

Transport in Porous Media

Spatial characterization of wetting in porous media using local lattice-Boltzmann simulations --Manuscript Draft--

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Abstract:	<p>Wettability is one of the critical parameters affecting multiphase flow in porous media. The wettability is determined by the affinity of fluids to the rock surface, which varies due to factors such as mineral heterogeneity, roughness, ageing, pore-space geometry, etc. It is well known that wettability varies spatially in natural rocks, it is still generally considered a constant parameter in pore-scale simulation studies. The accuracy of pore-scale simulation of multiphase flow in porous media is undermined by such inadequate wettability models.</p> <p>The advent of in-situ visualization techniques, e.g., X-ray imaging and microtomography, enables us to characterize the spatial distribution of wetting more accurately. There are several approaches for such characterization. Most include the construction of a meshed surface of the interface surfaces in a segmented X-ray image and are known to have significant errors arising from insufficient resolution and surface-smoothing algorithms.</p> <p>This work presents a novel approach for spatial determination of wetting properties using local lattice-Boltzmann simulations. The scheme is computationally efficient as the segmented X-ray image is divided into subdomains before conducting the lattice-Boltzmann simulations, enabling fast simulations. To test the proposed method, it was applied to two synthetic cases with known wettability and three datasets of imaged fluid distributions. The wettability map was obtained for all samples using local lattice-</p>	

	<p>Boltzmann calculations on trapped ganglia and optimization on surface affinity parameters. The results were quantitatively compared with a previously developed geometrical contact angle determination method.</p> <p>The two synthetic cases were used to validate the results of the developed workflow, as well as to compare the wettability results with the geometrical analysis method. It is shown that the developed workflow accurately characterizes the wetting state in the synthetic porous media with an acceptable uncertainty, and is better to capture extreme wetting conditions.</p> <p>For the three datasets of imaged fluid distributions, our results show that the obtained contact angle distributions are consistent with the geometrical method. However, the obtained contact angle distributions tend to have a narrower span and are considered more realistic compared to the geometrical method.</p> <p>Finally, our results show the potential of the proposed scheme to efficiently obtain wettability maps of porous media using X-ray images of multiphase fluid distributions. The developed workflow can help for more accurate characterization of the wettability map in the porous media using limited experimental data, and hence more accurate digital rock analysis of multiphase flow in porous media.</p>
<p>Response to Reviewers:</p>	<p>The reply the reviewers is uploaded as a pdf file, with color coding. It is also glued in here:</p> <p>“Spatial characterization of wetting in porous media using local lattice-Boltzmann simulations” Hamidreza Erfani, Reza Haghani, James McClure, Edo Boek, Carl Fredrik Berg* November 5, 2023</p> <p>Dear Prof Blunt, We appreciate the critical comments and suggestions mentioned by the respected reviewers and editor. All comments raised by the reviewers have been addressed in the manuscript or responded to. We hope that this new revision will further clarify these issues for the readers and make this manuscript a beneficial contribution to TiPM readers. In the end, we would like to thank you for your consideration and the constructive review process. Sincerely, Carl Fredrik Berg</p> <p>Reviewer #1: The manuscript presents novel approach for spatial characterisation of contact angle from pore-scale images of natural rocks with two fluid phases by performing isolated lattice-Boltzmann simulations for each particular ganglion. I find this manuscript interesting and believe that it has scientific contributions, however, I also believe that some important aspects should be clarified before recommending it for publication. The most important thing to me is that authors claim that proposed method is more “accurate” (or more “realistic”) than geometrical method previously published to which results are compare to. However, I do not find any strong scientifically valid proof that this is actually the case. Thus, I recommend a major revision where I believe that authors should provide more relevant validation with comparison between their method and geometrical analysis that was used in the manuscript. For example, using a single high-fidelity ganglia image, or creating artificial verification geometry, could be direction to take. However, I do believe that their statements regarding accuracy between two methods should be supported by some additional proof, or something that more clearly states pros and cons of each approach. Response: As suggested by the reviewer, we have now included a new section where</p>

we have generated an artificial image with known affinities. This new example does show that the proposed workflow yields more accurate results for extreme wetting cases, as indicated by the previously considered experiments.

Additional comments are as follows:

1. There are multiple statements that proposed LBM method is computationally more efficient than existing geometrical method, yet there are no any comparison of computational cost or further discussion why this is the case.
Response: In the developed scheme the LBM model is running on individual ganglia independent of other ganglia, thus the scheme is easy to parallelize. As we are running the LBM on a subdomain, it is fast and not RAM-demanding. Since both methods are highly dependent on the implementation, we will not include numbers in the manuscript. For our given implementations the proposed workflow is however significantly faster and requires much less memory.
2. Wettability map is obtained by LBM only for existing ganglia in the segmented images, while linear interpolation is used to obtain contact angles for whole domain. What is physical relevance of doing so? Understanding that obtaining wettability for whole domain is not trivial task to deal with and that this is improvement considering that most studies still assume constant contact angle thought domain, I still think that it would be valid point to discuss further and bring into context of realistic heterogeneity of wettability.
Response: We appreciate the reviewer's comment. In this study we use direct numerical simulations using LBM on isolated ganglia to obtain the wettability. Each simulation gives information about the wettability of the considered three-phase contact line. So, aggregating all the information for a snapshot (corresponding to a saturation) can give information for three-phase contact line voxels of the sample. That we only obtain wettability along the three-phase lines holds irrespective of the method, thus this is the case for both the geometrical method and our proposed method. Even though, if the information of different μ CT snapshots is assimilated we get more information, still we miss the wettability of the majority of rock voxels. There are many options for how to calculate the surface affinity for the voxels not located at the three-phase contact line, e.g., using the closest neighbour or doing spatial 3D linear interpolation. The latter seems more realistic as it uses more spatial information from neighbouring voxels.
3. In section 2, during LBM method description, explanations of some variables in equations are missing. Moreover, I find order of writing an equation and calling them in text difficult to follow.
Response: We appreciate the comment. Section 2 has been rewritten, with definitions of missing parameters added to this section. We hope that the reviewer finds the updated section more easy to follow.
4. Line 211-212 states that having scalar affinity parameter is beneficial from optimisation standpoint. Why is this? What would be beneficial comparing to other computational methods that

directly
 use contact angle value?
 Response: It is beneficial that the wettability is controlled by a single parameter. If wettability was controlled by more parameters the optimization might not have a unique answer and more iterations might be needed to converge to the global optimum answer for each ganglion. Optimizing on affinity or the corresponding contact angle is of no consequence, as we simply relate them through $\phi_s = \cos(\theta)$.

5. Even that is existing published method, section 2.3 should contain more details regarding geometrical method as it is directly used for comparison with proposed LBM approach. Additionally, more discussion regarding the uncertainty associated with geometrically measured contact angle is needed as it is mentioned later without proper explanation.
 Response: More explanation of geometrical contact angle measurement is provided in the revised manuscript, while the original reference is also provided for deeper discussions. To avoid repetition, we think the reader should be referred to the original reference for detailed formulation as the algorithm is fully adopted from [1].

6. Line 245-247. I think that investigating a non-uniform wetting property of single ganglia could be significant contribution to this work and potentially show additional benefit of proposed method.
 Moreover, from this work, it is not clear if geometrical method uses uniform or non-uniform contact angle for single ganglia.
 Response: In this study, we only deal with assuming a single surface affinity for each ganglion to show the proof of concept. Assuming spatial variation of wettability for a single ganglion is out of the context of this study as it adds much more complexity, computational load and uncertainty regarding the non-uniqueness of the optimization results.
 To address this concern and show the improvement of the developed workflow over the geometrical wettability characterization we added 2 numerical validation cases in Section 3. We think that this provides a clearer view of improvement over the geometrical method as well as the accuracy of each of the methods.
 The geometrical method calculates a contact angle for each mesh-boundary along the three-phase contact line (approximately on for each grid-cell along the three-phase contact line), thus it yields different contact angles along the three-phase for a single ganglion. This has now been specified in the text, see section 2.3.

7. Line 310 - add explanation of subscripts for oil and brine.
 Response: This has now been added to the manuscript.

8. Section 2.5 - explain difference between unaltered and altered samples.
 Additionally, injection rates in real units could be stated.
 Response: This has now been added to the manuscript, see updated section 2.5.

9. Line 350 has typo: oil-wetting should be water-wetting.
 Response: This is revised in the manuscript text.

10. Line 384 has typo: altered should be unaltered.
 Response: We believe this should be "altered", the alteration renders the sample more oil-wetting.

11. Line 391-393. I don't see point of this statement. What is relevant here?
 Response: This part compared the altered and unaltered contact angles obtained from both schemes all together to compare them in the same plot. We expect the altered sample to show a shift toward oil wetting behaviour compared to the unaltered sample. Surprisingly the unaltered sample CA distribution from geometrical analysis is the same as the altered sample CA distribution from the developed scheme.

12. Line 393-397 I would like see more discussion here as this is important point to make considering the aim of the manuscript.
 Response: We appreciate the reviewer's comment. Generally, we added more discussions to the results and discussion section and also added section 33.1 with 2 cases for validating and comparing the developed algorithm and geometrical analysis method results.

