# The bright end of the galaxy luminosity function at $z \simeq 7$ from the VISTA VIDEO survey

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# ABSTRACT

We have conducted a search for  $z \simeq 7$  Lyman-break galaxies over 8.2 deg<sup>2</sup> of near-infrared imaging from the Visible and Infrared Survey Telescope for Astronomy (VISTA) Deep Extragalactic Observations (VIDEO) survey in the *XMM–Newton*-Large Scale Structure (XMM-LSS) and the *Extended Chandra Deep Field*-South (ECDF-S) fields. Candidate galaxies were selected from a full photometric redshift analysis down to a Y + J depth of 25.3 (5 $\sigma$ ), utilizing deep auxiliary optical and *Spitzer*/Infrared Array Camera (IRAC) data to remove brown dwarf and red interloper galaxy contaminants. Our final sample consists of 28 candidate galaxies at  $6.5 \le z \le 7.5$  with  $-23.5 \le M_{\rm UV} \le -21.6$ . We derive stellar masses of  $9.1 \le \log_{10}(M_\star/M_{\odot}) \le 10.9$  for the sample, suggesting that these candidates represent some of the most massive galaxies known at this epoch. We measure the rest-frame ultraviolet (UV) luminosity function (LF) at  $z \simeq 7$ , confirming previous findings of a gradual decline in number density at the bright end ( $M_{\rm UV} < -22$ ) that is well described by a double power law (DPL). We show that quasar contamination in this magnitude range is expected to be minimal, in contrast to conclusions from recent pure-parallel *Hubble* studies. Our results are up to a factor of 10 lower than previous determinations from optical-only ground-based studies at  $M_{\rm UV} \le -23$ . We find that the inclusion of *YJHK*<sub>s</sub> photometry is vital for removing brown dwarf contaminants, and  $z \simeq 7$  samples based on red optical data alone could be highly contaminated ( $\gtrsim 50$  per cent). In comparison with other robust z > 5 samples, our results further support little evolution in the very bright end of the rest-frame UV LF from z = 5-10, potentially signalling a lack of mass quenching and/or dust obscuration in the most massive galaxies in the first Gyr.

Key words: galaxies: evolution - galaxies: formation - galaxies: high-redshift.

# **1 INTRODUCTION**

Observing the formation and evolution of galaxies at very high redshifts is vital for understanding the build-up of early structures and all subsequent galaxy evolution to the present day. The Lyman-break technique has been used successfully for over three decades to search for high-redshift galaxies (Guhathakurta, Tyson & Majewski 1990; Steidel et al. 1996) via the redshifted Lyman  $\alpha$  break at  $\lambda_{rest} = 1216$  Å, by utilizing the strong spectral break that appears at the wavelength of the Lyman  $\alpha$  line due to absorption by the intergalactic medium (IGM) along the line of sight. The selection of high-redshift ( $z \ge$ 7) Lyman-break galaxies (LBGs) has been revolutionized thanks to the unparalleled near-infrared (NIR) capabilities of the James Webb Space Telescope (JWST). Early results have identified and spectroscopically confirmed galaxies in the range z = 10.6-13.1(e.g. Arrabal Haro et al. 2023; Bunker et al. 2023; Curtis-Lake et al. 2023) and rest-frame optical spectroscopy has revealed the gas-phase properties for the first time (e.g. Cameron et al. 2023; Curti et al. 2023; Fujimoto et al. 2023). One surprising result from the first samples of galaxy candidates found by *JWST* has been the discovery of an unexpected number of luminous ( $M_{UV} \sim -21$ ) sources at z > 8 (e.g. Finkelstein et al. 2022; Naidu et al. 2022) that have challenged models of galaxy formation (Ferrara, Pallottini & Dayal 2023).

Galaxies form and reside within dark matter haloes, and the luminosity function (LF, number density as a function of brightness) provides a connection to the dark matter halo mass function and early galaxy formation since galaxy formation efficiency has a strong dependency on halo mass (e.g. Vale & Ostriker 2006; Behroozi, Wechsler & Conroy 2013; Wechsler & Tinker 2018). This makes the LF a key measurement in galaxy evolution. As at  $z\gtrsim 6$  LBGs are typically selected and studied in the NIR, which provides access to the rest-frame ultraviolet (UV) emission from young stars, many studies have derived the rest-frame UV LF (hereafter UV LF) from photometric samples at z > 4 (e.g. McLure et al. 2013; Finkelstein et al. 2015; Stefanon et al. 2019; Bouwens et al. 2021). The shape of the UV LF gives insight into feedback properties within galaxies and ionizing properties of galaxies in the early Universe (e.g. Benson et al. 2003; McLure et al. 2010; Bradley et al. 2012; Bowler et al. 2015). A common form used to fit the rest-frame UV LF in the local Universe is the Schechter function (Schechter 1976),  $\phi(L) dL = \phi^* (L/L^*)^\alpha e^{-L/L^*} d(L/L^*)$ , where  $\alpha$  is the faint-

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end slope,  $L^*$  is the characteristic luminosity, and  $\phi^*$  is the density normalization. The faint end follows a power law and a rapid decline is seen in the bright end, necessitating an exponential cut-off beyond the knee at  $L^*$ . This suggests the growth of the brightest, most massive galaxies is quenched in lower redshift LBG samples (e.g. van der Burg, Hildebrandt & Erben 2010; Stevans et al. 2018; Adams et al. 2023), since the dark matter halo mass function has a shallower highmass slope and is more akin to a double power law (DPL). Probing the bright end of the rest-frame UV LF requires galaxy searches to be conducted in degree-scale surveys, since the number density of bright objects declines rapidly. For example, LBGs at  $z \simeq 7$  have a surface density of a few per deg<sup>2</sup> at  $M_{\rm UV} < -22$  (Bowler et al. 2014). The widest survey fields imaged by the Hubble Space Telescope (HST) only cover 1136 arcmin<sup>2</sup> (e.g. Bouwens et al. 2021), proving only weak constraints on the LF at  $M_{\rm UV} \lesssim -21$ . Prior to the launch of space missions capable of imaging a wide area (e.g. Euclid and Roman), ground-based NIR surveys provide the only insight into the number density of the very bright galaxy population at z > 6. The first determinations of the  $z \simeq 7$  UV LF with degree-scale surveys showed an excess of bright galaxies compared to the commonly used Schechter function, with a DPL fit preferred (Bowler et al. 2012, 2014). The DPL form has also been identified in optical-only studies from the Hyper Suprime-Cam (HSC) GOLDRUSH program (Ono et al. 2018; Harikane et al. 2022). This bright-end excess appears to continue out to  $z \simeq 10$  (Stefanon et al. 2019; Bowler et al. 2020). Additionally, there is evidence that whatever mechanism causes the excess in bright galaxies at  $z \simeq 7$  has a rapid onset, with a clear steepening of the bright-end slope found to occur in the 400 Myr between z = 7 and z = 5 (Bowler et al. 2015; Adams et al. 2023)

Interestingly, many simulations do not predict intrinsic UV LFs with a DPL shape at z > 5, with a steeper bright end only appearing after the addition of significant dust obscuration (e.g. Cai et al. 2014, see compilation in Bowler et al. 2015). While at lower redshifts the steep decline in luminous/massive galaxies (e.g. Yang, Mo & van den Bosch 2009) can be attributed to mass quenching (Peng et al. 2010), attenuation, or a lack of attenuation is predicted to be the more significant effect in shaping the UV LF at high redshifts (e.g. Cai et al. 2014; Vijayan et al. 2021). Additionally, as active galactic nuclei (AGN) begin to grow and turn on in the early Universe, energetic feedback from their accretion disc may suppress star formation in massive galaxies (e.g. Bower, Benson & Crain 2012; Davé et al. 2019; Lovell et al. 2022). While quasars at  $z \simeq 7$  are extremely rare (e.g. Mortlock et al. 2011; Bañados et al. 2018; Wang et al. 2021), the quasar LF appears to rapidly increase to lower redshifts, which leads to the faint end of the AGN LF being comparable to the bright end around  $M_{\rm UV} \simeq -23$  at  $z \simeq 5$  (Cano-Díaz et al. 2012; Adams et al. 2023; Harikane et al. 2022). By determining the functional form of the LF at different redshifts, the onset of potential quenching and/or dust obscuration can be investigated. Comparison of the bright end of the galaxy UV LF with simulations, plus follow-up observations with the JWST and Atacama Large Millimeter/submillimeter Array (ALMA), can then uncover the dominant mechanism responsible for the shape of the UV LF.

A key step in searching for  $z \simeq 7$  galaxies is the removal of red, dusty galaxies at  $z \sim 1-2$  and cool Galactic M, L, and T dwarf stars (e.g. Stanway et al. 2008; Bowler et al. 2015) that have photometry that can mimic a Lyman break. The selection of robust samples requires the use of deep optical, NIR, and *Spitzer/Infrared Array Camera* (IRAC) photometry via a colour– colour selection or spectral energy distribution (SED) fitting (e.g. Bowler et al. 2014; Bouwens et al. 2022). In this work, we conduct a search for  $z \simeq 7$  galaxies within the Visible and Infrared Survey Telescope for Astronomy (VISTA) Deep Extragalactic Observations (VIDEO) in *XMM–Newton*-Large Scale Structure (XMM-LSS) and *Extended Chandra Deep Field*-South (ECDF-S) fields. Crucially, these fields are covered by a range of deep multiwavelength data including the optical, red optical, and NIR filters extending to 2.2 µm (and up to 5 µm with *Spitzer*/IRAC data) allowing the robust removal of contaminant populations. This work comprises the largest ground-based optical + NIR search for  $z \simeq 7$  galaxies to date, providing the most robust constraints on the bright end of the rest-frame UV LF prior to the launch of wide-area space missions.

This paper is structured as follows. In Section 2, we outline the multiwavelength data sets, and in Section 3, we present the selection process of our  $z \simeq 7$  LBG sample. We present our candidate  $z \simeq 7$  galaxies and their properties in Section 4, and in Section 5, we measure the rest-frame UV LF at  $z \simeq 7$  and present our results. We discuss contamination and potential contribution by AGN in Section 6. We then present our conclusions in Section 7. In the Appendix, we present SED fitting and postage stamps of all candidate objects. We also present an alternate calculation of the UV LF at  $z \simeq 7$  using a more inclusive sample. All magnitudes are reported in the AB system (Oke & Gunn 1983). We assume a standard  $\Lambda$  cold dark matter ( $\Lambda$ CDM) cosmology,  $H_0 = 70$  km s<sup>-1</sup> Mpc<sup>-1</sup>,  $\Omega_M = 0.3$ ,  $\Omega_{\Lambda} = 0.7$ .

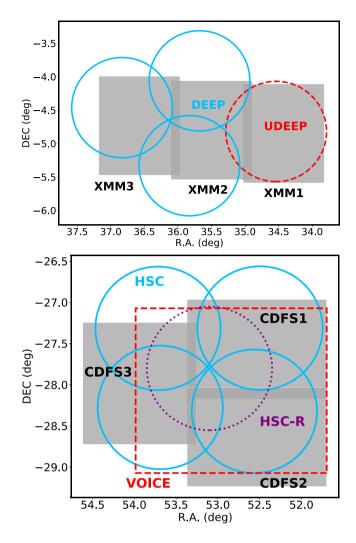
#### 2 DATA

In this study, we use multiwavelength data across two fields, XMM-LSS and ECDF-S. These fields are selected for their coverage by the VISTA VIDEO survey in the NIR *YJHK*<sub>s</sub> bands (Jarvis et al. 2013), in addition to having overlapping deep optical data. We use the final data release of the VISTA VIDEO survey, which is now publicly available.<sup>1</sup> The VIDEO data in the ELAIS-S1 field are not included due to a current lack of deep optical imaging, particularly in the red optical bands around the expected position of the Lyman break. The total overlapping optical and NIR area used in this work covers 8.22 deg<sup>2</sup>. The field footprints are shown in Fig. 1. The optical and NIR bands are used for SED fitting, and *Spitzer/IRAC* imaging in the mid-infrared is used to remove low-redshift dusty interlopers (see Section 3.4).

#### 2.1 XMM-LSS

The VIDEO data are split into three VISTA pointings, or tiles (XMM1, XMM2, and XMM3), of 1.5 deg<sup>2</sup> each as shown in Fig. 1. Optical data are taken from the Data Release 3 (DR3) of the Hyper Suprime-Cam Subaru Strategic Program (HSC-SSP; Aihara et al. 2022). XMM1 contains an 'ultradeep' HSC pointing with longer exposure than the 'deep' pointings. The DR3 data provide an increase in 5 $\sigma$  depth of ~0.1–0.7 mag on the Data Release 2 (DR2; Aihara et al. 2019). HSC-SSP also provides two narrow bands, NB816 and NB921, near the expected position of the Lyman break at  $z \simeq 7$ . This can aid in reducing errors on photometric redshifts. In addition, imaging from the *Spitzer* Extragalactic Representative Volume Survey (SERVS; Mauduit et al. 2012) is used along with the deeper *Spitzer* Extended Deep Survey (SEDS; Ashby et al. 2013) where available. The area containing overlapping VIDEO and HSC

<sup>1</sup>http://www-astro.physics.ox.ac.uk/~video/public/Home



**Figure 1.** XMM-LSS footprint (top) and ECDF-S footprint (bottom). The VIDEO tiles are shown in grey and the HSC pointings in blue circles. In ECDF-S, the VOICE pointing is shown as a dashed red square. The central HSC-*R* pointing is shown as a dotted purple circle. In XMM-LSS, the ultradeep HSC pointing is shown as a dashed red circle.

data is 4.33 deg<sup>2</sup>, accounting for masking of bright stars and artefacts such as stellar ghosts.

# 2.2 ECDF-S

Similarly to XMM-LSS, the VIDEO images in ECDF-S are split into three tiles (CDFS1, CDFS2, and CDFS3) of 1.5 deg<sup>2</sup> each as shown in Fig. 1. Optical data are composed of ancillary HSC data compiled by Ni et al. (2019), in the form of four pointings in *G*, *I*, and *Z* and a central single pointing of the *R* band. The HSC data are a combination of different observation conditions, leading to poor seeing in the *I* and *Z* bands. The full width at half-maximum (FWHM) of the point spread function (PSF) varies by 0.1 arcsec across the field in these bands (see Section 2.4). Furthermore, there is poor coverage in the *R* band. We therefore complement the HSC data with optical data from the VST Optical Imaging of the CDFS and ES1 Fields (VOICE; Vaccari et al. 2017), covering 4 deg<sup>2</sup>. *Spitzer/*IRAC imaging is taken from the Cosmic Dawn Survey (CDS; Euclid Collaboration 2022a). The area containing overlapping VIDEO, HSC, and VOICE data is 3.89 deg<sup>2</sup>.

**Table 1.** The  $5\sigma$  limiting magnitudes for each band used in this work, in each VISTA VIDEO tile (see Fig. 1). The local depths were measured by placing 2 arcsec diameter circular apertures on empty regions of the images. The depth quoted here is the mode of these local depths. IRAC depths were measured in 2.8 arcsec diameter circular apertures to account for the poorer resolution. The first section shows bands from VLT Survey Telescope (VST). The second section shows bands from HSC. The third section shows bands from VISTA. The final section shows bands from Spitzer/IRAC.

