## IMPROVEMENTS IN THE CHERS SYSTEM FOR DT EXPERIMENTS ON TETR

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## Improvements in the CHERS system for DT Experiments on TFTR

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#### **Abstract**

Improvements in the charge exchange recombination spectroscopy (CHERS) system have resulted in accurate measurements of  $T_i$  and  $V_\varphi$  profiles during DT experiments. These include moving the spectrometer detector array and electronics farther away from the tokamak to a low neutron flux location. This relocation has also improved access to all components of the system. Also, a non-plasma-viewing calibration fiber system was added to monitor the change in fiber transmission due to the high flux DT neutrons. Narrowband filtered light transmitted through the calibration fiber is now used as a reference for the  $V_\varphi$  measurement. At the highest neutron flux of  $\sim 2.5 \times 10^{18}$  neutrons/sec (fusion power  $\sim 6.2$  MW) a modest 5% decrease in fiber transmission was observed. Corrections for transmission loss are made and  $T_i$  (r,t) and absolute  $V_\varphi$  (r,t) profiles are automatically calculated within four minutes of every shot.

#### I. Introduction

Charge exchange recombination spectroscopy (CHERS) [1] is being used to measure ion temperature,  $T_i$  (r,t), and plasma toroidal velocity,  $V_{\varphi}$  (r,t), profiles during historic high power deuterium/tritium (DT) experiments now underway on TFTR. The measurements make use of the neutral beams used for plasma heating in order to obtain time and spatially resolved  $V_{\varphi}$  and  $T_i$  data. High concentrations of tritium have been used [2]. The injected power can consist of from all D° beams to

all T° beams. Although the tritium concentrations can be high, many of the interactions and physics phenomena can be fairly subtle. For example, detection of alphas and their effect on the plasma (alpha heating, TAE modes, etc.) or beam and thermal ion species effects on confinement and transport, and changes in velocity associated with supershot to H-mode transition [3] represent important physics issues; however, it may be difficult to make quantitative measurements of these effects. Accurate spatial and temporal determination of  $V_{\varphi}$  and  $T_i$  using the charge exchange recombination spectroscopy technique is essential in observing and quantifying these effects.

Basic features and characteristics of the CHERS system were described in detail in earlier reports [4]. However, several modifications and changes have been made which allow new and improved measurement of  $T_i$  (r,t),  $V_{\varphi}$  (r,t), and  $n_c$  (r,t) during reactor relevant DT operation. These include moving the system a long distance away from the tokamak to a location behind a concrete wall, and addition of a calibration fiber run parallel to the measurement which allows determination of the effect of neutrons on the fiber transmission [5] and a bonus capability of having a stable zero reference for the toroidal velocity measurement. A shakedown of the electronics was undertaken, and several sources of noise were eliminated, significantly improving the overall dynamic range of the system. Finally, the capability of fast high precision automatic analysis is now available and time resolved  $V_{\varphi}$  and  $T_i$  profiles with fiber transmission and  $V_{\varphi}$  reference taken into account can be generated within four minutes after every shot.

#### II. Description of Apparatus

A schematic drawing of the CHERS system including the calibration fiber and diagnostic fiber bundles are shown in Fig. 1 as they are located relative to TFTR. The main components of the system consist of the fiber optic bundles for viewing the plasma and transmitting the collected light to the spectrometer. At present, there are two fiber bundles, one coming from a toroidal location at Bay H and the other 180° away toroidally at the Bay P location. The fibers are 600 µm diameter strands of pure fused silica with plastic cladding. The two bundles have double passes of calibration fiber installed directly beside them so that they sample the same neutron flux and electromagnetic environment as they traverse the distance from the viewing window to the entrance of the spectrometer. The calibration fibers begin at a position near the spectrometer where they are illuminated by a stable filtered light source; a meter away, they are placed with the two diagnostic fiber bundles. The two bundles are enclosed in two separate and opaque one inch I.D. Corelok tubes (for protection) for the full length from viewing window to spectrometer.

