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MBE Growth of GaInAsSb p/n Junction Diodes for Thermophotovoltaic Applications

by

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Abstract

This paper reports recent progress in the development of quaternary III-V thermophotovoltaic (TPV) devices based on MBE grown $\text{Ga}_x\text{In}_{1-x}\text{As}_y\text{Sb}_{1-y}$. TPV is of great interest for a variety of applications^(1,2). The objective of this work is to develop a TPV cell which is "tunable" to the emission spectrum of a heated blackbody, at temperatures in the range of 1200 - 1473 K. One aspect of this "tuning" is to match the band gap, E_{gap} , of the photovoltaic device to the peak output of the heat source. An advantage of the quaternary III-V semiconductor systems is that devices can be fabricated by molecular beam epitaxy on a suitable binary substrate, such as GaSb or InAs, and the band gap and lattice constant can be adjusted more or less independently, to match requirements. Quaternary cells, with band-gaps in the 0.5 to 0.72 eV range, have been fabricated and tested. For 0.54eV devices we obtained $V_{\text{oc}} = 0.3$ V and $I_{\text{sc}} = 1.5$ amperes/cm² under infrared illumination of a 1200K blackbody. Under high illumination levels the V_{oc} and I_{sc} ranged from 0.5 V at 3 amperes/cm² for 0.72eV devices to 0.31 V at 1.2 amperes/cm² for 0.5eV devices, indicating good photovoltaic device characteristics over the range of bandgaps. The diode ideality factor for 0.54eV devices ranged from 2.45 at low illumination indicating tunneling-dominated dark current, to 1.7 at high illumination intensity indicating recombination-generation dominated dark currents.

Introduction

The elements of a TPV system include an emitter which is a heat source coupled to a blackbody radiator, a spectral control element (band-pass filter) and a photovoltaic converter which is matched to the blackbody and spectral control element. A schematic for a TPV system is shown in Figure 1.

Figure 2 compares the power emitted by a blackbody at various temperatures and compares the power at wavelengths shorter than the cut-off wavelengths for GaSb, and two lower bandgaps. The three lower curves represent the maximum power available for conversion by cells with those bandgaps without taking into consideration any factors limiting their conversion efficiency. It must be pointed out that a high efficiency PV diode is only one aspect of demonstrating an efficient TPV system, an effective spectral

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control device is critical to obtain high overall system efficiency. As an example, for a 1200K blackbody, almost 70% of the blackbody power density lies below the bandgap of a 0.5eV bandgap PV cell. This requires an efficient band-pass filter which reflects most of the radiation below the semiconductor bandgap out to 10- μ m back to the blackbody to recycle the photons and allows maximum transmittance above the semiconductor bandgap for the PV cell to convert.

Among the unique advantages of the quaternary GaInAsSb system are that the bandgap and the lattice constant can be adjusted to match the requirements. Figure 3 shows the alloy composition and the bandgap range of interest lattice matched to GaSb.

In this paper we present recent progress on obtaining large area (1cm x 1cm) variable bandgap GaInAsSb PV cell lattice matched to GaSb substrates with bandgaps between 0.72eV and 0.5eV.

Material Growth

The material used for the fabrication of the GaInAsSb TPV devices was grown using a Varian Gen II MBE system. 2" n-type ($1\text{-}2\text{e}17\text{cm}^{-3}$) exact (001) oriented GaSb substrates were used. The oxide on these wafers was desorbed in-situ in the MBE system. The wafers were heated to about 400 C, then the Sb₄ shutter was opened, the oxide was observed to come off at temperatures in the 525-535 C range as monitored by the pyrometer. An n+ ($\sim 5\text{e}18\text{ cm}^{-3}$) 100nm GaSb buffer layer was grown to smooth the surface and improve the RHEED pattern. During the GaSb growth the substrate temperature was dropped to 500 C and growth of GaInAsSb was initiated. This formed the base region of the TPV device and had a thickness of 5- μ m and an n-type doping density of $2\text{e}17\text{ cm}^{-3}$. Next the p-type emitter and p-type GaSb (300A) window layer were grown with a doping density of $1\text{-}2\text{e}18\text{ cm}^{-3}$. Figure 4 shows a typical device structure.

