

## CORE AND PATCH POSITION OPTIMIZATIONS FOR VIBRATION CONTROL OF PIEZOLAMINATED STRUCTURES

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### 1. Abstract

This paper deals with a finite element formulation based on the classical laminated plate theory, for active control of thin plate laminated structures with integrated piezoelectric layers, acting as sensors and actuators. The control is initialized through a previous optimization of the core of the laminated structure, in order to minimize the vibration amplitude. Also the optimization of the patches position is performed to maximize the piezoelectric actuator efficiency. The simulating annealing method is used for these purposes. The finite element model is a single layer triangular nonconforming plate/shell element with 18 degrees of freedom for the generalized displacements, and one electrical potential degree of freedom for each piezoelectric element layer, which can be surface bonded or imbedded on the laminate.

To achieve a mechanism of active control of the structure dynamic response, a feedback control algorithm is used, coupling the sensor and active piezoelectric layers. To calculate the dynamic response of the laminated structures the Newmark method is considered. The model is applied in the solution of an illustrative case and the results are presented and discussed.

**2. Keywords: Active Control, Sensors and Actuators, Adaptive Laminated Structures, Structural Optimization.**

### 3. Introduction

Advanced reinforced composite structures incorporating piezoelectric sensors and actuators are increasingly becoming important due to the development of adaptive structures. These structures offer potential benefits in a wide range of engineering applications such as vibration and noise suppression, shape control and precision positioning.

In the analysis of adaptive structures integrating piezoelectric material, one of the pioneering works is due to Allik and Hughes [1] who developed a solid finite element for vibration analysis. Vibration control of composite beams with embedded or surface bonded piezoelectric material had been studied by Crawley and de Luis [2]. Tzou and Tseng [3] presented a finite element formulation for plates and shells containing integrated distributed piezoelectric sensors and actuators applied to control advanced structures. Chen et al. [4] developed a finite element based on the first order displacement field for dynamic analysis of plates. Active control is obtained through actuators potential, which is given by an amplified signal of the sensors potential.

Using a higher order shear deformation theory Samanta et al. [5] developed an eight-noded finite element for the active vibration control of laminated plates with piezoelectric layers acting as distributed sensors and actuators. The active control capability is studied using a simple algorithm with negative velocity feedback. Lam et al. [6] and Moita et al. [7] developed finite element models based on the classical laminated theory for the active control of composite plates containing piezoelectric sensors and actuators using the Newmark method [8] to calculate the dynamic response of laminated structures. Reviews of the modelling and the design of composite structures with adaptive capabilities are given in Franco Correia et al. [9] and Benjeddou [10].

Batra and Liang [11] used a three-dimensional linear theory of elasticity to find the optimal location of an actuator on a simple-supported rectangular laminated plate with embedded PZT layers. The optimal design is obtained by fixing the applied voltage and the size of the actuator and moving it around in order to find the maximum out-of-plane displacement. Liang et al. [12] proposed a model for the optimization of the induced-strain actuator location and configuration for active vibration control. Correia et al. [13] presented refined finite element models based on higher order displacement fields applied to the optimal design of laminated composite plates with embedded or surface bonded piezoelectric actuators and sensors.

Most of past work in the area of adaptive structures has focused on the analysis of structures with sensors and actuators, and the corresponding associate control system. Very few works have focused on the development of methodologies for the optimization of laminated structures incorporating sensors and actuators, to enhance their performance. A review of the current work in design methodologies and applications of formal optimization methodologies for adaptive structures has been carried out by Padula [14] and Frecker [15].

The finite element used in the present work is a flat three-noded triangular element with 18 mechanical degrees of freedom and one electric degree of freedom per piezoelectric layer of the finite element and is based on the classical laminated plate theory. An integrated control is considered in order to achieve an active damping mechanism, with the amplified electrical potential of the sensors being used as electric potential input to the actuators through a negative velocity feedback control. The governing equations are solved by the Newmark method, which is a direct method for time integration. Prioity to the dynamic control, to minimize the (forced) vibration amplitude through active damping, an optimal lamination sequence of the structure core, as well as an optimal placement of piezoelectric actuators patches are carried out using a structural optimization methodology based on simulated annealing algorithm [16].

### 4. Displacement and strain fields.

The assumed displacement field is that of the classical laminated plate theory. The displacement components referenced to the local axis  $(x,y,z)$  of the laminate are:

$$\begin{aligned}
u(x, y, z) &= u_0(x, y) - z \theta_y(x, y) \\
v(x, y, z) &= v_0(x, y) + z \theta_x(x, y) \\
w(x, y, z) &= w_0(x, y)
\end{aligned} \tag{1}$$

where  $u_0$ ,  $v_0$ ,  $w_0$  are displacements of a generic point in the middle plane of the laminate referred to the local axes -  $x, y, z$  directions,  $\theta_x$ ,  $\theta_y$  are the rotations of the normal to the middle plane, about the  $x$  and  $y$  axis, respectively.