Reviewer #2:

1. It is recommended to start the abstract with general information about the topic.
 Response: The abstract in the revised version opens with general sentences highlighting the importance of the topic.
 "Wettability is one of the critical parameters affecting multiphase flow in porous media. The wettability is determined by the affinity of fluids to the rock surface, which varies due to factors such as mineral heterogeneity, roughness, aging, pore-space geometry, etc. Despite that it is known that wettability varies spatially in natural rocks, it is generally considered a constant parameter in pore-scale simulation studies. The accuracy of pore-scale simulation of multiphase flow in porous media is undermined by such inadequate wettability models."

2. It is suggested to discuss more about the findings of this study in the abstract.
 Response: We added more findings to the abstract in the revised version. The changes in the abstract can be followed in the annotated manuscript.

3. It is recommended to mention about the applications of this study at the end of abstract: The findings of this study can help for better understanding of . . .
 Response: We added this to the end of the abstract in the revised manuscript.

4. I strongly recommend the authors to add one paragraph discussing the difference between their work and the previously performed studies in literature. In other words, what is the novelty of this work? I offer the authors to revise the abstract and introduction in order to incorporate the novelty of their work. This change motivates the readers of "Transport in Porous Media" to study this work with interest.
 Response: This is thoroughly discussed throughout the results and discussion section. We compare our method extensively with the geometrical analysis method. We also added a new section to the results and discussion (section 3.1) that we utilize two synthetic cases for validation and quantitative comparison between the developed scheme and geometrical analysis.

5. It is recommended to include a paragraph at the end of introduction to present the steps of the work like: First, the methods and materials are described. Then, . . .
 Response: The last paragraph of the introduction serves this specific purpose in the revised

manuscript.

6. What are the advantages and disadvantages of this study? I recommend the authors to highlight this topic. What are the limitations of this study? I recommend the authors to highlight this topic.

Response: We appreciate the reviewer's comment. We believe that this purpose is served through the results and discussion section as well as section 2.4.

7. The quality of all the figures should be improved.

Response: All the figures are added to the manuscript as vector graphics to accommodate high-quality and infinite zoom features.

8. It is recommended to add minor ticks (or intervals) on horizontal and vertical axis of all the plots.

Response: This is now added to all figures in the revised manuscript.

9. The title for last section should be changed to "Summary and Conclusions".

Response: This is now changed in the revised manuscript.

10. It is suggested to add a nomenclature (including alphabetic letters, Greek letters, subscripts, and superscripts).

Response: We have done our best to define variables in their appropriate places. The nomenclature will be mostly LBM parameters which is not the core contribution of the manuscript, we therefore think that defining the parameters next to the equations will keep the main focus of the article better.

Reviewer #3:

This paper proposes a workflow to characterize local wetting distributions in porous media using X-ray images of two-phase fluid distributions. It first identifies isolated ganglia that are larger than a specific size. These ganglia are then considered as initial guesses for lattice Boltzmann simulations. The LB simulation is performed with different surface affinity values, and the single value that results in a geometry with the minimum difference compared to the initial guess, determined through voxel-by-voxel comparison, is considered the wettability of the contact line. This work is clear and to the point, addressing an interesting problem and proposing an alternative workflow to determine wettability. Although it is an interesting study, I still have a few concerns which I address below. The image token was acquired under dynamic flow conditions, while the simulation was performed under static conditions. This situation could lead to uncertainty in the determined surface affinity unless I am misunderstanding something here.

Response: We appreciate the comment. We used 2 different experimental datasets to validate the developed workflow. The first dataset is water injection into a core sample at residual gas saturation. The other oil-brine experimental data are from steady-state co-injection of fluids into a sample with 1:1 volumetric ratio. In both conditions, we assume that the trapped ganglia are not moving, which is realistic in a snapshot. As the surrounding fluid is moving around the ganglion, this can result in a slight deformation compared to both fluids being stagnant. However, the capillary forces are assumed strong compared to the viscous drag on the ganglion, we therefore do not expect large changes to a fully stagnant case.

While assuming a single value for the surface affinity of each ganglion might be reasonable for small ganglia, it could be inaccurate in cases where we have large ganglia that span a significant region.

Response: This is correct. It is a source of uncertainty and error in the current workflow, but for the time being the most viable option was assuming a single value for each LBM simulation. If we wanted to assume multiple local affinities for each ganglion the optimization calculations would be much more expensive, and it might result in several local optima for the optimization. This is, however, a problem that we will try to tackle in future work.

Could another method be employed to determine the spatial variation of wettability across the sample instead of interpolation? Methods that can consider the spatial variogram?

Response: Any method can be employed in the stage of populating local wettability information to the whole sample (all solid-fluid voxels). We chose simple 3D interpolation for the sake of simplicity in this post-processing stage to avoid complications which might make the interpretation of results more difficult.

Since all determined wettabilities are based on isolated ganglia, could this introduce bias in the simulation compared to the geometrical method? I didn't understand if the geometrical results are solely based on isolated ganglia or if any other three-phase lines are involved. Cause based on what was mentioned, the percolating clusters or ganglia that cover a wide range of pore space and hold more interest cannot be considered in this wettability analysis.

Response: This is absolutely correct. To avoid such bias the percolating ganglia were excluded from the geometrical analysis as well and both methods were applied to the same ganglia. More explanation is added to Section 2.4 for clarification.

Considering the method requires numerous lattice Boltzmann simulations, how does its computational efficiency compare to other geometrical methods?

Response: The LBM simulations were run on a subdomain including a single ganglion so LBM simulations are very RAM-efficient, but in the geometrical method, the whole solid-fluid and fluid-fluid surfaces need to be meshed. Moreover, as the LBM simulations are independent of each other they can be easily parallelized to be run at the same time. For our cases, the LBM method was faster. This is, naturally, strongly dependent on the implementation, and might change with a more efficient implementation of the geometrical method.

I might be mistaken, but in line 350, there appears to be a mistake. Do both distributions in Fig 4 display oil-wetting behavior? Given that it's a gas-water system, should it exhibit water-wetting behavior instead?

Response: We highly appreciate the reviewer's comment. This was a typo in the manuscript and is now revised. This sample was saturated with gas and water and the distributions show liquid wetting conditions as expected.

References
[1] Hamid Hosseinzade Khanamiri, Per Arne Slotte, and Carl Fredrik Berg. Contact angles in two-phase

1 Spatial characterization of wetting in porous media
2 using local lattice-Boltzmann simulations

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5 **Abstract**

6 Wettability is one of the critical parameters affecting multiphase flow in
7 porous media. The wettability is determined by the affinity of fluids to
8 the rock surface, which varies due to factors such as mineral heterogeneity,
9 roughness, ageing, pore-space geometry, etc. It is well known that wettabil-
10 ity varies spatially in natural rocks, it is still generally considered a constant
11 parameter in pore-scale simulation studies. The accuracy of pore-scale simu-
12 lation of multiphase flow in porous media is undermined by such inadequate
13 wettability models. The advent of in-situ visualization techniques, e.g., X-
14 ray imaging and microtomography, enables us to characterize the spatial dis-
15 tribution of wetting more accurately. There are several approaches for such
16 characterization. Most include the construction of a meshed surface of the in-
17 terface surfaces in a segmented X-ray image and are known to have significant
18 errors arising from insufficient resolution and surface-smoothing algorithms.
19 This work presents a novel approach for spatial determination of wetting
20 properties using local lattice-Boltzmann simulations. The scheme is compu-
21 tationally efficient as the segmented X-ray image is divided into subdomains
22 before conducting the lattice-Boltzmann simulations, enabling fast simula-
23 tions. To test the proposed method, it was applied to two synthetic cases
24 with known wettability and three datasets of imaged fluid distributions. The
25 wettability map was obtained for all samples using local lattice-Boltzmann
26 calculations on trapped ganglia and optimization on surface affinity param-
27 eters. The results were quantitatively compared with a previously developed
28 geometrical contact angle determination method. The two synthetic cases

29 were used to validate the results of the developed workflow, as well as to
30 compare the wettability results with the geometrical analysis method. It is
31 shown that the developed workflow accurately characterizes the wetting state
32 in the synthetic porous media with an acceptable uncertainty, and is better
33 to capture extreme wetting conditions. For the three datasets of imaged
34 fluid distributions, our results show that the obtained contact angle distri-
35 butions are consistent with the geometrical method. However, the obtained
36 contact angle distributions tend to have a narrower span and are considered
37 more realistic compared to the geometrical method. Finally, our results show
38 the potential of the proposed scheme to efficiently obtain wettability maps
39 of porous media using X-ray images of multiphase fluid distributions. The
40 developed workflow can help for more accurate characterization of the wet-
41 tability map in the porous media using limited experimental data, and hence
42 more accurate digital rock analysis of multiphase flow in porous media.

