



## Optical Spectroscopy and Photoionization Model of Planetary Nebula NGC 6572

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**Abstract.** We investigated NGC 6572 through optical spectroscopy to determine its kinematical and chemical properties. Two intermediate dispersion spectra ( $R\sim 5000$ ) centered around  $H\alpha$  and  $H\beta$  were used to derive the nebular expansion velocity from emission lines associated with  $H\alpha$ ,  $H\beta$ , [OIII], [NII], and [SII] ions. A low dispersion spectrum ( $R\sim 1000$ ) was used to determine the nebular electron temperature, density, and chemical composition. We performed photoionization modeling to construct a self-consistent nebular model, while intermediate-resolution spectral images showed the global elliptical structure of the nebula. The expansion velocity deduced from most of the emission lines is consistent with the typical expansion velocity of planetary nebulae, i.e., around  $15\text{-}20\text{ km s}^{-1}$ . The nebular physical properties align well with those determined by other studies. The nebular abundances were found to be lower than the solar abundances (except for oxygen) but still comparable with the abundances derived by other researchers. The photoionization model generated spectral lines that are consistent with the lines found in the observations. Further spectroscopic observations with higher resolution and wider range at various position angles will be very useful to reveal a more complete and detailed structure of the nebula and to improve the determination of the nebular physical properties.

**Keywords:** *chemical abundances; expansion velocity; photoionization model; planetary nebula; spectroscopy.*

### 1 Introduction

NGC 6572 is among the brightest planetary nebulae (PNe) in the night sky. It makes the observation of the PN quite easy in terms of instrumental requirements, i.e., one can make use of a small-class telescope to obtain data with an adequate signal-to-noise ratio (SNR). Therefore, NGC 6572 is a PNe that has already been studied extensively. It not only shows typical PNe characteristics; several studies

have revealed that NGC 6572 has variability in its brightness and in the intensity of its nebular emission lines and stellar spectra. This indicates that NGC 6572 possesses an internal activity such that it affects its nebular brightness and spectral features.

Hyung *et al.* [1] investigated the variability of NGC 6572 using detailed high-resolution spectra in the optical wavelength from 3650 to 10050 Å and high-dispersion mode ultraviolet (UV) spectra from the International Ultraviolet Explorer (IUE) satellite. Hyung *et al.* [1] found that the central star of NGC 6572 may have a cycle of the order of 70 years. Miranda *et al.* [2] investigated the morphology of NGC 6572 using high-resolution images and reported that it possesses a point-symmetric feature that is enveloped within a previously known elliptical structure. Using high-resolution long-slit spectra, Miranda *et al.* [2] found the existence of highly collimated outflows, which may take part in the formation of the point-symmetric signature.

In this work, we report our long-slit spectroscopic observations at intermediate and low resolutions. We were able to obtain the recombination lines of H $\alpha$ , H $\beta$ , and He 6678 Å in the intermediate-resolution spectra. We also obtained several forbidden lines of [OIII] 4959 Å, [OIII] 5007 Å, [NII] 6548 Å, [NII] 6584 Å, [SII] 6716 Å, and [SII] 6731 Å. Forbidden lines in PNe spectra are useful to help investigate the nature of a PNe, e.g., nebular electron temperature and electron density, which are required to derive its nebular chemical composition. Different lines correspond to different ionization layers in the nebula, along with differences in their physical properties. Particularly for [OIII] and [NII], the emission lines can be used to constrain the temperature of the central star of planetary nebulae (CSPN).

We utilized the intermediate-resolution spectra to calculate the nebular expansion velocities and investigate its kinematical characteristics. In addition to that, the low-resolution spectra, covering the wavelengths from 3800 Å to 7700 Å, were used to derive the physical and chemical properties of the PN. The lines present in the spectra were employed as a constraint in the photoionization modeling, to make sure that the model could approximate the observed data to a certain degree of satisfaction.

## 2 Observations and Data Reductions

### 2.1 Intermediate-Resolution Spectra

We successfully secured two intermediate-resolution spectra taken on 14 June and 10 August 2021. Both spectra were obtained using a spectrograph mounted on a 10-inch ( $F/9.8$ ) telescope with a CCD camera at Bosscha Observatory ITB,

Indonesia. The spectrograph was operated with a grating of 1200 grooves/mm and a slit with a width of 25  $\mu\text{m}$ . Using these instruments, NGC 6572 was observed with an exposure time of 2700 seconds on 14 June and 3600 seconds on 10 August. During both observations the slit was centered on the object at a position angle (PA) of  $0^\circ$ .

