

Cite this: DOI: 00.0000/xxxxxxxxxx

## Electron doping as a handle to increase the Curie temperature in ferrimagnetic $\text{Mn}_3\text{Si}_2\text{X}_6$ (X=Se, Te)

Lei Qiao,<sup>a,b</sup> Paolo Barone,<sup>c</sup> Baishun Yang,<sup>b</sup> Phil D.C. King,<sup>d</sup> Wei Ren,<sup>\*a</sup> and Silvia Picozzi<sup>\*b</sup>

Received Date

Accepted Date

DOI: 00.0000/xxxxxxxxxx

By analysing the results of ab-initio simulations performed for  $\text{Mn}_3\text{Si}_2\text{X}_6$  (X=Se, Te), we first discuss the analogies and the differences in electronic and magnetic properties arising from the anion substitution, in terms of size, electronegativity, band widths of  $p$  electrons and spin-orbit coupling strengths. For example, through mean-field theory and simulations based on density functional theory, we demonstrate that magnetic frustration, known to be present in  $\text{Mn}_3\text{Si}_2\text{Te}_6$ , also exists in  $\text{Mn}_3\text{Si}_2\text{Se}_6$  and leading to a ferrimagnetic ground state. Building on these results, we propose a strategy, electronic doping, to reduce the frustration and thus to increase the Curie temperature ( $T_C$ ). To this end, we first study the effect of electronic doping on the electronic structure and magnetic properties and discuss the differences in the two compounds, along with their causes. Secondly, we perform Monte Carlo simulations, considering from first to fifth nearest-neighbor magnetic interactions and single-ion anisotropy, and show that electron doping efficiently raises the  $T_C$ .

### 1 Introduction

$\text{Mn}_3\text{Si}_2\text{Se}_6$  (MSS) and  $\text{Mn}_3\text{Si}_2\text{Te}_6$  (MST)<sup>1,2</sup> represent two sister compounds, halfway between two-dimensional (2D) and three-dimensional (3D) materials; in fact, they are formed by *i*) a 2D  $\text{Mn}_2\text{Si}_2\text{X}_6$  (X = Se, Te) layer, in all respects equivalent to one of the prototypical 2D magnets, *i.e.*  $\text{Cr}_2\text{Ge}_2\text{Te}_6$ <sup>3</sup>; *ii*) a Mn intercalating layer, forming a triangular lattice, placed in the so called "van der Waals gap" of bulk 2D-magnets and granting a 3D behaviour to this class of materials. MSS and MST have been studied by numerous researchers in recent years<sup>1,4-7</sup>. MST features polaronic transport and very low values of thermal conductivity, which is further suppressed in magnetic field<sup>6</sup>. The resistivity decreases by 7 orders of magnitude when the MST is placed in a magnetic field greater than 9 T, leading to an insulator-metal transition at 130 K<sup>8</sup>. Zhang *et al* also proposed to regulate the metal-insulator transition by spin orientation and Se instead of Te<sup>7</sup>. Sala *et al* systematic studied the magnetic frustration state in the MST by combining experiments and calculations<sup>9</sup>. Wang

*et al* demonstrated that at a pressure of 10 GPa, the  $T_C$  of MST can be increased to 210 K while maintaining the ferrimagnetic ground state<sup>10</sup>. MSS and MST are isomorphic compounds, hence they are expected to feature similar properties. Among them, we recall the ground state of both MSS and MST being ferrimagnetic with 5  $\mu_B$  magnetic moments and opposite direction of Mn1 and Mn2 sublattices, resulting in a net magnetic moment of 1.6  $\mu_B/\text{Mn}$ <sup>1,2</sup>. The antiferromagnetic (AFM) coupling between Mn atoms and the competition between different nearest neighbors leading to the interesting ferrimagnetic configuration in MST<sup>2,11</sup> is also expected to play a role in MSS. There are also some differences between MSS and MST, mostly induced by the different ionicity and size of the anion, but also to the magnitude of spin orbit coupling (SOC) in the telluride and in the selenide. The current doubt lies in the magnetization direction of MSS and MST. Multiple experiments and calculations believe that the two materials have in-plane magnetization<sup>1,6-8,11-13</sup>. However, some researchers have proposed that the magnetization directions of MSS and MST have angles of around 35° and 10° with the *ab* plane respectively<sup>2,14</sup>. Our calculation results more support the former statement.

A magnet ideal for applications should possess several properties, such as strong magnetocrystalline anisotropy, quantum anomalous Hall effect, perfect crystalline order, high  $T_C$  etc<sup>15</sup>. In particular, great efforts were directed towards increasing the ordering temperature of magnetic materials whenever their  $T_C$  is far below room-temperature. Different routes were followed, among which the most common are represented by external fields<sup>16,17</sup>, intercalation<sup>18</sup>, strain engineering<sup>19</sup>, molecule absorption<sup>20-23</sup> and formation of heterostructures with other materials<sup>24-27</sup>. In

<sup>a</sup> Physics Department, International Center of Quantum and Molecular Structures, Materials Genome Institute, State Key Laboratory of Advanced Special Steel, Shanghai Key Laboratory of High Temperature Superconductors, Shanghai University, Shanghai 200444, China; E-mail: renwei@shu.edu.cn

<sup>b</sup> Consiglio Nazionale delle Ricerche (CNR-SPIN), Unità di Ricerca presso Terzi c/o Università "G. D'Annunzio", 66100 Chieti, Italy; E-mail: silvia.picozzi@spin.cnr.it

<sup>c</sup> Consiglio Nazionale delle Ricerche (CNR-SPIN), Area della Ricerca di Tor Vergata, Via del Fosso del Cavaliere 100, I-00133 Rome, Italy

<sup>d</sup> SUPA, School of Physics and Astronomy, University of St Andrews, St Andrews KY16 9SS, UK

† Electronic Supplementary Information (ESI) available: [details of any supplementary information available should be included here]. See DOI: 00.0000/00000000.

this respect, we note that the  $T_C$  of MSS is 67 K, which is slightly lower than 78 K of MST<sup>2,10,11</sup>. It is therefore obvious that such low  $T_C$ s limit the range of applications of these materials and increasing the  $T_C$  appears vital for their use in spintronic devices.

In our previous work<sup>28</sup>, we focused on magnetic properties in MST, by means of a joint theoretical and experimental study. In addition to pointing to a strong covalency between Mn and Te states, we highlighted the mechanism by which the magnetic frustration in MST eventually leads to a ferrimagnetic ground state, and demonstrate a crucial role of both exchange interactions extending beyond nearest-neighbours and of anti-symmetric exchange in dictating its  $T_C$ .