43 *Keywords:* Wettability, contact angle, lattice Boltzmann method, digital
44 rock modelling, multiphase fluid flow, porous medium

45 1. Introduction

46 Multiphase simulation of flow and transport across different scales is rele-
47 vant to a vast range of natural and industrial applications, e.g., enhanced oil
48 recovery [1, 2, 3], geological CO₂ and hydrogen storage [4, 5, 6, 7], optimisa-
49 tion of fuel-cells operation [8, 9], industrial energy storage [10, 11], etc. The
50 accuracy of continuum-scale multiphase modelling depends on constitutive
51 relations such as relative permeability and capillary pressure. These can be
52 obtained numerically using digital rock physics, or through time-consuming
53 and expensive experimental procedures (i.e., core flooding, two-phase cen-
54 trifuge, etc.) [12].

55 In recent decades, with improvements in X-ray tomography and other vi-
56 sualisation techniques, we can visualize inside porous media and characterise
57 events, fluid distribution, and in-situ distribution of phases in porous media.
58 Such advances helped us develop and verify digital rock physics techniques
59 to simulate multiphase flow in porous media [12]. Hence, pore-scale models
60 are a backbone of existing predictive models, either for phenomenological
61 studies, upscaling of constitutive relations or avoiding experimental routines
62 [13].

63 Recent developments in the realm of pore-scale modelling of multiphase

64 flow focus on simulating flow through a realistic digital representation of
65 complex pore space [14, 15]. It is well-understood that the pore structure
66 in combination with fluid-fluid and solid-fluid interfacial forces governs the
67 dynamics of multiphase flow and the geometry and morphology of phases
68 [16, 17, 18]. The pore structure of porous media can be obtained with good
69 accuracy using μ -CT X-ray imaging, additionally, there are well-established
70 experimental methods to accurately measure fluid-fluid interfacial forces. In
71 contrast, methods for assigning solid-fluid interfacial forces are not yet estab-
72 lished. Fluid-fluid interfacial forces can often be considered constant through-
73 out the simulation domain, and can therefore be described by simple models.
74 Solid-fluid interfacial forces are usually imposed into the model using the
75 wetting state of the solid wall and can have significant variation throughout
76 the simulation domain, complicating its description [19].

77 Wettability is the relative preference of one fluid to coat the solid sur-
78 face in the existence of another immiscible fluid, which is usually quantified
79 by the contact angle [20]. The pore surface wettability determines the local
80 balance of capillary forces and controls the local fluid distribution and mor-
81 phology. The assignment of wettability to pore-scale models has long been
82 recognised as having a strong impact on effective multiphase flow, and the
83 lack of established methods has been recognised as the most significant issue
84 for predictive multiphase modelling [21, 22]. This long-lasting challenge is
85 still an area of debate, and the in-situ characterisation of wettability is still
86 considered inaccurate [23].

87 Despite recent attempts, measuring in-situ contact angles on the fluid-
88 solid-fluid (three-phase contact) line exhibits significant uncertainty and error
89 due to contact angle hysteresis, effects of pore structure complexity, surface
90 roughness, noise in image capturing, processing and segmentation [24, 25, 26].
91 The pore-scale observations of two-phase fluid flow can be averaged into an
92 equilibrium contact angle for the whole sample [27], while it is well-known
93 that the wetting state can be spatially variable in the pore space due to
94 numerous reasons, e.g., mineralogy, clay coating, roughness, polar component
95 coating, etc. [12, 28, 29, 30, 31].

96 Various approaches were followed by different researchers to measure local
97 contact angles in porous media to make a wettability map instead of assigning
98 a constant wetting property to the whole sample with the aim of a more
99 accurate multiphase flow simulation. Most of the studies use geometrical
100 approaches to measure the local contact angle on the three-phase contact
101 line [32, 33, 34]. Deficit curvature of the fluid and solid interfaces have

102 also gained significant attention recently [35, 36, 37]. Some studies used
103 the extracted pore network of the sample to measure local wettability, e.g.,
104 Mascini et al. [38] used the extracted network of the sample to identify pore-
105 filling events (Haines jump) to calculate local wettability [39, 40]. Garfi
106 et al. [26] used the solid surface coverage in the pore regions to calculate the
107 local wettability. Moreover, Foroughi et al. [41] used a pore-by-pore scheme
108 to optimise the local wetting property using experimental dynamic X-ray
109 tomography of capillary-dominated displacement in the porous media.

110 In this paper, we introduce a new automated workflow to characterise
111 the local wetting distributions in porous media using X-ray images of two-
112 phase fluid distribution. For the multiphase simulation, we use the lattice
113 Boltzmann method (LBM), and the model of choice is called the colour-
114 gradient model which was originally proposed by Gunstensen et al. [42] and
115 modifications later on [43]. Yang and Boek [44] showed the capabilities of
116 colour-gradient LBM to simulate the flow of binary fluids with high viscosity
117 contrast and high numerical stability. In this model, the contact angle can be
118 simply defined by an affinity parameter [43]. Simulations are conducted on
119 isolated ganglia to get the wetting properties on the three-phase contact lines.
120 Consequently, we have one LBM simulation for each ganglion, reducing the
121 simulation domain and thereby improving the computational efficiency. We
122 used a derivative-free optimisation scheme to minimize the difference between
123 the fluid distribution from the LBM simulations and the imaged ganglia ge-
124 ometry by varying the surface affinity parameter in the colour-gradient LBM.
125 Local wetting properties information was populated into the whole sample by
126 a three-dimensional (3D) linear interpolation technique to construct a wet-
127 tability map for the whole sample. Then, the proposed workflow was utilised
128 to characterise the wettability of three publicly available datasets, and the
129 obtained results were compared with geometrical curvature analysis.

130 This manuscript is organised as follows: In the next section, Section 2,
131 we will present the methods used in this paper; the LBM method, wettabil-
132 ity determination from geometrical analysis, and wettability determination
133 through our introduced method using LBM simulations. We will also present
134 three experimental data sets used for testing our introduced methodology. In
135 Section 3, first we provide numerical validation, in which we compare the re-
136 sults of the developed workflow with the more traditional geometrical analysis
137 method for contact angle calculation for two synthetic cases. Then, we will
138 present the contact angle distribution obtained by applying the geometrical
139 method and the introduced LBM simulation method for the three experi-

140 mental cases. Further, we will compare and discuss the results. In the last
 141 section, Section 4, we summarize and conclude.

142 2. Methods and Materials

143 This section starts by introducing the methods used in this article and
 144 is followed by a presentation of the experimental data. The experimental
 145 data will be used in the next section to test our introduced methodology
 146 for wettability characterization. As our introduced methodology is based on
 147 LBM simulations, we start this section with a brief overview of the LBM
 148 method.

149 2.1. Colour-gradient lattice-Boltzmann model

150 The LBM is a computationally efficient and alternative approach to the
 151 classical computational fluid dynamics (CFD) methods to model flow and
 152 transport in complex geometries, which can be adapted to a wide range of
 153 applications. Models based on the LBM have been widely applied to study
 154 multiphase flow in porous media, with the colour-gradient model being the
 155 most popular implementation of two-fluid simulation [45, 46]. In this study,
 156 we used the LBPM open-source package [43], which is optimised to simulate
 157 incompressible immiscible two-fluid flows in μ CT images of porous media.
 158 LBPM has routines to simulate unsteady displacement, steady-state flow at
 159 fixed saturation, and mimic centrifuge experiments [43]. Here we provide a
 160 brief introduction to lattice Boltzmann (LB) formulation and wetting state
 161 implementation of the LBPM for the sake of completeness, and the readers
 162 are referred to McClure et al. [43] for more detailed formulation and in-depth
 163 discussion.

164 The colour-gradient LB model is defined based on three sets of lattice
 165 Boltzmann equations (LBE) to capture the interfaces between components
 166 a and b , and hydrodynamic properties. For incompressible immiscible mix-
 167 ture, the number density of two fluids, N_a and N_b , must be conserved. Two
 168 particle-distribution functions, A and B , are defined which are governed by
 169 the following LBEs:

$$A_q(\mathbf{x} + \boldsymbol{\zeta}_q \delta t, t + \delta t) = \omega_q N_a \left[1 + \frac{\mathbf{u} \cdot \boldsymbol{\zeta}_q}{c_s^2} + \beta \frac{N_b}{N_a + N_b} \mathbf{n} \cdot \boldsymbol{\zeta}_q \right] \quad (1a)$$

$$B_q(\mathbf{x} + \boldsymbol{\zeta}_q \delta t, t + \delta t) = \omega_q N_b \left[1 + \frac{\mathbf{u} \cdot \boldsymbol{\zeta}_q}{c_s^2} - \beta \frac{N_a}{N_a + N_b} \mathbf{n} \cdot \boldsymbol{\zeta}_q \right] \quad (1b)$$