Filter	XMM1	XMM2	XMM3	CDFS1	CDFS2	CDFS3
и	_	_	_	25.5	25.6	25.4
g	-	-	-	26.0	25.9	26.0
r	-	-	-	26.0	26.0	26.0
i	-	-	-	24.6	24.8	24.7
G	27.6	27.1	27.1	26.5	26.2	26.4
R	27.1	26.6	26.5	25.5	25.5	25.5
Ι	26.9	26.2	26.2	25.6	25.5	25.5
NB0816	26.0	25.4	25.4	-	-	-
Ζ	26.5	25.9	25.9	25.1	24.8	24.8
NB0921	26.0	25.3	25.4	-	-	-
у	25.6	24.7	24.7	-	-	-
Y	25.2	25.1	25.2	25.2	25.1	25.3
J	24.7	24.7	24.7	24.6	24.9	24.6
Y + J	25.3	25.3	25.3	25.3	25.3	25.2
Н	24.2	24.3	24.3	24.1	24.2	24.1
Ks	23.8	23.9	23.9	23.8	23.7	23.7
3.6 µm	24.3	24.3	24.4	25.0	24.9	24.9
4.5 µm	24.1	24.1	24.1	24.8	24.8	24.8

# 2.3 Image processing and depths

The astrometry of the VIDEO data is matched to Gaia (Gaia Collaboration 2018). If not already matched to Gaia, all other auxiliary data were shifted into this frame using SCAMP (Bertin 2006), and the pixel-scale was matched using SWARP (Bertin et al. 2002). The  $5\sigma$  depths were computed across the images to determine flux uncertainties. Circular apertures with a 2 arcsec diameter were placed on empty regions of the image. For the Spitzer/IRAC data we use a 2.8 arcsec diameter aperture to account for the poorer resolution. The SEGMENTATION map produced by SEXTRACTOR (Bertin & Arnouts 1996) was used to avoid foreground objects. To measure the standard deviation of the fluxes, the median absolute deviation (MAD) was used. This is more robust to outliers than Gaussian fitting, and is given by MAD = median(|flux - median|), with  $\sigma = 1.4826 \times MAD$ . Local depth maps in each image were produced by splitting the image into local regions with a size determined by taking the closest 300 apertures to each point. The global depths were taken as the mode of these local depths and are reported in Table 1.

# 2.4 Catalogue creation

At  $z \simeq 7$ , the Lyman break will cause objects to drop out of either the Z, y, or Y bands. The VIDEO Y band is deeper than the HSC y band (see Table 1), hence we produced fiducial catalogues by running SEXTRACTOR in dual image mode on a Y + J stacked image (parameters used are provided in Table 2). This probes the rest-frame UV (~1500 Å) at  $z \simeq 7$ . The photometry was performed in 2 arcsec diameter apertures, which encloses 70–80 per cent of the total flux of a point source. This provides high signal-to-noise ratio whilst balancing against the need for a large aperture correction. For the

 Table 2. Values of parameters used when running SEXTRACTOR in dual image mode.

Parameter	Value
DETECT_MINAREA	3
DETECT_THRESH	1.4
ANALYSIS_THRESH	1.4
DEBLEND_NTHRESH	32
DEBLEND_MINCONT	0.0001
BACK_SIZE	64
BACK_FILTERSIZE	9
BACKPHOTO_TYPE	LOCAL
BACKPHOTO_THICK	24

*Spitzer*/IRAC photometry, 2.8 arcsec diameter apertures were used to account for the broader PSF.

Since fluxes are measured in fixed apertures, the measurement must be corrected to a total flux by accounting for light beyond the aperture using the shape of the PSF. The FWHM of the *Y*-band imaging is 0.8 arcsec. However, as mentioned in Section 2.2, for ECDF-S the FWHM of the PSF varies markedly (by up to 0.1 arcsec) across the field for the *I* and *Z* bands. We thus used PSFEx (Bertin 2011) to generate a local PSF model to aperture correct our photometry by splitting each tile into a  $10 \times 10$  grid. For our analysis we only considered regions where VOICE, HSC, and VIDEO data overlap (see Fig. 1). Bright stars and regions of low signal-to-noise ratio at the edges of the VISTA VIDEO tiles are masked.

# **3 GALAXY SELECTION**

The selection of robust  $z \simeq 7$  candidates required several steps, which we detail in this section and present an overview in Table 3.

#### 3.1 Initial selection

The Y + J selected catalogues contained 1814513 and 1846027 objects in XMM-LSS and ECDF-S, respectively. Objects without overlapping VOICE + HSC + VIDEO imaging are removed, shown as the 'remove missing data' step in Table 3. Bright objects were selected by requiring detections at >5 $\sigma$  significance in Y + J for XMM-LSS. In the ECDF-S data, a 5 $\sigma$  cut was insufficient because Z is ~1 mag shallower than Y (see Table 1), making it challenging to select drop-out objects. We therefore applied a brighter 8 $\sigma$  cut in ECDF-S, corresponding to ~0.5 mag brighter than the 5 $\sigma$  cut in XMM-LSS.

We then require  $<2\sigma$  significance (i.e. a non-detection) in bands bluewards of and including *i*/*I*. In ECDF-S the HSC *R* band is not used for this due to its poor coverage and because *r* is deeper by 0.5 mag. Similarly, the *i* band is not used for this since HSC-*I* is  $\geq$ 0.7 mag deeper (see Table 1). The *u* band is also not used due to its significantly shallower depth. Note that all of these bands are later included for SED fitting in Section 3.2. The results of these steps on reducing the sample are shown in Table 3.

#### 3.2 SED fitting and photometric redshifts

We make use of all available bands to select LBGs at  $z \simeq 7$  by using a SED fitting analysis. SED fitting is done using LEPHARE (Arnouts et al. 1999; Ilbert et al. 2006), which works by minimising  $\chi^2$  to find the best-fitting redshift and galaxy templates. The redshift was allowed to vary between z = 0 and 9. Bruzual & Charlot (2003) stellar population models were used with metallicities of Z = [0.2, 0.4, **Table 3.** The results of our selection in XMM-LSS (top) and ECDF-S (bottom). The left-hand column describes the selection step and the right-hand column shows the number of objects remaining after this step. SED fitting is used to find candidates with photometric redshifts z > 6 and remove low-redshift interlopers, and visual checks are used to remove artefacts. The objects remaining after the removal of brown dwarfs are the final number of candidate galaxies in each field.

Selection step	Objects remaining
XN	IM-LSS
Initial catalogue	1814 513
Remove missing data	1707 431
$5\sigma Y + J$ cut	807 190
$< 2\sigma$ in I	11 821
$< 2\sigma$ in R	3398
$< 2\sigma$ in G	3217
SED fitting: $z > 6$	307
Visual selection	121
Dusty $z \sim 1-2$ galaxies	97
Brown dwarf removal	22
E	CDF-S
Initial catalogue	1846 027
Remove missing data	1291 653
$8\sigma Y + J$ cut	428 695
$< 2\sigma$ in I	2859
$< 2\sigma$ in r	1504
$< 2\sigma$ in g	1427
$< 2\sigma$ in G	1390
SED fitting: $z > 6$	186
Visual selection	70
Dusty $z \sim 1-2$ galaxies	61
Brown dwarf removal	6

1.0]  $Z_{\odot}$ . The star formation histories (SFHs) explored were constant, instantaneous bursts, and exponentially declining with time-scales in the range  $\tau = 0.05$ –10 Gyr. Stellar population ages could range from 10 Myr to 13.8 Gyr, limited by the age of the Universe at a given redshift. Dust reddening was prescribed by the Calzetti et al. (2000) dust law, with attenuation in the range  $0.0 \le A_V \le 4.0$  to allow for very dusty low-redshift interlopers. A Chabrier (2003) initial mass function is assumed, and absorption by the IGM was applied according to Madau (1995). *Spitzer/IRAC* photometry was not used initially to determine the photometric redshifts of objects due to the larger uncertainties and high rates of confusion due to a larger PSF. Instead, we included the IRAC photometry (when unconfused) for a separate check for low-redshift interlopers (see Section 3.4).

Lyman  $\alpha$  emission can increase the photometric redshift of objects by  $\Delta z \sim 0.5$  (Bowler et al. 2014) due to the addition of flux to the broad-band. This is accounted for by simultaneously fitting Bruzual & Charlot (2003) templates that have lines of equivalent widths (EWs)  $0 \leq EW_0 \leq 240$  Å added. To do this, the continuum level was measured from the mean flux between 1250 and 1300 Å. The photometric redshift constraints from this Lyman  $\alpha$  analysis are stronger in XMM-LSS due to many overlapping filters around the expected position of the Lyman break, and due to the availability of narrow-band filters that can very precisely pick out excess flux due to an emission line. It is highly degenerate in ECDF-S due to the gap between Z and Y.

Candidates were first required to have their best-fitting solution at z > 6. The fits then had to be sufficiently good, defined as  $\chi^2 < 11.3$  (12.9) for XMM-LSS (ECDF-S). These values correspond to  $2\sigma$  significance given 5 (6) degrees of freedom. We also required that the high-redshift solution was better than the low-redshift solution by a threshold  $\Delta \chi^2 > 4$ , which corresponds to a  $2\sigma$  significance. This removes 2910 objects in XMM-LSS and 1204 in ECDF-S, leaving 307 and 186 objects, respectively.

### 3.3 Visual selection

We conducted a visual selection of the remaining objects to remove artefacts that appear as good high-redshift candidates in the photometry, primarily in the form of cross-talk in VIDEO images and diffraction spikes. Cross-talk is an artefact produced in the readout from the Vista InfraRed CAMera (VIRCAM) instrument on VISTA that produces ghost images at multiples of 128 pixels in the native scale from bright stars (Bowler et al. 2017). Optical stacked postage stamp images are created to check for low-level optical flux that would indicate a low-redshift galaxy. The bands used for the optical stack are required to have similar  $5\sigma$  depths, leading to a *GRI* stack in XMM-LSS and a *grGI* stack in ECDF-S. Single-band detections are also removed. This step removes 186 objects in XMM-LSS and 116 in ECDF-S, leaving 121 and 70 objects, respectively.

Two objects in XMM-LSS (VIDEO\_z7\_8 and VIDEO\_z7\_12) appeared to have low-level flux in their optical stacks. This optical flux is offset from the position of the object in the Y + J detection filter by >1 arcsec, suggesting the optical flux originates from a foreground object and hence we retain them. One object (ID 1610530, RA 02:27:12.04, Dec. -04:32:05.05) appeared to be remarkably bright at  $M_{\rm UV} \simeq -24.4$ , but a marginal detection at 3.6  $\mu$ m around 2 mag fainter than in  $K_s$  and a non-detection at 4.5  $\mu$ m resulted in very unusual photometry that passed our selection criterion. For such a rest-UV bright high-redshift object, we would expect much brighter Spitzer/IRAC flux. We checked for this object in J-band imaging from the overlapping WIRCam Deep Survey (WIRDS, Bielby et al. 2012), where it was not present. However, a fainter object closely separated to the south is visible. This object is thus likely either a VIDEO artefact or a transient, supported by the faint 3.6 µm detection. The SED fitting and postage stamp images of this object are presented in Appendix C.

#### 3.4 Removing contaminants

Dusty galaxies at  $z \sim 1-2$  with Balmer breaks masquerading as Lyman breaks form one class of contaminant. These galaxies tend to have much redder continuum slopes than galaxies at  $z \simeq 7$ , as observed in the NIR. We use SED fitting with the *Spitzer/IRAC* bands to test for dusty low-redshift galaxies. If the SED fitting has a low-redshift solution that is preferred to the high-redshift solution when IRAC is included, the object is rejected if the IRAC photometry is unconfused. This step removes 24 objects in XMM-LSS and nine objects in ECDF-S, leaving 97 and 61 objects, respectively.

M, L, and T brown dwarf (BD) stars comprise a major contaminant in high-redshift searches around a magnitude  $m_{AB} \sim 25$  where their number density in extragalactic surveys peaks (Ryan et al. 2011). They have intrinsically red optical to NIR colours with heavy molecular spectral absorption complexes (e.g. Cushing et al. 2008; Marley et al. 2021). The removal of BDs must be done carefully since peaks in their SEDs can line up with the ground-based *YJHK*<sub>s</sub> filters, mimicking a flat NIR colour. We fit our photometry with stellar templates taken from the SpeX prism library (Burgasser 2014). The u, G, g, R, and r bands are excluded from the SED fitting since the templates do not contain any information at these wavelengths. The  $\chi^2_{BD}$  values are then compared to  $\chi^2_{galaxy}$  values, also fitted without these filters. In Bowler et al. (2015), a simulation of BDs in the Milky

Way shows that removing objects with good BD fits at  $\chi^2_{BD} < 10$ decreases the contamination rate to essentially zero. We therefore retain those high-redshift candidates that have  $\chi^2_{BD} > 10$ . This step removes around 80 per cent of objects (see Table 3) compared to around 20 per cent in Bowler et al. (2015). We use a similar model to Bowler et al. (2015) to compute the expected number of BDs in our fields, assuming a Galactic scale-height of 300 pc. Considering magnitudes down to the  $5\sigma$  (8 $\sigma$ ) depths in XMM-LSS (CDFS). we expect 480 and 220 dwarf stars, respectively. Restricting this to the most common stellar types found by our initial SED fitting, we expect 146 and 76 dwarf stars, respectively. We do not predict the exact number of dwarf stars because the Galactic scale-height is highly uncertain, with estimates varying between 300 (Ryan et al. 2017) and 400 pc (Sorahana, Nakajima & Matsuoka 2019). The differences can thus be attributed to a higher LBG number density at the fainter magnitudes and lower redshift probed by their work  $(-22.7 < M_{\rm UV} < -20.5 \text{ at } z = 6)$ , coupled with higher BD densities at brighter magnitudes probed by this work.

Bowler et al. (2015) state that keeping objects with  $\chi^2_{BD} > 10$  likely removes some genuine high-redshift galaxies, and they account for this with a completeness correction. We note that many of our candidates with  $\chi^2_{BD} < 10$  have a significantly better fitting high-redshift galaxy template. In an attempt to be complete to genuine  $z \simeq 7$  galaxies that may be removed from our primary robust sample due to being fit well as BDs, we conducted a more inclusive BD cut that retains these potential  $z \simeq 7$  galaxies, and is described in Appendix A.

#### 3.5 SED fitting with BAGPIPES

We repeat SED fitting on the final sample with Bayesian Analysis of Galaxies for Physical Inference and Parameter EStimation (BAG-PIPES; Carnall et al. 2018) to compare photometric redshifts with LEPHARE and determine physical parameters such as stellar masses. The advantage of BAGPIPES is the implementation of grids of nebular emission line models based on input stellar population models. We used a fiducial constant SFH model. The time since star formation switched on varies between the start of the Universe and the redshift being considered. We adopt a uniform prior on the redshift, 0 < z < z9, and the ionization parameter,  $-4 \le \log U \le -2$ . The metallicity is fixed at  $Z = 0.2 Z_{\odot}$ . Fixing of the metallicity at a value Z < ${
m Z}_{\odot}$  is motivated by recent spectroscopic measurements at  $z\gtrsim7$ that suggest galaxies tend to be metal poor (e.g. Langeroodi et al. 2022; Curtis-Lake et al. 2023; Fujimoto et al. 2023; Matthee et al. 2023). We also restrict the dust attenuation to the range  $0 \le A_V \le$ 0.5, limiting it to the largest attenuation we measure with LEPHARE (VIDEO\_z7\_25 and VIDEO\_z7\_26, shown in Table 5). If we do not restrict the range of  $A_V$ , large values occur for the objects with the most massive stellar masses:  $A_V \gtrsim 1$  when  $\log_{10}(M/M_{\odot}) \gtrsim 11$ . This is because dust extinction is degenerate with age (since we have fixed metallicity), so large values are a consequence of the constant SFH model rather than being a realistic estimate of the dust emission (e.g. Inami et al. 2022; Ferrara et al. 2023).

#### **4 THE FINAL SAMPLE**

The final sample consists of 28 objects, with photometry presented in Table 4. The results of SED fitting and properties of objects (e.g. photometric redshift,  $M_{\rm UV}$ , and dust attenuation) are shown in Table 5. Of the 22 primary candidates in XMM-LSS, 18 are newly identified. Five are within the ultradeep HSC pointing, and four have

**Table 4.** The photometry of sources in our final  $z \simeq 7$  sample. The top section shows objects in XMM-LSS, and the bottom section shows objects in ECDF-S. Objects are ordered by their photometric redshift. The first column shows the object ID, and the next two columns show the coordinates of the candidate. The remaining columns show the photometry in the bands available in each field. We require a  $<2\sigma$  detections bluewards of and including the *i/I* bands, so we only present bands redwards of this. The photometry is measured in a 2.0 arcsec diameter circular aperture apart from the *Spitzer/IRAC* bands where 2.8 arcsec diameter apertures are used to account for the broader PSF. The photometry is corrected to a total flux assuming a point-source correction. If the measured flux is  $<2\sigma$ , we present an upper limit. Objects with confused *Spitzer/IRAC* photometry are marked with an asterisk (\*).

