THE CRYSTAL AND MOLECULAR STRUCTURE OF 2,2'BIPYRIDYLGLYCINATOCHLORO COPPER(II) DIHYDRATE

THESIS

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By

Conrad J. Neitzel, B. A. Denton, Texas May, 1975 Neitzel, Conrad J., <u>The Crystal and Molecular Structure</u> of <u>2,2'-bipyridylglycinatochloro Copper(II)</u> <u>Dihydrate</u>. Master of Science (Chemistry), May, 1975, 54 pp., 11 tables, 6 figures, bibliography, 36 titles.

The three-dimensional x-ray structure of 2,2'-bipyridylglycinatochloro copper(II) dihydrate has been fully refined to a final R factor of 0.081. The bipyridyl and glycine ligands are arranged about the central copper atom in a square planar configuration while the chlorine atom is 2.635 angstroms above this plane directly over the copper atom. This unusually long distance is explained by the positioning of a glycine group on the opposite side of the square plane, resulting in a distorted octahedral arrangement. Also, the chlorine atom is linked to three oxygen atoms via hydrogen bonding, thus stabilizing the distorted octahedral complex.

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CHAPTER I

INTRODUCTION

It has been known for some time (8) that $3d^9$ copper(II) octahedral and square planar complexes undergo considerable Jahn-Teller (9) distortions. These distortions produce respectively distorted octahedral and distorted tetrahedral structures (1). Dutta and associates observed that the magnetic moments of several $3d^9$ copper(II) heterochelates in aqueous solution are approximately 1.90 Bohr Magnetons (5), indicating the presence of an unpaired electron (4, p. 537).

2,2'-Bipyridiylglycinatochloro copper(II) dihydrate (BPGC) absorbs at 16,940(5)* to 15,880(5) cm⁻¹ in the ultraviolet region (5, p. 354). For a system with one free electron in its d orbitals, transitions in the ultraviolet region are expected to peak out at a wave-number of approximately 10,000 (4, p. 572). This is true also of the d⁹ system. The absorption band of copper(II) in aqueous media is complex and the bands tend to consist of many peaks over a rather scattered range, which is indicative of Jahn-Teller distorted systems (2; 4, p. 572). Most copper(II) compounds

^{*}Here and elsewhere in this manuscript, the standard error or estimated deviation will be specified in parenthesis following the parameter value. The digits of this number correspond to the least-significant digits of the parameter.

exhibit such extreme distortions that they are best described as distorted square planar. The distortions most frequently occur along the dz^2 axis (6, p. 322).

A comparison of some copper(II) pyridine bonded systems is shown in Table I. It should be noted that the square planar and octahedral complexes are distorted. Huheey has compiled a table of bond distances for Jahn-Teller distorted copper(II) complexes, part of which is shown in Table II (7, p. 325). From the table, it can be seen that, for octahedral distorted structures, certain bond lengths differ dramatically. These deviations could be, in some cases, explained by arguments consistent with the Jahn-Teller theorem.

Using this information, it is possible to say tentatively that BPGC will be distorted from an octahedral or tetrahedral structure. Merritt found that short hydrogenhydrogen contacts produced steric interaction, thereby resulting in distortions (11). The bonding sights for dipyridine in all cases observed have been at the nitrogen positions. The glycine molecule can bond at either oxygen or nitrogen positions, or both.

From this information, it is possible to obtain several chemically acceptable structures of BPGC. The crystal structure of this compound was attempted in order to observe

TABLE I

BOND LENGTHS (A) AND COORDINATION STRUCTURE OF SOME COPPER(II) COMPLEXES CONTAINING PYRIDINE

Copper(II) Coordination Structure	Cu – N	Cu - 0	Cu - Halide
Trigonal Bipyramid (13)	2.04, 2.01 2.07, 1.99	2.30	
Trigonal bipyramid (17)	1.98, 2.11 2.07, 1.96	2.05	
Trigonal bypyramid (2)	2.03, 2.10 1.96, 2.00		2.70
Square planar distorted (15)	1.98, 1.98		2.25
Square planar distorted (15)	1.98, 1.98		2.39
Square planar dístorted (16)	1.98, 1.98	1.99, 1.99	
Octahedral distorted (12)	1.98, 2.02 1.97, 1.99	2.45, 2.73	
Octahedral distorted (6)	1.98, 2.02	1.98, 1.95 2.40, 2.40	
Octahedral distorted (10)	2.00, 2.00	1.92, 1.92 2.38	
Octahedral distorted (3)	1.98, 1.98 2.03, 2.03	2.72, 2.72	<u></u>
	(

TABLE II

Compound	Short I	Distanco	es (Å)	Long Distances (Å)
	Cu-Cl	Cu-0	Cu-N	Cu-C1
CuCl ₂	2.30 2.30 2.30 2.30			2.95 2.95
CsCuCl ₃	2.30 2.30 2.30 2.30			2.65 2.65
CuCl ₂ ·2H ₂ 0	2.28 2.28	1.93 1.93		2.95 2.95
$CuCl_2 \cdot 2C_5H_5N$	2.28 2.28		2.02 2.02	3.05 3.05
Cu (NH ₃) ₂ C1 ₂			1.95 1.95	2.78 2.78 2.78 2.78 2.78

METAL-LIGAND DISTANCES OF SOME COPPER(II) COMPLEXES*

*Table from Huheey (7)

effects of Jahn-Teller distortions on the copper coordination as well as the types, if any, of hydrogen bonding.

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CHAPTER II

EXPERIMENTAL

Deep blue, long prismatic crystals of BPGC (1) were obtained from A. Syamal of the Department of Chemistry, North Texas State University, Denton, Texas. The crystals were recrystallized in a slightly basic aqueous solution in order to reduce the crystalline size to a suitably small size for the General Electric manual diffractometer. Flotation methods indicated a crystal density of approximately 1.55 g cm⁻³. Using four molecules per unit cell, the actual calculated density is 1.665 g cm^{-3} . These crystals were examined with the aid of a polarizing microscope, and several optically suitable crystals were selected. Single oscillation photographs revealed, after several tried, a non-twinned crystal. The crystal chosen was approximately 0.30 x 0.42 x 0.25 mm in size. The crystal was mounted along the long axis, later designated as the b axis, in an effort to minimize absorption problems. From the oscillation and Weissenberg photographs, the cell geometry was determined to be monoclinic. It was determined from the systematic absences of reflections,

> OkO absent when k is odd, hOl absent when l is odd,

that the unit cell is of the centric space group $P2_1/c$ (3). Easy alignment on the General Electric XRD manual diffractometer was possible due to the fact that the crystal was mounted along the b axis giving a plane of reflection at a chi (χ) of 90[°] and any setting of phi (ϕ).

Using cell constants calculated from Weissenberg photographs, preliminary settings of two-theta $(2\theta), \chi$, and ϕ were calculated for twelve reflections in the 20 range of between 0.00° and 23.00° . The intensities were maximized for these reflections, none of which were in the zero layer, and the settings of χ , ϕ , and 20 were recorded. Subsequently, these settings were used to refine the unit cell constants through a least squares fit of the original settings to the settings calculated from the indices. These results are shown in Table III.

