

# THE ANALYTICAL VIEW OF SOLUTION OF THE FIRST BOUNDARY VALUE PROBLEM FOR THE NONLINEAR EQUATION OF HEAT CONDUCTION WITH DEVIATION OF THE ARGUMENT

Yaroslav M. Drin\*, Iryna I. Drin†, Svitlana S. Drin‡ §

**Abstract.** In this article, for the first time, the first boundary value problem for the equation of thermal conductivity with a variable diffusion coefficient and with a nonlinear term, which depends on the sought function with the deviation of the argument, is solved. For such equations, the initial condition is set on a certain interval. Physical and technical reasons for delays can be transport delays, delays in information transmission, delays in decision-making, etc. The most natural are delays when modeling objects in ecology, medicine, population dynamics, etc. Features of the dynamics of vehicles in different environments (water, land, air) can also be taken into account by introducing a delay. Other physical and technical interpretations are also possible, for example, the molecular distribution of thermal energy in various media (solid bodies, liquids, etc.) is modeled by heat conduction equations. The Green's function of the first boundary value problem is constructed for the nonlinear equation of heat conduction with a deviation of the argument, its properties are investigated, and the formula for the solution is established.

**Key words:** heat nonlinear equation, boundary value problem, Green's function; deviation argument.

**2010 Mathematics Subject Classification:** 35K61.

*Communicated by Prof. A. V. Plotnikov*

## 1. Introduction

Thermal conductivity is the molecular distribution of thermal energy in various solids, liquids and gases due to the difference in temperature and due to the fact that the particles are in direct contact with each other. The process of heat conduction was first described by Jean Baptiste Joseph Fourier (1768 - 1830) in 1807 in the work "Equations with partial derivatives for heat conduction in solids". A description of the results of other scientists who studied and developed this theory is presented by T.N. Narasimhan [1]. Based on different criteria, models of heat conduction processes are divided into two groups of models

---

\*Department of Mathematical Problems of Management and Cybernetics, Institute of Physical, Technical and Computer Sciences, Yuriy Fedkovych Chernivtsi National University, 2, Kotsyubinsky av., Chernivtsi, 58012, Ukraine, [y.drin@chnu.edu.ua](mailto:y.drin@chnu.edu.ua)

†Department of Finance, Accounting and Taxation, Chernivtsi Trade and Economic Institute of the State Trade University, 7, Tsentralna Square, Chernivtsi, 58000, Ukraine, [iryndrin@gmail.com](mailto:iryndrin@gmail.com)

‡Department of Mathematics, National University of Kyiv-Mohyla Academy, 2 Skovorody vul., 04070 Kyiv, Ukraine, [svitlana.drin@ukma.edu.ua](mailto:svitlana.drin@ukma.edu.ua)

§Department of Statistics, School of Business, Örebro University, 2 Studentgatan st., 70182 Örebro, Sweden, [svitlana.drin@oru.se](mailto:svitlana.drin@oru.se)

using integral and fractional order derivatives. In this paper the solution of the first boundary value problem for the heat conduction equation with the variable diffusion coefficient and deviation of the argument is found.

If the evolution of the concentration of impurities, point defects, and the temperature field is studied, then the corresponding transfer coefficients are not constant values. Non-stationary models of one-dimensional heat conduction are described by the equation of heat conduction [2]. Different methods of solving this problem are described in [3], [4]. Applied aspects of such problems are described in [5], [6].

Processes with spatially dependent transmission coefficients or a desired thermal field are well studied and sufficiently describe processes in heterogeneous and nonlinear media [2], applied problems for modelling and research, whose transmission coefficients depend on time change, are also described here. At the same time, physically adequate modelling of thermal processes often requires their investigation in a semi-limited region [7].

In our paper, for the first time, we consider the first boundary value problem for a non-homogeneous nonlinear equation with a deviation of the argument and a variable diffusion coefficient in a semi-bounded domain, which generalizes the corresponding problem for [9].