ID	RA	Dec.	Ζ	NB921	у	Y	J	Н	Ks	[3.6]	[4.5]
VIDEO_z7_1*	02:22:56.06	-05:28:00.37	$26.3^{+0.5}_{-0.3}$	$24.9^{+0.2}_{-0.2}$	$25.3^{+0.6}_{-0.4}$	$24.7^{+0.2}_{-0.2}$	$24.4^{+0.2}_{-0.2}$	$24.9^{+0.7}_{-0.4}$	$23.7^{+0.2}_{-0.2}$	$22.2_{-0.1}^{+0.1}$	$22.0^{+0.1}_{-0.1}$
VIDEO_z7_2	02:25:51.03	-04:03:16.98	$26.3_{-0.4}^{+0.7}$	$25.5^{+0.4}_{-0.3}$	$25.1_{-0.4}^{+0.7}$	$24.5_{-0.2}^{+0.2}$	$24.5_{-0.2}^{+0.3}$	$24.3_{-0.3}^{+0.4}$	>24.1	$23.4_{-0.2}^{+0.2}$	$23.7_{-0.2}^{+0.3}$
VIDEO_z7_3	02:26:39.93	-04:01:09.51	$26.3_{-0.4}^{+0.5}$	$24.9^{+0.2}_{-0.1}$	$24.6_{-0.2}^{+0.3}$	$24.4_{-0.1}^{+0.2}$	$24.0^{+0.2}_{-0.2}$	$24.4_{-0.3}^{+0.4}$	>24.2	$23.3_{-0.2}^{+0.2}$	$23.8^{+0.2}_{-0.2}$
VIDEO_z7_4	02:27:38.62	-04:51:36.34	>26.2	$25.2^{+0.2}_{-0.2}$	>24.6	$24.7^{+0.2}_{-0.2}$	$25.0\substack{+0.5 \\ -0.4}$	$24.9_{-0.4}^{+0.7}$	>24.2	$23.5\substack{+0.2 \\ -0.2}$	$23.5^{+0.2}_{-0.2}$
VIDEO_z7_5*	02:26:30.02	-04:20:32.23	>26.2	$25.9^{+0.5}_{-0.3}$	$24.8_{-0.3}^{+0.4}$	$24.4_{-0.2}^{+0.2}$	$24.4_{-0.2}^{+0.3}$	$24.2_{-0.3}^{+0.4}$	$24.2_{-0.3}^{+0.5}$	$21.8^{+0.2}_{-0.2}$	$21.8^{+0.2}_{-0.2}$
VIDEO_z7_6	02:16:32.44	-05:30:05.82	>25.9	$24.9^{+0.2}_{-0.2}$	$24.6^{+0.2}_{-0.2}$	$24.7^{+0.2}_{-0.2}$	>24.6	>24.1	$24.3_{-0.4}^{+0.5}$	$24.5^{+0.3}_{-0.2}$	>25.1
VIDEO_z7_7	02:27:30.78	-04:25:12.11	$26.4_{-0.3}^{+0.4}$	$24.7^{+0.1}_{-0.1}$	$24.6_{-0.2}^{+0.3}$	$24.5_{-0.1}^{+0.1}$	$25.2^{+0.5}_{-0.3}$	$24.4_{-0.3}^{+0.4}$	>23.9	$23.8^{+0.3}_{-0.2}$	$24.2_{-0.4}^{+0.6}$
VIDEO_z7_8	02:22:05.84	-04:54:10.85	>26.0	>25.5	$25.0^{+0.6}_{-0.4}$	$24.6_{-0.2}^{+0.2}$	$24.8^{+0.4}_{-0.3}$	>24.4	$24.4_{-0.4}^{+0.5}$	$23.9^{+0.3}_{-0.2}$	>24.4
VIDEO_z7_9	02:23:47.30	-05:21:25.49	$25.4^{+0.2}_{-0.2}$	$25.0\substack{+0.2\\-0.2}$	$24.4_{-0.2}^{+0.2}$	$24.5_{-0.1}^{+0.1}$	$24.3_{-0.2}^{+0.2}$	$24.3_{-0.3}^{+0.3}$	$24.3_{-0.3}^{+0.5}$	$22.9^{+0.1}_{-0.1}$	$23.2^{+0.1}_{-0.1}$
VIDEO_z7_10*	02:16:36.51	-04:54:50.62	>26.8	>26.4	$24.9^{+0.2}_{-0.2}$	$24.8^{+0.3}_{-0.2}$	$24.6_{-0.2}^{+0.3}$	$24.7^{+0.6}_{-0.4}$	$24.3_{-0.4}^{+0.6}$	$22.1_{-0.2}^{+0.2}$	$22.2_{-0.2}^{+0.2}$
VIDEO_z7_11	02:24:59.99	-04:43:54.27	>26.6	>26.0	$24.7^{+0.3}_{-0.2}$	$24.7^{+0.3}_{-0.2}$	$24.9_{-0.3}^{+0.4}$	$24.7_{-0.3}^{+0.4}$	>24.2	$23.6\substack{+0.2\\-0.2}$	>24.3
VIDEO_z7_12*	02:17:00.07	-04:12:26.78	>26.6	>26.6	$25.1_{-0.2}^{+0.2}$	$24.8^{+0.2}_{-0.2}$	$25.3_{-0.4}^{+0.6}$	>24.7	>23.8	$23.6_{-0.2}^{+0.2}$	$24.4_{-0.2}^{+0.3}$
VIDEO_z7_13	02:25:17.59	-04:38:54.33	>26.2	>25.4	$24.8^{+0.4}_{-0.3}$	$24.7\substack{+0.2\\-0.1}$	$24.7^{+0.3}_{-0.2}$	$24.6_{-0.3}^{+0.4}$	$24.7\substack{+0.6 \\ -0.4}$	>24.0	$24.0\substack{+0.4\\-0.3}$
VIDEO_z7_14	02:27:37.26	-04:10:52.04	>26.6	>25.8	$25.2_{-0.3}^{+0.5}$	$24.8\substack{+0.2\\-0.2}$	$25.2\substack{+0.6 \\ -0.4}$	$25.0\substack{+0.7 \\ -0.4}$	$24.7\substack{+0.7 \\ -0.4}$	>24.6	>24.2
VIDEO_z7_15	02:21:29.47	-05:15:27.24	>26.5	>25.7	$25.5_{-0.4}^{+0.6}$	$25.2^{+0.3}_{-0.2}$	$25.3_{-0.3}^{+0.4}$	>24.5	>24.2	>24.1	>23.8
VIDEO_z7_16	02:22:27.88	-05:13:41.40	>26.5	>25.9	$25.5^{+0.7}_{-0.4}$	$24.9^{+0.2}_{-0.2}$	>25.1	$25.1\substack{+0.6 \\ -0.4}$	>24.6	>24.6	>24.3
VIDEO_z7_17	02:25:30.47	-05:18:01.46	>26.6	>25.5	$25.3_{-0.4}^{+0.7}$	$24.9\substack{+0.2 \\ -0.2}$	$25.2^{+0.4}_{-0.3}$	>24.6	>24.2	>24.0	>24.5
VIDEO_z7_18	02:21:16.56	-04:33:20.28	>26.7	>25.8	$25.2^{+0.3}_{-0.2}$	$24.4_{-0.1}^{+0.1}$	$24.6^{+0.2}_{-0.2}$	$24.3^{+0.3}_{-0.2}$	$24.6\substack{+0.6 \\ -0.4}$	$22.6\substack{+0.2 \\ -0.2}$	$23.1_{-0.2}^{+0.2}$
VIDEO_z7_19	02:18:06.79	-04:25:30.56	>27.1	>26.5	$25.3_{-0.2}^{+0.2}$	$24.7\substack{+0.2\\-0.2}$	$25.0\substack{+0.4\\-0.3}$	>24.7	>24.3	$24.9\substack{+0.4\\-0.3}$	>24.3
VIDEO_z7_20	02:15:31.44	-05:09:07.95	>26.6	>26.4	$25.1\substack{+0.3 \\ -0.2}$	$24.6\substack{+0.2 \\ -0.2}$	$24.7^{+0.4}_{-0.3}$	$24.6\substack{+0.7\\-0.4}$	>24.0	$24.1_{-0.2}^{+0.2}$	$25.0\substack{+0.5 \\ -0.3}$
VIDEO_z7_21*	02:26:46.18	-04:59:53.51	>26.4	>25.7	>24.5	$24.6_{-0.1}^{+0.2}$	$24.5_{-0.2}^{+0.3}$	$24.5_{-0.3}^{+0.4}$	$24.4_{-0.4}^{+0.6}$	$24.1_{-0.3}^{+0.4}$	$23.4_{-0.2}^{+0.2}$
VIDEO_z7_22*	02:16:25.10	-04:57:38.56	>27.2	>26.2	$26.5^{+0.7}_{-0.4}$	$24.7\substack{+0.2 \\ -0.2}$	$24.1^{+0.2}_{-0.1}$	$24.1_{-0.2}^{+0.2}$	$23.6\substack{+0.3 \\ -0.2}$	$23.1\substack{+0.2 \\ -0.2}$	$22.9\substack{+0.2 \\ -0.2}$
						• · • +0 1	• · · +0 2	• · •+0.4	a a =±0 3		•• <b>-</b> +0 2
VIDEO_z7_23*		-28:14:20.71	>25.1	-	-	$24.3^{+0.1}_{-0.1}$	$24.4^{+0.2}_{-0.2}$	$24.5^{+0.4}_{-0.3}$	$23.7^{+0.3}_{-0.2}$	$22.7^{+0.2}_{-0.2}$	$22.5^{+0.2}_{-0.2}$
VIDEO_z7_24*		-27:49:32.15	>25.3	-	-	$24.3^{+0.1}_{-0.1}$	$24.6^{+0.3}_{-0.2}$	$24.9^{+0.7}_{-0.4}$	>24.0	>24.8	$24.8^{+0.5}_{-0.3}$
VIDEO_z7_25		-28:02:51.72	>25.5	-	-	$24.2^{+0.1}_{-0.1}$	$23.7^{+0.1}_{-0.1}$	$23.2^{+0.1}_{-0.1}$	$22.8^{+0.1}_{-0.1}$	$22.5^{+0.2}_{-0.2}$	$22.4^{+0.2}_{-0.2}$
VIDEO_z7_26		-27:59:28.38	>25.5	-	-	$24.5^{+0.1}_{-0.1}$	$24.2^{+0.2}_{-0.1}$	$23.5^{+0.1}_{-0.1}$	$23.1^{+0.1}_{-0.1}$	$22.4^{+0.2}_{-0.2}$	$22.1^{+0.2}_{-0.2}$
VIDEO_z7_27		-27:17:30.39	>25.6	-	-	$24.2^{+0.1}_{-0.1}$	$24.3^{+0.2}_{-0.2}$	$24.7^{+0.6}_{-0.4}$	$23.4^{+0.2}_{-0.2}$	$23.8^{+0.2}_{-0.2}$	$24.3^{+0.2}_{-0.2}$
VIDEO_z7_28	03:34:29.33	-28:13:00.47	>25.1	-	-	$24.1^{+0.1}_{-0.1}$	$23.5_{-0.1}^{+0.1}$	$24.4_{-0.3}^{+0.5}$	$24.2^{+0.4}_{-0.3}$	$23.7^{+0.2}_{-0.2}$	$22.4_{-0.2}^{+0.2}$

a best-fitting Lyman  $\alpha$  template (and have  $z_{No Ly\alpha} < 6.5$ ). The bestfitting rest-frame EWs for these objects are noted in Table 5. In ECDF-S, all candidates are presented here for the first time. One object in ECDF-S had a best-fitting Lyman  $\alpha$  template, but was rejected due to having a good low-redshift fit. Therefore, all Lyman  $\alpha$ objects lie in XMM-LSS. The absolute UV magnitudes,  $M_{\rm UV}$ , are computed by placing a top-hat filter at 1500 Å with width 100 Å on the rest-frame best-fitting SED. The sample has a mean redshift  $\bar{z} = 6.74$ , and spans nearly 2 mag in the rest-UV,  $-23.5 \le M_{\rm UV} \le$ -21.6. The distribution of photometric redshifts against UV absolute magnitude, split by field, is shown in Fig. 2. Candidates in ECDF-S are brighter on average than those in XMM-LSS simply by nature of the different selection thresholds that we imposed because of the shallower optical data available in this field. Uncertainties on the photometric redshifts tend to be lower in XMM-LSS compared to ECDF-S due to the proximity and overlap between Z, NB921, y, and Y around the Lyman break. The grey points in Fig. 2 are candidate galaxies from Bouwens et al. (2021), which use 1136 arcmin<sup>2</sup> of *Hubble* imaging to derive the rest-UV LF at  $z \simeq 7$ . Their brightest candidate is  $M_{\rm UV} \simeq -22.1$ . Only 10 of our candidates are fainter than this, showing that these ground-based data better probe the bright end of the rest-UV LF.

The SED fitting and stamps of a candidate galaxy from each field are shown in Fig. 3. The SED fitting and stamps for all candidates in this work are shown in Appendix B. The [3.6] and [4.5] Spitzer/IRAC bands can be contaminated by nebular emission lines at z > 6.5 (primarily H $\alpha$ , H $\beta$ , [OII], and [OIII]). The [3.6] - [4.5] colour will accordingly change with redshift as these lines move through the filters (Smit et al. 2014, 2015). By accounting for nebular emission during SED fitting, we can obtain more precise photometric redshifts. An interesting example is the brightest object in our sample, VIDEO\_z7\_28, which has a red IRAC colour of [3.6]  $-[4.5] = 1.3 \pm 0.4$  and a photometric redshift of  $z = 7.38^{+0.03}_{-0.06}$ . Bowler et al. (2020) extend the expected IRAC colours derived by Smit et al. (2014) out to higher redshifts, predicting that galaxies with strong H $\beta$  + [O III] emission (EW<sub>0</sub> = 500–2000 Å) can have colours as high as  $[3.6] - [4.5] \simeq 1.0$  at  $z \simeq 7.5$ , in agreement with our photometry for this object. As we discuss later in Section 4.1,

Table 5.       The 28 LBG candidates in XMM-LSS (top) and ECDF-S (bottom), ordered by photometric redshift as in Table 4. The first column shows the object
ID. The next four columns show properties of the high-redshift solution: photometric redshift, absolute magnitude in a top-hat filter at 1500 Å with width 100 Å,
dust attenuation, and $\chi^2$ value as derived by LEPHARE. The next two columns show the secondary photometric redshift and $\chi^2$ . The following two columns
show the stellar type for fitting to standard MLT brown dwarf spectra and $\chi^2_{BD}$ . The next two columns show the equivalent width (EW) for objects that have a
best-fitting SED with Lyman $\alpha$ emission. The final three columns show properties derived from SED fitting with BAGPIPES, namely the photometric redshift z,
dust attenuation, and stellar mass. Objects with IRAC non-detections have upper limits for their stellar masses. Object IDs marked with an asterisk (*) have
confused IRAC photometry, meaning stellar mass estimates may be overestimated. Object IDs marked with superscripts 1 and 2 have been previously identified
in Bouwens et al. (2021) and Endsley et al. (2021), respectively.

