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# Applications of Ab Initio Modeling to Materials Science: Grain Boundary Cohesion and Solid State Diffusion

GA Young, R Najafabadi, W Strohmayer, J Vollmer, C Thompson,  
W Hamm, C Geller, E Wimmer, J Sticht, A Mavromaras and J Harris

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# Applications of Ab Initio Modeling to Materials Science: Grain Boundary Cohesion and Solid State Diffusion

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Angel Fire, NM 87710-1318*



# Background

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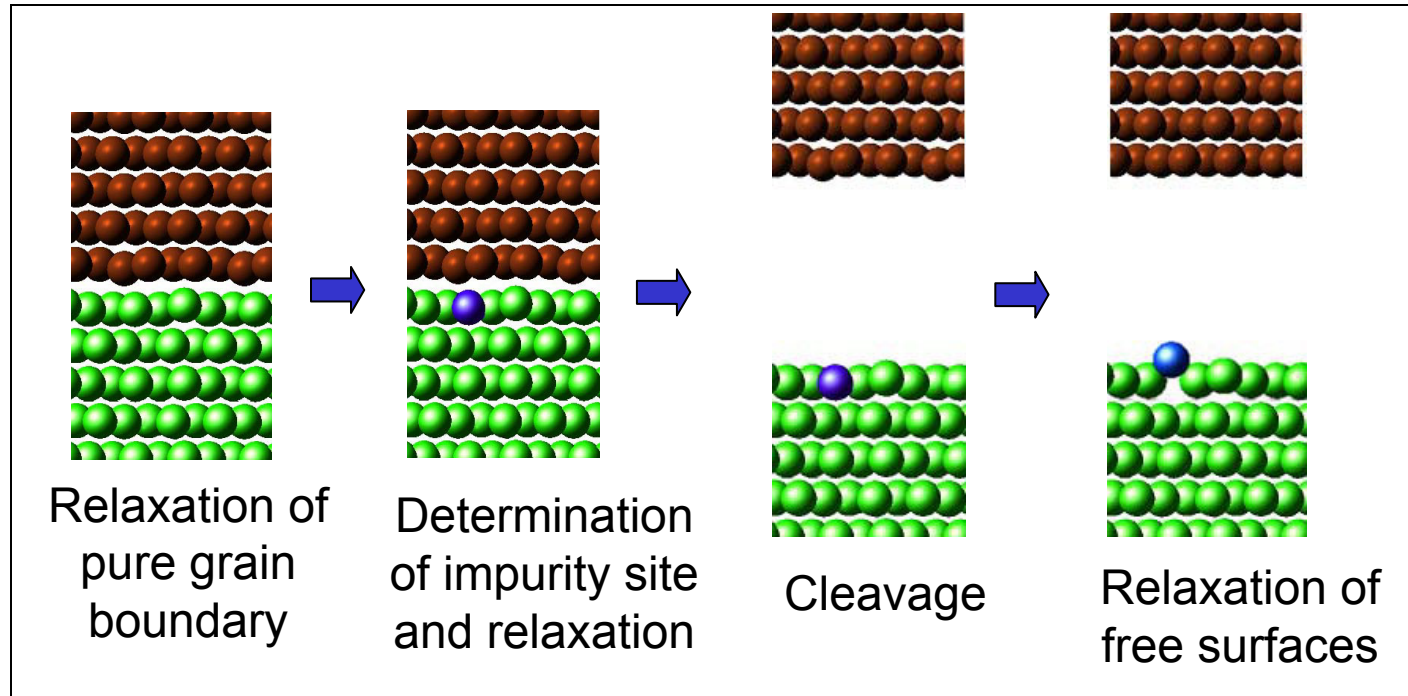
- I. Grain boundary cohesion

- Metallurgical understanding
- Account for multiple elements
- Effect of H concentration
- Mechanisms of embrittlement
  - He embrittlement
  - SCC resistance

- II. Solid State Bulk Diffusion

- Work of Wolverton *et al.*
- Proof of principle for H in nickel
- Determination of  $D_0$  and  $Q$  from first principles
- H diffusion in HCP Ti and Zr
- H diffusion in FCC Fe – effect of lattice parameter

# Computational Procedure: Pure nickel, $\Sigma 5$ - $\{100\}$ twist grain boundary, 0 K



$$\Delta E_{Griffith} = \left( E_{Surface} - E_{Surface}^{impurity} \right) - \left( E_{Grain Boundary} - E_{Grain Boundary}^{impurity} \right)$$

# Geometry of impurity atoms on nickel $\Sigma 5$ {100}

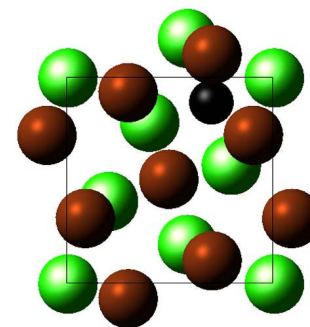
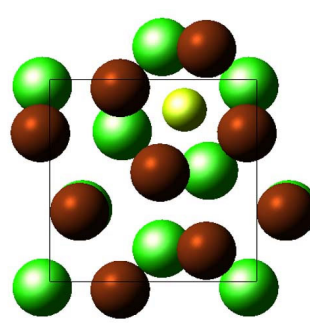
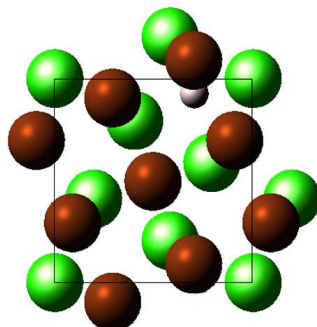
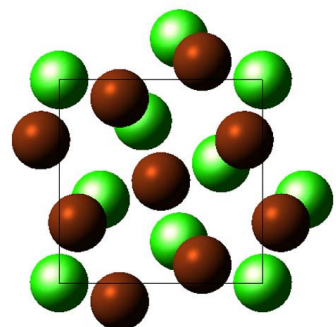
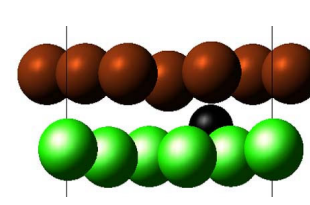
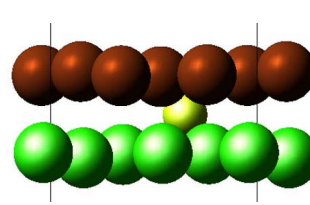
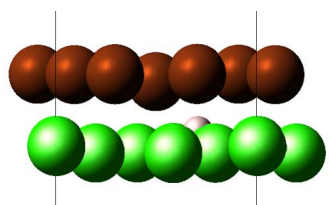
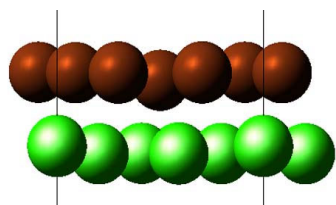
## Interstitial impurities

pure Ni

H

B

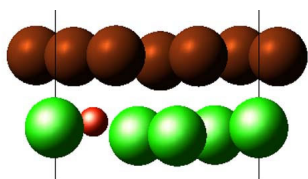
C



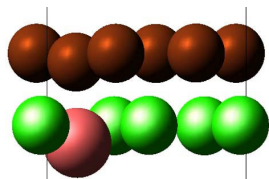
# Geometry of impurity atoms on nickel $\Sigma 5$ {100}

## Substitutional impurities

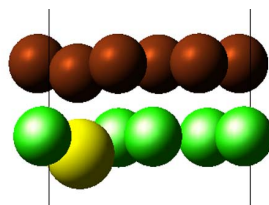
He



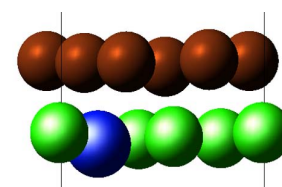
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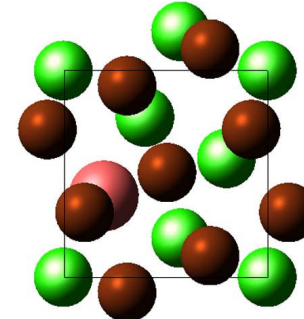
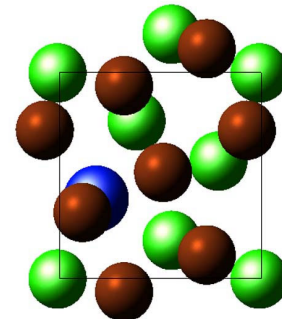
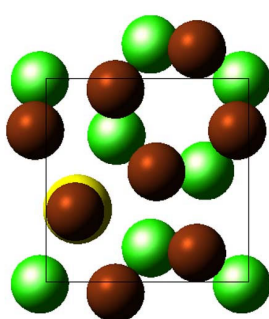
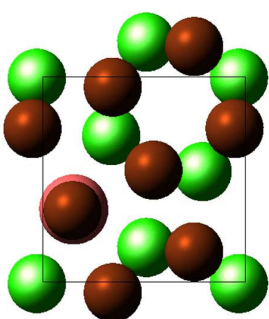
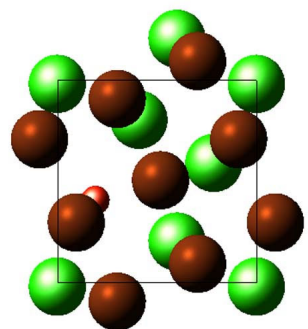
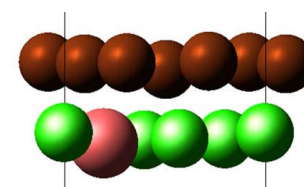
S



Cr

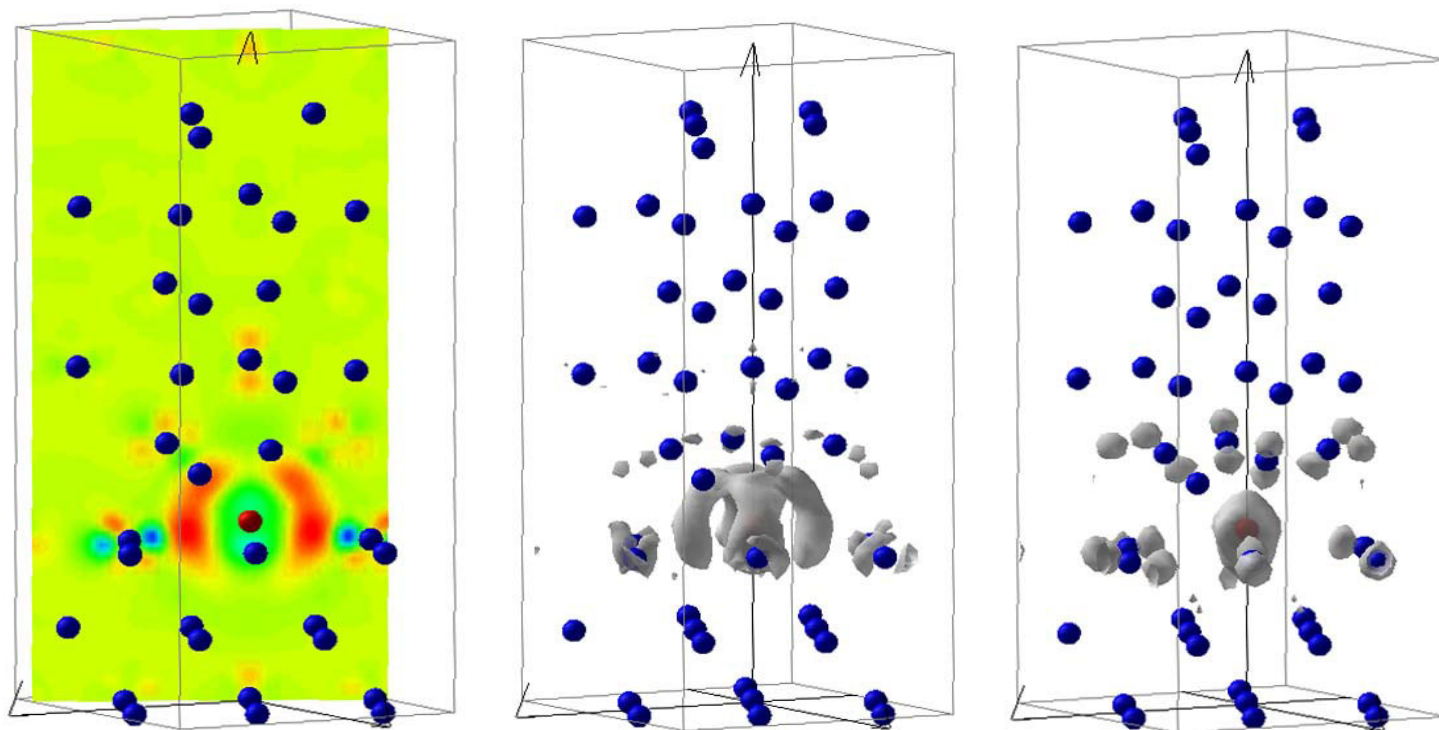


Fe



# Structure / energy mapping: degree and extent of electronic interactions

Electronic charge changes around a P impurity atom in Ni



(a)