In this work, we systematically investigate the electronic structure and the magnetic properties of MSS and MST, by discussing the differences between the two materials, mostly in terms of *i*) different ionicity and size of the anion; *ii*) SOC strengths. Importantly, we put forward electronic doping as an efficient tool to largely increase the  $T_C$  of MSS and MST; by decreasing the exchange frustration, our results show that electron doping stabilizes ferromagnetism in the Mn1 plane. The Mean Field (MF) theory and the Monte-Carlo (MC) simulations - based on a spin-Hamiltonian which also takes into account the single-ion anisotropy, with all parameters evaluated from first-principles - show that electron doping can effectively increase the  $T_C$ . Based on MC simulations, up to 250 K and 140 K are reached for MST and MSS, respectively, for a doping of 0.8 electrons/unit-cell.

## 2 Methods:

**DFT.** Density functional theory (DFT) calculations were performed within the generalized gradient approximation (GGA)<sup>29</sup> in the form proposed by Perdew-Burke-Ernzerhof (PBE), as implemented in the Vienna ab Initio Simulation Package (VASP)<sup>30</sup>. An additional on-site interaction was considered within a DFT+U approach in the form introduced by Dudarev *et al*<sup>31</sup>. From our previous study<sup>28</sup> focused on an accurate comparison between DFT and spectroscopy experiments,  $U = 2$  eV appeared to be the best choice in the description of electronic and magnetic properties, that is, we only use different  $U$  values for section 3. The projector augmented wave (PAW) pseudopotentials<sup>32,33</sup> were used, considering as valence states  $3p^63d^54s^2$  for Mn,  $3s^23p^2$  for Si,  $4s^24p^4$  for Se and  $5s^25p^4$  for Te. The plane-wave energy cutoff was fixed to 600 eV. The convergence criterion for total energy differences during the self-consistent cycle was set to less than  $10^{-6}$  eV and the threshold on the maximum force on each ion was less than 0.005 eV/Å. It should be noted that the lattice constants from experiments are used for MST, while the DFT-optimized lattice parameters of MSS (for which, to the best of our knowledge, no experimental data are available) are adopted:  $a = 6.544$  Å,  $c = 13.724$  Å for MSS and  $a = 7.028$  Å,  $c = 14.255$  Å for MST. We note that for MST the difference between DFT-predicted ( $a = 7.079$  Å,  $c = 14.332$  Å) and experimental values is smaller than 1% (*i.e.* a standard accuracy for DFT predictions) and we therefore expect a careful description of structural properties also for MSS. The magnetic anisotropic energy (MAE) was calculated by taking into consideration SOC. We chose the  $6 \times 6 \times 4$   $\Gamma$ -centered k-grid sampling<sup>34</sup> for the unit-cell. The four-state method<sup>35</sup>

was used for the evaluation of the exchange-coupling constants and single-ion anisotropy. In that case, for the  $2 \times 2 \times 1$  supercell used in the calculation of single ion anisotropy (SIA) and for the  $3 \times 3 \times 1$  supercell used in the calculation of exchange constants, we chose a  $2 \times 2 \times 2$  k-grid sampling. The electronic doping is modeled by artificially changing the number of electrons in the unit-cells and considering an homogeneous background-charge.

**Monte-Carlo.** The magnetic properties can be modeled via a spin model for classical spins, including a pair-wise spin-spin scalar Heisenberg coupling and a SIA term:

$$H = -\frac{1}{2} \sum_{ij} J_{ij} S_i \cdot S_j + \sum_i A_i \cdot S_i^2 \quad (1)$$

where  $J_{ij}$  is the magnetic exchange parameter between  $S_i$  and  $S_j$ . We used the VAMPIRE<sup>36</sup> code to perform MC simulations, assuming isotropic exchange constants up to the fifth nearest neighbor; anisotropic effects are modeled via an effective SIA term with coupling constant  $A_i$  extracted from MAE. The standard Metropolis algorithm<sup>37</sup> has been used for MC simulations with  $5 \times 10^4$  MC steps for equilibrium and  $1 \times 10^5$  MC steps for averaging. Since the VAMPIRE software can only simulate cuboid cells, we used  $1 \times \sqrt{3} \times 1$  cells, which contains eight Mn1 and four Mn2 sites. We performed calculations for a  $16 \times 8 \times 6$  supercell with  $N_s = 10368$  spins. The transition temperature can be estimated from the peaks that appear in the temperature evolution of specific heat, evaluated as:

$$C_v = \frac{k_B \beta^2}{N_s} [\langle E^2 \rangle - \langle E \rangle^2]$$

where  $E$  is the energy calculated using model Eq.(1),  $k_B$  is the Boltzmann constant and  $\beta = 1/k_B T$ , while  $\langle \dots \rangle$  indicates statistical averages.

## 3 Results and discussion for pristine MSS and MST compounds

### 3.1 Band structures

Before starting the discussion of the electronic properties, let us recall that the unit cell of MSS and MST, as shown in Fig.1(a), in the  $P\bar{3}1c$  phase (space group No. 163) comprises 6 Mn atoms, each carrying a local magnetic moment: there are two symmetry-inequivalent manganese atoms belonging to  $4f$  and  $2c$  Wyckoff positions. The Se (Te) atoms occupy  $12i$  Wyckoff positions in MSS (MST). Each Mn atom is at the center of a tilted octahedron composed of Se/Te, with Mn1-centered octahedra sharing edges with octahedra around Mn1 and sharing faces with the octahedra around Mn2.

Figure1(b) and (c) show the band structures and density of states (DOS) of MSS and MST, which are characterized as indirect bandgap semiconductors with bandgap of 1.19 and 0.45 eV, respectively. These values are obtained within DFT, so they are likely underestimated with respect to experimental values, due to the well known failure of DFT in treating excited states<sup>38</sup>. As expected from the larger electronegativity of Se with respect to Te, MSS shows a stronger insulating behaviour with respect to MST. To further understand the electronic structure, we show the band structures in Fig.1(b) and (c) taking into account their or-

