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Abstract

We demonstrate experimentally that magnetism in ZnO, TiO₂, CeO₂, and SnO₂ nanocrystals (NCs) has a fluctuant nature that varying with capping surfactant type and concentration. By developing a forced hydrolysis approach with additional post-processing for the synthesis and surfactant capping of these NCs, we effectively avoid the influence of size, shape, and magnetic impurities on the magnetic behavior of NCs, thus revealing the systematic influence of the capping surfactants on the NC magnetism. The x-ray photoelectron spectroscopy results and theoretical calculations clearly show that the magnetism fluctuation with surfactant concentration dependent electron transfer from surfactants to NCs. Our results not only explain the previously reported seemingly irregular magnetism induced by capping surfactants, but also provide an effective approach to tune or optimize the NC magnetism.

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I. INTRODUCTION

Developing ferromagnetism in semiconducting metal oxide nanocrystals (NCs) has been intensively investigated due to their potential applications in several areas such as spintronics, information storage, and biomedicine.¹⁻¹⁴ Ferromagnetism (FM) can be induced in non-magnetic metal oxide NCs by a variety of methods or factors. For example, intrinsic defects such as oxygen or cation vacancies in pure metal oxides (such as ZnO, CeO₂, or SnO₂) can lead to room temperature (RT) FM.²⁻⁵ Defects intentionally induced by doping nonmagnetic elements (such as Ti and V) into metal oxides can also produce RT FM.⁶ A more traditional method is doping magnetic impurities (such as Co, Ni, Mn, and Cr) to introduce magnetic moments directly.⁷⁻⁹ Recently, it was found that capping metal oxides with inorganic (such as NH_3^{10}) or organic (such as dedecanethiol, THIOL¹¹) molecules can effectively induce FM via changing electronic structure. This approach might be promising for synthesis of metal oxide semiconductor NCs with optimized FM because there are numerous molecules to be chosen for capping and modifying the oxide cores. Moreover, it can avoid the problem of clustering of magnetic elements.¹² However, the influence of the capping surfactants on magnetism has been found to vary with surfactants, and show inconsistent results in some similar systems. For example, Gamelin et al. reported no RT FM in dodecylamine (AMINE)-capped ZnO NCs,13 whereas Jiang et al. observed FM in AMINE-capped ZnO NCs and revealed FM to arise from the capping AMINE.¹⁴ Furthermore, Garcia *et al.* found the saturation magnetization (M_s) of undoped ZnO NCs to increase with the capping surfactants in the order of tryoctylphosphine oxide (TOPO) < AMINE < THIOL.¹¹

Although it is not completely clear how the capping molecules affect the magnetism of the metal oxide NCs, electron configuration variations induced by capping molecules have been considered to be a crucial factor.^{10,11,14} Besides the surfactant type, the electron configuration is also expected to vary with the capping surfactant concentration (SC). Namely, the capping SC might be able to change the NC magnetism, thus leading to the seemly irregular magnetism variations described above. However, until now, systematic studies of the influences of the capping SC on magnetism of metal oxide NCs have been very limited. It is possible that capping NCs with surfactants is often accompanied with size and shape variations of NCs and defect introduction, and therefore only a combined effect is observed.^{13,15,16} To reveal how the capping surfactants affect the magnetism of the oxide NCs, a forced hydrolysis approach with additional post-processing for the synthesis and surfactant capping of ZnO, TiO₂ SnO₂, and CeO₂ NCs has been developed here. This ensures that the size and shape of NCs remain unchanged and only the systematic influence of the surfactants on the NC magnetism is investigated. Our results clearly show that the NC magnetism is intimately related to both the type and amount of capping surfactants, and can be tuned simply by modifying their capping SC.

II. EXPERIMENTAL CONDITIONS

All the analytical reagent grade chemicals were used as received. For the synthesis of ZnO NCs, 400 mL of DEG containing 8 g of $Zn(Ac)_2 \cdot 2H_2O$ was heated to 180°C under stirring for 30 min. The resulting NCs were washed with ethanol 3 times via centrifugation and sonication. For the syntheses of TiO₂, SnO₂, CeO₂ NCs, nanopure water (4 mL for TiO₂, 8 mL for SnO₂ and CeO₂) was added to DEG (400 mL) containing 8 mL of Ti(OCH(CH₃)₂)₄ for TiO₂, or 8 g of Sn(Ac)₄ (for SnO₂) or Ce(Ac)₃·1.5H₂O (for CeO₂) at 80 °C, and then the system was heated at 180°C for 1 h. The resulting product was also washed with ethanol 3 times via centrifugation and sonication. The as-prepared product was dried at 50°C to obtain a powder sample, or dispersed in nanopure water (80 mL) for post-treating with surfactants. CeO₂ NCs were dried at 80 °C for 12 h before the surfactant capping.