We cleaned the spectrograms (subtracted them from biases, darks, and corrected the normalized flat) and performed spectra extraction (including wavelength and flux calibration) using long-slit IRAF (Image Reduction and Analysis Facility) routines. A Fe-Ne-Ar lamp was used as the calibration lamp to obtain wavelength calibration frames. The final spectra had a resolution of  $R \sim 5000$  and a spectral linear dispersion of 0.34  $\text{\AA}/\text{pixel}$ . The spectral wavelength coverages were from 4810-5090  $\text{\AA}$  for the 14 June observation and 6475-6740  $\text{\AA}$  for 10 August observation. Therefore, we obtained emission lines at  $\text{H}\beta$ , [OIII] 4959  $\text{\AA}$ , [OIII] 5007  $\text{\AA}$ ,  $\text{H}\alpha$ , [NII] 6548  $\text{\AA}$ , [NII] 6584  $\text{\AA}$ , HeI 6678  $\text{\AA}$ , [SII] 6716  $\text{\AA}$ , and [SII] 6731  $\text{\AA}$ .

## 2.2 Low-Resolution Spectra

We also obtained low-resolution spectra to determine and analyze the chemical properties of NGC 6572. On 10 July 2022, we observed NGC 6572 using a fast and low-resolution long-slit spectrograph (Malasan *et al.* [3]) mounted on an 11-inch ( $F/10$ ) telescope with a CCD camera at Bosscha Observatory ITB, Indonesia. With an exposure time of 300s, the slit was centered on the object at  $0^\circ$  PA. Data reduction was carried out in the same manner as for the intermediate-resolution spectra, including the calibration lamp that was used. The procedure resulted in a spectrum with a resolution of  $R \sim 1000$ , covering a wavelength range of 3875-7760  $\text{\AA}$  and with a spectral linear dispersion of 2.88  $\text{\AA}/\text{pixel}$ .

## 3 Nebular Kinematics

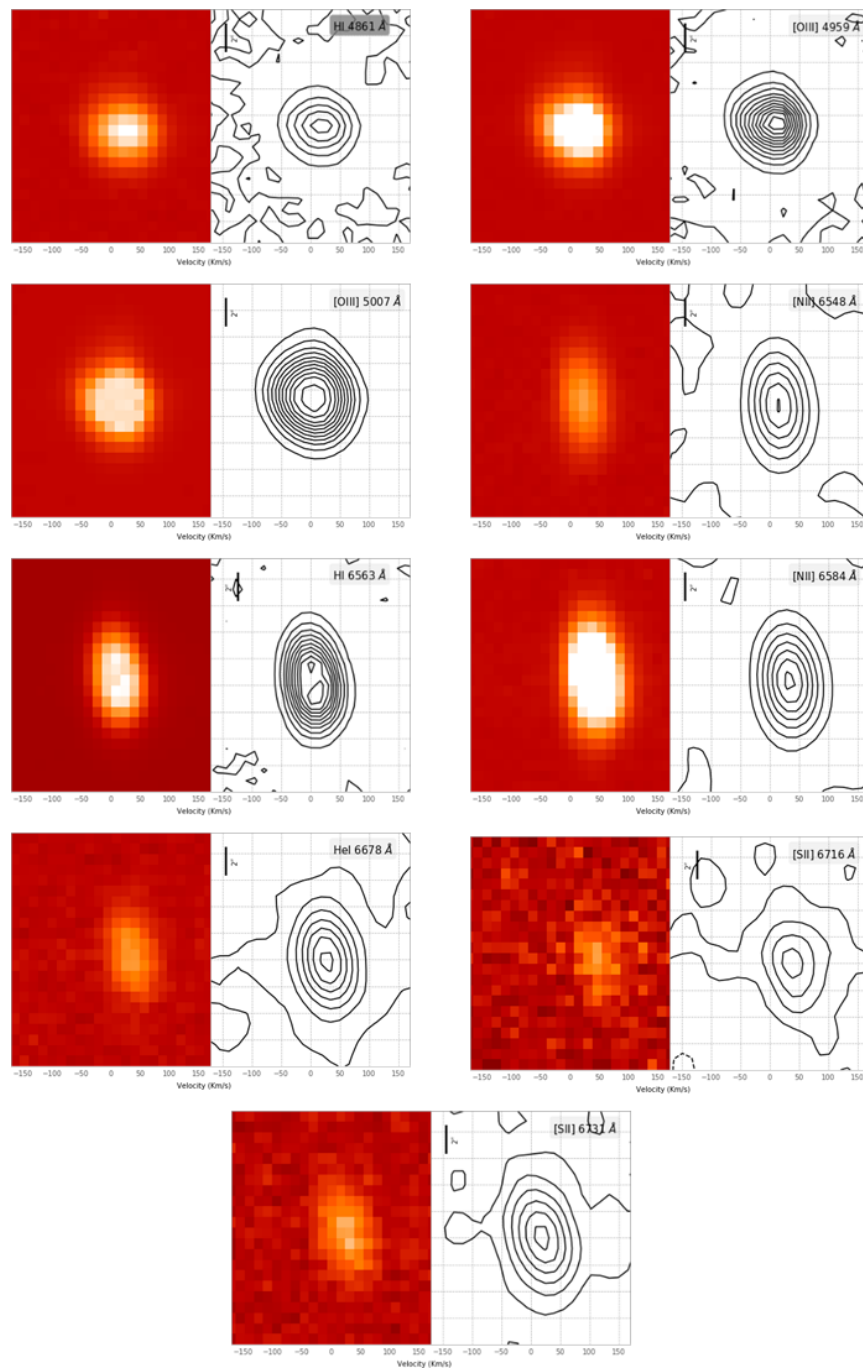
To analyze the nebular kinematics, we present the spectral images as well as the position-velocity (PV) contour maps in Figure 1. The PV maps were constructed from the reduced spectrograms of the intermediate-resolution spectra, where we converted the horizontal axis from wavelength into velocity using Doppler shift, with laboratory wavelengths as the reference. The spectral resolution from full width at half-maximum (FWHM) was  $\sim 15 \text{ km s}^{-1}$  and the spatial resolution was  $\sim 0.73''/\text{pixel}$ . From the PV contour maps, we can see the elliptical nebular shell of NGC 6572, but there is no sign of prominent and well-separated emission peaks. This is likely caused by the insufficient resolution to detect such a feature in the spectral images.

