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# A Physical Optics Simulator for Multireflector THz Imaging Systems

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Abstract—This work presents a Physical Optics-based simulator for the analysis of THz imaging systems. The simulation starts by calculating the electromagnetic interactions inside the multireflector system and the incident field that the focusing system creates on the target under inspection. In a second step, the electric field that the modeled target scatters back to the system receiver, is also calculated. This allows to predict the imaging behavior of the system for different targets before manufacturing. The simulator results are validated by using measurements from an existing 300 GHz standoff imaging system. This contribution aims to help in the development of better imaging systems for security applications in the near future.

*Index Terms*—Reflector antennas, submillimeter wavelength imaging, inverse scattering, radar imaging, physical optics.

#### I. INTRODUCTION

In recent years interest in personal security applications has markedly grown. In this context, active millimeter wave radar has been proved to provide excellent capabilities for the detection and imaging of concealed threat or contraband at security checkpoints [1]-[3]. In particular, the use of active submillimeter (THz) imaging systems is of great interest for the challenging problem of standoff observation. Different systems to achieve through-clothing imaging at standoff ranges of tenths of meters have been presented [4]–[7]. These systems are capable of penetrating opaque materials, preserve the privacy of the explored target and achieve a good image resolution. The area of interest is scanned in a pixel by pixel imaging strategy by using multireflector systems and mechanical fast beam scanning [8]. The theoretical and simulated focusing properties of these kind of systems have been profusely studied in the literature [9], [10]. However, these studies are mainly limited to the properties of the system in transmission, while

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the response of the imaging to different targets has been mostly based on small sets of measurements.

However, the development of better and more complex THz systems for security applications demands the use of reliable and efficient simulation tools to improve the design and optimization of novel radar configurations before manufacturing. Also, these simulators are of great interest to obtain datasets of realistic examples for automatic target recognition and allow for the use of model-based inversion algorithms [11] to improve the detection and identification abilities.

In [12] a thorough review of the available simulators for these kind of THz systems is presented. It concludes that most of the simulators for concealed object detection are based on ray-tracing and corrected using different scattering models as Kirchoff [13] or bidirectional reflectance distribution function (BRDF) [14].

In this work a Physical Optics (PO) based simulator for THz imaging systems is presented. PO [15], [16] is a wellknown asymptotic high frequency computational technique traditionally used to analyze the behavior of reflector antennas in terms of gain, secondary lobe level or cross-polar radiation. However, the use of PO to analyze imaging systems at THz frequencies is very challenging due to the large electrical size of the reflector surfaces and the repeated number of simulations that are needed to simulate the pixel-by-pixel operation at multiple frequencies. We have developed an inhouse simulator to overcome these limitations. In the simulator PO is used not only to predict the incident field that the focusing system creates on the target under inspection, but also to calculate the electric field that the target scatters back to the receiver through the reflector system. It could be integrated in an optimization framework to design imaging systems according the required specifications in terms of obtained field of view, resolution, dynamic margin, etc. The aforementioned optimization concept refers to design parameters, such as the geometry of each of the elements of the system. In this work, the simulation results are compared with measurements made by a system that already exists, in order to verify the robustness of the simulator. We believe the simulator is specially useful to verify the behavior of the system in terms of resolution, Field of View and other important parameters before fabrication, and also to detect possible problems once the system is manufactured. To our knowledge, these features are not available in a straightforward manner in current commercial simulation software, which is mainly devoted to the design and simulation of antennas in terms of their main transmission parameters.

The main contributions of this work are:

i) a PO-based simulator for THz imagers in both transmission and reception, including nearfield interactions between the reflector system and the target under test;

ii) qualitative and quantitative comparisons between simulated and measured results.

The structure of the paper is as follows. The imaging system that is simulated and used to validate the obtained solutions is described in Section II. The details of the implemented PO approach are presented in Section III. The simulation process is described in Section IV. The obtained results are shown in Section V and finally conclusions are provided in Section VI.

#### II. SIMULATED 300 GHZ IMAGER

The imaging system used to validate the proposed PO approach is based on a Bifocal Elliptical Gregorian Reflector System (BEGRS). The system is capable of generating high resolution images in real time at a standoff distance of 8 m and with a field of view (FoV) of 90 cm  $\times$  50 cm, focusing a 300 GHz beam [17]. A 27 GHz bandwidth is used in the system to ensure good range resolution.