To ensure that all the NCs capped with different surfactants are identical in their initial crystal quality, size and shape, thus ensuring that we are effectively studying only the influence of the surfactants on the NC magnetism, a surfactant post-treating method using the same NC source was developed. Briefly, the as-prepared ZnO NCs were dispersed in 80 mL of nanopure water to produce a stock suspension. An ethanol solution (20 mL) of surfactants (75 mM) was used to suspend the ZnO NCs that were fabricated by centrifuging 10 mL of stock suspension. After 50

min of sonication, the product was washed with ethanol twice and dried at 50 °C to obtain a powder sample. A small concentration (5 mM) was used for polyvinylpyrrolidone (PVP) due to its large molecular weight ($M_w = 58,000$) which makes it hard to prepare a high-concentration solution. An aqueous solution (20 mL) of NaCl (1 mM) and poly (acryl acid) (PAA, $M_w = 1800, 10 \text{ mg/mL}$) followed by twice washing with water was used to cap PAA.

The product morphology was examined using a JEOL JEM-2100HR transmission electron microscope (TEM) with an accelerating potential of 200 kV. Cryststallographic phases of the samples were studied by x-ray diffraction (XRD) using a Philips X'pert MPD diffractometer with a Cu K_{α} x-ray source in Bragg-Brentano geometry. X-ray photoelectron spectroscopy (XPS) data were recorded by a Physical Electronics Versa probe using a raster-scanned microfocused monochromatic Al K_{α} X-ray (1486.7 eV) source. High resolution XPS spectra were obtained with a pass energy of 23.5 eV and at a take-off angle of 45°. The trace magnetic impurity in samples was detected by an inductively coupled plasma mass spectrometry (ICP-MS, PQEXCe II). The magnetization hysteresis of the dried powder samples was acquired at RT using a vibrating sample magnetometer. The noise floor of the VSM magnetometer is 0.1microemu for a collection time of 1 second/point and 0.75 microemu for a collection time of 0.1s/pt, which ensures that the weak magnetic signal down to 1 microemu in our samples can be precisely measured.

III. RESULTS AND DISCUSSION

The TEM images disclose that both ZnO NCs post-treated with TOPO (Fig 1a) and PAA (Fig. 1b) show short nanorod-like shapes with an average diameter of ~9 nm and length of ~15 nm. The corresponding XRD patterns (Fig. 1c) indicate that all the ZnO NCs post-treated with different surfactants have the same pure wurtzite ZnO phase and a similar average size (~9.2 nm, obtained using Scherrer formula), which agrees well with the TEM observations. The results show that the capping surfactants do not affect the ZnO NC size and shape, and therefore these samples really allow a systematic and effective study of the capping-surfactant influence on the ZnO NC magnetism. As seen in Fig. 1d, all the ZnO NCs treated with different surfactants show a weak FM, which does not arise from the impurity in surfactants or the measurement artifacts because all the pure surfactants used here show their diamagnetic nature under the same measurement conditions (Fig. S1).¹⁷ The NC M_s increases with capping surfactants in the order of THIOL \leq PAA \leq PVP < AMINE < CTAB < TOPO, which is different from the results reported earlier.¹¹ As addressed below, this may be ascribed to the higher SC used here to cap NCs, because the capping SC can greatly affect the NC magnetism.