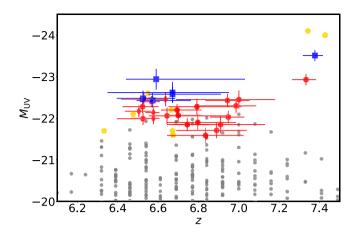
ID	z	$M_{ m UV}$	$A_V$	$\chi^2$	Zgal2	$\chi^2_{gal2}$	Stellar type	$\chi^2_{BD}$	EW <sub>0</sub>	z	$A_V$ (mag)	$\log_{10}(M_{\star}/M_{\odot})$
		(mag)	(mag)						(Å)	(BAGPIPES)	(BAGPIPES)	(BAGPIPES)
VIDEO_z7_1*	$6.50\substack{+0.07\\-0.06}$	$-22.2\pm0.1$	0.2	11.9	1.3	18.4	M8	18.2	40	$6.43\substack{+0.05 \\ -0.06}$	$0.41\substack{+0.07\\-0.13}$	$10.31\substack{+0.17 \\ -0.32}$
VIDEO_z7_2	$6.52\substack{+0.08 \\ -0.10}$	$-22.3\pm0.2$	0.0	11.3	1.2	36.8	Т3	15.8		$6.53\substack{+0.03 \\ -0.04}$	$0.32\substack{+0.12 \\ -0.15}$	$10.20\substack{+0.15 \\ -0.18}$
VIDEO_z7_3	$6.53\substack{+0.02 \\ -0.02}$	$-22.5\pm0.1$	0.2	5.4	1.3	42.8	M8	27.7		$6.49_{-0.05}^{+0.03}$	$0.27^{+0.15}_{-0.17}$	$10.17\substack{+0.24 \\ -0.33}$
VIDEO_z7_4	$6.53\substack{+0.03\\-0.13}$	$-22.0\pm0.2$	0.2	7.1	1.3	31.6	Т3	23.2		$6.48\substack{+0.04 \\ -0.06}$	$0.42^{+0.06}_{-0.10}$	$10.28\substack{+0.11\\-0.15}$
VIDEO_z7_5	$6.57\substack{+0.04 \\ -0.06}$	$-22.4\pm0.1$	0.0	6.3	1.5	21.9	Т3	12.5		$6.57_{-0.02}^{+0.25}$	$0.32\substack{+0.14 \\ -0.18}$	$10.24_{-0.28}^{+0.21}$
VIDEO_z7_61	$6.58\substack{+0.03\\-0.12}$	$-22.0\pm0.1$	0.0	7.9	1.3	46.2	M8	31.0	50	$6.46\substack{+0.04\\-0.07}$	$0.05\substack{+0.07\\-0.03}$	$9.12_{-0.26}^{+0.29}$
VIDEO_z7_7	$6.58\substack{+0.02 \\ -0.04}$	$-22.1\pm0.1$	0.0	8.0	1.3	74.9	M8	65.5	50	$6.44_{-0.04}^{+0.03}$	$0.17_{-0.12}^{+0.17}$	$10.03\substack{+0.22\\-0.37}$
VIDEO_z7_8	$6.64\substack{+0.31 \\ -0.08}$	$-22.1\pm0.2$	0.0	4.4	1.45	16.5	M6	14.2		$6.73_{-0.13}^{+0.18}$	$0.22\substack{+0.18 \\ -0.15}$	$10.04_{-0.34}^{+0.23}$
VIDEO_z7_9	$6.64^{+0.01}_{-0.22}$	$-22.5\pm0.2$	0.0	4.1	1.25	23.5	M8	17.7	240	$6.30\substack{+0.05\\-0.05}$	$0.47^{+0.02}_{-0.04}$	$10.55\substack{+0.05\\-0.08}$
VIDEO_z7_102*	$6.69^{+0.12}_{-0.07}$	$-22.2\pm0.2$	0.2	3.4	1.35	40.4	T4	25.2		$6.74\substack{+0.10 \\ -0.10}$	$0.46\substack{+0.03\\-0.07}$	$10.51\substack{+0.09 \\ -0.13}$
VIDEO_z7_11	$6.70\substack{+0.12\\-0.06}$	$-22.1\pm0.2$	0.0	2.7	1.5	26.9	T8	20.4		$6.70\substack{+0.11 \\ -0.07}$	$0.18\substack{+0.17 \\ -0.13}$	$9.93_{-0.34}^{+0.25}$
VIDEO_z7_12*	$6.74_{-0.10}^{+0.12}$	$-21.8\pm0.2$	0.0	8.6	1.5	30.0	M6	33.9		$6.68\substack{+0.10\\-0.06}$	$0.23^{+0.16}_{-0.13}$	$9.94^{+0.16}_{-0.18}$
VIDEO_z7_13	$6.79\substack{+0.22\\-0.14}$	$-22.3\pm0.2$	0.2	8.6	1.5	54.0	T8	25.9		$6.83\substack{+0.15 \\ -0.17}$	$0.13\substack{+0.16 \\ -0.09}$	$9.80_{-0.34}^{+0.24}$
VIDEO_z7_14	$6.80\substack{+0.19 \\ -0.15}$	$-21.9\pm0.2$	0.0	7.8	1.5	28.2	T8	18.9		$6.79_{-0.13}^{+0.13}$	$0.25\substack{+0.17 \\ -0.17}$	<9.90
VIDEO_z7_15	$6.83\substack{+0.19 \\ -0.18}$	$-21.6\pm0.2$	0.0	4.7	1.5	24.4	T8	12.7		$6.77_{-0.13}^{+0.19}$	$0.18\substack{+0.20 \\ -0.14}$	<9.70
VIDEO_z7_16	$6.89\substack{+0.15 \\ -0.20}$	$-21.7\pm0.2$	0.0	8.0	1.5	31.8	T8	23.5		$6.79_{-0.14}^{+0.14}$	$0.10\substack{+0.13 \\ -0.07}$	< 9.32
VIDEO_z7_17	$6.91\substack{+0.22 \\ -0.24}$	$-21.8\pm0.2$	0.0	7.8	1.15	35.1	T8	16.9		$6.79\substack{+0.15 \\ -0.12}$	$0.12\substack{+0.13 \\ -0.09}$	<9.35
VIDEO_z7_18	$6.94\substack{+0.12 \\ -0.10}$	$-22.4\pm0.2$	0.0	6.3	1.5	45.4	T3	23.6		$6.86\substack{+0.11\\-0.11}$	$0.44\substack{+0.04\\-0.09}$	$10.52\substack{+0.10 \\ -0.14}$
VIDEO_z7_19	$6.95\substack{+0.08 \\ -0.09}$	$-22.0\pm0.2$	0.0	9.4	1.25	88.2	T8	35.9		$6.76\substack{+0.11 \\ -0.09}$	$0.08\substack{+0.11 \\ -0.05}$	$9.40^{+0.28}_{-0.32}$
VIDEO_z7_20	$6.99\substack{+0.09\\-0.18}$	$-22.3\pm0.2$	0.0	3.7	1.45	38.7	Т8	16.0		$6.78\substack{+0.10 \\ -0.11}$	$0.14\substack{+0.15 \\ -0.09}$	$9.55_{-0.35}^{+0.24}$
VIDEO_z7_21 <sup>1*</sup>	$7.00\substack{+0.18 \\ -0.31}$	$-22.5\pm0.2$	0.1	4.5	1.5	41.8	T8	17.5		$7.00\substack{+0.10 \\ -0.15}$	$0.25\substack{+0.14 \\ -0.14}$	$10.18\substack{+0.14 \\ -0.19}$
VIDEO_z7_221*	$7.33\substack{+0.06 \\ -0.07}$	$-22.9\pm0.1$	0.0	11.4	1.6	33.6	T1	14.5		$7.29\substack{+0.06\\-0.06}$	$0.40\substack{+0.07\\-0.10}$	$10.59_{-0.13}^{+0.11}$
VIDEO_z7_23*	$6.52_{-0.13}^{+0.16}$	$-22.5 \pm 0.2$	0.0	6.7	1.45	12.2	M6	14.7		$6.49_{-0.12}^{+0.19}$	$0.45\substack{+0.03 \\ -0.07}$	$10.61^{+0.09}_{-0.13}$
VIDEO_z7_24*	$6.57^{+0.2}_{-0.08}$	$-22.4 \pm 0.2$	0.0	8.3	1.45	17.7	M7	19.7		$6.71^{+0.25}_{-0.21}$	$0.04^{+0.06}_{-0.03}$	$9.18^{+0.22}_{-0.20}$
VIDEO_z7_25	$6.59^{+0.44}_{-0.15}$	$-22.9\pm0.3$	0.5	4.0	1.25	14.6	L8	13.5		$7.31_{-0.05}^{+0.05}$	$0.46^{+0.03}_{-0.06}$	$10.92^{+0.08}_{-0.14}$
VIDEO_z7_26	$6.67^{+0.28}_{-0.32}$	$-22.6 \pm 0.3$	0.5	4.0	1.5	8.7	L8	14.7		$7.30^{+0.05}_{-0.06}$	$0.48^{+0.02}_{-0.03}$	$10.86^{+0.06}_{-0.09}$
VIDEO_z7_27	$6.67^{+0.24}_{-0.12}$	$-22.6 \pm 0.2$	0.0	5.6	1.45	29.4	M8	29.7		$6.79^{+0.15}_{-0.12}$	$0.20^{+0.18}_{-0.14}$	$9.70^{+0.21}_{-0.31}$
VIDEO_z7_28	$7.38^{+0.03}_{-0.06}$	$-23.5\pm0.1$	0.0	9.9	1.5	77.7	T8	11.1		$7.24^{+0.06}_{-0.09}$	$0.09^{+0.14}_{-0.07}$	$9.93^{+0.26}_{-0.50}$

we show that several previous  $z \sim 7$  candidates at around z = 7.4 are likely to be BDs. VIDEO\_z7\_28 passes our selection criterion with the BD fit showing  $\chi^2_{BD} = 11.1$ . It is slightly extended in the data with a FWHM of 1.5 arcsec in comparison to the PSF of the *Y* band. However, we cannot fully rule out a BD solution for this source, and further follow-up will be required to confirm this extremely luminous LBG.

## 4.1 Comparison to previous work

The Reionization Era Bright Emission Line Survey (REBELS; Bouwens et al. 2022) is a search for bright z > 6.5 galaxies in XMM-LSS and Cosmological Evolution Survey (COSMOS), with the aim of following up these galaxies with the ALMA. Within XMM-LSS they make use of VIDEO (Jarvis et al. 2013), the UKIRT Infrared Deep Sky Survey (UKIDSS) Ultra-Deep Survey (UDS; Lawrence et al. 2007), and optical data from HSC-SSP DR2 and the Canada– France–Hawaii Telescope Legacy Survey (CFHTLS) D1 field (Erben et al. 2009). Of the 40 REBELS objects, 15 are in XMM-LSS and we recover three out of four that are bright enough to meet our  $5\sigma Y + J$  cut, namely REBELS-01 (VIDEO\_z7\_22), REBELS-02 (VIDEO\_z7\_6), and REBELS-14 (VIDEO\_z7\_21). REBELS-10 is not recovered due to blending with a foreground object. VIDEO\_z7\_21 and VIDEO\_z7\_22 are spectroscopically confirmed at z = 7.084 and 7.177, respectively, in agreement with our photometric redshifts. The difference in objects selected is due to deeper UDS data used by REBELS. By making use of the UDS data they are able to select fainter candidates, and their selection also identifies candidates at z > 7.5.

Great Optically Luminous Dropout Research Using HSC (GOL-DRUSH; Harikane et al. 2022) uses HSC-SSP DR2 data (*GRIZy*) to search for LBGs at z = 4-7. This survey covers 20.2 deg<sup>2</sup> to



**Figure 2.** Rest-frame absolute UV magnitudes of the 28 LBGs in our sample against their photometric redshifts as measured by LEPHARE. Candidates from XMM-LSS are shown as red circles, and those with a Lyman  $\alpha$  best-fitting SED are shown as red diamonds. Candidates from ECDF-S are shown as blue squares. The grey circles are candidate galaxies from Bouwens et al. (2021), which use *Hubble* data. The yellow pentagons show the high-redshift solution of SED fitting on objects from Harikane et al. (2022) cross-matched with our catalogues.

depths comparable to the 'deep' and 'ultradeep' regions of XMM-LSS outlined in Section 2.1. There are eight objects that overlap with our XMM-LSS data, all in the 'ultradeep' HSC pointing. Three out of eight of these objects are recovered in our  $5\sigma Y + J$  cut. We find that the two brightest objects recovered at  $5\sigma$  in Y + J have preferred BD fits, with poor galaxy fits. The SED fits for these objects are shown in Fig. 4. These objects have coordinates (RA 02:18:11.39, Dec. -05:29:56.55) and (RA 02:17:00.95, Dec. -04:40:35.14). The third object (RA 02:19:40.78, Dec. -05:13:53.76)<sup>2</sup> recovered at  $5\sigma$  in Y + J appears to have low-level optical flux, best seen in a GRI stack, due to confusion with a nearby foreground object. Our SED fitting is also used on the remaining five objects that do not satisfy the  $5\sigma Y + J$  cut. Of these, two objects have  $z_{phot} < 6.5$ , and three objects have  $z_{phot} >$ 6.5. This appears to give a fiducial contamination rate of  $\sim$  50 per cent, but because we find BD fits for their brightest objects, it may be the case that their brightest bins are more significantly contaminated. We show the photometric redshifts of the  $z \ge 6$  solutions of SED fitting and resulting absolute UV magnitudes  $M_{\rm UV}$  for these objects in Fig. 2. The high-redshift solution for the two BDs places them at  $z \simeq 7.4$  with extremely bright absolute UV magnitudes of  $M_{\rm UV}$  $\simeq -24$ . The remainder of the objects are fainter, similar to the bulk of our sample. Contamination of the Harikane et al. (2022) sample, and the impact on the derived bright end of the UV LF, is discussed further in Section 6.1.

Endsley et al. (2021) select galaxies at z = 6.63-6.83 using overlapping data from HSC-SSP DR2, VIDEO, and the UDS. They search in XMM1 and COSMOS, finding nine objects in XMM1. XMM1-313310 in their study corresponds to VIDEO\_z7\_10 in this work. We find  $z_{phot} = 6.72$ , agreeing with their redshift range. XMM1-418672 in their study exists in our initial catalogue, but has  $\chi^2 = 15.3$  so is rejected since it fails the  $\chi^2 < 11.3$  cut in XMM1-LSS. The other seven objects were selected by Endsley et al. (2021) in the deeper UDS data, so are not recovered in this work.

<sup>2</sup>These have IDs 37484833881991283, 37485130234754131, and 37484567594038420, respectively, in the Harikane et al. (2022) catalogue.

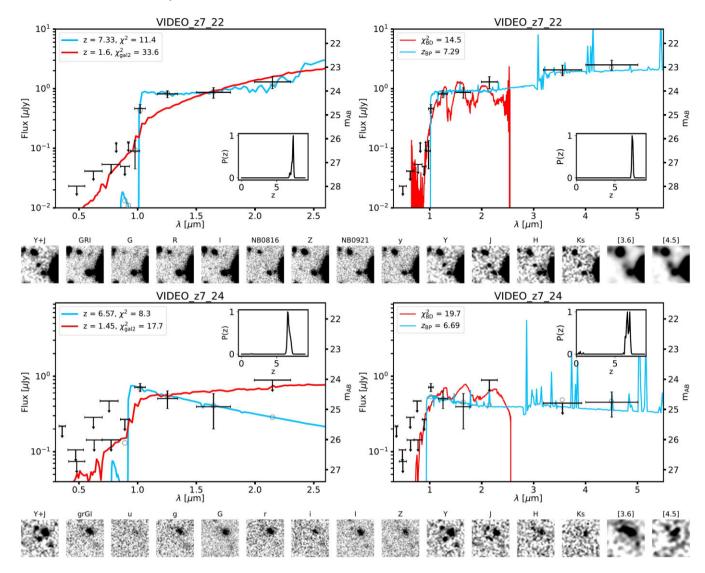
#### 4.2 Physical properties with BAGPIPES

Table 5 shows the results of our BAGPIPES fitting. All photometric redshifts measured by BAGPIPES agree with those measured by LEPHARE apart from objects VIDEO\_z7\_25 and VIDEO\_z7\_26 in ECDF-S, which agree within  $2\sigma$ , due to degeneracies in bestfitting redshift because of the gap between the Z and Y bands. IRAC detections probe the rest-frame optical, providing improved estimates of stellar masses. Measured stellar masses of the galaxies with unconfused Spitzer/IRAC photometry are in the range 9.1 < $\log_{10}(M_{\star}/M_{\odot}) < 10.9$ , with a mean stellar mass of  $\log_{10}(M_{\star}/M_{\odot}) =$ 10.0. Our stellar mass for VIDEO\_z7\_22 agrees with that found by Bouwens et al. (2022), but our stellar mass for VIDEO\_z7\_21 is an order of magnitude higher than theirs. This is because Spitzer/IRAC bands for this object are confused, leading to boosted flux due to the foreground object to the north. Bouwens et al. (2022) apply neighbour subtraction to their IRAC photometry, which we do not do. Consequently, we mark objects that are confused in their Spitzer/IRAC photometry with an asterisk in Table 5, indicating that these objects may have overestimated stellar masses. We also show confused IRAC photometry in green and with dashed error bars in the SED fitting plots for all candidates in Appendix B. Nearly half (10/23) of our galaxies with unconfused IRAC photometry have stellar masses  $\log_{10}(M_{\star}/M_{\odot}) > 10$ , suggesting that these galaxies represent some of the most massive galaxies at this epoch (compared to e.g. Bowler et al. 2014; Bouwens et al. 2022; Labbé et al. 2023). The lower end of our derived stellar masses is more uncertain because the rate of Spitzer/IRAC non-detections increases for  $\log_{10}(M_{\star}/M_{\odot}) \leq 9.4$ , meaning we cannot directly measure the rest-frame optical emission. Our masses are therefore upper limits.