The baseline design consists of two opposite-faced reflector antennas with a common focus. This configuration is similar to the offset confocal elliptical reflector (CEGRS) [9], but their surfaces are shaped to have a better scan capability. Fig. 1 shows the multireflector system configuration.



Fig. 1. Multireflector-based imaging system at 300 GHz used to validate the proposed PO simulator.

The principal elements of this configuration are [12]:

- (a) A main reflector aperture which focuses the 300 GHz beam to the target plane.
- (b) A parabolic subreflector that radiates the field from the mirror to the main reflector.
- (c) A rotating flat mirror placed between the feed reflector and the subreflector which is used to achieve mechanical beam scanning in both elevation and azimuth.
- (d) An auxiliar parabolic feed reflector to achieve the collimated beam to feed the subreflector.
- (e) A beamsplitter to separate the transmission (Tx) and reception (Rx) paths.

With the BEGRS it is possible to focus a 300 GHz beam at a standoff distance of 8 meters with a cross-range resolution  $\Delta CR = 1.6$  cm and a range resolution  $\Delta R = 1$  cm. The system operates as a monostatic active radar and the target is placed facing the main reflector at the standoff distance. The auxiliar parabolic reflector is fed by a horn antenna placed at its focus. The proposed simulator uses PO to calculate the incident field created by the feed (modeled as a cos-q) and the nearfield interactions between the reflector surfaces in both transmission (from transmitter to reflectors to target) and reception (from target to reflectors to receiver). The beam splitter effects are not considered in the simulation at this moment and both transmitter and receiver are simulated as placed at the same position

#### **III. PHYSICAL OPTICS APPROACH**

The developed code is based on the computation of the induced currents between two reflecting surfaces. First, the surfaces are imported as CAD models into the program and modeled as a set of triangular patches (see Fig. 2) with three points  $\vec{R}_{m_1}$ ,  $\vec{R}_{m_2}$  and  $\vec{R}_{m_3}$ , that define a central point  $\vec{R}_m = \frac{\vec{R}_{m_1} + \vec{R}_{m_2} + \vec{R}_{m_3}}{3}$  and two edge vectors  $\vec{v}_{m_{12}} = \vec{R}_{m_2} - \vec{R}_{m_1}$  and  $\vec{v}_{m_{13}} = \vec{R}_{m_3} - \vec{R}_{m_4}$ . From these vectors, the normal is calculated as  $\hat{n}_m = \frac{\vec{v}_{m_{12}} \times \vec{v}_{m_{13}}}{\|\vec{v}_{m_{12}} \times \vec{v}_{m_{13}}\|}$  and the area as  $A_m = \frac{\|\vec{v}_{m_{12}} \times \vec{v}_{m_{13}}\|}{2}$ .

The process starts by calculating the incident magnetic field  $\vec{H}_m$ , the incident electric field  $\vec{E}_m$  and the incident Poynting vector  $\hat{p}_{i_m}$  on the first reflector surface created by a feed modeled as a cos-q. The value of the parameter q is adjusted as a function of the actual horn antenna used to feed the reflector system. Once these three vectors are calculated, the induced currents in the next reflector surface are calculated as follows.

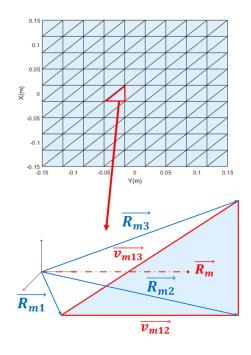


Fig. 2. Triangular mesh and vector notation used in the implementation of the PO simulation.

First, the electric current  $\vec{J}_m$  is calculated using the incident magnetic field at the center of each triangle:

$$\vec{J}_m = 2\hat{n}_m \times \vec{H}_m \tag{1}$$

PO nearfield equations are used to calculate magnetic and electric field  $\vec{H}_n$ ,  $\vec{E}_n$  at a point  $\vec{R}_n$  on the next reflector surface. The position vector between the triangles in both surfaces is  $\vec{R}_{mn} = \vec{R}_n - \vec{R}_m$ , the distance  $R_{mn} = \|\vec{R}_{mn}\|$  and the direction  $\hat{R}_{mn} = \frac{\vec{R}_{mn}}{\vec{R}_{mn}}$ .

