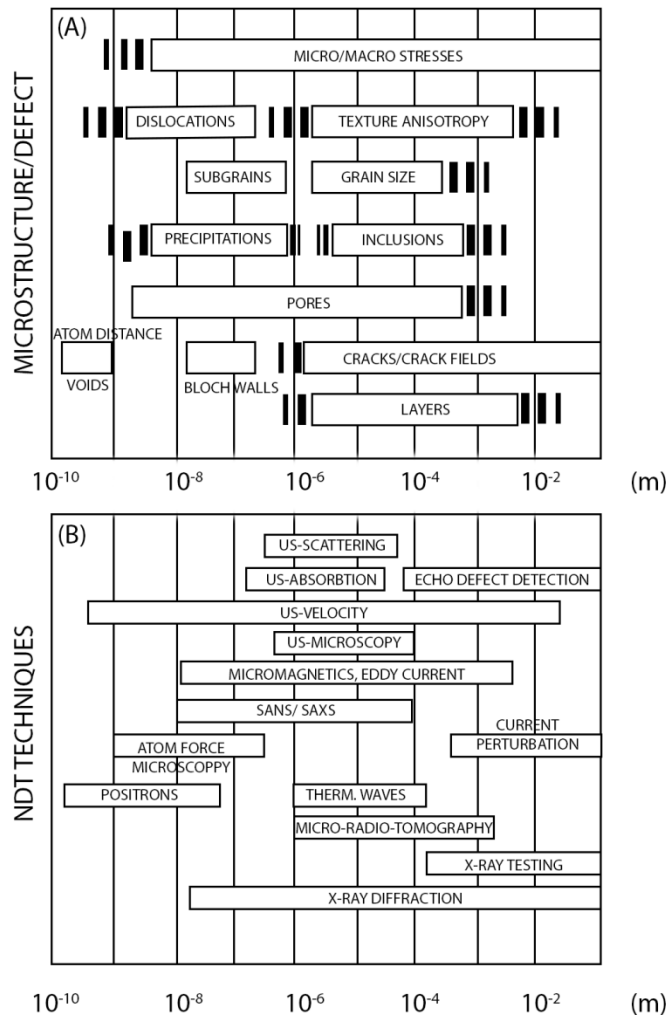


Figure 8. Microstructural defect length scales and applicable NDE techniques, after (Raj et al. 2003)

An added dimension to increasing sensitivity is the potential aging and degradation of the sensor itself, over long time periods. While this is generally not an issue for periodic in-service inspection technologies (differences in response due to sensor change-out can usually be calibrated out), the use of sensors for long-term condition monitoring in harsh environments is likely to result in a gradual change in the sensor response and sensitivity due to aging and degradation especially in regions of high temperatures and irradiation (neutron and gamma). Application of in-situ monitoring in nuclear reactors can conceptually be divided into two broad areas: (1) fuel and other **in-pile materials** and (2) **external structural materials** such as the reactor pressure vessel (RPV). The requirements on the sensor materials are quite different in

these two cases, as temperatures and radiation fields for in-pile materials are much higher. Aging effects in sensor materials may be compounded by similar effects in the cabling and instrumentation used for measurement. The level of drift and reduction in sensitivity is likely to be a function of the type of sensor (acoustic vs magnetic vs electrical) and the specific material used (for instance, PZT vs AlN for piezoelectric sensors (Parks and Tittmann 2011)). Gamma and/or neutron testing has been conducted for the production of radiation-hardened magnetic sensors for space applications, fusion reactors, and other nuclear applications. Similar studies have been conducted, though on a limited scale, for piezoelectric materials commonly used for ultrasonic measurements. While recent advances (Coble et al. 2012) may be used to monitor sensor drift, techniques to compensate for decreasing sensitivity may be needed to maintain the ability to monitor the materials/components over the long term. An assessment of the current state-of-the-art (SOTA) for detection of degradation and degradation precursors in various materials of interest to nuclear plants is shown in Table III. This table focuses on FeCr/Ni alloys, Zr-alloys, oxide dispersion strengthened (ODS) steels, and UO₂, as these constitute the dominant structural and fuel materials used (or planned for use) in current and future reactor designs. Listed are some phenomena of interest for degradation, field NDE SOTA, and microstructure changes of interest. Also listed are some examples of property changes of interest, such as magnetization (M), thermal conductivity (K), volume (V), oxygen level (Ox), and ductile-brittle transition temperature (DBTT).

Several NDE technologies have emerged as potential candidates to meet the requirements for early material degradation measurement. These include micromagnetic techniques such as magnetic Barkhausen noise and magnetic loop measurements. Substantial efforts to relate thermal, mechanical, and radiation embrittlement degradation to micro-magnetic responses are described by Raj et al. (Raj et al. 2003) and Dobmann et al. (Dobmann 2006). Nonlinear acoustic/ultrasonic techniques have been shown to be more sensitive to micro-damage than conventional linear ultrasonic techniques (Jhang 2009). The development of practical NDE technologies for early degradation monitoring is in its infancy and one critical gap is the ability to interpret the resulting measurement signals. Empirical correlative analysis (e.g., Figure 9) is often performed on signals obtained over a range of conditions to relate measurements to material condition. The development of physical-based models to allow determination of material state without knowledge of initial state and stressor history is an important need; addressing this need will allow more robust interpretations of material condition from engineering scale measurements. The insights gained from the multi-scale computational models described earlier can be used to address the needed inverse analysis to aid in characterizing the microstructural state of the material.

Table III. Candidate property changes for NDE monitoring of nuclear materials (see acronym list)

	FeCr/ Ni-Alloy	Zr-Alloy	ODS	UO₂
Microstructure	Gas bubble, Grain boundary segregation, Dislocation loop, Oxide film.	SSP, ZrH ₂ , Dislocations, Oxide film	Gas bubble, Grain boundary segregation, Dislocation loop, Oxide film; Stability of nano-particles and the impact on microstructure	Gas Bubble, Porosity, Grain coalescence, Re-crystallization
Property Changes	$\Delta H, \Delta V, \Delta DBTT$	$\Delta H, \Delta O_x, \Delta V, \Delta DBTT$	$\Delta H, \Delta O_x, \Delta V, \Delta DBTT$	$\Delta K, \Delta O_x, \Delta V, \Delta DBTT$
ΔM (ΔO_e)	Precipitate growth	Second phase particle behavior		N/A
ΔE (ΔMPa)	Precipitate Growth, Void Growth,	Dislocations, hydrogen		
(ΔMPa)	Oxide film thickness growth	Oxide film thickness growth		
Phenomena/ Degradation Mechanisms	SCC Irradiation Hardening DBTT	Material Loss Irradiation Hardening Hydrogen Pickup SCC DBTT	DBTT SCC	Cracking, Swelling, Porosity Fission Products
Field NDE: Detection SOTA	Cracking: 0.1 – 1 mm (acoustic, eddy current) best case in the field Hardening: Usually ABI or microindentation used DBTT shifts – not measured in field, surveillance specimens used to assess DBTT	Crack: ~0.1 mm (eddy current) Data on fabrication inspection capability difficult to locate. Field inspections (ISI) not done. Literature usually focused on hydriding and other impurity detection – research phase only currently.	N/A (material not used in current LWRs) Literature not very widespread on SCC/DBTT measurement	N/A – field NDE (ISI) methods not used Degradation detection PIE demonstrated in literature for porosity + fission gases. Cracking and swelling not measured in field

NOTES:

- H – magnetic field; V – volume; DBTT – ductile-brittle transition temperature; Ox – oxidation; SCC – stress-corrosion cracking; ISI – in-service inspection; PIE – post irradiation evaluation; ABI – automatic ball indenter

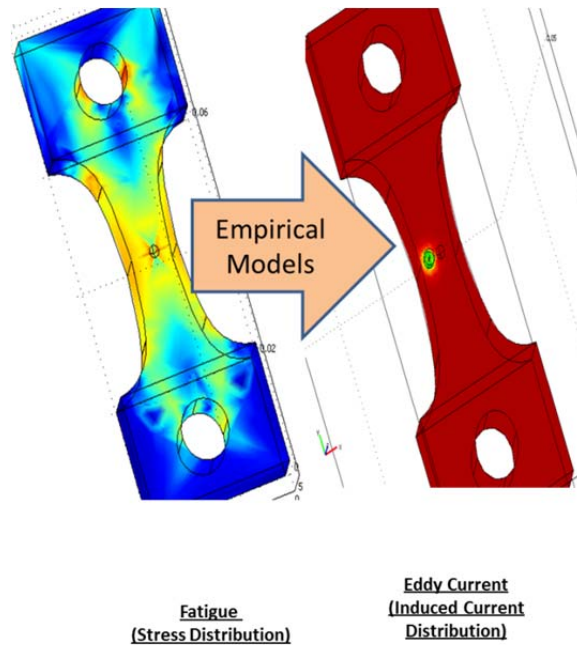


Figure 9. Sensor modeling through empirical correlations between fatigue stress and induced eddy currents

In summary, general detection technology gaps involve 1) measurement signal: robustness of measurement signal to noise, sensitivity to small-scale changes, and selectivity minimally impacted by other factors; 2) sensor technologies: robustness and long-term stability to harsh temperature, radiation, and chemical environments; high sensitivity over large bandwidths; and integrated sensors in structural materials; 3) electronics and instrumentation: robustness to harsh environment, low power or power-harvesting capable; and 4) analytics: multiscale multiphysics computational modeling to guide interpretation of measured data. Barriers to detection do not lie in the physics of the material being inspected; in other words, the relevant phenomenology is not too small. The main issues are the ability to *sample* large enough volumes to find rare features such as defects and crack precursors (impacts the fieldability of the NDE technique), the ability to generate *models of microstructures* with spatially varying electromagnetic and acoustic properties (impacts the development of the NDE technique), and the ability to generate *sensor models* which statistically predict detection based on spatially differing electromagnetic and acoustic properties (impacts interpretation of the measurement). Secondary issues include sensor- or probe-specific limitations regarding radiation hardness, high temperature performance, and frequency limitations.

A summary of the state of the art in in-pile and ex-core measurements is provided below, as a starting point for the gaps assessment.

3.5.1 In-pile measurements

Improved safety and reliability of nuclear fuels are important for the development of Gen IV systems and to the competitiveness of existing nuclear plants. The performance of these fuels must be verified prior to licensing under operating conditions in a fuel cycle test program using test reactors such as Idaho National Laboratory's Advanced Test Reactor (ATR). It has been determined that this fuel cycle test program should include in-pile instrumentation that allows real-time three-dimensional characterization of fuel during irradiation (Rempe et al. 2011). Characterization needs include dimensional changes in fuel pellets, microstructural changes (such as porosity and cracking), as well as characterization of fission gasses. Several recent reviews have focused on in-pile instrumentation for reactor monitoring (Kim et al. 2011; Solstad and Van Nieuwenhove 2011; Villard and Schyns 2011; Rempe et al. 2010). On-line continuous monitoring of in-pile materials requires **radiation-hardened sensor materials**. Gamma and/or neutron testing has been conducted for the production of radiation-hardened magnetic sensors for space applications, fusion reactors, and other nuclear applications.