Further issues arise in estimating stellar masses at these redshifts since it is difficult to tell if there is nebular emission (primarily  $[O III] + H\beta$  contaminating the IRAC photometry, outshining the stellar continuum and masking contributions from the older stellar population. Since we have fixed metallicity, this results in a degeneracy between age and dust attenuation. Whitler et al. (2023) have shown that properties derived from SED fitting depend on the SFH used. The dependence is strongest for derived age and weakest for stellar mass. This can be somewhat remedied by restricting the redshift range and using narrow- and intermediate-band photometry to compute more precise photometric redshifts (Endsley et al. 2021), limiting nebular emission to only one of the two Spitzer/IRAC bands and leaving the other free from contamination. Despite this contamination of the IRAC bands, we find that BAGPIPES provides reasonable estimates of the stellar masses in comparison to those derived by excluding contaminated bands entirely (e.g. Bowler et al. 2018).

#### **5 THE REST-FRAME UV LUMINOSITY FUNCTION**

We use the LBG candidates found in our selection to determine the UV LF at  $z \simeq 7$ . The  $1/V_{max}$  method is used to compute points on the LF (Schmidt 1968), where  $V_{max}$  is the maximum comoving volume a galaxy can occupy and still be included in our survey. The maximum redshift  $z_{max}$  is determined by redshifting the galaxy in steps of  $\Delta z = 0.01$  until it no longer satisfies the  $5\sigma$  ( $8\sigma$ ) detection threshold in Y + J in XMM-LSS (ECDF-S).  $V_{max}$  is then the comoving volume between z = 6.5 and  $z = z_{max}$ . An upper limit of z = 7.5, the maximum considered redshift, is placed on  $z_{max}$ . The UV LF is then



**Figure 3.** Two example candidate galaxies in XMM-LSS (top object, VIDEO\_z7\_22) and ECDF-S (bottom object, VIDEO\_z7\_24). For each object, the black points in both plots are the measured photometry, with non-detections replaced by  $2\sigma$  upper limits. Top-left: SED fitting with LEPHARE. The blue curve shows the best-fitting high-redshift solution, and grey circles are its model photometry. The red curve shows the low-redshift solution. The legend shows the redshift and  $\chi^2$  value of each solution. The wavelength range of this plot is limited to  $\lambda \simeq 2.5 \,\mu$ m to focus on the optical and NIR photometry around the Lyman break. The inset black curve shows the redshift probability distribution. Top-right: best-fitting BD template (red) along with the best-fitting BAGPIPES galaxy fit. The  $\chi^2$  value for the BD fit,  $\chi^2_{BD}$ , is computed without IRAC and without bands bluewards of *ill*. The grey circles show the model photometry of the BAGPIPES solution. The legend shows  $\chi^2_{BD}$  and the redshift of the BAGPIPES solution,  $z_{BP}$ . The wavelength range of this plot extends out to  $\lambda > 5 \,\mu$ m to show the *Spitzer/*IRAC photometry and the nebular emission component of the BAGPIPES fit. The inset black curve is the redshift probability found by BAGPIPES. Bottom:  $10 \times 10 \, \text{arcsec}^2$  postage stamps of the object in the filters used.

calculated as

$$\Phi(M)d\log(M) = \frac{1}{\Delta M} \sum_{i}^{N} \frac{1}{C(M_i, z_i)} \frac{1}{V_{\max, i}},\tag{1}$$

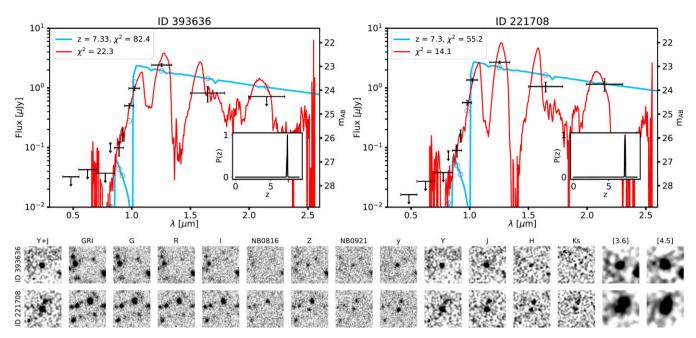
where  $\Delta M$  is the width of the magnitude bin and  $C(M_i, z_i)$  is a completeness correction, which depends on the magnitude and redshift of the object. Poissonian errors are given by

$$\delta \Phi(M) = \frac{1}{\Delta M} \sqrt{\sum_{i}^{N} \left(\frac{1}{V_{\max,i}}\right)^2}$$
(2)

following Adams et al. (2023). Bin widths  $\Delta M_{\rm UV} = 0.75$  are chosen for all bins to maximize the number of galaxies in each bin. The bins are centred at  $M_{\rm UV} = -22.175, -22.925,$  and -23.675. These bins contain 21, four, and one galaxies, respectively. The brightest bin only contains VIDEO\_z7\_28, so the binning is also chosen to ensure this object is roughly in the centre of the bin. The faintest bin was chosen to ensure we are significantly above where our sample becomes incomplete, as discussed in Section 5.1. We compute an upper limit on the number density of galaxies at  $M_{\rm UV} = -24.425$ with bin width  $\Delta M = 0.75$ . This is calculated by  $1/V/\Delta M$ , where V is the volume probed by our survey at  $6.5 \le z \le 7.5$ .

#### 5.1 Completeness correction

The usable area after accounting for foreground objects and artefacts is 83 per cent in both fields. This factor is incorporated into the calculation of  $V_{\text{max}}$ .



**Figure 4.** Top: SED fitting of the two brightest GOLDRUSH objects in XMM-LSS that overlap with our selection, ID 393636 (left) and ID 221708 (right). The black points are the measured photometry, with non-detections replaced by  $2\sigma$  upper limits. The dwarf star SED is shown in red, and the high-redshift solution in blue. The grey circles are the model photometry for the high-redshift solution. Inset is the redshift probability distribution for the galaxy fit. The legend shows the redshift of the galaxy template and  $\chi^2$  values. Both objects have preferred BD fits and poor galaxy fits: the rising slope in the blue end of the BD SED and the drop in flux at 1.6 µm are captured by the optical and *YJHK*s bands, respectively. Bottom: 10 × 10 arcsec<sup>2</sup> postage stamps each object. From left to right, we have the detection filter *Y* + *J*, the optical stack used to check for low-level flux in each field, and the filters available for objects in each field ordered by increasing wavelength, as in Fig. 3.

The focus of this work is the bright end of the rest-frame UV LF. We expect that the brightest bins are close to complete (excluding area lost due to foreground objects) in a  $5\sigma Y + J$  selection since they are much brighter than the detection threshold (>8 $\sigma$  at  $M_{\rm UV} \leq -22.4$ ). We estimate the completeness of our initial selection (see Section 3.1) by comparing the number of objects per deg<sup>2</sup> as a function of apparent magnitude in Y + J in XMM-LSS and ECDF-S with the 'ultradeep' stripes from the UltraVISTA Survey (McCracken et al. 2012) in COSMOS, reaching  $5\sigma$  depths of 26.3 in Y + J in the ultradeep stripes,  $\sim 1$  mag deeper than VIDEO. We find, as expected, that our Y + J selection is 95 per cent complete by  $M_{\rm UV} = -21.8$ , so we choose our faintest bin to cut-off at this absolute UV magnitude. This excludes two candidates from the LF calculation, VIDEO\_z7\_15 and VIDEO\_z7\_16.

SED fitting is in general more complete than a colour-colour selection (e.g. Adams et al. 2020), motivating its use in this work. To ensure our SED fitting selection is not overly conservative, we test the completeness of our SED fitting by estimating the fraction of genuine high-redshift galaxies incorrectly cut by our selection as a function of absolute UV magnitude and redshift. The values of  $M_{\rm UV}$ used reache 1 mag below the  $5\sigma$  limit of each of our fields to account for up-scattering of objects into the sample. We use BAGPIPES to generate z = 6-8 SEDs with ages varying between 50 and 700 Myr (or the age of the Universe at a given redshift). We vary the dust attenuation between 0 and 0.5 mag and fix the metallicity at 0.2  $Z_{\odot}$ . We generate 26000 galaxies for ECDF-S and each of XMM-LSS 'deep' and 'ultradeep'. Mock photometry in the available bands is generated by perturbing the model photometry using local depths within each field. SED fitting is performed on a catalogue of mock objects, where we apply the  $\chi^2$  cuts outlined in Sections 3.2 and 3.4.

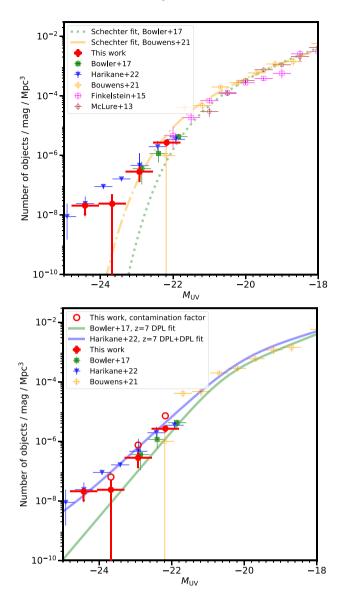
**Table 6.** The rest-UV LF points at  $z \simeq 7$  derived from the sample presented in this work. The first two columns show the absolute UV magnitude of the bin and the number of galaxies in the bin. All bins have a width of 0.75 mag. The third column shows the comoving number density of galaxies. The final column shows the median completeness in the bin.

Bin (mag)	$n_{\rm gal}$	$\phi (\mathrm{mag}^{-1}\mathrm{Mpc}^{-3})$	Completeness
-22.175	21	$2.70 \pm 0.66 \times 10^{-6}$	0.63
-22.925	4	$2.81 \pm 1.54  imes 10^{-7}$	0.76
-23.675	1	$2.37 \pm 2.50 \times 10^{-8}$	0.79
-24.425	0	$<\!2.07 \times 10^{-8}$	-

The completeness is then the fraction of galaxies recovered in each magnitude–redshift bin. We report the median completeness values for each LF bin in Table 6.

# 5.2 Cosmic variance

Galaxy surveys are a discrete sampling of the large-scale structure of the Universe. Substructures such as voids and filaments can thus impact measurements of the LF if the dimensions of the survey are small. This effect has come to be known as 'cosmic variance'. We make use of the Trenti & Stiavelli (2008) calculator to estimate additional uncertainties due to this. We find that the Poisson errors dominate: cosmic variance contributes just 7.1 per cent to the brightest bin at  $M_{\rm UV} \simeq -23.7$ , going down to a 5.7 per cent contribution to our faintest bin at  $M_{\rm UV} \simeq -22.2$ . We include these contributions in our LF errors by adding them in quadrature.



**Figure 5.** The rest-frame UV LF at  $z \simeq 7$  derived from 8.2 deg<sup>2</sup> of VIDEO data. The red filled circles in both panels show measurements of the number density from this work. We compute an upper limit on the number density of galaxies at  $M_{\rm UV} = -24.425$  as described in Section 5. Other measurements from Harikane et al. (2022), Bowler et al. (2017), McLure et al. (2013), Finkelstein et al. (2015), and Bouwens et al. (2021) are shown. The green dotted line and dot–dashed orange line show the best-fitting Schechter functions from Bowler et al. (2017) and Bouwens et al. (2021), respectively. The bottom panel shows the DPL fit to the rest-UV LF at  $z \simeq 7$  derived by Bowler et al. (2017), and the DPL + DPL fit derived by Harikane et al. (2022), shown as the green and blue lines, respectively. The red open circles show our LF points incorporating a contamination factor of 50 per cent, as derived by SED fitting of objects from Harikane et al. (2022) that overlapped with our catalogues.

#### 5.3 Results

In this section, we present our binned rest-frame UV LF measurement at  $z \simeq 7$  and compare to measurements and best-fitting functions from other studies. In Fig. 5, we show the measurement of the rest-frame UV LF at  $z \simeq 7$  from this work. The values of each bin are reported in Table 6.

## 5.3.1 Comparison of binned LF points

We compare our measurement of the  $z \simeq 7$  UV LF to current widearea studies from Harikane et al. (2022) and Bowler et al. (2017), and add studies of the faint end using *Hubble* from McLure et al. (2013), Finkelstein et al. (2015), and Bouwens et al. (2021). Bowler et al. (2017) do not probe as bright as this work because they were limited to a smaller survey area of 1.65 deg<sup>2</sup>. Harikane et al. (2022) use 20.2 deg<sup>2</sup> of survey data comparable to the depths of optical bands in this work, allowing them to provide the first constraints on the  $z \simeq 7$  UV LF beyond  $M_{\rm UV} \lesssim -23$ .

Our faintest point at  $M_{\rm UV} \simeq -22.2$  is in agreement with those found by Finkelstein et al. (2015), Bowler et al. (2017), Bouwens et al. (2021), and Harikane et al. (2022). Galaxies at  $M_{\rm UV} \simeq -22$ represent the brightest that can be found by the widest area Hubble surveys, and close to the faintest that can be found from VIDEO. The brightest candidate in the current widest area space-based survey (Bouwens et al. 2021) has an absolute UV magnitude  $M_{\rm UV} \simeq -22.1$ . Low number counts at this magnitude result in large error bars for their brightest bright bin, shown in yellow in Fig. 5. Our point at  $M_{\rm UV} \simeq -22.9$  is in agreement with Bowler et al. (2017) and Harikane et al. (2022), although errors between these three studies span over an order of magnitude. The Bowler et al. (2017) point lies roughly in the middle of our point and that of Harikane et al. (2022). Our brightest point in the bin at  $M_{\rm UV} \simeq -23.7$  is in tension with Harikane et al. (2022) at the  $2\sigma$  level, being four times lower than their determination at  $M_{\rm UV} \simeq -23.9$ . In Fig. 2, we showed that the high-redshift solution of SED fitting for the two likely BDs in their sample (see Fig. 4 and Section 4.1) puts them at very bright absolute UV magnitudes of  $M_{\rm UV} \sim -24$ , explaining this tension. We note that our brightest bin contains only one galaxy, VIDEO\_z7\_28. This source could be a BD, as it is particularly luminous and at a similar redshift to the Harikane et al. (2022) likely BDs (Fig. 2). If VIDEO\_z7\_28 is indeed shown to be a BD with later follow-up, this would change this point into an upper limit of  $2.07 \times 10^{-8}$  objects mag<sup>-1</sup> Mpc<sup>-3</sup>, leaving our conclusions unchanged. Finally, we compute an upper limit for the UV LF in a bin at  $M_{\rm UV}\simeq -24.4$  with width  $\Delta M = 0.75$ , shown in Fig. 5. The LF point derived by Harikane et al. (2022) at  $M_{\rm UV} = -24.42$  lies on our upper limit, within the errors.

In summary, our determination of the UV LF agrees well with previous studies at  $M_{\rm UV} \sim -22$ . At brighter magnitudes, however, we find a deficit of sources in comparison to Harikane et al. (2022). An analysis of the two brightest objects in their study suggests they are likely to be BDs. We therefore attribute this difference to contamination and emphasize the importance of NIR photometry in the selection of  $z \simeq 7$  LBGs.

#### 5.3.2 Schechter versus double power-law fitting

In the top panel of Fig. 5, we show the best-fitting Schechter function fits as derived by Bouwens et al. (2021) and Bowler et al. (2017). Our results appear to agree with the bright end of the Bouwens et al. (2021) Schechter fit, although this is due to large errors on our brightest bin. The Schechter function fit derived by Bowler et al. (2017) uses results from McLure et al. (2013). Similarly to Bowler et al. (2017) and Harikane et al. (2022), we find that the Schechter function does not fully reproduce the number density of bright galaxies, which our results show continues beyond  $M_{\rm UV} \simeq -23$ .