In order to obtain the correct bandgap, Ga and In ratios were adjusted. To obtain the lattice constant match with the substrate, As and Sb ratios were adjusted. The GaInAsSb devices were lattice matched to within 0.05% with the substrate. When the mismatch exceeded 0.1% the wafer was observed to bend. The wafers bent in a convex shape with a plus (larger than the substrate-excess Sb) mismatch. The wafers were observed to bend in a concave shape with a negative (smaller than the substrate-excess As) mismatch. This bending of wafers was much worse in the case of InAs substrates, because InAs is a softer material than GaSb. To relieve the lattice mismatch strain quaternary GaInAsSb appears to deform the layer instead of relieving strain by the formation of dislocations. This is presumably due to the alloy hardening effect of the second group V in the films. Due to this alloy hardening effect the layer finds it energetically difficult to produce and propagate dislocations and hence bends to relieve the strain. This bending of the wafers took place at the growth temperature and could be observed through the view port, this rules out the bending due to any co-efficient of thermal expansion mismatch. To further investigate the possibility of lattice mismatch related bending peculiar to the GaInSb material system we grew the ternary GaInSb (with the same Ga and In ratio as in the quaternary) on GaSb, we did not observe any bending despite a much larger mismatch ($>0.5\%$). This indicates that in the mixed arsenide-antimonide quaternaries strain relaxation proceeds differently as compared to the ternaries, and can have important device implications. Further investigations using TEM are needed on the GaInAsSb quaternary system.

Device Fabrication and Characterization

1cm x 1cm mesa diodes were fabricated using chemical etching, the mesa diodes had a single 1mm wide central busbar for making electrical contact to the devices. The busbar was connected to grid lines 0.5cm long 100- μ m apart and 10- μ m wide. Front metallization was done by evaporating Ti and Pt in an e-beam evaporator and then depositing a 5- μ m thick Au or Al to obtain a low resistance contact for high current densities. For bottom Ohmic contact a thick Sn/Pt or Sn/Au was used. A 0.2- μ m Silicon Nitride layer on top was used as an antireflection coating.

The diodes were tested under dark and under blackbody illuminated conditions. Figure 5 shows an external quantum efficiency vs wavelength for a typical device. The Figure shows a comparison for a typical GaSb and GaInAsSb device. The fall off at short wavelengths indicates high surface recombination, pointing to the need for an effective front surface passivation or a high bandgap window layer. A softer roll off in quantum efficiency at wavelengths near the band edge indicates the absence of an effective back surface field. Adding these two layers to the devices will considerably improve their performance.

Figure 6 shows comparison of a typical I-V characteristics from a GaInAsSb and GaSb PV diodes under similar low light level illumination conditions. As is to be expected the short circuit current is higher for the lower bandgap device. The I-V characteristics point to the need for better Ohmic contacts and this is an area which is under current investigation. Some of series resistance in the GaInAsSb devices could be due to a heterostructure barrier between GaInAsSb and GaSb⁽³⁾. The conduction and valence band discontinuities between GaSb and GaInAsSb are 0.24eV and -0.1eV respectively. This is particularly troublesome and needs to be reduced by degenerately doping the junction or by bandgap grading.

Figure 7 shows the PV short circuit current density vs the blackbody temperature. It can be observed that for low temperature (1200K) TPV applications the narrow gap (0.54eV) cells are the answer because they offer more than three times the current density as compared to the GaSb (0.73eV) cells.

Figure 8 plots the measured short circuit current density vs the open circuit voltage. Assuming the ideal diode relation

$$I_{sc} = I_0 [\exp(qV_{oc}/nkT) - 1]$$

where I_{sc} is the short circuit current, I_0 is the dark current, V_{oc} is the open circuit voltage, n is the diode ideality factor, k is the Boltzmann constant and T is the temperature.