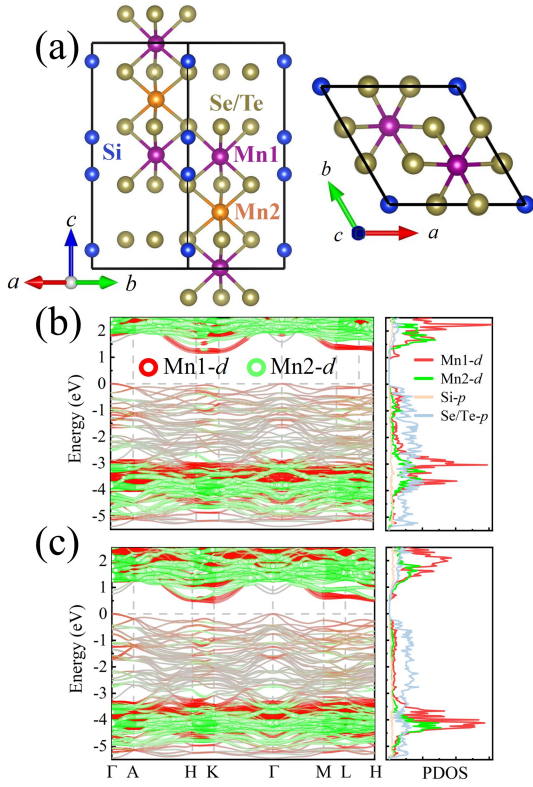


Fig. 1 (a) MSS and MST structures. Side (top) view on the left (right) panel. Legend for atomic colors: Mn1- purple, Mn2- orange, Si- blue, Se or Te- yellow. Band structure and corresponding projected density of states of (b) MSS and (c) MST, respectively (calculated with effective  $U = 2$  eV).

bit character; in particular, Mn1- and Mn2-derived states are represented by red and green, respectively. Se/Te gives rise to most of the states near the valence band maximum, while the  $d$  electrons of Mn mostly occupy somewhat deeper energy levels, *i.e.*, the energy range of -5 to -3 eV. Interestingly, we find that the conduction band minimum (CBM) is mostly contributed by Mn1, with Si and Se/Te also contributing partly, and only Mn2 not contributing at all. The difference in electronic structure between Mn1 and Mn2 suggests that, when doping with electrons, these latter are more likely to move in the plane containing Mn1 octahedra, rather than in the plane of Mn2 octahedra. This behaviour will be particularly important when discussing below the dependence of the exchange parameters upon doping (see Section 4).

As noted previously for MST<sup>11</sup>, the band gap can be largely modulated by the magnetic moments orientation, as shown in Fig. S1. Due to the nodal-line degeneracy protected by the underlying symmetry, the valence band at the  $\Gamma$  point splits into two bands when MSX have out-of-plane magnetization. This effect is present in both MSS and MST, with a splitting size of 0.16 and 0.38 eV for MSS and MST, respectively. The smaller magnitude in MSS suggests the SOC strength to be the key parameter in the band gap modulation, consistently with previous reports<sup>11</sup>.

	method	$U$ (eV)	$J_1$	$J_2$	$J_3$
MSS	$a$	0	-3.71	-2.3	-2.36
	$a$	1	-2.70	-1.60	-1.69
	$a$	2	-2.04	<b>-1.15</b>	-1.20
	$b$	2	-1.92	-0.69	-1.11
MST	$a$	0	-5.72	-1.84	-2.27
	$a$	1	-4.59	-1.36	-1.72
	$a$	2	-3.53	<b>-0.88</b>	-1.20
	$b$	2	-3.19	-0.38	-1.01
			$J_4$	$J_5$	$J_2 + J_5$
MSS	$b$	2	-0.08	-0.28	<b>-0.97</b>
MST			-0.08	-0.26	<b>-0.64</b>

$a$  = energy-mapping method,  $b$  = four-state method.

Table 1 Exchange parameters (in meV) of MSS and MST, calculated with different effective  $U$  values and different methods (energy-mapping vs four-state). Note that  $a$  and  $b$  represent the energy-mapping method and four-state method for solving the exchange interaction. Negative values denote an antiferromagnetic interaction.

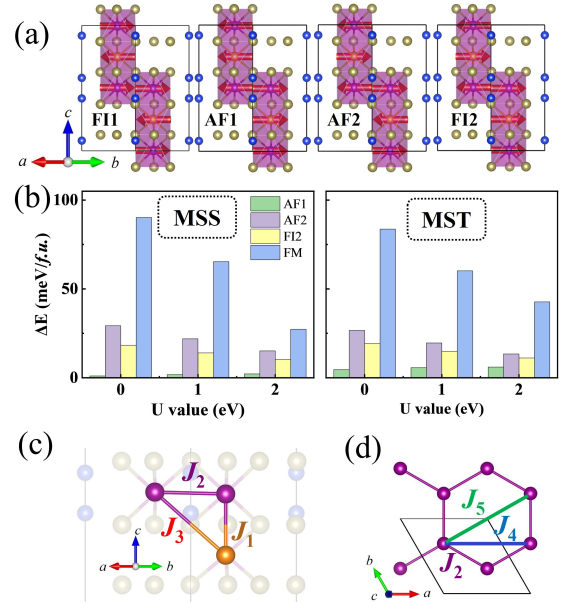


Fig. 2 Schematic representation of the different magnetic configurations used in energy-mapping method (FM not shown here). (b) Energy difference between different magnetic configurations of MSS (left panel) and MST (right panel), respectively. As the effective  $U$  value increases, the energy difference decreases. (c) and (d) Schematic diagrams of first nearest neighbor to fifth nearest neighbor. Panel (c) shows a side-view, whereas (d) shows a top-view of the Mn1 plane.

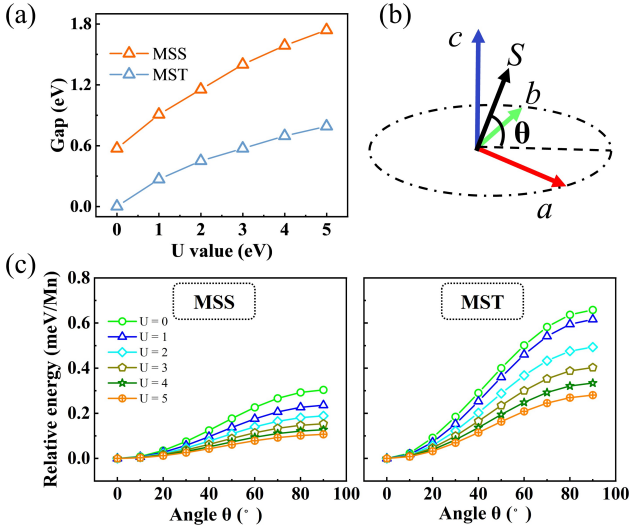


Fig. 3 (a) Dependence of the DFT band gap on the effective  $U$  values. (b) Schematic diagram of magnetic moments direction. An angle  $\theta$  equal to  $0^\circ$  and  $90^\circ$  represent the direction of magnetization lying in the  $xy$  plane and along the  $z$  axis, respectively. (c) Angular dependence of the magnetocrystalline anisotropy energy of MSS (left panel) and MST (right panel). Different color lines are calculated with different effective  $U$  values. It should be noted that during the process of magnetization rotation from in-plane to out-of-plane, the antiparallel configuration between Mn1 and Mn2 magnetic moments is preserved.