With the given spectral resolution, the emission lines of each detected ion do not show the double-peaked feature but rather single and broad emissions. The nature of the object, i.e., a compact planetary nebula with high surface brightness, adds to the underlying reason for this lack of detection. Thus, to determine the nebular expansion velocity we assumed that the nebula is an expanding spherical shell of gas that produces emission lines that can be represented by a Gaussian distribution. This introduces a separation velocity ( $\Delta v^2$ ) that corresponds to the part of the nebular shell that is approaching and receding towards the observer. The separation velocity, which reflects the intrinsic velocity of the expanding nebula, can be calculated by subtracting the observed velocity from the observed instrumental and thermal line (Eq. 1).

$$\Delta v^2 = v_{observed}^2 - v_{instrumental}^2 - v_{thermal}^2 \cdot \quad (1)$$

The comparison lamp line measurements provide the  $v_{instrumental}$  ( $c \times \frac{FWHM}{\lambda}$ ). The instrumental line broadening (converted into velocity) of the lines H $\alpha$ , [NII], HeI, and [SII] was  $\sim 61$  kms $^{-1}$ , while the instrumental line broadening for H $\beta$  and [OIII] was  $\sim 60$  kms $^{-1}$ . The thermal broadening ( $v_{thermal}$ ) corresponding to the elements H, He, O, N, and S, was 21.26 kms $^{-1}$ , 10.63 kms $^{-1}$ , 5.31 kms $^{-1}$ , 5.68 kms $^{-1}$ , and 3.76 kms $^{-1}$ , respectively (Choi *et al.* [4]). In Table 1, we present the nebular expansion velocities from every detected line. We found that the expansion velocities derived from blue spectra (H $\beta$  and [OIII]) were somewhat larger compared to those derived from red spectra (H $\alpha$ , [NII], HeI, and [SII]), while the red spectra expansion velocities were close to the typical expansion velocity of planetary nebulae ( $\sim 14$ -25 kms $^{-1}$ ).