Gain (red), loss (green)

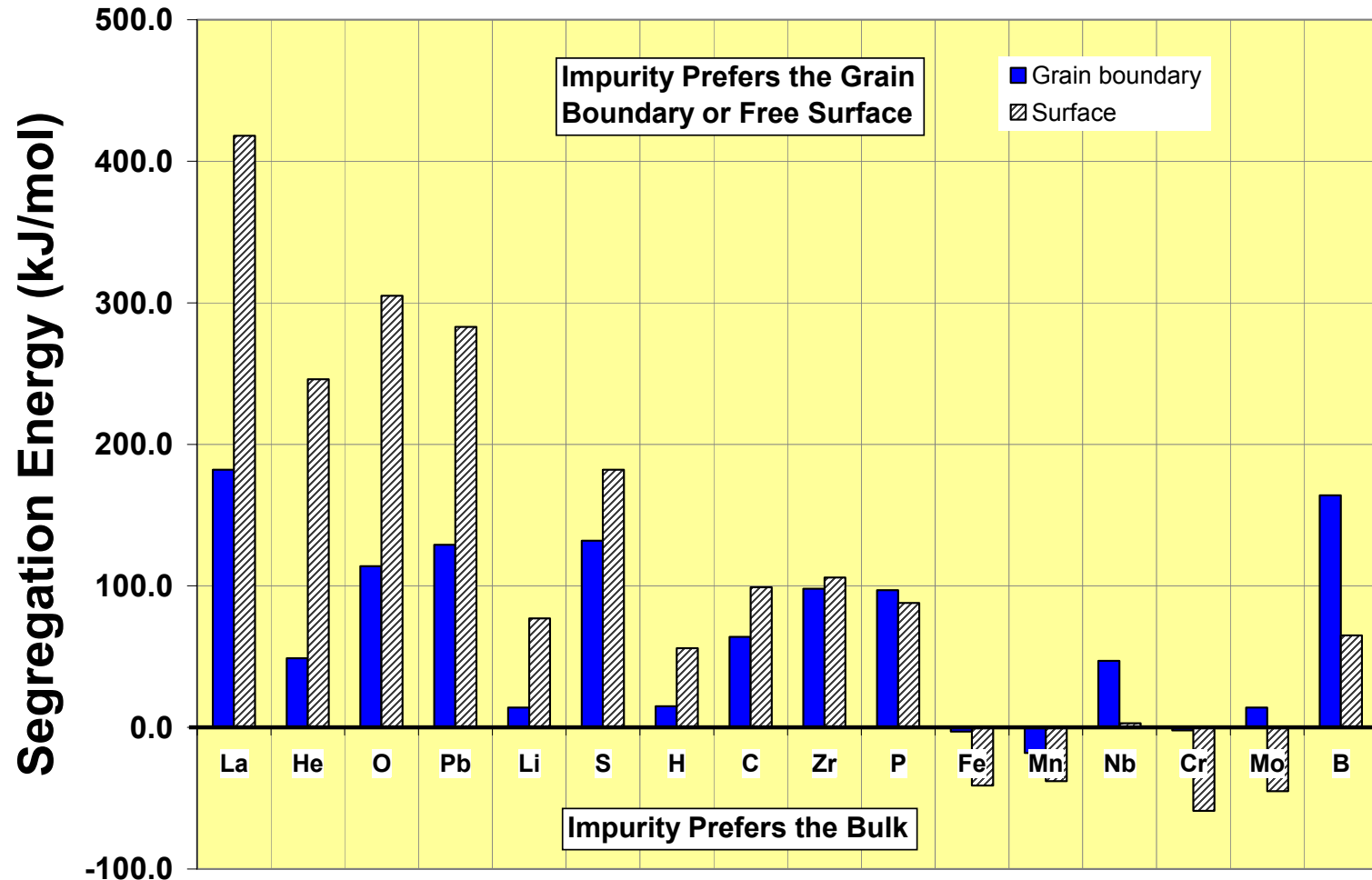
(b)

Gain

(c)

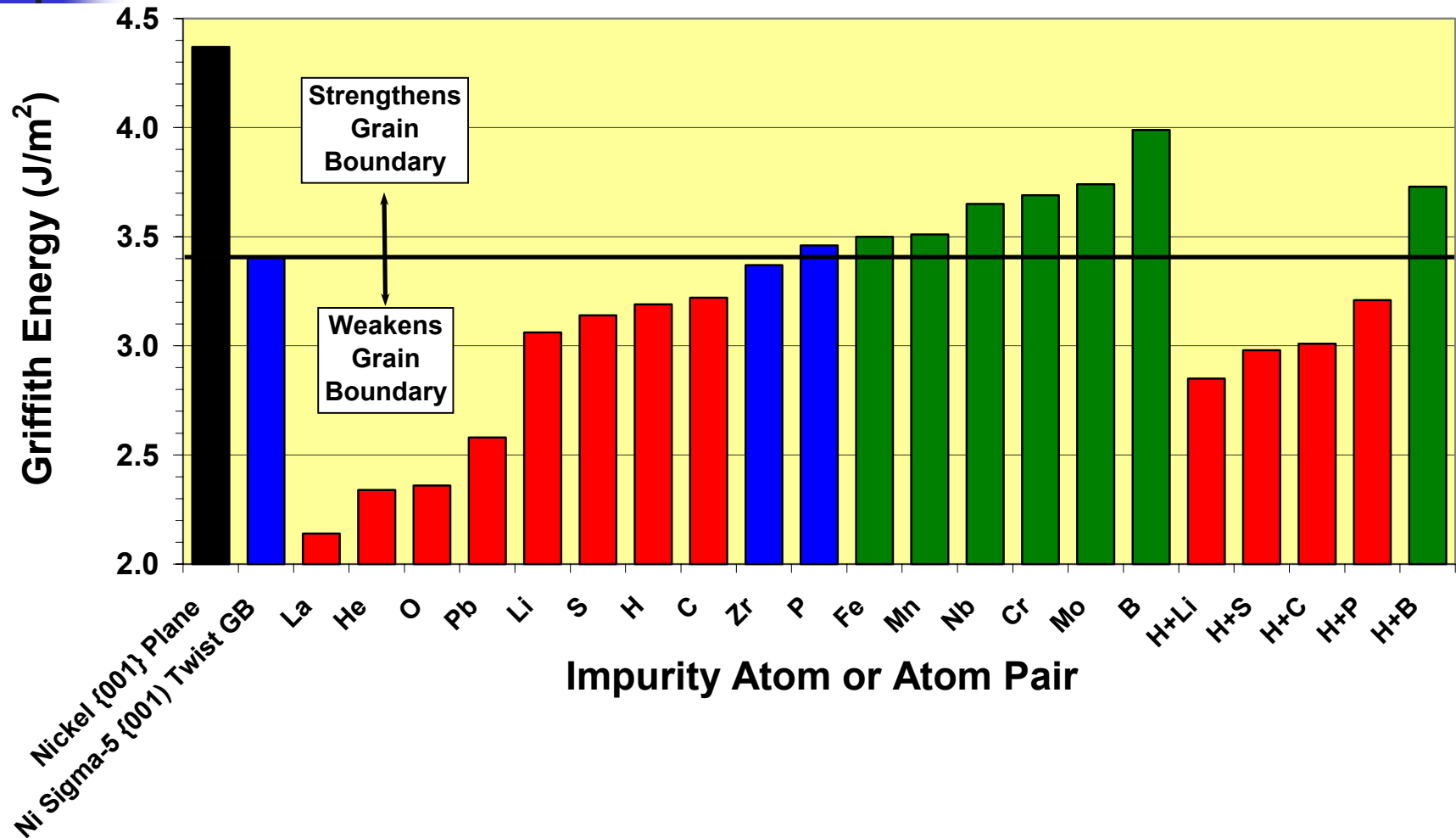
Loss

# Tendency to migrate to grain boundary, surface, or bulk (Note boron)





# Fundamental effects of impurities on grain boundary strength



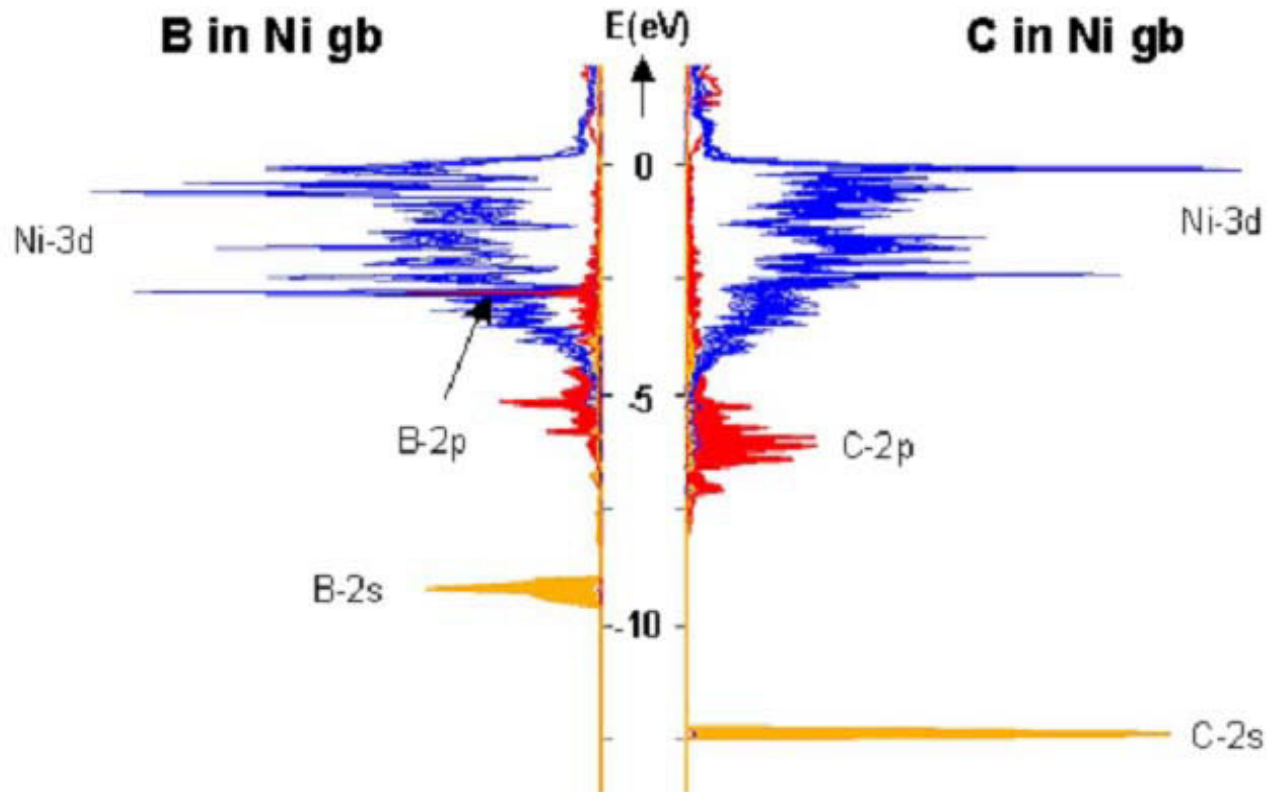


## Metallurgical strengthening of grain boundaries

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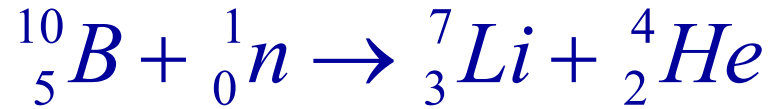
- Boron a known strengthener of Ni grain boundaries
- Theories (Donachie from Superalloys Source Book):
  - Boron alters grain boundary precipitate structure
  - Boron has a beneficial interaction with a deleterious element (*e.g.* ties up sulfur ?)
  - Boron reduces the grain boundary diffusivity (slows S segregation ?)