Some of the key enabling sensor technologies identified for in-pile monitoring are acoustic-based. Typical acoustic transducer elements are composed of **piezoelectric** ceramics, most of which are known to degrade severely at high temperatures and in high neutron flux, resulting in loss of signal. High Curie temperature (T_C) is important for in-pile sensor applications, as predictions for fuel cladding temperatures are $\sim 400^\circ\text{C}$ for light water reactors and $\sim 650^\circ\text{C}$ for sodium fast reactors (Rempe et al. 2011). Radiation-induced thermal spikes in neutron irradiated permanent magnets in the Ohio State University Research Reactor (OSURR) have been estimated to reach 275°C or higher for only 2×10^{13} n/cm² fluence (Liu et al. 2007). For reference, Curie temperatures of conventionally used ferroelectrics for piezoelectric acoustic transducers are $\sim 365^\circ\text{C}$ (PZT5A) (Zhang and Yu 2011), while that of more advanced non-commercial ones can reach only 570°C (PbNb₂O₆) (Holbert et al. 2003). LiNbO₃, while having a high T_C of 1145°C , suffers from oxygen loss and activation of ⁶Li making it less useful than other piezoelectric ferroelectrics. Higher T_C piezoelectrics tend to have lower electromechanical conversion factors. Some non-ferroelectric materials have been considered as very high temperature piezoelectrics, such as AlN and GaPO₄, but these materials typically have very low intrinsic transduction capability (Zhang and Yu 2011).

An alternative technology for acoustic transduction relies on **magnetostrictive** elements to create the acoustic wave in the material to be inspected (e.g., fuel cladding) by NDE. Ultrasonic magnetostrictive transducers have been developed for a number of products, including surgical tools, underwater sonar, and chemical and materials process monitoring. Magnetostrictive devices are numerous and include ultrasonic cleaners, adaptive optics positions, active vibration control, and linear motors. Proposals for in-core diagnostics using magnetostrictive transducers have thus far relied on propagating acoustic signals, such as with guided wave modes, from outside the core, thus shielding the transducer from the highest temperatures and neutron fluxes (Guers 2010). Currently available magnetostrictive acoustic

sensors are limited to low acoustic frequencies (<500 kHz and normally ~10 kHz) due to induced eddy currents in conductive metallic magnetostrictive materials. These low frequencies limit the resolution of acoustic measurements below ~10 μm necessary for inspecting fuel porosity, where acoustic frequencies of ~10 MHz are required. Magnetostrictive transducer materials with high T_C should at least be able to compete with ferroelectric piezoelectric transducers for thermal stability, and may offer superior transduction to current materials considered for high-temperature piezoelectric acoustic transducers.

One important issue constraining in-pile measurement methods is that sensor materials may become altered with radiation such that their transduction function is changed. Experiments are needed to assess real-time transient changes in the sensor materials, such that the change in the sensor itself can be decoupled from the changes of the acoustic signal of the fuel bundle or other component under investigation.

3.5.2 Ex-core measurements

Monitoring of ex-core structural components such as the RPV, on the other hand, allows the introduction of many other types of sensors because the materials used for ex-core components are metallic and therefore allow a wider range of measurement physics. In addition, the environmental conditions are often less severe than in-core conditions. Acoustic, ultrasonic and magnetic sensors are of primary interest when dealing with metallic pressure boundary components. In this case, magnetic sensors measure magnetic field changes near degradation and cracking by means of flux leakage or the Barkhausen effect caused by the discontinuous jump of magnetization due to pinned domain wall movement. These methods have shown great promise in other, less demanding, applications related to inspection of metallic components. Acoustic and ultrasonic measurements could potentially be implemented using a variety of sensors including those based on piezoelectric, magnetostrictive, and electromagnetic acoustic transduction (EMATs). Laser excitation of ultrasound may also be a possible means for inspecting components when access to the surface is available.

In addition to environmental factors, sources of noise due to reactor operation pose a challenge to real time monitoring of structural components. For instance, acoustic noise caused by the flow of reactor coolant has a significant impact on the ability to monitor component degradation via acoustic emission. Electrical noise issues caused by running pumps and motors may cause similar issues for magnetic or electrical based measurements.

Achieving sufficient coverage is another challenge posed by real time monitoring of early degradation in structural components in nuclear reactors. Achieving full component coverage is not practical for the near term and the detection of damage precursors requires detection in the earliest stages of degradation. Hence, optimal placement of sensors for real time monitoring of damage precursors in structural components is a significant challenge to overcome.

4 Magnetic Case Study

The overall goal of expanded capability is to 1) **image** the relevant NDE property at a relevant length scale to the degradation phenomena, 2) **cross-correlate** static computational microstructures with physical and NDE property images acquired with advanced characterization techniques, 3) **predict** microstructural evolution of radiation-induced degradation of a particular type, 4) **design and characterize** radiation-hardened sensors materials, 5) **perform in-situ measurements** of NDE properties in a representative radiation environment, 6) **interpret** NDE measurements with the aid of physics-based models of signal development (sensor model), and 7) **estimate** the remaining life of the component based on interpreted NDE signals and models.

In the case of magnetic sensors, the primary signal is based on magnetization of the sample (e.g. structural steel component such as RPV) as a function of applied current/magnetic field. Magnetization occurs via various processes, among them movement of magnetic domain walls. Domain wall motion is impeded by defects such as dislocations, precipitates, voids, and second phases (Bida 2004). Irradiation produces and/or destroys these types of defects, thus changing the local coercive fields of the material. Analogous to the paragraph above, it is desirable to 1) image defects impeding domain wall motion, 2) compute the effects of defects on domain wall motion and compare these results to bulk measurements like major loop hysteresis and NDE measurements like Barkhausen noise, 3) compute the evolution of these defects as a function of irradiation, 4) produce radiation-hardened Barkhausen noise or other magnetic sensors, 5) test these rad-hard sensors in a representative neutron environment, and 6) show that the measured Barkhausen signals can be interpreted in terms of a nano/micro-scale distribution or density of defects of a certain type, and 7) predict the remaining life of the monitored component based on measured current state and models of defect evolution. A schematic of this process is shown in Figure 10. Some components of this approach have been shown in the literature, as will be discussed below. For the purposes of pedagogical illustration, the following discussion will begin at the sensor and phenomenological level and proceed toward the nanoscale fundamental level.

4.1 *Signal interpretation and sensor modeling*

This section will focus on only one type of magnetic sensor, the Barkhausen noise sensor. Other candidate magnetic NDE sensors are possible, but Barkhausen sensors are sufficiently developed such that much of the physics is well-understood and many steel alloys, including irradiated alloys, have been at least preliminarily investigated. First, the physics of the Barkhausen effect and related magnetic properties are reviewed. Second, the literature on the effects of radiation on magnetic properties of steels is reviewed.

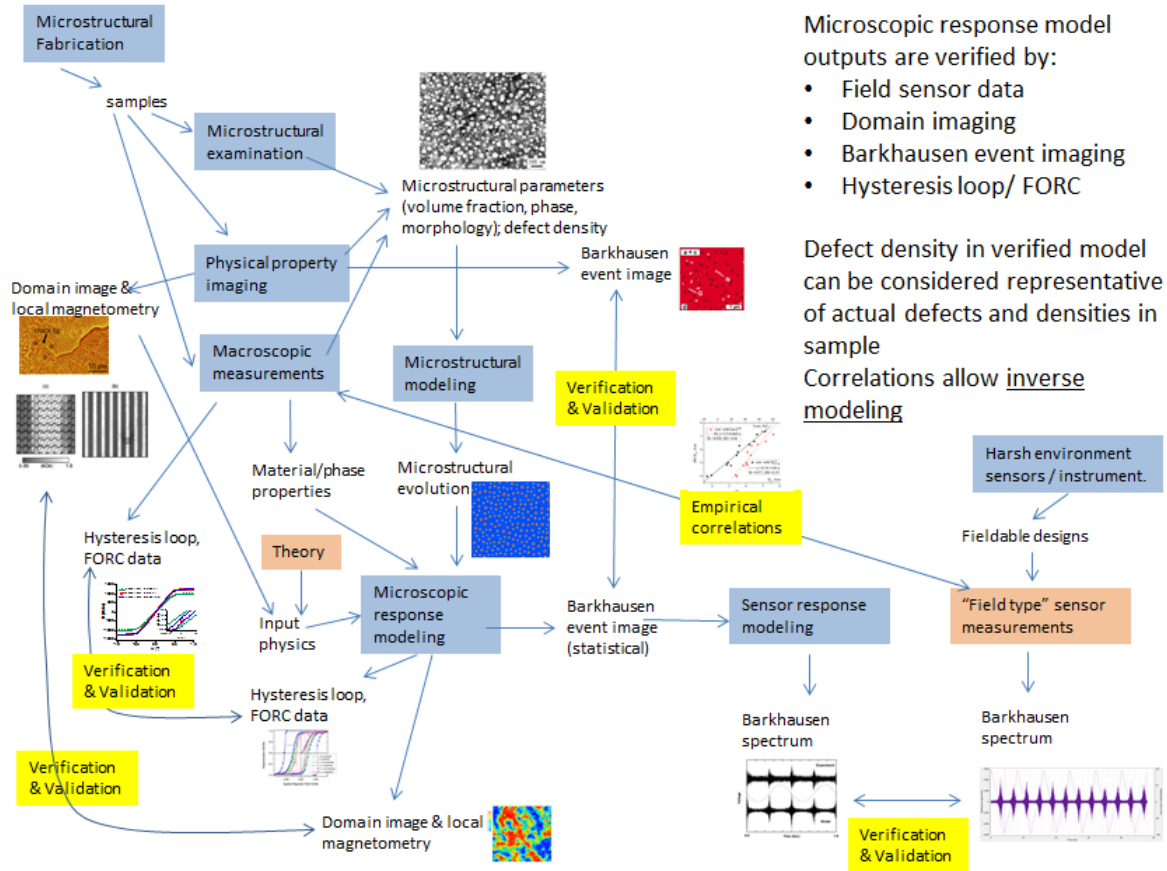


Figure 10. Conceptual relationships among components of mesoscale modeling and measurement for magnetic sensing

4.1.1 Barkhausen noise

Barkhausen noise is a result of the magnetic hysteresis in ferromagnetic materials (Jiles and Atherton 1986; Jiles 2000; Stupakov et al. 2008). The magnetic flux density (B) in ferromagnetic materials placed in an externally applied magnetic field is a function of the applied magnetic field intensity (H) and the magnetic permeability (μ): $B = \mu H$. For ferrous materials, the B - H (or magnetization) curve is sigmoid in shape. Close inspection of the B - H curve reveals that changes in the induced magnetization are not continuous. Instead, they manifest as discrete steps arising from the movement of magnetic domain walls as the applied magnetic field intensity is increased. The magnetic Barkhausen effect is caused by abrupt changes in magnetization with imposed magnetic field, principally due to interactions of 180° domain walls with dislocations (Ranjan et al. 1987b; Ranjan et al. 1987a; Krause et al. 1994). As the imposed magnetic field strength (H) increases, larger numbers of magnetic domains align in the applied field direction. This realignment is, however, not a continuous process since the presence of dislocations, voids, inclusions or other microscopic damage sites results in domain wall pinning.

Increasing the applied field strength results in abrupt realignment of some domains, and is accompanied by a release of energy which may be detected using a sensing coil or piezoelectric sensor. The magnetic signal from a sensing coil is known as magnetic Barkhausen noise (MBN) while the acoustic signal from a piezoelectric sensor is known as magnetic acoustic emission (MAE). MBN signals are detected only from near-surface regions, and are sensitive to 180° domain wall motion which is the largest change in magnetic moment. MAE, on the other hand, originates from the bulk material and is sensitive to 90° domain wall motion which has magnetostrictive strain (Park et al. 1999d; Ranjan et al. 1987a). The number of Barkhausen counts is given by (Ranjan et al. 1987b)

$$N_{MB} = c' \int \rho_{180}(H) \bar{V}_{180}(H) \overline{\Delta B} dH \quad (1)$$

where c' is a constant which depends on the time constant of the pickup coil, permeability, and conductivity of the sample; $\rho_{180}(H)$ is the density of 180° domain walls at field H ; $\bar{V}_{180}(H)$ is the average critical velocity of a 180° domain wall when it is released from pinning sites; and $\overline{\Delta B}$ is the average change in the local magnetic induction due to unit displacement per unit area of domain walls. Note that two Barkhausen bursts are present – one for the positive magnetization and the other for negative magnetization. Numerous models have been developed to predict Barkhausen response to microstructural defects in steels such as grain boundaries and second phase precipitates (Pérez-Benitez et al. 2005; Kameda and Ranjan 1987; Moorthy et al. 1997).