In the bottom panel of Fig. 5, we show the DPL derived by Bowler et al. (2017) and the DPL + DPL (galaxy and AGN

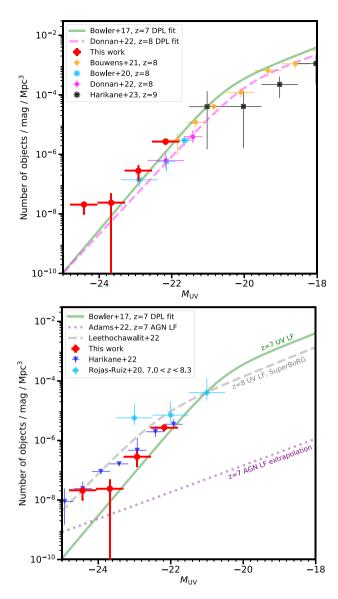
component) derived by Harikane et al. (2022). These functions provide a qualitative better fit to the UV LF measurements from our study. Furthermore, the reduced  $\chi^2$  value we compute for the DPL of Bowler et al. (2017) is  $\chi^2_{red} = 2.3$ , which is preferred to their Schechter fit ( $\chi^2_{red} = 4.7$ ), as well as the Schechter fit of Bouwens et al. (2021) ( $\chi^2_{red} = 4.3$ ), providing further evidence for a DPL as the form of the rest-UV LF at  $z \simeq 7$  continuing out to  $M_{\rm UV} \simeq$ -24. The DPL + DPL of Harikane et al. (2022) is in excess of our results, with  $\chi^2_{red} = 6.2$ . Our faintest two bins are in agreement with both Bowler et al. (2017) and Harikane et al. (2022). In particular, the faintest bin is in closer agreement with Harikane et al. (2022) and the middle bin lies closer to the Bowler et al. (2017) DPL. The differences are more pronounced at brighter magnitudes, where our results are in tension with the Harikane et al. (2022) fit at the  $2\sigma$  level. The Bowler et al. (2017) DPL is a factor of 5 lower at  $M_{\rm UV}\simeq -22.9$ , and a factor of 10 lower at  $M_{\rm UV}\simeq -23.7$  than Harikane et al. (2022). Our upper limit at  $M_{\rm UV} \simeq -24.4$  lies just above the Harikane et al. (2022) DPL + DPL. Assuming that AGN do not contribute to the bright end of rest-frame UV LF at  $z \simeq 7$ (see Section 6.2) such that the number density continues to decline rapidly at brighter magnitudes, we expect that a measurement of the LF at this magnitude would lie closer to the DPL of Bowler et al. (2017), which is 20 times lower than predicted by Harikane et al. (2022).

As shown in Section 4.1 and Fig. 4, the brightest objects from Harikane et al. (2022) that lie in XMM-LSS are best fit by BD templates when  $YJHK_s$  photometry is included. This suggests contamination in their brightest bins. The impact of this contamination on the LF is discussed further in Section 6.1. By utilising VIDEO photometry we have thus provided a robust measurement of the restframe UV LF at  $z \simeq 7$  out to  $M_{\rm UV} \simeq -24$ .

#### 5.4 Comparison to $z \simeq 8$ results

In this section, we compare our results to various measurements of the rest-frame UV LF at  $z \simeq 8$  to ascertain whether we find any evidence for evolution in the LF from  $z \simeq 7-8$ . In the top panel Fig. 6, we plot results from other ground-based studies (Bowler et al. 2020; Donnan et al. 2023) who both use deep optical and NIR imaging from VIDEO, UDS, and UltraVISTA, in a similar manner to this work, to identify robust  $z \simeq 8$  galaxy candidates. We also plot space-based results from Bouwens et al. (2021) of the faint end of the  $z \simeq 8$  LF. As expected, our results at  $z \simeq 7$  are in excess of these results at  $z \simeq 8$ . We show the DPLs measured at  $z \simeq 7$  and  $z \simeq 8$  by Bowler et al. (2017) and Donnan et al. (2023), respectively. Little evolution is expected in the bright end between these epochs (Bowler et al. 2020; Harikane et al. 2023), and this is reflected in the small differences between the DPLs and the slight excess of our results to those at  $z \simeq 8$ .

In the bottom panel of Fig. 6, we show recent results from *Hubble* pure-parallel studies. Leethochawalit et al. (2022) measure the  $z \simeq 8$  LF in 0.41 deg<sup>2</sup> of the Search for the Brightest of Reionizing Galaxies and Quasars in *Hubble Space Telescope (HST)* Parallel Imaging Data (SuperBoRG), and Rojas-Ruiz et al. (2020) use similar data to derive comparable results. Both studies are in tension with this work. The DPL derived by Leethochawalit et al. (2022) is a factor of 2.5–8 higher than our results, with the excess more pronounced at brighter magnitudes. At  $M_{\rm UV} \simeq -24.4$ , their DPL lies on the upper limit derived in this work. Similarly, the brightest LF point measured by Rojas-Ruiz et al. (2020) is 20 times higher than our measurement at  $M_{\rm UV} \simeq -23.7$ .



**Figure 6.** The top panel shows a comparison of our  $z \simeq 7$  rest-UV LF results (red circles) and the  $z \simeq 7$  DPL derived by Bowler et al. (2017) with various  $z \simeq 8$  results. The blue squares show the  $z \simeq 8$  rest-UV LF results from Bowler et al. (2020). The magenta hexagons show  $z \simeq 8$  rest-UV LF results from Donnan et al. (2023). Both of these studies use deep ground-based optical and NIR imaging. The magenta dashed line is the DPL fit derived by Donnan et al. (2023). The yellow points are  $z \simeq 8$  rest-UV LF results from Bouwens et al. (2021), derived from *Hubble* data. The bottom panel again shows our  $z \simeq 7$  rest-UV LF results (red circles) and  $z \simeq 7$  DPL derived by Bowler et al. (2017). We include the  $z \simeq 7$  rest-UV LF points found by Harikane et al. (2022) as dark blue triangles. Here we compare to  $z \simeq 8$  rest-UV LF results derived from *Hubble* pure-parallel programs, shown as light blue diamonds (Rojas-Ruiz et al. 2020) and the grey dashed DPL fit derived by Leethochawalit et al. (2022). We also plot the  $z \simeq 7$  extrapolation of the AGN LF found at z = 3-5 by Adams et al. (2023) as the purple dotted line.

#### 6 DISCUSSION

Our results agree with other ground-based deep optical + NIR studies at  $z \simeq 7-8$ , but work using deep optical-only ground-based imaging and pure-parallel *Hubble* imaging derive an excess compared to our findings by up to an order of magnitude. In this section, we discuss the potential contributions of quasars and BD contaminants, and implications of our results on the evolution of the bright end of the rest-frame UV LF. We also consider how upcoming space-based wide-area missions can aid in building a clearer picture of bright LBG abundances at high redshift.

#### 6.1 Contamination in the bright end of the LF

In Section 5.3 and Fig. 5, we showed that our results are in agreement with the steeper DPL found by Bowler et al. (2017) and are in tension with the flatter decline found by Harikane et al. (2022). Without the  $YJHK_s$  bands, most  $z \sim 7$  candidates found in HSC imaging appear as single-band detections, or at most have a detection in the Z and y bands. Thus only a spectral break can be detected, leaving such samples more vulnerable to contamination. Moreover, by including the HSC narrow-bands our objects in XMM-LSS further gain the advantage of the NB921 and Y filters that overlap with the Z and y filters, respectively, allowing for better determination of the nature of the break of an object. We find this to be powerful in detecting the slope of BD SEDs.

As discussed in Section 4.1, when the  $YJHK_s$  data are included in the SED fitting analysis, we found that four out of eight objects from the GOLDRUSH catalogue that overlaps with our data have poor z $\simeq$  7 SED fits. The fainter two of these objects have best-fitting SEDs at z < 6.5, while the brighter two have best-fitting BD templates, as shown in Fig. 4. This gives us a fiducial contamination rate of 50 per cent. We apply this contamination factor to our derived LF points, shown as open red circles in Fig. 5, to assess the impact of contamination on our robust sample. For our faintest and middle bins, these contaminated points lie slightly above and on the fit derived by Harikane et al. (2022), respectively. Our upper limit also suggests that their DPL + DPL is overestimated at  $M_{\rm UV} \simeq -24.4$ . This implies that not using YJHKs results in an overestimation of the bright end of the rest-UV LF at  $z \simeq 7$ . Note that our brightest contaminated LF point still lies slightly below the Harikane et al. (2022) value at  $M_{\rm UV} = -23.92$ . This suggests that the fiducial contamination rate of 50 per cent could be higher at brighter magnitudes, with brighter bins preferentially affected by BDs or low-redshift galaxies when only optical data are included. This is in tension with the contamination rate estimated by Harikane et al. (2022), who claim this is negligible based on results from their spectroscopic catalogue. They are also unable to use SED fitting to remove  $z \sim 2$  interlopers at z > 6 due to lack of bands redwards of the Lyman break. This contamination likely causes the unphysical crossing of the UV LFs as derived in their study at  $z \simeq 6$  and  $z \simeq 7$  in the range  $-25 \leq M_{\rm UV} \leq -23$  (see fig. 8 in Harikane et al. 2022).

#### 6.2 AGN contribution

In Section 5.3, we showed that the results of pure-parallel *Hubble* studies (Rojas-Ruiz et al. 2020; Leethochawalit et al. 2022) at  $z \simeq$  8 are significantly in excess of our results at  $z \simeq 7$ . Leethochawalit et al. (2022) attribute this excess to a contribution by AGN. The motivation for this comes from Harikane et al. (2022) who suggest that the faint end of the AGN LF dominates at  $M_{\rm UV} \le -24$  at z = 4-7. Adams et al. (2023) measure the AGN LF in the same fields as this study at z = 3-5, as well as in the COSMOS field, and find that the AGN LF evolves more rapidly towards higher redshift. Matsuoka et al. (2018) find that this evolution continues to  $z \simeq 6$ : their AGN LF crosses the  $z \simeq 7$  DPL of Bowler et al. (2017) at  $M_{\rm UV} \simeq -23.6$ . Thus, if this strong decline in number density of faint UV-selected AGN continues to  $z \simeq 7$  (AGN are rare at this epoch; Mortlock et al. 2011; Bañados et al. 2018; Wang et al. 2021), it is unlikely AGN

contribute to the  $z \simeq 7$  galaxy LF at the brighter magnitudes probed by this work and the pure-parallel studies from HST. Furthermore, in the bottom panel of Fig. 6, we show the extrapolation of the AGN LF to z = 7 as estimated by Adams et al. (2023). We extrapolate the z = 4.8 AGN LF to our mean redshift assuming a conservative value of k = -0.82, providing a likely upper limit on the number density of faint  $z \simeq 7$  AGN. Their extrapolation is  $\simeq 10$  times lower than the DPL computed by Bowler et al. (2017) at  $M_{\rm UV} = -23$ . Adams et al. (2023) conclude that the AGN number density at  $z \simeq 7$  is insufficient to contribute to the galaxy LF as found by Bowler et al. (2017). Instead, our results are consistent with little evolution in the bright end between  $z \simeq 7$  and 8, as determined by studies also using deep ground-based optical and NIR imaging (Bowler et al. 2020; Donnan et al. 2023). Contamination in the Leethochawalit et al. (2022) sample could be responsible for the excess of objects found in their work: for two-thirds of their sample, only a single filter (F350LP) is available bluewards of the break. This is the same filter used in Rojas-Ruiz et al. (2020) to assert non-detections bluewards of the break, likely leading to imperfect removal of low-redshift galaxies and BDs. The true overlap between the galaxy and AGN LFs at  $z \simeq 7$  therefore likely occurs somewhat brighter, at  $M_{\rm UV} < -24$ . We note that this discussion pertains to unobscured AGN – highly obscured AGN at z  $\simeq$  7 have been discovered in the COSMOS field (Endsley et al. 2023).

#### 6.3 Astrophysical interpretation

Robust measurements of the bright end of the rest-UV LF at  $z \simeq$ 7 are vital for understanding the evolution of the most star-forming and arguably massive galaxies between this epoch and the present day. The growth of massive galaxies must somehow be suppressed to account for the change in shape of the UV LF over cosmic time, with a DPL observed at high redshift to the Schechter form seen at low redshift. One mechanism that may govern the shape of the bright end of the LF is quenching by AGN, where feedback is driven by energetic release from their accretion disc (e.g. Davé et al. 2019; Lovell et al. 2022), removing gas reservoirs from galaxies and limiting star formation. The AGN number density increases rapidly between z = 6 and 5, thus we know that AGN activity becomes more common (e.g. Matsuoka et al. 2018; Niida et al. 2020; Adams et al. 2023). However, our results suggest that the contribution from AGN at  $z \simeq 7$  to the rest-frame UV LF is minimal. This indicates that quenching of star formation by AGN is unlikely to impact bright galaxies at  $z \simeq 7$ .

Alternatively, the build-up of dust may lead to significant obscuration of UV light from galaxies, modifying the intrinsic rest-frame UV LF and causing a steeper decline in the observed bright end of restframe UV LF. Star formation enriches the interstellar medium (ISM) with dust, thus preferentially impacting the most massive galaxies. For example, Vijayan et al. (2021) use a simple model linking dust attenuation with ISM metal content and find that dust attenuation becomes important for galaxies with  $M_{\rm UV} \lesssim -21$  at  $z \simeq 7$ . Our results suggest dust obscuration does not significantly impact the bright end of the rest-frame UV LF at this epoch. Dust attenuation values are low on average (see Table 5) with no dependence on  $M_{\rm UV}$ . These results suggest that a lack of dust attenuation at this epoch could explain the observed DPL shape and the almost constant number density of bright  $z \simeq 7$  and  $z \simeq 8$  sources (e.g. as shown in Ferrara et al. 2023). However, we note that the selection of rest-frame UV-bright galaxies can be blind to dust-obscured star formation, and recent studies are beginning to reveal the importance of this process (e.g. Bowler et al. 2022; Inami et al. 2022; Algera et al. 2023) and constrain dust properties at  $z \gtrsim 6.5$  (e.g. Inami et al. 2022; Sommovigo et al. 2022).

#### 6.4 Outlook for upcoming surveys

At  $z \simeq 7$  there is a need to bridge the gap between ground-based and space-based studies around the knee of the UV LF. A mismatch can be seen in Fig. 5 between the brightest bins of Bouwens et al. (2021) and the faintest bins of Bowler et al. (2017) and this work. At the knee, the former suffers from low number counts, and the latter suffers from a lack of depth. Euclid will be able to address this effectively: it will provide 50 deg<sup>2</sup> of YJH imaging down to 5 $\sigma$  depths of  $m_{AB}$  $\simeq$  26.4 in its deep fields (van Mierlo et al. 2022), comparable to the depths of UltraVISTA (covering 1.5 deg<sup>2</sup>). It will also be instrumental in providing much more accurate measurements of the ultrabright end of the  $z \simeq 7$  LF. It will provide 15 000 deg<sup>2</sup> of YJH imaging down to  $m_{AB} \simeq 24.5$  (Euclid Collaboration 2022b), slightly deeper than VIDEO. Simply providing thousands more sources will vastly improve number statistics. Additionally, with much greater spatial resolution that comes with space-based observatories, blending will be less of an issue, and will allow a more accurate deconfusion of the Spitzer/IRAC data. The morphology of objects may even be used to distinguish between high-redshift galaxies (extended) and BDs or AGN (point-like). Since AGN are very rare at  $z \simeq 7$ , BD contaminants that appear as point sources could be efficiently removed from LBG samples via a size or morphology cut. We may also be able to finally probe bright enough to see contributions by the  $z \simeq 7$  AGN LF, settling the debate about where the overlap between the AGN LF faint end and galaxy LF bright end begins.

Although *JWST* is finding many 'bright' candidates at ultrahigh redshifts (e.g. Naidu et al. 2022), it may struggle to probe the very bright end ( $M_{\rm UV} < -23$ ) of the UV LF during the Epoch of Reionization since the area it can survey is more limited than ground-based observations. In the top panel of Fig. 6, we plot the results of Harikane et al. (2023) who measure the  $z \simeq 9$  UV LF in ~90 arcmin<sup>2</sup> of NIRCam imaging. The smaller area compared to this work means they are unable probe brighter than  $M_{\rm UV} \simeq -22$ . Wider area studies such as COSMOS-Web (Casey et al. 2022) will be able to cover up to 0.5 deg<sup>2</sup>, but this is still significantly smaller than what can be achieved from ground-based observatories and with *Euclid* and *Roman*. It is clear that within the next decade, the combination of *JWST* with wide-area space observatories will provide a revolution in the measurement of the z > 6 LF over a broad magnitude range.

