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Cyanide-bridged Fe(III)-Co(II) bis double zigzag chains with a slow relaxation of the magnetisation[†]

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One of the fields developing rapidly in the area of molecular magnetic materials is the chemistry and physics of molecular systems with a slow relaxation of the magnetisation. Their interest is twofold: (a) fundamental, since classical properties of macroscopic magnets coexist with quantum effects (tunnelling of magnetisation, quantum phase interference;^{1,2} (b) applied, since they can lead to the ultimate density—molecular—of magnetic information storage and quantum calculations.³ They are often named, "single molecule magnets" (SMMs).⁴

A one-dimensional class, of slow relaxing magnetisation compounds appeared more recently, named either magnetic nanowires^{5,6} or single-chain magnets:⁷ one ferrimagnetic cobalt(II)–organic radical chain,⁵ two ferromagnetic bimetallic μ -cyano Fe(III)–Co(II) double chains,⁶ and a ferrimagnetic Mn(III)–Ni(II) chain.⁷ The anisotropy energy barrier is higher than in the SMMs, reaching values around 150 K,^{5–7} increasing the hope of magnetic information storage at reasonable temperatures.

Using the stable low spin iron(III) $[Fe^{III}(bipy)(CN)_4]^-$ as a ligand,⁸ we obtained two new isostructural bimetallic chains, $[{Fe^{III}(bipy)(CN)_4}_2M^{II}(H_2O)]$ ·MeCN·0.5H₂O [M = Co (1) and Mn (2); bipy = 2,2'-bipyridine], which result from the condensation of two double zigzag chains⁶ and whose structure and preliminary magnetic investigation are presented here.

Single crystals of **1** and **2** were grown in a MeCN : H_2O 90 : 10 (v/v) mixture by a slow-diffusion method using an Hdouble-tube glass vessel in the dark. The starting solutions were solutions of PPh₄[Fe(bipy)(CN)₄]·H₂O (0.05 mmol) in one arm and of Co(NO₃)₂·6H₂O (**1**) or Mn(NO₃)₂·4H₂O (**2**) (0.1 mmol) in the other one. After a few weeks orange plates (**1**) and red prisms (**2**) were formed together with an unidentified dark green powder. The crystals were collected and dried on filter paper. Yield: *ca.* 20% for **1** and **2**.[‡]

1 and 2 are isostructural.§ Their structure is made up of neutral one-dimensional [$\{Fe(bipy)(CN)_4\}_2M(H_2O)$] units running parallel to the *b* axis (Fig. 1, top) and crystallisation water and MeCN molecules. The crystallographically independent unit (Figs. S1 and S2)[†] contains two types of iron [Fe(1) and Fe(2)] and one cobalt [Co(1), 1] or manganese [Mn(1), 2] atoms, the latter being connected to five iron atoms through cyanide bridges (Fig. 1, bottom). This leads to a corrugated ladder-like chain with regular alternating Fe(1) and M(1) atoms along the edges. Each rung is defined by an Fe–M pair, where

† Electronic supplementary information (ESI) available: perspective views of the asymmetric units of **1** and **2** with the atom numbering (Figs. S1 and S2); a plot of the temperature dependence of $\chi_M T$ for **2** (Fig. S3); and a summary of the main bond lengths and angles for **1** and **2**. See http://www.rsc.org/suppdata/cc/b3/b302182n/

each pair of adjacent M atoms is connected through iron [Fe(2)]. A cyano group bridges each Fe–M pair. The two crystallographically independent iron atoms (Fe(1) and Fe(2)) exhibit the same distorted octahedral surrounding FeN₂C₄ as in previous reports.^{6,8} The difference between [Fe(1)(bipy)(CN)₄]⁻ and [Fe(2)(bipy)(CN)₄]⁻ is that the former acts as a trismonodentate ligand toward the M atoms through three cyanides in *fac* position whereas the latter adopts a bismonodentate coordination mode (two cyanides in *cis* position). The M atom is six-co-ordinated with five cyanide-nitrogen atoms and a water molecule forming a distorted MN₅O octahedron.

The value of the dihedral angle between adjacent mean planes of the corrugated ladder-like motif is 90° [angle at the Fe(1)…M(c) hinge in Fig. 1 (bottom)], whereas those between the Fe(1)MM(a) plane and Fe(1)M(a)Fe(1c)M(c) and Fe(1)M-Fe(1d)M(c) mean planes are 87 and 93° (1 and 2). The values of the Fe…M distances vary between 4.963(1) and 5.217(1) Å for 1 and 5.007(1) and 5.313(1) Å for 2. The shortest interchain metal-metal separations are 8.085(1) (1) and 8.090(1) Å (2) [Fe(2)…Fe(2e); (e) = 1 - x, -y, 2 - z]. The acetonitrile molecule is well isolated in the structure whereas hydrogen bonds link the coordinated and crystallization water molecules [O(10)…O(1) = 2.39(2) (1) and 2.43(2) Å (2)] and the



Fig. 1 Top: A view of the structure of **1** and **2** along the *b* axis. Solvent and hydrogen atoms are omitted. The metals are yellow (Fe) and pink [Co (**1**) or Mn (**2**)]. Bottom: Schematic view of the metallic frame of **1** (M = Co) and **2** (M = Mn). Selected bond lengths and angles are in Table S1 in ESI.†

crystallization water and one of the terminal cyanide nitrogens $[O(10)\cdots N(8b) = 3.00(2)$ (1) and 2.97(1) Å (2)].