### 3.2 Exchange coupling constants

Here, we start discussing the magnetic properties of MSS and MST. We denote as  $S_i$  and  $s_j$  in Eq.(1) for Mn1 and Mn2, respectively (with  $i=1, \dots, 4$  and  $j=1, 2$ ), the classical spins that describe the magnetic moments of the two sets of Mn. Assuming up to third nearest neighbor interactions and an isotropic Heisenberg model for classical spins [i.e., neglecting the anisotropic term in Eq. (1)], the energy per cell reads:

$$E = -\frac{1}{2} \{ 2J_1 \cdot [s_1 \cdot (S_1 + S_3) + s_2 \cdot (S_2 + S_4)] + 6J_2 (S_1 \cdot S_2 + S_3 \cdot S_4) + 6J_3 [s_1 \cdot (S_2 + S_4) + s_2 \cdot (S_1 + S_3)] \} \quad (2)$$

As displayed in Fig.2(c),  $J_1, J_3$  describe interlayer coupling between  $S_1$  and  $s_j$  spins, whereas  $J_2$  is the intralayer coupling between  $S_i$  spins. It should be noted that we consider  $S = s = 5/2$  in this work. By considering five different spin configurations, as shown in Fig.2(a), we obtain the following equations:

$$\begin{aligned} E_{F11} &= +4J_1 - 6J_2 + 12J_3 + E_0 \\ E_{AF1} &= -4J_1 + 6J_2 + 12J_3 + E_0 \\ E_{F12} &= +6J_2 + E_0 \\ E_{AF2} &= +4J_1 + 6J_2 - 12J_3 + E_0 \\ E_{FM} &= -4J_1 - 6J_2 - 12J_3 + E_0 \end{aligned} \quad (3)$$

The best-fit solutions of exchange parameters from Eq.3 are shown in Table 1. For both MSS and MST, all the interactions between Mn spins are negative, *i.e.* corresponding to an AFM interaction. As shown in Table 1, the magnetic interaction decreases as the  $U$  value increases. This is consistent with superexchange theory, according to which the exchange interaction goes as  $t^2/U$ , where  $t$  represents a hopping parameter. Within this simple picture, the larger the  $U$  value, the more localized the electrons, the smaller the exchange interaction. The exchange parameters of MST are basically the same as previously reported<sup>1,7,9,11</sup>. The in-plane AFM exchange interactions indicate that the exchange paths between different nearest neighbors are frustrated in the ferrimagnetic ground state. In fact, the  $J_1$  and  $J_3$  drive antiferromagnetic interactions between Mn1 and Mn2 atoms, while  $J_2$  would lead to antiferromagnetic interactions between Mn1 atoms and competing with the desired action of  $J_1$  and  $J_3$ . The final ferrimagnetic ground state however shows that the ‘‘competition is won’’ by interplanar AFM exchange interactions, since both  $J_1$  and  $J_3$  are larger than  $J_2$ .

Heisenberg spin exchange parameters can also be evaluated by using the four-state method<sup>35</sup>. As shown in Table 1, we calculate even longer distance exchange parameters, that is from  $J_1$  to  $J_5$ , as shown in Fig.2(c) and (d). It is clear that for  $J_1$  and  $J_3$  the two methods give similar results. However, for  $J_2$ , the results of the two methods are quite different, with the results of MSS and MST differing by 67% and 130%. Actually, this difference comes from the used methodology itself. In fact, the hexagon plane formed by the Mn1 atoms contains three  $J_2$ , six  $J_4$  and three  $J_5$  interactions for each Mn1, which in turn is shared among three hexagons, and the energy per cell reads:

$$\begin{aligned} E = -\frac{1}{2} \{ & 2J_1 \cdot [s_1 \cdot (S_1 + S_3) + s_2 \cdot (S_2 + S_4)] + \\ & 6J_2 \cdot (S_1 \cdot S_2 + S_3 \cdot S_4) + \\ & 6J_3 \cdot [s_1 \cdot (S_2 + S_4) + s_2 \cdot (S_1 + S_3)] + \\ & 6J_4 \cdot (S_1^2 + S_2^2 + S_3^2 + S_4^2) + \\ & 6J_5 \cdot (S_1 \cdot S_2 + S_3 \cdot S_4) \}. \end{aligned} \quad (4)$$

On the other hand, when describing the hexagon plane formed by the Mn1 atoms by a single  $J_2$ , the latter effectively contains also longer-ranged exchange interactions (not explicitly taken into account). Clearly, the  $J_4$  exchange constant cannot be evaluated from the energy-mapping method, as it gives a constant contribution for any magnetic configuration that can be described within the unit cell. By comparing Eqs. (2) and (4), it is also evident that the  $J_2$  estimated within the energy-mapping method effectively includes contribution from fifth nearest neighbor and should be compared with  $J_2 + J_5$  as obtained from the four-state method.

Based on the above relations and Table 1, we find that the difference between the ‘‘effective’’  $J_2$  (from the mapping-method) and  $J_2 + J_5$  (from the four-state method) is indeed quite close. From the perspective of the exchange constant, for both MSS and MST we also get the same conclusion as our previous work, that

is, the importance of explicitly including  $J_4$  and  $J_5$ , when calculating magnetic properties<sup>28</sup>.

### 3.3 Magnetic anisotropy energy and Band-gap: dependence on Hubbard-U parameter

In addition to the magnetic exchange interactions, other relevant properties of MSS and MST, such as MAE and band gap, are also sensitive to  $U$ . As shown in Fig.3(c), the MAE of MSS and MST decreases significantly as the value of  $U$  increases. The MAE of MSS is smaller than that of MST, as expected from the weaker SOC effect of Se compared to Te. This is similar to what obtained in the comparison of MAE in CrBr<sub>3</sub> and CrI<sub>3</sub><sup>39</sup>. According to our DFT predictions, both MSS and MST prefer in-plane magnetization. Zhang *et al* claims that the system maintains the in-plane magnetization during the substitution of Se for Te up to 30% concentration, which somewhat implies that MSS will also be in-plane magnetization<sup>7</sup>. From this point, our calculations are consistent. However, May *et al* claims the magnetization direction of MSS is rotated by 35° with respect to the  $xy$  plane<sup>2</sup>. The reason for the discrepancy is unclear.