A tungsten lamp is the light source for the calibration fiber system. A narrow band pass optical filter which transmits at a center wavelength of 5292Å with bandwidth ~ 4.3 Å is placed between the lamp and the inlet end of the calibration fibers. The filter is placed within a well controlled oven which maintains it at a constant temperature (compensating for changes in atmospheric temperature which affect the spectral transmission of the lens). The exit ends of the calibration fibers are placed in the same vertical plane as their respective diagnostic fiber bundles. A fast switching mirror directs the light from the selected array into the spectrometer. The mirror can be programmed to switch from one bundle (Bay H) to the other (Bay P) to obtain CHERS data at two toroidal positions. A co-rotating plasma moves toward the fiber array viewing the plasma at Bay P and away from the fiber array viewing Bay H. During normal operation, the mirror is usually fixed to view the Bay H array since it provides the widest radial coverage of the plasma. The images of the diagnostic fiber array along with that of the illuminated calibration fiber are directed by the fast

switching mirror onto the entrance slit of a 0.6 m, f/5.7 Czerny-Turner spectrometer. At the output image plane of the spectrometer, the spectrally resolved light is imaged onto a two-dimensional intensified photodiode array (128 × 128 pixel, Reticon MC9128), thus providing the spatial resolution. The system stores up to 64 frames of data during a plasma shot. Typically, the frames are taken every 20 ms for 1 sec NBI heating pulses. Integration times from 10 to 100 ms have been used depending on the heating pulse length and the allowable time resolution.

#### III. Application and Results

The spectrometer itself, the photon/neutron sensitive diode array, and all electronics are now in a remote location where the neutron flux is reduced. The signals from the spectrometer are then transmitted to the control room where they are recorded and archived. In the original CHERS system, all of these components were located directly below the tokamak (about 5 m from the vacuum chamber) in the TFTR Test Cell Basement. Longer fiber bundles are now used so that the CHERS system is located ~25 meters away. The relocation of the CHERS system to a more well shielded remote location reduced the absorbed neutron dose-equivalent by greater than a factor of 70 [6]. This means significantly reduced neutron noise during a shot and reduced accumulative radiation damage to the detectors and electronics.

Another advantage of the new location is that the access to the system components, electronics, spectrometer, etc., is almost unlimited. In the old location, access was very restricted. Now under actual DT plasma operating conditions, debugging, optimization, and testing can be carried out by personnel right at the CHERS system location. When the spectrometer gain is increased for low plasma signal, the calibration fiber light may be too bright and saturate the

detector. The light to the calibration fibers is adjusted accordingly by inserting neutral density filters in front of the lamp. Presently, during DD or DT operation, the filters are inserted manually.

The image or spectral/spatial data for one frame for the Bay H fiber array is shown in Fig. 2. The image for the calibration fiber is at the bottom of the photograph. Data for the 18 diagnostic fibers are clearly distributed vertically upward in the photograph. The fibers view the plasma through a quartz window at the midplane with each fiber aimed so as to view at 18 different tangency radii or about every 6 cm (from a minimum tangency radius of 2.53 m to a maximum of 3.51m). The Shafranov shifted plasma center (magnetic center) is typically at  $\sim$ 2.70 m. Figure 2 is a very good visual image of most of the imformation contained in the CHERS raw data. The vertical position represents the plasma radial position viewed, the brighter colors represent the higher light intensities, the horizontal shift of an elliptical pattern relative to the reference fiber image is indicative of  $V_{\phi}$  and the horizontal width of each ellipse represents  $T_i$ . The plasma light emission is well fitted by a single Gaussian. The Doppler broadening and spectral shift are largest for the fiber at Y-pixel 100, which happens to view the center of the plasma. Thus for Fig. 2, the highest  $\,T_{i}$  and  $V_{\varphi}$  are at the center of the plasma. A shift to the left of the reference fiber image indicates counter toroidal rotation. Charge exchange recombination emission from the C<sup>5+</sup> 5292Å line (n = 7 - 8) due to the interaction of the heating beam with the plasma is normally used for the CHERS measurement on TFTR.

Time resolved calibration fiber transmission data is shown in Fig. 3 for two of the highest integrated DT neutron fluxes achieved on TFTR. Also shown are the respective neutron source rates. In the first case with a 1 sec NBI pulse, the DT neutron rate is high, with a peak at  $\sim 2.5 \times 10^{18}$  neutrons/sec. For the second case, the peak value is less than half this but the beam pulse is about twice as

long. The fluctuations in the transmission data, the source of which has yet to be determined, are the same independent of whether the plasma is DD or DT. As can be seen from the figure, the total decreases in fiber transmission due to the total integrated fluxes are about the same for the two cases, ~ 5%. The peak neutron rate for the high flux example corresponds to a fusion power of 6.2 MW. For longer pulse length and/or high fusiom power, proportionally, greater transmission losses will result.