## 2. Statement of the First Boundary Problem

Let  $a > 0, h > 0$  be real numbers;  $x \in R^+, t \in R^+$ , are independent variables;  $f, \varphi, \mu, D > 0$  are known continuous functions;  $u(x, t)$  is the desired function that describes the evolution of the system defined on the semi-axis  $x \in R^+$  for all  $t \in R^+$ . We will study the problem

$$u_t = D(t)u_{xx} + f(x, t, u(x, I_h(t))), x > 0, t > h, \quad (2.1)$$

$$u(x, t)|_{0 \leq t \leq h} = \varphi(x, t), x \geq 0 \quad (2.2)$$

$$u(0, t) = \mu(t), t \geq h \quad (2.3)$$

which is the first boundary value problem, where the functions  $\varphi(x, t) \in C(R^+ \times \{0 \leq t \leq h\})$  is initial function,  $\mu(t) \in C(R_n^+)$  is boundary function,  $R_h^+ \equiv \{t; t \geq h\}$ ,  $R^+ \equiv \{x; x \geq 0\}$ ,  $f(x, t, u) \in C(R^+ \times R_h^+ \times R)$  is the inhomogeneity of equation (2.1) is well known. If a smooth solution of the problem (2.1)–(2.3) is sought up to the limit, then the initial and marginal functions must be consistent  $\varphi(0, h) = \mu(h)$ .

## 3. The Steps Method

Let  $x \in R^+$ , then  $u(x, I_h(t)) = \varphi(x, t)$ , and from (2.1)–(2.3) we get the problem:

$$u_t = D(t)u_{xx} + f(x, t, \varphi(x, t)), \quad x > 0, h < t < 2h, \quad (3.1)$$

$$u(x, t)|_{t=h} = \varphi(x, h), x \geq 0, \quad (3.2)$$

$$u(0, t) = \mu(t), t \geq h \quad (3.3)$$

with the conditions of agreed  $\varphi(0, h) = \mu(h)$ .

We will solve a problem (3.1)–(3.3) in the form of sum of three functions

$$u(x, t) = u_1(x, t) + u_2(x, t) + u_3(x, t), \quad (3.4)$$

where  $u_i, 1 \leq i \leq 3$ , respectively, take into account the influence only initial condition, the boundary condition and the inhomogeneity of the, that is, they are the solutions of such problems.

**Problem 1.** Find a function  $u_1(x, t)$  that satisfies the conditions

$$\frac{\partial u(x, t)}{\partial t} = D(t) \frac{\partial^2 u(x, t)}{\partial x^2}, x > 0, h < t < 2h, \quad (3.5)$$

$$u(x, h) = \varphi(x, h), x \geq 0, \quad (3.6)$$

$$u(0, t) = 0, h \leq t \leq 2h, \quad (3.7)$$

moreover,  $\varphi(0, h) = u(0, h) = 0$  is a condition of agreement.

**Problem 2.** Find a function  $u_2(x, t)$  that satisfies equation (3.5) and conditions

$$u(x, h) = 0, x \geq 0, \quad (3.8)$$

$$u(0, t) = \mu(t), h \leq t \leq 2h, \quad (3.9)$$

moreover,  $\mu(h) = 0$  is a condition of agreement.

**Problem 3.** Find the function  $u_3(x, t)$  that satisfies equation (2.1) and conditions (3.7), (3.8), which are agreed.

### 3.1. Solving problem 1

Let's expand the domain of definition of equation (3.5) and the initial condition to  $x \in R, h \leq t \leq 2h$  and solve it by separating of variables method ( $u_1(x, t) = X(x)T(t)$ ). and after rearrangement in (3.5) and separation of variables, we obtain that  $T(t) = C(\lambda)e^{-\lambda^2 I_h(t)}$ ,  $I_h(t) = \int_h^t D(\tau) d\tau$ ,  $X(x) = e^{i\lambda x}$ , where  $\lambda$  is the variable separation parameter. Then the solution is  $u_1(x, t, \lambda) = C(\lambda)e^{-\lambda^2 I_h(t) + i\lambda x}$ ,  $\lambda \in R$  and to take into account all  $\lambda \in R$  we create a function