For the band gap, the larger the  $U$  value, the larger the band gap. In closer detail, in the range of  $U$  from 0 to 5 eV, the DFT band gap of MSS changes in the range from 0.61 to 1.74 eV; MST goes from a metal to semiconductor with a band gap of 0.82 eV. The behaviour of the band gap as a function of the Hubbard  $U$  parameter is expected from the situation of a half-filled  $d$  shell, as the one we have here for Mn  $d^5$ . Indeed, as the (empty) CBM is predominantly contributed by Mn1  $d$  states, it shifts up in energy as  $U$  increases, thereby leading to a band-gap opening. The experimental band gap of MST is 1.39 eV<sup>40</sup>, while the band gap of MSS has not yet been experimentally published. Based on the difference in electronegativity between Se and Te element, it can be predicted that the experimental band gap of MSS is larger than that of MST.

### 3.4 Mean-field analysis

Next, we derive the effective Weiss field generated on each magnetic moment in the cell (representative of a given sublattice) by all surrounding atoms to obtain the critical temperature. The Weiss fields are,

$$\begin{aligned}
 B_1 &= J_1 m_1 + 3J_2 M_2 + 3J_3 m_2 \\
 B_2 &= J_1 m_2 + 3J_2 M_1 + 3J_3 m_1 \\
 B_3 &= J_1 m_1 + 3J_2 M_4 + 3J_3 m_2 \\
 B_4 &= J_1 m_2 + 3J_2 M_3 + 3J_3 m_1 \\
 b_1 &= J_1 (M_1 + M_3) + 3J_3 (M_2 + M_4) \\
 b_2 &= J_1 (M_2 + M_4) + 3J_3 (M_1 + M_3)
 \end{aligned} \tag{5}$$

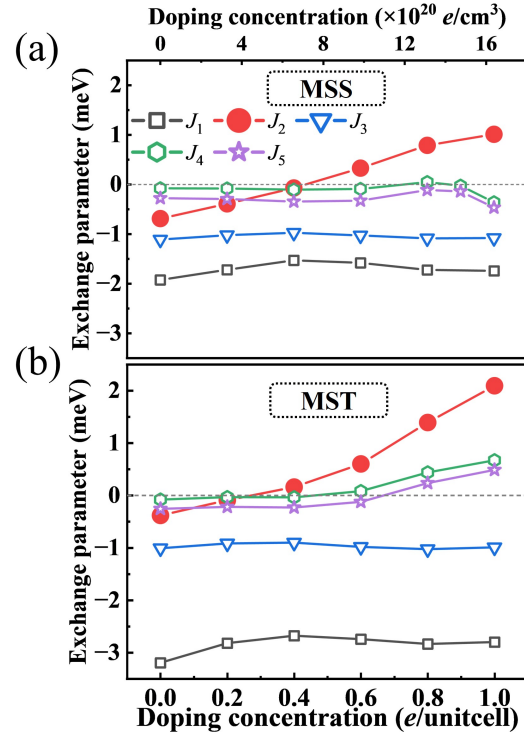


Fig. 4 Electronic doping effect on exchange parameters of (a) MSS and (b) MST. Black, blue, green and purple triangles show  $J_1$ ,  $J_3$ ,  $J_4$  and  $J_5$ , respectively.  $J_2$  is shown in red filled circles, to highlight its importance related to the sign change from AFM to FM upon doping.

where each expectation value  $M_i = \langle S_i \rangle$  and  $m_i = \langle s_i \rangle$  satisfies the self-consistent equation for a classical spin:

$$M_i = \coth(\beta B_i) - \frac{1}{\beta B_i} \approx \frac{\beta B_i}{3} \tag{6}$$

with  $\beta = 1/(k_B T)$ . The right-hand side approximation holds in the limit of  $\beta B_i \rightarrow 0$  when  $T$  is close to  $T_C$ . Combining Eq.(5) and (6), the following equations for the sublattice magnetizations  $M = \sum_{i=1}^4 M_i$  and  $m = \sum_{j=1}^2 m_j$  can be obtained:

$$\begin{aligned}
 M &= \frac{\beta}{3} [3J_2 M + (2J_1 + 6J_3)m] \\
 m &= \frac{\beta}{3} (J_1 + 3J_3)M.
 \end{aligned} \tag{7}$$

The following quadratic equation is readily derived:

$$0 = M[1 - J_2 \beta - \frac{2\beta^2}{9} (J_1 + 3J_3)^2] \tag{8}$$

using which the critical temperature is obtained as:

$$k_B T_C = \frac{|J_2|}{2} \left[ \text{sgn} J_2 \pm \sqrt{1 + 8 \left( \frac{J_1 + 3J_3}{3J_2} \right)^2} \right], \tag{9}$$

where  $\text{sgn} J_2 = J_2/|J_2|$  denotes the sign of the  $J_2$  exchange parameter. The physically meaningful solution of a Curie temperature larger than 0 implies that only the positive result from Eq.(9) is retained. Taking into account interactions for longer distances,

one can still use MF theory to estimate the critical temperature. By repeating the above derivation process, one obtains:

$$k_B T_C = \frac{|J_{\text{intra}}|}{2} \left[ \text{sgn} J_{\text{intra}} \pm \sqrt{1 + 8 \left( \frac{J_{\text{inter}}}{3J_{\text{intra}}} \right)^2} \right], \quad (10)$$

where we defined intra- and inter-sublattice effective exchange parameters  $J_{\text{intra}} = J_2 + 2J_4 + J_5$  and  $J_{\text{inter}} = J_1 + 3J_3$ , respectively.

## 4 Results and discussion on the effects of electronic doping

Our recipe for the increase in  $T_C$  starts from the following consideration. As previously discussed, in both MSS and MST  $J_2$ ,  $J_4$  and  $J_5$  predict an AFM interaction between Mn1, *i.e.* they have an opposite effect with respect to  $J_1$  and  $J_3$ ; in other words, the competing effects driven by  $J_2$ ,  $J_4$  and  $J_5$  reduce the stability of the ferrimagnetic configuration and thus the  $T_C$ . We mentioned that CBM is contributed by Mn1 in band structure part, in the case of light doping, electrons will first move in the Mn1 plane, which is likely to change the exchange parameters of Mn1 plane, that are  $J_2, J_4$  and  $J_5$ , thus change the magnetic frustration state. Here we propose a route to increase the  $T_C$  by weakening the competition between the magnetic interactions and let them cooperate one with the other to improve the stability of the ferrimagnetic state: tuning the intensity of exchange parameters by electron doping. The latter obviously induces a change from semiconducting to metallic behaviour, along with other less-trivial property changes, which we are going to discuss in detail in this section.