The strains components associated with the displacements in equation (1) are

$$\begin{aligned}
\varepsilon_{xx} &= \frac{\partial u_0}{\partial x} - z \frac{\partial \theta_y}{\partial x} \\
\varepsilon_{yy} &= \frac{\partial v_0}{\partial y} + z \frac{\partial \theta_x}{\partial y} \\
\gamma_{xy} &= \left( \frac{\partial u_0}{\partial y} + \frac{\partial v_0}{\partial x} \right) + z \left( -\frac{\partial \theta_y}{\partial y} + \frac{\partial \theta_x}{\partial x} \right)
\end{aligned} \tag{2}$$

### 5. Piezoelectric Laminates. Constitutive Equations.

Assuming that a piezoelectric composite plate consists of several layers, including the piezoelectric layers, the constitutive equation for an orthotropic layer of the laminate substrate, is

$$\bar{\mathbf{s}} = \bar{\mathbf{Q}} \bar{\mathbf{e}} \tag{3}$$

and the constitutive equations of a deformable piezoelectric material, coupling the elastic and the electric fields are given by, Tiersten [17]

$$\bar{\mathbf{s}} = \bar{\mathbf{Q}} \bar{\mathbf{e}} - \bar{\mathbf{e}} \bar{\mathbf{E}} \tag{4}$$

$$\bar{\mathbf{D}} = \bar{\mathbf{e}}^T \bar{\mathbf{e}} + \bar{\mathbf{p}} \bar{\mathbf{E}} \tag{5}$$

where  $\bar{\mathbf{s}} = [\bar{\sigma}_{xx} \ \bar{\sigma}_{yy} \ \bar{\sigma}_{xy}]^T$  is the elastic stress vector and  $\bar{\mathbf{e}} = [\varepsilon_{xx} \ \varepsilon_{yy} \ \gamma_{xy}]^T$  the elastic strain vector,  $\bar{\mathbf{Q}}$  the elastic constitutive matrix,  $\bar{\mathbf{e}}$  the piezoelectric stress coefficients matrix,  $\bar{\mathbf{E}} = [\bar{E}_x \ \bar{E}_y \ \bar{E}_z]^T$  the electric field vector,  $\bar{\mathbf{D}} = [\bar{D}_x \ \bar{D}_y \ \bar{D}_z]^T$  the electric displacement vector and  $\bar{\mathbf{p}}$  the dielectric matrix, in the element local system ( $x, y, z$ ) of the laminate.  $\bar{Q}_{ij}$ ,  $\bar{e}_{ij}$ ,  $\bar{p}_{ij}$  are functions of ply angle  $\alpha$  for the  $k^{\text{th}}$  layer, and are given explicitly in Reddy [18].

The electric field vector is the negative gradient of the electric potential  $\Phi$ , which is assumed to be applied in the thickness  $t_k$  direction, where it can vary linearly, i.e.

$$\bar{\mathbf{E}} = -\nabla \phi \tag{6}$$

$$\bar{\mathbf{E}} = \{0 \ 0 \ E_z\}^T \tag{7}$$

where

$$E_z = -\phi/t_k \tag{8}$$

Thus, we can define the strain vector for electro elasticity as follows

$$\hat{\mathbf{e}} = \left\{ \begin{array}{c} \bar{\mathbf{e}} \\ -\bar{\mathbf{E}} \end{array} \right\} \tag{9}$$

where the mechanical strain vector can be represented as follows:

$$\bar{\mathbf{e}} = \bar{\mathbf{e}}_m + z \bar{\mathbf{e}}_b \tag{10}$$

Thus, the constitutive equations (4) and (5) can be written in the form:

$$\hat{\mathbf{s}} = \begin{Bmatrix} \bar{\mathbf{s}} \\ \bar{\mathbf{D}} \end{Bmatrix} = \begin{bmatrix} \bar{\mathbf{Q}} & \bar{\mathbf{e}} \\ \bar{\mathbf{e}}^T & -\bar{\mathbf{p}} \end{bmatrix} \begin{Bmatrix} \bar{\mathbf{e}} \\ -\bar{\mathbf{E}} \end{Bmatrix} = \hat{\mathbf{C}} \hat{\mathbf{e}} \tag{11}$$

### 6. Finite element formulation.

The flat triangular finite element model has three nodes and six degrees of freedom per node, the displacements  $u_{0i}$ ,  $v_{0i}$ ,  $w_{0i}$  and rotations  $\theta_{xi}$ ,  $\theta_{yi}$ ,  $\theta_{zi}$ . The introduction of fictitious stiffness coefficients  $K_{\theta_z}$ , corresponding to rotations  $\theta_z$ , to eliminate the problem of a singular stiffness matrix, for which the elements are coplanar or near coplanar, is required. The element local displacements, slopes and rotations are expressed in terms of nodal variables through shape functions  $\mathbf{N}_i$  given in terms of area coordinates [19].