TOPO, PVP, and PAA were selected for a systematic study of the influence of capping SC on the ZnO NC magnetism because capping by these molecules leads to the maximum or minimum M_s . The high resolution XPS results clearly show that the actual surfactant amount deposited onto ZnO NCs is proportional to the capping SC. As seen in Table SI,¹⁷ with increasing SC, the number of surfactant molecules adsorbed onto per ZnO NC increases, but tends to be saturated at high SC, which may be ascribed to the reduction of the surface Zn atoms of NCs available for the surfactant bonding, and the steric hindrance effect among surfactant molecules. According to the calculation results based on the NC surface area, surfactant molecule size, and the adsorbed surfactant molecule number per NC, it is deduced that only a monolayer of surfactant molecules is formed on the ZnO NCs after treating with all the surfactants. As shown in Fig. 2a, with increasing SC, both TOPO and PVP capping initially increases the saturation magnetization M_s of ZnO NC abruptly by a factor of about 7, followed by a dramatic decrease of M_s . With further increasing SC, the M_s variation trend repeats itself, but at a much lower frequency. The ZnO NC magnetism also varies with the capping-PAA concentration in a similar manner, but its variations induced by the capping of high concentrations of PAA (which will apparently dissolve and change the ZnO NCs) have not been investigated here. As shown in Tables SI-II,¹⁷ the maximum adsorbed surfactant/ZnO weight ratio and magnetic impurity (including Cr, Mn, Fe, Co, and Ni) concentration in NCs are below 3.5 % and 3.5 ppb, respectively. Therefore, the influence of the small difference in the adsorbed-surfactant weight and the negligible amount of magnetic impurities on the above large FM variations can be excluded here.

To check whether the above magnetism variation with capping SC is a general trend among metal oxide NCs, we also investigated the SC influences of TOPO and PAA on the magnetism of SnO₂, TiO₂, and CeO₂ NCs made in a similar way as the ZnO NCs. TEM and XRD results (Figs. S2-4)¹⁷ show that the as-prepared SnO₂, TiO₂, and CeO₂ NCs have the short nanorod shapes with an average diameter of ~4 nm, and high-purity tetragonal rutile, anatase, and fluorite structures, respectively. The capping of both TOPO and PAA has no detectable influences on their crystalline phase, size, and shape, thus excluding the effect of factors other than the concentration of the capping surfactant on the NC magnetism. As expected, FM in SnO₂, TiO₂, and CeO₂ NCs is also greatly enhanced after the

initial capping of TOPO and PAA, and then fluctuates with increasing SC as illustrated in Fig. 2b-c and Fig. 3. Like in the case of ZnO NCs, the undoped and uncapped SnO₂, TiO₂, and CeO₂ NCs also show very weak FM after subtracting the diamagnetic background, which might arise from intrinsic defects such as oxygen vacancies.^{2,3,5} The strongest FM ($M_s = 8.7$ memu g⁻¹) is observed in SnO₂ NCs treated with 40 mg mL⁻¹ of TOPO, while the largest FM enhancement (~40 fold) is found in TiO₂ NCs treated with 10 mg mL⁻¹ of PAA.

The above magnetism fluctuation with SC can be ascribed to an alteration of electronic configuration^{10,13,14} induced by the electron transfer between NCs and surfactants,¹⁸⁻²⁰ which is revealed by the XPS spectra. The typical XPS spectra of ZnO and CeO2 NCs before and after post-treating with TOPO, PVP, and PAA are shown in Fig. 4, and Fig. S5,¹⁷ respectively. Treating ZnO NCs with TOPO, PVP, or PAA, apparently shifts the Zn 2p_{3/2} XPS peak to a low binding energy (Fig. 4), indicating that ZnO receives electrons from the surfactants.²⁰ Similarly, the TOPO and PAA molecules that are linked to the CeO₂ NCs also donate electrons, and this leads to the appearance of Ce³⁺ XPS peaks²¹(Fig. S5). As shown in Fig. 4, with the increase of the capping SC, the Zn $2p_{3/2}$ XPS peak of ZnO NCs continuously shifts to a lower binding energy, but the relative shift gradually slows down. This is consistent with the expected variations in the number of the adsorbed surfactant molecules per ZnO NC as SC increases. The relationship between the electron transfer and the NC magnetism is clearly demonstrated in Fig. 5, showing the ZnO NC M_s plotted as a function of the Zn $2p_{3/2}$ XPS peak shift induced by the capping of TOPO, PVP, and PAA at different concentrations. As shown in the figure, with the increase of the XPS peak shift to lower binding energ, (i.e., the number of electrons transferred), the NC M_s initially increases to a maximum followed by decreasing to a minimum, and then repeats the variation trend as the XPS peak shift increases further. These results clearly reveal the fluctuation nature of the NC magnetism with the number of the extra electrons donated by the capping surfactants. We also noted that the fluctuation rate of the NC magnetism with the SC or the number of electrons transferred varies with surfactants. This can be attributed to the variation in the interaction strength or bonding structure between NCs and surfactants, depending on the capping molecules. As shown in Figs. 4-5, compared with the functional O atoms of PAA, the functional P or N atoms of TOPO or PVP have relatively lower electronegativity and donate electrons to NCs more easily, thus leading to a larger Zn $2p_{3/2}$ XPS peak shift towards low binding energies. The different NC-surfactant interactions should affect the kinetic energy as well as the interaction or coupling among the extra surface electrons donated by the capping surfactants,^{18,19} thus lead to the different fluctuation rates of the NC magnetism with SC.