171

172 where \mathbf{x} and t are the space and time, respectively, δt is the time step and
 173 \mathbf{u} is the flow velocity computed from hydrodynamic equations. The β pa-
 174 rameter controls the thickness of the interface of the two components, and
 175 \mathbf{n} is the unit normal vector of the interface. In the above equations, $\boldsymbol{\zeta}_q$ and
 176 ω_q are the microscopic velocity and the weighting factor in the q -direction
 177 of a lattice model, respectively. For our three-dimensional (3D) modelling, a
 178 seven-lattice velocity model ($Q=7$) is chosen for the above-mentioned equa-
 179 tions, i.e., D3Q7. As such, the weights are $\omega_0 = \frac{1}{3}$, $\omega_{1,\dots,6} = \frac{1}{9}$ and the mi-
 180 croscopic velocities are the first seven directions in Eq. (6), i.e., $q = 0, \dots, 6$.
 181 The speed of sound for this lattice is given by $c_s = \frac{\sqrt{2}}{3}$. The number density
 182 of the two fluids is computed by the zeroth-moments of the corresponding
 183 particle distribution functions:

$$N_a = \sum_{q=0}^{Q-1} A_q \quad (2a)$$

184

$$N_b = \sum_{q=0}^{Q-1} B_q \quad (2b)$$

185 A phase indicator field, ϕ , is defined to locate the interface by the density
 186 of the two fluids:

$$\phi = \frac{N_a - N_b}{N_a + N_b} \quad (3)$$

187 The unit normal vector of the colour gradient is calculated as

$$\mathbf{n} = \frac{\nabla \phi}{|\nabla \phi|} \quad (4)$$

188 In addition to the mass transport equation, momentum needs to be trans-
 189 ported by an LBE to solve for two-fluid flow in porous media. A new particle
 190 distribution function, f , is defined which is governed by the following LBE
 191 equipped with the multi-relaxation-time (MRT) collision operator:

$$f_q(\mathbf{x} + \boldsymbol{\zeta}_q \delta t, t + \delta t) - f_q(\mathbf{x}, t) = \sum_{k=0}^{Q-1} M_{q,k}^{-1} S_k M_{q,k} (f_k^{eq} - f_k) \quad (5)$$

192 where M and M^{-1} are the orthogonal transformation matrix and its inverse,
 193 respectively [47], and S is the diagonal relaxation matrix containing the re-
 194 laxation rates for each moment. Among the relaxation rates, the relaxation

195 rate τ is related to the kinematic viscosity of the fluid with $\nu = c_s^2(\tau - 0.5)$,
 196 while the rest of the relation times are determined based on accuracy and
 197 numerical stability. Generally, the MRT collision operator is employed to
 198 perform the collision in a moment space instead of a discrete velocity space
 199 which results in higher stability in lower kinematic viscosities. The equilib-
 200 rium distribution function, f_k^{eq} , and matrices M , M^{-1} , and S are defined in
 201 detail in McClure et al. [43], and their definitions are not mentioned here
 202 for the sake of brevity.

203 The hydrodynamic LBE (Eq. (5)) is solved on a popular D3Q19 lattice
 204 which is used by the LBPM package. For this lattice, the velocity set is

$$\zeta_q = \begin{cases} \{0, 0, 0\}^T & \text{for } q = 0 \\ \{\pm 1, 0, 0\}^T & \text{for } q = 1, 2 \\ \{0, \pm 1, 0\}^T & \text{for } q = 3, 4 \\ \{0, 0, \pm 1\}^T & \text{for } q = 5, 6 \\ \{\pm 1, \pm 1, 0\}^T & \text{for } q = 7, 8, 9, 10 \\ \{\pm 1, 0, \pm 1\}^T & \text{for } q = 11, 12, 13, 14 \\ \{0, \pm 1, \pm 1\}^T & \text{for } q = 15, 16, 17, 18 \end{cases} \quad (6)$$

205 and the weights are $\omega_0 = \frac{1}{3}$, $\omega_{1,\dots,6} = \frac{1}{18}$, and $\omega_{7,\dots,18} = \frac{1}{36}$. The speed of
 206 sound for D3Q19 lattice is $c_s = \frac{1}{\sqrt{3}}$. By solving the LBE (Eq. (5)), one can
 207 determine the flow velocity based on the first moment of the distribution
 208 function:

$$\mathbf{u} = \frac{\sum_{q=0}^{Q-1} f_q \zeta_q}{\sum_{q=0}^{Q-1} f_q} \quad (7)$$

209 This velocity is used in the mass transport equations, Eqs. (1).

210 2.2. Modelling wettability with lattice-Boltzmann models

211 One of the essential pieces of the multiphase flow simulation puzzle in
 212 porous media is the role of wettability, which affects flow on all scales. There-
 213 fore, it is important to allow for various implementations of wetting condi-
 214 tions to enable the simulation of different scenarios to be studied. There are
 215 various approaches for imposing wetting behaviour into LBMs [45]. In these
 216 models, the wetting characteristic can be applied as an upscaled property
 217 instead of taking the impact of different role-playing parameters like rough-
 218 ness, surface charge, film existence, etc. Note that some researchers used
 219 LBMs to study these effects on the upscaled behaviour of contact lines [23].

220 In the colour-gradient LBM, the wetting condition can be defined in the
 221 form of a scalar affinity value. It is demonstrated that the affinity value is
 222 equivalent to a pseudo-phase indicator field, ϕ_s , ranging from -1 to 1. It
 223 should be noted that the subscript s stands for the solid phase. This simple
 224 method imposes the expected contact line behaviour for the stationary and
 225 moving contact lines [48]. On a flat surface with a well-defined contact line,
 226 the equilibrium contact angle (θ_{eq}) is related to ϕ_s by

$$\cos \theta_{eq} = \phi_s \quad (8)$$

227 Note that ϕ_s equals the thermodynamically based wettability index $\omega_i =$
 228 $(\sigma_{bs} - \sigma_{as})/\sigma_{ab}$ introduced in [49], where σ_{as} and σ_{bs} are the surface tension
 229 between the solid phase and fluid a and b , respectively, while σ_{ab} is the
 230 interfacial tension between fluid a and b .

231 Figure 1 shows the relation between the surface affinity parameter and
 232 contact angle for a ganglion. In this study, we stick to the commonly used
 233 terminology and refer to the local wetting state as $\phi_s = 1$ being a strongly
 234 water-wet state and $\phi_s = -1$ being a strongly oil-wet state.

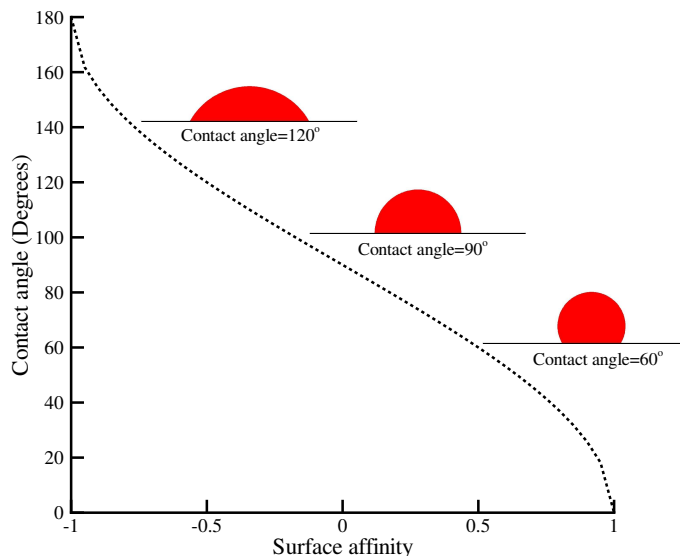


Figure 1: Contact angle (degrees) versus surface affinity parameter (ϕ_s) for a ganglion.

235 2.3. Geometrical contact angle determination

236 For geometrical contact angle measurements, we followed the automated
 237 method proposed by Khanamiri et al. [32] to calculate the contact angle

238 along the three-phase contact line in the segmented experimental X-ray to-
239 mography data of two-phase fluid distributions. To calculate local curvature
240 and contact angles, first, a triangulated mesh was applied to the fluid-fluid
241 and solid-fluid interfaces of the segmented image. After correction for some
242 artefacts in the generated mesh due to imaging resolution limitations, e.g.,
243 self-touching interfaces, the mesh was smoothed. The smoothing was ap-
244 plied following the method suggested by AlRatrouf et al. [50] with some
245 modification, in which the vertices are displaced to minimize the curvatures
246 while keeping the global phase volumes approximately constant [51]. In the
247 smoothing procedure, some tuning parameters are used which may result in
248 the divergence of meshes from the intended geometry. As such, the final
249 mesh may not imitate porous structures in low-resolution images, and as a
250 result, the measured properties such as the measured contact angle and the
251 curvature can deviate [32].

252 After smoothing the mesh, the contact angle between the solid-fluid and
253 fluid-fluid meshes on a vertex on the three-phase contact line is calculated by
254 the dot product of the unit normal vectors of the solid-fluid and fluid-fluid
255 interfaces. Note that the geometrical contact angle method yields different
256 contact angles along the three-phase contact line. As will be described in
257 the next section, our proposed method yields a single affinity value for each
258 ganglion.

259 Also, the local average curvature for an interface is calculated based on
260 the dot product of the Laplace-Beltrami operator [51] and the unit normal
261 vector. Further details of the geometrical method can be found in Khanamiri
262 et al. [32]. They showed that to have a good estimation of the mean contact
263 angle, the fluid clusters with at least a few thousand vertices at the fluid-fluid
264 interfaces should be considered for the computation. By investigating two
265 analytical examples with known contact angles and curvatures, they found
266 that the computed point-wise contact angles, average mean curvature, and
267 interfacial area converge by increasing the grid resolution (or increasing the
268 size of clusters) while the point-wise mean curvature does not converge.