$$\vec{H}_n = \sum_m \vec{H}_{mn} \tag{2}$$

$$\vec{H}_{mn} = \frac{e^{-jkR_{mn}}}{4\pi R_{mn}^3} \left(1 + jkR_{mn}\right) \left(\vec{J}_m \times \vec{R}_{mn}\right) I_{mn}$$
(3)

$$I_{mn} = 2A_m e^{-j\frac{\alpha_{mn}+\beta_{mn}}{3}} \left[\gamma_{mn}\right] \tag{4}$$

where

$$\gamma_{mn} = \left[\frac{\alpha_{mn}e^{j\beta_{mn}} - \beta_{mn}e^{j\alpha_{mn}} + \beta_{mn} - \alpha_{mn}}{(\alpha_{mn} - \beta_{mn})\,\alpha_{mn}\beta_{mn}}\right] \quad (5)$$

$$\alpha_{mn} = k \vec{v}_{m_{12}} \cdot \left( \hat{R}_{mn} - \hat{p}_{i_m} \right) \tag{6}$$

$$\beta_{mn} = k\vec{v}_{m_{13}} \cdot \left(\hat{R}_{mn} - \hat{p}_{i_m}\right) \tag{7}$$

This process is repeated for all the couples of triangles in both reflector surfaces. To obtain the Poynting vector  $\hat{p}_{i_n} = Re\{\frac{\vec{E}_n \times \vec{H}_n^*}{\|\vec{E}_n \times \vec{H}_n^*\|}\}$  it is necessary to calculate also the electric field:

$$\vec{E}_n = \sum_m \vec{E}_{mn} \tag{8}$$

$$\vec{E}_{mn} = -\frac{j\eta e^{-jkR_{mn}}}{4\pi kR_{mn}^3} \left[ \vec{J}_m \left( -1 - jkR_{mn} + (kR_{mn})^2 \right) + \right]$$

$$+ \hat{R}_{mn} \left(3 + 3jkR_{mn} - (kR_{mn})^2\right) \left(\hat{R}_{mn} \cdot \vec{J}_m\right) \left[I_{mn} \quad (9)\right]$$

If the distance between both surfaces are much greater than the wavelength, equations (3) and (9) can be simplified as

$$\vec{H}_{mn} = \frac{j}{2\lambda} \frac{e^{-jkR_{mn}}}{R_{mn}} \left( \vec{J}_m \times \hat{R}_{mn} \right) I_{mn} \tag{10}$$

$$\vec{E}_{mn} = \eta \vec{H}_{mn} \times \hat{R}_{mn} \tag{11}$$

This process is sequentially followed for all the reflector surfaces until the induced electric currents on the main reflector are obtained. Next section describes how the interaction with the target under test is obtained.

#### IV. SIMULATION METHOD

#### A. Simulation process

As stated before, the imaging system is based on a pixelby-pixel scanning of the target area placed at a prescribed distance from the antenna. Such scanning is obtained by the rotation and tilting of a flat mirror. In the simulation, a rectangular grid in the target area is defined and, by using a polynomial approximation, each pixel position is mapped to the corresponding rotation and tilting angles of the mirror [8]. For each position of the mirror the induced currents on its surface are calculated according to the equations in previous section. The process is repeated until the electric currents on the main surface for steering the beam at each particular position on the target plane are calculated ( $\vec{J}_2$  in Fig. 3). To complete the simulation of the system in transmission mode, the incident field on the target under test, modeled as a triangular mesh, is computed by using again the same equations described in Section III.

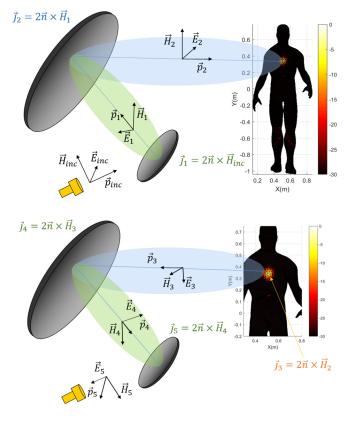


Fig. 3. Description of the PO-based simulation process in transmission (top) and reception (bottom).

The simulation of the imaging system response for the modeled target starts by calculating the induced electric currents on the target,  $\vec{J}_3$ , from the induced magnetic field  $\vec{H}_2$ . This part of the process models the different response to the imaging system of different areas of the target depending on their shape, which is a critical factor in these kind of systems [12]. The calculated currents are used to radiate back to the main reflector surface where a new set of currents  $\vec{J}_4$  are induced. These currents are used to sequentially radiate to the next reflector surfaces until the electric field at the receiving point (the focal point of the auxiliar reflector) is calculated according to Section III.