This can also be seen from the more rounded shape of the spectral images and the PV diagram of the blue spectra compared to the red spectra. Miranda *et al.* [2] in 1999 suggested that NGC 6572 possesses two velocities, i.e., 17 kms $^{-1}$  and 50 kms $^{-1}$ , therefore, it is possible that during the observation on 14 June, we observed a nebular region that exhibits this velocity. This means that NGC 6572 possesses several expansion velocities, where we can infer that several episodes of nebular ejection occurred during the formation of NGC 6572. The more rounded shape of the blue region spectra could also be the subject of the location where the lines are produced. H $\beta$  especially is related to a location close to the center of the nebula, hence the CSPN. At such a location, the nebula does not interact so much with its surrounding environment, allowing it to keep its spherical shape.



**Figure 1** Spectral images and PV diagrams of the emission lines of NGC 6572.

**Table 1** Expansion velocity from several ions.

Ion	$\lambda$ (Å)	$v_{\text{exp}}$ ( $\text{kms}^{-1}$ )
H $\beta$	4861.33	50.37
H $\alpha$	6562.82	15.22
HeI	6678.15	15.03
[NII]	6548.03	15.47
	6583.45	16.73
[OIII]	4958.92	48.31
	5006.84	48.72
[SII]	6716.16	17.76
	6730.85	17.89

Table 2 presents the variation of the mean of the expansion velocity from the red spectra, as well as the ionization potential of each ion. PNe stratification implies that high excitation lines (from ions with high IP) are emitted at a location closer to the CSPN. Meanwhile, low excitation lines are emitted at a location further out from the CSPN. If the geometry is spherical, we should be able to see that the speed of the nebular gas expansion increases with the distance from the center (Wilson [5]). Since NGC 6572 is an ellipsoidal PN, there must be additional factors for the outward acceleration. The internal point-symmetric structure would likely contribute to this matter. From Table 2 we can also see that the measured velocities do not strictly follow the simple rule of a linear increasing velocity, given the increase in the ionization potential. This may also be caused by the internal structure that shapes the nebula.

**Table 2** Expansion velocity vs ionization potential (IP).

Ion	IP (eV)	$v_{\text{exp}}$ ( $\text{kms}^{-1}$ )
[SII]	10.4	17.82
H $\alpha$	13.6	15.22
[NII]	14.5	16.10
HeI	24.6	15.03

#### 4 Nebular Properties

The intensity of the emission lines from the low-resolution spectrum were measured and used to deduce the nebular physical properties. We used NEAT (Nebular Empirical Analysis Tool) code [6] to determine the electron temperature and electron density from standard diagnostics. The code requires measured line fluxes, which are then corrected based on the interstellar extinction using the ratios of the intensity of the Balmer lines. The code was also used to calculate the ionic abundances from the flux-weighted averages of the emission lines of each

species. The total elemental abundances were calculated by incorporating an ionization correction factor (ICF). In this work, we applied the Howarth extinction law [7] for the interstellar extinction correction, and we used the ICF scheme from Delgado-Inglada *et al.* [8].

**Table 3** Nebular physical properties.

	$c(\text{H}\beta)$	$\text{Log } F_{\text{H}\beta}$	$T_e[\text{NII}]$	$T_e$ [OIII]	$T_e[\text{ArIII}]$	$N_e[\text{SII}]$	$(\text{N1+N2})/\text{H}$ $b$
This work	0.37	-11.33	12200 K	10900 K	9400 K	3400 K	19.56
Hyung <i>et al.</i> [1]	0.40	-9.82		11000 K		...	...

The nebular properties we derived are summarized in Table 3. The reddening,  $c(\text{H}\beta)$ , and electron temperature from collisionally-excited lines were in relatively good agreement with the values from Hyung *et al.* [1]. We also found  $\log F_{\text{H}\beta}$  and the density diagnostic from [SII] to be 0.37 and 3400 K, respectively. The line strength of 19.56 indicates that NGC 6572 has a moderate line strength, which means that the CSPN of NGC 6572 has a relatively moderate temperature and luminosity for a white dwarf star.

The elemental abundances (Table 4) calculated from the available emission lines were mostly lower than those from Hyung *et al.* [1], which may be caused by the limited number of emission lines that were taken into account in the calculation of the ionic abundances. The overall discrepancy from values determined by other researchers is subject to various aspects, e.g., a different nebular region observed as well as the data processing and calculation of physical properties procedures could have contributed to the departure from the resulted values.