Using the unit cell constants thus obtained, a new set of reflection settings was generated to a 20 of 60°. Intensity data were collected, using the peak height method. Molybdenum K α_1 radiation (λ = 0.70926 Å) was used with a zirconium filter. The x-ray tube was operated at 40,000 volts and twenty milliamps. The takeoff angle on the x-ray tube was 2.00°. Originally, data were taken using a forty second peak top count and two background counts of twenty seconds each. This procedure was modified to a twenty second peak top count and ten second background counts after approximately twenty percent of the data had been

TABLE III

Cell Constants	Film Data	Least Squares Fit Data
А	10.473 Å	10.474(3) Å
В	18.525	18,507(13)
С	7.723	7.724(6)
β	103.87 ⁰	103.918(59) ⁰

THE EXPERIMENTAL AND CALCULATED CELL CONSTANTS FOR THE MONOCLINIC UNIT CELL OF BPGC

This modification was made due to the excessive time taken. being consumed in data collection. Initially, the (800) reflection was used as a standard. When counting time was reduced, it was decided to add a thicker zirconium filter. This modification reduced the (800) standard to below 1000 counts per twenty second timing period; therefore, a new standard, the (830) reflection was chosen. The standard was measured at the end of each fifteen intensity collec-When the standard reflection intensity began to tions. fluctuate over ten percent, data collection was stopped for that day under the assumption that the diffractometer power supply stability was beginning to vary with increasing temperature. All data were collected using the same crystal.

The raw data were interpolated between standards and then corrected for peak spreading as 2θ increased. The

constants used to account for peak spreading were empirically determined by graphing as the abscissa,

and 20 as the ordinate for numerous reflections along the entire 20 range. A least squares best line was drawn through these points and the constant values taken from this line. This graph is shown in Figure 1. There were initially 3649 reflections, ranging through a 20 of 55[°].

The data thus obtained were multiplied by 0.9 to minimize computation errors. The resultant intensities, I_{obs} , were corrected for Lorentz and polarization effects (6, p. 196), which led to the observed structure factor, F_{obs} , for each hkl reflection. The structure factors were put on an absolute basis by the equation,

where K is the scale factor calculated by the statistical methods of Wilson (5) and $\left|F_{abs}\right|$ is the observed structure factor scaled to an absolute basis. In the Wilson method, a value for K is obtained from the equation,

$$\ln \frac{1}{\sum_{f_{\bullet,i}}^{1}} = \ln \frac{1}{K^2} - \frac{2B\sin^2\theta}{\lambda^2}$$

where B is the overall isotropic temperature factor and f_{o_i} is the atomic scattering factor for the ith atom. This latter quantity is summed over all N atoms in the unit cell. B and K are determined from the slope and intercept,





respectively, of the Wilson plot. The values of B and K were found to be 3.20 $Å^2$ and 0.7645, respectively.

Normalized structure factors, E's, were computed according to the equation (5),

$$E = \frac{KF_{obs} \exp \left[B(\sin\theta/\lambda)^2\right]}{\epsilon \sum_{i=1}^{N} f_{e_i}^2}$$

The factor \in is a statistical factor determined by symmetry which allows for the fact that reflections which are systematically absent have average intensities which are integral multiples of the average of the general reflections (6). For the space group P2₁/c, \in is equal to two for hol and 0k0 reflections, and equal to one for all other reflections. The experimental values for the E distribution are found in Table IV. This again confirms the centric space group.

TABLE IV

THE EXPERIMENTAL DISTRIBUTION OF E VALUES FROM DIFFRACTOMETER INTENSITY DATA FOR BPGC COMPARED WITH THE THEORETICAL CENTRIC AND ACENTRIC DISTRIBUTIONS

E Value Distribution	Experimental	Theoretical (Centric)	Theoretical (Acentric)
Average magnitude of E	0.814	0.798	0.886
Average magnitude of ${\rm \textbf{E}}^2$	1.003	1.000	1.000
Average magnitude of (E ² -1)	0.931	0,968	0.736
Percentage of E values greater than 1.00	32.86	32.00	37.00
Percentage of E values greater than 2.00	4.59	5.00	1.80
Percentage of E values greater than 3.00	0.26	0.30	0.01

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CHAPTER III

STRUCTURE DETERMINATION AND REFINEMENT

It has been shown that when a molecule possesses an atom which comprises approximately thirty percent of the total molecular weight of the molecule, heavy atom methods are preferred over a direct method solution. Heavy atom methods were used to determine the structure of BPGC due to the fact that the copper composition is approximately twenty-one percent.

At this point, in the absence of knowledge of the signs of the structure factors, it was not possible to deduce directly the actual positions of atoms contained within the unit cell. However, it is theoretically possible to deduce the interatomic vectors between the atoms. It has been shown by Patterson (1) that a Fourier synthesis, using the values of the observed structure factors squared, as Fourier coefficients, will produce a density map which yields the coordinates of the ends of these vectors according to the equation (1),

$$(u, v, w) = \frac{1}{\sqrt{h}} \frac{h}{k} \frac{1}{k} \left| F_{oss}(h, k, 1) \right|^2 \cos 2\pi (hu + kv + 1w).$$

Here, $\bigwedge^{O}(u, v, w_{c})$ represents the magnitude of an interatomic vector peak at the point (u, v, w). If one considers the

fact that groups of atoms are spherically symmetric (9), implying complete scattering, then the calculated structure factor has the equation

$$F_{calc} = \sum_{j=1}^{N} f_j \exp(2\pi i S \cdot r_j).$$

Here, r is the position vector of atom j and is mathematically defined as $hx_j + ky_j + lz_j$. S is the reciprocal lattice point at which the function is evaluated. One can derive the square for the structure factor by simply multiplying the factors together for the peak involved. For two atoms, j and k, the vectors r_j and r_k enter into the relation

$$\left|F_{h}\right|^{2}i,n=\sum_{j=1}^{N}\sum_{k=1}^{N}f_{j}f_{k}\exp(2\pi is\cdot(r_{k}-r_{j})),$$

where $|F_h|^2$ i,n represents a structure in which "atoms" of scattering factors $f_j f_k$ are located at $r_k - r_j$ (10, p. 228).

In 1936, Harker showed that for three-dimensional Patterson maps certain symmetry related conveniences were present (4). These conveniences were restricted to structures possessing certain symmetry operators (10, p. 232). For the BPGC monoclinic system, the Harker sections correspond to the P2/m point group (7). The P2/m symmetry lacks the glide plane and screw axis required for $P2_1/c$.

Using Harker symmetry points, it was possible to find the copper coordinates, from which a set of structure factors was calculated according to the equation (7)

 $F_{calc} = \sum_{j=1}^{N} f_{j}^{o} \cos 2\pi (hx_{j} + ky_{j} + lz_{j}) \exp(-B \sin^{2}\theta/\lambda^{2}).$ In the above equation, f_{j}^{o} is the atomic scattering factor at zero degrees kelvin for the j'th atom and the exponential term is the temperature correction for isotropic vibration. The scattering power curves were obtained from the averaged set of scattering curves in the International Tables for X-ray Crystallography (7, pp. 216-227).

Having obtained the calculated structure factors (F_{calc}) and the previously defined observed structure factors (F_{obs}) , it was possible to calculate the residual index (R). This index, N [, . .]

$$R = \frac{\sum_{i=1}^{N} |F_{obs}| - |F_{calc}|}{\sum_{i=1}^{N} |F_{obs}|}$$

where i represents an individual reflection and N represents the total number of reflections, is a measure of the reliability of a proposed structure solution. It has been shown (11, p. 372) that the theoretical statistical residual index should have the following values for centric and acentric random structures:

> $R_{random, centric} = 0.83$ $R_{random, acentric} = 0.59.$

Using only the copper atom found from the Patterson

synthesis, the residual index was 0.568. This index value was below the theoretical random centric value.

