

An atom probe perspective on phase separation and precipitation in duplex stainless steels*

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Abstract

Three-dimensional chemical imaging of Fe–Cr alloys showing Fe-rich (α)/Cr-rich (α') phase separation is reported using atom probe tomography techniques. The extent of phase separation, i.e., amplitude and wavelength, has been quantitatively assessed using the Langer–Bar-on–Miller, proximity histogram, and autocorrelation function methods for two separate Fe–Cr alloys, designated 2101 and 2205. Although the 2101 alloy possesses a larger wavelength and amplitude after annealing at 427 °C for 100–10 000 h, it exhibits a lower hardness than the 2205 alloy. In addition to this phase separation, ultra-fine Ni–Mn–Si–Cu-rich G-phase precipitates form at the α/α' interfaces in both alloys. For the 2101 alloy, Cu clusters act to form a nucleus, around which a Ni–Mn–Si shell develops during the precipitation process. For the 2205 alloy, the Ni and Cu atoms enrich simultaneously and no core–shell chemical distribution was found. This segregation phenomenon may arise from the exact Ni/Cu ratio inside the ferrite. After annealing for 10 000 h, the number density of the G-phase within the 2205 alloy was found to be roughly one order of magnitude higher than in the 2101 alloy. The G-phase precipitates have an additional deleterious effect on the thermal embrittlement, as evaluated by the Ashby–Orowan equation, which explains the discrepancy between the hardness and the rate of phase separation with respect to annealing time (Gladman T 1999 *Mater. Sci. Tech. Ser.* **15** 30–36).

Keywords: atom probe tomography, thermal embrittlement, Cu cluster, G-phase, Fe–Cr alloy, phase separation

(Some figures may appear in colour only in the online journal)

1. Introduction

Inexpensive Fe–Cr based duplex stainless steels (DSSs), composed of ferrite and austenite, are considered promising structural materials because of a combination of good mechanical properties and corrosion resistance [2, 3]. These properties have enabled DSSs for potential applications in power generation systems. However, long-term exposure to service temperatures between 300 °C and 538 °C leads to severe thermal embrittlement [4–6]. Several possibilities

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Table 1. Microhardness of alloy 2101 and 2205 in the ferrite grains for different aging times at 427 °C. The test was performed with a 10 gf load.

Hardness (HV)	1 h	100 h	1000 h	10 000 h
2205 alloy	265.1 ± 5.0	291.4 ± 10.5	404.0 ± 34.4	469.7 ± 72.4
2101 alloy	210 ± 10.3	209 ± 8.5	271 ± 13.4	335 ± 17.9

associated with the nanostructure and chemical changes within ferrite have been considered as the origin of the embrittlement. Firstly, the ferrite decomposes into Fe-rich (α) and Cr-rich (α') phases due to the miscibility gap of the Fe–Cr system. The Cr concentration and spacing of the α' phase increases concurrently with the aging time [4, 7, 8], leading to embrittlement. Secondly, Ni–Mn–Si-enriched G-phase prefers to form on the α/α' phase boundaries [8–10]. Previous transmission electron microscopy experiments pinpointed that G-phase has a face centered cubic structure with a complex unit cell of 116 atoms and a lattice parameter between 1.09 and 1.11 nm [10]. However, it is still under debate as to whether the G-phase plays an additional role in embrittlement [11]. A quantitative assessment on these two aspects is needed to tailor the alloying, design, and reliable life-span prediction of DSS. In this work, an atomically resolved chemical imaging technique, atom probe tomography (APT), was used to quantify the extent of α/α' phase separation and G-phase formation to better understand the origin of thermal embrittlement of Fe–Cr alloys.

2. Methods and experimental procedure

2.1. Material preparation

Two plates of the two Fe–Cr alloys, 2101 and 2205, with a thickness of 2.54 cm were hot-rolled, solution-annealed above 1038 °C, and water quenched. A detailed chemical composition has been reported elsewhere [12]. After the production of the steel, five samples from each alloy were cut and aged isothermally in an air furnace at a temperature of 427 °C for 1, 100, 1000, and 10 000 h, followed by air cooling. The desired service temperatures for DSSs in power plant applications are near 300 °C. Due to the sluggish nature of the phase transformation accelerated tests are necessary. The temperature of 427 °C was intentionally chosen to be above the service conditions to expedite the phase separation, so that extreme aging times were not required. Furthermore, this temperature was chosen to ensure it was below the 475 °C nose of the embrittlement curve so that phase transformation mechanisms would be the same as lower temperature applications [5]. To perform the metallographic analysis and distinguish ferrite from austenite, the samples were mechanically polished and etched in 40% NaOH for 10 s. Vicker's microhardness testing was performed in the ferrite grains with a 10 gf load. The average hardness values for each sample with respect to aging time are shown in table 1. Clearly, after annealing at 427 °C for 100 h, alloy 2205 demonstrates significantly more hardening than the 2101 alloy. Though microhardness testing does not provide precise embrittlement behavior, it has been shown