**Table 4** Nebular elemental abundances ( $\log X/\text{H}$ ) measured in this work and Hyung *et al.* [1].

Element	This work	Hyung <i>et al.</i> [1]
N	-4.05	-4.208
O	-3.28	-3.463
Ne	-4.15	-4.244
Ar	-5.51	-5.724
S	-5.81	-5.666
Cl	-6.79	-7.171
He	-0.82	0.0044

#### 4.1 Photoionization Model

Calculating nebular physical properties and elemental abundances using an empirical method is limited by the availability of emission lines in the spectra. A

theoretical model of the nebular ionization structure constructed through photoionization modeling can provide more complete data and information of the lines emitted by the nebula. We employed the photoionization code Cloudy (version 17.02, [9]) to construct a self-consistent nebular structure.

Cloudy utilizes a set of input parameters that define the ionizing source as well as the nebula and then solves the ionization balance and energy balance equations at each point of the nebula. The code calculates the radiation transfer and generates the model spectrum at the end. The emitted spectra of the model then can be compared to the ones found in the observation.

**Table 5** The input parameters for the photoionization modeling.

<b>Central star</b>	
$T_{\text{eff}}$	59979 K
$\log g$	4.2
$\log L/L_{\odot}$	36.75
<b>Nebula</b>	
Geometry	Sphere
$\log N_H$	$4.01 \text{ cm}^{-3}$
$\log r$	17.19 cm
Density	$N_0(r/r_0)^{-2}$
<b>Abundances (<math>\log X/H</math>)</b>	
N	-4.05
O	-3.28
Ne	-4.15
Ar	-5.51
S	-5.81
Cl	-6.79
He	-0.82

The input parameters listed in Table 5 consist of parameters regarding the ionizing source (CSPN), the nebular physical description, and the initial nebular chemical abundances that were previously determined empirically. The ionizing source is characterized with the SED from the stellar atmosphere that is defined by its luminosity, temperature, and gravity, all in a grid of stellar atmosphere models adopted from Hubeny [10].

The nebular parameters included the description of its geometry, density, radius, and chemical abundances. The spherical geometry was adopted so that the ionizing source would be located at the center of the nebula, and we did not employ a realistic complex geometry. The density was taken from the electron density with a decreasing distribution profile with a factor of -2. The initial elemental abundances were set to the values calculated previously using an empirical method.



To achieve the best-fitted model, we treated the density, radius, and the abundances as free parameters and varied them until the observables were satisfactorily reproduced by the model. Thus, the code will adjust for the line strength of different elements. Goodness of fit was obtained by calculating the root mean square (*RMS*) value using Eq. (2).

$$RMS = \sqrt{\frac{1}{N} \sum_i \left[ 1 - \frac{M_i}{O_i} \right]^2}. \quad (2)$$

$M_i$  and  $O_i$  are the modeled and observed values of the  $i$ -th observable and  $N$  denotes the total number of observables. The model is considered acceptable if the *RMS* is less than unity.

The comparison between the observed and the modeled parameters are given in Table 6. The observed parameters are lines measured from the spectra and have been dereddened, and the electron temperature was calculated from the [OIII] lines. From the model, we obtained the line intensities of various ions that were not detected in the observation. The *RMS* calculation was based on the lines measured in the observation, and the *RMS* value was found to be quite far below unity.

The measured emission lines from the low-resolution spectra were close to those measured by Hyung *et al.* [1] but there were still differences in the intensities, especially on the edge of the spectra. The photoionization model on the other hand was mostly successful in generating spectra close to the observed lines. Discrepancies in the modeled spectra were also apparent, especially at lines that were not available in the observations. This was somewhat expected, since we did not provide constraints apart from the lines observed in the optical spectra and in our case the code could not successfully generate the lines to be close to those measured by Hyung *et al.* [1].

The elemental abundances resulted from the modeling are presented in Table 7. The values were still lower than those measured by Hyung *et al.* [1]. This could indicate that the nebula had undergone changes, especially compared to the previous study Hyung *et al.* [1]. It may also be a sign of the spectral variability apparent in nebula, but further investigation is needed to confirm this.