$$u_1(x, t) = \int_{-\infty}^{\infty} C(\lambda)e^{-\lambda^2 I_h(t) + i\lambda x} d\lambda, x \in R, \quad h \leq t \leq 2h,$$

which satisfies condition (3.6). Then we get that

$$\varphi(x, h) = \int_{-\infty}^{\infty} C(\lambda)e^{i\lambda x} d\lambda,$$

$$C(\lambda) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \varphi(\xi, h)e^{-i\lambda\xi} d\xi, \quad \lambda \in R,$$

$$u_1(x, t) = \int_{-\infty}^{\infty} \frac{1}{2\pi} \left\{ \int_{-\infty}^{\infty} e^{-\lambda^2 I_h(t) + i\lambda(x-\xi)} d\lambda \right\} \varphi(\xi, h) d\xi.$$

The inner integral calculated

$$\frac{1}{2\pi} \int_{-\infty}^{\infty} e^{-\lambda^2 I_h(t) + i\lambda(x-\xi)} d\lambda = \frac{1}{2\sqrt{\pi I_h(t)}} e^{-\frac{(x-\xi)^2}{4I_h(t)}}$$

is denoted by  $G(x - \xi; I_h(t))$  and is the fundamental solution of equation (3.5). Then

$$u_1(x, t) = \int_{-\infty}^{\infty} G(x - \xi; I_h(t)) \varphi(\xi, h) d\xi, \quad x \in R, \quad h < t < 2h. \quad (3.10)$$

We use formula (3.10) to construct a solution to problem 1. For this, instead of equation (3.5), we consider equation

$$\frac{\partial U(x, t)}{\partial t} = \frac{\partial^2 U(x, t)}{\partial x^2}, \quad x \in R, t > h \quad (3.11)$$

with conditions (3.6), (3.7), extending in condition (3.6) the initial function  $\varphi(x, h)$  for  $x < 0$  undefined, and we set the condition (3.7) as follows:

$$U(x, h) = \Psi(x, h) = \begin{cases} \varphi(x, h), & x \geq 0, \\ -\varphi(-x, h), & x < 0 \end{cases} \quad (3.12)$$

Then, according to formula (3.10), the solution of problem (3.11), (3.12), (3.7) is

$$U(x, t) = \int_{-\infty}^{\infty} \{G(x - \xi; I_h(t)) - G(x + \xi; I_h(t))\} \varphi(\xi, h) d\xi.$$

In the integral where  $\xi < 0$ , we replaced  $\xi = -\xi$ . Simplifying the difference of the exponents included in the expression for the function  $G$ , we obtain that

$$u_1(x, t) = \frac{1}{\sqrt{\pi I_h(t)}} \int_0^{\infty} \varphi(x, h) e^{-\frac{x^2 + \xi^2}{4I_h(t)}} \operatorname{sh} \frac{x\xi}{2I_h(t)} d\xi,$$

where  $x > 0, h < t < 2h$ . Using the method of mathematical induction, we prove that in case  $x \geq 0, kh < t < (k+1)h$  the solution to problem 1 takes the form

$$u_1(x, t) = \frac{1}{\sqrt{\pi I_{kh}(t)}} \int_0^{\infty} \varphi(\xi, kh) e^{-\frac{x^2 + \xi^2}{4I_{kh}(t)}} \operatorname{sh} \frac{x\xi}{2I_{kh}(t)} d\xi. \quad (3.13)$$

Let's mark

$$G_1(x, y, I_{kh}(t)) = \frac{1}{\sqrt{\pi I_{kh}(t)}} e^{-\frac{x^2 + y^2}{4I_{kh}(t)}} \operatorname{sh} \frac{xy}{2I_{kh}(t)} \quad (3.14)$$

$x \geq 0, y > 0, kh < t < (k+1)h, k \in N$ .

