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## EFFECT OF $N_2H_4$ INTERCALATION ON THE TRANSITION TEMPERATURE AND ELECTRON TRANSPORT IN ANISOTROPIC SUPERCONDUCTOR: $TaSe_3$

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We have intercalated  $TaSe_3$  with  $N_2H_4$ . The room temperature resistance  $R_b(RT)$  along the  $b$ -axis decreases by 30% during the reaction. The ratio of  $R_b(RT)/R_b(4.2\text{ K})$  varies from 60 in  $TaSe_3$  to 5 in the intercalated samples. An anomaly in  $R_b$  is observed near 100 K. It is suggested that this anomaly is due to the formation of charge-density-wave (CDW). The superconducting transition temperature,  $T_c$ , is suppressed from 2.2 to 1.5 K. The correlation between the suppression of superconductivity and the formation of CDW is discussed as compared with the result obtained in the layered dichalcogenides.

### 1. Introduction

The trichalcogenides of niobium and tantalum,  $MX_3$ , have been investigated extensively because of their low-dimensional characters. Although their crystal structures are similar, their electron transport properties show significant differences [1].  $NbSe_3$  shows two charge-density-wave (CDW) transitions at 142 and 85 K [2] and  $TaS_3$  shows Peierls transition at 220 K [3]. No anomalies associated with low-dimensional phase transition have been observed in  $TaSe_3$  [1]. In fact,  $TaSe_3$  is a superconductor with a  $T_c$  of  $(2 \pm 0.5)\text{K}$  [1, 4]. The reasonable interpretation for the absence of the structural phase transition in  $TaSe_3$  is that the interchain overlap of electron wave functions is larger than that of other  $MX_3$  compounds.

As shown in fig. 1,  $TaSe_3$  has a structure made up of infinite chains of trigonal prisms,  $TaSe_6$ , extending parallel to the  $b$ -axis. Ta in  $TaSe_3$  is surrounded by eight Se atoms forming a bicapped trigonal prism. According to this point of view we have to consider layers made of coupled  $TaSe_3$  fibers. Between them lies a Van der Waals gap with dimensions close to those of the layered dichalcogenides,  $MX_2$ .  $TaSe_3$ , therefore, can be considered as a two-dimensional system with a

large anisotropy rather than a one-dimensional system.

$MX_2$  is well known to be intercalated with organic molecules and hydrogen [5, 6]. The enhancement of  $T_c$  and the suppression of the CDW formations were observed in many intercalated compounds of  $MX_2$  [5, 6]. It is very interesting to investigate the intercalation effect on  $T_c$  and the electron transport in  $TaSe_3$  and to compare them with the effect in  $MX_2$ . The measurements of X-ray reflecting power, elec-

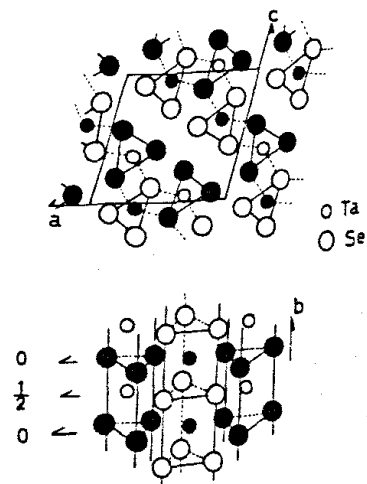


Fig. 1. The crystal structure of  $TaSe_3$ . The small (large) circles identify Ta (Se) atoms.

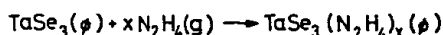
nal resistivity, and  $T_c$  in TaSe<sub>3</sub>, when this is exposed to N<sub>2</sub>H<sub>4</sub>, are reported in this paper.  $T_c$  is found to be suppressed on intercalation of N<sub>2</sub>H<sub>4</sub> and an anomaly in the resistivity is observed near 100K. The correlation between the suppression of  $T_c$  and an anomaly in resistivity is compared with that of intercalation on MX<sub>2</sub> and discussed.