### 4.1 Trends of exchange constants vs doping

Figure 4 (a) and (b) shows the response of the magnetic exchange parameters of MSS and MST to electron doping. Remarkably,  $J_2$  is predicted to be extremely sensitive to the doping concentration and changes from AFM to FM as the doping concentration increases, while  $J_1$  and  $J_3$  remain basically unchanged; this situation exactly corresponds to the weakening of the competing antiferromagnetic interactions and to the emergence of ferromagnetic interactions between Mn1 atoms. In other words,  $J_2$  no longer competes with  $J_1$  and  $J_3$ , but rather helps  $J_1$  and  $J_3$  to establish the ferromagnetic order in the Mn1 plane, which is expected to increase the  $T_C$ . It should be noted that  $J_4$  and  $J_5$  - which are also interactions occurring within the Mn1 plane - have a weaker effect on the  $T_C$  with respect to  $J_2$  in the undoped case, and their sensitivity to the doping concentration is far lower than that of  $J_2$ .

Additionally, it is noteworthy that the responses of  $J_4$  and  $J_5$  of MSS and MST diverge when the doping concentration exceeds 0.8  $e/\text{unitcell}$ . Given the similarities between MSS and MST, we think that this discrepancy may stem from the limitations of our approximations (*i.e.*, homogeneous electron doping, validity of Heisenberg-like short-ranged interactions, etc.), which may only be valid for low levels of electron doping. Consequently, all subsequent calculations are considered up to a doping concentration of 0.8  $e/\text{unitcell}$ .

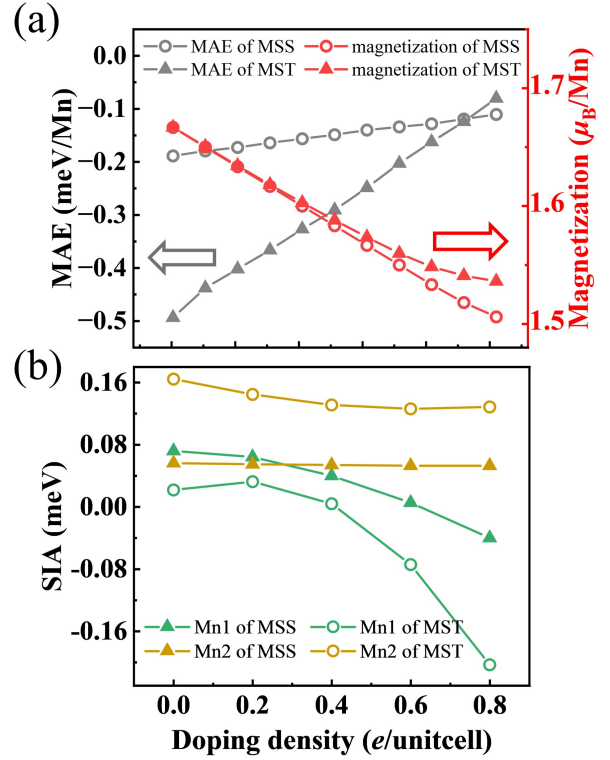


Fig. 5 The electronic doping effect in MSS and MST on (a) MAE (left y-axis) and magnetic moment (right y-axis), (b) single ion anisotropic energy in Mn1 and Mn2.

### 4.2 Trends of total magnetization per unit cell vs doping

As known from the previous discussion of the band structures (cfr Fig. 1), the doped electrons will fill minority Mn1 levels, thereby reducing the moments on Mn1. This will globally result in a linear decrease of the total magnetic moment per unit-cell from the value in the undoped case, *i.e.* 1.667  $\mu_B/\text{unitcell}$ , as shown in Fig.5(a). Notably, the magnetic moments of MSS and MST change linearly with the same trend for small doping levels. However, there is a difference between MSS and MST when the electron concentration becomes larger than 0.3  $e/\text{unitcell}$ , leading to a non-linear behaviour in MST.

One can understand the non-linear behavior from the band structure projected on Mn1 and Mn2 at different doping concentrations in Fig. S2. Electron doping is essentially raising the energy corresponding to the Fermi level. In MSS, the Fermi energy level with different doping concentrations only crosses the bands due to Mn1, *i.e.* the doped electrons reduce the magnetic moments of Mn1 atoms, keeping the Mn2 magnetic moments constant; as a result, there is a linear reduction of the total magnetization per unit cell due to the ferrimagnetic ordering, to be ascribed to the reduction of Mn1 moments. On the other hand, in MST, the Fermi energy level crosses the bands of both Mn1 and Mn2 even for small doping concentrations. The doped electrons are therefore reducing the magnetic moment of both Mn1 and Mn2; due to the ferrimagnetic arrangement and to the complex band structure effects with different filling of Mn1 and Mn2 minority-spin bands, the change in the magnetic moment per unit

cell is no longer linear when increasing the doping concentration.

### 4.3 Trends of magnetic anisotropy

As shown in Fig.5(a), electron doping also strongly weakens the magnetic anisotropy, *i.e.* it greatly reduces the energy difference between in-plane magnetization and out-of-plane magnetization. In principle, given the observed trend, electron doping could even lead to a switch from in-plane to out-of-plane anisotropy for large doping levels, as in  $\text{Mn}_2\text{Ge}_2\text{Te}_6$  with similar phenomena<sup>41</sup>.

We note from Fig.5(a) that the rate of reduction of the MAE energy difference is higher for MST than for MSS. We recall that MAE is composed by two different contributions, one related to the SIA and the other to symmetric anisotropic exchange, the latter being about 15% of the total MAE in undoped MST<sup>28</sup>. However, the discussion here will be limited to the behaviour of SIA upon doping, as the evaluation of anisotropic exchange (*i.e.* of off-diagonal components of the exchange tensor evaluated estimated from large supercells) would be too CPU-time-consuming to perform for all the doping levels. Based on the different site-symmetry of Mn1 and Mn2, we focus on the SIA of both separately, with the results shown in Fig.5(b). There are two interesting points to notice. Firstly, the SIA of Mn1 in MSS is bigger than that of Mn1 in MST, while the opposite is true for Mn2 in MSS and MST. The second and interesting point is that the SIA of Mn1 is much more sensitive to the doping concentration than that of Mn2. Remarkably, the SIA of Mn1 gradually tends toward out-of-plane magnetization upon increasing electron doping, which contributes to explaining the decrease of the magnetic anisotropy energy. Again, the stronger dependence of SIA of Mn1 with respect to Mn2 is likely due to the fact that Mn1 electronic levels are first filled by doping.

### 4.4 Trends of $T_C$

As well known, the MF theory usually overestimates the  $T_C$ . However, one can at least expect the MF theory to predict the right trend linking  $T_C$  to the electron doping concentration. In other words, MF is quantitatively expected to fail, but it is expected to work from the qualitative point of view.