This copper position was refined isotropically through a block diagonal least squares, minimizing the quantity

$$\sum_{h,k,l} w(h,k,l) (|F_{obs}(h,k,l)|^2 - |k F_{calc}(h,k,l)|^2)^2$$

where k is a factor scaling the F_{calc} to the F_{obs} . This method of refinement (4) allows the inclusion of "accidentally absent" reflections which are too weak for observation. Usually, the minimum observed structure factor can be estimated in the part of the spectrum where the reflection resides. Thus, in a centrosymmetric space group, such a $P2_1/c$, the expected observed structure factor for the reflection can be estimated as (1) 0.5 F_{obs} min and its variance estimated as $Q^2 = 0.08 F_{obs}$ min, where F_{obs} min is the minimum observable structure factor. In the above equation, w is represented as $W = Q^{-2} = 12 / (F_{obs} min)^2$.

One cycle of least squares isotropic refinement produced calculated structure factors which were used to produce an electron density map, using the equation (6, p. 329) $P(x,y,z) = \frac{4}{V} \left(\sum_{h,k,1}^{k+1=2N} [A - B] + \sum_{h,k,1}^{k+1=2n+1} [D - C] \right)$ where P(x,y,z) is the electron density at the point (x,y,z)and $A = [F(h,k,1) + F(h,k,\overline{1})] \cos 2\pi hx \cos 2\pi hx \cos 2\pi lz$, $B = [F(h,k,1) - F(h,k,\overline{1})] \sin 2\pi hx \cos 2\pi hx \sin 2\pi lz$,

 $C = [F(h,k,1) - F(h,k,\overline{1})] \sin 2\pi hx \sin 2\pi ky \sin 2\pi lz, and$ $D = [F(h,k,1) - F(h,k,\overline{1})] \cos 2\pi hx \sin 2\pi ky \sin 2\pi lz.$

From the electron density map produced, two nitrogen and one oxygen atom were located. Subsequent isotropic least squares refinement produced more precise coordinates for these atoms which led to a residual index of 0.413. The electron density map produced from these coordinates revealed the positions of one oxygen atom and twelve carbon atoms. Once again, an isotropic least squares refinement was carried out and resulted in a new residual index of 0.28.

At this point, a difference Fourier map was calculated using the equation (12, p. 356) $\triangle P(x,y,z) = \frac{1}{V} \sum_{h \ k \ l} \sum_{h \ k \ l} (F_{obs} - F_{calc}) \cos 2\pi (hx + ky + lz),$ where $\triangle P(x,y,z)$ represents the difference between the density calculated with the observed magnitudes and the trial structure (3, p. 198). This method revealed one more oxygen atom which, when included in an isotropic least squares cycle, lowered the residual index to 0.26.

At this point, six cycles of anisotropic block diagonal least squares calculations were performed on the structure resulting in a lowering of the residual index to 0.154. The calculated structure factor expression for anisotropic temperature factors is given in the expression $F_{calc_r} = \sum_{r} f_r^o \cos \left[2 \pi (hx_r + ky_r + lz_r) \right] \exp(-C_r)$

(12, p. 450), where C_r is defined as $C_r = B_{11}h^2 + B_{22}k^2 + B_{33}l^2 + 2B_{23}kl + 2B_{13}hl + 2B_{12}$ for the rth atom. In the above equation, f_r^o is the atomic scattering factor at zero degrees kelvin and the exponential term is a correction for the anisotropic thermal vibration.

Using the calculated and observed structure factors, a graph was constructed of $\sin^2\theta$ verses relative index for all data collected. From the graph, it was possible to determine that the data were inconsistent above forty-five degrees in 20. The computer program used as an aid in generating Figure 2 is contained in the appendix.

Using the forty-five-degree cutoff point reduced the number of usable reflections to 1896. Seven cycles of anisotropic refinement reduced the residual index to 0.101. At this point, theoretical atomic positions for the hydrogen atoms were calculated for the rings, as well as the tetrahedral positions (3). The hydrogen atoms coordinated to the two free water oxygen atoms were not determined. The temperature factors for the hydrogen atoms were the same as those for the atoms to which they were coordinated; however, they remained fixed during refinement. The hydrogen atomic positions were included in the refinement, but were not allowed to vary. An additional three anisotropic refinements reduced the residual index to 0.091. With the change in the atomic positions of the non-hydrogen atoms,



Fig. 2--The graph of relative index verses $\sin^2\!\theta$ required to choose a data cutoff point.

it was necessary to recalculate the hydrogen positions. Using these new hydrogen positions, one last structure factor calculation resulted in a final residual index of 0.081.

The final atomic and thermal parameters for BPGC are listed in Table V. A listing of the final observed and calculated structure factors for the 1896 reflections used is given in Table VI.

TABLE V

FINAL ATOMIC AND THERMAL PARAMETERS OF THE CRYSTAL STRUCTURE OF BPGC

Atom	Fra	actional Coordinates X 10 ⁵	
	Х	v	N
Cu	13724(14)	8656 (8)	32998 (21)
c1	5819(32)	17508(18)	5809 (47)
0(1)	-1780 (69)	2380 (69)	28201 (110)
0(2)	78371(75)	1034 (49)	33829 (118)
Q(3)	26719(90)	38262 (50)	35056 (127)
0(4)	85258(113)	25980 (64)	22973 (154)
N (1)	24339(81)	1606 (47)	22446 (119)
N(2)	31456 (86)	13577(51)	40760 (126)
N (3)	4982 (86)	13920 (51)	50080 (134)
c(1)	19540 (108)	-4250(67)	12910 (166)
C (2)	27740(117)	-9400(68)	7340(165)
C (3)	41180(121)	-7980 (71)	11970 (177)
C(4)	46100(106)	-0164 (68)	21640 (169)
C (5)	37389 (104)	2950 (62)	26473 (147)
C(6)	41289 (105)	9932 (59)	36345 (148)
C(7)	54146(105)	12548 (67)	40872 (169)
C(8)	56440(112)	19076 (70)	48709(181)
C(9)	46412 (127)	22790 (69)	52700(187)
c(10)	33925 (123)	19946 (69)	48396 (184)
c(11)	-8910(115)	11340(67)	46526 (194)
C(12)	-11260(113)	4446 (68)	35219 (164)
H(1)	90355	44903	40610
H(2)	76044	35829	50127
H(3)	84710	15540	39493
H(4)	44128	49570	25017

VContinued
TABLE

F					5	
ALOM			Fractio	nal Coordinat	ces X 10 ⁻	-
		x		λ		N
H(7)	61	852		9295		38016
H(8)	65	667		21304		51808
H(9)	48	167		22052		9235
H(10)	26	414		23034		51348
H(11)	88	1667		10271	eta . , a	59083
H(11)	53	499		38230	-	41857
H (N)	5	1831		37206		13169
H (N)	Ω	177		30546	-	-2040
			Anisotropic	Temperature	Factors X 10	D
	B(1,1)	B(2,2)	B(3,3)	B(1,2)	B(1,3)	B(2,3)
Cu	387 (15)	172 (5)	I797(35)	-32(16)	795 (36)	-193(25)
Cl	905(41)	198 (12)	2166 (87)	19(37)	(96) [96)	156 (53)
0(1)	353 (86)	226 (32)	2410(225)	-238 (86)	803 (222)	-406 (138)
0(2)	487 (95)	376 (40)	2723 (251)	-257(102)	1107 (250)	-366(163)
0(3)	1368 (138)	327 (40)	2914(272)	-35(122)	2215(317)	77 (169)
0(4)	2095 (195)	568 (57)	3770 (353)	582 (180)	1114(423)	64 (237)
(1) N	473(106)	149 (34)	1451 (222)	1 (99)	908 (250)	-278(142)
N (2)	505(113)	221 (39)	1711(242)	-108(107)	746 (265)	-122 (157)
N (3)	507(112)	198 (38)	2255 (267)	11(107)	999 (285)	527 (164)
c(1)	452 (136)	263 (51)	2026 (328)	110(136)	682 (339)	-47 (209)
c(2)	905 (162)	231 (50)	1968 (324)	47(150)	1447 (372)	168(213)
C(3)	934 (165)	273 (53)	2452 (359)	-197(161)	1883 (397)	115(233)
C(4)	401(134)	259 (52)	2162 (333)	-3 (135)	757(341)	-128(211)
C (5)	485(131)	214(46)	1220 (265)	-11(127)	722 (300)	-13(180)
c (e)	647(140)	129(41)	1320 (267)	66(124)	446 (309)	310(170)
c(7)	257(126)	252(49)	2295 (327)	-114(128)	165 (322)	89(207)