The reduction in the transmission of the diagnostics fibers should be exactly the same as that for the calibration fiber. Since there is a double loop of calibration fiber, the transmission for a single strand and thus, the single length of diagnostic fiber is  $T \propto \sqrt{I}$ , where I is the lamp light intensity. For each time or nth frame, the signal in the diagnostic fiber is  $S_n' = S_n/T_n$  where  $T_n$  is the transmission,  $S_n$  is the apparent signal and  $S_n'$  is the actual transmission corrected signal for the  $n^{th}$  time (frame). Error analysis for the transmission data is done at the same time and in the same manner as the  $V_{\phi}$  and  $T_i$  data [7]. The actual standard deviation is given by

$$\sigma_{S_{n}} = \sqrt{S_{n}^{'2} \left( \frac{\sigma_{S_{n}}^{2}}{S_{n}^{2}} + \frac{\sigma_{T_{n}}^{2}}{T_{n}^{2}} \right)}$$

where  $\sigma_{S_n}$  and  $\sigma_{T_n}$  are the apparent standard variations of the signal and transmission, respectively.

Corrected  $T_i$  and  $V_{\varphi}$  profiles which make use of the calibration fiber transmission and velocity reference information are plotted in Fig. 4 along with values using an edge fiber as the reference. Alternatives to the calibration fiber as  $V_{\varphi}$  reference are to use 5292 A line emission from the pre-beam ohmic plasma or for a fiber viewing the plasma edge, where  $V_{\varphi}$  is expected to be small, during beam

heating. However, depending on the plasma radius, the hard mounted fiber array may not have a fiber viewing directly at the edge (the exact location of the edge is uncertain in any case) and the light emission from the ohmic plasma is weak and noisy. In Fig. 4 the reference fiber viewing the plasma edge at 3.41 m was chosen as the reference and the resulting data are represented by the dashed curve. The solid curve is the  $V_{\varphi}$  profile using the calibration fiber as the reference and are absolute rather than relative velocities. There is a significant displacement between the two curves. There are other corrections to  $V_{\varphi}$  which may be more important, such as for variations in the background light emission [8] and the velocity dependence of the charge exchange recombination rate. The 5% decrease in fiber transmission for the plasma of Fig. 3 was found to result in a maximum correction at  $T_i \sim 27$  keV of  $\sim 0.5$  keV or 2%. Similarly, as for  $T_i$ , corrected profiles for  $V_{\varphi}$  and the plasma carbon light intensity are also available.

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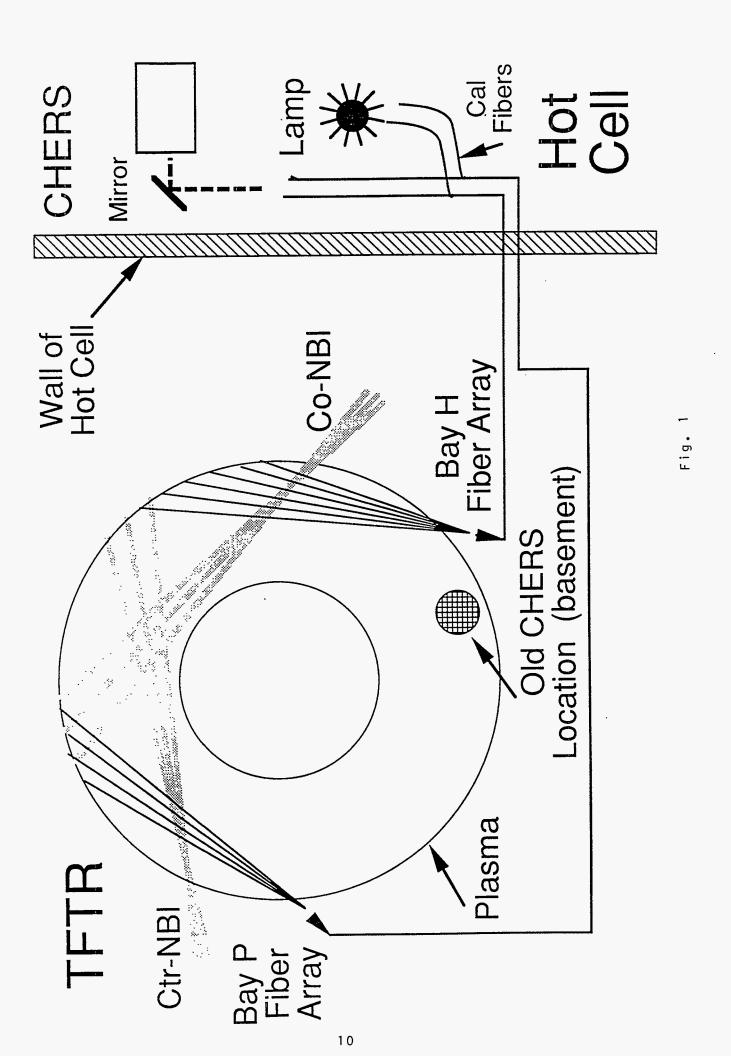
#### Figure Captions '