The displacement field can be represented in matrix form as

$$\mathbf{u} = \mathbf{Z} \left( \sum_{i=1}^3 \mathbf{N}_i \mathbf{d}_i \right) = \mathbf{Z} \mathbf{N} \mathbf{a} \quad (12)$$

$$\mathbf{d} = \sum_{i=1}^3 \mathbf{N}_i \mathbf{d}_i = \mathbf{N} \mathbf{a} \quad (13)$$

with

$$\mathbf{d}_i = \{u_{0i} \ v_{0i} \ w_{0i} \ \theta_{xi} \ \theta_{yi} \ \theta_{zi}\}^T \quad (14)$$

$$\mathbf{Z} = \begin{bmatrix} 1 & 0 & 0 & 0 & -z & 0 \\ 0 & 1 & 0 & z & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 \end{bmatrix} \quad (15)$$

and the strain field as follows

$$\{\bar{\epsilon}\} = \{\bar{\epsilon}_m + z \bar{\epsilon}_b\} = \left\{ \sum_{i=1}^3 (\mathbf{B}_i^m + z \mathbf{B}_i^b) \mathbf{d}_i \right\} = \mathbf{B}^{mb} \mathbf{a} \quad (16)$$

The electric field is given by

$$\mathbf{E} = -\mathbf{B}^\phi \phi \quad (17)$$

where  $\mathbf{B}^\phi$  is the electric field – potential relations matrix given by:

$$\mathbf{B}^\phi = \begin{bmatrix} 1/t_1 & \mathbf{L} & 0 \\ \mathbf{M} & \mathbf{O} & \mathbf{M} \\ 0 & \mathbf{L} & 1/t_{NPL} \end{bmatrix} \quad (18)$$

The dynamic equations of a laminated composite plate can be derived from the Hamilton's principle, which is given as follows:

$$\int_{t_1}^{t_2} \left\{ \sum_{K=1}^N \left( \int_{A^e} \int_{h_{k-1}}^{h_k} \delta \hat{\epsilon}^T \hat{C}_k \hat{\epsilon}^L dz dA^e - \int_{A^e} \int_{h_{k-1}}^{h_k} \delta \mathbf{u}^T \rho_k \mathbf{u} dz dA^e \right) - \left( \int_V \delta \mathbf{f} dV + \int_S \delta \mathbf{t} dS + \sum_i \mathbf{F}_i \delta \mathbf{u}_i + \int_S \mathbf{Q} \delta \phi dS \right) \right\} dt = 0 \quad (19)$$

Entering the equations (11) to (15) into equation (17), we have

$$\int_{t_1}^{t_2} \left\{ \sum_{K=1}^N \left( \int_A \int_{h_{k-1}}^{h_k} \delta \{a\}^T \begin{bmatrix} \mathbf{B}^{mb} & 0 \\ 0 & \mathbf{B}^\phi \end{bmatrix}^T \begin{bmatrix} \bar{\mathbf{Q}} & \bar{\mathbf{e}} \\ \bar{\mathbf{e}}^T & -\bar{\mathbf{p}} \end{bmatrix} \begin{bmatrix} \mathbf{B}^{mb} & 0 \\ 0 & \mathbf{B}^\phi \end{bmatrix} \{a\} dz dA - \int_A \int_{h_{k-1}}^{h_k} \delta \left\{ \begin{matrix} \mathbf{u} \\ \phi \end{matrix} \right\}^T \mathbf{N}^T \mathbf{m} \mathbf{N} \left\{ \begin{matrix} \mathbf{u} \\ \phi \end{matrix} \right\} dz dA \right) + \int_V \delta \{a\}^T \mathbf{N}^T \mathbf{f} dV + \int_S \delta \{a\}^T \mathbf{N}^T \mathbf{T} dS + \delta \{a\}^T \mathbf{F}_c + \int_S \mathbf{Q} \delta \{\phi\} dS \right\} dt = 0 \quad (20)$$

To the first and second terms of first member of Eq. (20), corresponds the element stiffness and mass matrices, respectively, which are defined by

$$\mathbf{K} = \begin{bmatrix} \mathbf{K}_{uu} & \mathbf{K}_{u\phi} \\ \mathbf{K}_{\phi u} & \mathbf{K}_{\phi\phi} \end{bmatrix} = \sum_{K=1}^N \int_A \int_{h_{k-1}}^{h_k} \begin{bmatrix} \mathbf{B}^{mb} & 0 \\ 0 & \mathbf{B}^\phi \end{bmatrix}^T \begin{bmatrix} \bar{\mathbf{Q}} & \bar{\mathbf{e}} \\ \bar{\mathbf{e}}^T & -\bar{\mathbf{p}} \end{bmatrix} \begin{bmatrix} \mathbf{B}^{mb} & 0 \\ 0 & \mathbf{B}^\phi \end{bmatrix} dz dA \quad (21)$$

$$\mathbf{M} = \int_A \mathbf{N}^T \left( \sum_{k=1}^n \rho_k \int_{h_{k-1}}^{h_k} \mathbf{Z}^T \mathbf{Z} dz \right) \mathbf{N} dA \quad (22)$$

To the third term of Eq. (20), correspond the applied electric charge vector  $\mathbf{F}^{ele}$  to be defined later, and the external mechanical force vector, which is defined by:

$$\mathbf{F}_{ext}^{mec} = \int_V \mathbf{N}^T \mathbf{f} dA + \int_S \mathbf{N}^T \mathbf{t} dS + \mathbf{F}_c \quad (23)$$

where  $\mathbf{f}$ ,  $\mathbf{t}$ ,  $\mathbf{F}_c$  are body, surface, and concentrated force vectors.