# Why does boron promote grain boundary strength (but carbon doesn't) ?

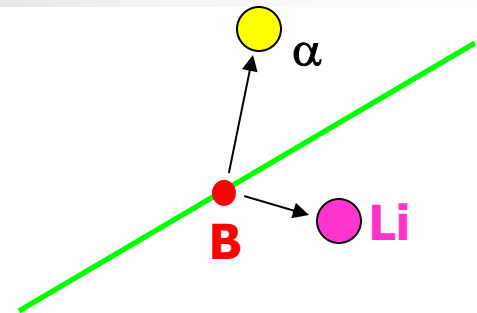


B-2p electrons overlap with Ni-3d → bonding

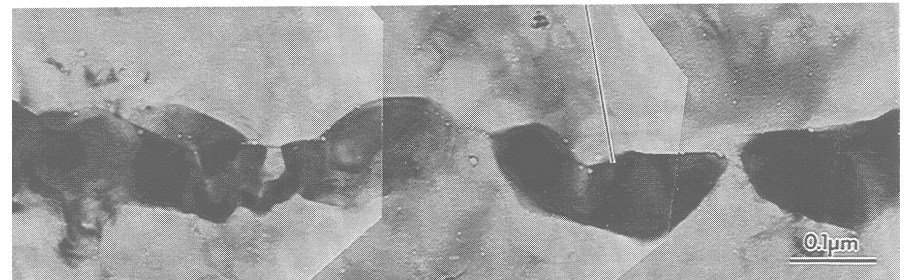
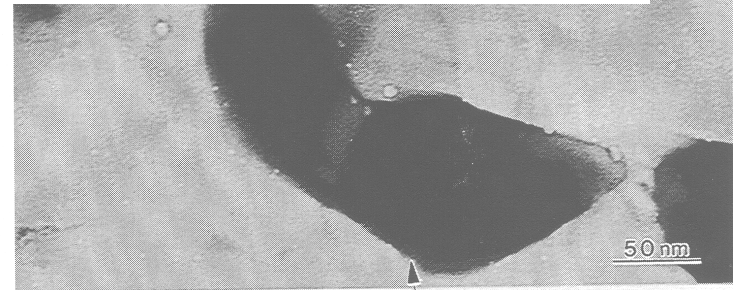
# Mechanism of "He Embrittlement"



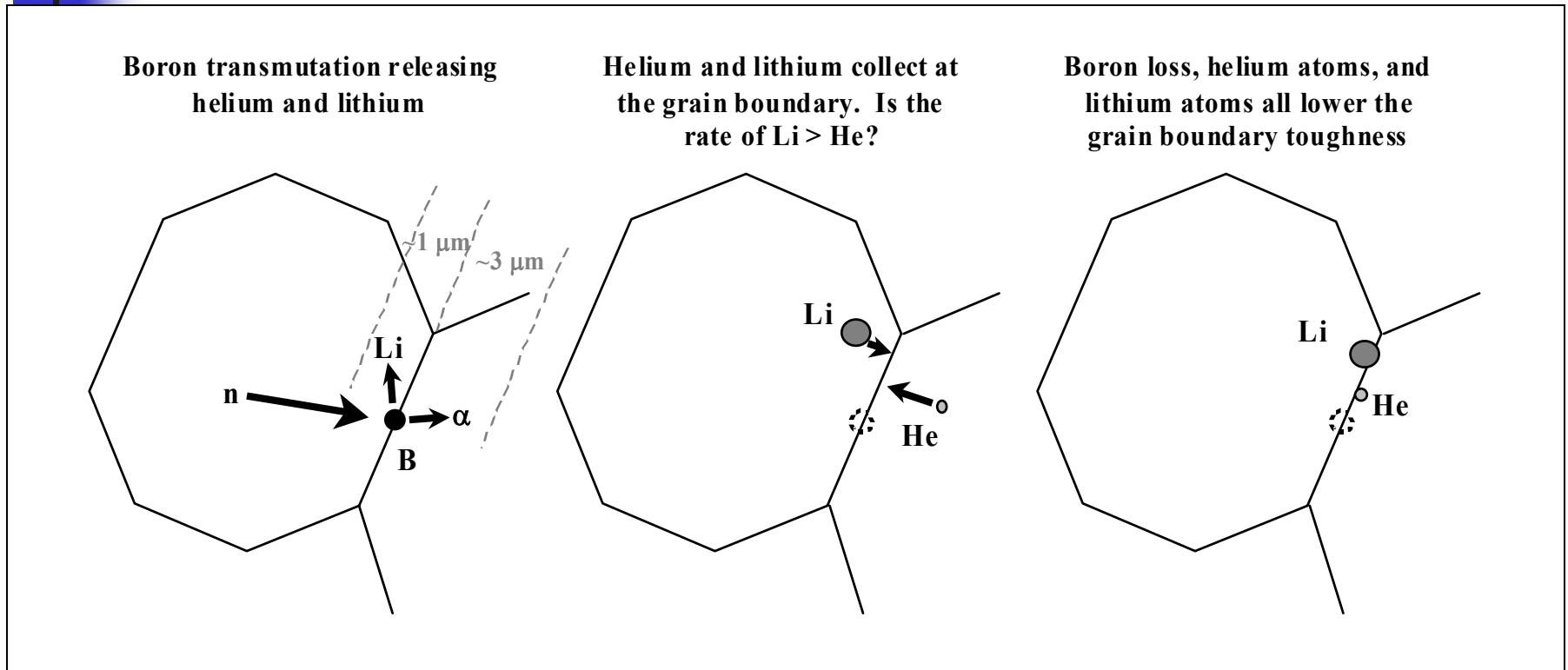
- After irradiation *and annealing*, He bubbles on nickel grain boundaries have been observed
- *But* you embrittle before you observe physical bubbles – intrinsic He embrittlement (see Mills *et al.*)
- Need to assess B loss and Li embrittlement, kinetics of He and Li diffusion