-282 (232) 457 (228) -1148 (234) 115 (208) -397 (235) B(2,3) 105 B(1,3) Anisotropic Temperature Factors X -274(145) 402(157) -292(152) 194(144) -12(138) B(1,2) . . 2634 (373) 2742 (388) B(3,3) 279(55) 204(52) 242(53) 275(53) 193(48) 219 219 219 219 219 219 219 219 B(2,2) 219 219 219 219 219 (1, 1)726 726 മ Atom H (N) H (N)

TABLE VI

FINAL OBSERVED AND CALCULATED STRUCTURE FACTORS FOR BPGC

		-			-									•	· •				• •
4	ro		Ç	4	10	FC	L	FO	FC	្រុះ	· FO	- PC	L	PO	_ PC	j 6	- 70	j 20	
	- 11 - 11 - 11 - 11 - 11 - 11 - 11 - 1				a		1	528	-527	le n		64 2	-6	116	+111		310		
Ha	÷.	英 申	0		VA N	- · ·	2	21	-44	• • .	• •		-7	102	-117	7	72	320	
			-	1	1.6.6	145	3	1.67	172	ن ا	325	266					· • •		:
· 2	123		3	5	715	-718		64	-60	- Ī.	224	-158	<u>þim</u>	1.	X= 6	He	1.	K# 10	
	214		12	3	367	-371	-5	222	+228	-	190	-116					- •		
-0	40		54 .	-4	200	118				3	531	-473	0	74	32	o	117	-62	
-9				÷5	380	342)H#	٥.	K= 15	-4	£2	-13	- j - 1	342	-122	ī	157	179	
44 m	o.	K.e.	Ł	-6	ى ئ	57				5	99	139	2	31	-47	2	- 11	35	
-			•	-1	64	60	· 1	200	+208	6	211	247	· 3	55	- 16	3	742	803	
1	240	-7	55				2.	404	393	7	127	-144	. 4	568	513	. 4	281	-296	
ž	802	a	56	24 H	0, 1	(* 🕴	٤	250	-251	-1	\$056	1113	5	24	-39	5	155	-171	
3	662	6	30					1.40	112	-2	357	304	. 6	255	-219	6	20	5	
	453	4	12	9	Zud	261		101	114	-3	337	-305				1	519	+506	
-5	403	-3	56	1	95	35	Ma	٥.	Ko 14	-4	152	-170	-1	1444	. 231	+2	469		
- 4	141	1	64	. 2	383	+147		•••			- 163. 17.1	9 2. 13 3	- 15	1940	2/11	-3	205	181	
-1	. 167	- 1.	01	1	134	+03	0	245	-215	-0	209	2017		561	-513	-4.	152	166	
-8	- 61	-	46		136	- 464	- Î-	271	-291		300	230 231		20			95	-142	
	•		_		117	-112	2	304	299	-0	•••			133	-115	- 7	52 188	- 89	
H=	0.	Я, н	- 2		176	-177	. 3	136	143 -	H#	1.	X + 3	· - 7	20	13		1.03		
		·	÷.,		•••	•••	, -4	37	57		••	•• •	•		•••		1.	M	
	1000	10	76	Ma.	0 . 1	K# 9				0	713	-647	M -		***		1.1	N.M. 11.	
	1031		-0	••	••		ļe 🗰	0,	K+ 17	· 1	758	*/3a	1		K4 7	a	30		
· · ·	1,303	- 19 5	90 20	Ľ	105	-635				. 2	984	1052	0	140	5.1.5		367	- 161	
	- 2 74		80	ž	201	+280	1	280	283	.́.Э	376	321	- ĭ	\$70	54.3	. 5	50	50	
	111		92	Ĵ	94	-73	2	59	64	- 4	309	+305		379	- 147	ĩ	144	184	
	-147	1	47	-4	449	-+14	3	230	-2+9	5	33	-12	· .	356	- 1 - 1		21	+14	
-7	231	2	16	->	9	1	**	20	-15	ં ઠ	115	-131	- T			5	÷ 40	46	
- 6	26	_	26	-6	3	-4 (1	159	134	2 -	520			148	128	ľ
-				· -1	21	31	FH	0.	K= 18	·-1	743	-733				- L	495	459	
14m	0.	X =	3	jąm.	. 0.	K= 10 .				-2	637	045	7	117	140	-2	164	186	
							0	372	-354	- 3	130	- 532	· • •	5.61	-570	- 3	313	- 472	
1	359	3	28	1 0	819	743	L	31	49		277	222	-2	173	-147	· - 4	222	-205	
2	1051	20	57	1 1	212	255	Z	214	228		400			710	+610	-5	198	-218	
ڌ	. 58	فتا	22	2	1 80	-192	3	137	-124	- 7	400	-110	-4	318	= 340	-6	20	19	
-4	551	5	10	3	143	196 .	· · ·			- 8	17	-5.3	-5	231	226	·			
. -	108	1	89		\$3	-101	H.	0,	K+ 18				-6	541	538	_ H# _:	, ir	K# 12	
-0	18	•	21 -	-5	155	194				He	1.	K= 4	-7	137	-159				
-1	99	1	14	~0	148	110	Ļ.	52	- b T	•	••					0	104	508	
	·			4 4m	0.	e		203	~291	0	413	-114	He	۱.	K= 8		340	- 30 8	
		N -	-	61 m				١.	# = ` A	1	83	-11					107		
	£ 2 5	-5	ao	1	187	155			- v	2	1045	-10-1	0	383	- 143	1	144		
	263	-7	23	;	232	205	n	A12	n 63. N	3	162	-122	1	341	-307		30		
2	\$ 54	ب ا	34		45L	448	2	64	-114	4	130	191	2	103	. 15	6	· · · · ·	- 2	
	77	7	59	-4	89	-107	4	510	-487	5	379	ذاذ	,	444	915	-1	108	124	
	633	5 5	81	-5	242	-245	6	169	182	6	- 51	-100	- 2	470	-141	-2	233	-219	
-5	. 492	2 -4	18	-6	116	-106	-2	433	- 161		47	-52	6	10	~ 27	-3	- 70	- 84	
-6	10	r -1	71				- 4	572	- 514	- ř	1013	1035	7	35	36		232	268	
-1	12	} 1	27) ja ka	0,	¥# 12	-6	270	511		979	414	-1	589	-526	-5	116	118	
	·						- 8	258	-203		053. 874	- 0 L 0 	ž	66	-46	+6	92	92	
71+	Ű,	I. №¶	5	0	546	324	A 1	-		-5	374	-113	-3	503	479	M=	1.	K+ 13	
		*		Ť	992		M #	1+	я» Ļ	-6	2.10	-219	-4	47	39		•		
	31		1415		341	-294		300	4.3.4	· -7	96	92	· =5	42	53	0	230	-250	
2	. OU:	י פ י- ג		د مہ	- 22 L L GR	-125		654	-1124 	-8	105	115	- 6	159	127	1	516	-508	
	2.4	1 1	936		1.62	85	,	100	-1120				-7	100	-91	2	201	236	
- 5	- 1 4	• '	q 2		154	165	1	614		N.	Le	K- 2				3	177	192	
6		<i>4</i> .	-46			_		239	-269	^	4 0 C		H =	1+	к* 9	•	- 223	-212	
- 1	2	φ.	-26	H#	0.	K= 13	5	214	~ 243	, ,	100	12.4	~			5	91	. 90	
•							6	146	- 11	1	781	724	U 1	220	489	− ↓	53	93	
M#	0	, Ka		⊧ 1	120	-142	1	89	-137	1	325	2.51	2	417	-430		224	199	
				2	1 20	172	÷1	634	633		64	-63	. в а	142	162	-3	417	223	
0	- 25	9 -	285	3	358	373	-2	212	214	5	231	~198	· 2	140	122		210	285	
- e - 1	24	1 -	234	- 4	. 444	472	- 3	698	6 ja	6	46		5	97	-109		120	-100	
é	47	1	445	-5	<u>جر</u> ا	-61		588	0 J2	7	163	167	6	258	260	-0		-00	
1	30	4	240	-6	- 19	23	- 5	166	-352	+1	519	533	- Ē	255	-221	He	. 1.	X = 34	
	01	d .	921 96		~		6	139	-126	, - 2	487	-+16	-2	177	+136				'
*3	• ¥	¥ 1 -	14	51 H	U a		-7	- 24	34	-3	19	-jj	-3	293	-226	. 0	187	-173	
	,)0 , 7	2	976 24		1 2 2	a	~0	- 24		-4	298	-231		583	-528	ī	216	-212	
							· · ·			5	589	576	5	70	-48	· .	•	2.8	