**Table 6** Comparison of the measured line intensities from this work and Hyung *et al.* [1] and Cloudy model fit.

Parameter	Obs.	Hyung <i>et al.</i> [1] (1)	Hyung <i>et al.</i> [1] (2)	Model
$\lambda 3727$	...	16.88	13.25	246.85
$\lambda 3868$	16.8	95.31	87.733	16.12
$\lambda 4101$	37.30	30.89	27.91	25.37
$\lambda 4340$	59.50	47.49	35.04	46.35
$\lambda 4363$	8.71	8.31	8.90	10.56
$\lambda 4471$	5.01	5.30	5.61	2.91
$\lambda 4686$	...	0.38	0.39	11.35
$\lambda 4740$	...	2.26	2.48	1.00
$\lambda 5007$	1530.00	1107.54	1077.36	1526.88
$\lambda 5200$	...	0.07	0.09	0.19
$\lambda 5518$	...	0.16	0.17	0.99
$\lambda 5538$	...	0.38	0.40	1.44
$\lambda 5755$	2.16	1.46	1.71	0.78
$\lambda 5876$	15.80	16.93	16.66	8.46
$\lambda 6300$	9.82	3.05	3.82	30.95
$\lambda 6312$	4.81	0.66	0.80	0.31
$\lambda 6563$	359.00	307.54	280.48	292.06
$\lambda 6584$	...	37.39	44.64	40.89
$\lambda 6678$	...	3.86	4.22	2.28
$\lambda 6717$	...	0.53	0.47	1.94
$\lambda 6731$	3.65	1.07	1.19	3.29
$\lambda 7006$	0.28	10.23	9.68	0.04
$\lambda 7136$	28.40	15.37	16.32	30.12
T[OIII]	10900 K		11000 K	10100 K
RMS				0.45
$\lambda 1549$	...		1.50	3.74
$\lambda 1909$	...		>11.369	59.55
$\lambda 2326$	...		11.82	13.32
$\lambda 9850$	...		0.02	0.34
$\lambda 1486$	...		...	0.07
$\lambda 1750$	...		2.59	1.04
$\lambda 2424$	...		...	0.44
$\lambda 8047$	...		0.22	0.41
$\lambda 9532$	...	17.03	13.28	12.53

Note: The emission line strength is relative to  $H\beta = 100$ . (1) and (2) correspond to data collected in 1990 and 1991, respectively.

**Table 7** Comparison of the modeled chemical abundances in this work, Hyung *et al.* [1], and solar photospheric abundances (Asplund [11]).

Element	This work	Hyung <i>et al.</i> [1]	Model	Sun
N	-4.05	-4.208	-4.6239	-4.17
O	-3.28	-3.463	-3.1384	-3.31
Ne	-4.15	-4.244	-4.6667	-3.94
Ar	-5.51	-5.724	-5.5649	-5.62
S	-5.81	-5.666	-6.0295	-4.88
Cl	-6.79	-7.171	-6.6979	-6.69
He	-0.82	0.0044	-1.1680	-1.086

## 5 Conclusions

From the intermediate-resolution long-slit spectra we can see an overall image of a nebular expansion pattern that corresponds to the ellipsoidal feature of the nebular morphology. The limited resolution of the observations hampered us to detect the internal point-symmetric structure as well as the collimated flows. Observations with higher resolution would provide a better image of the existing internal structure, while different slit orientations are needed to obtain the kinematics from another nebular region. The calculated expansion velocities are typical planetary nebulae expansion velocities for H $\alpha$ , [NII] 6548 Å, [NII] 6584 Å, HeI 6678 Å, [SII] 6716 Å, and [SII] 6731 Å emission lines, but the value is somewhat higher for blue lines H $\beta$ , [OIII] 4959 Å, and [OIII] 5007 Å. We suggest that the high velocity observed in the blue spectra may be a result of observing a different region within the nebula compared to the region observed in the red spectra. This, in turn, could indicate the presence of different velocity components in the nebula.

The nebular physical properties agree in terms of reddening and electron temperature, but differences in the measured line intensities appeared to affect the log FH $\beta$  and electron density. Several lines measured in our observation differed from those in Hyung *et al.* [1], which may suggest spectral variability in NGC 6572. By utilizing the measured lines and parameters from both the nebula and the CSPN, we successfully constructed a photoionization model with an RMS below unity. The overall chemical abundances, derived through both the empirical method and the photoionization model, were lower than those in Hyung *et al.* [1]. This difference may also indicate that the nebula may have evolved since Hyung *et al.* [1].

The complete kinematical maps with adequate resolution combined with high spatial resolution images can be further utilized to construct a 3D model of the nebula. Combining spatial and kinematical models with a photoionization model would produce a more comprehensive picture of the system. The complete picture

of the nebula would allow for the deep investigation of the formation of the object.

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