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## Universal multifractal description of an hourly rainfall time series from a location in southern Spain

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### RESUMEN

El formalismo multifractal de turbulencia ha sido usado para llevar a cabo el análisis de la estructura temporal, para escalas desde 1 hora hasta casi 6 meses, de la serie de datos de lluvia horaria registrada durante veinticuatro años en Córdoba, localidad situada en el sur de España. Los parámetros del modelo multifractal universal fueron estimados y se obtuvo la función teórica de los momentos estadísticos. Se encontró un buen ajuste a la función empírica para un intervalo de valores de momentos, demostrándose que el modelo multifractal universal resulta adecuado para describir estadísticamente la serie temporal de lluvia registrada en Córdoba.

### ABSTRACT

Multifractal turbulence formalism has been used to perform an analysis for scales from 1 hour to almost 6 months of the time structure of the hourly rainfall series recorded during twenty-four years in Córdoba, a location in southern Spain. The parameters of the universal multifractal model were estimated and the theoretical moments scaling exponent function was obtained exhibiting an acceptable agreement with the empirical function for a range of moments. The universal multifractal model shown itself to be a suitable tool for describing the statistics of the rainfall series recorded in Córdoba.

**Keywords:** Rainfall, multifractal analysis, universal multifractal model.

## 1. Introduction

Rainfall is the main factor in many processes with a significant environmental impact such as floods, droughts, runoff, soil erosion or spread of pollutants. However, the great temporal (and spatial) non-linear variability of precipitation limits the use of many existing models to describe this process.

On certain space-time scales and, specifically, in very small catchments and in urban areas, rainfall intensities present larger fluctuations, and therefore higher values than on meteorological model scales. Thus, the reliability of a forecast provided by meteorological models also depends on the efficiency of downscaling models.

The invariance of properties in any process over a range of scales is called scaling. Scaling systems can be described by fractal and multifractal theories, the latter being an evolution of the former. Fractal theory (Mandelbrot, 1982) deals with simple scaling and can describe complex phenomena by using few parameters. However, multifractal theory deals with multiscaling that allows to generalize the scaling properties of a process (Grassberger, 1983; Hentschel and Procaccia, 1983; Deidda, 2000; Deidda *et al.*, 2006).

The multifractal analysis of rainfall has been widely used to describe its temporal and spatial distribution (Schertzer and Lovejoy, 1987; Fraedrich and Larnder, 1993; Ladoy *et al.*, 1993; Tessier *et al.*, 1993, 1996; Over and Gupta, 1994; Svensson *et al.*, 1996; de Lima and Grasman, 1999; Kiely and Ivanova, 1999; Sivakumar, 2001; Labat *et al.*, 2002; Veneziano and Furcolo, 2002; Olsson and Burlando, 2002; Kantelhardt *et al.*, 2006; Langousis and Veneziano, 2007). The universal multifractal model (UM) which uses the formalism of turbulence, was developed by Schertzer and Lovejoy (1987) to model the variability of rainfall by a multiplicative cascade process, in which the flux of water is transferred from larger to small regions (scales) of the atmosphere (Over and Gupta, 1994), in a similar way to what happens in turbulence models, where the transfer of energy from larger to smaller scales is assumed. This multifractal approach has been used to carry out multiple analyses of rainfall and related topics (Hubert *et al.*, 1993; Ladoy *et al.*, 1993; Tessier *et al.*, 1993, 1996; Svensson *et al.*, 1996; de Lima and Grassman, 1999; Lilley *et al.*, 2006; Lovejoy and Schertzer, 2006; Lashermes and Foufoula-Georgiou, 2007).

The objective of the present work was to perform an analysis of the long-term hourly rainfall series recorded in Córdoba, a location in southern Spain, by using the multifractal turbulence formalism to check whether the results were consistent with those found in earlier studies for other locations. This was done as a contribution to the consolidation of the existing knowledge on the multifractal nature of precipitation, and to confirm the possible use of the universal multifractal model to generate synthetic rainfall data in the cited location.