The magnetic Barkhausen method is predominantly a near-surface measurement, with the skin depth, δ , decreasing with increasing frequency defined as $\delta = 1/\sqrt{\pi f \mu \sigma}$, where f is the Barkhausen jump frequency, μ is the magnetic permeability of the material, and σ is the electrical conductivity. For low magnetization frequencies (<10Hz), magnetization is nearly uniform throughout the sample volume. Barkhausen jumps arising from distances greater than the skin depth will be severely attenuated because the pickup coil is positioned on the sample surface. For non-ferritic steel (such as 316L), the skin depth at 1 kHz is about 13.1 mm.

The Barkhausen method has found wide application including as a nondestructive tool for characterizing microstructural changes, tensile properties and fracture toughness in a variety of steels (Ranjan et al. 1987b; Kameda and Ranjan 1987; Moorthy et al. 1997; Wilson et al. 2009; Gurkunov 2000; Dobmann et al. 2001; Raj et al. 2003).

MBN can be used for studying the effects of heat treatments and defects. In some instances, MBN increases due to an increased number of domain walls moving at a given time and/or increased mean free path of movement due to reduction in dislocation density (Kim et al. 1999). The shape and volume fraction of carbide precipitates, for instance, may result in increase or decrease of MBN depending on the thermal treatment, since large carbides create spike domains and decrease MBN, yet larger precipitates, even if they are fewer in number, have increased MBN compared to systems with smaller and fewer pinning sites (Kim et al. 1999). Thus MBN measurements tied to microstructural evidence can provide detailed NDE

information about precipitates in irradiated and heat treated steels. In some cases, multiple peaks in the MBN spectrum indicate domain wall pinning sites of different character, such as grain boundaries and precipitates (Moorthy et al. 1997). It has been demonstrated that the main microstructural influences on MBN is due to nucleation and annihilation of 90° domain walls at grain boundaries and jumps of 180° domain walls pinned by precipitate particles (Pérez-Benitez et al. 2005). Additionally, position and peak of the MBN signal has been correlated to martensite volume fraction in ferrite-martensite steels (Kleber et al. 2004). Peaks due to ferrite versus martensite can also be distinguished in some cases (Vincent et al. 2005).

4.1.2 Other magnetic methods

It is desirable to be able to distinguish different pinning methods of domain walls, in order to assess radiation-induced damage mechanisms non-destructively. Barkhausen techniques have been shown to offer this distinction, provided that the average critical fields for irreversible domain wall motion for different defects are sufficiently separated (Moorthy et al. 1997). Other micromagnetic methods sensitive to domain wall pinning over a wider range of applied magnetic fields are needed to further refine understanding of the reversible and irreversible processes affecting bulk measurements like Barkhausen noise emission. First-order reversal curve (FORC) measurements appear to be a suitable for generating additional information that can be used to probe the details of these domain wall pinning effects.

While it has been frequently shown that mechanical changes in steels, such as those resulting from neutron irradiation, have effects on the major loops of magnetization (e.g., (Park et al. 2003; Park et al. 2004a; Park et al. 1997a)), Barkhausen noise (e.g., (Chang et al. 2002; Altpeter et al. 2001; Park et al. 1999d; Park et al. 1999b; Sipahi et al. 1994)), impedance spectroscopy (Park et al. 2004b; Park et al. 2004a), and ferromagnetic resonance (Park et al. 2004b; Park et al. 2003), only one group has looked at the effects of mechanical changes and precipitation on minor hysteresis loops where the magnetic history is not merely between positive and negative saturations (Takahashi et al. 2006; Kobayashi et al. 2009, 2008; Takahashi et al. 2005). First order reversal curves (FORCs) are similar to minor loops but consist of only part of the loop. Each FORC is obtained by 1) saturating at a positive field, 2) ramping down to a reversal field H_r , and 3) measuring magnetization M at increasing fields H starting from H_r and increasing (see Figure 11).

By choosing a series of reversal points (H_r) with appropriate spacings, a major loop hysteresis curve can be “filled in” by a series of FORC curves which can then be transformed using the FORC function (below) to represent a density function which depends on magnetic switching behavior. As described below, it is typically plotted as a function of microcoercivity (H_c) and bias field (H_b). As explained recently, the FORC function can describe the distribution of pinning strengths for magnetic domain wall motion and hence be a sensitive measure of defect population (McCloy et al. 2012).

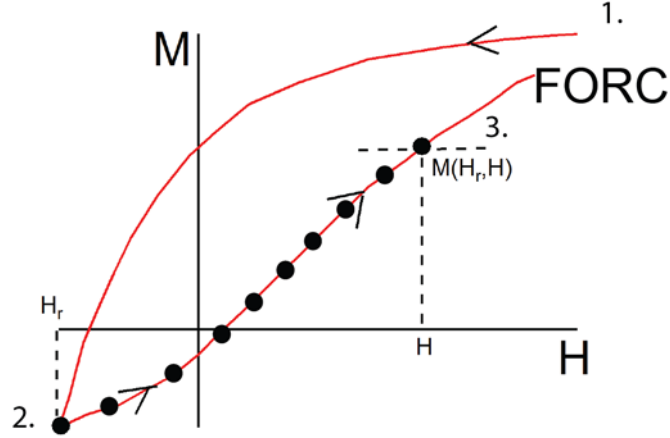


Figure 11. Schematic of First Order Revers Curve (FORC) measurements

The FORC function can be described for FORCs taken increasing after saturation at positive field as (Newell 2005)

$$\rho(H_r, H) = -\frac{1}{2} \frac{\partial^2 M(H_r, H)}{\partial H_r \partial H}, \quad (1)$$

where this function, which can take on positive or negative values, is well-defined for $H > H_r$. In processing the raw data, a smoothing factor is used to create a polynomial surface of the magnetization as a function of the two field parameters (H, H_r) (Pike et al. 1999).

As a first approximation, the FORC distribution can be thought of as a Preisach diagram which deconvolves the hysteresis loop into a series of “hysteron” or square loops each with a “local coercivity” (H_c) (also called a pinning field (Durin et al. 2000)) and a “local bias field” or interaction field (H_b) (Roberts et al. 2000). In practice, many features observed in FORC diagrams do not fit the Preisach model, but it is still useful to conceive of the FORC diagram this way (Stancu et al. 2003a; Newell 2005). In a FORC diagram, H_c is plotted on the x-axis from zero to $H_{c,max}$, and H_b is plotted with both positive and negative values. H_c and H_b are obtained by coordinate transform of the H and H_r data as:

$$H_c = \frac{H - H_r}{2} \text{ and } H_b = \frac{H + H_r}{2}. \quad (2)$$

Any hysteretic phenomenon can be described using a FORC diagram. FORC has been much used for studying the magnetization behavior of geological samples (Chen et al. 2007; Acton et al. 2007), nano-magnets (Pike et al. 2005; Davies et al. 2004; Kou et al. 2011), and spin transitions (Katzgraber et al. 2002; Davies et al. 2005; Katzgraber et al. 2004), but has also been used to examine magnetoresistance (Pomeroy et al. 2009), ferroelectric switching (Stancu et al. 2003b; Fujii et al. 2010), and metal-insulator transitions (Ramírez et al. 2009) by considering a different set of input (e.g., electric field, current) and response (e.g., capacitance, voltage).

It has been shown that FORC diagrams can be modeled as collections of components (Stancu et al. 2003a; Pike et al. 2005; Carvallo et al. 2003), suggesting that the individual contributions of multiple phases could be distinguished. Distinguishing between single domain and multi-domain behavior is straightforward using FORC diagrams (Pike et al. 2001). A Day plot (Roberts et al. 2006; Day et al. 1977) is frequently used in geology to display this data, where M_{rs}/M_r is plotted versus H_{cr}/H_{cv} where M_{rs} is the remnant magnetization after saturation, M_r is the saturation magnetization, H_{cr} is the coercivity of remanence, and H_{cv} is the major loop coercivity (distinguished here from H_c the local coercivity) (McCloy et al. 2012).

4.1.3 Irradiated steels – Barkhausen and other magnetic properties

It is the purpose of this section to review some of the published results on Barkhausen and other magnetic measurements of neutron-irradiated steels.

As far back as 1953 the use of Barkhausen noise and coercive force measurements were used to interrogate steels and other alloys in terms of microstructural state, particularly with regard to composition and strengthening precipitates (Anon 1953; Geisler 1953; Gordon 1953). Early work on the effects of irradiation and radiation damage on magnetism and micromagnetics of materials demonstrated that the majority of materials were sensitive in some degree to the changes in microstructure due to irradiation (Anon 1959), in one case demonstrating a linear response to neutron dose (Alley 1959). More recently, concerns regarding embrittlement of RPV steels encouraged research into NDE methods for mechanical property changes due to neutron irradiation and magnetic measurements were explored with some success. One of the first investigations used MAE and MBN to investigate radiation-induced changes in A302B plate and A533B weld steels (Kwon et al. 1988). The analysis concluded that MAE waveform analysis was a promising NDE technique for monitoring the microscopic changes in the steel components subjected to neutron irradiation. Other magnetic measurements, including remanence, coercivity, and permeability were also used to study radiation damage in these same steels. Magnetic remanence and maximum permeability were found to be highly sensitive to neutron dose (Stubbins et al. 1993). It was hypothesized that atomic disorder due to neutron irradiation caused large-scale magnetic domains to breakdown, while the resulting smaller domains more easily rotated to form closed flux loops, thus, decreasing remanence. The magnetic signals were more sensitive to these microstructural changes than mechanical properties, such as microhardness, in this study (Shong et al. 1993).

Others realized that micromagnetics could be used for NDE of both nuclear pressure vessel steels (Hsu and Chang 1996; Blaszkiewicz 1996; Ara et al. 1996; Dobmann 1994; Kee-Ok et al. 1999; Park et al. 1999a; Park et al. 1999c; Sukegawa et al. 2000; Park et al. 2000; Kikuchi et al. 2006; Park et al. 1999d; Park et al. 1999b) and power plant steels (Ginsztler et al. 1994). Consequently, these steels were soon being routinely investigated using these methods and correlations were developed for hardening and aging. The use of Barkhausen signals for

monitoring stress relief of welds was also pioneered during this time (Rodriguez and Raj 1997), as was thermal recovery monitoring of RPV steels. At the same time, some innovations in these magnetic methods were also being explored, such as the Pulse MAE method (Maeda et al. 1999), which demonstrated signatures increasing monotonically with neutron dose. In another study, a SQUID sensor was used to demonstrate that magnetic coercivity could be used as a signature for radiation hardening with good sensitivity (Otaka et al. 1999). Giant magneto-impedance measurements were also shown to reveal similar signatures (Kim et al. 2002). More general uses for micromagnetics measurements were also emerging, such as monitoring tensile and fatigue strain as a measure of degradation (Sukegawa and Uesaka 2001). Strong correlations are realized between neutron irradiation mechanical property changes and micromagnetic signatures when multiple methods are combined using statistical techniques (Dobmann et al. 2006b; Dobmann et al. 2006a). These methods are also extended to fatigue and tensile damage correlations (Dobmann et al. 2006b; Dobmann et al. 2006a). More recently, these methods have almost become routine for RPV materials in terms of monitoring hardening and recovery (Takaya et al. 2010; Baek et al. 2009; Park et al. 2008; Kikuchi et al. 2007). However, such micromagnetic methods cannot address all such embrittlement problems, such as phosphorus segregation when no hardening is involved (Debarberis et al. 2007).