Bill Mills *et al.* 7th Env. Deg. Proceedings

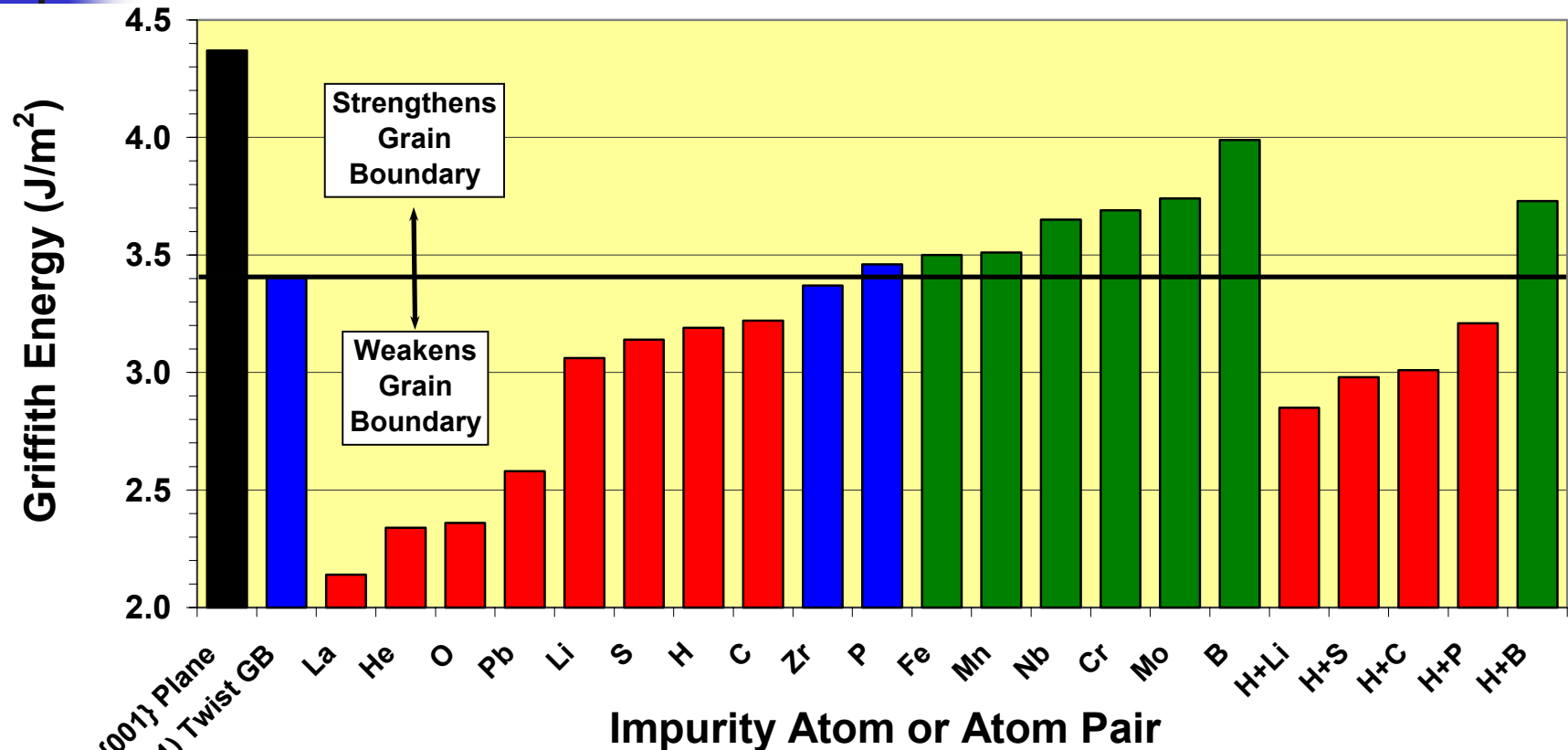


# B loss + Li embrittlement + He embrittlement



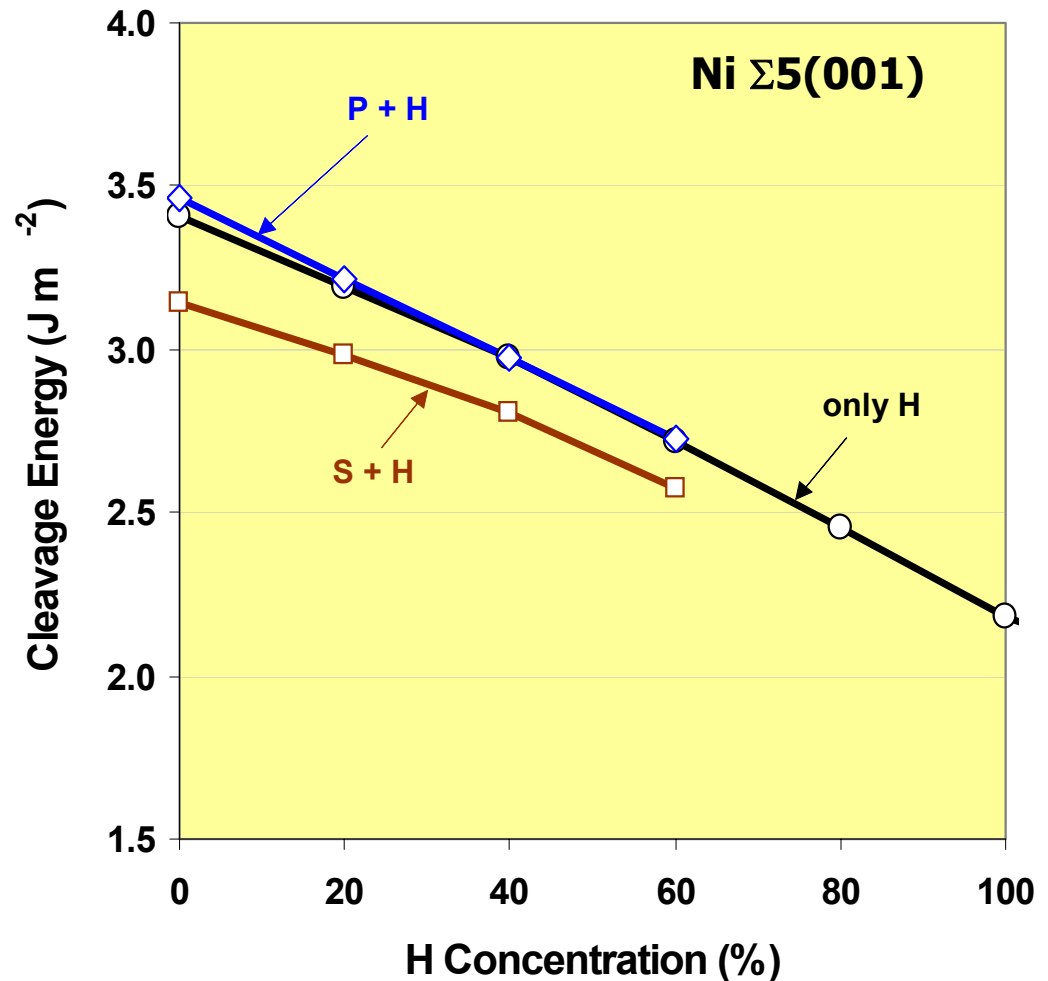
→ Need to understand kinetics of Li and He migration back to grain boundary

# He and Li embrittlers, B strengthener

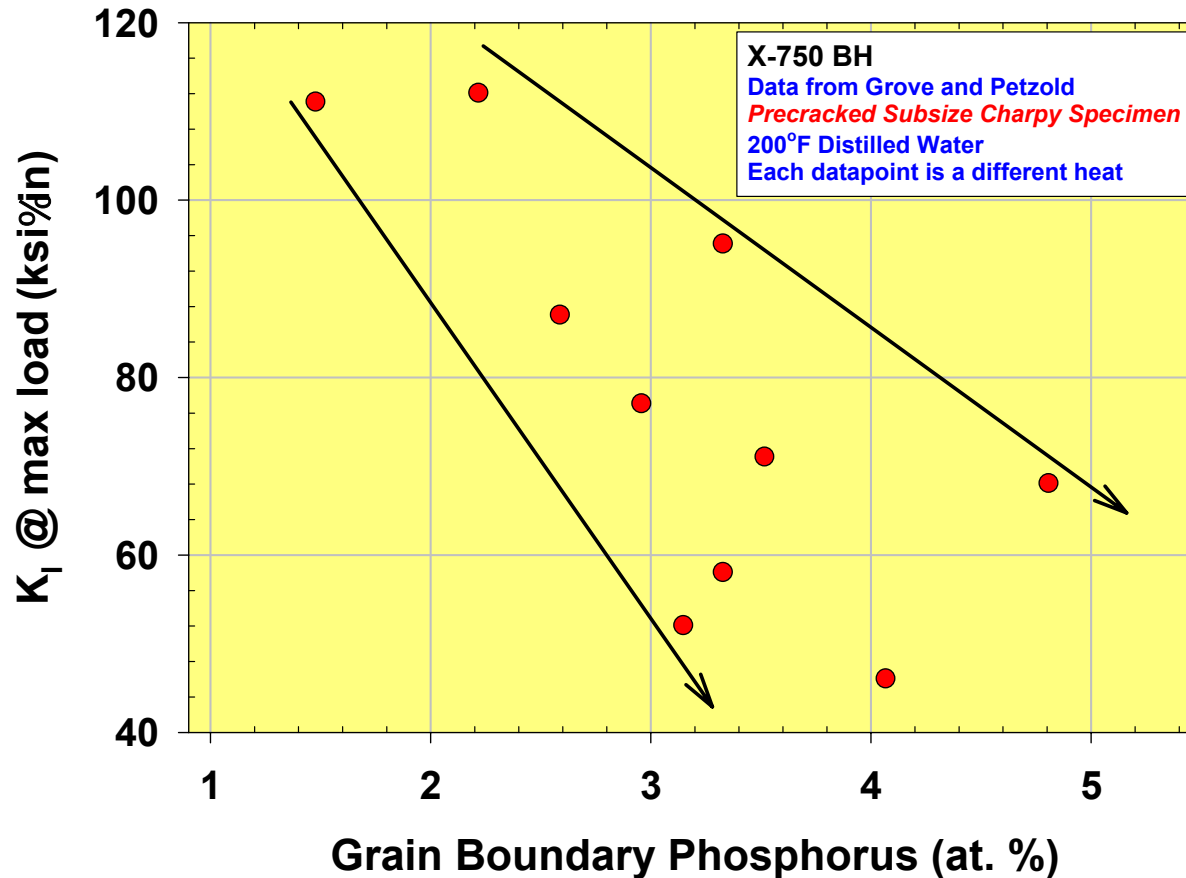


# Effect of H concentration on cleavage energy with and without S and P impurities

- Linear superposition good first approximation
- No significant synergistic effects
- S and H act additively



# Need to account for multiple effects: (1) GB strength and (2) effect on H uptake



$$\text{Relative Grain Boundary Strength} = -4.86[\text{He}] - 3.77[\text{Pb}] - 1.58[\text{Li}] - 1.00[\text{H}] - 0.86[\text{C}] - 0.18[\text{Zr}] \\ + 0.23[\text{P}] + 0.41\{\text{Fe}\} + 0.45[\text{Mn}] + 1.09[\text{Nb}] + 1.27[\text{Cr}] + 2.64[\text{B}]$$





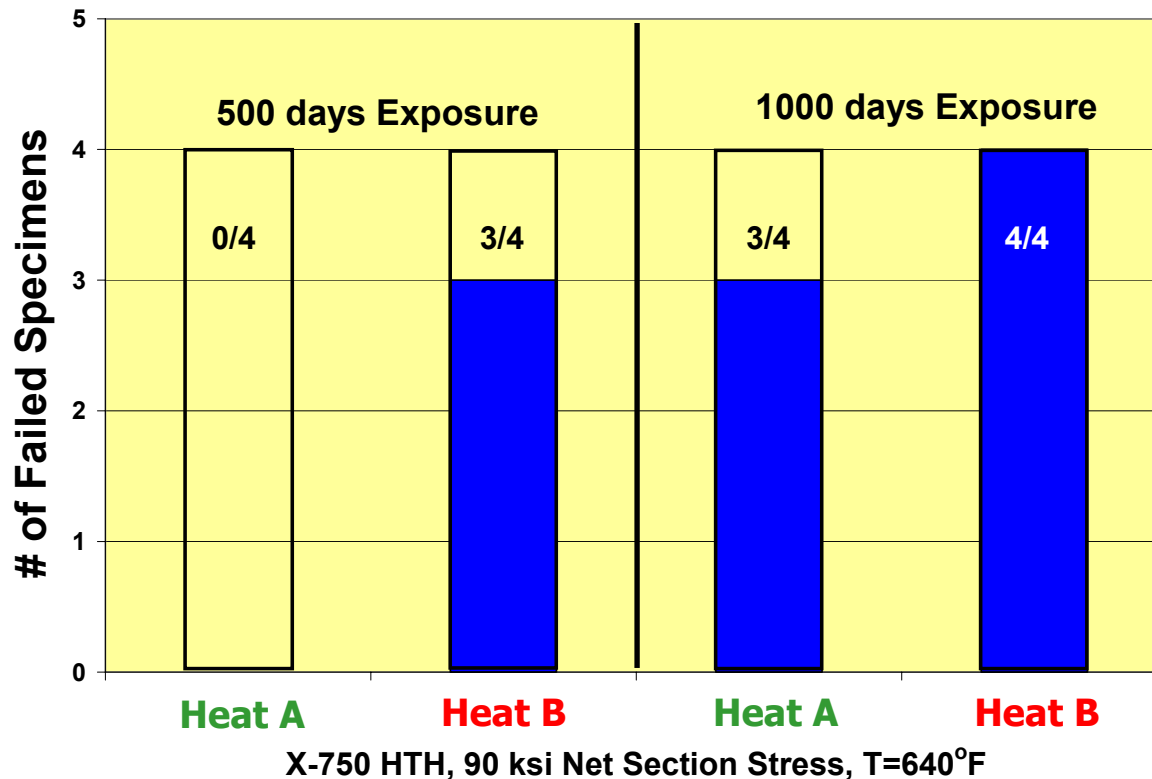
# SCC of X-750: Similar heats with different response to SCC initiation and growth

## **X-750 Condition HTH (2025°F/1hr + 1300°F/20 hrs)**

<b>Heat</b>	<b>Ni</b>	<b>Cr</b>	<b>Fe</b>	<b>Ti</b>	<b>Al</b>	<b>Mn</b>	<b>C</b>	<b>B</b>	<b>P</b>	<b>S</b>	<b>Nb+Ta</b>
<b>A</b>	72.43	15.54	7.93	2.60	0.79	0.07	0.039	0.0037	0.007	0.001	0.86
<b>B</b>	71.47	15.25	8.15	2.66	0.73	0.15	0.043	0.0022	0.002	0.002	0.97

<b>Heat</b>	<b>YS (ksi)</b>	<b>UTS (ksi)</b>	<b>% El</b>	<b>%RA</b>
<b>A (Good)</b>	118	172	27	32
<b>B (Bad)</b>	117	176	26	36