**Definition 3.1.** A function  $G_1(x, y, I_{kh}(t))$  is called a Green's function of problem (2.1), (2.2), (2.3) if it satisfies the following conditions:

1. the function  $G_1(x, y, I_{kh}(t))$  is continuous on  $x, y, t$ , continuously differentiable on  $t$  and twice continuously differentiable on  $x, y$  when  $x > 0, y > 0, kh < t < (k + 1)h, k \in \mathbb{N}$ , and possibly with the exception in the point  $x = y, t = kh$ ;
2. the function  $G_1(x, y, I_{kh}(t))$  by variables  $x$  and  $y$  satisfies the equation  $\frac{\partial G_1}{\partial t} = \frac{D(t)\partial^2 G_1}{\partial x^2}$  everywhere except in the points  $x = y, t = kh, k \in \mathbb{N}$ ;
3. the function  $G_1(x, y, I_{kh}(t))$  satisfies the boundary condition  $G_1(0, y, I_{kh}(t)) = 0$ .

The Green's function satisfying this definition is constructed above and takes the form (3.14)

$$G_1(x, y, I_{kh}(t)) = G_1(y, x, I_{kh}(t)).$$

### 3.2. Properties of the solution of the problem 1

Given that

$$G_1(x, y, I_{kh}(t)) = \frac{1}{2\sqrt{\pi I_{kh}(t)}} \left\{ e^{-\frac{(x-\xi)^2}{4I_{kh}(t)}} - e^{-\frac{(x+\xi)^2}{4I_{kh}(t)}} \right\}$$

we get from (3.13), when  $|\varphi(\xi, kh)| \leq M$ ,

$$\begin{aligned} |u_1(x, t)| &\leq M \frac{1}{2\sqrt{\pi I_{kh}(t)}} \left\{ \int_0^\infty e^{-\frac{(x-\xi)^2}{4I_{kh}(t)}} d\xi - \int_0^\infty e^{-\frac{(x+\xi)^2}{4I_{kh}(t)}} d\xi \right\} \\ &\equiv M \{I_1 - I_2\}. \end{aligned}$$

In the integral  $I_1$  we will do replacement  $\alpha = \frac{\xi-x}{2\sqrt{I_{kh}(t)}}$ , and in the integral  $I_2$   $\alpha = \frac{\xi+x}{2\sqrt{I_{kh}(t)}}$ . Then

$$I_1 = 2 \int_{-z}^\infty e^{-\alpha^2} d\alpha, I_2 = 2 \int_z^\infty e^{-\alpha^2} d\alpha$$

where  $z = \frac{x}{2\sqrt{I_{kh}(t)}}$  and we get an estimate

$$|u_1(x, t)| \leq M \operatorname{erf} \left( -\frac{x}{2\sqrt{I_{kh}(t)}} \right), \tag{3.15}$$

$x > 0, kh < t < (k + 1)h$

So, the following theorem is proved.

**Theorem 3.1.** *If there exists a number  $M > 0$  such that the initial function  $\varphi(x, kh)$  is bounded when  $x > 0, h > 0, k \in \mathbb{N}, |\varphi(x, kh)| \leq M$ , then the function  $u_1(x, t)$  (3.13) when  $x > 0, kh < t < (k + 1)h$  is also bounded and the estimate (3.15) is true for it.*

If  $\varphi(\xi, kh) = \varphi_0$ , where  $\varphi_0$  is a number, then

$$u_1(x, t) = \varphi_0 \operatorname{erf} \left( \frac{x}{2\sqrt{I_{kh}(t)}} \right)$$

$x > 0, kh < t < (k+1)h$ ,  $\operatorname{erf}(x) = \frac{2}{\sqrt{\pi}} \int_0^x \exp\{-\xi^2\} d\xi$  is the error function.