2 Experiments

TaSe<sub>3</sub> were grown by the transport method, as reported in [4]. The resistance ratios, RRR,  $R(\text{room temperature})/R(4.2\text{K})$ , and  $T_c$  of the crystals were confirmed to be the same values as observed in previous works [1, 4]. The resistivity of samples was measured by a four-probe method along the *b*-axis with a dc current. Superconductivity was detected by the resistivity measurements. The X-ray diffraction and transport measurements of TaSe<sub>3</sub> exposed to N<sub>2</sub>H<sub>4</sub> vapor in the absence of air were carried out using the method already developed for TaS<sub>2</sub> [7]. The measurements of the temperature dependence of the resistivity and  $T_c$  in the samples were carried out under He gas atmosphere to avoid the reaction with air.

1 Result and discussion

X-ray diffraction by the as-grown TaSe<sub>3</sub> was consistent with the reported structure [1]. The reflecting power of (400) plane in the TaSe<sub>3</sub> disappears on exposure to N<sub>2</sub>H<sub>4</sub> with 12 Torr at 24°C for 90 min. This suggests that N<sub>2</sub>H<sub>4</sub> must be intercalated in between the layers along the diagonal of the *a*- and *c*-axes. The stoichiometry of the complex formed TaSe<sub>3</sub>(N<sub>2</sub>H<sub>4</sub>)<sub>x</sub> was determined by weight gain measurements [8]. The dependence of the rate of weight gain vs. pressure of N<sub>2</sub>H<sub>4</sub> is shown in fig. 2. Different kinetics for the pressure of N<sub>2</sub>H<sub>4</sub> suggests that the conditions of the intercalation in the low pressure regions are more gentle than those in the high pressure regions. Within minutes of exposure to H<sub>2</sub>H<sub>4</sub>,  $x = 1/3$ , but after 48 h under 12 Torr of N<sub>2</sub>H<sub>4</sub> the value of  $x$  is equal to one. When the sample is evacuated,  $x$  is reduced to 1/3 immediately and after 48 h under 10<sup>-5</sup> Torr of N<sub>2</sub>H<sub>4</sub>,  $x = 1/3$ . However, X-ray diffraction by the



$$\left(\frac{\partial x}{\partial t}\right)_p = k(x_\infty - x)$$

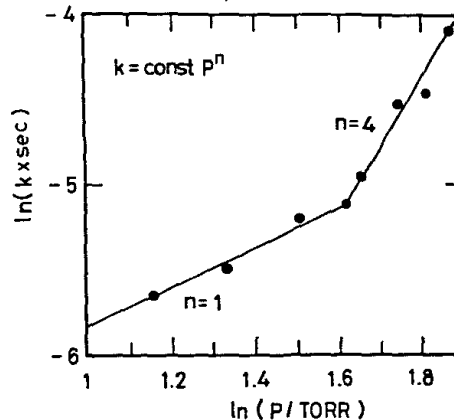


Fig. 2. Kinetics of intercalation vs. pressure of N<sub>2</sub>H<sub>4</sub>.  $x$  is the ratio of moles of N<sub>2</sub>H<sub>4</sub> to the moles of TaSe<sub>3</sub>.  $x_\infty$  is the extrapolated value as time,  $t, \rightarrow \infty$ . Here the same sample was intercalated and deintercalated several times for a given pressure and each point indicates an average.

intercalated sample, TaSe<sub>3</sub>(N<sub>2</sub>H<sub>4</sub>)<sub>x</sub>, does not allow a structure determination because the lines are very broad, as shown in fig. 3. This suggests the formation of complex superstructures.

On exposure to N<sub>2</sub>H<sub>4</sub> vapor (12 Torr, 24°C) the resistance decreases with the time to a constant value 30% less than the initial value. This value remained when the pressure was reduced to 10<sup>-4</sup> Torr. The result of the change in the resis-