## 2. Material and methods

### 2.1 Universal multifractals

To carry out a rainfall time structure analysis by applying multifractal turbulence formalism, the data series should be divided into non-overlapping intervals of a certain resolution. The ratio of

the maximum scale of the field to this interval is termed “scale ratio”,  $\lambda$  (Svensson *et al.*, 1996). Considering  $\varepsilon_\lambda$  as the rainfall intensity on each time interval at the scale ratio  $\lambda$  divided by the average of the sample corresponding to the largest scale of interest ( $\lambda = 1$ ), the ensemble average  $q$ th moment  $\langle \varepsilon_\lambda^q \rangle$  is related to the empirical moments scaling exponent function  $K(q)$  in such a way that (Schertzer and Lovejoy, 1987):

$$\langle \varepsilon_\lambda^q \rangle \approx \lambda^{K(q)} \quad (1)$$

The scaling behaviour expressed by Eq. (1) can be investigated by plotting  $\langle \varepsilon_\lambda^q \rangle$  as a function of  $\lambda$  in a log-log diagram for different values of  $q$ . The curves obtained will exhibit an approximately linear behaviour with slopes that are estimations of  $K(q)$  if Eq. (1) is valid. If  $K(q)$  versus  $q$  is a straight line, the data set is monofractal. However, if  $K(q)$  versus  $q$  is a convex function, the data set is multifractal (Parisi and Frisch, 1985).

The universal multifractal (UM) model was proposed by Schertzer and Lovejoy (1987) assuming that the generator of multifractals was a random variable with an exponentiated extremal Lévy distribution. Thus, the theoretical scaling exponent function  $K(q)$  for the moments  $q \geq 0$  of a cascade process is obtained according to:

$$K(q) = qH + \begin{cases} C_1 (q^\alpha - q)/(\alpha - 1) & \alpha \neq 1 \\ C_1 q \log(q) & \alpha = 1 \end{cases} \quad (2)$$

in which  $\alpha \in [0, 2]$  is the Lévy index and indicates the degree of multifractality (i.e. the deviation from monofractality).  $\alpha = 0$  and  $\alpha = 2$  correspond to the monofractal and log-normal cases, respectively.  $C_1 \in [0, d]$ , with  $d$  being the dimension of the support ( $d = 1$  in this case), describes the sparseness or inhomogeneity of the mean of the process. Parameter  $H$  determines the deviation from conservation times ( $\langle \varepsilon_\lambda^q \rangle = \lambda^{-H}$ ). For conserved processes,  $H = 0$ .

The estimation of the parameters  $C_1$  and  $\alpha$  can be made by applying the double trace moment (DTM) technique (Lavallée *et al.*, 1993). Thus, the intensity  $\varepsilon_\lambda$  associated with the finest resolution  $\Lambda$  is first raised to the power  $\eta$ , degraded to the scale ratio  $\lambda$  and the  $q$ th power of the result averaged over the available data. The double moment exponent function  $K(q, \eta)$  is derived from:

$$\langle (\varepsilon_\lambda^\eta)_\lambda^q \rangle = \lambda^{K(q, \eta)} \quad (3)$$

by using log-log plots, as described for the  $K(q)$  function.

The parameter can be estimated as the slope of the linear regression fitted to the plot of  $|K(q, \eta)|$  versus  $\log(\eta)$  for fixed  $q$  (Tessier *et al.*, 1996). The estimation of  $C_1$  is derived by considering that  $K(q, 1)$  is the intercept of the linear regression with  $\log(\eta) = 0$ . Alternatively, it can be obtained from the empirical moments scaling exponent function  $K(q)$ , considering  $C_1 = [dK(q)/dq]_{q=1}$ .

From the estimated values of  $\alpha$  and  $C_1$ , and taking the value of exponent  $\beta$  that characterizes the energy spectrum of the conserved process  $E(\omega) \approx \omega^{-\beta}$ , with  $\omega$  being the frequency, the parameter  $H$  is given by (Lavallée *et al.*, 1993):

$$H = 0.5 (\beta - 1 + K(2)) = 0.5 (\beta - 1 + C_1 (2^\alpha - 2)/(\alpha - 1)) \tag{4}$$

2.2. Rainfall data

The hourly rainfall data analyzed in this work were collected at Córdoba airport weather station (37.85° N, 4.85° W) from 1980 to 2003. The altitude is 117 m and the climate of this area can be defined as a mixture between Mediterranean characteristics and Continental effects. The mean annual temperature is 17.7 °C and the amplitude of the average temperatures is 19.6 °C. The annual average potential evaporation is 1500-2000 mm and the mean precipitation is around 600 mm. The rainy period mainly occurs from October to May with the highest amount of precipitation recorded during November-December and February. The general impression for the area is that it is semi-arid with hot, dry summers and relatively temperate winters.