Based on this premise, by using Eq.(10) and the magnetic exchange parameters for different concentrations of electron doping, we obtained the trend of  $T_C$ , as shown in Fig. S3. The  $T_C$  of both MSS and MST first decreases slightly as the doping concentration increases. This is likely because, although the antiferromagnetic interaction given by  $J_2$  becomes smaller in magnitude,  $J_1$  and  $J_3$  likewise become smaller. As the doping concentration exceeds 0.4  $e/\text{unitcell}$ , the  $T_C$  rises significantly. This can be well explained by the change in sign of  $J_2$ : indeed, the parallel configuration within the Mn1-hexagonal sublattice, already favoured by the strong AFM interaction  $J_{\text{inter}}$  between the Mn1 and Mn2 sublattices, is further stabilized by a ferromagnetic  $J_2$  interaction. As both  $J_4$  and  $J_5$  are found to weakly depend on doping, the  $T_C$  enhancement can be understood essentially in the same way when considering the effective  $J_{\text{intra}} = J_2 + 2J_4 + J_5$ . The results of the MF theory show that electron doping can raise the  $T_C$  from 143 K to 200 K for MSS and from 185 K to 312 K for MST. Since

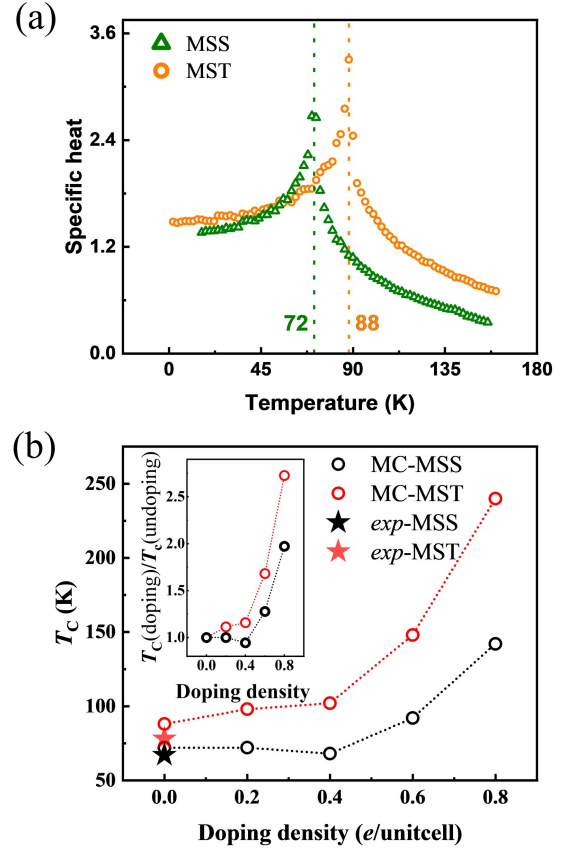


Fig. 6 (a) Specific heat of MST (yellow circles) and MSS (green triangles) from MC simulations. (b)  $T_C$  of MSS (black symbols) and MST (red symbols) as a function of doping concentration. The solid starts show the experimentally measured  $T_C$  of MSS and MST without doping. The inset is the comparison of relative critical temperatures  $T_C(\text{doping})/T_C(\text{undoping})$  as a function of doping density.

MF is known to be not quantitatively reliable, MC simulations were performed for the  $J_1 - J_5$  model to confirm the above trends and conclusions. Before discussing the trends vs electron doping, we show in Fig.6(a) the specific heat of undoped MSS and MST as a function of temperature, the peak in the response function signalling the ordering temperature. It can therefore be easily deduced that the  $T_C$  of undoped MSS and MST are 72 and 88K, which are very close to the experimental values of 67 and 78K, respectively<sup>19,21</sup>. As shown in Fig. S4, where we determined  $T_C$  from the peak of susceptibility, MC simulations vs electron doping predict the same trend as that predicted by MF theory (cfr Fig. S3 and its inset figure): the  $T_C$  increases from 72 to 142K and 88 to 240K for MSS and MST, respectively.

## 5 Conclusions

In this work, we investigated the properties of the magnetically-frustrated ferrimagnetic materials MSS and MST using first principles calculations. In particular, we proposed electron doping as a working recipe to increase the Curie temperature. We quantitatively show that the magnetic frustration can be reduced by electronic doping, which for certain levels even changes the sign of

$J_2$  and can therefore strengthen the ferrimagnetic configuration. The higher  $T_C$  brought about by electronic doping was confirmed by both MF theory and MC simulations. In addition, the electronic and magnetic properties of MSS and MST are discussed, finding that they are qualitatively similar, but not quantitatively the same. In particular, the two compounds do not respond in the same way to electronic doping. We have therefore discussed the mechanisms underlying the MSS and MST different responses in detail and interpreted the differences in terms of larger electronegativity and smaller SOC-strength of the anion and of different energy position of Mn states in MSS compared to MST.

## Conflicts of interest

There are no conflicts to declare.

## Acknowledgements

This work was supported by National Natural Science Foundation of China (Grants No. 12074241, 11929401, 52130204, 12311530675), Science and Technology Commission of Shanghai Municipality (Grants No. 22XD1400900, 22YF1413300, 20501130600, 21JC1402600, and 21JC1402700), High Performance Computing Center, Shanghai University, and Key Research and Development Program of Shanxi Province (No. 2021GXLH-Z-065), High Performance Computing Center, Shanghai Technical Service Center of Science and Engineering Computing. L.Q. acknowledges the support of China Scholarship Council. P.B. and S.P. acknowledge financial support from the Italian Ministry for Research and Education through PRIN-2017 projects ‘Tuning and understanding Quantum phases in 2D materials—Quantum 2D’ (IT-MIUR Grant No. 2017Z8TS5B) and ‘Towards ferroelectricity in Two dimensions- TWEET’ (IT-MIUR Grant No. 2017YCTB59), respectively. P.K. and S.P. acknowledge support from the Royal Society through the International Exchange grant IEC\R2\222041. PK acknowledges support from The Leverhulme Trust via Grant No. RL-2016-006, and the UK Engineering and Physical Sciences Research Council via Grant No. EP/X015556/1.