There are four potential dark current mechanisms: diffusion ($n=1$), bulk recombination/regeneration ($n=2$), surface recombination ($n=2$) and band to band tunneling ($n>2$). In Figure 8 we plotted the I_{sc} vs V_{oc} and fitted the ideal diode equation assuming a single dominant dark current mechanism. We found it useful to compare the data for our MBE grown GaInAsSb and GaSb diodes and JX Crystal LPE grown cells. At low current injection conditions MBE grown devices show a diode ideality factor of 2.43, indicating band to band tunneling, LPE devices showed a combination of diffusion and generation/recombination dark current. This difference could be due to a lower base doping for the LPE grown cells. Lower base doping leads to a lower V_{oc} . MBE cells show a higher V_{oc} at a lower I_{sc} (this is very desirable) as compared to LPE cells. This data leads us to conclude that it is much more desirable to have an undoped intrinsic region between the p and n layers to reduce the tunneling component of the dark

current. Some of the dark current could also be due to a lack of surface passivation. Use of an effective passivant or a high bandgap window layer will also reduce the dark current.

Figure 9 plots V_{oc} vs T for MBE grown GaInAsSb cells. This is done for various I_x values. This data also helps determine the differences in the dominant dark current mechanisms. From low to high illumination levels the V_{oc} decreases linearly with temperature. The difference in the slope dV_{oc}/dT between very low injection levels and higher injection levels indicate that at lower current levels the current is a combination of tunneling and some other mechanism, for higher injection levels once the tunneling barrier has been overcome the dark current mechanism is almost entirely due to tunneling⁽⁶⁾, this is indicated by a relatively small change in the dV_{oc}/dT slope. It will be interesting to measure dV_{oc}/dT for even higher injection levels to see if other dark current mechanisms become dominant.