that microhardness changes with thermal aging correlate well to the loss of toughness seen in Charpy impact toughness testing [5, 6, 12]. APT was performed to rationalize a clear structure–property relationship within the ferrite for these two alloys.

2.2. Atom probe tomography

APT offers both three-dimensional (3D) imaging with sub-nm resolution and chemical identity of single atoms. APT can be used to quantify parameters such as the wavelength (λ) and the change in Cr concentration (ΔCr) between the α/α' phases that traditional high-resolution techniques can not resolve due to its ability to image all chemicals in 3D with equal probability. It is impossible to trace the structural evolution of interconnected ferrite by techniques such as electron backscatter diffraction (EBSD) because the initial stage of phase decomposition of the Fe–Cr alloys occurs at the nanoscale, as shown in figure 1(a). Using a FEI Nova 200 focused ion beam (FIB) system and micromanipulators (figure 1(b)), the ferrite region of interest was extracted as a wedge, and mounted to a Si post array. A series of annular milling patterns were used to further sharpen the specimen into a needle shape with a diameter of less than 100 nm (figure 1(d)) [13]. After a 5 kV low voltage milling step, the Si coupon with the sharpened needles was then transferred into the atom probe analysis chamber at a base temperature of 50 K. The needle was then positioned within $\sim 40\ \mu\text{m}$ of the local electrode, in which a standing and pulsed voltage are induced to enable the field evaporation of ions at specific moments in time, such that a time-of-flight mass spectrum can be produced. For these experiments, the voltage was pulsed with a frequency of 200 kHz, such that the atoms at the needle surfaces had enough energy to be pushed over the Schottky hump and were field evaporated and projected onto a position sensitive detector consisting of a microchannel plate and a cross-delay line. The evaporation rate used was $\sim 0.5\%$. This setup allows for detection of both the time-of-flight of each ion (resulting in the mass-to-charge ratio), and their initial x – y position with respect to the needle tip shape. The z -position is determined by the sequence of detection, which allows for an atom-by-atom 3D reconstruction of the decomposed ferrite with sub-nm resolution. The APT experiments were conducted using a CAMECA Instruments LEAP 4000X HR. At least 5 million ions were collected for each sample to ensure adequate data statistics. The datasets were reconstructed and analyzed using the IVAS 3.6.8 software (CAMECA Instruments) [14].

Table 2. The average composition of ferrite grains analyzed by APT.

Atomic %	Fe	Cr	Ni	Mn	Cu	Si	Mo	Al	P	C
2101 alloy	68.12 ± 0.05	23.69 ± 0.05	1.13 ± 0.01	4.48 ± 0.02	0.19 ± 0.01	1.94 ± 0.01	0.14 ± 0.01	0.03 ± 0.002	0.07 ± 0.003	0.06 ± 0.003
2205 alloy	64.07 ± 0.04	25.76 ± 0.04	4.22 ± 0.02	1.65 ± 0.01	0.24 ± 0.01	1.08 ± 0.01	2.11 ± 0.02	0.004 ± 0.001	0.08 ± 0.003	0.08 ± 0.002