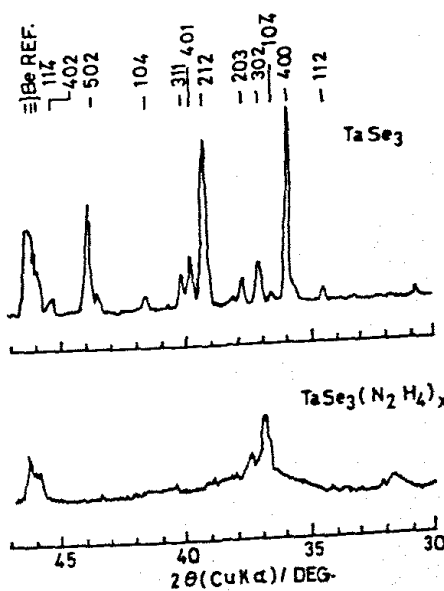


Fig. 3. X-ray reflecting power of TaSe<sub>3</sub> and its intercalated sample, TaSe<sub>3</sub>(N<sub>2</sub>H<sub>4</sub>)<sub>x</sub>.

tance is consistent with that in X-ray diffraction by the intercalated samples. The temperature dependence of the resistivity has been measured on four intercalated samples. Two of them (A and B) have been intercalated under a pressure of N<sub>2</sub>H<sub>4</sub> at 12 Torr and other samples (E and F) have been intercalated under a pressure of N<sub>2</sub>H<sub>4</sub> at 1 Torr. As shown in fig. 4, the dependence of the resistivity vs. temperature in TaSe<sub>3</sub>(N<sub>2</sub>H<sub>4</sub>)<sub>x</sub>-F is almost the same as that of TaSe<sub>3</sub>, i.e., resistance decreases monotonically as temperature decreases and shows negative curvature. The same result is also obtained for E. Also, the onset temperature for superconducting transition is the same as that of TaSe<sub>3</sub>, as shown in fig. 5. Thus, no significant changes in the resistivity and  $T_c$  are observed in the E- and F-samples exposed to N<sub>2</sub>H<sub>4</sub> under low pressure. On the other hand, the temperature dependence of resistivity in TaSe<sub>3</sub>(N<sub>2</sub>H<sub>4</sub>)<sub>x</sub>-B intercalated under the pressure of N<sub>2</sub>H<sub>4</sub> at 12 Torr, which is the same as that of the A-sample, is quite different from that obtained in TaSe<sub>3</sub> [1], as shown in fig. 6. The main differences are as follows. (1) An anomaly near 100 K is observed in the intercalated samples. (2) The positive curvature against tem-

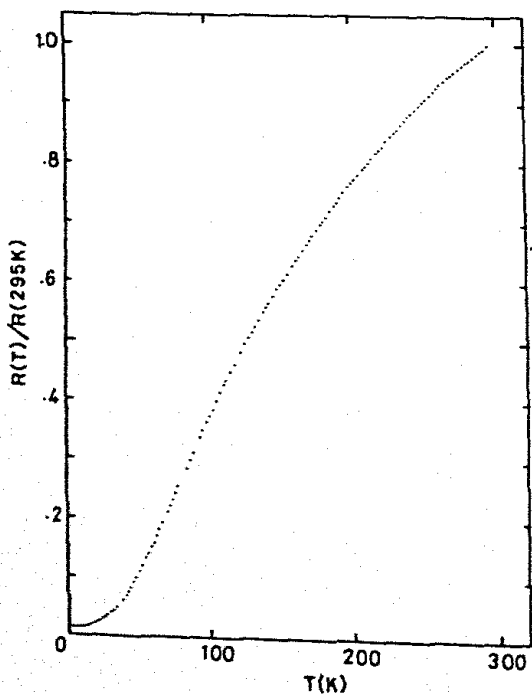


Fig. 4. The normalized resistance as a function of temperature in TaSe<sub>3</sub>(N<sub>2</sub>H<sub>4</sub>)<sub>x</sub>-E.

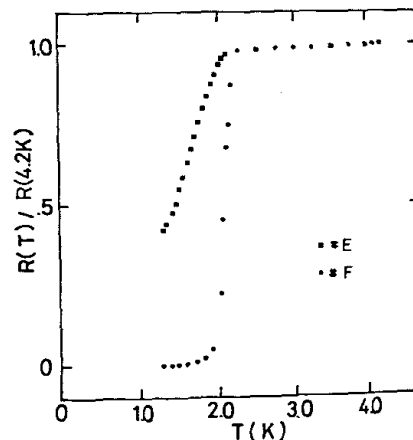


Fig. 5. The superconducting transition curves of TaSe<sub>3</sub>(N<sub>2</sub>H<sub>4</sub>)<sub>x</sub>-E and -F.