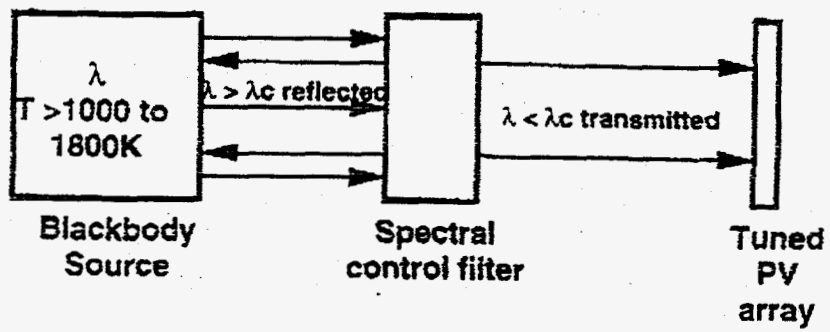
Conclusions

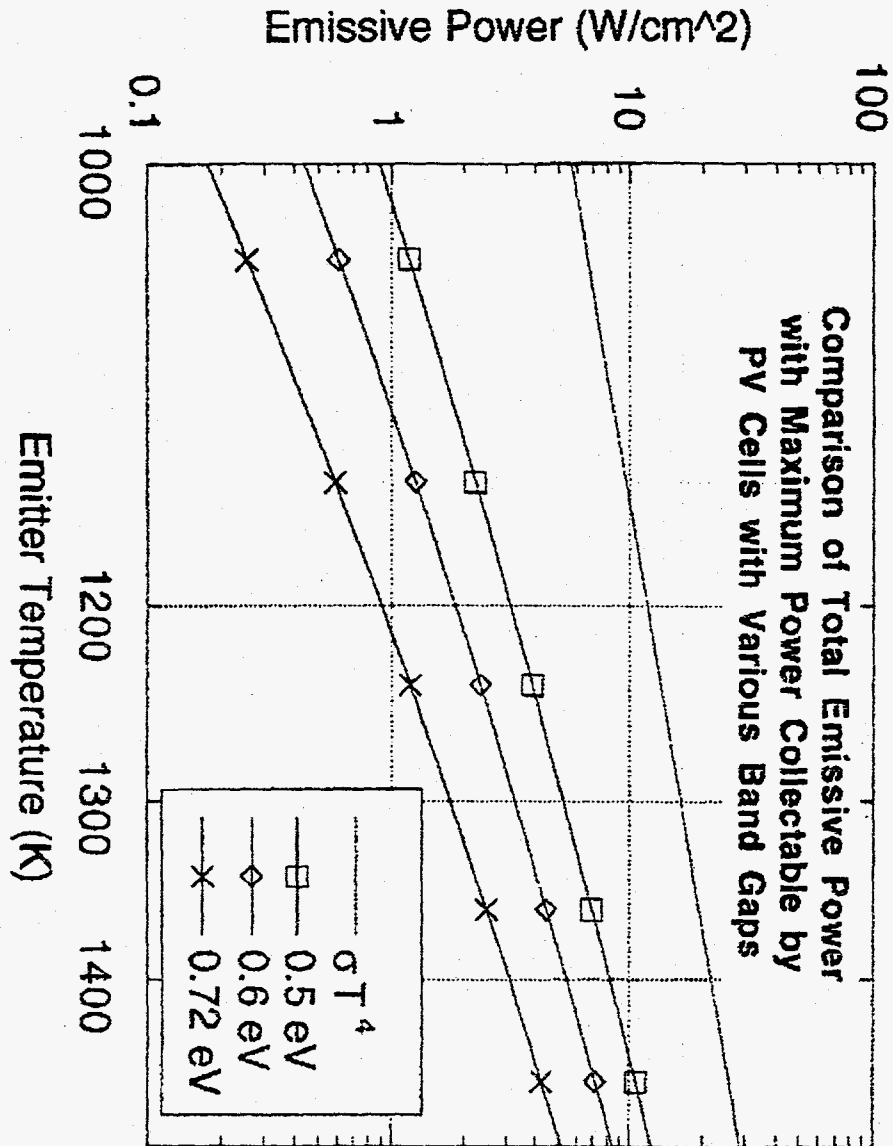
We have presented data on the current status of lattice matched GaInAsSb TPV devices and pointed out the areas in which improvements or device design changes are needed. Quaternary lattice matched TPV devices are very promising for thermal power conversion and offer superior performance due to potentially higher long term reliability and higher open circuit voltages and currents than lattice mismatched InGaAs/InP devices.

References

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TPV based upon Blackbody Emission