This process is independently repeated for multiple equally spaced frequencies in the radar bandwidth according the desired non-ambiguous range. The simulation process ends when all the pixels on the target plane have been scanned. The output of the program is a matrix with as many rows as pixels and as many columns as frequencies.

#### **B.** Simulation optimization

For a large number of pixels and frequencies, the simulation process described above can take a substantial amount of time. Two different approaches have been used to reduce the computational time of the simulator. First, if the triangles of the simulated surfaces are small in terms of wavelength Equation (4) can be simplified as [18]–[20].

$$I_m \approx A_m \quad , \quad \forall n \tag{12}$$

This approximation is used in the radiation between the main reflector and the target being simulated, and it reduces the total computation time about one third.

Second, for further time reduction, the PO codes have been implemented using a GPU through the set of libraries provided by Matlab. Efficient implementation of the code in GPU requires to create two intermediate matrices in order to reduce the number of calculations:

$$c_{mn} = \frac{I_{mn}e^{-jkR_{mn}}}{R_{mn}} \tag{13}$$

with  $I_{mn}$  obtained from expression (4) or (12), and

$$\vec{a}_{mn} = c_{mn} \vec{J}_m \times \hat{R}_{mn} \tag{14}$$

$$\vec{H}_n = \frac{j}{2\lambda} \sum_m \vec{a}_{mn} \tag{15}$$

$$\vec{E}_n = \frac{j\eta}{2\lambda} \sum_m \vec{a}_{mn} \times \hat{R}_{mn} \tag{16}$$

The GPU is used through the matlab arrayfun tool, even knowing that full performance is not obtained. The use of CUDA in this application is not possible due to the size of the matrices that must have all the same size. On the other hand, in Matlab is possible to mix vectors with matrices, using significantly less memory than using CUDA.

The computation time of each pixel at a single frecuency is 3 seconds.

#### C. Signal processing for imaging

As explained in Section I, the imaging system uses a pixelby-pixel observation strategy. For each observation pixel, the system captures the received field for all studied frequencies, what is defined as burst. Each of the burst is stored in a received data matrix whose rows correspond to the number of bursts and the columns with the number of used frequencies. By applying the Inverse Fast Fourier Transform (IFFT) to each burst, High Resolution Synthetic Range profiles (HRR) are achieved, building a high resolution profile matrix whose columns correspond to the range cells [21].

Range images are created by locating the position of the maximum value of each HRR and the reflectivity ones are achieved by using that maximum HRR value, as shown in Fig. 4. For its representation, the reflectivity image is normalized by its maximum value.

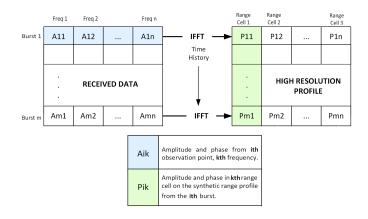


Fig. 4. Description of the imaging algorithm applied to the simulation results.

#### V. VALIDATION AND RESULTS

In this Section, to validate the performance of the presented simulator, the simulation results are compared to measurements obtained using the imaging system described in Section II. Two different targets placed at a standoff distance of 8 m have been considered. The first case is a metallic flat plate which serves as a canonical example to study the response of the imaging system to tilted surfaces. In the second example a female mannequin torso is used to evaluate the range and shape accuracy obtained by using the simulation results. A bandwidth of 27 GHz and an observation pixel size of 1.5 cm have been used. Additive Gaussian noise 40 dB below the maximum signal level has been added to the obtained simulation results before creating the images. Imaging is based on applying an IFFT to the obtained frequency response at each pixel.

The simulations have been performed using the hardware described below:

- $2 \times Intel^{\textcircled{B}}$  Xeon<sup>B</sup> E5-2690 v2 @3GHz processors.
  - RAM: 512 GB
  - Number of Cores: 10 (2 logical cores per physical)
- + 1  $\times$  NVIDIA GeForce GTX TITAN X
  - Memory: 12 GB (GDDR5 SDRAM)
  - Number of CUDA Cores: 3072

#### A. Example 1: metallic flat plate

One of the main purposes of the presented simulator is to study the effect of the target shape on the obtained images. Since, at these frequencies, the target response is mainly given by the specular reflection of the incident beam, these type of configurations can misrepresent certain areas of the target if the response of these areas is below floor noise. This first experiment aims to study this effect.