The trends with regards to hardening and neutron irradiation effects are clear, however, and magnetic methods that are sensitive to domain wall pinning are suitable for NDE of magnetic materials. Regardless of the alloy, magnetic Barkhausen noise is almost always decreased with irradiation (Sipahi et al. 1994; Park et al. 1997b) due to domain wall pinning, and typically coercivity and permeability are increased, though not always (Park et al. 1997a; Park et al. 1999b). Irradiation processes that produce point defects and small defect clusters can be effectively probed as a function of dose for small doses using these methods and linear responses are found initially. Embrittlement thresholds can be determined using statistical regression methods and calibrated samples. However, these methods cannot distinguish between hardening events and researchers often suggest that additional work is required to sort out complex effects.

4.2 Magnetic sensor materials and in-situ measurements

Compared with piezoelectric materials, very little is known about the effects of radiation, particularly neutron irradiation, on the properties of magnetic sensor materials. Neutron damage studies of magnetic sensor materials are few and far between (see Appendix for an incomplete summary), and results indicate a wide range of behaviors, including amorphization (Chukalkin et al. 1983), cation site mixing (Chukalkin et al. 1975), superparamagnetism (Parkhomenko et al. 1976), anisotropy change (Chukalkin et al. 1981), loss of remnance (Anderson et al. 2005), or complete loss of magnetization due to radiation-induced thermal spikes (Liu et al. 2007). In addition to their use in magnetostrictive transducers, magnetic materials are integral to the operation of linear variable differential transformers (LVDTs), which will also find future use as in-pile diagnostic instruments (Rempe et al. 2011).

Some magnetic sensor types have been shown to be relatively rad-hard, including Hall (Bolshakova 1999; Bolshakova et al. 2009), SQUID (Nagaishi et al. 2003), magnetoresistive (anisotropic magnetoresistance, AMR (Díaz-Michelena 2009; Sanz et al. 2012); giant magnetoresistance, GMR (Carroll 2010); and magnetic tunnel effect magnetoresistance, TMR (Persson et al. 2011) type), coils (Nemoto et al. 2011), and fluxgate (Nielsen et al. 1997) sensors. These effects are, of course, temperature, dose, dose rate, flux, and radiation-type dependent.

Based on existing research on the characteristics of radiation-tolerant ceramics (Sickafus et al. 2007; Sickafus et al. 2000), it may be possible to find magnetic oxides or other magnetic ceramic materials which possess the tolerance for lattice disorder necessary for radiation-tolerance, in addition to good high temperature properties and good magnetostriction. Since one limitation of current magnetostrictive materials (see 3.5.1) is high electrical conductivity and hence large eddy currents, low conductivity oxide materials are candidates for this application.

A list of magnetostrictive materials is shown in Table IV. The materials currently considering as “conventional” magnetostrictive materials include Ni, Co, CoFe_2O_4 , $\text{Fe}_{81}\text{Si}_{13.5}\text{B}_{13.5}\text{C}_2$ (Metglas), and $\text{Tb}_{0.3}\text{Dy}_{0.7}\text{Fe}_{1.92}$ (Terfenol-D). Terfenol-D is a giant magnetostrictive material used in sonar. A related Fe-Ga alloy, known as Galfenol, having somewhat smaller magnetostriction but much higher Curie temperature, is under current development. Next-generation, rad-hard magnetostrictive materials will likely come from the class of materials that is very tolerant of disorder, or is already disordered itself. For example, bulk metallic glasses (BMG) such as FeDyBSiNb (Li 2011) as well as commercial “Metglas” Fe-Si-B based amorphous metals have been shown to have reasonable magnetostriction coefficients. Usually these materials are used for mechanical applications or as soft magnets, and optimization for magnetostriction and radiation damage is possible.

Some work has been conducted on radiation damage of Metglas alloys. Reduction of permeability with neutron irradiation is much slower than with conventional crystalline magnetic materials such as Permalloy (Brown et al. 1984). Decrease in permeability for Metglas appears to be a result of a short-range ordering due to defect accumulation at low fluences and pinning of domain walls at higher fluences (Brown et al. 1984). Neutron irradiation decreases permeability at low frequencies due to impeding of domain wall motion, but increases permeability at higher frequencies due to increased importance of rotational magnetization (Kim et al. 2000). Metglas alloys crystallized to a nanograin structure are known as FINEMET (Škorvánek and Gerling 1994). These have an order of magnitude lower magnetostriction than Metglas. It has been shown that neutron irradiation results in residual stress centers which react with domain walls through magnetostriction. Thus FINEMET alloys have lower sensitivity to radiation damage than do the chemically similar Metglas alloys due to magnetostriction, even when neutron irradiation partially amorphizes the FINEMET alloys (Škorvánek and Gerling 1994). In a similar study using pulsed laser irradiation, it was shown that laser-induced crystallization of amorphous alloys was strongest in materials with larger magnetostriction constants (Sorescu 1998). Ion irradiation studies have shown the importance of in-plane versus out-of-plane anisotropy in

determined the magnetic properties of amorphous magnetic alloy films (McCord et al. 2005; Kopcewicz and Dunlop 2001). On the other hand, neutron irradiation can restore ductility in thermally-embrittled amorphous Metglas, while nanocrystalline FINEMET do not undergo such radiation-induced annealing (Škorvánek and Gerling 1994). Thus it is not immediately clear whether amorphous magnetostrictive alloys necessarily be advantageous for radiation hardened sensors, as there appear to be some disadvantages associated with high magnetostriction. Recent work studying electron-induced crystallization has shown that stability against crystallization can be assessed from standard thermal tests (Nagase et al. 2012). Driving force for defect-induced crystallization in amorphous metallic alloys is based on free energy considerations, local atomic structure due to radiation-induced defects, and radiation enhanced diffusion (Nagase et al. 2012). There appears to be more room for material tailoring in these systems, and thus more work is warranted in this area.

Crystalline metallic systems known as “high entropy alloys,” (Zhang et al. 2008) have been recently identified as candidates for radiation-tolerant components due to their potential for “self-healing” through recrystallization via a glassy state. The magnetostriction properties of these alloys (e.g., $Al_xCoCrFeNi$) are virtually unexplored, and at least some compositions are suggested to have Curie temperatures as high as ~ 600 °C (Chou et al. 2009). Additionally, some room-temperature ferromagnetic oxides in the double perovskite family have been shown to be very tolerant to disorder in their cation sites (Serrate et al. 2007), and may be additional candidates for rad-hard magnetic materials if materials with high enough T_c can be found. The saturation magnetostriction and Curie temperatures from the literature for some of these materials are listed in Table IV. Galfenol, high entropy alloys, and double perovskites are not commercially available materials.

Table IV: Magnetostrictive materials

Material	Alternative name	Magnetostriction (ppm)	T_c (°C)	Resistivity (Ω -m)	Ref.
$Fe_{81}Si_{3.5}B_{13.5}C_2$	Metglas 2605SC	+20	350	5.0E-3	6
Ni		-33	358	7.1E-8	1
Co		-62	1130	5.6E-8	1
49Co-49Fe-2V	Rem/perm-endur	+70	932	4.2E-7	1,7,12
$CoFe_2O_4$		-110	520	1.0E+2	1,10
$Tb_{0.3}Dy_{0.7}Fe_{1.92}$	Terfenol-D	+1640 to +2000	380	6.0E-7	2,4,11
$Fe_{81.6}Ga_{18.4}$	Galfenol	+400	670	8.0E-3	3,9
$Al_2CoCrFeNi$	High entropy alloy	Unknown	596	1.0E-2	5
Ca_2FeReO_6	Double perovskite	~ 200	250	2.0E+0	8

1. (Lide 2006)

2. (Zhang and Chen 2005)

3. (Staley and Flatau 2005)

4. (Moffett et al. 1991)

5. (Chou et al. 2009)

6. <http://www.metglas.com>

7. (Pinnel and Bennett 1973)

8. (Serrate et al. 2005)

9. (Clark et al. 2002)

10. (Özgür et al. 2009)

11. <http://www.etrema-usa.com/core/terfenold/>

12. <http://www.hightempmetals.com/techdata/hitempPermendurdata.php>

4.3 Characterization of magnetic microstructure and nanostructure

Characterization of magnetic materials can generally be broken up into seven categories. It is beyond the scope of this report to review these techniques in detail, but the reader is referred to a recent review (Lau and Shaw 2011). Characteristics of some of the following potential methods are summarized in Table V.

- 1) X-ray methods include x-ray transmission microscopy plus x-ray transmission microscopy + x-ray circular dichroism (XMCD); and photoelectron emission microscopy (PEEM) + XMCD which can provide time-resolved imaging
- 2) Neutron methods such as neutron tomography, which provides 3-D imaging
- 3) Electron microscopy + polarization analysis includes scanning electron microscopy polarization analysis (SEMPA); spin-polarized scanning tunneling microscopy (SP-STM); ballistic electron magnetic microscopy (BEMM), which uses a variation of STM for thin films on semiconductor substrates; and Lorentz mode transmission electron microscopy (TEM)
- 4) Orthogonal fluxgate method (wires)
- 5) Magnetic particles/ Bitter patterns
- 6) Magnetic force microscopy (MFM), scanning Hall probe microscopy (SHPM), superconducting quantum interference device (SQUID) microscopy, and similar techniques
- 7) Magneto-optical techniques, including magneto-optical Kerr effect (MOKE) microscopy; Faraday effect (magnetic circular birefringence) on transparent materials; and magneto-optic versions of conventional optical microscopy, including confocal optical microscopy or scanning near-field optical microscopy (SNOM).

In comparing these methods, it is important to consider the components of magnetization of the sample under test for which the particular technique is sensitive. For example, some methods are sensitive to the magnetization of the whole sample (Lorentz TEM, Faraday effect). Other methods are only sensitive to surface magnetization (MOKE microscopy, SEMPA). Still others, particularly scanning methods, are sensitive to stray fields at the surface (MFM, Scanning SQUID, SHPM). Domain walls at surfaces tend to exert fields of ~ 10 mT and have a spatial extent of 20 – 50 nm. Finally, some techniques are sensitive to individual atomic moments, such as synchrotron techniques of XMCD, which can assess site occupancy the ferromagnetic atoms Fe, Co, and Ni.

Given the inherent spatial resolution and instrumental constraints, a subset of these methods was down-selected for further consideration, as shown in Table VI. Applicable materials that can be analyzed, details on the physical property measured, and bulk measurement equivalencies are listed. These bulk measurements are listed as a means of considering the correlation with bulk NDE methods. Table VII further assesses these techniques, examining pros and cons of each technique, applicable temperature ranges, and equipment manufacturers, if applicable.

The bottom line of this exercise was that the magnetic imaging techniques of most interest for near-term work on steels are magnetic force microscopy (MFM), magneto-optical Kerr effect (MOKE), and possibly scanning Hall probe microscopy (SHPM). Of the three, MFM gives the best spatial resolution, ~25 nm, sufficient for imaging domain walls. The disadvantage is that only a small area can be scanned at any one time. MOKE, on the other hand, has much lower resolution, ~300 – 500 nm, as it depends on visible light; however, MOKE can image large areas simultaneously and can thus illustrate grain texturing and magnetic domain structure. SHPM has similar resolution to MOKE, ~300 nm, but in theory is more reliable than MFM for true magnetization measurements since it measures stray fields directly, not through interaction of a tip with the sample fields as does MFM.