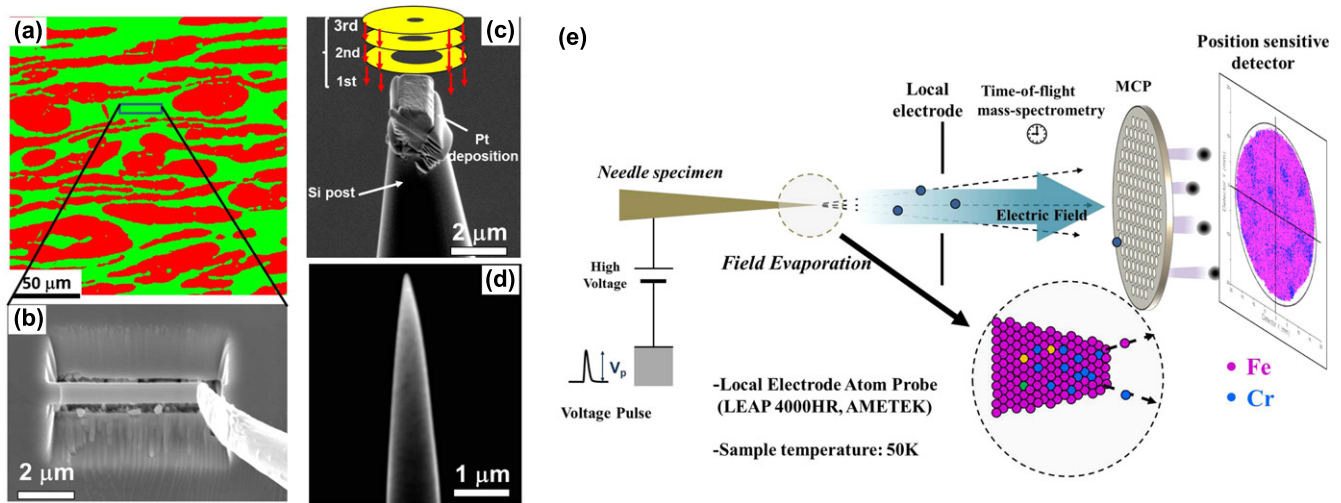


Figure 1. Phase-specific APT tip preparation in the ferrite phase in Fe–Cr alloys: (a) EBSD phase map showing the ferrite (green) and austenite (red); (b) SEM image showing the lift-out procedure by using the micromanipulator in the FIB; (c) series of annular milling steps by FIB; (d) final shape of APT needle after 5 kV ion cleaning; (e) schematic description of a local electrode atom probe.