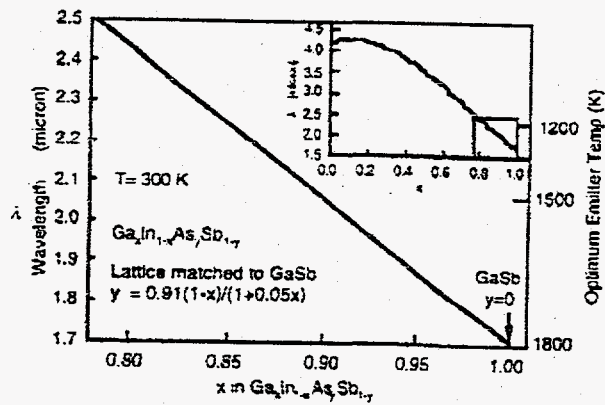
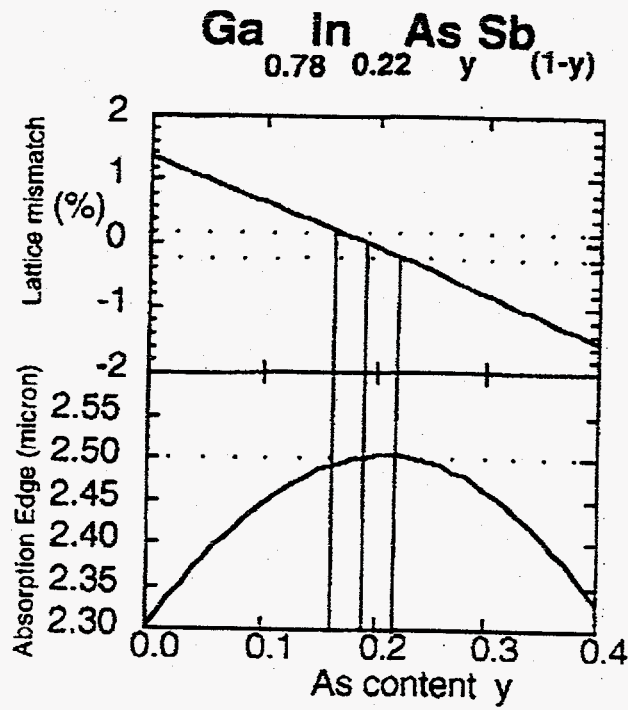