The rainfall data were recorded using a Hellmann rain gauge (World Meteorological Organization standard), with a horizontal opening of 200 cm<sup>2</sup> at a height of 1.2 m. The temporal resolution of this continuously-recorded rain gauge was 1 hour and its accuracy, expressed as the minimum amount of rain that it was able to record, was 0.1 mm.

3. Results

Following Eq. (1), a wide range of moments of the rainfall intensity  $\varepsilon_\lambda$  on the time scales from 1 hour ( $\lambda = 4096$ ) to almost 6 months ( $\lambda = 1$ ) have been calculated (Figs. 1a and 1b).

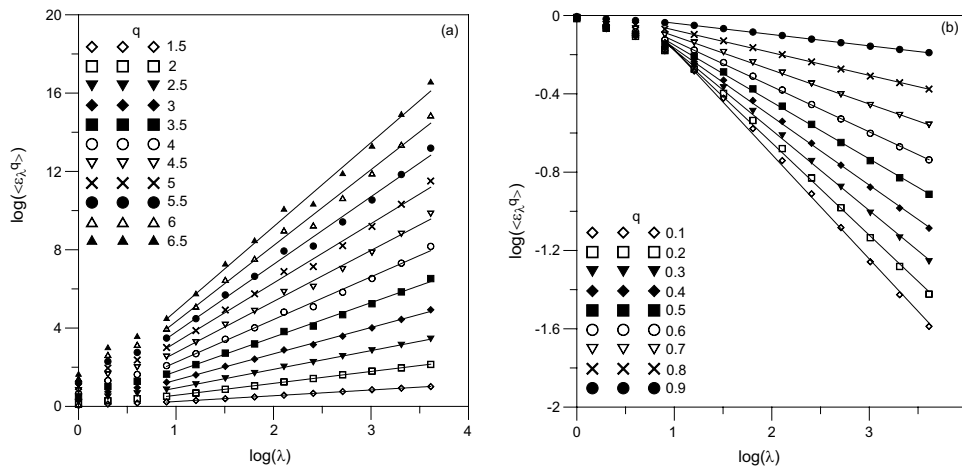


Fig. 1. Log-log plot of the  $q$ th moments of the rainfall intensity on the time scales from 1 hour to almost 6 months versus the scale ratio  $\lambda$ . (a) For moments larger than 1; (b) for moments smaller than 1.

Straight lines in Figs. 1a and 1b can be fitted for all the  $q$ th moments, showing the scaling behavior from 1 hour to almost 21 days. The accuracy of the rain gauge (0.1 mm) and the value of large observations in the rainfall data set, affect the results obtained for the  $q$ th moments (de Lima and Grasman, 1999). This influence can be observed for small and high values of  $q$ , for which the experimental data (symbols in Figs 1a and 1b) are not aligned. This fact is a consequence of the limitations of the data set.

The scaling break at 21 days shown in Figures 1a and 1b is presumably a manifestation of the “synoptic maximum”, which is the typical lifetime of planetary scale atmospheric structures (Kolesnikov and Monin, 1965). Different values of the “synoptic maximum” for the rainfall have been found between 11 (de Lima and Grasman, 1999) and 16 days (Ladoy *et al.*, 1993; Tessier *et al.*, 1996), and one month (Fraedrich and Larnder, 1993; Svensson *et al.*, 1996).

Figure 2 shows the empirical exponent function  $K(q)$  that describes the scaling of the moments from 1 hour to 21 days. The shape of  $K(q)$  is mainly convex, meaning that the rainfall time series is multifractal. Nevertheless, this function exhibits a linear behaviour for  $q < q_{min} \approx 1.75$  and  $q > q_{max} \approx 2.9$ .