The measured object is a 47 cm  $\times$  18 cm flat metal plate placed in front of the imaging system. The plate is first measured with its normal pointing to the center of the main reflector, which corresponds to the specular reflection direction. Then, it is rotated around the vertical axis in increments of 1 degree.

For the simulation, a CAD model of the plate is used. An observation grid of 47 cm  $\times$  18 cm in the target area is defined and discretized every 1.5 cm. The imaging system is synthetically pointed to each pixel as described in Section IV-A. Fig. 5 shows the measured and simulated results for different tilt angles. In these plots, the maximum amplitudes of the IFFT of the measured and simulated signal for each pixel are represented. As the plate rotates, the amount of received power decreases rapidly, due to the fact that the antenna is no longer fully capturing the specular reflection from the plate. By comparing the results of the measurements and the simulations a high concordance can be verified.

- Figs. 5(a) and 5(b) show a great fit. In the simulated image it can be seen that the similarity in the central part of the plate is very high, but closer to the corners, the power level received in measurements is slightly lower. These looseness arise due to inaccuracies of position (in practice the target position adjustment at the standoff distance is a really complex task).
- Figs. 5(c), 5(d), 5(e) and 5(f), as in the previous case, show a great fit but although the adjustment is very high in the area of the maximum, due to small inaccuracies in the real position of the target, when moving axially the power drop is slightly higher in the measurements compared to the simulations.
- In Figs. 5(g) and 5(h) apart from observing a good fit between both responses, an edge effect in simulation can be appreciated. This effect is due to the inacuracy of PO simulations when there is an edge highly iluminated.

Fig. 6 shows the difference between the images recovered by measurement and those recovered by simulation. The difference between the two images is calculated by subtracting the simulated image and the measured image (both normalized by their maximum). In certain areas of the recovered images there are higher error values (over 12 dBs). In the lower left corner of the Figs. 6(a) and 6(c) the high error value observed is due to mismatches in the positioning of the OUT. On the left side of the Fig. 6(b) and on the right side of the Fig. 6(d), high error values are also observed, as already explained above, this effect occurs due to the inacuracy of PO simulations when there is an edge highly iluminated.

To quantify the correlation between the simulations and the measurements, the mean-squared error (MSE) for each of the studied metal plate positions, was calculated. Table I shows the obtained results using the following expression:

$$MSE = \frac{1}{n} \sum_{i=1}^{n} (i_{sim} - i_{meas})^2$$
(17)

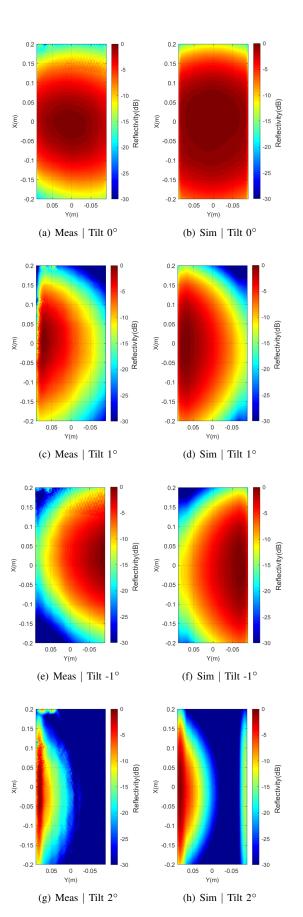


Fig. 5. Measured (Meas) VS Simulation (Sim) reconstructed images. Reflectivity in dB relative to the maximum of each image for different tilt angles of the metallic plate is depicted.

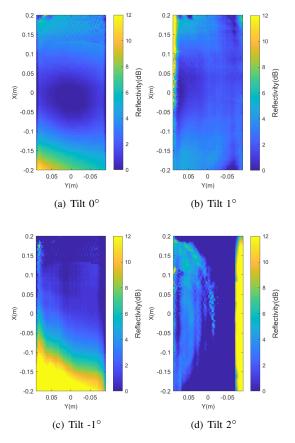


Fig. 6. Subtraction between simulation and measurement recovered images for the different analized tilt angles.

being *n* the number of elements of each of the images,  $i_{sim}$  the elements of the simulation recovered image and  $i_{meas}$  the elements of the measured recovered image. For the MSE calculation, the value of each pixel of the images are expressed in natural units.