# Heat "B" shows shorter initiation times and faster crack growth rates



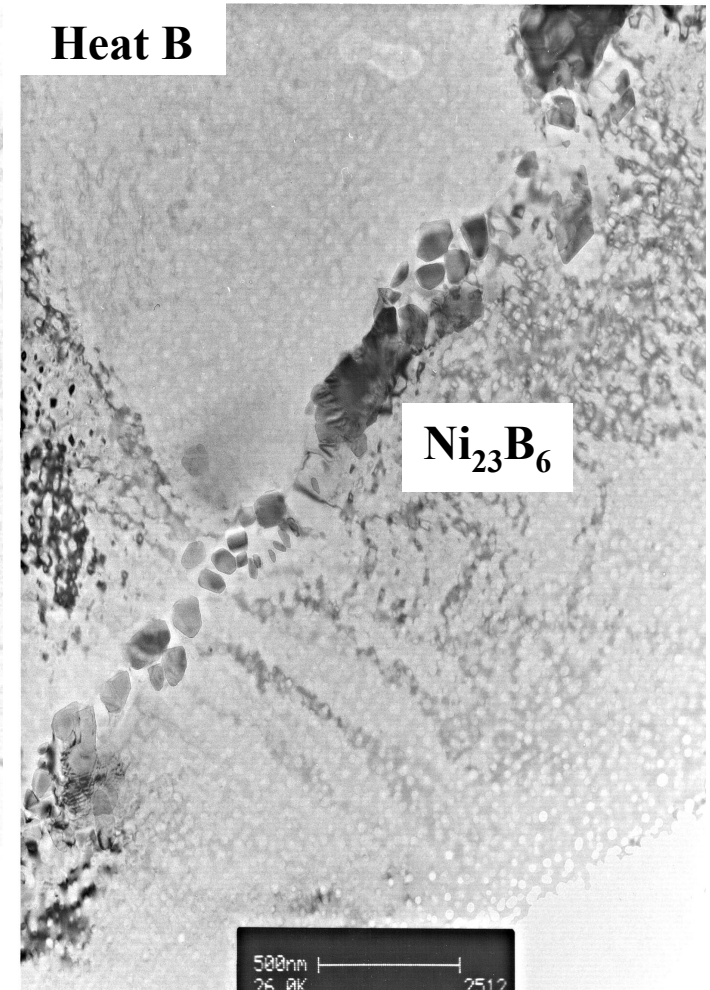
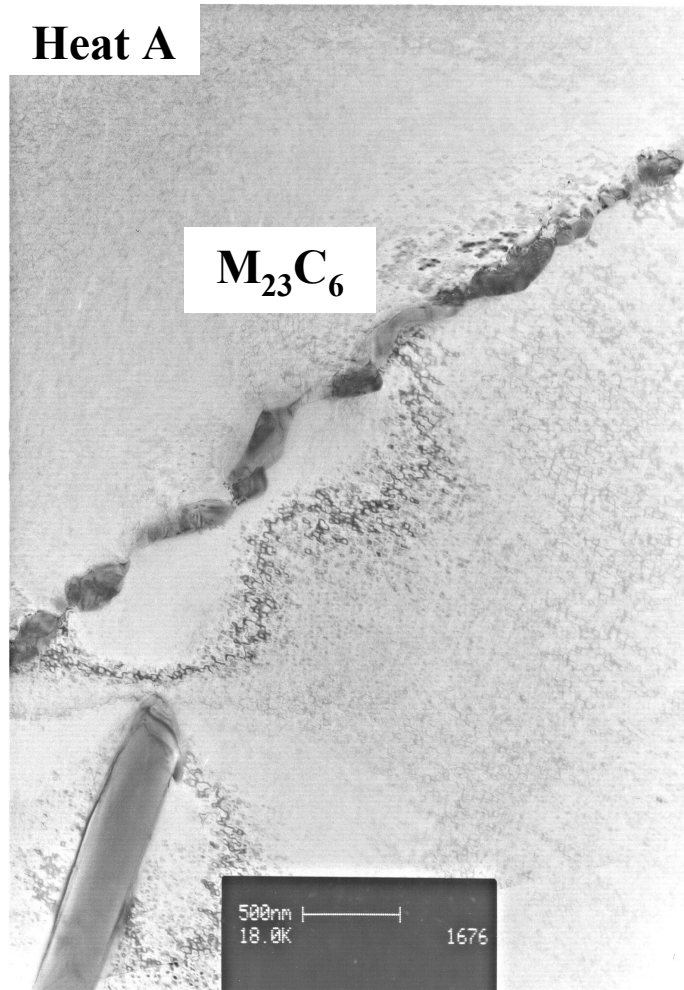
Heat	CGR at 680°F K=35 ksi√in
A	1.1 mils/day
B	1.8 mils/day

See Young *et al.* in 11<sup>th</sup> International Symposium on Environmental Degradation of Materials in Nuclear Power Systems – Water Reactors, Skamania, WA, August 2003.

Heat A - typical  $\text{Cr}_{23}\text{C}_6$  carbides

Heat B - unusually high # of  $\text{Ni}_{23}\text{B}_6$

- Loss of atomic boron
- No Cr depletion around borides
- Faster SCC initiation and growth





## Summary

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- Ab Initio atomistic modeling gives unique insight into metallurgical effects
  - Explain alloying effects: Boron intrinsically strengthens Ni grain boundaries by helping to fill the Ni 3d orbital
  - De-convolute complex embrittlement phenomena: boron transmutation / stress corrosion cracking
  - Quantitatively assess the effects of multiple grain boundary impurities

**Embrittling** ← **Neutral** → **Strengthening**

Al He O Pb Li S H C Zr P Fe Mn Nb Cr Mo B



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# Ab Initio Modeling of Solid State Diffusion: Hydrogen in Structural Metals



## Ab initio modeling of solid state diffusion

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- Phonon capability in MedeA/VASP is a significant advance, allowing accurate determination of things like enthalpy, and entropy *as a function of temperature*  $\rightarrow$  *free energy*
- Phonons enable first principles studies of diffusion, solubility, etc. that are often experimentally difficult and subject to significant controversy, *e.g.* H diffusion in Al, Ti, O solubility in Ni, etc.
- Examples
  - Chris Wolverton on H in Al
  - Present work on H in Ni, Ti, Zr, and Fe



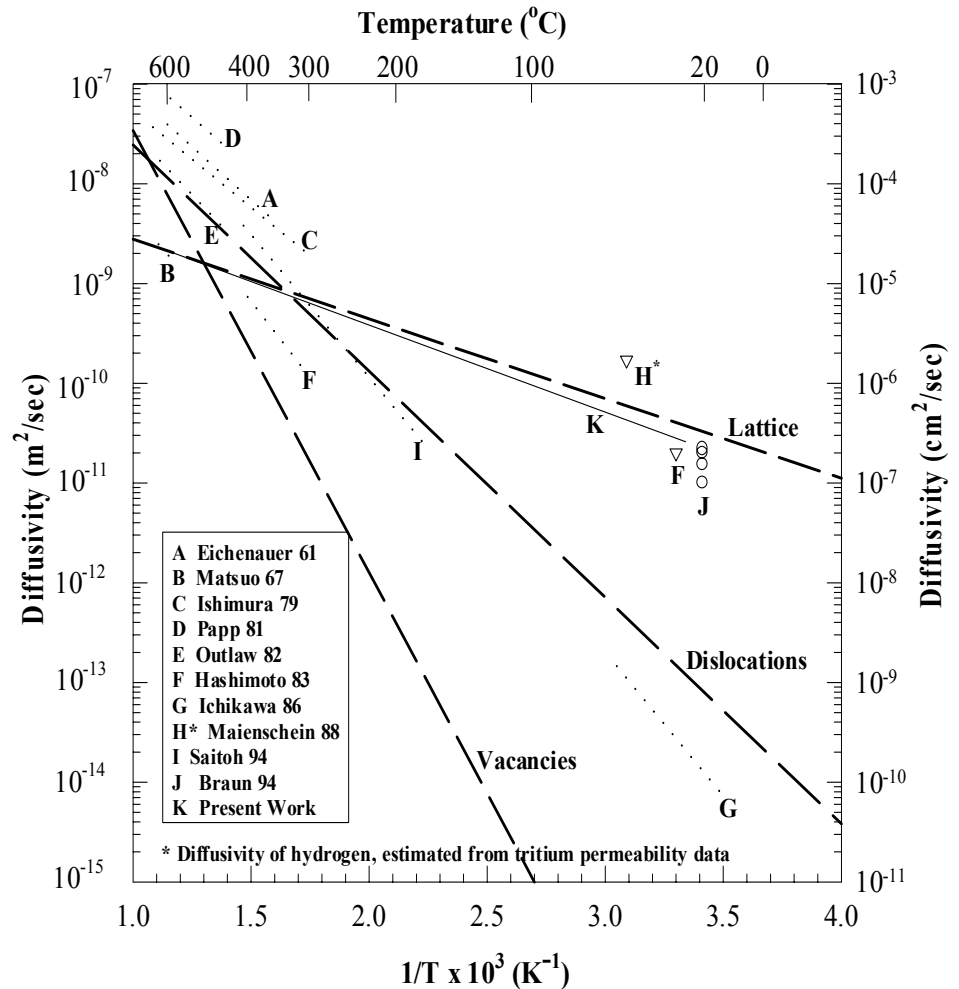
## Experimental difficulties in diffusion

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- Often large noise / small signal
- Oxide or other surface films often diffusion barriers
- Effects of traps often not considered
- May be difficult to ensure “lattice diffusion” control
- Electrochemical methods: charging solutions can degrade sample (pitting)
- Vacuum methods: time lag between charging and measurement
- etc., etc.

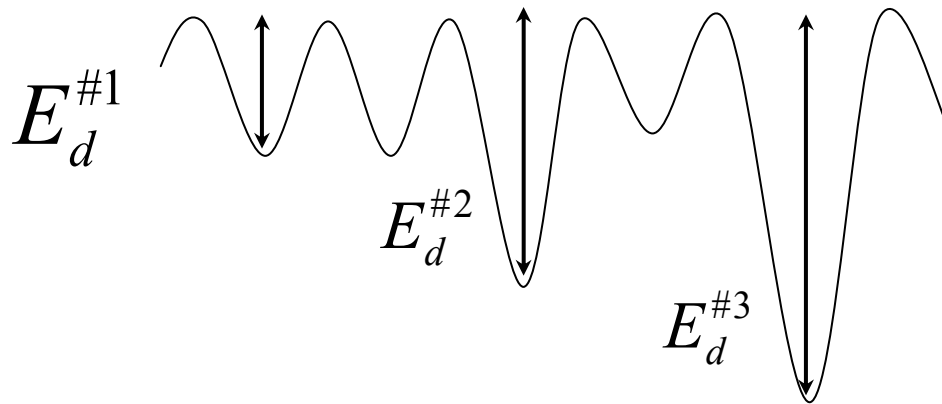
# Hydrogen diffusion in aluminum

- H diffusion both fundamental and technically important to Al and Al alloys
- Large controversy in literature
- See:
  - Young and Scully, *Acta Mat.* Vol. 46, No. 18, pp. 6337-6349
  - Wolverton *et al.*, *Phys. Rev. B*, 69, 144109, 2004



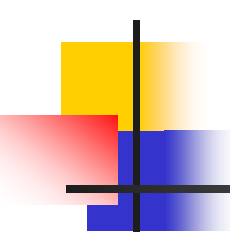


Barrier for lattice diffusion is relatively low,  
vacancies are strong trapping states



### Desorption Energies in Aluminum (kJ/mol)

Researcher	Lattice	Dislocation	Vacancy
Young (experimental)	15.3 $\pm$ 4.8	43.5 $\pm$ 17.5	84.8 $\pm$ 32.2
Wolverton ( <i>ab initio</i> )	17	---	52



Diffusion Equation – see Wert and Zener Phys. Rev. Vol. 76, No. 8, Oct. 15, 1949, pp. 1169-1175.

$$D = \underbrace{n \alpha a^2 \nu \exp\{\Delta S / R\}}_{D_0} \exp\{-\Delta H / RT\}$$

- $n$ : number of nearest neighbor jump sites
- $\alpha$ : numeric coefficient that depends on the location of the interstitial positions
- $a$ : lattice parameter (net jump distance,  $l$ )
- $\nu$ : vibrational frequency
- $\Delta S$ : activation entropy
- $\Delta H$ : activation enthalpy ( $Q$ )
- $R$ : gas constant,  $T$ : temperature