- FIG. 1. Schematic drawing showing the CHERS system in its new remote location.

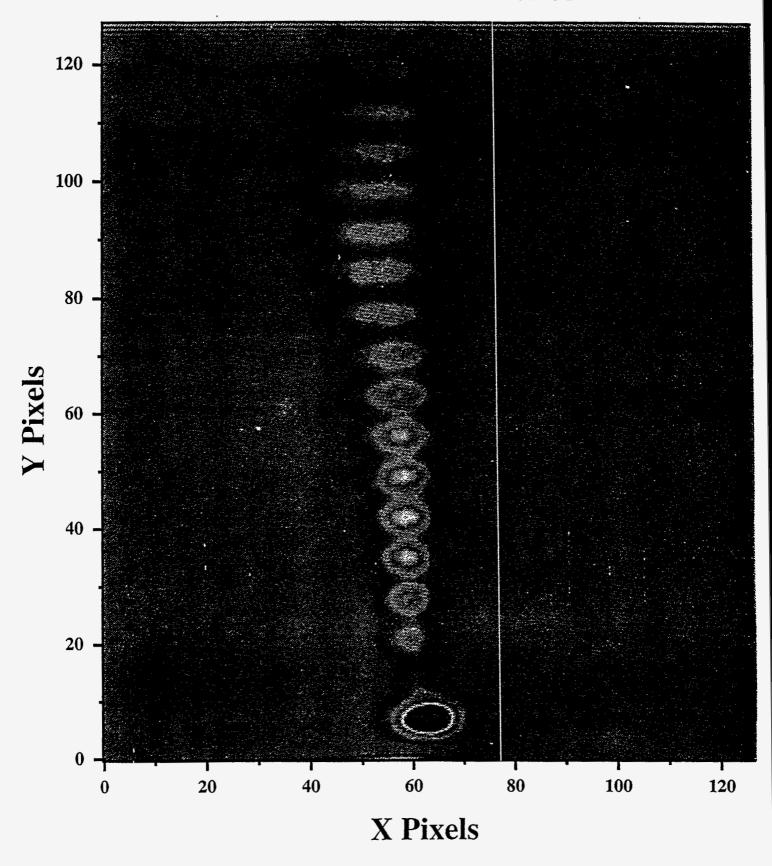
  The two fiber array bundles and their sightlines as they intersect heating beam paths are indicated.
- FIG. 2. Image in fiber position vs. wavelength space of the output from the 2-D intensified photodiode array located at the exit end of the Czerny Turner spectrometer. The calibration fiber is at the bottom and the 18 diagnostic fibers are distributed vertically upward on the screen.
- FIG. 3. (a) DT neutron emission for two different plasma heating pulses, a short how power pulse ( $P_{fusion} \sim 6.2 \text{ MW}$ ) and a long medium pulse (b) give evolution of the fiber transmission for the two cases.
- FIG. 4. T<sub>i</sub> profile at t = 3.75 s for the high power / high neutron rate DT plasma of Fig. 2.

  The correction in the peak Ti (~27 keV) for a 5% decrease in fiber transmission is

  500 ev; extrapolation to 10 and 20 % decreases in fiber transmission indicate
  corrections of 1 and 2 keV respectively.