Tilt (°)	MSE (dB)
0	-27.79
1	-34.39
-1	-31.03
2	-43.36
TABLE I	

MSE VALUE IN DB FOR EACH OF THE TILTED PLATE CASES UNDER STUDY.

To further explore this specular reflection behavior, the measured and simulated received power at the imaging system receiver as a function of the angle is depicted in Fig. 7. In this case, only the power received from the central pixel, placed at (X = 0, Y = 0) is considered. There is a good agreement between the measured and simulated results for small tilting angles, and in both cases it can be observed how quickly the received power decreases when the plate is no longer pointing to the main reflector center.

#### B. Example 2: Human body torso

In this second example, a more realistic case is considered. The target of this experiment was a female mannequin placed

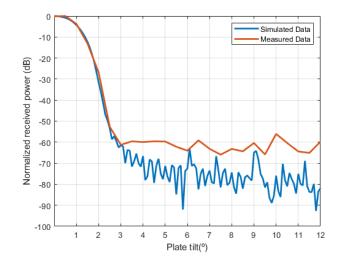


Fig. 7. Normalized power obtained at the receiver as a function of the tilt of the metallic plate for the central pixel placed at (X = 0, Y = 0) for measurements and simulation.

once again at 8 m standoff from the imaging system. This target presents a more complex geometry due to the curvature of the human body. As previously noted (see Section V-A), the response of the system gets much worse when the response of the OUT moves away from the direction of the specular reflection. For this reason, this example is very interesting, because it allows analyzing the performance of the system in more complicated and realistic conditions.

In order to accomplish the simulation of the same target, a 3-D model of the mannequin shown in Fig. 8 was obtained using a consumer infrared structured light depth sensor that is moved around the surface of the mannequin [22]. Fig. 9(a) shows the CAD model imported into the simulator. In Fig. 9(b) the same mesh is coded in range from the antenna in order to compare the obtained range results using both measurements and simulations. The simulations are performed in a 90 cm  $\times$  50 cm observation grid in the target area. According to the results in [17], at these frequencies the plastic material of the mannequin reflects most of the incident electric field which makes the approximation good enough in terms of relative received power and range estimation.

Range images obtained using simulations and measurements are presented in Fig. 10. Range is obtained from the corresponding peak of the IFFT for both the measurement and simulation of each scanned pixel in the target area. Fig. 10(a) presents results for the simulated case while Fig. 10(b) depicts the measured one. The agreement between the measurements and simulations is good, and both results match well the expected range profile of the actual mannequin as depicted in Fig. 9(b).

#### VI. CONCLUSION

A Physical Optics-based simulator for the analysis of the THz imaging systems based on multireflector configurations has been presented. The simulator is able not only to model the incident field that the focusing system creates on the target



Fig. 8. Picture of the measured female mannequin used to validate the PO simulator.

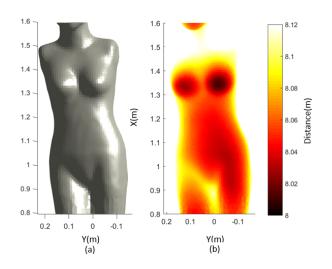


Fig. 9. (a) Imported CAD model of the mannequin; (b) Model coded in depth from the antenna: each pixel represents the distance from the central point of the main reflector to the center of each facet of the model.

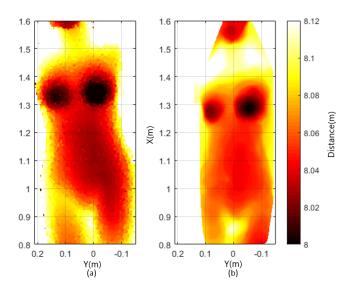


Fig. 10. Depth map (distance from the center of the main reflector to each of the pixels of the imaging plane obtained (a) from simulations and (b) from measurements.

under inspection, but also to calculate the electric field that the target, imported as a CAD model, scatters back to the system receiver. The simulation results have been compared with measurements from an existing real prototype to validate the results of the simulator and good agreement has been obtained. This work is a step forward to improve the design and optimization of novel imaging configurations based on THz radar. Further work will be related with the simulation of targets composed of different materials (e.g. dielectrics).

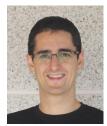
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