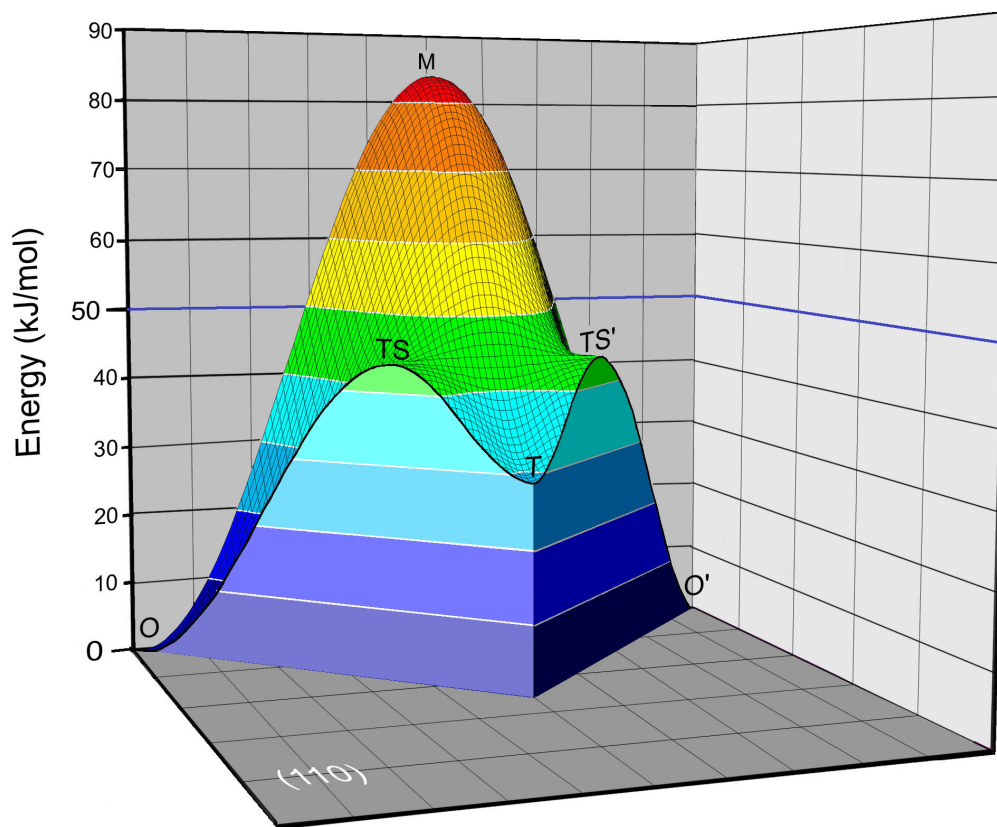
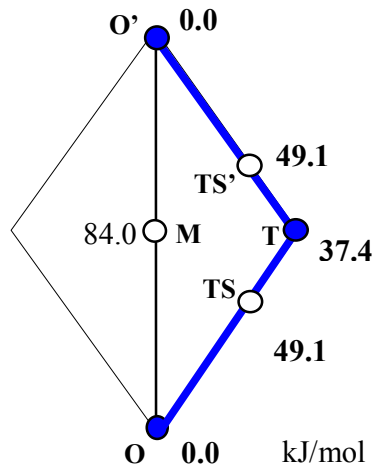
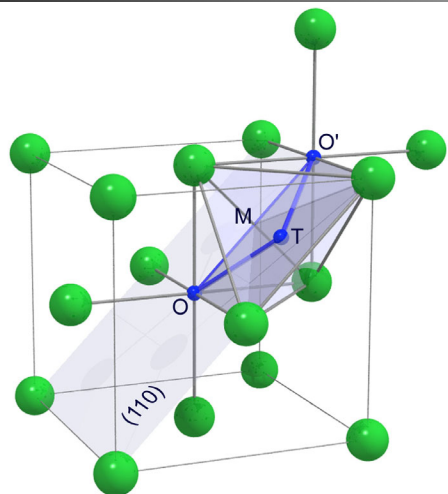
Diffusion procedure: consider temperature dependence of  $\Delta S$  and  $\Delta H$  (and  $\nu$ )

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$$D = n \alpha l^2 \nu \exp\{\Delta S(T) / R\} \exp\{-\Delta H(T) / RT\}$$

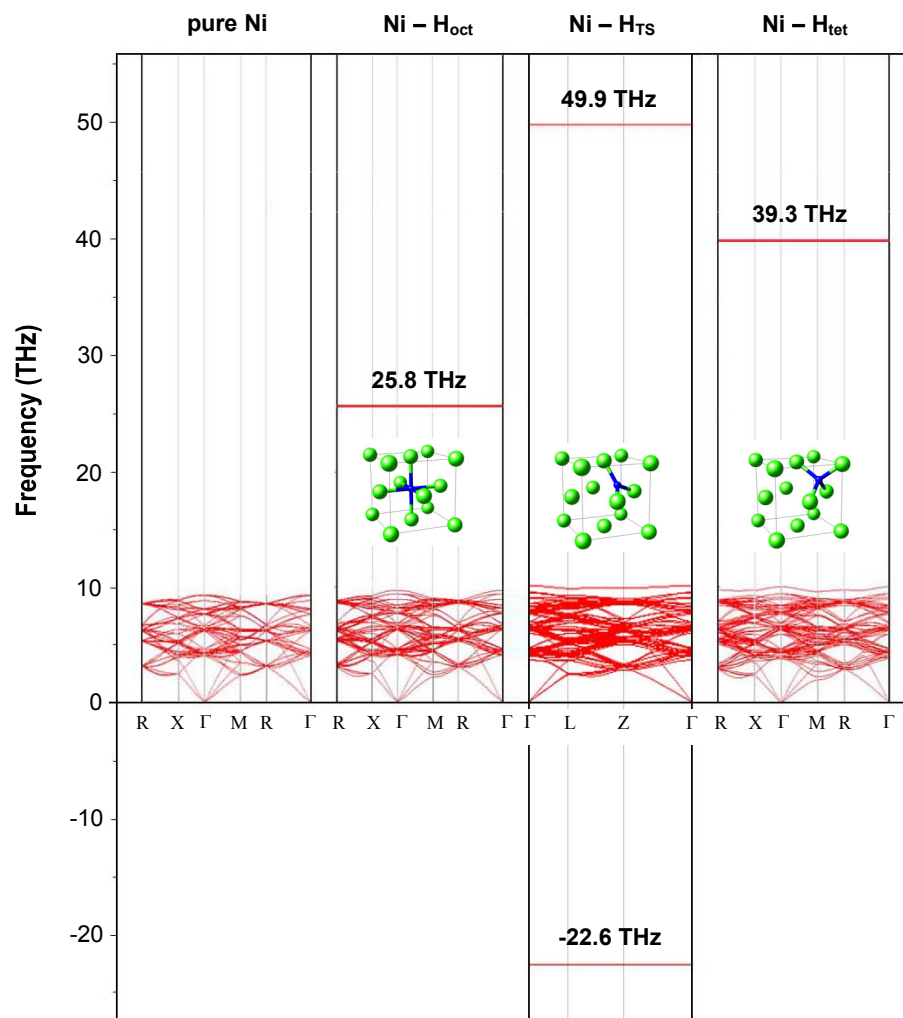
- Determine low energy site
  - Determine low energy path
  - Vibrational frequency (relatively  $T$  independent)
  - Entropy change  $f(T)$
  - Enthalpy change  $f(T)$
- }  $n, \alpha, \text{ and } l$
- }  $\nu$
- }  $\Delta S$
- }  $\Delta H$

# Octahedral site and indirect diffusion path via transition state

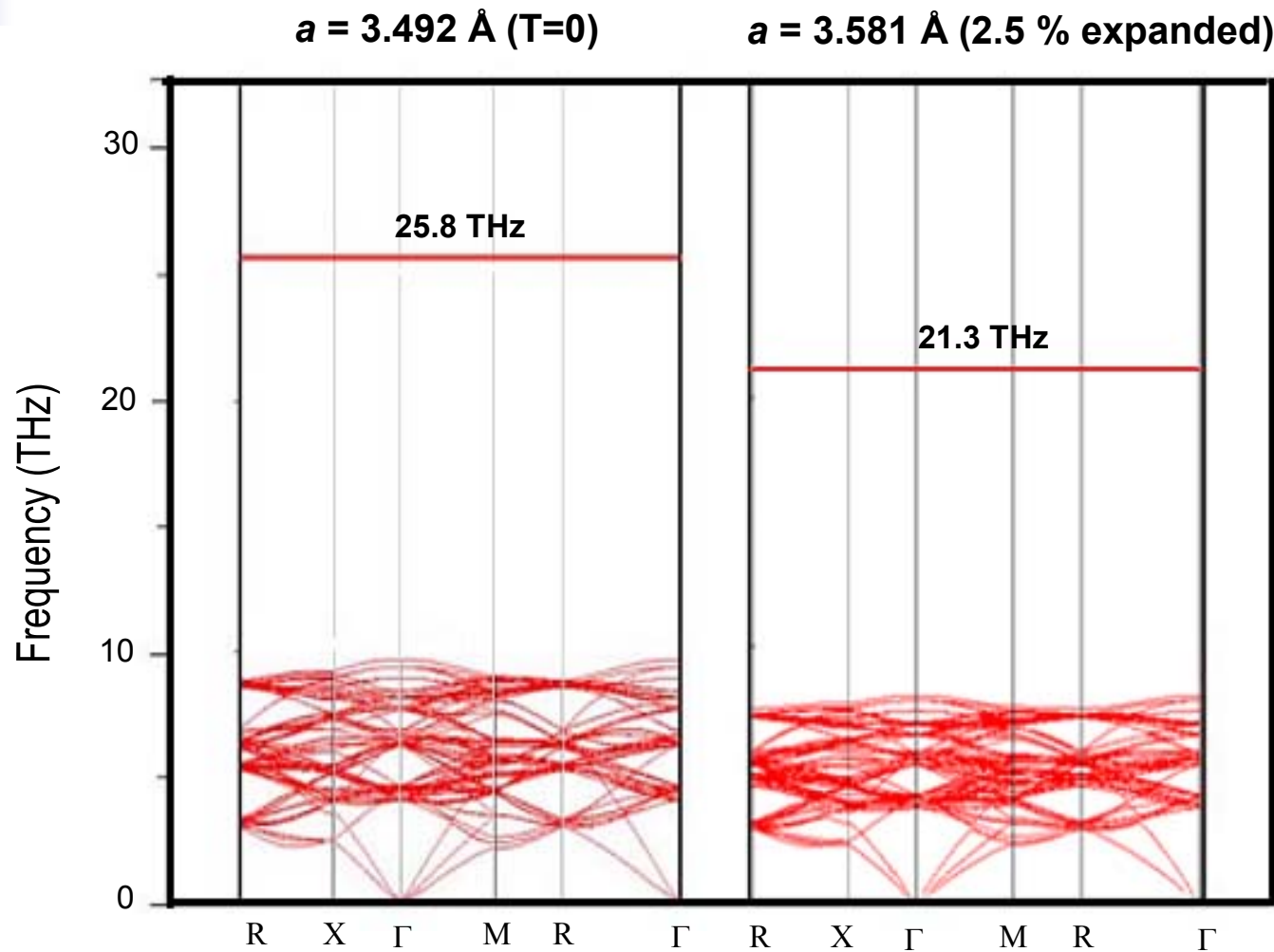


# Vibrational frequency via VASP/Phonon

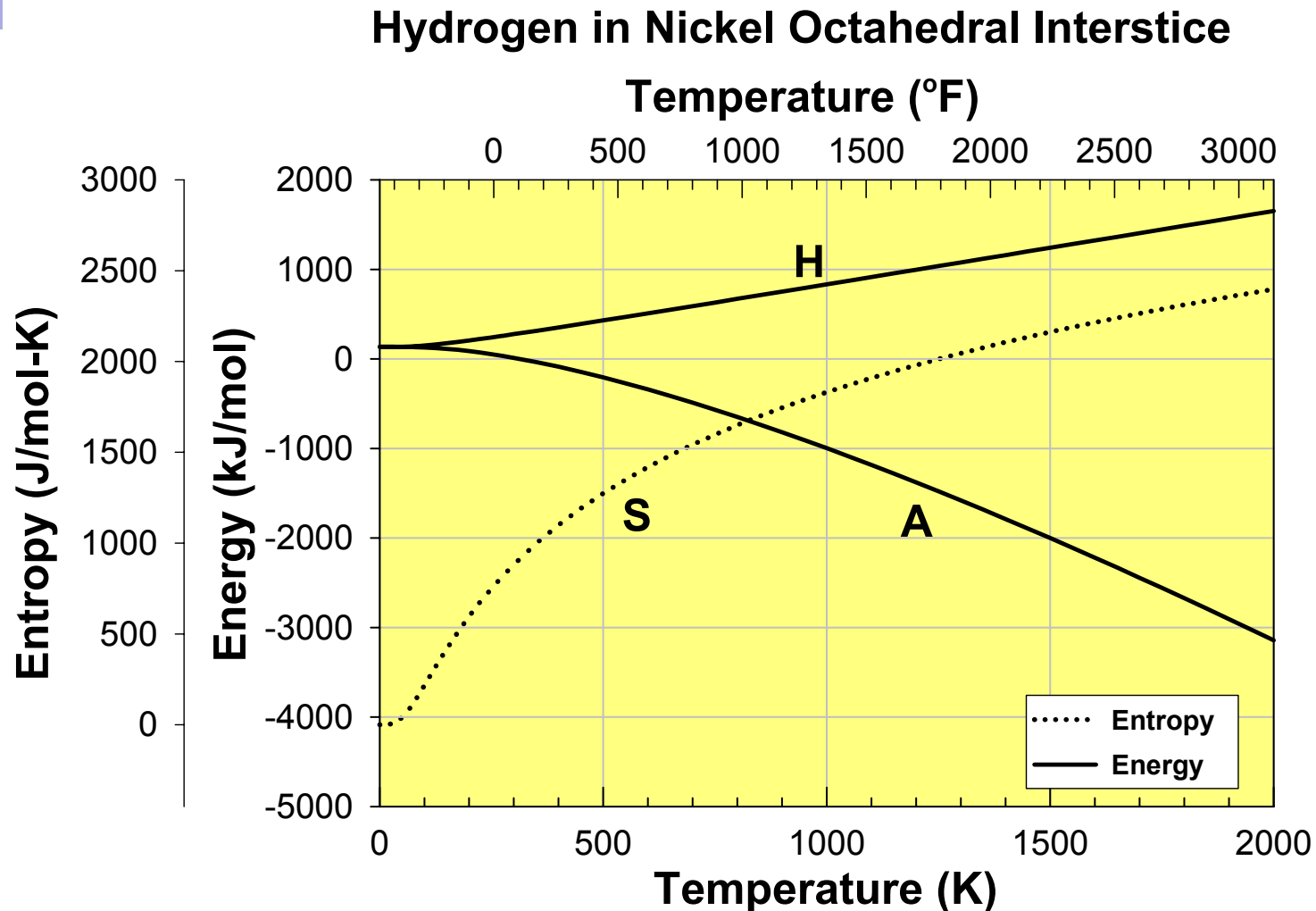
- Frequency at stable octahedral site = 25.8 THz
- Frequency in tetrahedral site = 39.9 THz (less space than octahedral interstice)
- At the transition state, negative eigenvalue  $\rightarrow$  imaginary frequency