#### 7 CONCLUSIONS

We have conducted a wide-area search for  $z \simeq 7$  LBGs using deep NIR photometry from the VIDEO survey combined with deep optical data and *Spitzer*/IRAC photometry over an area of 8.2 deg<sup>2</sup>. Candidates were selected in a Y + J stack reaching  $5\sigma$  depths of  $m_{AB} = 25.3$  in XMM-LSS and  $8\sigma$  depths of  $m_{AB} = 24.8$  in ECDF-S, with non-detections required in bands bluewards of the Lyman break. We show that the inclusion of NIR data from VISTA and *Spitzer* enables a robust removal of low-redshift galaxies and Galactic BD contaminants.

We found 28 galaxy candidates with a mean redshift  $\bar{z} = 6.74$ and UV absolute magnitudes in the range  $-23.5 \le M_{\rm UV} \le -21.6$ . We recovered two spectroscopically confirmed sources from the REBELS survey (Bouwens et al. 2022). The BAGPIPES SED fitting code was used to confirm the photometric redshifts and derive galaxy properties. We derived stellar masses of  $9.1 \le \log_{10}(M/M_{\odot}) \le 10.9$ for objects with unconfused *Spitzer/*IRAC photometry, suggesting that these galaxies are some of the most massive at this epoch.

We measure the UV LF at  $z \simeq 7$  using our candidates and compare to the current widest area searches: Harikane et al. (2022) who use optical-only data in an area of 20.2 deg<sup>2</sup>, and Bowler et al. (2017) who use optical and NIR data in an area of 1.65 deg<sup>2</sup>. Our results agree with the DPL fit of Bowler et al. (2017) to brighter magnitudes, however they lie significantly below the results of Harikane et al. (2022). Through SED fitting of the Harikane et al. (2022) candidates, utilizing the VIDEO *YJHK*<sub>s</sub> photometry and *Spitzer/IRAC*, we find that the brightest galaxy candidates identified by Harikane et al. (2022) are likely to be BDs. Our results suggest that using optical data alone to select  $z \simeq 7$  galaxies leads to an overestimate by a factor of 2 at  $M_{\rm UV} \simeq -23$  and by a factor of 10 at  $M_{\rm UV} < -24$  due to contamination.

Extrapolating findings from Adams et al. (2023) predict a negligible contribution of unobscured AGN at the magnitudes we probe. This is in tension with the conclusions of recent pure-parallel *Hubble* results at z = 8 (Rojas-Ruiz et al. 2020; Leethochawalit et al. 2022), which we find to be in excess of our z = 7 LF by a factor of >10. Our results provide a robust measure of the bright end of the  $z \simeq$ 7 UV LF, which does not evolve significantly from  $z \simeq 8$ . This suggests a lack of dust attenuation and/or mass quenching between these epochs. Upcoming wide-area space missions such as *Euclid* and *Roman* will provide much larger samples extracted from thousands of deg<sup>2</sup> of NIR imaging to depths comparable to VIDEO, allowing for determinations of the UV LF beyond  $M_{\rm UV} < -24$  at z > 6, whilst also providing superior resolution to ground-based imaging, permitting the removal point-like BD contaminants.

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# DATA AVAILABILITY

All imaging data were obtained from original sources in the public domain. Catalogues of the photometry of samples used in this study can be provided on request.

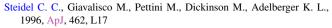
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# APPENDIX A: ADDING INCLUSIVE CANDIDATES TO THE LF CALCULATION

In this appendix, we present a more inclusive selection of candidates, which includes sources that have a plausible BD fit but are still

**Table A3.** The rest-UV LF points at  $z \simeq 7$ , calculated by incorporating the inclusive candidates. The first two columns show the absolute UV magnitude of the bin and the width of the bin. The third column shows the number of galaxies in the bin. The fourth column shows the comoving number density of galaxies we calculate with equations (1) and (2). The final column shows the median completeness in each bin.

Bin (mag)	Bin width (mag)	n <sub>gal</sub>	$\phi \ (mag^{-1} Mpc^{-3})$	Completeness
-22.05	0.5	18	$2.57 \pm 0.97 \times 10^{-6}$	0.68
-22.55	0.5	10	$8.40 \pm 4.51 \times 10^{-7}$	0.74
-23.175	0.75	5	$2.02 \pm 1.19 \times 10^{-7}$	0.79
-23.925	0.75	1	$2.37 \pm 2.50 \times 10^{-8}$	0.82

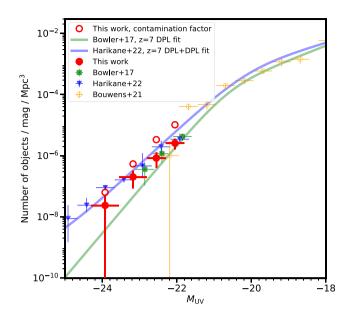
better represented by a high-redshift galaxy model. Bowler et al. (2015) state that removing objects with a good BD fit  $\chi^2 < 10$ , likely removes genuine high-redshift galaxies. We retain objects that have good BD fits and significantly better galaxy fits,  $\chi^2_{star} - \chi^2 > 4$ . These candidates are not robust due to the well-fitted BD SEDs. However, they pass our other selection criteria. We find four inclusive candidates in XMM-LSS, and four in ECDF-S. Their photometry is

**Table A1.** The photometry of our  $z \simeq 7$  inclusive candidates. The top section shows objects in XMM-LSS, and the bottom section shows objects in ECDF-S. Objects are ordered by their photometric redshift. The first column shows the object ID, and the next two columns show the coordinates of the candidate. The remaining bands show the photometry in the bands available in each field. We require  $<2\sigma$  detections bluewards of and including the *i*/*I*-bands, so we only present bands redwards of this. The photometry is measured in a 2.0 arcsec diameter circular aperture apart from the *Spitzer*/IRAC bands where 2.8 arcsec diameter apertures are used to account for the broader PSF. The photometry is corrected to a total flux assuming a point-source correction using a  $10 \times 10$  grid of PSFs measured by PSFEx across each VIDEO tile.

ID	RA	Dec.	Ζ	NB921	У	Y	J	Н	Ks	[3.6]	[4.5]
170389	02:16:38.75	-04:28:50.98	>26.7	>26.4	>25.8	$24.8^{+0.3}_{-0.2}$	$24.6^{+0.3}_{-0.2}$	$24.3^{+0.4}_{-0.3}$	>24.0	>24.5	$24.8^{+0.6}_{-0.4}$
1415071	02:25:49.21	-04:02:38.00	>26.2	$25.7^{+0.6}_{-0.4}$	$24.8_{-0.3}^{+0.5}$	$24.8^{+0.3}_{-0.2}$	$24.6^{+0.3}_{-0.2}$	$24.8^{+0.7}_{-0.4}$	>24.0	$23.7^{+0.2}_{-0.2}$	$24.1_{-0.3}^{+0.5}$
1017553	02:22:52.86	-05:04:37.97	>26.4	>26.0	>24.6	$24.7^{+0.3}_{-0.2}$	$24.7_{-0.3}^{+0.4}$	$24.5_{-0.3}^{+0.4}$	>23.9	$23.3_{-0.1}^{+0.2}$	$24.1_{-0.4}^{+0.5}$
914913	02:22:07.26	-05:24:02.32	>26.3	>25.5	>24.9	$25.1_{-0.2}^{+0.2}$	$25.2\substack{+0.5 \\ -0.3}$	$25.1\substack{+0.7 \\ -0.4}$	$24.8\substack{+0.7 \\ -0.4}$	$24.0\substack{+0.2 \\ -0.2}$	$24.5\substack{+0.7 \\ -0.4}$
1495402	03:35:25.99	-28:38:19.13	>25.1	_	_	$23.9_{-0.1}^{+0.1}$	$22.9_{-0.1}^{+0.1}$	$22.6^{+0.1}_{-0.1}$	$22.4_{-0.1}^{+0.1}$	$22.5^{+0.2}_{-0.2}$	$23.0^{+0.2}_{-0.2}$
705214	03:30:30.53	-27:23:16.47	>25.4	-	-	$24.6^{+0.1}_{-0.1}$	$24.6^{+0.3}_{-0.2}$	$24.8^{+0.6}_{-0.4}$	>24.1	$24.1_{-0.2}^{+0.2}$	>24.7
1090797	03:32:28.28	-27:48:02.25	>25.5	-	-	$24.2_{-0.1}^{+0.1}$	$23.8^{+0.2}_{-0.1}$	$24.1_{-0.2}^{+0.3}$	$24.2_{-0.3}^{+0.5}$	$24.3_{-0.2}^{+0.2}$	$24.7^{+0.3}_{-0.2}$
834371	03:31:09.95	-28:43:55.05	>25.0	-	-	$25.1_{-0.2}^{+0.3}$	$23.5_{-0.1}^{+0.1}$	$22.9^{+0.1}_{-0.1}$	$22.6_{-0.1}^{+0.1}$	$21.8^{+0.2}_{-0.2}$	$21.6^{+0.2}_{-0.2}$

**Table A2.** Inclusive candidates in XMM-LSS (top) and ECDF-S (bottom). These are defined as having good BD fits,  $\chi^2_{BD} < 10$ , yet have significantly better galaxy fits,  $\chi^2_{BD} - \chi^2 > 4$  when fitted without *GR* in XMM-LSS and *ugrGR* in ECDF-S. The first column shows the object ID. The next four columns show properties of the high-*z* solution: photometric redshift, absolute magnitude in a top-hat filter at 1500 Å with width 100 Å, extinction, and  $\chi^2$  value. The next two columns show the secondary photometric redshift and  $\chi^2$ . The final two columns show the stellar type for fitting to MLT dwarfs and  $\chi^2_{BD}$ .

ID	z	M <sub>UV</sub> (mag)	A <sub>V</sub> (mag)	$\chi^2$	z <sub>gal2</sub>	$\chi^2_{gal2}$	Stellar type	$\chi^2_{BD}$
170389	$7.34_{-0.14}^{+0.10}$	$-22.6\pm0.2$	0.0	5.3	9.0	26.5	T8	9.8
914913	$6.53\substack{+0.64\\-0.10}$	$-21.7\pm0.3$	0.0	4.3	1.2	14.4	Т3	8.1
1017553	$6.93\substack{+0.26\\-0.34}$	$-22.2\pm0.2$	0.0	0.9	1.5	13.5	Т3	6.2
1415071	$6.58\substack{+0.03 \\ -0.10}$	$-22.1\pm0.1$	0.2	1.6	1.4	18.4	Т3	9.0
705214	$6.61\substack{+0.66\\-0.18}$	$-22.3\pm0.3$	0.2	3.8	1.45	11.5	Т8	9.8
834371	$6.99\substack{+0.56\\-0.07}$	$-23.4\pm0.2$	0.0	1.8	1.7	12.2	L7	6.4
1090797	$7.33_{-0.19}^{+0.07}$	$-23.3\pm0.2$	0.0	3.2	1.5	38.4	Т8	6.8
1495402	$7.39_{-0.46}^{+0.04}$	$-24.1\pm0.2$	0.0	1.6	1.6	25.8	L4	5.7



**Figure A1.** The UV LF at z = 7, computed including the inclusive sample. The red circles show measurements of the number density from this work. Other measurements from Harikane et al. (2022), Bowler et al. (2017), and Bouwens et al. (2021) are shown. The solid lines show the best-fitting double power laws (DPLs) from Harikane et al. (2022) (blue) and Bowler et al. (2017) (green). The red open circles show our LF points incorporating a contamination factor of 50 per cent, as derived by SED fitting of objects from Harikane et al. (2022) that overlapped with our catalogues.

shown in Table A1. The results of SED fitting analysis with LEPHARE are shown in Table A2. In ECDF-S the candidates are much brighter than our primary sample: four have absolute UV magnitudes  $M_{\rm UV}$ < -23, and one has  $M_{\rm UV}$  < -24. The SED fitting and stamps for the inclusive candidates are shown in Figs B4 and B5. We add these objects to the primary sample and conduct a separate calculation of the rest-UV LF at  $z \simeq 7$ , as outlined in Section 5. The result is shown in Fig. A1. We also plot the DPL fits derived by Bowler et al. (2017) and Harikane et al. (2022), and their LF points. We include Bouwens et al. (2021) LF points to represent the faint end. The bins

are centred at -22.05, -22.55, -23.175, and -23.925, with bin widths of  $\Delta M = 0.5$  for the two fainter bins and  $\Delta M = 0.75$  for the two brighter bins. The bins contain 18, 10, 5, and 1 candidate galaxies, respectively. The LF points are summarized in Table A3. If we include the contamination factor of 50 per cent derived in Section 4.1, our three faintest bins lie slightly above the DPL + DPLderived by Harikane et al. (2022). The brightest bin is consistent with both DPLs, although the Poisson errors are large. The value of the LF in our brightest bin is a factor of 4 lower than the Harikane et al. (2022) bin at  $M_{\rm UV} = -23.92$ . The contamination factor pushes this bin onto their DPL. To summarize, the fiducial UV LF we present in Fig. 5 is our best estimate of the number density of genuine  $z \simeq$ 7 sources. However, even if we relax our selection criterion to allow sources that have a plausible BD fit into the sample, we still find the UV LF to be significantly lower than that found in optical-only surveys for our brightest bins.

#### **APPENDIX B: OBJECT STAMPS AND SED FITS**

In this appendix, we present postage stamps and SED fitting of all candidates in this work (Figs B1–B10). The first two sections are primary candidates in ECDF-S and XMM-LSS. The third section is Lyman  $\alpha$  candidates, all in XMM-LSS. The final two sections are inclusive candidates in ECDF-S and XMM-LSS.