269 *2.4. Wettability characterisation workflow*

270 In this study, we propose a novel workflow to characterise the local wetting
271 properties in a sample under two-phase conditions using local LBM simula-
272 tion on individual ganglia. The procedure starts with a segmented image of
273 a sample containing two immiscible phases. We assume that all interfaces
274 are stagnant, either under steady-state conditions (i.e., the flow takes place

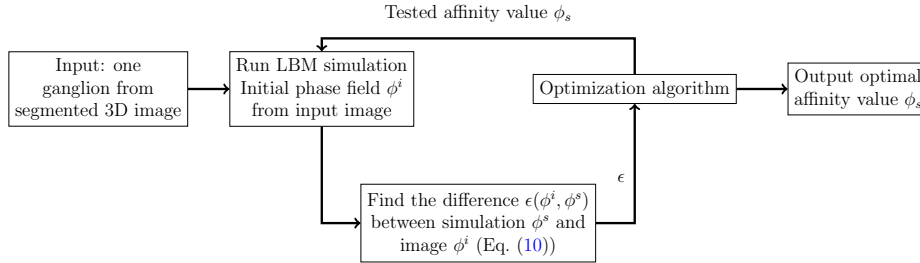


Figure 2: A flowchart of the optimization loop, finding the affinity value ϕ_s that minimizes the difference ϵ in the phase field between the imaged ganglion ϕ^i and the simulated one ϕ^s .

275 in connected pathways of phases from the inlet to the outlet), or the injec-
 276 tion of fluids is stopped and the image is taken after equilibration of phases
 277 inside the sample. Moreover, we assume that the wetting property along the
 278 three-phase contact line is uniform, so that each ganglion is associated with a
 279 single tuning parameter (i.e., surface affinity) for a more robust optimisation
 280 procedure.

281 The first step in our workflow is the identification of all trapped ganglia
 282 of the phase with lower saturation in the sample and disconnected from the
 283 boundaries of the image. In the gas-water sample, we found the trapped
 284 ganglia of the gaseous phase, while in the oil-water samples, we found the
 285 trapped ganglia of the oil phase. In both cases, we then identified the trapped
 286 ganglia of the non-wetting phase. When we identify the ganglia of the non-
 287 wetting phase, we consider two voxels connected if they share a face. This
 288 ganglia labelling was conducted using the `scipy` Python-library, giving each
 289 ganglion a specific number. We only considered ganglia of size larger than
 290 10 voxels, removing all ganglia smaller than this cut-off value from our set
 291 of ganglia. This cut-off value is arbitrary, however, including smaller ganglia
 292 than 10 voxels seemed to increase the noise in our results.

293 The optimization procedure for finding the wetting of a single ganglion is
 294 outlined in Fig. 2. In this optimization procedure we first find the smallest
 295 rectangular cuboid encapsulating the ganglion under consideration, and then
 296 enlarge this domain by five voxels in each direction to ensure that the ganglion
 297 is not interacting with the domain boundaries. Then we set all fluid voxels
 298 outside this ganglion to the wetting phase, so that the ganglion contains all
 299 non-wetting phases inside the domain. We then initialise the LBM from this
 300 fluid distribution, so that we have the initial phase field $\phi^i = 1$ for all wetting

301 phase voxels and $\phi^i = -1$ for non-wetting phase voxels. In other words, the
 302 phase field is -1 inside the considered ganglion and 1 outside.

303 The optimisation parameter for each set of simulations (i.e., each gan-
 304 glion) was the surface affinity parameter, ϕ_s , and the optimisation was per-
 305 formed with a derivative-free algorithm (Nelder-Mead). For a given surface
 306 affinity value ϕ_s provided by the optimization algorithm, we associate all
 307 solid surface faces with this affinity value. We then performed an LBM sim-
 308 ulation inside the domain with no flow boundary conditions until we achieved
 309 a relaxed system. From the relaxed system, we binarised the phase field as:

$$\phi^s = \begin{cases} 1 & \text{if } \phi \geq 0 \\ -1 & \text{if } \phi < 0 \end{cases} \quad (9)$$

310 Note that the initial phase field is already -1 for voxels containing the non-
 311 wetting phase and 1 for voxels containing the wetting phase. We then cal-
 312 culated an error function from a voxel-by-voxel comparison of the simulated
 313 ganglion ϕ^s with the initial X-ray image geometry of the ganglion as given
 314 by ϕ^i :

$$\epsilon = \epsilon(\phi^i, \phi^s) = \sqrt{\frac{1}{|\Omega|} \int_{\Omega} (\phi^i - \phi^s)^2 dV} \quad (10)$$

315 where Ω is the pore space of the rectangular cuboid domain. After conver-
 316 gence of the optimisation loop, the obtained affinity value is assigned to the
 317 3-phase contact line voxels in the X-ray image before moving to the next
 318 ganglion.

319 After optimising the surface affinity values for all ganglia, we obtained
 320 surface affinity values along all three-phase contact lines for these ganglia.
 321 These surface affinity values can then be translated to contact angles, using
 322 Eq. (8), for comparison with contact angle values obtained using the geomet-
 323 ric method. Noteworthy, to make the comparison between the geometrical
 324 analysis and the developed workflow consistent, the percolating ganglia were
 325 excluded from the geometrical analysis as well. For this purpose, a new do-
 326 main only including the analysed ganglia in the wettability characterization
 327 workflow was generated and fed into the geometrical analysis model.

328 For multi-phase flow simulations, a wettability map on all solid surfaces
 329 is required. While this is typically given as a single value, our workflow en-
 330 ables the distribution of wettability according to the obtained affinity values
 331 from the workflow outlined in Fig. 2. For this end, we populate all solid-fluid

332 voxels of the segmented X-ray image using three-dimensional linear inter-
333 polation between the solid-fluid voxels on the three-phase contact lines, as
334 these surfaces on the three-phase contact lines already have assigned surface
335 affinity values from the workflow above. For the pore surface close to the
336 boundaries of the domain, where the three-dimensional linear interpolation
337 has insufficient data, the surface affinity value was chosen equal to the nearest
338 point with an assigned value.

339 2.5. Experimental datasets

340 We used three publicly available experimental μ -CT datasets of imaged
341 two-phase fluid distributions to assess the results of the developed wettability
342 characterisation workflow. These datasets were also used by other researchers
343 for similar purposes which gave a good benchmark to compare our obtained
344 results.

345 **Gas-water Bentheimer sample:** This dataset was originally obtained
346 by Sun et al. [35] and was used to get the wetting state using their theoret-
347 ical development of geometrical analysis. A bench-top helical μ -CT scanner
348 was used to image the 3D configuration of air-water immiscible fluids under
349 ambient conditions in an untreated Bentheimer sandstone sample (4.9 mm
350 in diameter and 10 mm long). The sample was flooded with brine with an
351 injection rate of 3.3×10^{-7} m/s until irreducible air saturation was obtained
352 ($S_w = 0.93$). The images were acquired with a resolution of $4.95 \mu\text{m}$ and we
353 used a sub-volume of $720 \times 891 \times 891$ voxels from the segmented image.

354 **Oil-water Bentheimer samples (unaltered and altered):** These
355 datasets were previously created from laboratory observations and discussed
356 by Lin et al. [52, 53], and recently used by Garfi et al. [26] for the spa-
357 tial distribution of wettability. These datasets consist of extensive sets of
358 quasi-static co-injection at different fractional flows of oil (decalin) and water
359 phases in cylindrical rock samples (diameter of 6.1 mm). The brine fraction
360 of flow $f_b = Q_b / (Q_b + Q_o)$ was defined as the ratio of the brine injection flow
361 rate, Q_b to the total injection rate, $Q_b + Q_o$, where Q is the injection flow
362 rate and b and o subscripts denote brine and oil phases, respectively. For
363 the unaltered sample, the oil phase (decalin) drainage was performed into
364 the brine-saturated sample using centrifugation to reach irreducible water
365 saturation. The imbibition was performed right after drainage in capillary
366 dominant condition [54]. Brine imbibition was conducted in seven steps, in-
367 creasing the fractional flow of brine from 0 to 1 ($f_b=0, 0.05, 0.15, 0.30, 0.50,$
368 $0.85, 1$). The injection was continued until reaching the steady state, and

369 once the pressure difference across the sample was equilibrated the μ -CT
370 images were acquired for each fractional flow with a resolution of $3.58 \mu\text{m}$.
371 For the altered sample, the fractional flows were performed for $f_b=0, 0.02,$
372 $0.06, 0.24, 0.50, 0.80, 0.90, 1$.