4.3.1 Magnetic imaging as input to microstructure simulations

An additional advantage of the down-selected techniques is that point-by-point local hysteresis loops can be generated, showing differences in local magnetization behavior. This has the benefit of allowing direct comparison with meso-scale phase field computer models, which compute “voxel-by-voxel” magnetization. Magnetometry (or susceptometry) is obtainable from both MFM (Babcock et al. 1996; Zhu et al. 2003; Sorop et al. 2003) and MOKE microscopy (Choe et al. 2002; Sato and Ishibashi 2008) data. MOKE microscopy/ magnetometry is routinely used for combinatorial screening of magnetic compounds (Zhao et al. 2004).

MFM and MOKE have also been used to elicit detailed information about the mechanisms of domain wall dynamics. Comparison of subsequent MFM images obtained as a function of magnetic field has been used to show the difference between Barkhausen jump volumes grown by domain wall motion or nucleated in isolated regions (Schwarz et al. 2004). A similar experiment using a MOKE microscope known as a magneto-optical microscope magnetometer (MOMM) has been used to plot probabilities of Barkhausen jump sizes as a function of NiO thickness on Fe, showing that mechanisms of domain wall creep can be distinguished from domain wall pinning and depinning (Barkhausen) (Lee et al. 2011). Finally, MOKE magnetometry has been used to study the effects of AC magnetic fields on domain walls, allowing the differentiation between domain boundary creep, slide, and switching as a function of AC field excitation frequency, and allowing the accurate measurement of depinning fields and domain wall mobility (Kataja and van Dijken 2011). In summary, magnetic imaging methods such as MFM and MOKE provide a critical experimental tool to obtain numerical parameters (such as depinning fields, domain wall mobilities, etc.) for domain wall dynamics, as affected by defects, which can be put into microstructural models of magnetic materials to study the evolution of magnetic signatures as a function of changes of defect concentration and magnetic field.

Table V. Summary of some magnetic imaging techniques

Category	Technique	Spatial resolution	Time resolution	Notes
Electron	Lorentz TEM	100 nm / 1 nm	ms to s, except pump-probe	Fresnel mode, For (u)ncorrected/ (c)orrected aberrations
Electron	Differential phase contrast	20 nm		STEM mode, uncorrected aberrations
Electron	Electron holography	0.5 nm / 0.1 nm		For (u)ncorrected/ (c)orrected aberrations
Electron	SEM with polarized analysis (SEMPA)	20 nm		
Electron	Spin-polarized low-energy electron microscopy (SP-LEEM)	10 nm		Good for in-situ spin reorientation (like LEED)
Electron/ X-ray	Photo-emission electron microscopy (PEEM)	50 nm	100 ps	Circularly polarized x-ray incident on sample
X-ray	X-ray magnetic circular dichroism (XMCD)		> 70 ps for synchrotron; poss. better for other	Generate with synchrotron, free-electron laser, or high-harmonic generation laser; resolution in principle much better than M-O due to smaller wavelength of EUV/ soft x-ray
Magneto-optical	Magneto-optical Kerr effect (MOKE)	300-500 nm		Visible laser, spot size few μm very easy; smaller spots diffraction limited
Magneto-optical	Time-resolved MOKE (TR-MOKE)	~ 50 nm	< 1 ps	pump-probe; Fourier transform time-resolved data to get frequency-resolved
Scanning probe	Magnetic force microscopy (MFM)	25 nm		Mostly limited to perpendicular field with normal tip; use in vacuum improves SNR 2-3 orders magnitude
Scanning probe	Ballistic electron magnetic microscopy (BEMM)			Variety of STM
Scanning probe	Ferromagnetic resonance force microscopy (FMRFM)	\sim nm		Resolution determined by FMR linewidth and field gradient, not tip radius
	Micro-Brillouin light scattering (BLS) + AFM near field imaging	~ 85 nm		

Table VI. Down-select list of magnetic imaging methods relevant for NDE

Technique	Spatial resolution	Applicable materials	Physical parameter measured	Bulk property analogue
X-ray magnetic circular dichroism (XMCD) microscopy	~10 nm	Those containing Fe, Co, or Ni	Difference spectrum of x-ray absorption taken in magnetic field with right-hand and left-hand circularly polarized light (x-ray, 700-900 eV); can look at ferromagnetic atoms Fe, Co, Ni; tune x-ray energy for atoms of interest	Magnetic hysteresis loops, saturation, coercivity, permeability/susceptibility; can get element specific hysteresis data
Magneto-optical Kerr effect (MOKE)	300-500 nm	ferro-, ferri-, and anti-ferromagnetic materials	Reflection of polarized light (UV, VIS, or IR) by material sample subjected to magnetic field (rotation, ellipticity, and intensity change)	Magnetic hysteresis loops, saturation, coercivity, permeability/susceptibility
Magnetic force microscopy (MFM)	25 nm	Ferritic steel, martensitic steel, deformed austenitic steel (304, 316), other ferro-, ferri-, and anti-ferromagnetic materials, superconductors	Magnetic force gradient by shift in phase or resonant frequency of oscillating cantilever probe; measures interaction between sample and tip	Magnetic hysteresis loops, saturation, coercivity, permeability/susceptibility
Scanning Hall Probe Microscopy (SHPM)	300 nm	ferro-, ferri-, and anti-ferromagnetic materials, superconductors	Hall voltage in semiconductor sensor induced by local magnetic field	Magnetic hysteresis loops, saturation, coercivity, permeability/susceptibility; Provides true quantitative assessment of local magnetic field of sample
Scanning SQUID microscope (SSM)	~2.5 μm	ferro-, ferri-, and anti-ferromagnetic materials, superconductors	Sample stray magnetic field effect on the Josephson junction superconductor circuit	Magnetic hysteresis loops, saturation, coercivity, permeability/susceptibility; Provides true quantitative assessment of local magnetic field of sample

Table VII. Down-select trade study for magnetic imaging methods

Technique	Pros	Cons	Temperature ranges	Notes	Manufacturers
X-ray magnetic circular dichroism (XMCD) microscopy	Can obtain individual atomic moments; Get chemical environment “for free” with XAS; some leverage with chemical imaging initiative; some in-house expertise (C. Pearce); can get tomography	Currently need synchrotron; in-house capability probably 5 years out	LT/RT/HT?	May eventually be feasible with compact x-ray light source	N/A Everything is custom
Magneto-optical Kerr effect (MOKE) - optical	In-house expertise in optical microscopy	Only near surface information; resolution not great	LT/RT/HT	Integrates over visible wavelengths	Evico
Magneto-optical Kerr effect (MOKE) - scanning	In-house expertise in optics			Use low-cost lasers and optical table	Durham Magneto Optics (through Quantum Design)
Magnetic force microscopy (MFM) – PPMS based	Takes advantage of existing equipment for low T and high field measurements; can get Hall probe sensor as well for low incremental cost	Only near surface information; data relies on interaction between tip and sample; only sees domain walls; applied fields influence tip and sample; tip magnetization may affect sample	LT/RT only	Can get various sensor heads to upgrade probe	Nanomagnetics (through Quantum Design) Attocube
Magnetic force microscopy (MFM) – stand alone	Easy to perform low and high T measurements in modest fields		LT?/RT/HT		Nanomagnetics instruments, Asylum
Scanning SQUID microscope (SSM)	Can detect very small transition metal impurities in insulators; the most sensitive of the instruments	Overkill in sensitivity for steels	LT/RT/HT?	Could get widespread use in FCSD to characterize substrates	Neocera Magma C20

NOTES: LT – low temperature, RT – room temperature, HT – high temperature

4.3.2 Existing literature on magnetic imaging of steels

It is worthwhile to review a few relevant studies on defects in steels using magnetic imaging methods, in order to assess what is possible from an experimental point of view for studying defects in steel.

Stress-corrosion cracking susceptibility of 304 stainless steel was evaluated by MFM (Takaya et al. 2004). It was shown that transformation from austenite to martensite due to chromium depletion along grain boundaries caused the magnetic phase seen with MFM and corroborated with phase identification by electron backscatter diffraction (EBSD). Strain-induced martensite forming at crack tips of 304 and 310S austenitic steels has been imaged using MFM (Zhang et al. 2010). Fatigue crack tips have also been imaged in bearing steel using SHPM since the magnetic flux density is different at the crack tip (Kida et al. 2012).

A related force microscopy technique has been developed, in which it is the electrical conductivity which is measured by means of eddy currents (Nalladega 2009). So-called eddy current force microscopy (ECFM) allows the simultaneous probing of skin depth changes due to changes in conductivity and magnetic permeability.

Perhaps most relevant is the very recent work from the German group (Batista et al. 2012) showing the changes in structure and phase of austenitic twinning induced plasticity (TWIP) steels as a function of tensile stress. Additionally, these authors show the evolution of the domain structures of cementite particles in Fe-0.8%C unalloyed steel using MFM, and correlate their orientations with EBSD. These methods could be used to understand the growth habits of radiation-induced precipitates (Maziasz 1989) and their behavior under applied magnetic fields such as used in NDE methods.

4.4 *Phase field modeling of the effect of microstructures on magnetic response*

Irradiation causes microstructure evolution such as grain boundary solute segregation, precipitation, void formation, dislocation networking, and grain growth, hence, the material property degradation. To monitor the material property degradation, we need a fundamental understanding of the effect of microstructures on material responses, i.e., signal physics. In this case study, we present a phase-field model for simulating the effect of nonmagnetic particles/precipitates on magnetic domain nucleation and magnetic response. The simulations are to demonstrate the method capability in interpreting NDE measurements and understanding signal physics.

FeCr alloy is taken as the model system in the case study. The nonmagnetic particles/precipitates could be Cr rich precipitates which are formed during irradiation-induced phase separation in FeCr alloys. The magnetization vector $\mathbf{M}(\mathbf{r}, t)$ is used to describe the spatial and time distribution of magnetic domains. The magnetic domain evolution is simulated by solving the Landau–Lifshitz–Gilbert (LLG) equation, (Hubert and Schaefer 1998)

$$(1 + \alpha^2) \frac{\partial \mathbf{M}}{\partial t} = -\gamma_0 \mathbf{M} \times \mathbf{H}_{eff} - \frac{\gamma_0 \alpha}{M_s} \mathbf{M} \times (\mathbf{M} \times \mathbf{H}_{eff}), \quad (3)$$

where $\mathbf{M}(\mathbf{r}, t)$ is the magnetization vector, M_s is the saturation magnetization of the magnetic phase, γ_0 is the gyromagnetic ratio, α is the dimensionless damping constant, and \mathbf{H}_{eff} is the effective magnetic field, which can be represented as a variational derivative of the total free energy of the system with respect to magnetization, $\mathbf{H}_{eff} = -\frac{1}{\mu_0} \frac{\delta F}{\delta \mathbf{M}}$,

where μ_0 is the permeability of vacuum and F is the total energy including magnetocrystalline anisotropy energy, exchange energy, magnetostatic energy, external field energy, and elastic energy [see Hu et al. “360° Bloch domain wall stability and its interaction with nonmagnetic particles in Fe alloys: computer simulation of magnetic hardening” submitted to *IEEE Magnetics Letters* (2012)].