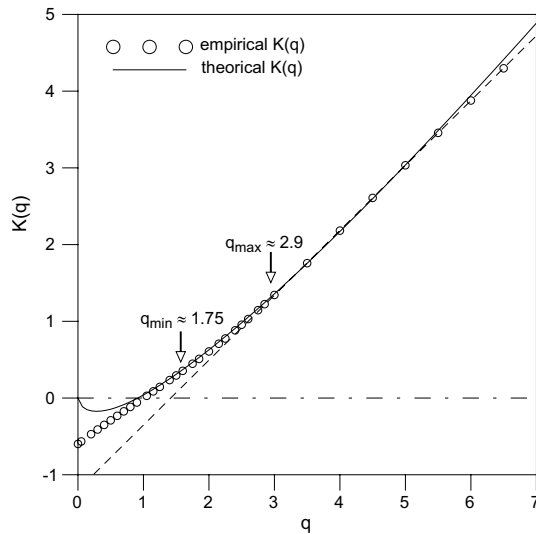


Fig. 2. Empirical and theoretical moments scaling exponent functions represented by the black dots and the solid line, respectively. The empirical  $K(q)$  function exhibits a linear behavior for  $q < q_{min} \approx 1.75$  and  $q > q_{max} \approx 2.9$ .

Fig. 3a shows the log-log plots of  $\langle (\varepsilon_\lambda^n)_\lambda^q \rangle$  versus  $\lambda$  for  $q = 2$ , considering different values of  $\eta$ . The log-log linearity in these plots demonstrates the scaling nature of the rainfall series. The corresponding plots of  $\log |K(q, \eta)|$  versus  $\log(\eta)$ , for  $\eta$  values of  $10^{-1}$  to  $10$  at  $10^{0.05}$  increments, are shown in Figure 3b for some of the investigated  $q$ -moments, all of them within the range from  $q_{max} = 1.75$  up to  $q_{max} = 2.9$ ; these values being those that delimit the scaling behavior (i.e. convex shape) in the empirical scaling exponent function  $K(q)$  described in Figure 2. The slope of the linear segments located at the middle parts of these curves permitted to obtain a value of  $\alpha = 0.732 \pm 0.011$ . According to the estimated value for  $\alpha$ , the rainfall process belongs to the universality

class with  $1 < \alpha < 0$  (Lovejoy and Schertzer, 1990) that corresponds to (log) Lévy processes with bounded singularities, denominated as conditionally hard multifractals. This result is in agreement with those found by authors such as Tessier *et al.* (1993) for a global network of 1000 weather stations, Hubert *et al.* (1993) for Reunion Island (Indian Ocean), Ladoy *et al.* (1993) for Nimes (France) and de Lima and Grasman (1999) for Vale Formoso (Portugal).