By direct verification, it is possible to make sure that the Green's function (3.14) satisfies the homogeneous heat conduction equation (item 2 of the definition). When formally differentiating the function (3.13) under the sign of the integral, we obtain expressions

$$\frac{1}{(I_{kh}(t))^r} \int_0^\infty \varphi(\xi, kh) |x \pm y|^m e^{-\frac{(x \pm y)^2}{4I_{kh}(t)}} dy,$$

$x > 0, kh < t < (k+1)h$ , where integrable functions are majored by an expression of the form  $M|\xi|^m e^{-\xi^2}$  that is integrable on the entire numerical axis. This ensures uniform convergence of the integrals obtained after differentiation under the sign of the integral. Then the Poisson integral (3.13) is a continuous function, differentiable of arbitrary order with respect to  $x$  and  $t$  when  $x > 0, kh < t < (k+1)h, k \in \mathbb{N}$ , bounded with a bounded initial function, satisfying the homogeneous heat conduction equation (3.5), since the Green's function (3.14) satisfies equation (3.5). The implementation of the initial condition (3.6) and the boundary condition (3.7) is carried out. Let us prove the uniqueness theorem of the solution to problem 1.

**Theorem 3.2.** *Let there be a number  $M > 0$  such that in the domain  $x \geq 0$  and  $kh \leq t \leq (k+1)h, k \in \mathbb{N}$  the functions  $u_1(x, t)$  and  $u_2(x, t)$  are bounded, that is  $|u_i(x, t)| < M, i = 1, 2$ , satisfy the equation (3.5) and condition*

$$u_1(x, kh) = u_2(x, kh), \quad x \geq 0, k \in \mathbb{N},$$

then

$$u_1(x, t) = u_2(x, t), \quad x \geq 0, kh \leq t \leq (k+1)h$$

Consider the function

$$v(x, t) = u_1(x, t) - u_2(x, t),$$

which is continuous, equation (3.5), bounded by

$$\begin{aligned} |v(x, t)| &\leq |u_1(x, t)| + |u_2(x, t)| < 2M, \\ x \geq 0, \quad kh \leq t \leq (k+1)h, \quad v(x, kh) &= 0. \end{aligned}$$

Consider the domain  $0 \leq x \leq L, kh \leq t \leq (k+1)h$ , where  $L$  is a real number and a function

$$V(x, t) = \frac{4M}{L^2} \left( \frac{x^2}{2} + (I_{kh}(t)) \right)$$

for which

$$\frac{\partial V}{\partial t} = \frac{4M}{L^2}, \quad \frac{\partial V}{\partial x} = \frac{4Mx}{L^2}, \quad \frac{\partial^2 V}{\partial x^2} = \frac{4M}{L^2}$$

and which satisfies the thermal conductivity equation (3.5), as well as

$$\begin{aligned} V(x, kh) &\geq v(x, kh) = 0, \\ V(\pm L, t) &\geq 2M \geq |v(\pm L, t)| \end{aligned} \quad (3.16)$$

For each limited region  $0 \leq x \leq L$ ,  $kh \leq t \leq (k+1)h$ ,  $k \in N$ , the principle of the maximum value is true. The functions  $\underline{u} = -V(x, t)$ ,  $u = v(x, t)$ ,  $\bar{u} = V(x, t)$ , taking into account (3.17), we obtain that

$$-\frac{4M}{L^2} \left( \frac{x^2}{2} + I_{kh}(t) \right) \leq v(x, t) \leq \frac{4M}{L^2} \left( \frac{x^2}{2} + I_{kh}(t) \right). \quad (3.17)$$

We fix  $(x, t)$  and use the fact that  $L$  is arbitrary and can be increased indefinitely. Passing to the limit at  $L \rightarrow \infty$ , we obtain that  $v(x, t) \equiv 0$  for  $x \geq 0$ ,  $kh \leq t \leq (k+1)h$ . Theorem 2 is proved.

Therefore, the following theorem is true.