## Notes and references

- 1 A. F. May, Y. Liu, S. Calder, D. S. Parker, T. Pandey, E. Cakmak, H. Cao, J. Yan and M. A. McGuire, *Physical Review B*, 2017, **95**, 174440.
- 2 A. F. May, H. Cao and S. Calder, *Journal of Magnetism and Magnetic Materials*, 2020, **511**, 166936.
- 3 J. Kanamori, *Journal of Physics and Chemistry of Solids*, 1959, **10**, 87–98.
- 4 L. Martinez, C. Saiz, A. Cosio, R. Olmos, H. Iturriaga, L. Shao and S. Singamaneni, *MRS Advances*, 2019, **4**, 2177–2184.
- 5 L. Martinez, H. Iturriaga, R. Olmos, L. Shao, Y. Liu, T. T. Mai, C. Petrovic, A. R. High Walker and S. Singamaneni, *Applied Physics Letters*, 2020, **116**, 172404.
- 6 Y. Liu, Z. Hu, M. Abeykoon, E. Stavitski, K. Attenkofer, E. D. Bauer, C. Petrovic *et al.*, *Physical Review B*, 2021, **103**, 245122.
- 7 Y. Zhang, L.-F. Lin, A. Moreo and E. Dagotto, *Physical Review B*, 2023, **107**, 054430.
- 8 Y. Ni, H. Zhao, Y. Zhang, B. Hu, I. Kimchi and G. Cao, *Physical Review B*, 2021, **103**, L161105.
- 9 G. Sala, J. Lin, A. Samarakoon, D. Parker, A. May and M. Stone, *Physical Review B*, 2022, **105**, year.
- 10 J. Wang, S. Wang, X. He, Y. Zhou, C. An, M. Zhang, Y. Zhou, Y. Han, X. Chen, J. Zhou *et al.*, *Physical Review B*, 2022, **106**, 045106.
- 11 J. Seo, C. De, H. Ha, J. E. Lee, S. Park, J. Park, Y. Skourski, E. S. Choi, B. Kim, G. Y. Cho *et al.*, *Nature*, 2021, **599**, 576–581.
- 12 Y. Liu, C. Petrovic *et al.*, *Physical Review B*, 2018, **98**, 064423.
- 13 R. Olmos, J. A. Delgado, H. Iturriaga, L. M. Martinez, C. L. Saiz, L. Shao, Y. Liu, C. Petrovic and S. R. Singamaneni, *Journal of Applied Physics*, 2021, **130**, year.
- 14 F. Ye, M. Matsuda, Z. Morgan, T. Sherline, Y. Ni, H. Zhao and G. Cao, *Physical Review B*, 2022, **106**, L180402.
- 15 S. Xu, F. Jia, X. Yu, S. Hu, H. Gao and W. Ren, *Materials Today Physics*, 2022, **27**, 100775.
- 16 E. S. Morell, A. León, R. H. Miwa and P. Vargas, *2D Materials*, 2019, **6**, 025020.
- 17 C. Gong, L. Li, Z. Li, H. Ji, A. Stern, Y. Xia, T. Cao, W. Bao, C. Wang, Y. Wang *et al.*, *Nature*, 2017, **546**, 265–269.
- 18 M. Yang, Q. Li, R. V. Chopdekar, C. Stan, S. Cabrini, J. W. Choi, S. Wang, T. Wang, N. Gao, A. Scholl *et al.*, *Advanced Quantum Technologies*, 2020, **3**, 2070041.
- 19 X.-J. Dong, J.-Y. You, B. Gu and G. Su, *Physical Review Applied*, 2019, **12**, 014020.
- 20 M. Lu, Q. Yao, Q. Li, C. Xiao, C. Huang and E. Kan, *The Journal of Physical Chemistry C*, 2020, **124**, 22143–22149.
- 21 J. He, G. Ding, C. Zhong, S. Li, D. Li and G. Zhang, *Journal of Materials Chemistry C*, 2019, **7**, 5084–5093.
- 22 K. Wang, K. Ren, Y. Cheng, M. Zhang, H. Wang and G. Zhang, *Physical Chemistry Chemical Physics*, 2020, **22**, 22047–22054.
- 23 C. Song, W. Xiao, L. Li, Y. Lu, P. Jiang, C. Li, A. Chen and Z. Zhong, *Physical Review B*, 2019, **99**, 214435.
- 24 N. Liu, S. Zhou and J. Zhao, *Physical Review Materials*, 2020, **4**, 094003.
- 25 L. Zhang, X. Huang, H. Dai, M. Wang, H. Cheng, L. Tong, Z. Li, X. Han, X. Wang, L. Ye *et al.*, *Advanced Materials*, 2020, **32**, 2002032.
- 26 M. Mogi, A. Tsukazaki, Y. Kaneko, R. Yoshimi, K. Takahashi, M. Kawasaki and Y. Tokura, *Apl Materials*, 2018, **6**, 091104.
- 27 L. Zhang, L. Song, H. Dai, J.-H. Yuan, M. Wang, X. Huang, L. Qiao, H. Cheng, X. Wang, W. Ren *et al.*, *Applied Physics Letters*, 2020, **116**, 042402.
- 28 C. Bigi, L. Qiao, C. Liu, P. Barone, M. C. Hatnean, G.-R. Siemann, B. Achinuq, D. A. Mayoh, G. Vinai, V. Polewczyk *et al.*, *arXiv preprint arXiv:2303.00294*, 2023.
- 29 J. P. Perdew, K. Burke and M. Ernzerhof, *Physical Review Letters*, 1996, **77**, 3865.
- 30 G. Kresse and J. Hafner, *Physical Review B*, 1993, **47**, 558.
- 31 V. I. Anisimov, J. Zaanen and O. K. Andersen, *Physical Review B*, 1991, **44**, 943.
- 32 P. E. Blöchl, O. Jepsen and O. K. Andersen, *Physical Review B*, 1994, **49**, 16223.



- 33 G. Kresse and D. Joubert, *Physical Review B*, 1999, **59**, 1758.
- 34 H. J. Monkhorst and J. D. Pack, *Physical Review B*, 1976, **13**, 5188.
- 35 H. Xiang, C. Lee, H.-J. Koo, X. Gong and M.-H. Whangbo, *Dalton Transactions*, 2013, **42**, 823–853.
- 36 R. F. Evans, W. J. Fan, P. Chureemart, T. A. Ostler, M. O. Ellis and R. W. Chantrell, *Journal of Physics: Condensed Matter*, 2014, **26**, 103202.
- 37 P. Asselin, R. F. L. Evans, J. Barker, R. W. Chantrell, R. Yanes, O. Chubykalo-Fesenko, D. Hinzke and U. Nowak, *Physical Review B*, 2010, **82**, 054415.
- 38 J. M. Crowley, J. Tahir-Kheli and W. A. Goddard III, *The Journal of Physical Chemistry Letters*, 2016, **7**, 1198–1203.
- 39 W.-B. Zhang, Q. Qu, P. Zhu and C.-H. Lam, *Journal of Materials Chemistry C*, 2015, **3**, 12457–12468.
- 40 Q. Song, C. A. Occhialini, E. Ergeçen, B. Ilyas, D. Amoroso, P. Barone, J. Kapeghian, K. Watanabe, T. Taniguchi, A. S. Botana *et al.*, *Nature*, 2022, **602**, 601–605.
- 41 Z. An, Y. Su, S. Ni and Z. Guan, *The Journal of Physical Chemistry C*, 2022, **126**, 11330–11340.