τ.	T O	-	7.	50	845	*	=0	-				· _	-		-	·	
-			. .				- - -	1.6	- -		- FC	L.	, LÔ	rc	L	n no) / / / C
3	207	-325	H +	2.	Ka 1-	H=	2.	K=: 5	1	388	- 382	- H#	Z e	K= 14		b 11	
· •	. 56	63			-				2	319	#306				-		
. 5	84	15	n	150	-110	n	373	200	3	91	137	0	- 197	-159	- Mar	1.	
- 1	544	560	ĩ		-317	÷	1000	1007	. 4	264	= 24.4	1	271	261	1		
و س	260	-263		413	- 39J	÷.	5080	1041	5	123	1.4.9	,	115	125			
	36.2	- 164		105	-517	· · ·	201	243		21.0	144		217		9		- 164
	434	-240	3	- 42	70)	356	-317		. 222	60.4	2	311		1	. 201	-178
		-112	4	18	-58	· •	229	\$37	-1	222	-324		10	-12	2	408	-420
5	14	30	5	195	162	5	449	-44	2	006	-568	2	• (53	1	.433	-455
			6	152	-185	6	10	-19	3	405	-388	_1	85	-76	. 4	132	140
利用	1.	K= 15	7	174	-197	7	115	1 15	4	235	-219		325	-309		220	255
			-1	140	- 178	- 1	114	-120	∴ -5	92	-108	3	259	-256			
0	311	-329		100	-120		1.30	~1.30	-6	· 3Ē	-15	-4	26	-5	1	27	- 10
. 1	72	71	-2	305	303		Ç 3	- /1	-7	66	4.2	-5	144	229		2 0	-19
;	444	461		10	~>		508	452		••		•			- <u>-</u> i		702
	151	-171		396	359	- 4	309	292	94 m	5.	8- 10	11 -	1	No. 18	-2	40	16
	12	-121	- 5	157	-347	- 5	276	274		~ •	M- 10	**	49	V. 13	3	F 492	2 -479
	< 0 7	~17	-6	70	-87	-6	241								-4	24	-227
		10	+7	- 260	262	-7	203	-199		140	m134	U		-473		281	-240
+1	265	+293	- 8	122	161	-8	89	97	. <u>1</u>	93	+119	. 1	64	i 195	-6	51	-32
-21	52	-64							Z	360	365	- 2	265	263	-1	279	1 270
- 3	98	#114	He	2.	K= 2	He	2.	K= ~	3	536	r 531	. 3	14	-19		174	140
. н ђ.	70	99		- 1				• •	- 4	265	-262	4	130	- 118			
-5	42	67		A 10	-604			437	· 5	34	-7 :	1	115	-121	H-	1.	
	2			10.74 10.7	- 270	Ŷ	407		6	35	-19	-2	195	194		31	
H=	1.	Ke IA	1			1		-14	-1	117	- 9 A	-3	193	=179			
	• • •		2	209	-146	Z	793	-/50	+2	405	+387	-4	160	175	u	+95	-+52
•	117	-147	3	844	+836	3	89	105	-1	347	151		44	-58	1	272	243
, v	34	-176	4	68	. 76	4	401	19A		234	214				Z	289	232
		00	- 5	168	193	5	137	-181		667	-40*	64-	3	K. 14	3	615	-635
2	24	~21	6	7	15	6	120	-125		210		***	4 9	V= 10	- 4	217	263
3	93	95	1	90	98	. 7	42	36	-0	43	-69				5	125	+124
4	165	- 141	-1	1.60	95	-i	314	-279	- 7	76	-109				. 6	74	- 74
-1	389.	414	,	56.4	-571		1101	1166	•				923		. 7	214	284
-2	248	265		104	- 375		116	141	M	•	F	Z	109	-111	+1		
- 3	232	+761		132	-224		112			6 8	N- 11	3	13	~1		400	-041
	102	-120.		114	10		203	-244	à			- 4	61	80	4	204	70.04
•	•			242	504	+5	110	-137	0	486	460	-1	. 23	7.		102	-010
Mat	1.	Ke 17	-6	- 78	-97	-6	95	84 -	1	197	-186	-2	301	320		365	374 -
•••=	• •		- # 7 -	177	162	_ - 7	, 21	6	2	100	78	3	202	-223	~ ⁵	742	711
	1.12		- 8	24.	-27	. Ma	2.	K.e. 1	3	26	-57	-4	112	+161	· •6	354	+3681
Ŷ	167	-130							4	72	-64	•		• • •	-1	- 29	- 13
	742	241	N/m	2.	K# 3	•		340	5	158	13.6	M-	3.	-	- 5	34	30
2	29.	. ~1 4				U U	213	249	Ā	141	146		4 ¥.	N= Lr			
- 3	- 41	- 84	n	A 34	- 423		027	010			874				Hu	1.	K . 2
-1	45	-42	Ĩ	4. G . H	401	5	567	-516		10	-14	0	124	-119			
- 2	0	14	<u></u>	770	4.71	3	95	-57	- 2	10		. 1	186	205	1. 6	1.16	
- 1	119	+126		- 7 D		<u> </u>	428	-406		44	30	2	38	56	Ŷ	1.22	
-	73	22	1	21	-04	- 5	285	-273		8 r	-72	3	0	-7	. 1	114	115
•			4	203	176	6	205	215	-5	64	-48	_ −1	69	-95	· 2	353	-117 -
Ma	1.	R= 18	5	165	167	1	95	98	-6	48	-39	-2	26	-38	3	217	194
			6	195	-214							-3	153	167	- 4	447	425
0	55	-90	7	149	169		. 007	-001	He	2.	K= 12 -	-4	93	40	5	31	16
1	165	=191	-1	450	-459		587	-522							. 6	172	-167
;	10		- 2	613	669	-3	122	166	0	44	-6	. M=	9 .	Xa 14	1	112	124
		34	-1	54	50	4	- 82	15	i	124	-151		, 4 1	LO	-1	3.9	-316
ر ا ــــ	40 AL		- 6	25	· •••	-5	128	1,24	5	617	599	: a	71	91	-2	1001	977
		-121		71	- 50	6	° 181	174	2	114	-12	ĭ	21		-1	87	<u>, , , , , , , , , , , , , , , , , , , </u>
	552	101		434	-304	-7	166	-164		247			104	-202	-	547	-104
- 9	191	[97			-375					202	-317	- 1	1 4 1			103	
			-[67	102	H	2.	K= 4	2			-1	1.41	-130		104	-128
14 a	1.	K= 19		124	166	••	•••	14 - 14	- I	65	-62	-2	316	371	-0	. 1.44	-165
						•	2.34	- 200	- Z .	436	-415	- 3	104	124	. - I	83	9[
0	- 74	89	He	2+	к= 4	· •	31	-219	-3	138	121				- 8	201	217
1	1.61	155				- V 🛔	230	-249	-4	385	350	H=	2.	K= 19			
2	212	-222	0	54	+100	2	157	-190	-5	86	92				Heri	3.	K= 4
-1	143	-177	1	504	533	· 3	524	506	-6	17	- 7	. 0	61	91			
 	74	- <u>6</u> 7 7 44.	2	476	-480	4	229	230	1.	• •		ĩ	67	102	- 0 .	185	-178
e	14	20	3	264	+271	5	224	-242	He	2.	K= 13	- - i	112	+104	. 1	998	986
.	*			294	322	6	133	109					162	-113	. 2	118	-121
44 -	<+ >	4 - 0		21	107	-1	298	-257	0	467	-469		102	~	. ī	70	
-				40		-ż	185	225	1	218	-215	M-			Í.	101	104
ø	331	361				-3	775	711	· 2	49	-34	99 8	3+	K# 0	. 2	101	100
2	7.38	727		72	0)		110	-118	- ī	A.2	-95	, · · _ ·			· 7	22	0 3
4	254	-254	- L	123	339		216		· 4	<u> </u>	-10	0	655	-664	. 0	100	108
6	- 24	- 36	- Z	121	129	-	14			24	O	<u>,</u> 2	577	586	Ţ	119	1.10
- Z	594	.613	-3	466	-453					644	. OL .	4	221	-217	— L ∶	127	100
	LL64	1129		531	-498	4	. 76	392	: <u> </u>	44(776	6	322	-341	2	601	-549
-6	124	+121	- 5	50	64	3 4	-		-4	58		- 2	953	940	-3	677	-610
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-1	261	-288	H=	31	K= 9	·	202	-186		387	366	7	399	- 195	0	309	-292
~ 4	1.0	-14					1.30	132		121	-114	0	100	104		7 2 1	473
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6	79	61		363	-364	- ÷	47	52	- 5	6	-11	0	21	21	- 6	105	÷81
7	24	25	. **	1 40	181		6	- 24	6	62	54	-1	199	-114	-7	97	~104
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- · •	- 19	-13		330	111	-4	204	-231	4	222	570	+1	25	34	. ~ 7	- 11	64
5.5	· 85	-102	·	2.30	233		91	-121		268	-234			- 86.4			
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-6	356	363	1	47	-60	3	25	-20	- á	154	~198	ben -	4.	Ka 7.	1. 🕺		- 422
-7	90	-104	· 2	191	191	-1	115	125	-7	241	-206			•	5	Z.9	29
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U 1	- 24	-35				-4	282	-375	4	101	-118	-1	282	317	-4	256	-2#3
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-3	25	- 30		45	-93	· 1	175	104	-6	152	141		42	-17	5 i	228	234
-5	59	81	5	131	136	2	160	-157	-7	114	-126	- 4	371	-384	2	22	~16
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1	187	-202	-5	227	709	· ~]	89	-117	· 2	130	-178		• •77 • 14	493		197	211
2	191	-218	-1	241	267	-5	204	220	J	337	349	-6	64	-55	5 . - 5	234	263
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d a	488	447	-2	- 47	-211		157	-158	H u	6,	K= 3	-1	391	373	- i	- <u>\$</u> .	• €3€ ♥
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-3	337	324	0	- 77	-71	-2	155	-151	-2	518	-523		· •	- 36	3	177.	117
· #\$. * _\$	16	-106	2	241	-245		195	-22A	- - ₩	15	-18	ĩ	301	403	-z	243	-275
-6	160	167	3	62	73				-5	211	-265	2	211	221	-3	63	65
- 7	229	242	- 4 5	119	89 118	₩ n	- 5.	Ke 15	-7	90	-100	- 4	219	-217	-5	154	22
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3	51	37	-6	63	78	- 3	208	219	2	237	301 - 21 A	*8	19L.	-210		64 . 244 -	. 84. in:
: *	292	-272	Ha	5,	8#1 11		107	234	÷ 4	59	-83	Ha	6		-2	151	
-1	25	23	, 		- -	He	5,	Kn 17	5	165	-191		• • •		-3	279	303
-2	535	-522	0	149	-136	٥	135	141		65	-58	1	391 278	-368		172	-198
	- 	503	2	149	176	. 1	261	-315	-3	533	527	2	251	237	Ha	6. K	• 15
-5	106	97	3	1 89	-209	-1	171	187 -	-	243	247	3	101	115	•		
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	354	- 14	-5	391 511	378	0	119	-104	•••	4.	सन्तः हो। •्	-9	69	-104	-3	50	-83
1.	505	-496	-6	82	-95	2	790	-613	0	450	449	- 6	190	-224		48	45
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- 1	120	-127		±1	-07	-4	- 211	+229	7	96	107		16	3 180	- 2	140	121	È.
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- 2	441		-3	57	-100	0	337	335 -				· •4		-27	ji s	91	K.a.	4
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+1	389	384	· 2.	337	-361	0	180	193	. – 🕊			·	124	124	-0		141	
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CHAPTER IV