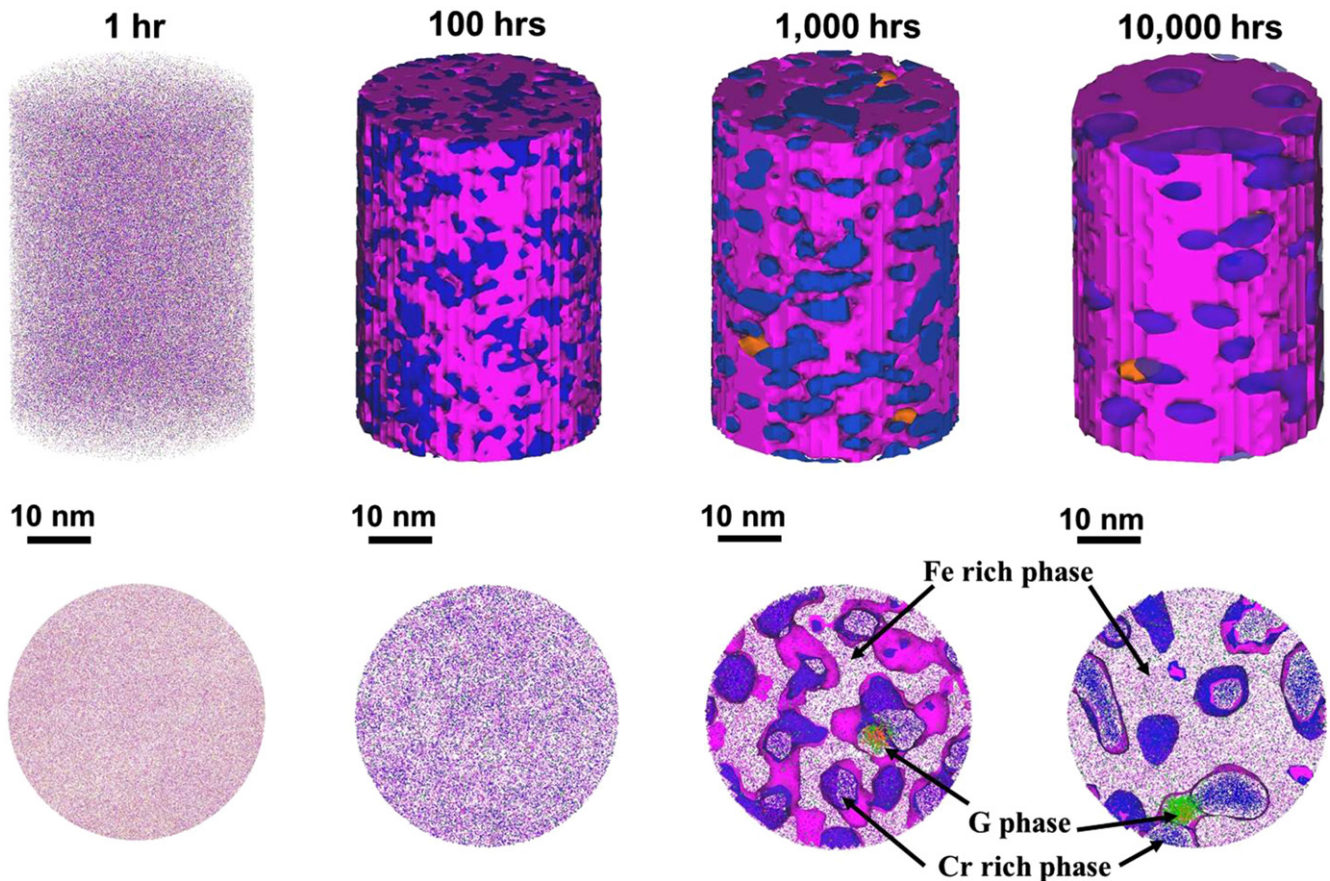


Figure 2. Evolution of phase separation as a function of annealing time in 2101 DSS. APT shows the main metallic elements distributions: Fe (pink), Cr (blue), Ni (green) and Cu (orange). The isoconcentration surfaces of 20 at% Cr (blue) and 65 at% Fe (pink) highlights the location of the α and α' phases. Also shown are 10 at% Cu isosurfaces identifying the location of small G-phase precipitates. The lower images were 2.5 nm slices extracted from the reconstruction volume.