We simulated magnetization switching in a single crystal of iron with a simulation cell of $64l_0 \times 128l_0 \times 64l_0$. l_0 is a real spatial length that should be smaller than the exchange length $l_{ex} = \sqrt{2A_c / (\mu_0 M_s^2)} = 2.86 \text{ nm}$ (Rave et al. 1998; Bertotti 2008). Periodic boundary conditions were applied in the three directions. The material constants of body-centered cubic α -phase single crystal iron are listed in

Table VIII. The damping constant is usually a small value of $\alpha \leq 0.1$. In our simulation $\alpha = 0.05$. In addition, the magnetostrictive constants are relatively small so the elastic interaction energy is ignored.

Error! Reference source not found.-13 summarize the predictive magnetic hysteresis loops and magnetic domain evolution in a single crystal of iron with or without nonmagnetic particles, respectively. The nonmagnetic particle is described by an order parameter field η . $\eta=1$ represents the magnetic phase while $\eta=0$ for the nonmagnetic phase. η varies smoothly from 1 to 0 across the interface. Therefore, $\mathbf{M}(\mathbf{r}, t) = M_s \mathbf{M}^*(\mathbf{r}, t) = M_s \eta(\mathbf{r}) \mathbf{m}(\mathbf{r}, t)$, where $\mathbf{m}(\mathbf{r}, t)$ is a unit vector, and $\eta(\mathbf{r})$ does not evolve with time. The applied external magnetic field $\mathbf{H}_{ex} = (H_{ex1}, 0, 0)$ varies from $-40 \text{ kAm}^{-1} \sim +40 \text{ kAm}^{-1}$ with an increment or decrement of $\Delta H_{ex1} = 0.1 \text{ kAm}^{-1}$ for each 26 ps. The simulation time step is $\Delta t = 0.1 t_0 = 0.26 \text{ ps}$ where $t_0 = (1 + \alpha^2) (\gamma_0 M_s)^{-1} = 2.6 \text{ ps}$.

For a perfect single crystal, i.e., without any nonmagnetic particle, the hysteresis loop started from point (A) where multiple domains coexist as shown in Fig. 11(b)(A). When the applied magnetic field H_{ex1} reached its maximum, a 360° domain wall was formed and remained (see Fig. 11(b)(B)). When unloading (i.e., decreasing the field), the 360° domain wall thickened as shown by Fig. 11(b)(C). With the applied magnetic field H_{ex1} switched its direction, an anti-direction domain of $(-M_s, 0, 0)$ nucleated and grew by splitting the 360° domain wall into two 180° domain walls. From Fig. 11(b)(B) to Fig. 11(b)(H), it is seen that a process of emerging and splitting of a 360° domain wall occurred repeatedly during the domain switching. Consequently, the corresponding hysteresis loop is repeatable. Both the 360° domain wall and 180° domain wall

are Bloch domain walls as illustrated by Fig. 11(c). The result demonstrates that the 360° domain wall can be stable with the applied magnetic field and is a nucleation site of an anti-direction domain during domain switching.

Table VIII. Material properties of iron (Fe) used in the simulations.

Symbol	Name	Units	Value	Ref.
α	damping constant	Unitless	0.05	*
M_S	saturation magnetization	Am^{-1}	1.71×10^6	(Craik 1995)
$\mu_0 M_S$		$\text{NA}^{-1} \text{m}^{-1}$	2.16	(d'Aquino 2004)
K_1	anisotropy constants	J/m^3	4.80×10^4	(Cullity and Graham 2009)
K_2	anisotropy constants	J/m^3	5.0×10^3	(Cullity and Graham 2009)
A_c	exchange stiffness	J/m	1.5×10^{-11}	(Bertotti 1998)
λ_{100}	magnetostrictive coefficient	Unitless	2.10×10^{-5}	(Cullity and Graham 2009)
λ_{111}	magnetostrictive coefficient	Unitless	1.57×10^{-5}	(Cullity and Graham 2009)
$(\gamma_0 M_S)^{-1}$		ps	2.6	(d'Aquino 2004)
c_{11}	elastic constant	GPa	230	(Adams et al. 2006)
c_{12}	elastic constant	GPa	134	(Adams et al. 2006)
c_{44}	elastic constant	GPa	116	(Adams et al. 2006)

*this value of α was used in simulations of (d'Aquino 2004). Much larger values have been used in other phase field models (Zhang and Chen 2005), but their physical basis is questionable (Chun et al. 2012). Assuming the spectroscopic splitting factor (g) of Fe is 2.094 (Craik 1995), the gyromagnetic ratio γ_0 is 1.84×10^7 rad/s/Oe or 29.316 GHz/T since $\gamma = -g^*e/(2m_e)$ in SI units, where e is the charge and m_e is the mass of the electron. The damping parameter (α) is then calculated from the measured ferromagnetic resonance linewidth (ΔH) and the resonant frequency f_0 for Fe in (Frait and Heinrich 1964) as $\alpha = \gamma_0 \Delta H / (2\omega_0)$, where $\omega_0 = 2\pi f_0 = \gamma_0 H_0$ and H_0 is the internal field, sometimes taken as the anisotropy field. The result of this analysis yields $\alpha = 0.005$, which is similar to values for magnetic garnets and other microwave materials (Chun et al. 2012). In the phase field modeling work, it was shown that $\alpha = 0.05$ and $\alpha = 0.005$ resulted in the same loop, but that $\alpha = 0.005$ required a smaller time step. For ease of computation, $\alpha = 0.05$ was used in the simulations.

The effect of nonmagnetic particles on the stability of 360° domain wall and magnetic response can be seen from Figs. 12-13 where only particles with spherical shape were considered for simplicity. The particle size is counted by its volume fraction with respect to the simulation cell. It is seen from Fig. 12(a) that a small nonmagnetic particle (e.g., $V_f=0.02336\%$) does not bring any influence on the magnetic hysteresis loop comparing with the case without any particle. The domain morphologies shown in Fig. 12(b) also illustrated the same conclusion. There is almost no difference between the domain morphologies in Fig. 11(b) with $V_f=0$ and Fig. 12(b) with $V_f=0.02336\%$. However, when the particle size increases, its impact on the magnetic hysteresis loop and domain morphologies is significant and can be seen clearly from Figs. 12 and 13. The impact is reflected in two aspects. On one hand, the presence of the nonmagnetic particle prevents the formation or reformation of the 360° domain wall as shown by Fig. 12(b) (D') and (E) with $V_f=0.10285\%$. On the other hand, the particle location becomes the nucleation site of any new anti-direction domain as demonstrated by Fig. 13(b) (N1) and (N2). In such cases, the larger the particle, the easier the nucleation and the smaller the coercive field. The coercive field is the applied field value at $M_1=0$. However, the corresponding coercive field cannot be smaller than the case with a 360° domain wall as displayed by Fig. 13(c). When the particle is too small (e.g., $V_f=0.10285\%$), it is very difficult to switch back as depicted by Fig. 12(a).

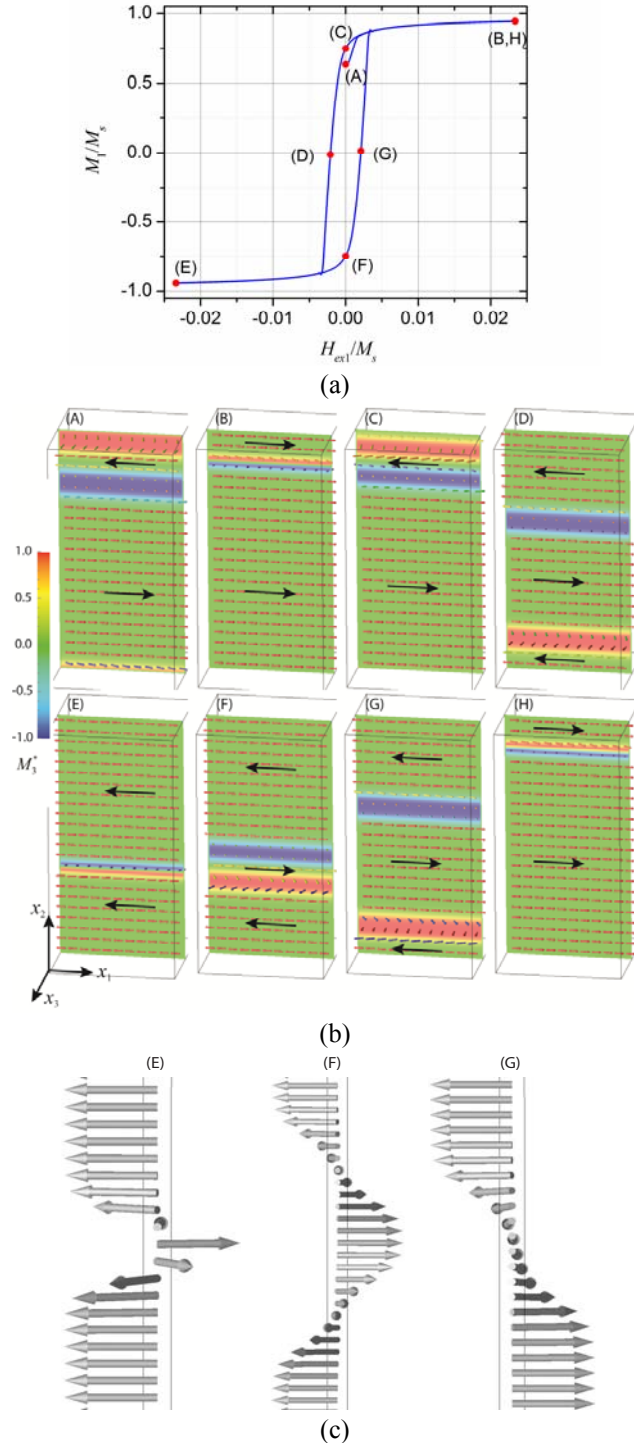
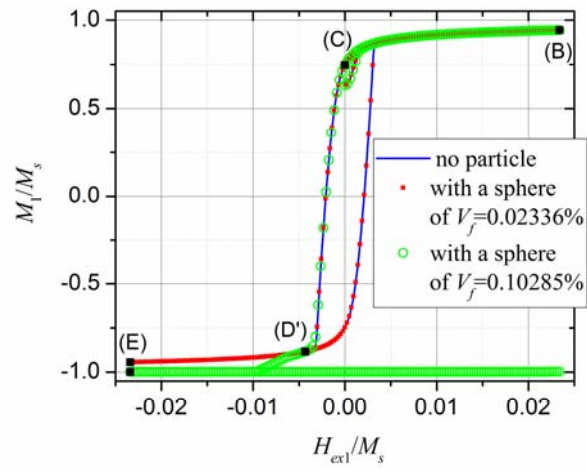
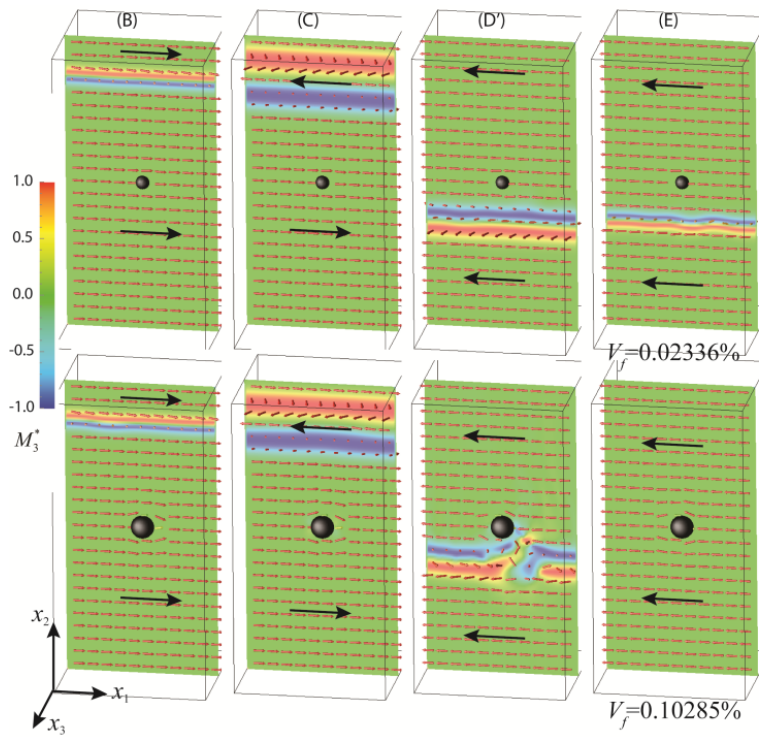


Figure 11. (a) Magnetic hysteresis loop of a single crystal iron; (b) Magnetic domain morphologies of the central x_1x_2 plane at the labeled points of (a). The outside box is the 3D simulation cell, the black arrows are the magnified magnetization vectors shown in the corresponding domains, and the color illustrates the magnitude of the magnetization component ($M_3^* = M_3 / M_s$) as depicted by the color bar (e.g., blue is into the page and red is out of the page). (c) Arrowed Bloch domain walls at the given points (E,F,G). (E) represents a 360° wall, (F) shows two split 180° walls, and (G) shows a single 180° wall.