DISCUSSION AND CONCLUSION

Thermal Motion

All atoms in a molecule undergo a certain amount of This vibration increases with increasing temperavibration. tures; thus, it is called thermal motion. The thermal motion of an atom results from the superposition of all the normal modes of vibration in the crystal and yields a very complicated result (1). The net effect of vibration is to spread out the electron density of the atom over a finite volume. The spreading of electron density over an atom's equilibrium position results in a volume element consistent with the bonding for the complex involved. Since a vibrating atom possesses a scattering power less than that of a stationary atom, it is necessary to determine the vibrational contributions to the scattering factor through x-ray methods. A first approximation would be to assume that each sphere of vibration possesses a constant radius and is, therefore, isotropic. This would produce an individual temperature parameter for each atom, $B = 8\pi^2 \bar{u}^2$, where \bar{u}^2 represents the root-mean-square amplitude of atomic vibration. As mentioned in the previous chapter, this produces an approximate model for a tentative structure,

but it represents each atom as a centroid. A more accurate description of the temperature factor is given by considering the thermal motion as anisotropic. The anisotropic model contains all of the information normally obtained on thermal motion. The anisotropic temperature factor is expressed as,

 $\exp[-(B_{11}h^2 + B_{22}k^2 + B_{33}l^2 + 2B_{12}hk + 2B_{13}hl + 2B_{23}kl)],$ where h, k, and l are the indices of a reflection and the B terms are components of a symmetry second-order tensor, representing the vibrational modes for each atom of the molecule (3; 8, p. 44). The six B terms are convenient to use in least squares refinement.