The element stiffness and mass matrices as well as external load vector are initially computed in the local coordinate system attached to the element. To solve general structures, local - global transformations are needed, Zienkiewicz [19]. After these transformations the assembled system of equations is:

$$\begin{bmatrix} \mathbf{M}_{uu} & 0 \\ 0 & 0 \end{bmatrix} \begin{Bmatrix} \mathbf{q} \\ \phi \end{Bmatrix} + \begin{bmatrix} \mathbf{K}_{uu} & \mathbf{K}_{u\phi} \\ \mathbf{K}_{\phi u} & \mathbf{K}_{\phi\phi} \end{bmatrix} \begin{Bmatrix} \mathbf{q} \\ \phi \end{Bmatrix} = \begin{Bmatrix} \mathbf{F}_{ext}^{mec}(t) \\ \mathbf{F}^{ele}(t) \end{Bmatrix} \quad (24)$$

Assuming that piezoelectric sensors as well as actuators are bonded or embedded in the structure, the electric potential vector is subdivided in a sensor component  $\phi^{(S)}$  and an actuator component  $\phi^{(A)}$ .

The external applied electric charge at the sensors is zero. Separating the actuator and sensor components, the system of Eq. (24) take

the following form:

$$[M_{uu}]\{\ddot{q}\} + [K_{uu}]\{q\} + [K_{u\phi}^{(S)}]\{\phi^{(S)}\} = \{F_{ext}^{mec}(t) - K_{u\phi}^{(A)}\phi^{(A)}\} \quad (25)$$

$$[K_{\phi u}^{(A)}]\{q\} + [K_{\phi\phi}^{(A)}]\{\phi^{(A)}\} = \{F^{ele}(t)\} \quad (26)$$

$$[K_{\phi u}^{(S)}]\{q\} + [K_{\phi\phi}^{(S)}]\{\phi^{(S)}\} = \{0\} \quad (27)$$

From the last equation, the induced sensory electric potentials  $\phi^{(S)}$  are obtained as follows:

$$\{\phi^{(S)}\} = -[K_{\phi\phi}^{(S)}]^{-1} [K_{\phi u}^{(S)}]\{q\} \quad (28)$$

Thus the equation (25) takes the form

$$[M_{uu}]\{\ddot{q}\} + \left( [K_{uu}] - [K_{u\phi}^{(S)}][K_{\phi\phi}^{(S)}]^{-1}[K_{\phi u}^{(S)}] \right)\{q\} = \{F_{ext}^{mec}(t) - K_{u\phi}^{(A)}\phi^{(A)}\} \quad (29)$$

### 6.1 Dynamic Analysis

The sensor output can be obtained as follows (Reddy [20]). The charge output of each sensor, with poling in the z direction, can be expressed in terms of spatial integration of the electric displacement over its surface, taking into account that the converse piezoelectric effect is negligible. Thus we have

$$Q^{(S)}(t) = \int_A D_z(t) dA = \frac{1}{2} \left[ \int_{A(z=z_k)} D_z(t) dA + \int_{A(z=z_{k+1})} D_z(t) dA \right] \quad (30)$$

From Eq. (5), the last equation can be written as follows

$$Q^{(S)}(t) = \int_A \bar{e}^T \bar{\epsilon} dA \quad (31)$$

or in the discretized form

$$Q^{(S)}(t) = \left( \int_A \bar{e}^T B^{mb} dA \right) \{q\} = [K_{\phi u}^{(S)}]\{q\} \quad (32)$$

The current on the surface of the sensor is given by

$$I(t) = \frac{dQ^{(S)}}{dt} \quad (33)$$

When the piezoelectric sensor is used as strain rate sensor, the current can be converted into the open circuit sensor voltage output  $\phi^{(S)}$  as

$$\phi^{(S)} = G_c \frac{dQ^{(S)}}{dt} \quad (34)$$

where  $G_c$  is the constant gain of the amplifier, which transforms the sensor current to voltage.

The sensor output voltage can be feed back through an amplifier to the actuator with a change of polarity. Thus, we have for the actuator voltage

$$\phi^{(A)} = -G_i G_c \frac{dQ^{(S)}}{dt} \quad (35)$$

where  $G_i$  is the gain of the amplifier to provide feedback control?

The actuator voltage written in the discretized form, is then given by

$$\phi^{(A)} = -G_i G_c [K_{\phi u}^{(A)}]\{q\} \quad (36)$$

Use of Eq. (36) into Eq. (29) introduces an equivalent negative velocity feedback, and the motion equations become:

$$[M_{uu}]\{\ddot{q}\} - G_i G_c [K_{u\phi}^{(A)}][K_{\phi u}^{(S)}]\{q\} + \left( [K_{uu}] - [K_{u\phi}^{(S)}][K_{\phi\phi}^{(S)}]^{-1}[K_{\phi u}^{(S)}] \right)\{q\} = \{F_{ext}^{mec}(t)\} \quad (37)$$

Considering Rayleigh type damping, we can write:

$$[M_{uu}]\{\ddot{q}\} + (C_R + C_A)\{\dot{q}\} + \left( [K_{uu}] - [K_{u\phi}^{(S)}][K_{\phi\phi}^{(S)}]^{-1}[K_{\phi u}^{(S)}] \right)\{q\} = \{F_{ext}^{mec}(t)\} \quad (38)$$

with

$$C_R = \alpha M_{uu} + \beta K_{uu} \quad (39)$$

where  $\alpha$  and  $\beta$  are Rayleigh's coefficients, to account for inherent structural damping, and the damping effect due to the active control is given by:

$$C_A = -G_i G_c [K_{u\phi}^{(A)}][K_{\phi u}^{(S)}] \quad (40)$$

The solution of Eq. (38) is carried out using Newmark direct method of time integration, Bathe [8].