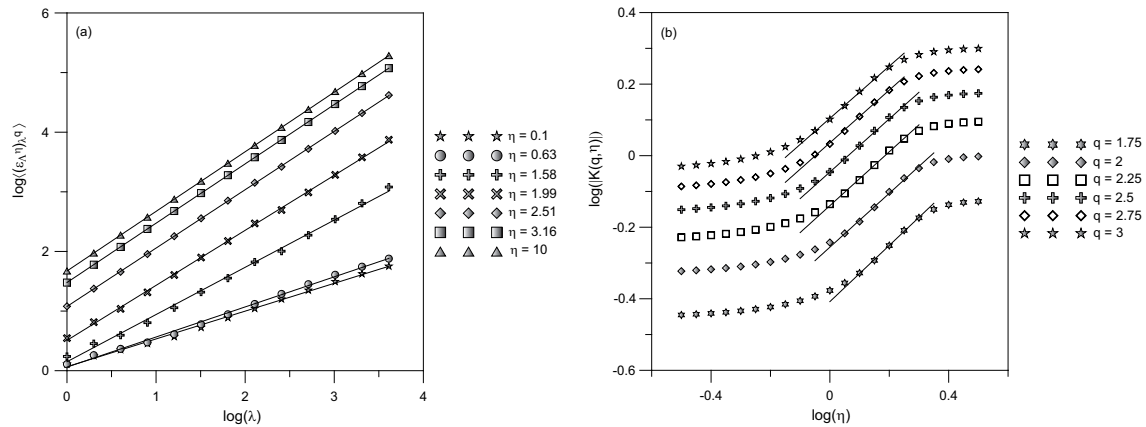


Fig. 3. Double trace moments (DTM) technique for the estimation of the parameters of the universal multifractal model  $\alpha$ ,  $C_1$  and  $H$ : (a) log-log plots of  $\langle (\epsilon_\lambda^\eta)_{\lambda^q} \rangle$  versus  $\lambda$  for  $q = 2$ , considering different values of  $\eta$ . The exponent function  $K(q, h)$  is determined as the slope of the linear trends of these plots; (b) plots of  $\log |K(q, h)|$  versus  $\log(\eta)$  for some of the investigated  $q$ -moments. The solid straight line in each plot indicates the fitted linear regression.

From Figure 3b the value of  $C_1$  can be yielded, this being  $C_1 = 0.434 \pm 0.005$ . Working with the empirical scaling exponent function  $K(q)$  shown in Figure 2, a similar value is obtained, this being  $C_1 = [dk(q)/dq]_{q=1} = 0.458$ . Parameter  $C_1$  is the measurement of the sparseness of the data. The larger  $C_1$ , the more the mean rainfall is inhomogeneous.

The estimation of  $H$  requires the spectral exponent  $\beta = 0.54$  estimated from the log-log plot of the energy spectrum for the rainfall series analyzed (Fig. 4). The energy spectrum was calculated by using a Tukey-Hanning window of 4096 hours in length with a 50% overlap. Applying Eq. (4) and taking into account the value  $K(2) = 0.561$ ,  $H = 0.04$  was yielded. The theoretical scaling exponent function is then obtained by replacing in Eq. (2) the estimated parameters  $\alpha$ ,  $C_1$  and  $H$  (Fig. 2, solid line). The agreement between this function and the empirical  $K(q)$  is acceptable for a range of values limited by the small and very high orders of the moments. The disparity between both functions for the smaller orders of moments was a consequence of the presence of zeros in the time series analyzed, because the universal multifractal model implicitly assumes zero-free data.

#### 4. Conclusions

The analysis of the 23-year long-term series of hourly rainfall recorded in Córdoba has shown that the multifractal framework is suitable for describing the temporal structure of the rainfall process at this location. Scales from 1 hour to almost 6 months have been explored in this work, and a scaling behaviour of rainfall has been detected for a range from 1 hour to almost 21 days. Although factors such as the accuracy of the rain gauge or the maximum data recorded may affect the results, the scale break interpreted as the synoptic maximum is comparable to values found in the literature for different places.

The theoretical moments scaling exponent function has been obtained by applying the universal multifractal model. It is very similar to the empirical function, with the exception of the small and very high moments. The good agreement between both functions suggests that universal multifractals can be used in generating synthetic hourly rainfall data in Córdoba.

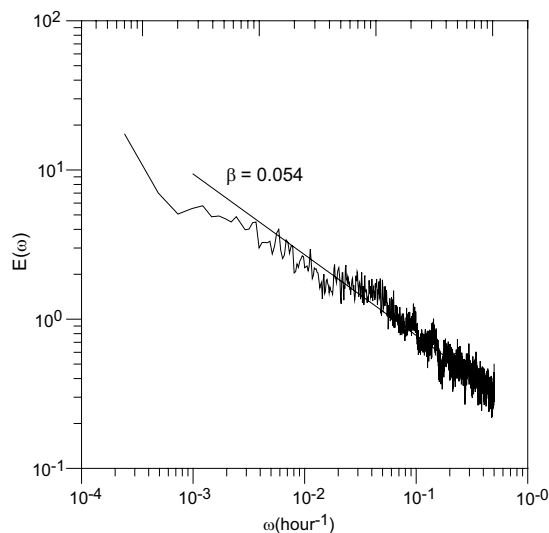


Fig. 4. Energy spectrum of the rainfall series analyzed. The spectral exponent  $\beta$  is used to determine the parameter  $H$ .

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