One can compute the directions of each principal axis and the root-mean-square displacement along each axis of the thermal ellipsoid from a principal axis transformation done on the exponential forms containing the B terms. Table VII gives these root-mean-square displacements for BPGC. Figure 3 shows a CalComp plot of the molecule with fifty percent probability ellipsoids (4, p. 220; 5). From the figure, one can see that the principal axes with the longest root-mean-square amplitude for some atoms are not perpendicular to the plane of the bipyridyl ring. This would indicate that some poor data exist in the reflection list, which resulted in some difficulty in the least squares convergence. These slight abnormalities are not

TABLE VII

ROOT-MEAN-SQUARE AMPLITUDES (Å) OF VIBRATION IN THE PRINCIPAL AXES DIRECTIONS FOR THE ATOMS IN BPGC

Atom	Axis l	Axis 2	Axis 3
Cu	0.1283	0.1672	0.2310
Cl	0.1821	0.2097	0.2520
0(1)	0.1121	0.1917	0.2735
0(2)	0.1388	0.2411	0.2937
0(3)	0.2155	0.2403	0.3095
0(4)	0.2750	0.3289	0.3660
N(1)	0.1190	0.1590	0.2174
N(2)	0.1504	0.1929	0.2270
N(3)	0.1352	0.1735	0.2702
C(1)	0.1437	0.2159	0.2417
C(2)	0.1747	0.1968	0.2569
C(3)	0.1472	0.2318	0.2807
C(4)	0.1372	0.2090	0.2511
C(5)	0.1406	0.1902	0.1941
C(6)	0.1312	0.1806	0.2102
C(7)	0.1128	0.2109	0.2638
C(8)	0.1352	0.2291	0.2860
C(9)	0.1507	0.2389	0.2921
C(10)	0.1305	0.1969	0.3288
C(11)	0.1308	0.1820	0.3324
C(12)	0.1667	0.2238	0.2268



Fig. 3--The fifty percent probability thermal ellipsoids from the thermal parameters of BPGC. View is twisted ten degrees along the Cu-N(1) axis and also the Cu-N(2) axis with copper as the origin.

reflected to any major extent in the bond distances, however.

Molecular and Crystal Structures

Figure 4 shows the entire unit cell and atoms located in the region twenty percent beyond the cell in each direction. The viewing direction is rotated twenty degrees along the c axis and twelve degrees along the a axis. The calculated hydrogen positions are not shown. Bond distances are given in Table VIII. Valance angles for BPCG are given in Table IX.

TABLE VIII

		l
Atom 1	Atom 2	Distance Apart (Å)
Cu	0(1)	1.957(8)
Cu	N(3)	2.025(10)
Cu	N(1)	2,010(9)
Cu	N(2)	2.027(10)
Cu	C1	2.635(4)
0(1)	C(12)	1.298(15)
N(1)	C(1)	1,339(16)
N(1)	C (5)	1.350(14)
N(2)	c (6)	1.342(15)
C(1)	C(2)	1,417(18)
C(2)	C(3)	1,392(18)
C(3)	C(4)	1,426(18)
C(4)	C(5)	1,358(17)
C(5)	C (6)	1,506(16)
C(6)	C(7)	1,394(16)
C(7)	C (8)	1,347(18)
C(8)	C (9)	1.351(18)
C(9)	C(10)	1.374(19)
C(10)	N(2)	1.316(16)
C(11)	N (3)	1.492(16)
C(12)	C(11)	1.533(18)
C(12)	0(2)	1.244(18)

BOND LENGTHS (Å) IN BPGC





TABLE IX

VALENCE ANGLES (DEGREES) IN THE BPGC MOLECULE

Atoms Forming Angle

Δna	
LT I'U	10

Cu Cu CU CU Cu Cu	O(1) N(1) N(1) N(2) N(2) N(3)		000000	(12 (1) (5) (6) (10 (1)	2) - - 	•	•	• • •	• • • • •	• • • •	• • • •	• • • • •	• • • •	• • • •	• • • •	• • • • • • • • • • • • • • • • • • • •	• • • • •	• • • •	• • • •	- - - -	115.5(7) 125.3(7) 114.5(7) 113.8(7) 127.1(8) 107.5(7)
C1 C1 C1 C1	 Cu Cu Cu Cu		O N N N	(1) (1) (2) (3)	•	•	•	•	• • •	• • •	•	• • •	• • •	• • •	• • •	• • •	• • •	•	• • •	• • •	97.5(6) 100.3(6) 92.7(6) 96.6(6)
0(1) 0(1) 0(1) 0(1) 0(2)	Cu Cu C(12) C(12) C(12)		N N O C C	(1) (3) (2) (1) (1)	• • L)	• • •	•	• • •	• • •	• • • •	•	• • •	• • • •	• • • •	•	• • •	• • •	• • • •	• • • •	• • •	93.0(3) 86.7(3) 123.9(12) 116.7(10) 119.4(11)
N(1) N(1) N(1) N(2) N(2) N(2) N(2) N(3)	Cu C(1) C(5) C(5) Cu C(6) C(6) C(10) C(11)		NCCCNCCCC	(2) (2) (4) (6) (3) (5) (7) (9) (12		• • • •	•	• • • • • •	•	• • • • • • • • • • • • • • • • • • • •	• • • • • • • • •	• • • • • • • • • • • • • • • • • • • •	• • • • • • • • • •	•	• • • • • •	• • • • • •	• • • • • • • • • • • • • • • • • • • •	• • • • • •	• • • • • • • •	• • • • • • • •	81.6(3) 122.5(10) 122.2(10) 114.3(9) 96.7(3) 115.4(9) 121.1(10) 121.6(12) 112.7(10)
C(1) C(2) C(3) C(4) C(5) C(6) C(6) C(7) C(8)	N(1) C(2) C(3) C(4) C(5) C(6) N(2) C(7) C(8) C(9)		000000000000000000000000000000000000000	<pre>(5) (3) (4) (5) (6) (7) (10 (8) (9) (10</pre>		• • • • • • •	• • • • •	• • • • •	• • • • • • •	• • • • • • • •	• • • •	• • • • •	• • • • •	• • • • • • • • •	• • • • •	•	• • • • • • •	• • • • • • •	• • • • • • •	• • • • • •	120.1(9) 116.4(11) 120.5(11) 118.4(11) 123.6(10) 123.6(10) 119.0(10) 118.8(11) 119.6(12) 119.8(12)

One can easily see that the coordination about the copper atom is distorted from octahedral. One would predict from Table I that the most probable structure would be distorted octahedral.

From the least squares mean plane calculation, Table X, one can see that the copper atom is slightly above the plane of the four atoms coordinated to it. This could be due to steric interaction of the glycine group and the chlorine atom with the central copper atom.

It has been postulated by Hamilton (2, p. 13) that the O - H...Cl bonding distance is 3.10 Å to 3.20 Å. In the same article, he also postulates the O-H...O bonding distance to be from 2.70 Å to 2.80 Å.

The chlorine atom of BPGC is surrounded by three oxygen atoms which are approximately equidistant to it. From the bond distances found, it is apparent that these bonds are most probably hydrogen bonds. The proposed hydrogen bonded distances are 3.15, 3.20 and 3.19 Å for the two 0(4) - H...Cl and 0(3) - H...Cl bonds, and are shown in Figure 5. The 0(2) and 0(3) atoms also indicate, by their distance apart, that they could possess some hydrogen bonding between them. The proposed hydrogen bonded distance for 0(2)...H - 0(3) is approximately 2.76 Å, and can be seen in Figure 4.