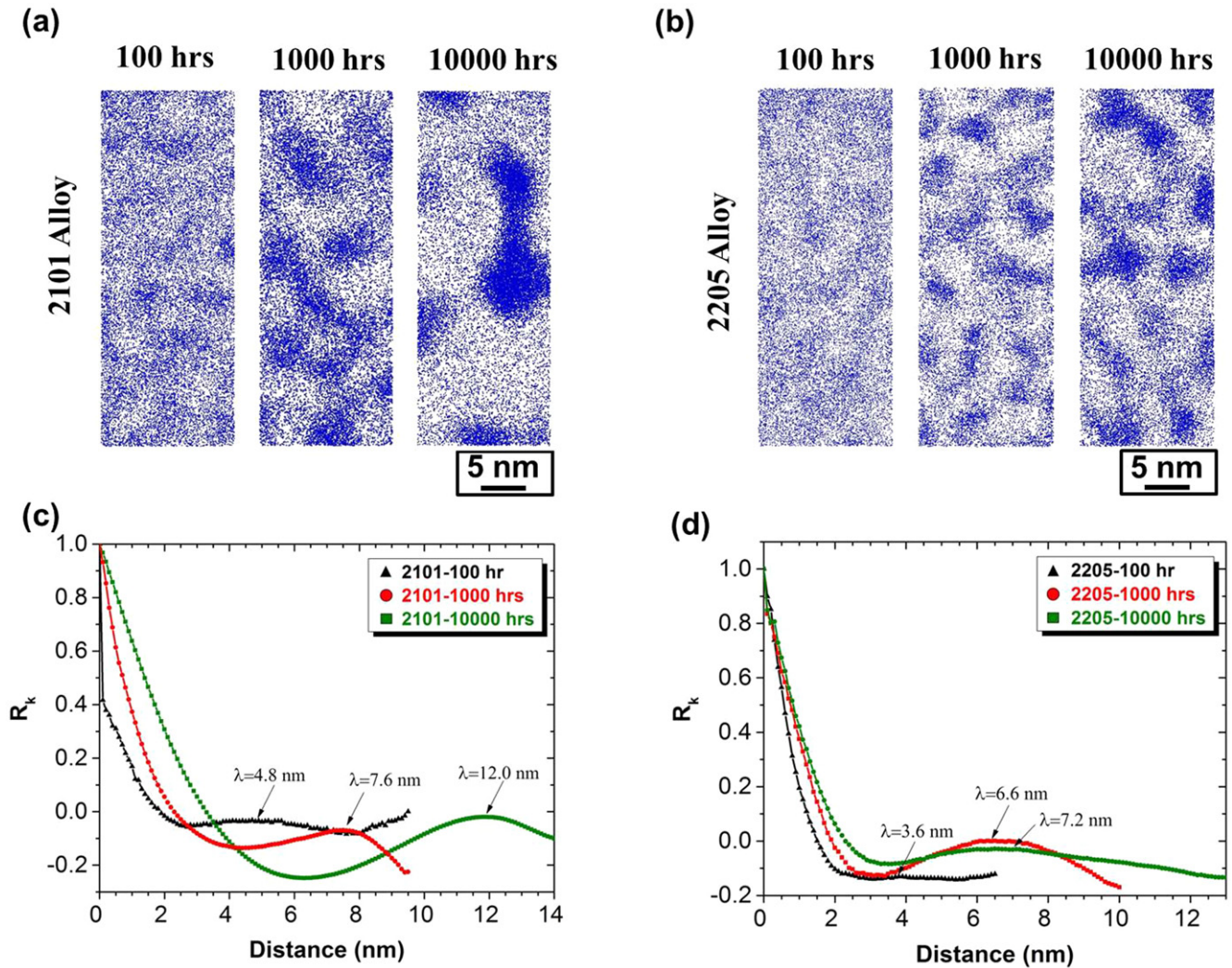


Figure 3. $40 \times 15 \times 5 \text{ nm}^3$ Cr atom maps for the 2101 alloy (a) and 2205 alloy after aging at 427°C for 100, 1000, and 10 000 h. The corresponding autocorrelation function analysis is shown in (c) and (d).

3. Result and discussions

3.1. Nanostructure evolution during phase decomposition

The APT reconstruction of the fine scale microstructure evolution during the phase separation process (1, 100, 1000, and 10 000 h anneals at 427°C) is shown in figure 2. Fe, Cr, Cu, and Ni atoms are presented in pink, blue, orange, and green colors, respectively. The other elements are omitted for clarity. The lower images in figure 2 correspond to 2.5 nm thick slices extracted from the reconstruction, elucidating the nanoscale structure evolution from a rather homogeneous elemental distribution to a clear α/α' phase separation. Rather than having an interconnected α' phase [15], the α' phase forms isolated particles at 10 000 h due to the relatively low Cr content of $\sim 23.7\%$. In addition, Ni- and Cu-enriched phase formed in the vicinity of the α and α' interfaces. The extent of the phase separation process is quantitatively evaluated using several different techniques described below that have their advantages and disadvantages for quantifying early- and late-stage phase separation.

3.2. Wavelength determination of Fe–Cr phase separation

As shown in figure 2, the Cr-rich phase's composition and spacing continue to develop throughout the phase separation process. Since the spatial distribution of the α' phase is quite homogeneous and symmetric, a sinusoidal concentration variation fits well with the decomposition process. The wavelength (λ) is determined by a measurement of the distances between Cr composition maxima within the compositional sine wave. The autocorrelation function (ACF) is a versatile mathematical tool for quantifying the wavelength of the repeated sinusoidal Cr compositional profile, which is ideal for understanding the phase separation in three dimensions. The 3D radial ACF, R_k , can be expressed as [16]:

$$R_k = \frac{1}{\sigma^2} \sum_{r=0}^{r_{\max}-k} (C_r - C_0)(C_{r+k} - C_0), \quad (1)$$

where C_r is the concentration of a spherical shell at radius r from the chosen center point, r_{\max} is the maximal radius over which the analysis is taken, σ^2 is the variance of the