## 7. Optimal Design

A general structural optimization problem can be stated as:

$$\begin{aligned} \min \{ \Omega(\mathbf{b}) \} \quad & \text{subject to: } b_i^l \leq b_i \leq b_i^u \quad i = 1, \dots, ndv \\ & \Psi_j(\mathbf{q}, \mathbf{b}) \leq 0 \quad j = 1, \dots, m \end{aligned} \quad (41)$$

where  $\Omega(\mathbf{b})$  is the objective function,  $\mathbf{b}$  is the vector of design variables  $b_i$ ,  $\Psi_j(\mathbf{q}, \mathbf{b})$  are the  $m$  inequality behavioral constraint equations,  $b_i^l$  and  $b_i^u$  are respectively, the lower and upper limits of the design variables and  $ndv$  is the total number of design variables.

If the objective function and/or the constraint equations are continuous functions of the design variables, mathematical programming techniques [21,22] requiring only the computation of  $\Omega(\mathbf{b})$ ,  $\Psi_j(\mathbf{q}, \mathbf{b})$  and their gradients, provide a general, flexible and efficient formulation for engineering design problems.

For discrete variable structural problems, a variety of methods including simulated annealing can be used [14,23,24]. As pointed out by Correia et al [24,25], among others, the main advantage of this method, in comparison with gradient-based methods, is the ability to overcome the premature convergence towards a local optimum. By other hand, the main disadvantage is related with the computational cost, because of the high number of objective function evaluations usually required to reach the optimal solution, which is especially relevant when the objective function evaluation is computationally expensive. The implemented simulated annealing procedure employs a random search that generates feasible sets of design variables, accepting not only changes in the design variables that decrease the objective function but also changes that increase it. The latter changes are accepted with a certain probability. The basic functioning of the simulated annealing algorithm can be easily described as follows [16,23]:

- a) Generate a random perturbation on the design variables and obtain the change in the objective function  $\delta\Omega(\mathbf{b})$ .
- b) If  $\delta\Omega(\mathbf{b}) \leq 0$ , the new design is better and the new design is accepted.
- c) If  $\delta\Omega(\mathbf{b}) > 0$ , accept the new design with a probability given by

$$p = \exp(-\delta\Omega/T) \quad (42)$$

In equation (42),  $T$  is a control parameter, which is called the system temperature, based on the analogy to the physical process of annealing a metal. The initial temperature  $T_0$  and the rate at which it is lowered, usually referred to as the cooling schedule, have great influence on the performance of the algorithm. This temperature must be set high enough so that, at least initially, all states proposed are accepted. Some authors use the acceptance ratio  $\mu$  of accepted moves to proposed moves, obtaining the initial temperature from the relation

$$\mu = \exp\left(\frac{-\delta\Omega(\mathbf{b})}{T_0}\right) \quad (43)$$

where  $\mu$  is set to a chosen value. The temperature decrement can be given by

$$T_{k+1} = \alpha \cdot T_k \quad (44)$$

where  $T_k$  and  $T_{k+1}$  represent the system temperatures at  $k$  and  $k+1$  successive iterations and  $\alpha$  is the cooling parameter usually taken in the range 0.8-0.95. Note that with decreasing temperatures the probability given by equation (42) will decrease, hence as the simulated annealing iterations proceed, it becomes increasingly unlikely that a worse design can be accepted. The search is halted when no improvement in the objective function is found combined with the acceptance ratio falling below a specified value.

In the next section an illustrative example of the forced response with active feedback control of a composite adaptive laminated plate is presented, where firstly, it is designed for maximum stiffness being the design variables the orientation angles of the reinforcement fibers in the orthotropic layers. Secondly, the optimal location of the piezoelectric actuator discrete patches is found in order to achieve maximum piezoelectric actuator performance.

The first optimization problem can be regarded as a continuous design variable problem and can be solved by using a gradient-based optimization technique. However if we consider that because of manufacturing constraints, normally the fiber reinforcement directions assume discrete values in a limited set of different possible angles we can regard also this problem as a discrete design variable problem. In this case we solve this problem by using the simulated annealing method. The objective here is to maximize the elastic strain energy of the laminated composite or to minimize the elastic displacements in specific locations of the structure, being the design variables the orientation angles of the orthotropic axes in each layer.