**Theorem 3.3.** *If  $|\varphi(x, h)| \leq M$ ,  $x \geq 0$ ,  $M > 0$ ,  $h > 0$ , then the solution of problem (3.5), (3.6), (3.7) exists, is unique and is determined by formula (3.13).*

### 3.3. Solving the problems 2 and 3

It is necessary to solve equation (3.5) when the zero initial condition (3.8) and the general boundary condition (3.9) are met. First, let's solve the auxiliary problem of cooling a heated rod, at the boundary of which a constant zero temperature is maintained. Then, for equation (3.5), the Cauchy condition and the boundary condition are given as follows:

$$V_1(x, t_0) = T, \quad v_1(0, t) = 0, \quad x > 0, \quad t > h.$$

Then, according to formula (3.13), we get that

$$\bar{v} = T \operatorname{erf} \left( \frac{x}{2\sqrt{I_{t_0}(t)}} \right), \quad x \geq 0, \quad t > t_0, \quad (3.18)$$

Let  $\mu(t) = \mu_0 \equiv \text{const}$  in condition (3.9). Then, according to (3.18), the function

$$\bar{v} = \mu \operatorname{erf} \left( \frac{x}{2\sqrt{I_{t_0}(t)}} \right), \quad x \geq 0, \quad t > t_0,$$

is a solution of problem (3.5), (3.8), (3.9). Then the function

$$v(x, t) = \mu_0 - \bar{v}(x, t) = \mu_0 \left[ 1 - \operatorname{erf} \left( \frac{x}{2\sqrt{I_{t_0}(t)}} \right) \right], \quad x > 0, \quad t > 0. \quad (3.19)$$

We denote the expression in parentheses of formula (3.19) by  $U(x, I_{t_0}(t))$ , which makes sense when  $t > t_0$ . If for  $t < t_0$  the value of this function is extended by zero, then this definition is consistent with the zero value of the function at  $t = t_0$ . The limit value of this function at  $x = 0$  is a step function equal to zero at  $t < t_0$  and equal to 1 at  $t > t_0$ . The constructed function is often found in applications and is an auxiliary link in constructing the solution to problem 2.

The second auxiliary task is to find a solution of the equation (3.5) under the following conditions:

$$v(x, t_0) = 0, x \geq 0, \\ v(0, t) \equiv \mu(t) = \begin{cases} \mu_0, t_0 < t < t_1, \\ 0, t > t_1. \end{cases}$$

It is directly verified that  $V(x, t) = \mu_0 [U(x, I_{t_0}(t)) - U(x, I_{t_1}(t))]$ ,  $x \geq 0, t > t_0$ . If

$$\mu(t) = \begin{cases} \mu_0, t_0 < t \leq t_1, \\ \mu_1, t_1 < t \leq t_2, \\ \dots\dots\dots, \\ \mu_{n-1}, t_{n-2} < t \leq t_{n-1}, \\ \mu_{n-1}, t_{n-1} < t \leq t_n, \end{cases},$$

and then the solution of the corresponding problem can be written in the form

$$u(x, t) = \sum_{i=0}^{n-2} \mu_i [U(x, I_{t_i}(t)) - U(x, I_{t_{i+1}}(t))] + \mu_{n-1} U(x, I_{t_{n-1}}(t)).$$

Using the theorem on finite increments, we get

$$u(x, t) = \sum_{i=0}^{n-2} \mu_i \left. \frac{\partial U}{\partial t} (x, I_\tau(t)) \right|_{\tau=\tau_i} + \mu_{n-1} U(x, I_{t_n}(t)), \tag{3.20}$$

where  $x \geq 0, t_i \leq \tau_i \leq t_{i+1}$ .

The approximate solution of problem 2 can be obtained by formula (3.20), if replace the function  $\mu(t)$  with a piecewise-constant function.

Heading to the limit when the interval of constancy of the auxiliary function decreases, we obtain that the limit of the sum (3.20) will take the form

$$\int_0^t \frac{\partial U}{\partial t} (x, I_\tau(t)) \mu(\tau) d\tau$$

because when  $x \geq 0$ , we have

$$\lim_{t-t_{n-1} \rightarrow 0} \mu_{n-1} U(x, I_{t_{n-1}}(t)) = 0.$$



If we consider

$$\frac{\partial U}{\partial t}(x, t) = -2 \frac{\partial G}{\partial x}(x, 0, t) = 2 \frac{\partial G}{\partial \xi} \Big|_{\xi=0}$$

then we will get the final result

$$u_2(x, t) = \frac{1}{2\sqrt{\pi}} \int_{kh}^t \frac{x}{[I_\tau(t)]^{3/2}} \times \exp \left\{ -\frac{x^2}{4I_\tau(t)} \right\} \mu(\tau) d\tau, \quad (3.21)$$

$x > 0$ ,  $kh \leq t \leq (k+1)h$ .