perature in resistivity vs. temperature is observed above 120 K, in contrast with the negative curvature obtained in TaSe<sub>3</sub> [1]. (3) The value of RRR varies from 60 to 5. Furthermore,  $T_c$  of the intercalated samples are observed to decrease from 2.2 to 1.5 K, as shown in fig. 7. The resistivity anomaly observed here is similar to that associated with the CDW formation observed in dichalcogenides and their intercalated com-

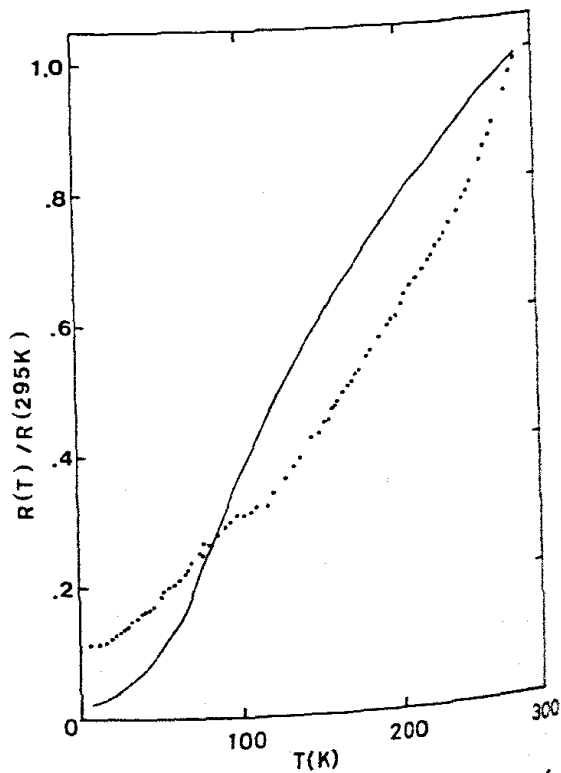


Fig. 6. The normalized resistance as a function of temperature in TaSe<sub>3</sub>(N<sub>2</sub>H<sub>4</sub>)<sub>x</sub>-B. Solid line shows the temperature dependence of the resistance in TaSe<sub>3</sub>.

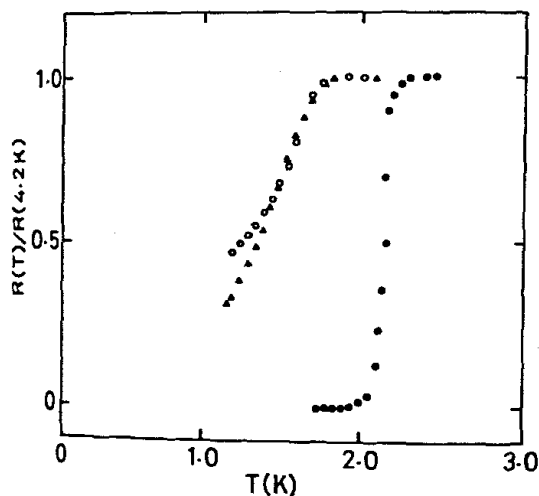


Fig. 7. The superconducting transition curves of TaSe<sub>3</sub>(N<sub>2</sub>H<sub>4</sub>)<sub>x</sub>-A (○); -B (▲); and TaSe<sub>3</sub> (●).

pounds [5-7]. The interlayer overlap of electron wave function in TaSe<sub>3</sub> can be expected to be reduced by the intercalation. Thus, the resistivity anomaly observed in the intercalated samples (A and B) suggests that CDW are produced and that these in turn suppress the superconductivity. This result is consistent with the fact observed in MX<sub>2</sub> in which the enhancement of superconductivity leads to suppression of the CDW formation. However, much more work, such as an observation of the formation of periodic lattice distortions and measurements of Hall effect and magnetic susceptibility, will be necessary to confirm the CDW formation in TaSe<sub>3</sub>(N<sub>2</sub>H<sub>4</sub>)<sub>x</sub>. Also, it is important to investigate if the anomaly observed here is due to the weakly pinned CDW which had been already observed in NbSe<sub>3</sub> [9].

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