For the optimal location of the piezoelectric patches we consider that the patches can only assume the positions corresponding to the finite element mesh discretization. This is in essence a discrete variable optimization problem, which is solved by using the simulated annealing method. In this case the objective function can be for example the maximization of the displacement in a specific point of the structure and the design variables are the discrete locations of the piezoelectric patches. Here we are interested to obtain the best combination of patches in discrete locations, which give the maximum actuation displacement.

## 8. Numerical applications

### 8.1 Forced response with active feedback control of a simply-supported square plate, in free vibration and forced vibration under sinusoidal loading.

A simply-supported square ( $a \times a$ ) laminated plate, having the initial lamination sequence of  $[0^\circ/90^\circ/0^\circ]$ , integrating piezoelectric actuator and sensor layers or patches made of PZT, bonded on upper and lower surfaces, is considered. The material properties of the substrate layers are:  $E_1 = 172.5$  GPa,  $E_2 = 6.9$  GPa,  $G_{12} = 3.45$  GPa,  $\nu_{12} = 0.25$ ,  $\rho = 1600$  kg/m<sup>3</sup>. The material and piezoelectric properties of PZT are:  $E_1 = E_2 = 63$  GPa,  $G_{12} = 24$  GPa,  $\nu_{12} = 0.30$ ,  $\rho = 7600$  kg/m<sup>3</sup>,  $e_{31} = e_{32} = C/m^2$ ,  $p_{33} = 1.5 \times 10^{-8}$  F/m.

The side dimension is  $a = 0.18$  m and the thickness of the substrate layers and PZT are 0.002 m and 0.0001 m, respectively. The plate is modeled by a (6x6) element mesh (72 triangular elements). First we search for the optimal core lamination sequence, which leads to the maximum fundamental natural frequency of the plate, by using the simulated annealing optimization method. For the simulated annealing problem, the design variables are chosen from a discrete set of possible ply angles defined as:  $S_{dv} = \{0^\circ, \pm 15^\circ, \pm 30^\circ, \pm 45^\circ, \pm 60^\circ, \pm 75^\circ, 90^\circ\}$ . The cooling schedule parameters used in the simulated annealing process were:  $T_0 = 200$  for the initial temperature and  $\alpha = 0.9$  for the rate of temperature reduction. The optimal lamination sequence is found to be  $[45^\circ/-45^\circ/45^\circ]$ . The same optimal design is obtained using gradient optimization.

The amplitudes and the natural frequencies of the plate in free vibrations, enabled by an applied uniform distributed transverse load  $q = 10000$  N/m<sup>2</sup> and then removed, are shown in Figure 1 for both the initial and final lamination sequences. As expected, for the final design, there is a decreasing in the central amplitude and the fundamental natural frequency increases.

Next we pretend to investigate the optimal position of the piezoelectric actuators patches, in order to maximize the control of the plate, in this case measured by the amplitude of the deflection in the center of the plate, and compare it with the control obtained by using an entire piezoelectric layer. In this application, 4 actuators patches covering 8 triangular elements are introduced. The plate is modeled by a (6x6) element mesh (72 triangular elements). Using the simulating annealing algorithm the optimal positions of the patches is obtained as represented in Figure 2 a), i.e. the central plate elements (*Patches 1*). In Figure 2 a) and b) two different patches positions are shown, and the corresponding central line deflections, as well as those obtained with an entire actuator layer, are shown in Figure 3. Figure 4 illustrates the responses for central amplitude of the plate, and compares these responses in three different situations. The controlled responses, obtained using a time step  $\Delta t = 0.000125$  s in the Newmark method, with  $G_i G_c = 8000$ , clearly demonstrates the damping effect of the piezoelectric actuators on the free vibration of the plate with an applied uniform distributed transverse load  $q = 10000$  N/m<sup>2</sup>, which was then suddenly removed.

For an applied distributed transverse harmonic load  $q(t) = q \sin 2\pi f t$  with a magnitude  $q = 10000$  N and frequency  $f = 10$  Hz, Figure 5 shows the uncontrolled and controlled responses, where the effect of negative velocity feedback control,  $G_i G_c = 3.2 \times 10^7$ , is evident.

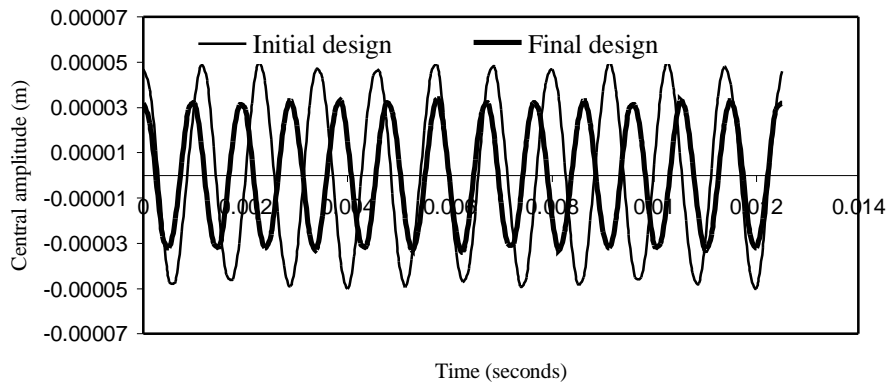


Figure 1. Central amplitude for the initial and final lamination sequence.