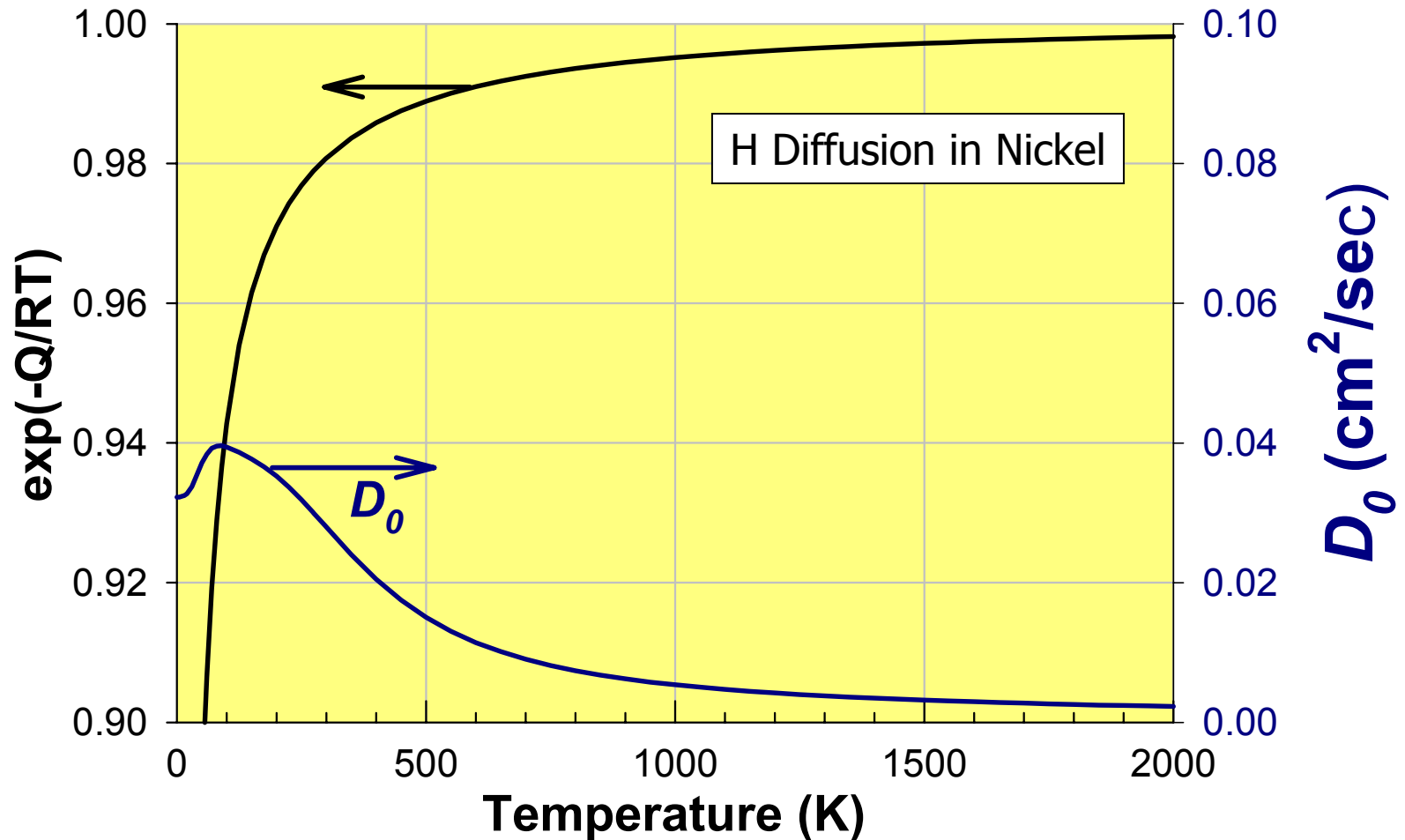
# Temperature has little effect on the vibrational frequency of H in nickel (0 K to $\sim 1000$ K)



$\Delta S$  and  $\Delta H$  vary with temperature. Typical assumption of constant  $D_0$  is an approximation

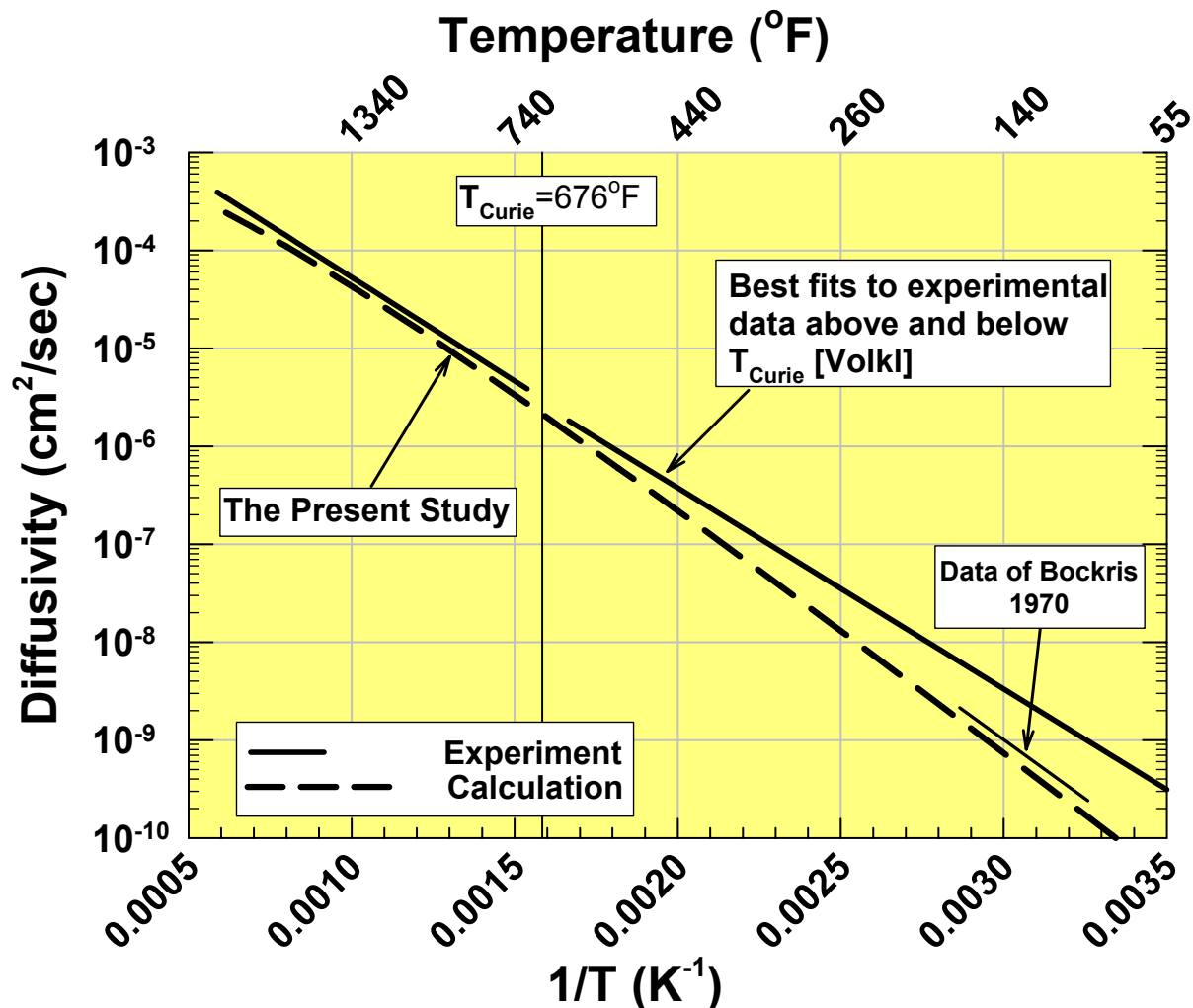


Variation of  $Q$  and  $D_0$  with temperature: note "balance" at temperatures  $>200$  K