373 The aging process applied to the modified Bentheimer dataset occurred
374 between the drainage stage and the co-injection fractional flow stages of oil
375 and brine. Following the drainage phase, the sample was placed in crude
376 oil, causing decalin to be gradually replaced by crude oil through diffusion.
377 The sample was then kept in crude oil for 30 days at a temperature of 80
378 $^{\circ}\text{C}$. Following this modification process, the sample was immersed in decalin
379 to replace the crude oil. Afterwards, the coreflooding experiment was con-
380 ducted. The composition of brine and more detail regarding the aging process
381 for the altered sample can be found in Lin et al. [52, 53]. For both datasets,
382 our analysis was performed on a cubic region of 800^3 voxels corresponding to
383 2.86^3 mm^3 , of the image acquired for the fractional flow of $f_b=0.5$.

384 3. Results and Discussion

385 In this section, we provide the obtained results and discuss the differences
386 with geometrical analysis of the oil ganglia to obtain the wetting properties of
387 the natural rocks filled with two fluid phases. Firstly, we present a simplified
388 example for numerical validation, then we provide the results of Bentheimer
389 sandstone with gas-water fluids, and finally, we discuss the results of water-
390 wet and altered-wet (mixed-wet) sandstone samples with oil-water pair of
391 fluids.

392 3.1. Numerical validation

393 In this section, we provide numerical experiments for the validation of
394 the developed workflow as well as compare the accuracy with the geometri-
395 cal analysis contact angle measurement. For this purpose, two different cases
396 were generated using two-phase colour LB simulations with surface affinities
397 of $\phi_s = 0.5, 1$, where each case consisted of the same set of initial ganglia in
398 a porous medium of 100^3 voxels as shown in Figure 3(a). The domain was
399 then relaxed by LBM under no-flow conditions for the two surface affinities.
400 To mimic the added noise and uncertainty of the imaging and segmentation
401 process of experimental imaging of porous media, a median filter was applied
402 to the simulated domains that introduce differences compared to the higher
403 resolution LBM simulation results. Figure 3 depicts the initial domain with

404 patches of the non-wetting phase, as well as the obtained numerical simula-
 405 tion results for $\phi_s = 0.5$ and $\phi_s = 1$. Finally, the wettability of the results was
 406 characterized by the developed workflow as well as the geometrical method to
 407 quantitatively evaluate the accuracy of each method based on initial input.

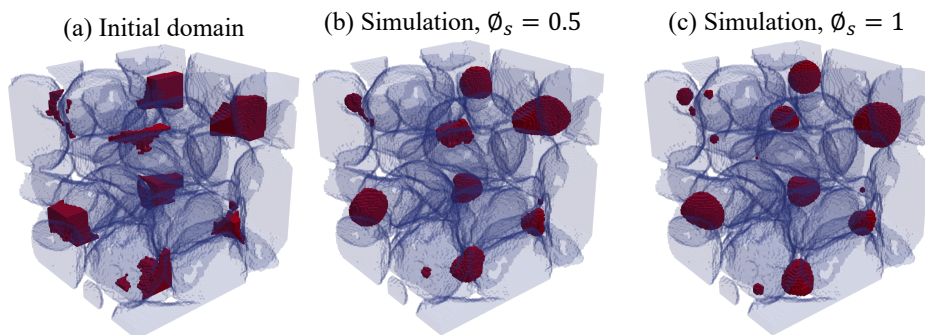


Figure 3: Numerical simulation of the two-phase static condition using the LBM. (a) Initial condition, a porous media of 100^3 voxel size, saturated with fluid 1 with patches of fluid 2 (non-wetting fluid, red colour). (b) final result for relaxation of the domain using $\phi_s = 0.5$, and (c) final results using $\phi_s = 1$.

408 Figure 4 shows the results of both the developed workflow and the geometrical
 409 analysis for these two cases. As it can be seen the geometrically
 410 obtained contact angles have a wider range. This is as expected due to the
 411 central limit theorem since the proposed method yields one single value for a
 412 ganglion, which corresponds to a value for each grid cell along the three-phase
 413 contact line for the geometrical method. The geometrical method yields a
 414 fair estimate for the $\phi_s = 0.5$ model but struggles to capture the affinity
 415 values for the $\phi_s = 1$ case. The results of the developed workflow are more
 416 representative of the high-affinity case of $\phi_s = 1$.

417 For the case with the surface affinity of 1 (i.e., $\phi_s = 1$, strongly non-
 418 wetting case) the geometrical analysis contact angles results range mostly
 419 from 0 to 60 degrees with an average of 37.5 degrees. For the case with
 420 $\phi_s = 0.5$ the obtained geometrical contact angles range from 37 to 90 degrees
 421 with an average of 60 degrees. On the other hand, the developed workflow
 422 gave a range of affinities for each case with averages of $\phi_s = 0.95$ and $\phi_s =$
 423 0.70 for these two cases, corresponding to contact angles of 18 and 46 degrees,
 424 respectively. The presented results show that the wettability characterization
 425 results can be more representative using the developed workflow than the
 426 geometrical analysis method for the presented cases.

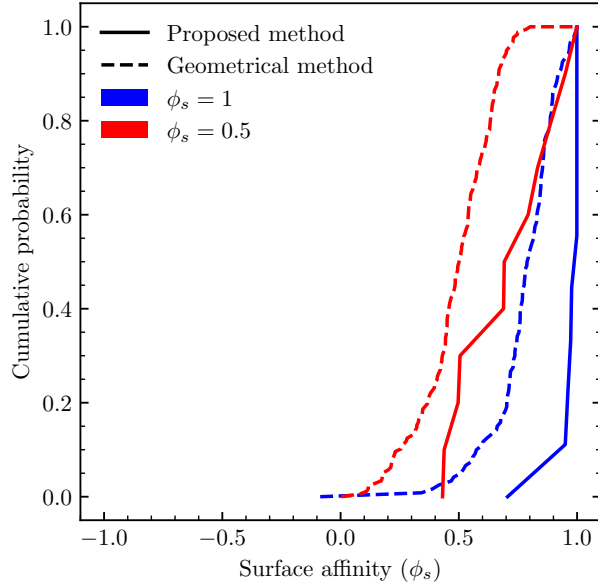


Figure 4: Obtained affinity value distributions for both the proposed method and the traditional geometrical analysis for the two simulated domains using lattice-Boltzmann simulations ($\phi_s = 0.5$ and $\phi_s = 1$, corresponding to contact angles of 60 and 0 degrees).

427 *3.2. Gas-water Bentheimer sample*

428 The first experimental dataset that we used to examine the proposed
 429 scheme for wettability characterization of two-phase fluid distribution was
 430 the Bentheimer sample saturated with water and gas fluid pair under no-
 431 flow boundary conditions in ambient pressure and temperature. The exist-
 432 ence of water and gas together is beneficial as it is well-known that gas is
 433 usually the strongly non-wetting phase in such conditions. Figure 5(a) shows
 434 the obtained distribution of surface affinity parameter (ϕ_s) for this sample.
 435 For the proposed wettability characterization scheme in this research, each
 436 ganglion gives a data point as the local simulation was performed on each
 437 disconnected gas ganglion and the surface affinity parameter was optimized
 438 based on how well the lattice-Boltzmann model described the geometry of
 439 that specific ganglion. The obtained distribution is significantly shifted to-
 440 ward the liquid-wetting end of the spectrum with most of the ganglia being
 441 well described with $\phi_s=1$ which is the strongly non-wetting condition for the
 442 gas phase.

443 The geometrical algorithm gives multiple calculations along the three-

444 phase contact line, to examine such effect in the proposed scheme surface
 445 affinity distribution we also plot the distribution of surface affinity weighted
 446 by the ganglion size in Figure 5(b). As can be seen, compared to Figure 5(b)
 447 weighting the surface affinity by the size of the ganglion does not materially
 448 change the obtained distribution for the surface affinity. So, in the remaining
 449 part of the paper we stick to unweighted distributions for the proposed wet-
 450 tability characterization scheme to highlight the differences with geometrical
 451 analysis.

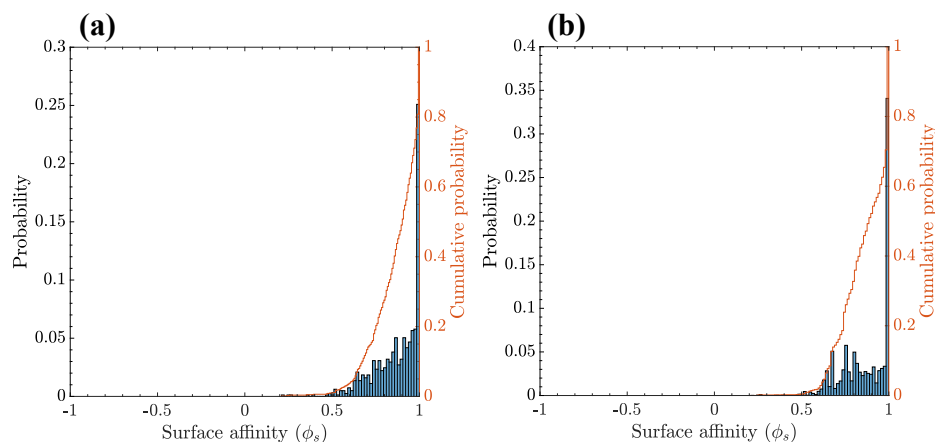


Figure 5: The distribution of surface affinity parameter (ϕ_s) for different ganglia in gas-water Bentheimer sample. (a) Unweighted distribution, (b) weighted distribution by the size of ganglia.