The thermal variation of $\chi_{\rm M}T$ per Fe^{III}₂Co^{II} unit of a powder sample of **1** is shown in Fig. 2. At room temperature, $\gamma_{\rm M}T$ is equal to $5.0 \text{ cm}^3 \text{ mol}^{-1} \text{ K}$, which corresponds to one high spin Co(II) $(S_{Co} = \frac{3}{2})$ and two low spin Fe(III) $(S_{Fe} = \frac{1}{2})$ ions with significant orbital contributions. Upon cooling, $\chi_{\rm M}T$ increases smoothly in a monotonous way [in line with an intrachain ferromagnetic coupling between Co(II) and Fe(III)] until T < 50K where it sharply increases to reach a maximum of ca. 94 cm³ mol^{-1} K at 10 K and further decreases with T (as the magnetisation becomes field dependent). A maximum of susceptibility is observed around $\overline{7}$ K for applied fields lower than 600 G (Fig. 2, top, inset), a sign of a weak antiferromagnetic coupling between ferromagnetic entities. The maximum disappears for H > 600 G, suggesting a fieldinduced transition from an antiferromagnetic to a ferromagnetic ground state. This magnetic behaviour is consistent with the structure of 1. The sigmoidal shape of the *M* versus *H* plot at 2 K (Fig. 2, bottom) agrees with a metamagnetic behaviour. The magnetic energy for H = 600 G is *ca*. 0.06 cm⁻¹ and gives the order of magnitude of the antiferromagnetic interchain interaction. In an applied field above 600 G, the interchain interaction is overcome and 1 behaves as ferromagnetic Fe^{III}₂Co^{II} chains. The lack of λ -peak in the heat capacity measurements on powder samples of 1 in the temperature range 290–1.8 K both at H = 0 and 800 G, is a strong indication of the absence of three-dimensional magnetic order in 1. The magnetisation at saturation per Fe^{III}₂Co^{II} unit, at 2 K and 5 T, is $M_{\rm S} = 4.1 \ \mu B \ [ca. 1 \ \mu B \ per \ Fe^{\rm III}$ and about 2 \ \mu B \ per \ Co^{\rm II}



Fig. 2 Magnetic behaviour of **1**. Top: Temperature dependence of $\chi_M T$ per Fe^{III}₂Co^{II} unit under an applied magnetic field of 100 G. The solid line is an eye guide. The inset shows the field dependence of the magnetic susceptibility. Bottom: Magnetisation *vs. H* plot at 2 K. Inset: χ'' ac susceptibility [1 G oscillating field (0.1–1000 Hz), 800 G applied dc field].

(assuming an effective $S_{\rm Co} = \frac{1}{2}$ and $g_{\rm Co} = 4$)], in agreement with an intrachain ferromagnetic coupling. The magnetic behaviour of **1** above $H \ge 600$ G is similar to that of the Fe^{III}₂Co^{II} double zigzag chains,⁵ which exhibit slow magnetic relaxation and hysteresis effects not associated with a threedimensional magnetic order. The frequency dependence of the out-of-phase component (χ_M'') of the molar ac susceptibility of **1** under an applied field of 800 G is shown in Fig. 2. Between 2 and 7 K, the relaxation times follow an Arrhenius law [$\tau = \tau_0$ exp($E_a/k_{\rm B}T$)] with $E_a/k_{\rm B} = 152$ K and $\tau_0 = 1.5 \times 10^{-17}$ s. The energy barrier $E_a/k_{\rm B}$ is close to those found in the previous nanowires.^{5–7} Instead, the τ_0 prefactor is much smaller: obtained under an applied magnetic field, it deserves more study.

The magnetic behaviour of 2 (susceptibility and magnetisation at 2 K, Fig. S3) is the one of two ferrimagnetic Fe₂Mn chains coupled by a weak antiferromagnetic interaction.

The flexibility of the system allows the substitution of Fe and M by other metal ions, the replacement of bipy by other chelating ligands and the tailoring of new cyanide-bridged architectures and magnetic properties. Indeed, our new example of magnetic nanowire demonstrates the possibility of expanding the range of slow magnetic relaxing materials. The Glauber's model⁹ is probably too simple to fully describe the new magnetic behaviours.

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Notes and references

‡ Characterisation of 1 and 2: IR (solid, KBr pellets), v_{CN} 2160, 2149 and 2115 (1) and 2152, 2143 and 2114 cm⁻¹ (2). Elemental analysis (C, H, N) is in agreement with their formulae.

§ Crystal and structure refinement parameters for 1 and 2. Compound 1: C₃₀H₂₂Fe₂CoN₁₃O_{1.5}, *T* = 295 K, *P*2₁/*n*, *Z* = 4, *a* = 19.848(7) Å, *b* = 7.365(2) Å, *c* = 21.143(7) Å, *β* = 96.67(3)°, *V* = 3070(2) Å³, *d*_{calc} = 1.64 g cm⁻³, *R*₁ = 0.0367, *wR*₂ = 0.0449. Compound 2: C₃₀H₂₂Fe₂MnN₁₃O_{1.5}, *T* = 295 K, *P*2₁/*n*, *Z* = 4, *a* = 19.956(5) Å, *b* = 7.464(3) Å, *c* = 21.215(5) Å, *β* = 96.42(3)°, *V* = 3140(1) Å³, *d*_{calc} = 1.60 g cm⁻³, *R*₁ = 0.0406, *wR*₂ = 0.0494. Data were collected on a CAD-4 Enraf-Nonius diffractometer with graphite-monochromated Mo Kα radiation (λ = 0.71069 Å), and were corrected for Lorentz, polarisation and absorption effects. The structures were solved by direct methods through SHELXS-86 and subsequently refined by Fourier recycling. The residual maxima and minima in the final Fourier difference maps were 0.57 and -0.47 e Å⁻³ for 1 and 0.69 and -0.57 e Å⁻³ for 2. CCDC 204864 (1) and 204865 (2). See http:// www.rsc.org/suppdata/cc/b3/b302182n/ for crystallographic data in .cif format.

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