(a)



(b)

Figure 12. (a) Magnetic hysteresis loops for the cases of $V_f=0.0$ (no particle), $V_f=0.02336\%$ and $V_f=0.10285\%$, respectively. V_f is the volume fraction of the nonmagnetic sphere with respect to the simulation cell. (b) Magnetic domain morphology snapshots at the corresponding points of the hysteresis loops for both cases with nonmagnetic particle (black sphere).

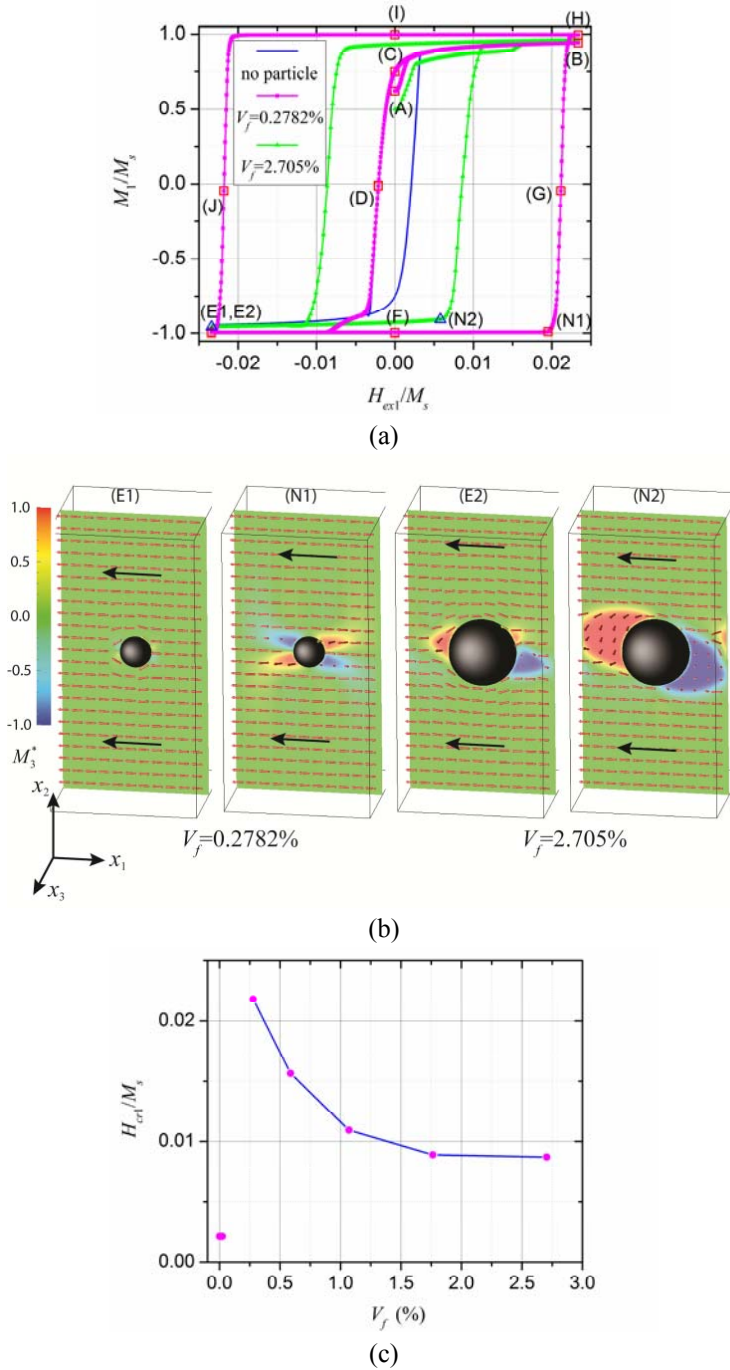


Figure 13. (a) Magnetic hysteresis loops for the cases of $V_f=0.0$ (no particle), $V_f=0.2782\%$ and $V_f=2.705\%$, respectively. V_f is the volume fraction of the nonmagnetic sphere with respect to the simulation cell. For the case of $V_f=0.2782\%$, the loop follows the sequence of (ABCDEFN1GHIJ). (b) Magnetic domain morphology snapshots at the indicated points (E1,N1; E2,N2) of the hysteresis loops. The points (E1,E2) are corresponding to the maximum applied external magnetic field. The points (N1,N2) are corresponding to the nucleation sites of new anti-direction domains. (c) Summary of the coercive field versus the volume fraction of a nonmagnetic sphere.

In summary, we have developed a phase-field model to study the interaction between magnetic domain walls and nonmagnetic particles based on the Landau–Lifshitz–Gilbert equation. From simulations, we found that both nonmagnetic particles and 360° domain walls are the nucleation sites of an anti-direction domain during domain switching. But an anti-direction domain nucleates much easier at a 360° domain wall through splitting the 360° domain wall into two 180° domain walls, which results in the smallest coercive field. The appearance of any defect, such as a void or nonmagnetic precipitate, could result in larger coercive field, i.e., magnetic hardening, if it can prevent the formation of 360° domain walls during domain switching.

More importantly, the developed model enables investigation of the interaction mechanisms between nonmagnetic particles and magnetic domain walls and improves the understanding of the physics behind observed signals in NDE, a critical component for the development of advanced NDE and sensors for monitoring material degradation under thermal and irradiation environments.

In irradiated materials there are complex microstructure including initial microstructures such as martensitic phases, grain boundaries, and dislocations in ferritic-martensitic steels, and radiation-induced defects such as precipitates, dislocation loops, and voids. Material property degradation depends on the types of defects, defect densities, and spatial and temporal evolution of defect morphology and distributions. Our ultimate goal of modeling is to quantitatively study the effect of different microstructures on material response to applied fields, understand the interaction mechanisms and signal physics, validate the simulation results with experimental measurements including domain nucleation at defects, magnetic domain structures, first order reversal curves (FORCs), and hysteresis loop, statistically analyze the correlation between distributed defects and response, and provide insights for the development of advanced monitoring methods and sensor materials.

5 Acoustic Case Study

Another in-situ approach to detecting materials degradation is through the use of acoustic methods to characterize and monitor material performance. The intent of this document is not to describe all of the available ultrasonic measurement methods for materials characterization. The behavior of acoustic waves in solids (and the consequent measurable parameters that may be used to correlate with microstructural changes) are described in detail elsewhere (for instance, (Ensminger and Bond 2012; Krautkramer and Krautkramer 1990; Goebbels 1994; Doctor et al. 1989; Raj et al. 2000)). Here, we summarize some of the ultrasonic measurement methods that have been applied for materials characterization:

- **Ultrasonic velocity and attenuation:** These techniques use classical ultrasonic measurement methods to measure the velocity and attenuation, potentially as a function of frequency, and have been shown to be sensitive to a range of degradation mechanisms.

- **Bulk wave** measurements use body or bulk waves to measure the velocity and attenuation, usually as a function of wave mode (longitudinal or shear).
- **Guided wave** measurements have been proposed as a potential technique for long-range examination of components such as piping. The method uses the structure or component as a waveguide to propagate the applied stress waves over long distances. Sensitivity to large-scale cracking has been shown in prior studies.
- **Acoustic emission (AE)**: AE is a passive technique that “listens” for stress waves initiated by the onset of cracking or other dislocation movement. AE has been shown in prior work to be sensitive to crack growth (Hutton et al. 1993). A potential challenge is to discriminate crack growth signals from environmental noise.
- **Nonlinear acoustics** (Ogi et al. 2001; Cantrell and Yost 2001): This class of techniques attempts to measure the relative change in the nonlinear elastic wave response of the material due to accumulated damage. Techniques included in this method can employ:
 - Bulk measurements (Cantrell and Yost 2001)
 - Rayleigh wave measurements (Shui et al. 2008)
 - Guided wave measurements (Bermes et al. 2008)
- **Acoustic microscopy**: This class of techniques (mostly laboratory-based) uses high frequency (typically in excess of 20 MHz) focused ultrasound to image variation in elastic properties in materials. A number of enhancements are possible, including surface imaging, sub-surface imaging, leaky Rayleigh wave, and nonlinear acoustic microscopy measurements.

In general, resolution for linear acoustics measurements is proportional to the wavelength in the material, and for direct detection of degradation precursor states in steels, the wavelength needed is likely to be on the order of 10-25 microns (corresponding to a frequency of 250-500 MHz). While this is possible, the higher frequency range will push the limits of current generation of piezoelectric probes. Also at higher frequencies, there are potential issues with attenuation of the sound field deep into the material, so it may be only possible to image near surface regions. Gaps in this respect include: localized measurement interior to specimens, radiation performance of probes not known (but assumed to be a function of the composition/material, fabrication process and impurities).

In recent years, nonlinear ultrasonics (NLU) has seen increased interest as a means of characterizing the internal damage state of the material early in the fatigue process (Cantrell and Yost 2001; Matlack et al. 2012b). Conventional ultrasonic methods, currently in use in ISI, apply high-frequency (in excess of 500 kHz, typically around 2.25 MHz) acoustic energy and measure the resulting response due to scattering and reflection of the energy at interfaces. The presence of cracking is detected by means of a reflection from the crack surface. Other forms of damage may

be detected by making use of velocity and attenuation measurements through one or more measurement configurations. However, most such measurements are not sensitive to earlier stages of damage. In contrast, NLU methods rely on the generation of harmonics from an initially monochromatic input. The generation of harmonics is due to nonlinearities in the elastic constants associated with the material (Zarembko and Krasil'nikov 1971). The second harmonic is of particular interest, and the resulting nonlinear material parameter is represented by β (Kyung-Young 2000). A schematic of a typical single-sided measurement setup for nonlinear acoustics measurements is shown in Figure 12. Alternative setups that use a transmitting probe to transmit acoustic energy through the specimen and a receiving probe on the opposite side of the specimen (through-transmission mode) are also possible, as are methods that generate surface waves.