452 To be able to directly compare the proposed method for wettability char-
 453 acterization using LBM with the geometrical analysis we converted the ob-
 454 tained ϕ_s distribution to the contact angle, shown in Figure 6. Both distri-
 455 butions show the water-wetting behaviour of the sample as expected while
 456 the proposed scheme characterizes the wettability as strongly liquid wetting,
 457 with most of the contact angle distribution being in the range of 0-45 de-
 458 grees. On the other hand, the geometrical analysis contact angle distribution
 459 provides a wider span with the peak of the distribution at around 60 degrees,
 460 which is unexpected for the water and gas in the porous media in ambient
 461 conditions. As the geometrical analysis is a local numerical approach that
 462 calculates the contact angle based on the orthogonal vector of solid-fluid and
 463 fluid-fluid surfaces it is more prone to local artefacts, roughness, and abnor-
 464 malities in imaging or natural rock surface. Moreover, Figure 7 shows one

465 of the observed ganglia in the porous media together with the simulated re-
 466 gion for different affinities. In this specific case, the optimization algorithm
 467 converged to $\phi_s=0.724$ which provides the best description of the cluster.

468 Figure 8 depicts the obtained wettability distribution map (i.e., ϕ_s) for
 469 all solid-fluid voxels in the gas-water Bentheimer sample. The voxels are not
 470 shown for the top half of the sample for the sake of better visualization of
 471 the spatial distribution of the trapped gas ganglia in the sample, which gives
 472 the wetting properties of the sample.

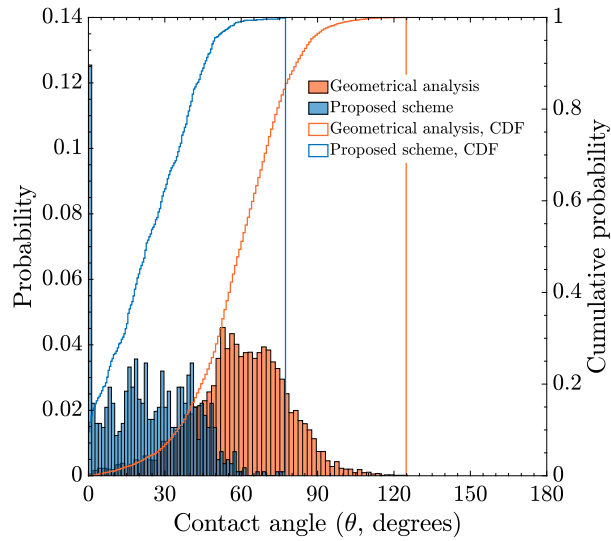


Figure 6: Obtained contact angle distributions for geometrical analysis and the proposed scheme for the gas-water Bentheimer sample.

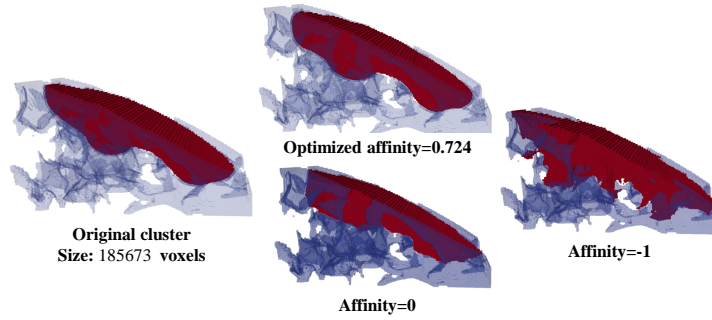


Figure 7: One of the trapped gas ganglions in the gas-water Bentheimer sample and simulation cases for different surface affinity values using local LB simulations. For this cluster, the final optimized value of surface affinity is $\phi_s=0.724$ which describes the cluster the best.

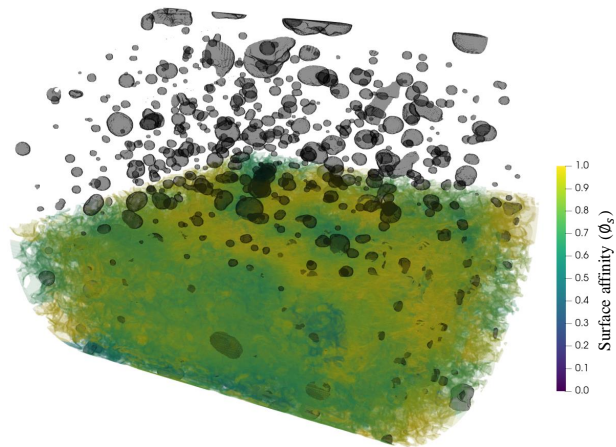


Figure 8: Surface affinity parameter (ϕ_s) distribution map for all solid-fluid voxels in the gas-water Bentheimer sample.

473 *3.3. Oil-water Bentheimer samples (unaltered and altered)*

474 In this section, we describe the results from running both our proposed
475 wettability characterization workflow and from geometrical analysis on two
476 Bentheimer samples, altered and unaltered, saturated with oil and brine.
477 We expect the altered wettability sample to be mixed wet and the unaltered
478 sample to be water wet [26, 53]. Figure 9 shows the obtained surface affinity
479 distributions for altered and unaltered samples using local LB simulation on
480 trapped ganglia. The wettability distribution of samples shows an obvious
481 distinction between the two samples, with a shift toward oil-wetness for the
482 altered sample, as well as a wider range showing the mixed-wet behaviour.
483 Figure 10 depicts the obtained contact angle distributions for the altered
484 and unaltered samples using geometrical analysis. The difference between
485 the two samples is lower compared to Figure 9. Moreover, both distribution
486 profiles show a wide range of obtained contact angles with a slight difference
487 between the peak contact angle values.

488 Figure 11 compares all obtained contact angle distributions for altered
489 and unaltered Bentheimer samples under oil-water fluid pair. For the un-
490 altered sample, in which we expect a water-wet behaviour from clean Ben-
491 theimer sandstone the results of the proposed wettability characterization
492 scheme show a stronger water-wetting behaviour compared to geometrical
493 analysis. The same difference applies to the altered sample. The local LBM
494 simulations show a peak contact angle distribution of around 70 degrees
495 while the geometrical analysis distribution peaks around 90 degrees, showing
496 a neutral wetting behaviour. In the presented results, the contact angle dis-
497 tribution of the unaltered sample obtained by geometrical analysis matches
498 the obtained contact angle distribution for the altered sample using local
499 LBM simulations. Based on the provided results we believe that the pro-
500 posed scheme provides more accurate and representative results compared to
501 geometrical analysis, which might be due to smoothing the interface along
502 the contact line and calculation of the normal vector locally.

503 Finally, Figure 12 shows the obtained surface affinity parameter (ϕ_s) map
504 for all solid-fluid voxels for half of the altered and unaltered Bentheimer
505 samples. The solid-fluid voxels are not shown for the top half of the sample
506 for better visualization and to be able to visualize the wetting distribution
507 inside the sample as well as the spatial distribution of three-phase contact
508 lines. There is a significant difference between the affinity maps with a shift
509 toward oil-wetness for the mixed wet sample, which is expected.

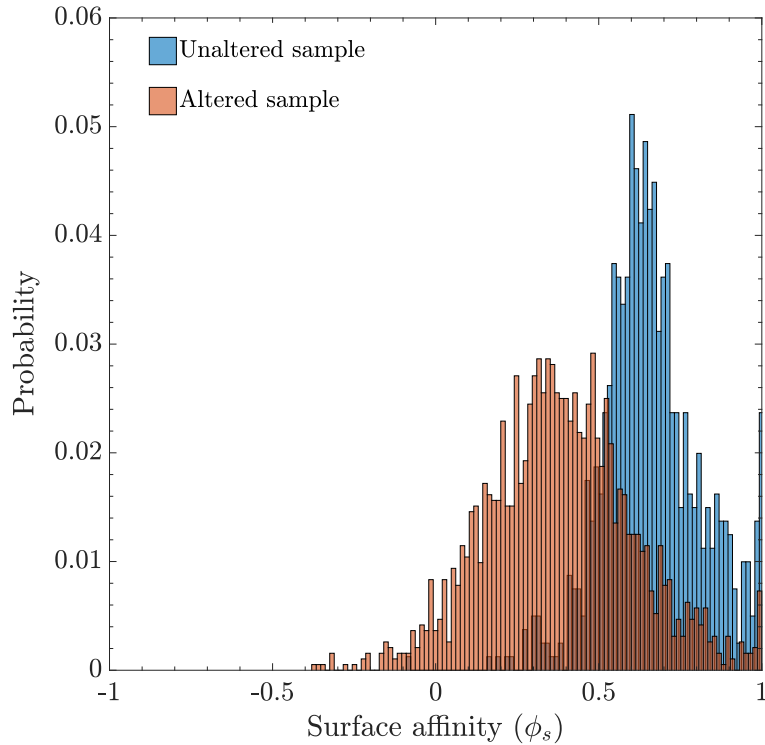


Figure 9: The distribution of obtained surface affinity parameter (ϕ_s) for altered and unaltered Bentheimer sandstone samples saturated with oil and brine phases.