To explain the abovementioned fluctuation of the magnetism of NCs with SC theoretically, we took a ZnO NC with an adsorbed TOPO molecule as an example to carry out the first-principle calculations and structure optimization. The first-principle calculations were performed using the generalized gradient approximation (GGA) of the Perdew, Burke, and Ernzerholf (PBE) form²² under package CASTEP,²³ in which a plane-wave norm-conserving pseudopotential method²⁴ was used. We used a kinetic energy cutoff of 400 eV to represent the single-particle wave functions. The NC had a diameter of 9.5 Å, and adsorbed a TOPO molecule. To form a periodic structure for the calculation with CASTEP package, the unit cells of size $25 \times 20 \times 16$ Å³, each of which containing a NC, a TOPO molecule, and vacuum space, were periodically placed along three directions. The geometry of the configuration was optimized using BFGS minimizer in the CASTEP package with default convergence tolerances: 2×10^{-5} eV for energy change, 0.05 eV/Å for maximum force, and 0.002 Å for maximum displacement.²⁵ Fig. 6a shows the optimized structure of the NC with an adsorbed TOPO molecule and the spatial distribution of electron density. It can be seen that with increase in the NC-TOPO interaction, several oxygen atoms in the NC are pulled towards the TOPO molecule. These oxygen atoms carry a dense electron cloud due to the charge transfer from the TOPO molecule. By assigning the charge to every atom, we found the average charge of an oxygen atom in the NC to be -0.787e (e is the elementary charge), while that of a zinc atom was only 0.657e. So, even though there are 24 Zn atoms and only 21 O atoms in the NC, the total charge in the NC is still negative, -0.759e, indicating the net electron transfer from TOPO molecule to NC. This is consistent with our experimental result shown in Fig. 4a, where the shift of the Zn 2p_{3/2} XPS peak towards low binding energies induced by the TOPO capping indicates the transfer of electrons to ZnO NCs from TOPO.

It was shown theoretically that the total orbital momentum of the Fermi level in non-magnetic NCs can oscillate with increasing the number of extra electrons donated by the capping surfactant, with a maximum occurring at one-half occupation.^{18,19} The above magnetism fluctuation of NCs with SC may arise from the variations of the number of electrons transferred to NCs from the capping surfactants. To confirm this, we performed the first-principle calculations of the spins in NCs with different numbers of extra electrons. At this stage we did not include the surfactant molecules so that we could deal with larger NC clusters, 14 Å in diameter. The effect of

electron transfer between NCs and surfactant molecules was simulated by setting different charges of NCs. In Fig. 6b-f, we plot the integrated densities I(E) of spins for NCs with different charges Z, where I(E) is defined as

$$I(E) = \int_{-\infty}^{E} d\varepsilon \left[\rho_{\uparrow}(\varepsilon) - \rho_{\downarrow}(\varepsilon) \right] \text{ with } \rho_{\uparrow}(\varepsilon) \text{ and } \rho_{\downarrow}(\varepsilon) \text{ being the partial densities of states of spin up and}$$

down electrons in the NC, respectively. So I(E) corresponds to the total spin of electrons with energies below E in the NC. The value of I(E) at the Fermi energy E = 0 is just the total spin of electrons as the states with energies below zero are occupied. Fig. 6g shows the variation of I(E = 0) with charge number. As expected, the charge variation in NCs not only shifts the Fermi level, but also largely alters the density of spin states. With increasing the charge from zero, the NC spin becomes nonzero and the magnitude oscillates. These calculation results explain the general trend of the measured NC magnetization variations with SC.