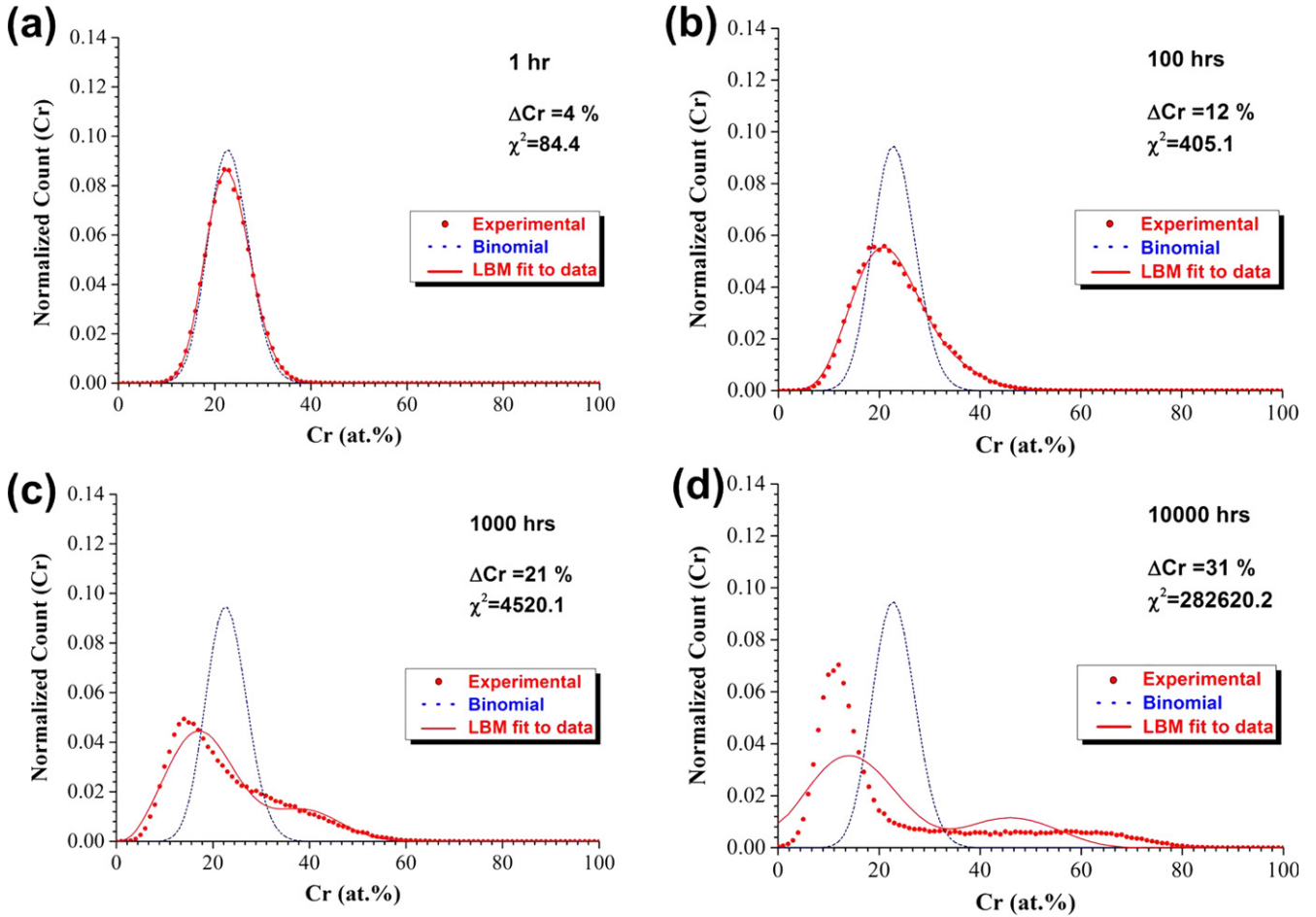


Figure 4. Experimental and binomial composition frequency distributions and the LBM fit for the 2101 alloy after aging at 427 °C for between 1 and 10 000 h.

compositions given by $\sum_0^{\max}(C_r - C_0)^2$, and C_0 is the mean concentration of the element of interest. The first maximum peak represents the wavelength of the α' phase, which is the distance where the Cr composition is most correlated with the central location [17].

Figures 3(a) and (b) display small regions of interest of the APT volumes, which do not represent the average α - α' separation distance for the datasets as a whole. On the other hand, the ACF was applied to the entire APT volume for each sample, and represents the average distance between α' phases with a higher accuracy. In these data, a 0.1 nm shell thickness and a random sampling of 20% of the total ions were used. The data beyond the first maximum peak have little significance and are omitted for clarity. When aged at 427 °C between 100 and 10 000 h for the 2101 alloy, the wavelengths denoted by the first maximum peak are 4.8, 7.6, and 12.0 nm (figure 3(c)). For the same annealing conditions, the wavelengths for the 2205 alloy are 3.6, 6.6, and 7.2 nm, respectively (figure 3(d)). The 2205 alloy exhibits a more sluggish phase decomposition, as confirmed by the comparison of Cr atom maps shown in figures 3(a) and (b). After annealing for 10 000 h, the Cr enriched regions tend to form as isolated Cr-rich clusters for the 2101 alloy, whereas the

interconnected Cr-rich regions are still maintained for the 2205 alloy.