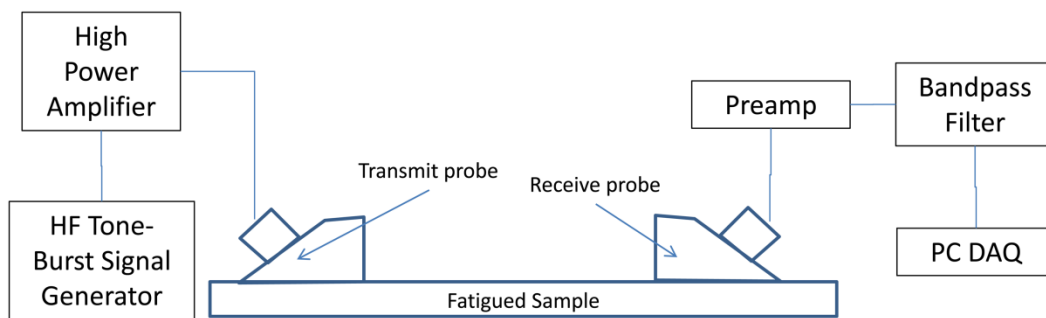


Figure 12. Schematic of NLU Measurement System

NLU has been applied to the characterization of a range of damage mechanisms, including fatigue (Kyung-Young 2000; Cantrell and Yost 2001), irradiation embrittlement (Matlack et al. 2012b), SCC (Matlack et al. 2012a; Shintaku et al. 2010), and corrosion pitting (De et al. 2010).

5.1 Acoustic Sensor Materials and Irradiation Studies

The most common sensor materials for ultrasonics are piezoelectric materials. Of the many different piezoelectric materials, PZT (lead zirconate titanate) is among the most widely used, primarily due to its relatively high sensitivity. Magnetostrictive materials are also used for ultrasonic measurements, as are EMAT probes and lasers (both of which enable non-contact measurements).

Ultrasonic measurements for detecting microstructural changes representative of irradiation-related embrittlement in steels have been performed by a number of researchers (Doctor et al. 1989; Blaszkiewicz 1996; Change et al. 2001; Dobmann et al. 2001; Ishii et al. 2002; McHenry 1998). Most studies to date have examined ultrasonic velocity and attenuation of different wave modes and correlated the resulting measurements with some metric of

embrittlement (usually fluence or microhardness). Results to date appear to indicate some trending; however, the results vary from study to study and in many cases, are based on a limited sample set. In general, the results are likely to be a function of the specific alloy, the starting microstructure, and fluence and flux levels. Recent studies have also examined ultrasonic absorption for sensitivity to irradiation damage; again, the results are encouraging though based on a limited sample set. Recent studies have also examined the capability of nonlinear ultrasonic measurements to assess irradiation damage (Matlack et al. 2012b). While promising, the results are preliminary and require further validation.

A complementary issue is that of sensor material irradiation hardening. Studies to date have shown that several piezoelectric materials may be relatively radiation-hardened, though the fluence levels achieved in these tests is generally rather low ($\sim 10^{18}$ n/cm²). It is likely that long-term exposure, particularly in advanced reactor designs, may result in significantly higher dose, and therefore a recent project to examine the impact of increased irradiation levels on transducer performance was started at Penn State, using the MIT Materials Test Reactor (MITR). The issue of high temperature operation is related, as the sensor is expected to encounter temperatures in excess of $\sim 300^\circ\text{C}$ in advanced reactor concepts as well as close to the core. As a result, materials with Curie temperatures less than the expected operating temperature will not survive. On the other hand, materials with Curie temperatures in excess of the operating temperature may still encounter aging and degradation in sensor response over the long term, and this effect needs to be quantified.

5.2 Acoustic Sensor Modeling and Signal Interpretation

Acoustic sensor modeling at the macroscale is a relatively well-developed field, with a number of commercial software packages available, including ANSYS and CIVA. However, these packages require modification to handle propagation issues arising from complex materials such as large anisotropic grains. The inverse problem of signal interpretation and analysis at the macroscale has also seen a significant level of attention in the research community and a summary of some of the approaches to addressing the macroscopic signal inversion problem are documented in reference (Daw et al. 2012). Modeling and analysis to address elastic behavior at the microscale are limited.

5.3 Imaging

Ferritic steel of grain size ~ 17 μm has been imaged as a function of depth with scanning acoustic microscopy (SAM) at 967 MHz (Hirse Korn et al. 2011). Atomic force microscope (AFM) and atomic force acoustic microscopy (AFAM) images of these steels can be combined with EBSD to give a complete picture of the local crystallographic and acoustic properties. In another example, cementite and ferrite phases in Fe-0.8%C unalloyed steel can be distinguished using AFAM contact resonance spectra (Batista et al. 2012). Studies of this nature imply that defects could be imaged and local acoustic property maps could be measured, allowing the nano-

/micro-structural investigation of the effects of voids, second phases, and grain boundaries on physical properties.

6 Related work at other institutions

Thought leaders in the aerospace NDE community have already called for the use of integrated materials property assessment in structural materials as a means to achieve state-of-health determination (“materials state awareness”) (National Research Council 2008; Banerjee et al. 2011). The term “electromagnetic materials state awareness monitoring” has been coined to describe the use of NDE methods that result in an electrical signal produced by an electrical, magnetic, or thermal stimulus (Nagy 2010). Integrated monitoring and prognosis using techniques with sufficient sensitivity and selectivity allow various aspects of material and component state to be assessed through changes in electrical conductivity, magnetic permeability, thermal conductivity, and thermoelectric power, for instance. Through the measurement of these properties, microstructure evolution, phase transformation, hardening, embrittlement, and elastic strain can be assessed. For example, electrical conductivity in metal alloys will change as a function of elastic strain.

Scientists at the Japanese Atomic Energy Agency have recently been developing NDE techniques to detect irradiation-assisted stress corrosion cracking (IASCC) susceptibility in austenitic stainless steels, using eddy current and AC magnetization methods (Nemoto et al. 2011).

Recent work by the Fraunhofer Institute for Non-destructive Testing (IZFP), has combined electron backscatter diffraction (EBSD) imaging with scanning acoustic microscopy (SAM) and scanning probe techniques such as AFM, AFAM, and MFM to show the potential of these techniques for studying the detailed mechanisms of grain and second phase formation in steels (Batista et al. 2012; Hirsekorn et al. 2011).

Researchers at Arizona State University have been evaluating an integrated approach to modeling and measurements with the goal of predicting remaining life of fatigued aerospace materials. They use first principles computational models to study the evolution of damage in the material, and measurements at the bulk scale along with computational approaches to inversion (Soni and Chattopadhyay 2010; Soni et al. 2009).

7 Acknowledgements

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8 Appendix: literature review of effects of neutrons on magnetic materials

	flux (n/cm ² /s)	fluence (n/cm ²)	Ref, notes	Material effects
Mn-Zn ferrite, Ni-Zn ferrite ceramics	5 x 10 ¹³	10 ¹⁸ to 2.3 x 10 ²⁰	(Chukalkin et al. 1975) Fast neutr E>0.8 MeV; temperature 70 °C, thermal neutrons removed with Cd	Above 10 ¹⁸ /cm ² dose, cations redistribute into disordered nanoregions (1.2-3.5 nm), T _c increases, max magnetization (low T) decreases; at highest fluence, Ni-Zn conductivity changes from p-type to n-type (Chukalkin et al. 1974)
ZnFe ₂ O ₄		0.25 x 10 ¹⁹ to 10 ²⁰	(Parkhomenko et al. 1976) Temperature 70 °C, thermal neutrons removed with Cd	Zn ferrite, normally a paramagnet at room temperature, has been shown to become a ferromagnet at the highest neutron doses, and a superparamagnet with ferrimagnetic clusters at lower doses due to replacement cascades of atoms in the spinel sites. Changes to ferromagnet at highest doses by change of Fe ³⁺ to Fe ²⁺
BaFe ₁₂ O ₁₉		2 x 10 ¹⁹ to 3.4 x 10 ²⁰	(Chukalkin et al. 1981) Fast neutrons E>1 MeV in reactor	With increasing fluence, saturation magnetization decreases, does not totally saturate at high magnetic field, and has kink in curve; magnetic structure changes to helicoid (and ultimately isolated block) structure possibly due to removal of Fe ³⁺ from 2b sites that normally result in collinear structure
Y ₃ Fe ₅ O ₁₂		5 x 10 ¹⁸ to 1.6 x 10 ²⁰	(Chukalkin et al. 1983) Fast neutrons E>1 MeV in reactor	Amorphization of YIG at highest fluences; sample becomes a spin glass with T _f ~70 K and shows bifurcation in FC/ZFC magnetization; summary discussion on spinels and hexaferrite crystal structures not changes by high fluences but garnets changed due to large ion size difference between rare earths and transition metals resulting in frozen disorder.
Y ₃ Fe ₅ O ₁₂ powdered and polycrystalline	10 ¹³		(Podsekin and Zaitsev 1982) fast neutrons E>0.1 MeV	As dose rate increased, unit cell period increased and Curie temperature decreased
Y ₃ Fe ₅ O ₁₂	8 x 10 ¹² to 1.5 x 10 ¹³	10 ¹⁸ to 10 ²⁰	(Podsekin et al. 1982) fast neutrons E>0.1 MeV	Frequencies up to 10 GHz, complex magnetic permeability changes with fast neutron fluxes >10 ¹⁸ /cm ² ; permeability in predispersion region (up to 100 MHz) decreases and increases in RF dispersion region; resistivity and sat magnetization decreases

Appendix: literature review of effects of neutrons on magnetic materials (continued)

	flux (n/cm ² /s)	fluence (n/cm ²)	Ref, notes	Material effects
Y ₃ Fe ₅ O ₁₂ and YIG:La epitaxial films		10 ¹⁷ to 10 ¹⁹	(Ubizskii et al. 1993) fast neutrons >0.2 MeV, average energy 2 MeV; results compared to electron irradiation (1.3 MeV (Ubizskii et al. 1993) and 3.5 MeV (Ubizskii et al. 1996))	At fluences >5 x 10 ¹⁷ /cm ² , YIG lattice constant increases (restored to pre-irradiated value after annealing at 1250 K in air), saturation magnetization decreases (~8% at 10 ¹⁹ /cm ²), ferromagnetic linewidth increases (factor of 2 at 10 ¹⁹ /cm ²), and shortwave absorption edge broadens; data are explained by assuming disordered nanoregions of 6-8nm created from displacement cascades and resulting in amorphous areas which spread out the energy of transitions at the optical edge.
Sm-Co permanent magnets			(Ito et al. 2001) 200 MeV protons 10 ² -10 ⁵ Gy	No change in magnetic flux with irradiation
Nd-Fe-B permanent magnets			(Ito et al. 2001) 200 MeV protons 10 ² -10 ⁵ Gy	Magnetic flux decreased after irradiation; increase of permanence coefficient and coercive force improved radiation hardening; operating point on demagnetization curve important
Nd-Fe-B and SmCo permanent magnets	4.2 x 10 ¹⁰ (fast neutrons, 1 MeV equivalent)		(Spencer and Volk 2003) data on Hadrons needed for accelerators; (Anderson et al. 2005) fast neutron test at UC Davis McClellan Nuclear Reactor Center, suppressed thermal neutrons and gammas.	(Anderson et al. 2005) Induced radioactivity discussed; differential damage in Gauss/Gy negative (demagnetization)
Nd-Fe-B and SmCo permanent magnets		10 ¹⁶ (fast neut, 1 MeV equiv)	(Liu et al. 2007) fast neutron test at Ohio State Research Reactor	Radiation-induced thermal spike caused Nd ₁₃ Dy ₂ Fe ₇₇ B ₈ to lose 100% of magnetic flux with 10 ¹⁶ n/cm ² of neutron flux. Sm(Co,Fe,Cu,Zr) ₂ magnets do not show noticeable change. Best Sm-Co magnets have T _{curie} up to 750-900 C, so can operate up to ~550 C.
Fe ₂ O ₃				(Bakkaloglu and Thomas 1992) change in Morin T; (Donbaev et al. 1993) creates low T weak FM phase

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