TABLE X

EQUATIONS OF THE LEAST SQUARES BEST PLANES THROUGH THE TWO RING SYSTEMS OF BIPYRIDYL IN THE MOLECULE OF BPGC AND THE DEVIATION OF ATOMS FROM THE BEST PLANES

	Ring	Equations of Best Planes
Ring	1*	0.0397X + 0.4478Y - 0.8933Z = -1.3879
Ring	2*	0.1186X + 0.5054Y - 0.8547Z = -1.0212
Ring	1 + Ring 2	0.0789X + 0.4835Y - 0.8717Z = -1.1533

Atoms	Ring l	Ring 2	Ring l + Ring 2
N(1)	0.2022(89)	-0.0143(89)**	0.0003 (89) **
N(2)	-0.0162(94)**	-0.0199(94)	-0.0929(94)**
C(1)	0.3433(125)	0.0112(125)**	0.0742(125)**
C(2)	0.3278(125)	0.0002(125)**	0.0541(125)**
C(3)	0.1865(132)	-0.0082(132)**	-0.0171(132)**
C(4)	0.0781(127)	0.0053(126)**	-0.0545(126)**
C(5)	0.0945(111)	0.0058(111)**	-0.0404(111)**
C(6)	0.0208(117)**	0.0531(111)	-0.0425(111)**
C(7)	-0.0147(124)**	0.1578(124)	-0.0028(124)**
C(8)	0.0040(135)**	0.2762(134)	-0.0755(135)**
C(9)	0.0009(139)**	0.2359(137)	0.0584(137)**
C(10)	0.0051(137)**	0.1006(136)	-0.0929(94)**

Distances of Atoms from Best Planes (\hat{A})

Rings are defined in figure 3.

**Atom used to define the least squares plane.



Fig. 5--The hydrogen bonded environment surrounding the chlorine atom, with symmetry related glycine group opposite to the chlorine. (See text for further explanation.)

TABLE XI

EQUATIONS OF THE LEAST SQUARES BEST PLANES THROUGH THE FIVE ATOM CENTRAL COORDINATION AND GLYCINE LIGAND OF BPGC AND THE DEVIATION OF ATOMS FROM THE BEST PLANES

Equations of Best Planes Plane 1. -0.1372X + 0.5794Y - 0.8034Z = -1.4082 Plane 2. -0.1365X + 0.5792Y - 0.0837Z = -1.3615

· · · · · · · · · · · · · · · · · · ·		
Atoms	Plane 1	Plane 2
Cu	0.2358(15)	0.1886(15)*
N(1)	-0.0641(89)*	-0.1099(89)*
N(2)	0.0610(94)*	0.0146(94)*
N(3)	-0.0594(99)*	-0.1079(99)*
0(1)	-0.0625(81)*	0.0146(81)*
0(2)	-0.1216(89)	-0.1710(89)
C(11)	0.0683(138)	0.0190(138)
C(12)	-0.0146(124)	-0.0343(124)

Distances of Atoms from Best Planes(A)

*Atom used to define the least squares plane.

From Table XI one can see that the least squares best planes through the bipyridyl system indicate that the two bipyridyl rings are not coplanar. The dihedral angle between the planes is approximately six degrees. It is possible to see this puckering of the two rings in Figure 4. This is unusual since other copper structures coordinated to 2,2'-bipyridyl show the pyridine ring as planar (6, 9). The distance between the C(5) and C(6) atoms indicates a small amount of π bonding, if any, between them, as compared with the other bond distances of other atoms composing the rings. Merritt (7) has shown that the free bipyridyl system is in the least sterically hindered position, with the two nitrogen atoms opposing each other by approximately 180 degrees relative to the center of the molecule. The N(1) - Cu - (N(2) bond angle is only eighty-one degrees. These facts could cause a weakening of the C(5)-C(6) bond by causing a bending stress on the bond.

The glycine ligand could be a stabilizing factor in the tetrahedral distortion due to the sp^3 nature of the C(11) and N(3) atoms. These sp^3 bonds cause a slight misalignment of the sp^2 bonded O(2), C(12), and O(1) atoms. The steric configuration of the molecule is best seen in Figure 6.



Fig. 6--The unit cell of BPGC. Viewing direction is rotated twenty degrees along the a axis and twelve degrees along the c axis.

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APPENDIX

An IBM 360 computer program was used in the determination of diffractometer linearity. This program produces a series of printer plotted graphs of R index verses $\sin^2 \Theta$, containing a data point for each reflection. From these graphs, it is possible to observe the contribution of individual reflections as well as small groups of reflections to the overall R index. The input to this program consists of $\sin^2 \Theta$ and individual R index for each reflection observed.

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   Ň.
                       FROM THESE GRAPHS IT IS POSSIBLE
      PREVIOUS GRAPH.
С
      TO VISUALLY DETERMINE THE CUTDEE POINT FOR THE HIGH
   ★
                                                            *
С
   龙
      AND LOW TWO THETA RANGE.
                                THIS IS ACCOMPLISHED BY
C
      DRAWING TWO LINES, ONE ON EITHER SIDE OF THE DIS-
   *
                                                            *
С
      CONTINUOUS LINE OF AVERAGES.
   *
                                    WHEN THE TWO LINES
C
      BEGAN TO DIVERGE OR THE AVERAGE OF THE LINES BEGINS
                                                            末
C
      TO FALL THE DATA AT THAT POINT IS CONSIDERED BAD.
   ×
   *******
C
      REAL A(28), BAR, BLIN, GRAPH(60, 45), B(28), DOT, N, VERT(45)
      READ(5,105)VERT
      REWIND 1
      N = 450
      DATA A /* *,*A*,*B*,*C*,*D*,*E*,*F*,*G*,*H*,*I*,*J*.*
    2 *K*,*L*,*M*,*N*,*O*,*P*,*Q*,*R*,*S*,*T*,*U*,*V*,*W*,
     3 * X*, * Y*, * Z*, * **/
     DATA BAR, BLIN/ !!.
                         • /
     DATA COT/ .../
     DO 112 I = 1,28
 112 B(I) = I - 1
1000 DC 1 J=1,45
     DO 1 I=1.60
   1 \text{ GRAPH}(I,J) = A(1)
     DC 2 I=1,45
     GRAPH(1,I) = BAR
   2 \text{ GRAPH(60,I)} = BAR
     DO 3 I=1,60
     GRAPH(I,1) = BLIN
   3 \text{ GRAPH}(1,45) = \text{BLIN}
   4 REAC(1,20,END=100)SINSQ,RINDX
     I = N * (SINSQ*.5) + 1.
     J = RINDX / 4. + 1.
     IF((I.GT.60).OR.(J.GT.45))GO TO 4
     IF(GRAPH(I,J).EQ.A(28))GO TO 4
     DC 5 INDX = 1,28
```

```
IF(BAR.EQ.GRAPH(I,J))GO TO 6
    IF (BLIN. EQ. GRAPH(I, J)) GO TO 6
    IF(A(INDX).EQ.GRAPH(I,J))GC TO 6
  5 CENTINUE
    GO TO 4
  6 GRAPH(I,J) = A(INDX + 1)
    GC TO 4
100 WRITE(6,103)
    DO 110 I=1,60
    XNUM = 0.0
    XDEN = 0.0
    DC 111 J=1.45
    DC 111 L=1,28
    IF(GRAPH(I_J) = Q = A(L)) \times NUM = \times NUM + B(L) + J
111 IF(GRAPH(I,J) \cdot EQ \cdot A(L)) \times DEN = \times DEN + B(L)
    IF(XCEN.EQ.O.)GO TO 110
    NAVG = XNUM / XDEN + 1.
    GRAPH(I, NAVG) = DOT
110 CONTINUE
    DO 101 J=1+45
    JJ = (J-1) * 4
101 WRITE(6,102)(GRAPH(I,J),I=1,60),JJ,VERT(J)
    DC 7 I=1,61
    XC = I
    SINSQ = XC/N
    SI = SQRT(SINSQ)
    SII = ARSIN(SI) * 229.183
  7 WRITE(6,104)I, SINSQ, SII
    N = N + 200.
    REWIND 1
    IF (N.GT.2200.) STOP
    GC TO 1000
 20 FCFMAT(21X,F10.4,30X,F10.4)
102 FCRMAT(3)X,60A1,1X,I3,2X,1A1)
103 FOFMAT(1H1,//////,25X, SINE **2 THETA*,/)
104 FORMAT(14,2F10.4)
105 FCRMAT(43A1)
    END
```

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