3.3. Amplitude determination of Fe–Cr phase separation

The averaged peak-to-trough Cr composition (ΔCr) was used to evaluate the amplitude of the decomposition. The Langer-Bar-on-Miller (LBM) method is one of the most common approaches to quantify ΔCr [18], and works well for low values of ΔCr ; however, the technique fails as ΔCr reaches values over ~ 20 at%, which will be discussed in more detail later. For the LBM method, the data are split into bins containing a set number of atoms, and the local composition is calculated within each bin (here we use 100 atoms/bin). The result is a frequency distribution histogram of the Cr concentration of each bin, which, for a dataset with two distinct phases, can be fit with a sum of two Gaussian distributions with the same compositional variance σ :

$$P(x) = \frac{1}{(x_2 - x_1)\sqrt{2\pi\sigma^2}} \left\{ (x_2 - x_0) \exp\left[\frac{-(x - x_1)^2}{2\sigma^2}\right] + (x_0 - x_1) \exp\left[\frac{-(x - x_2)^2}{2\sigma^2}\right] \right\}, \quad (2)$$

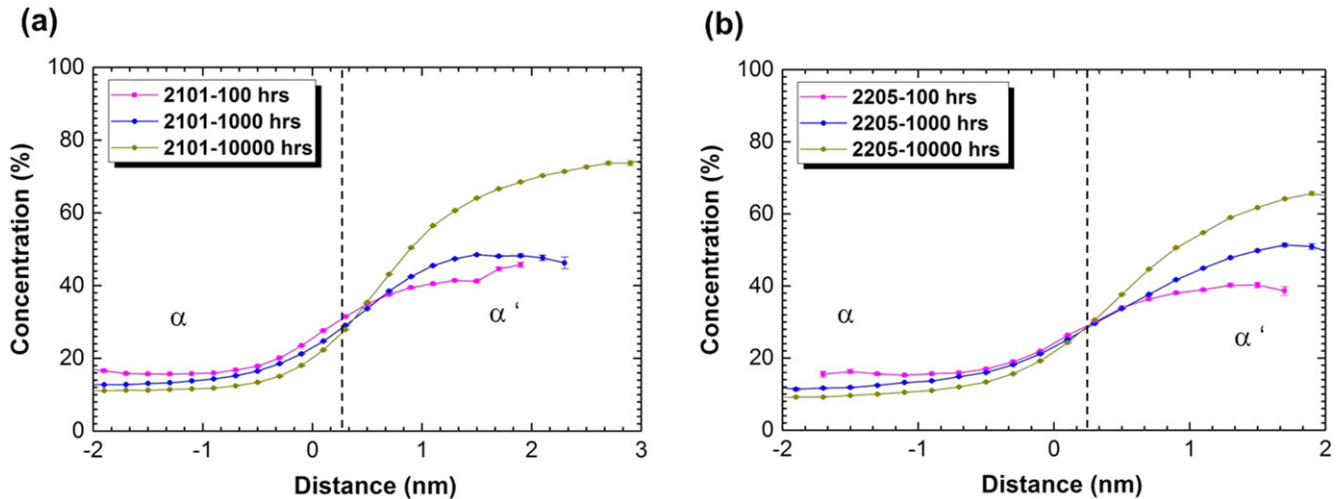


Figure 5. Proximity histogram of Cr atoms at all the Cr-rich and Fe-rich interfaces in 2101 and 2205 duplex stainless steels. A drastic increase in the Cr content of the α' region with aging time is shown.

where x_1 and x_2 are the element of interest compositions for the two Gaussian peaks, and x_0 is the average composition in the analyzed volume of the sample. The experimental histogram is thus represented by three parameters: x_1 , x_2 , and σ , and the composition amplitude (ΔCr) is $x_2 - x_1$.