# CHERS $V_{\phi}$ and $T_{i}$ Data for TFTR Shot # 75963



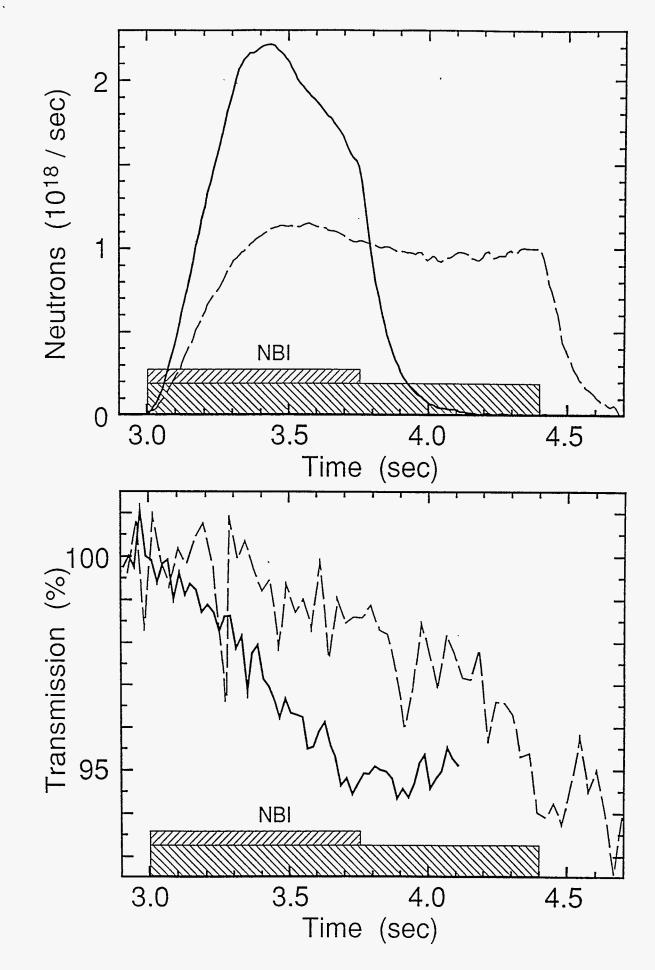
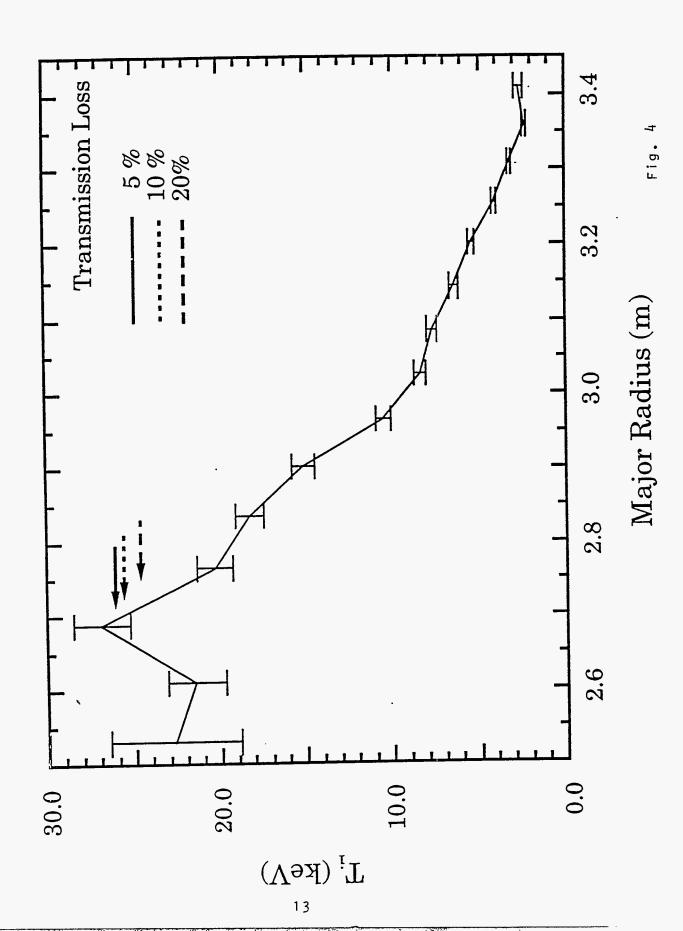


Fig. 3



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