510 4. Summary and Conclusions

511 This study presented a new approach to spatial characterisation of wetta-
 512 bility in porous media. The presented workflow is computationally efficient,
 513 as it uses local lattice-Boltzmann (LB) simulations, and upscales well as it
 514 is performed on the segmented structure of rock and fluids in porous media.
 515 The developed workflow works on trapped ganglia of the non-wetting phase
 516 individually and conducts local lattice-Boltzmann simulations to obtain the
 517 most appropriate wetting parameter which describes that specific ganglion
 518 as close as possible to the observed geometry in the porous medium. We
 519 used a colour-gradient lattice-Boltzmann model to simulate two-phase fluid
 520 distribution in porous media.

521 To determine the local wettability in the porous medium, we isolated each
 522 disconnected ganglion of fluids and ran one LB simulation for each ganglion

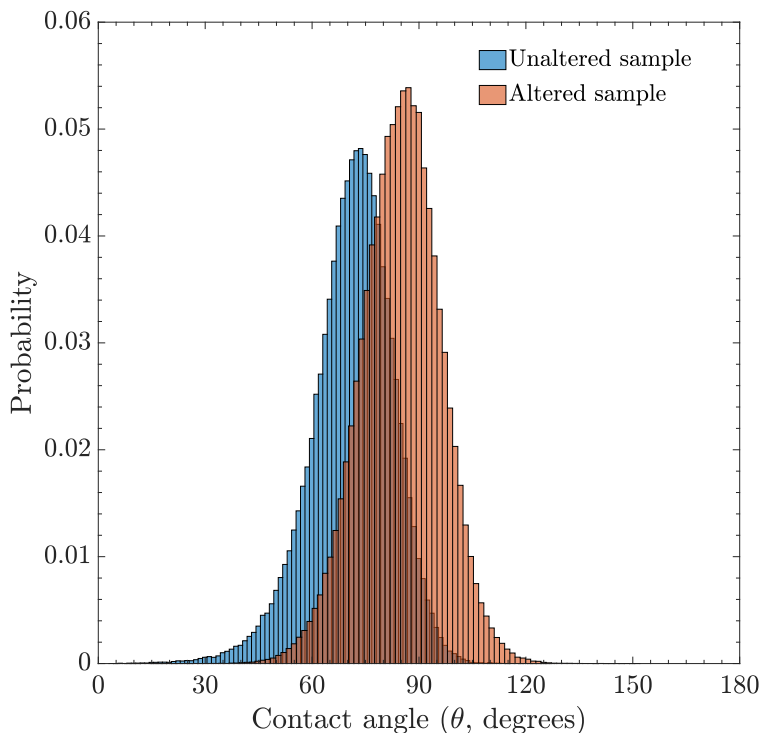


Figure 10: The contact angle distribution for altered and unaltered Bentheimer sandstone samples saturated with oil and brine phases. The contact angle distribution is obtained by geometrical contact angle determination.

523 in the porous medium. We optimized the surface affinity parameter (through
 524 which the wettability is imposed in the solver) for each ganglion with the aim
 525 of the best ganglion shape replicating the LB results with the imaged ganglion
 526 geometry. Using this approach, we obtained a surface affinity parameter for
 527 each disconnected ganglion and assigned it to the three-phase (fluid-fluid-
 528 solid) contact line. Computing sequentially on a single ganglion makes the
 529 method computationally advantageous and RAM-efficient, moreover it can
 530 be parallelized very efficiently to run the algorithm on multiple ganglia at
 531 the same time. After obtaining the surface affinity for all three-phase contact
 532 line voxels, the surface affinity of all solid-fluid voxels of the porous medium
 533 was estimated using 3D spatial interpolation.

534 First, the results of the developed workflow were compared with a geo-
 535 metrical contact angle (CA) determination scheme proposed by Khanamiri

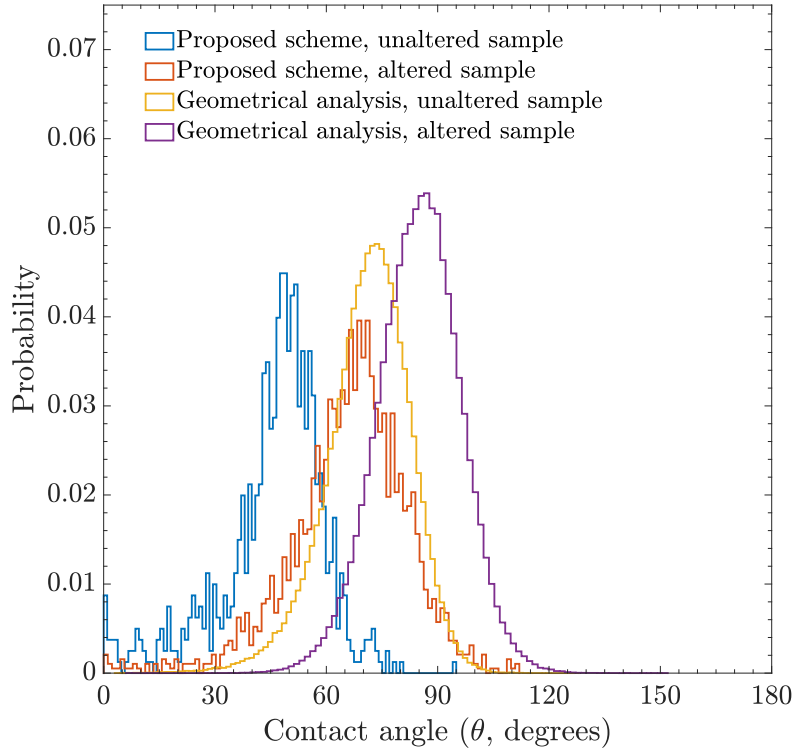


Figure 11: Contact angle distributions obtained from local LB simulations and geometrical contact angle determination schemes for altered and unaltered Bentheimer samples saturated with oil and brine phases.

536 et al. [32] on two synthetic cases, generated by two-phase flow simulation in
 537 a porous medium. Our analysis revealed that the results obtained from the
 538 developed workflow accurately represent both strongly wetting and interme-
 539 diate wetting states, whereas the geometrical analysis struggles to capture
 540 the extreme wetting case and yields a broader spectrum of contact angles,
 541 leading to increased uncertainty. Then, three sets of experimental data were
 542 used to quantitatively compare the results of the proposed scheme with geo-
 543 metrical analysis. The first set of experimental data was from a Bentheimer
 544 sandstone saturated with air and water in ambient conditions. Both the pro-
 545 posed scheme and the geometrical analysis showed water wetness, with the
 546 former showing more water-wetting behaviour and a narrower span of con-
 547 tact angle distribution. It is known that natural surfaces show a strong liquid
 548 wetting behaviour in the presence of gas [55], from which we conclude that

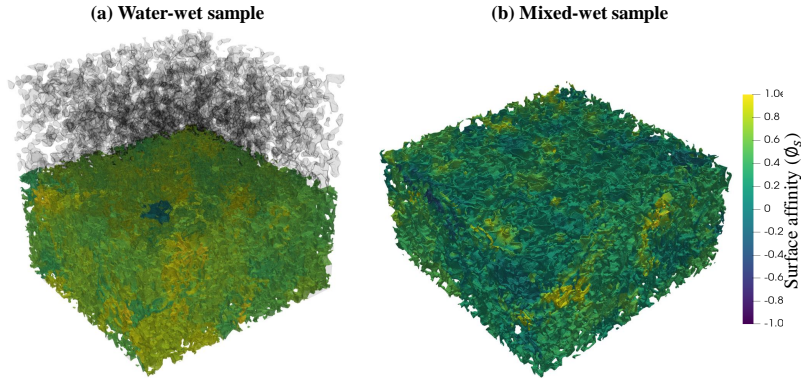


Figure 12: Surface affinity parameter (ϕ_s) distribution map for all solid-fluid voxels in the (a) unaltered and (b) altered Bentheimer samples.

549 the local LBM simulations provide a more realistic wettability description
 550 for this sample.

551 The second and third sets of experimental data were oil-water saturated
 552 unaltered and altered (aged with crude oil) Bentheimer samples under steady-
 553 state flow ($f_w = 0.5$). We expect the unaltered sample to show a water-wet
 554 behaviour while the altered sample is expected to be mixed-wet. The dif-
 555 ference in contact angle distribution for these two samples was small when
 556 the geometrical contact angle determination was applied, while the proposed
 557 scheme showed a larger difference between the two datasets. The proposed
 558 scheme showed a water-wet behaviour for the unaltered sample while the
 559 geometrical contact angle determination showed contact angle distributions
 560 peaking around a neutral wetting state ($75 < CA < 105$). Based on the pro-
 561 vided results we believe that the proposed scheme based on local LBM sim-
 562 ulations provides more accurate and representative results compared to the
 563 traditional geometrical analysis.

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