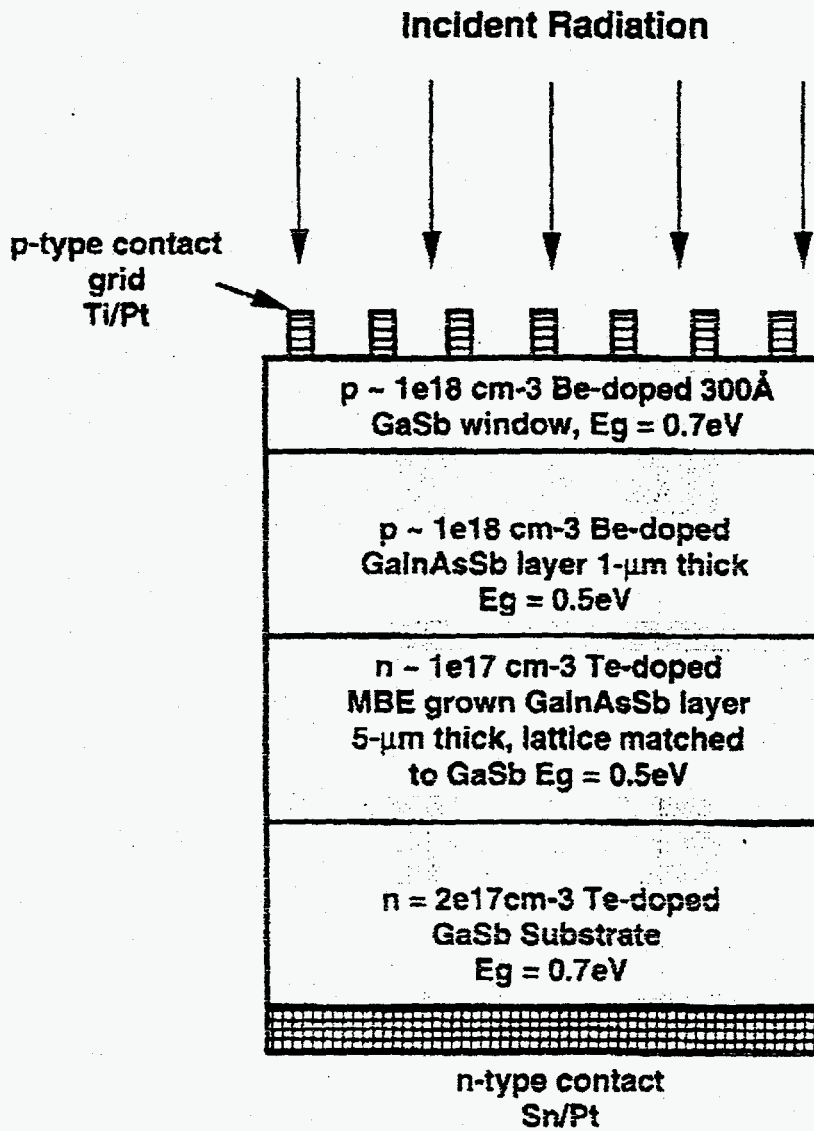
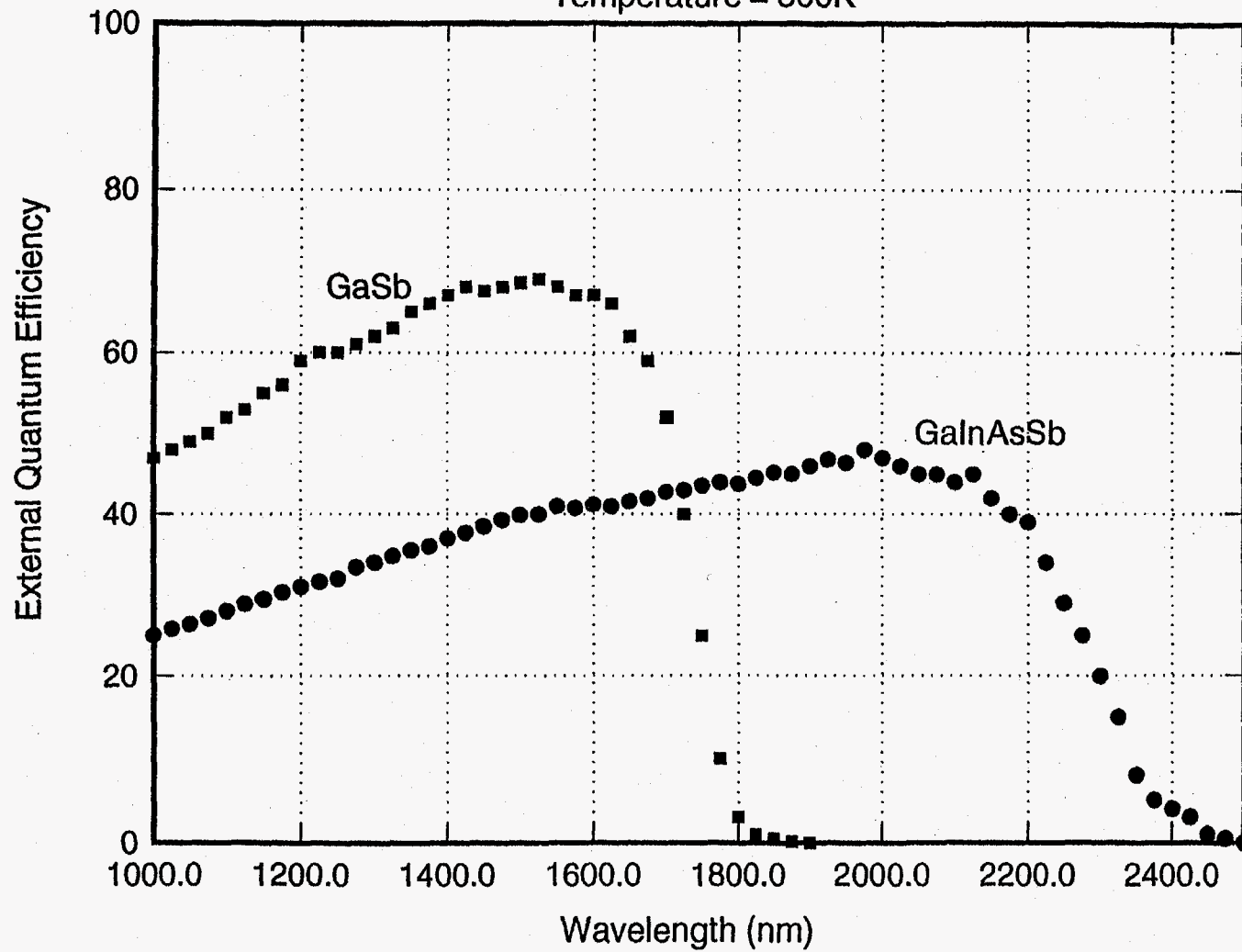