IV. CONCLUSIONS

In summary, we have developed a forced hydrolysis approach with additional post-processing to synthesize and coat uniform and size-homogeneous ZnO, TiO₂, CeO₂, and SnO₂ NCs with different types and concentrations of surfactants in order to investigate the systematic influence of the capping-surfactant on NC magnetism. As confirmed by the TEM, XRD, XPS, and ICP-MS results, the influence of the NC size and shape, and any role of magnetic impurities such as Cr, Mn, Fe, Co, and Ni on the magnetic properties of NCs have been effectively ruled out. Thus, the magnetism fluctuation with the capping-surfactant type and/or concentration in these NCs has been demonstrated. The XPS results clearly show that the amount of electron transfer from the capping surfactants to NCs is dependent on the SC-, and this causes the NC magnetism to fluctuate. The theoretical calculations confirm the electron transfer from the capping molecules to NCs, and clearly show that the electron transfer to NCs can lead to a periodic variation of spins, thus the observed NC magnetism fluctuation. Our results reasonably explain the seemingly irregular magnetism induced by capping surfactants in previous reports, and provide an effective approach to tune or optimize the NC magnetism. Compared with modifying the NC magnetism by introducing defects and/or dopants, which depends on the preparation conditions and procedures, we find that varying the surfactant type/concentration via post-processing to be a simple and highly reproducible method of tuning the NC magnetism.

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- ¹⁷See supplementary material at [*URL will be inserted by publisher*] for Fig. S1 (hysteresis loops of pure CS), Figs. S2-4 (TEM and XRD results of SnO₂, TiO₂, and CeO₂ NCs before and after CS capping), Fig. S5 (XPS spectra of
- CeO₂ NCs treated by CS), Table SI (Carbon/Zn atomic ratio, adsorbed surfactant/ZnO weight ratio, and adsorbed surfactant molecule number per NC of ZnO NCs treated by CS), Table SII (ICP-MS results of ZnO CNs treated by CS).
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Figure captions

Color online FIG. 1. Typical TEM images of ZnO NCs post-treated with TOPO (a) and PAA (b); XRD patterns and hysteresis loops of ZnO NCs treated with different surfactants are shown in (c) and (d), respectively. The diamagnetic background of the measurement system including the plastic capsules used to pack samples in (d) was subtracted for all the curves, and 1memu g^{-1} is defined as 10^{-3} emu g^{-1} .

Color online FIG. 2. (a) M_s from the ZnO NCs as function of the capping-surfactant concentration. (b) M_s from the SnO₂, TiO₂, and CeO₂ NCs as functions of the capping-TOPO concentration. For clarity, the M_s of CeO₂ NCs has been multiplied by a factor of 20. (c) Typical hysteresis loops of SnO₂ NPs treated by TOPO at different concentrations. The diamagnetic background was subtracted for all the curves.

Color online FIG. 3. Typical hysteresis loops of SnO_2 (a), TiO_2 (b), and CeO_2 (c) NPs treated by PAA at different concentrations. The diamagnetic background was subtracted for all the curves. (d) M_s as function of the capping-PAA concentration.

Color online FIG. 4. Zn $2p_{3/2}$ XPS spectra of the selected ZnO NC samples treated by TOPO (a), PVP (b), and PAA (c) with different concentrations. The binding energy scales have been charge-referenced using the C 1s peak at 284.8 eV.

Color online FIG. 5. ZnO NC M_s versus the Zn $2p_{3/2}$ XPS peak shift induced by the capping of TOPO, PVP, and PAA at different concentrations.

Color online FIG. 6. (a) Optimized structure and electron density distribution of a ZnO NC with an adsorbed TOPO molecule, taken from the first-principle calculation. The gray-blue, red, violet, gray, and white balls represent Zn, O, P, C, and H atoms, respectively (for print in grayscale: the big black, small black, big grey, moderate size black, and small white grey balls represent Zn, O, P, C, and H atoms, respectively). (b)-(f) Integrated density of spin of a ZnO NC with different transferred charges Z given in the legend. The Fermi level is at the zero energy. (g) The variation of spin in the NC with transferred charge.







Figure 3, Zhang et al.



Figure 4, Zhang et al.