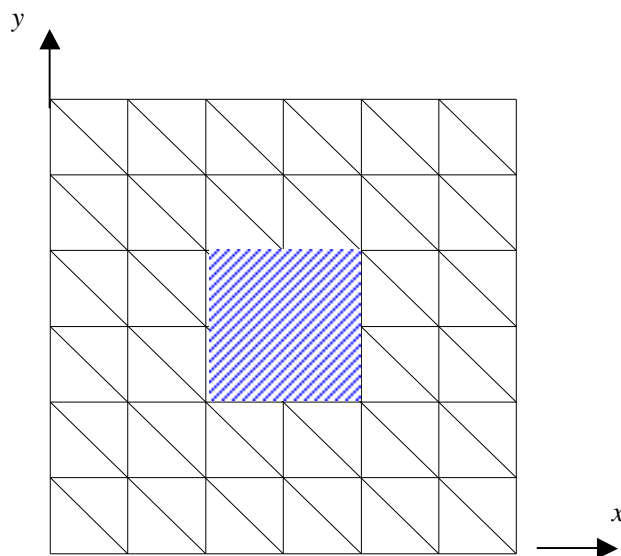


Figure 2 a). Patch locations – *Patches 1*.

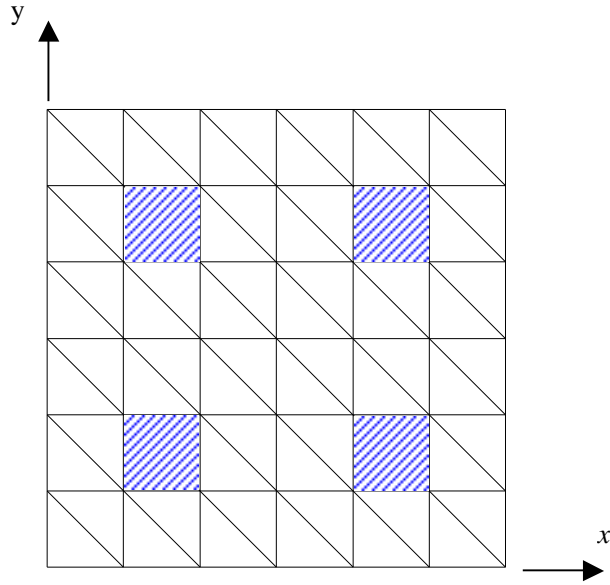


Figure 2 b). Patch locations – Patches 2.

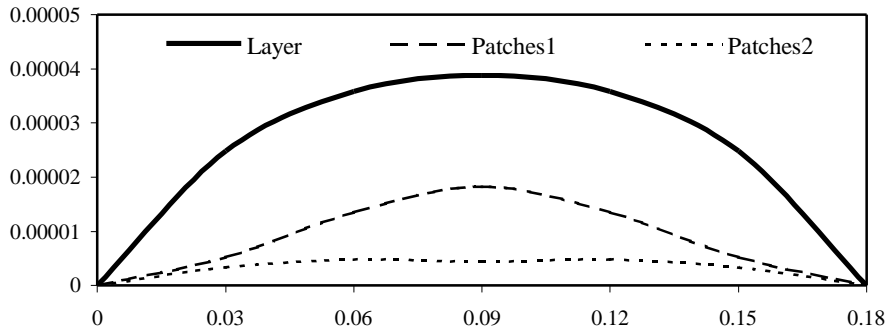


Figure 3. Central line deflections.

Figure 6 shows the relation of the amplitude of plate vibration versus feedback control gain  $G$ . It can be observed that the amplitude of plate vibration decreases quickly when the control gain increases up to approximately  $G = G_1 G_c \approx 40 \times (1.6 \times 10^6)$ , but then it decreases very slowly for higher values of  $G$ .

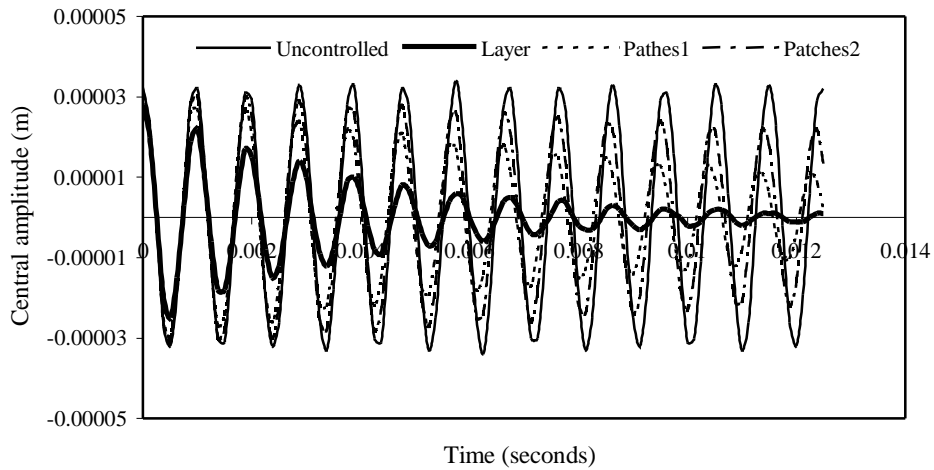


Figure 4. Uncontrolled and controlled responses on the plate central amplitude.