#### **B1** Primary candidates

B1.1 XMM-LSS B1.2 ECDF-S **B2 Lyman α candidates** B2.1 XMM-LSS

# **B3** Inclusive candidates

B3.1 XMM-LSS

B3.2 ECDF-S

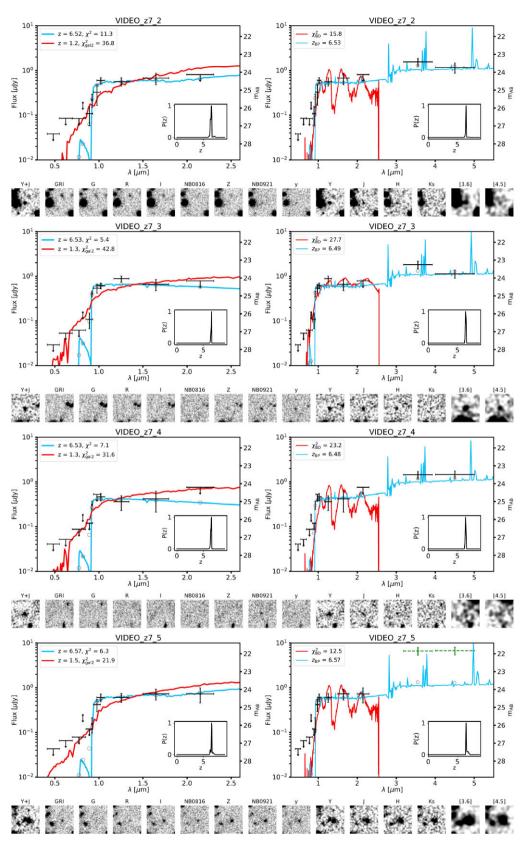
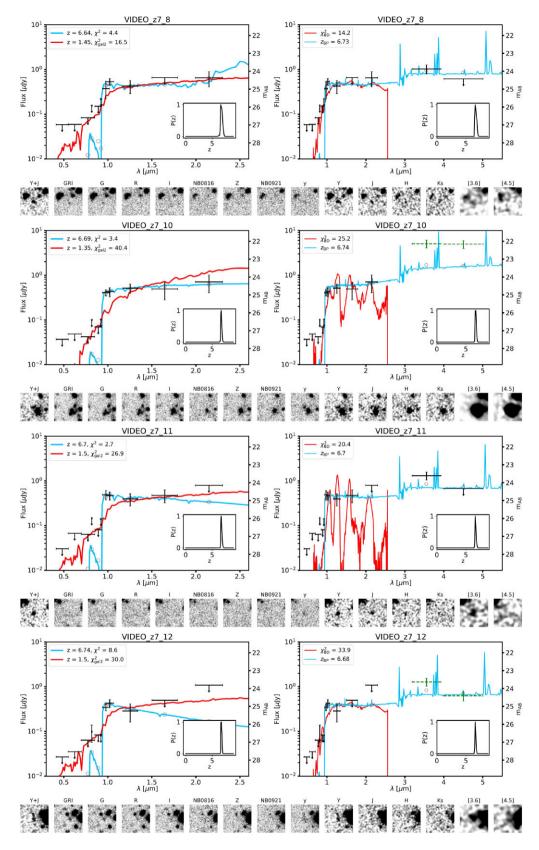
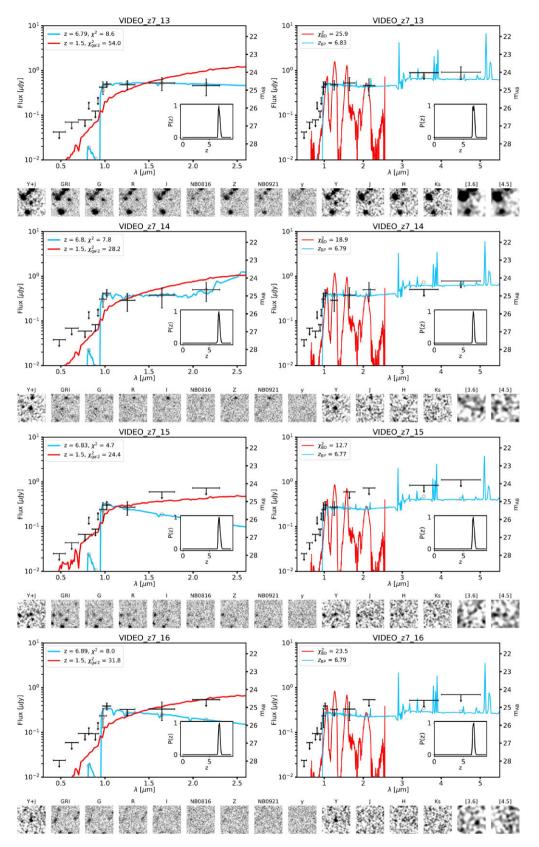


Figure B1. Candidate galaxies in XMM. The plots and stamps are the same as Fig. 3. Confused IRAC photometry is shown in green with dashed wavelength error bars.



**Figure B2.** – *continued* 



**Figure B3.** – *continued* 

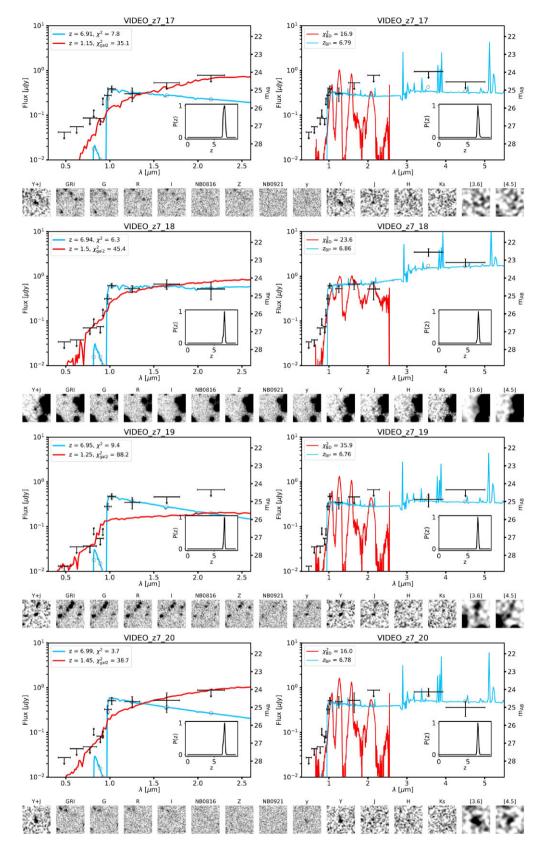


Figure B4. – continued

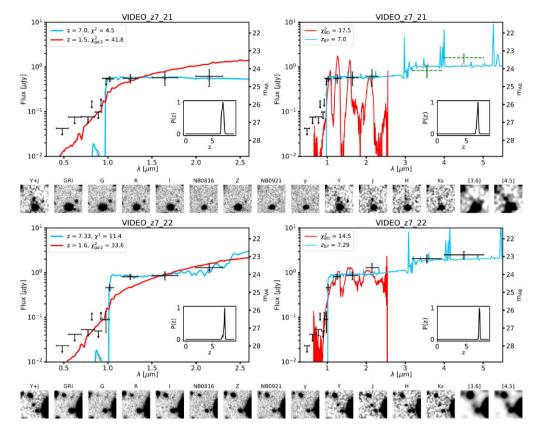


Figure B5. – continued

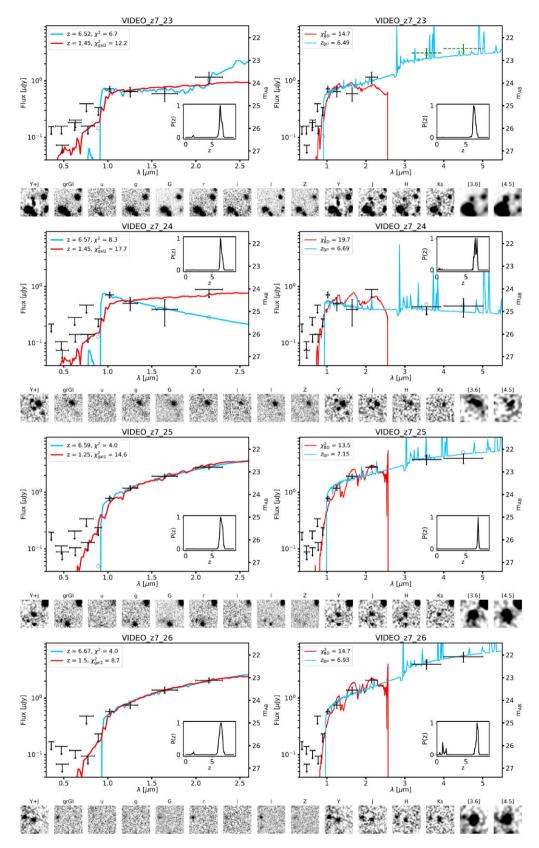
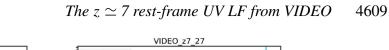
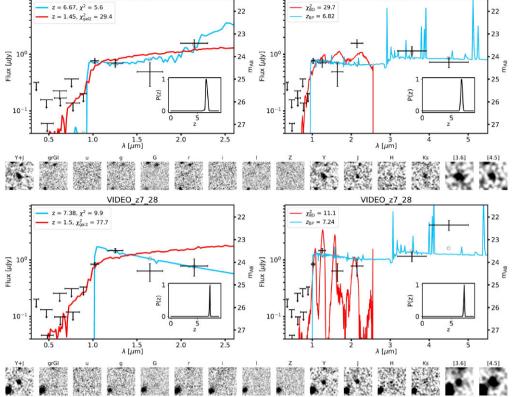


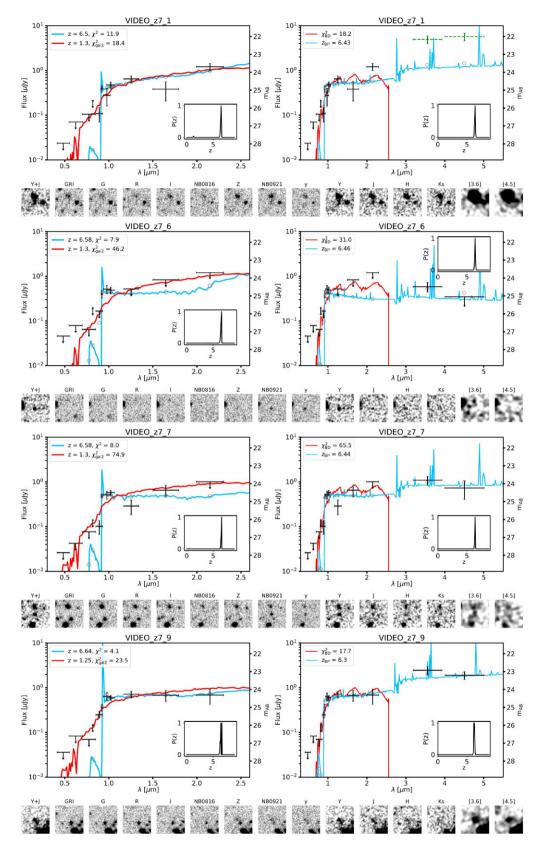
Figure B6. Candidate galaxies in CDFS. The plots and stamps are the same as Fig. 3. Confused IRAC photometry is shown in green with dashed wavelength error bars.



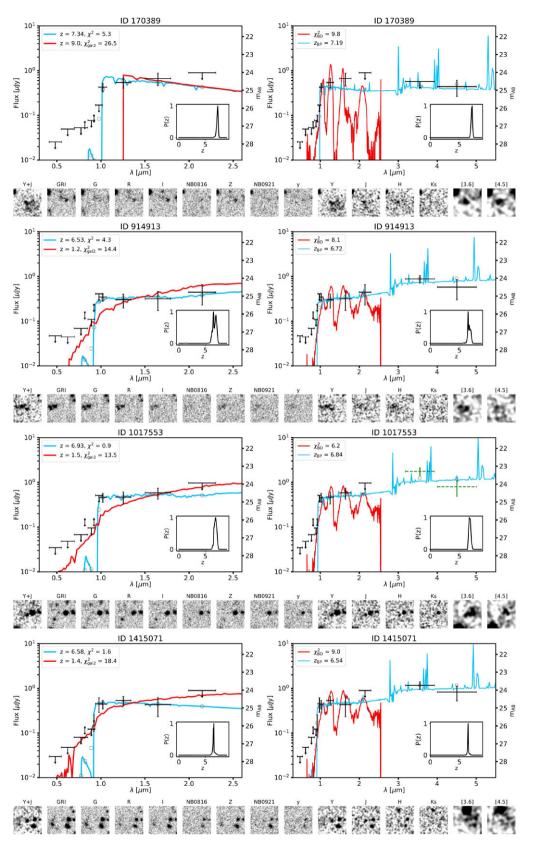


VIDEO\_z7\_27

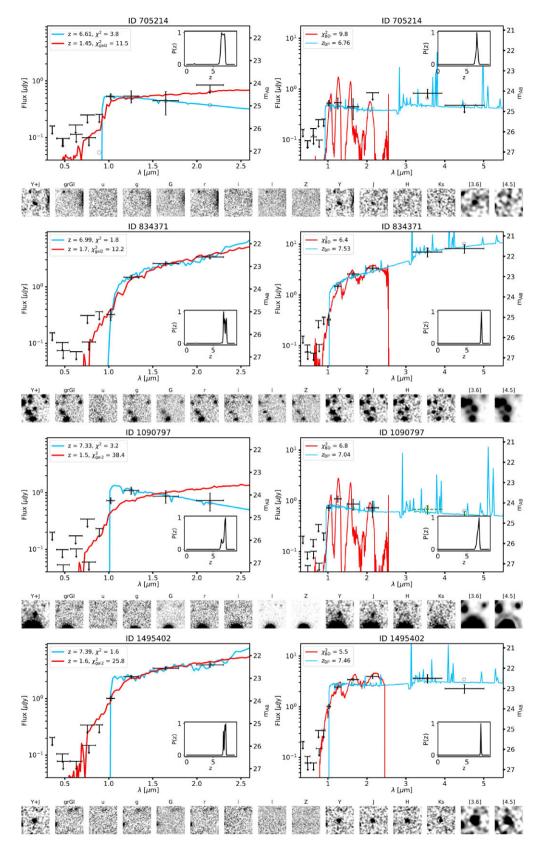
Figure B7. – continued



**Figure B8.** Candidate galaxies in *XMM* that require a Lyman  $\alpha$  emission line to be included in our sample. The plots and stamps are the same as Fig. 3. Confused IRAC photometry is shown in green with dashed wavelength error bars.



**Figure B9.** Inclusive candidates in XMM-LSS. These are defined as having good BD fits,  $\chi^2_{star} < 10$ , yet have significantly better galaxy fits,  $\chi^2_{star} - \chi^2 > 4$  when fitted without *GR*. The plots and stamps are the same as Fig. 3. Confused IRAC photometry is shown in green with dashed wavelength error bars.

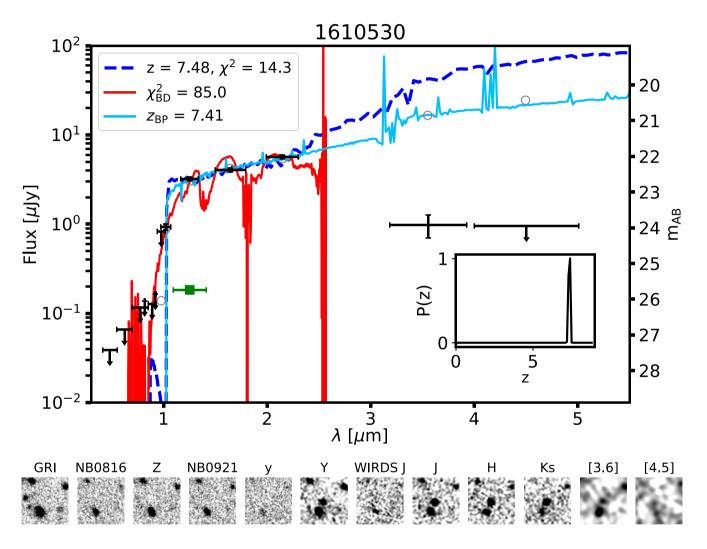


**Figure B10.** Inclusive candidates in ECDF-S. These are defined as having good stellar fits,  $\chi^2_{star} < 10$ , yet have significantly better galaxy fits,  $\chi^2_{star} - \chi^2 > 4$  when fitted without *ugrGR*. The plots and stamps are the same as Fig. 3. Confused IRAC photometry is shown in green with dashed wavelength error bars.

# APPENDIX C: UNUSUAL OBJECT IN XMM

In this appendix, we present the SED and postage stamp images of ID 1610530, an unusual object noted in Section 3.3. If this object is at high redshift, the rest-frame UV is very bright ( $M_{\rm UV} = -24.4$ ), but the *Spitzer/IRAC* photometry appears unusually faint, around 3.5–4 mag fainter than suggested by the model. When *Spitzer/IRAC* 

photometry is included in the SED fitting, the fit becomes very poor with  $\chi^2 = 167.6$ . A cross-talk artefact is perhaps ruled out by the  $3.5\sigma$ detection in the [3.6] band. Additionally, the object is not detected in WIRDS-J. In Fig. C1, we present a measurement of the WIRDS-J flux in a 2 arcsec diameter aperture at the position of the object. The measured flux is ~3 mag lower than measured in VIDEO-J.



**Figure C1.** SED fitting (top) and postage stamp images (bottom) of ID 1610530, the unusual object discussed in Section 3.3. In the top panel, the black points in both plots are the measured photometry, with non-detections replaced by  $2\sigma$  upper limits. The blue dashed curve shows the best-fitting LEPHARE solution. The solid light blue curve shows the best-fitting BAGPIPES solution, with its model photometry shown as the grey circles. The red curve shows the best-fitting BD solution. The legend shows the redshift and  $\chi^2$  of the best-fitting LEPHARE solution (without IRAC),  $\chi^2_{BD}$  (fit without *G*, *R*, and IRAC) and the redshift of the BAGPIPES solution,  $z_{BP}$ . The inset black curve is the redshift probability found by BAGPIPES. The bottom panel shows the  $10 \times 10 \operatorname{arcsec}^2$  postage stamps of the object in the filters used. We also show the WIRDS *J* band, where the object is not detected. The green square point in the top panel corresponds to the photometry measured in a 2 arcsec diameter circular aperture in WIRDS-*J*.

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