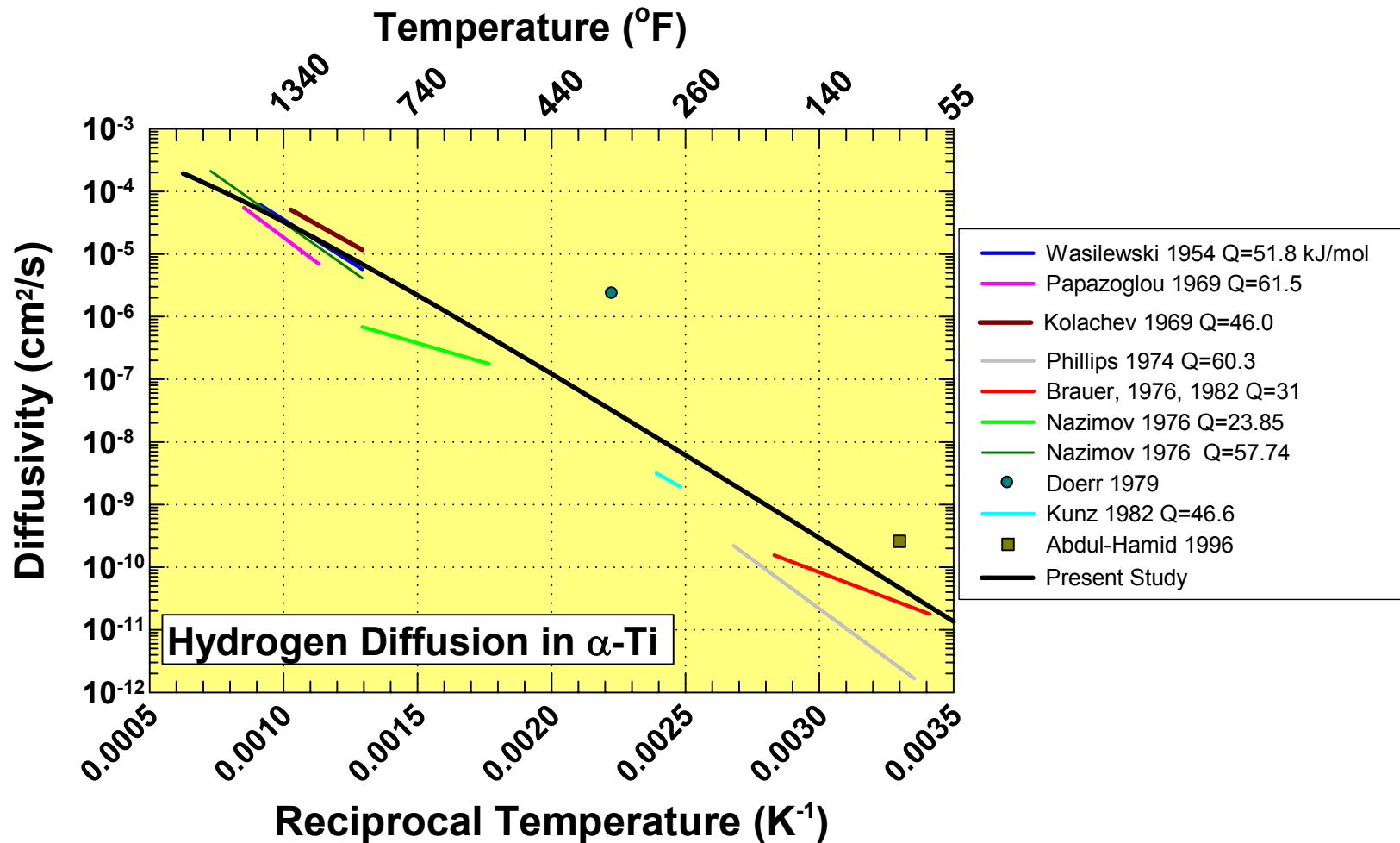




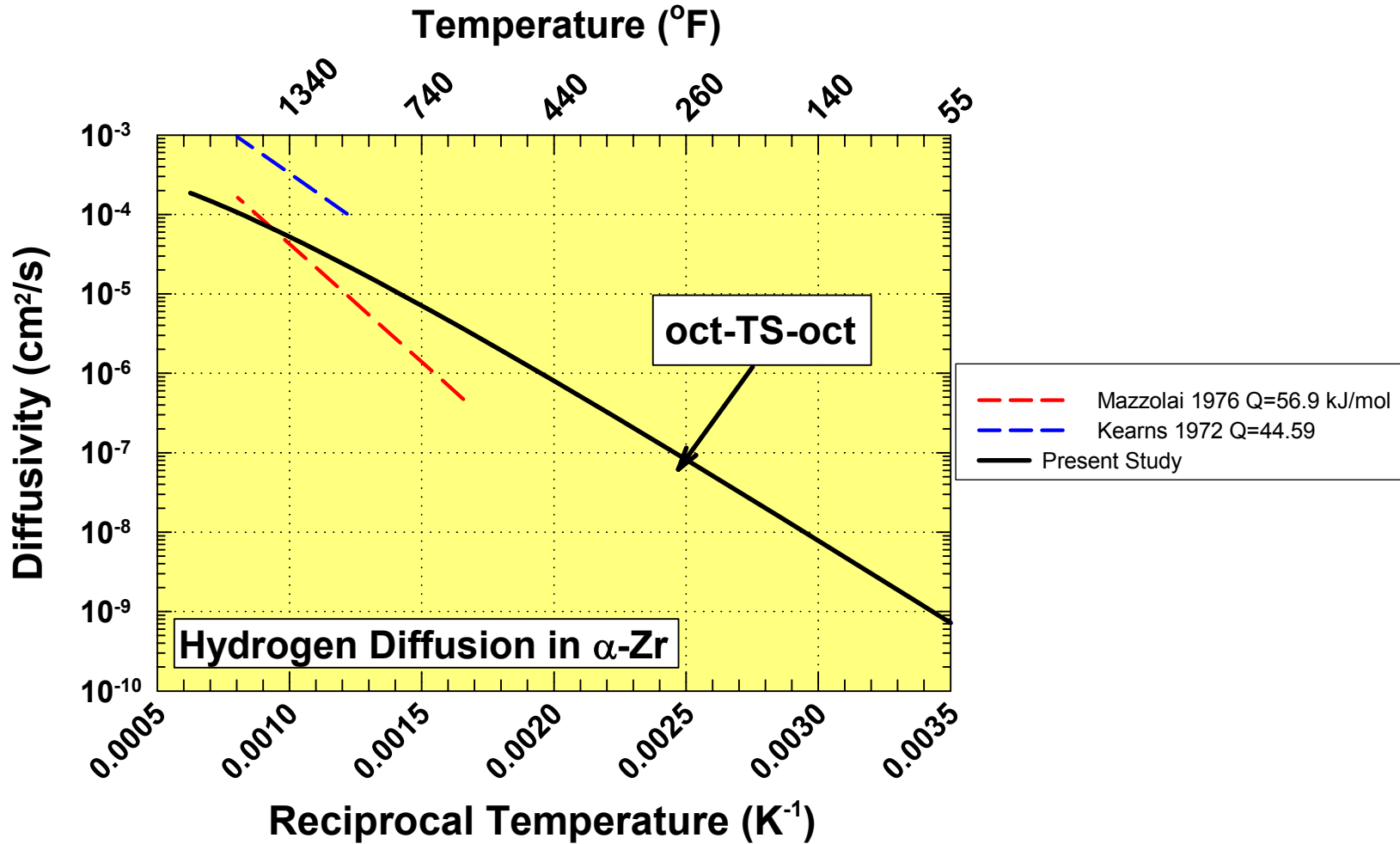
# Excellent agreement between first principles calculations and experimental methods

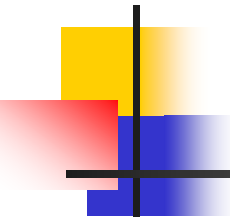


# Large disagreement in experimental data for $\alpha$ -Ti. Modeling line gives an impressive "best fit"



# Limited data for H in $\alpha$ -Zr. Results for H in Ni and $\alpha$ -Ti give confidence that modeling is accurate

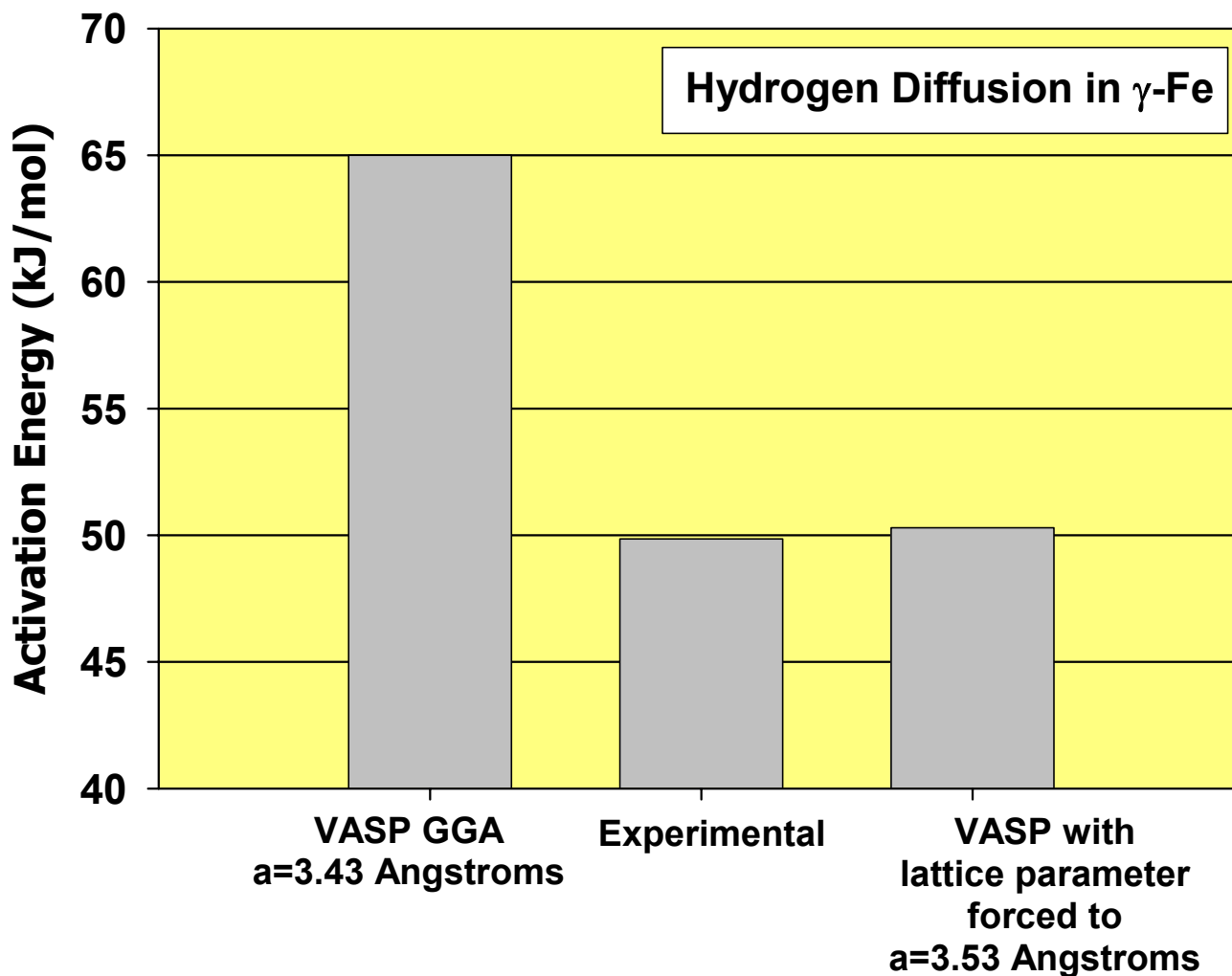




For  $\gamma$ -Fe: Larger than typical error between calculated and experimental lattice parameter

<b>Metal</b>	<b>Computed Lattice Parameters (Å)</b>	<b>Experimental Lattice Parameters (Å)</b>	<b>Deviation (%)</b>
<b>Ni</b>	a=3.492	3.5239	-0.9
<b>Ti</b>	a=2.904 c=4.652	2.950 4.683	-1.6 -0.7
<b>Zr</b>	a=3.213 c=5.210	3.233 5.148	-0.6 +1.2
<b><math>\gamma</math>-Fe</b>	a=3.433	3.6599	<b>-6.2</b>

# Calculations for H diffusion in $\gamma$ -Fe highlight the strong effect of the lattice parameter, $a$



# Summary of parameters

Table I. Summary of Calculated Diffusion Parameters

Metal	Crystal Structure	Interstitial Site	n	$\alpha$	Approximate $l^\#$ (cm)	$\nu$ (THz)		$\Delta S$ (J/mol-K)		$\Delta H$ (kJ/mol)	
						0 K	1000 K	0 K	1000 K	0 K	1000 K
Ni	fcc	octahedral	24	1/12	$2.5 \times 10^{-8}$	25.8	21.3 $^\#$	0	-14.9	49.1	40.2
Ti	hcp	octahedral	$(n\alpha) \approx 1$		$2.3 \times 10^{-8}$	27.7	Not Calculated	0	-8.3	52.7	42.4
Zr	hcp	octahedral			$2.6 \times 10^{-8}$	21.8	Not Calculated	0	-16.0	40.7	30.9
		tetrahedral			$1.7 \times 10^{-8}$	37.1	Not Calculated	0	-5.5	35.5	30.8
Fe*	fcc	octahedral	24	1/12	$2.5 \times 10^{-8}$	28.4	Not Calculated	0	0.2	50.3	43.7

\*Results based on lattice parameter of  $a=3.53$  Å

$l^\# = (a\sqrt{2}/2)$  for the fcc oct.-oct. transition,  $\sim c/2$  for the hcp oct.-oct. transition and  $\sim c/3$  for the hcp tet.-tet. transition

$^\#$ Estimated from expanding the lattice parameter from  $a=3.492$  Å to  $a=3.581$  Å, since the effect of temperature is relatively small; the 0 K attempt frequencies are used in the subsequent calculations.



## Summary

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- Excellent agreement between first principles calculation of H diffusion in nickel
- Calculations help resolve controversy in systems where experimental data are in disagreement:
  - H in  $\alpha$ -Tiand guide predictions where there are sparse data:
  - H in  $\alpha$ -Zr
- Results for H in  $\gamma$ -Fe show *strong influence of lattice parameter*
- Calculations give new insight into diffusion paths and temperature dependencies
- Techniques have broad applicability to Materials Science, *e.g.* diffusion, solubility, entropy, enthalpy, free energy, ...