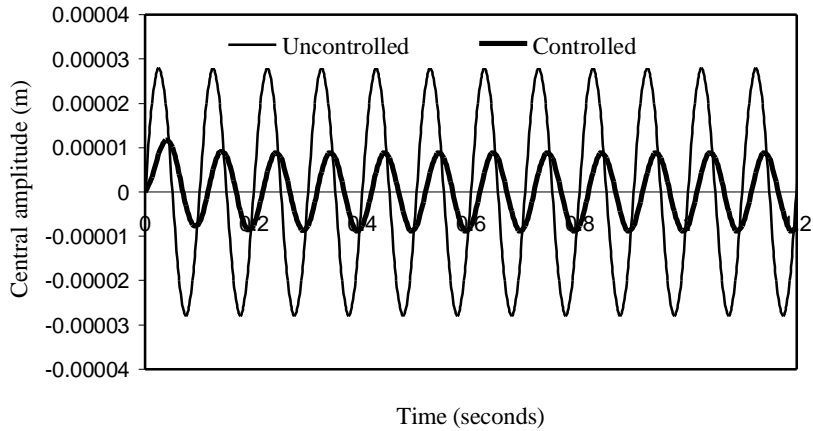


Figure 5. Uncontrolled central deflection,  $f = 10$  Hz

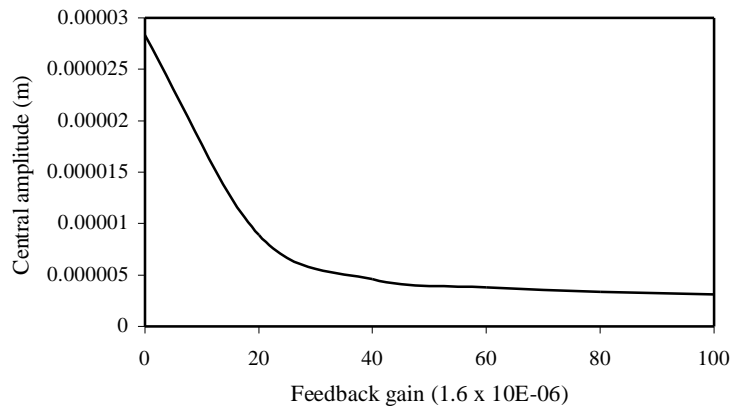


Figure 6. Amplitude vs. Feedback gain

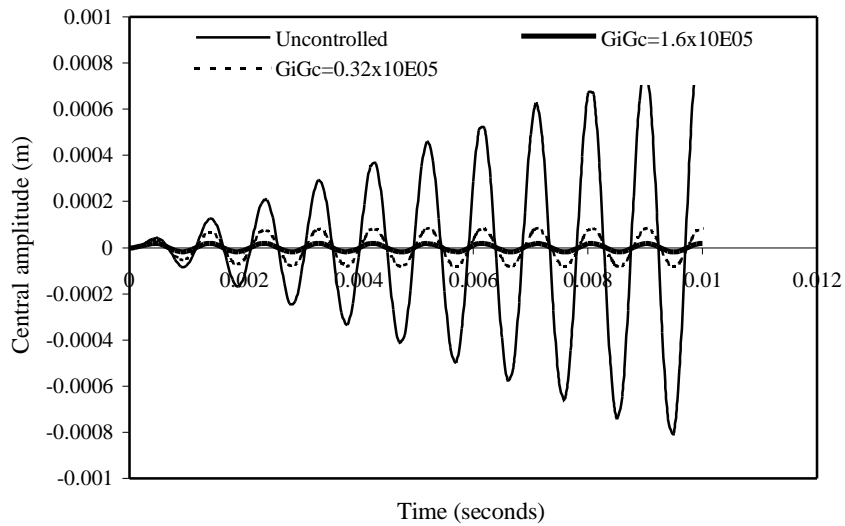


Figure 7. Effect of negative velocity feedback control on the central deflection

For the same applied transverse harmonic load but now considering a frequency of  $f = 1050$  Hz, which is very close to the first natural frequency, Figure 7 illustrates the uncontrolled and controlled responses for central deflection  $w$ . The controlled responses for gains  $G_i G_c = 0.32 \times 10^5$  and  $G_i G_c = 1.6 \times 10^5$ , clearly demonstrates the action of the piezoelectric actuator layer, which increases the overall damping of the system.



## 9. Conclusions

The active control capability of composite structures covered with piezoelectric layers or patches is investigated, using the finite element method. A finite element based on the Kirchhoff classical theory, has been developed. The present model has been validated in Moita *et al.* [26], where the solutions for deflection and sensed voltage in a bimorph beam, are compared with the solutions obtained by other authors. Here, the results obtained, show that the negative velocity feedback control algorithm used in this model is effective for an active damping control of vibration response. Also the core optimization had been performed in order to minimize the vibration amplitude and maximization of first natural frequency. The patch position optimization had been also performed in order to maximize the effect a defined set of actuators. Both of these optimizations used the simulated annealing method.

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