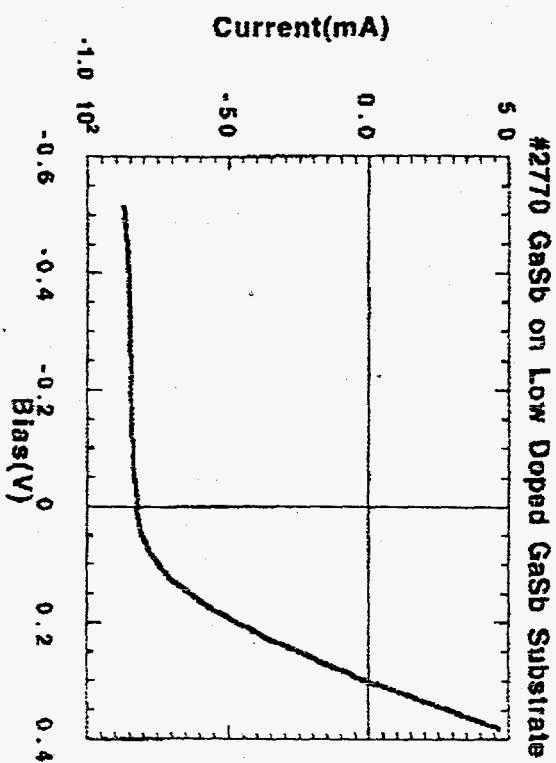
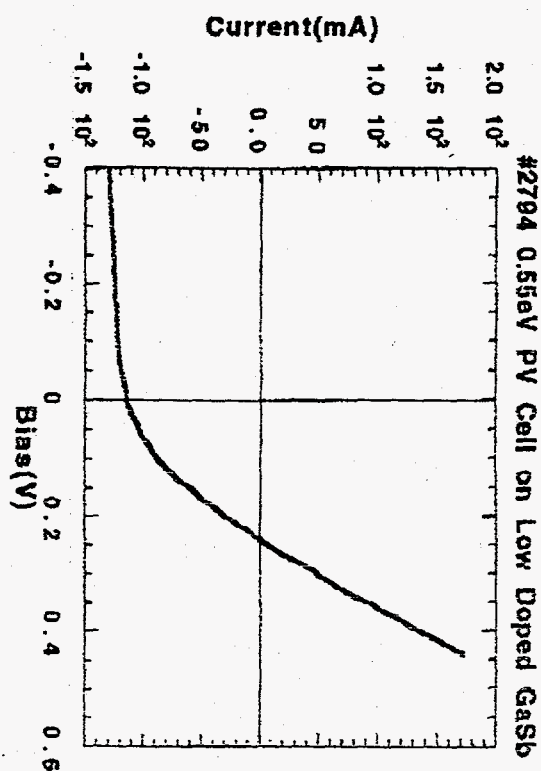
Fig 3