The solution of problem 3 using the Green's function (3.14) can be written in the form of a Poisson integral

$$u_3(t, x) = \int_{kh}^t d\tau \int_0^\infty f(y, \tau) G_1(x, y, I_{kh}(t)) dy \quad (3.22)$$

$x > 0$ ,  $kh \leq t \leq (k+1)h$ ,  $k \in \mathbb{N}$ , for the existence of which the function  $f(x, t)$  must be such that the improper integral in formula (3.22) coincides.

So, the following theorem is proved.

**Theorem 3.4.** *The solution of problem (3.5), (3.8), (3.9) is determined by formula (3.22). The solution of problem (3.1), (3.2), (3.3) is determined by formula (3.4), where the terms  $u_1$ ,  $u_2$ ,  $u_3$  are the solutions of problems 1, 2 and 3 respectively.*

The first, second and third initial-boundary problems for the heat conduction equation with inversion of the argument and  $D(t) \equiv a^2 > 0$ , constant are considered in [9], [10], [11].

### References

1. T.N. NARASIMHAN, *Fourier's heat transfer equations: History, influence and connections*, *Review ws of Geophysics*, **37** (1) (1999), 151–172.
2. L.C. EVANS, *Partial Differential Equations*, *Grad. Stud. Math.*, Amer. Math. Soc., Providence, RI, **19** (2010), 44–65.
3. IAN N. SNEDDON, *Elements of Partial Differential Equations*, *Dover Books of Mathematics*, 2006, 352 p.
4. A.B. TAYLER, *Mathematical Models in Applied Mechanics*, Oxford, 2008, 288 p.
5. I. KAUR, Y. MISHIN, W. GUST, *Fundamentals of Grain and Interphase Boundary Diffusion*, John Wiley & Sons Ltd, Chichester, 1995, 512 p.
6. J. BIAZAR, Z. AYATI, *An Approximation to the Solution of Parabolic Equation by Adomian Decomposition Method and Comparing the Result with Crank-Nicolson Method*, *International Mathematical Journal*, **39** (2006), 1925–1933.
7. COLE KEVIN D., BECK JAMES V., HAJI-SHEIKH A., LITKOUHI BAHAN, *Heat conduction using Green's functions*, *Series in Computational and Physical Processes in Mechanics and Thermal Sciences (2nd ed.)*, Boca Raton, FL: CRC Press, 2011.
8. R.K.M. THAMBYNAYAGAM, *The Diffusion Handbook: Applied Solutions for Engineers*, McGraw-Hill Professional, 2011.

9. Y.M. DRIN, I.I. DRIN, S.S. DRIN, Y.P. STETSKO, *The first boundary value problem for the nonlinear equation of heat conduction with deviation of the argument*, in *Proc. The 12th International Conference on Electronics, Communications and Computing*, 20-21 October, 2022, Chisinau, Republic of Moldova.
10. Y.M. DRIN, I.I. DRIN, R.Y. DRIN, *The analytical view of solution of the second boundary value problem for the nonlinear equation of heat conduction with deviation of the argument*, in *Proc. The Eleven International Conference on "Informatics and computer technique problems" (PICT - 2022)*, 10-13 October 2022, Chernivtsi, Ukraine, 11–18.
11. YA. DRIN, I. DRIN, R. DRIN, *The third initial-boundary value problem for the nonlinear equation of heat conduction with deviation of the argument*, 2022 *International Conference on Innovative Solutions in Software Engineering (ICISSE)*, Vasyl Stefanyk Precarpathian National University, Ivano-Frankivsk, Ukraine, Nov. 29-30, 2022, 282–285

*Received 10.08.2023*