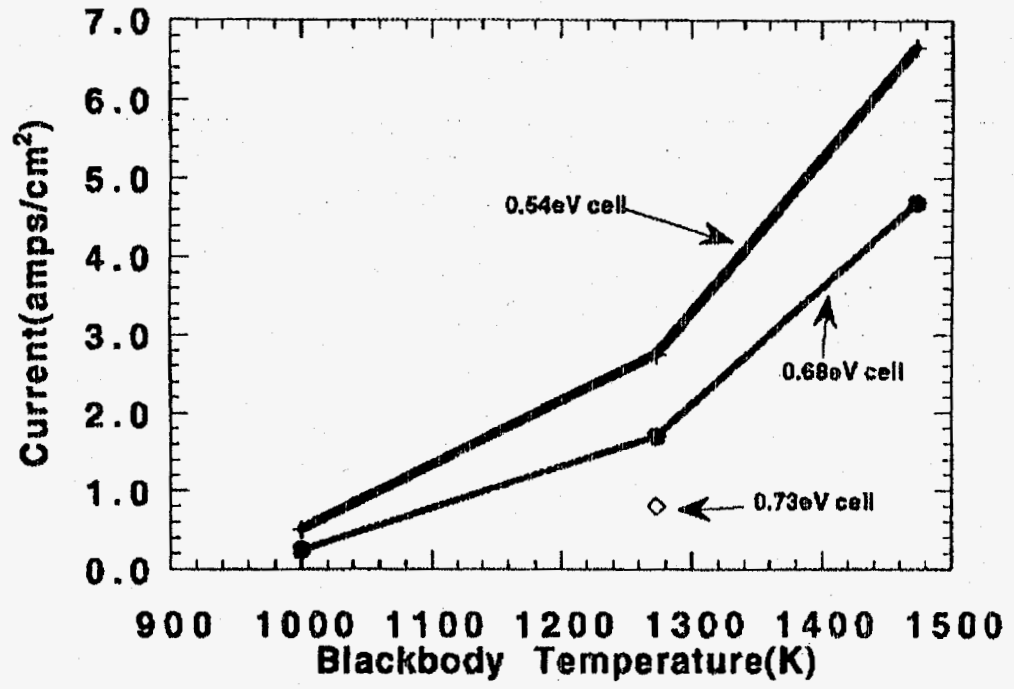
InGaAsSb Quantum Efficiency

Temperature = 300K



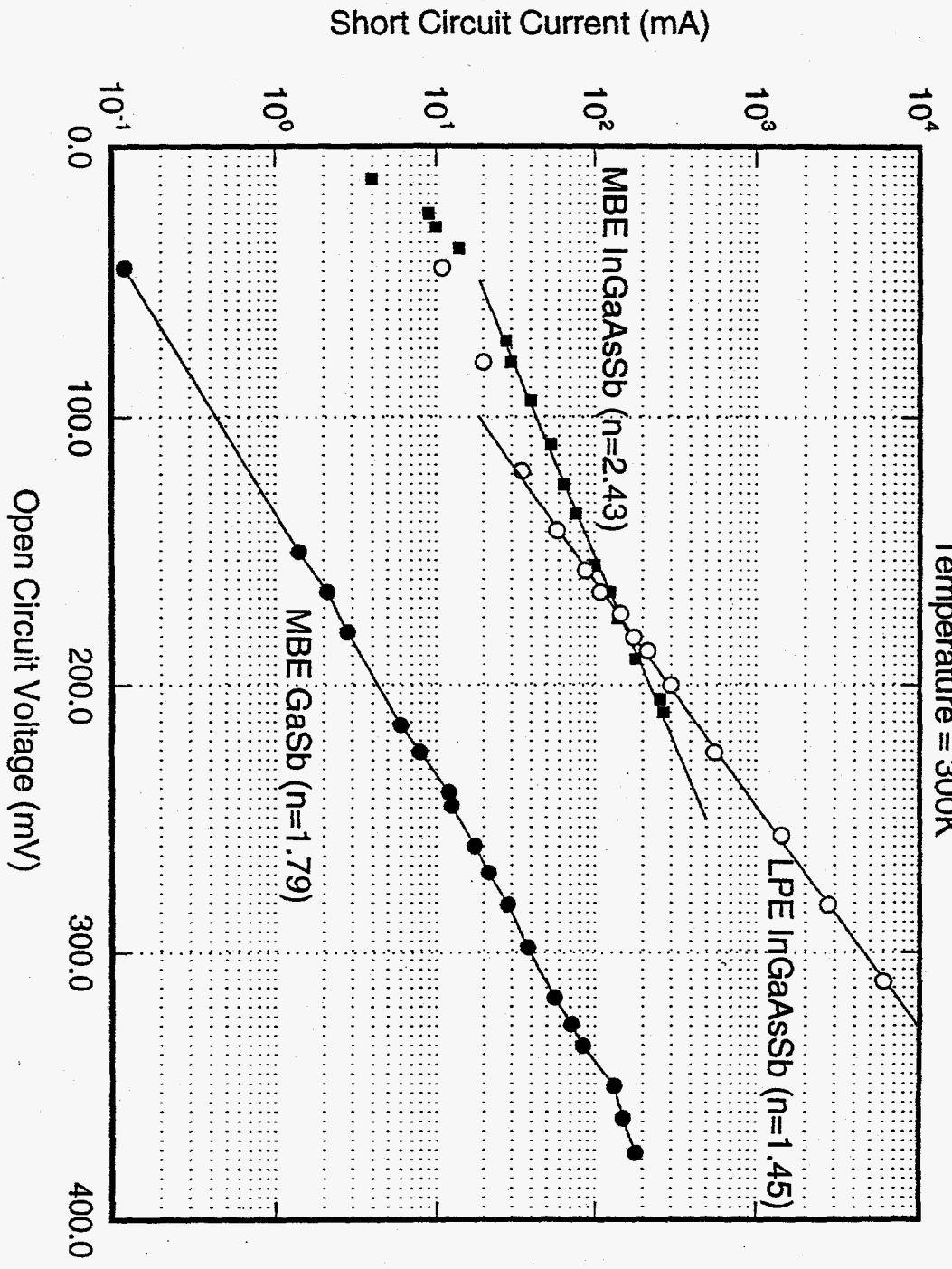


PV Cell Current Density vs Blackbody Temperature



InGaAsSb V_{oc} vs. I_{sc}

Temperature = 300K



InGaAsSb V_{oc} Temperature Dependence