Figure 4 shows the experimental frequency distribution overlaid with the LBM fit and the binomial distribution (the frequency distribution if Cr were randomly distributed throughout the dataset) of Cr atoms for alloy 2101. As the aging time increases, the phase separation can be visualized in the frequency distribution histograms as two Gaussian functions separating from each other, indicating that the 100 ion bins are detecting Cr-rich and Cr-poor phases. The extended Cr tails from the experimental data in the 1000 h and 10 000 h aged conditions (figures 4(c) and (d)), which are a result of 100 ion bins encompassing phase borders, deviates the frequency distribution histogram from two Gaussian curves. Previous APT study of Fe–Cr binary alloys also showed that the amplitude values obtained from the LBM method are 5–15 at% Cr underestimated compared to the radius distribution function estimate [19]. For these late stage phase separation conditions, the LBM method is not ideal and another method must be applied to the data to quantify ΔCr . To this end, a proximity histogram [20] of Cr isoconcentration surfaces was performed for the same APT dataset to obtain more accurate ΔCr values for late stage phase separation. For this method, 20 at% Cr isoconcentration surfaces were defined within the APT dataset to border the Cr-rich and Cr-depleted interfaces, which can be seen in figure 2. The proximity histogram was then calculated for these surfaces, which calculates a concentration profile with respect to distance from the isoconcentration surface for all solute atoms, and therefore, the compositional profile direction follows the surface normal. The results of the proximity histogram using 20 at% Cr isosurfaces for different aging times are summarized in figure 5. For the 10 000 h annealing condition of the 2101 alloy, the maximum concentration of the α' region calculated using the proximity histogram

reaches 73.8 ± 1.5 at% compared to 46.0 at% calculated from the LBM method, which confirms the underestimation of ΔCr for the LBM fit. It should be noted as well that the proximity histogram does not calculate ΔCr accurately for early-stage phase separation due to the challenge of distinguishing interfaces between Fe-rich and Cr-rich regions. Also, interfaces can be drawn between ‘random’ statistical variations, which would overestimate the ΔCr value. For instance, defined interfaces using isoconcentration surfaces for randomly distributed data would reveal compositional variations in the proximity histograms. These compositional variations are due to interfaces being defined between random fluctuations in local compositions that fit to a binomial distribution, and are not statistically significant. Hence, these two approaches, LBM and proximity histogram analysis, should be jointly considered when assessing the exact amplitude of phase decomposition depending on the stage of the phase separation.

Figure 6 summarizes the evolution of the amplitude, ΔCr , with respect to aging time for both alloys. Initially, the ΔCr difference is small for both alloys; e.g., the Cr distribution is fairly homogeneous. For 100 h aging, ΔCr nearly doubles. The decent fit to the data shown in figures 4(a) and (b) demonstrates that it is more appropriate to use data obtained from LBM. Hence, alloy 2101 shows a slightly higher ΔCr (12 at%) than alloy 2205 (11 at%). For 1000 h aging and beyond, the proximity histogram method is more accurate for describing the actual amplitude, which shows no difference between these two alloys. However, for a 10 000 h annealing, ΔCr is 63.0 at% for the 2101 alloy, which is significantly higher than that (56.0 at%) for the 2205 alloy.

The APT characterization reveals that both the amplitude and wavelength between the α and α' phases increases with aging time, as expected. However, the phase separation process of the 2101 alloy is much more pronounced than that for the 2205 alloy, as manifested by the larger values of λ (figure 3) and ΔCr (figure 6), especially after 10 000 h

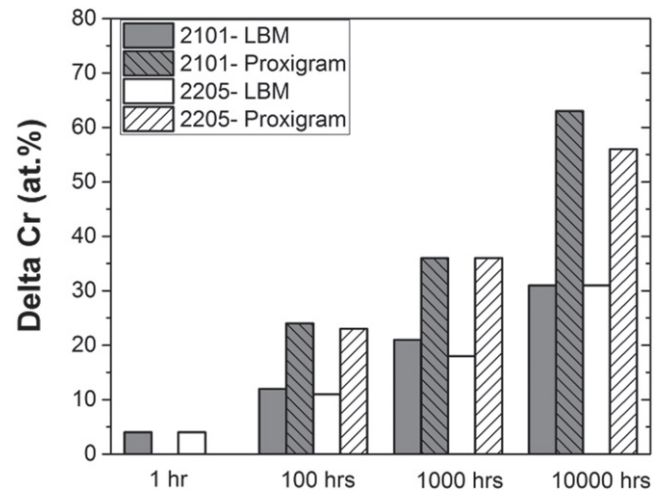
Table 3. Compositions of the α and α' phases after aging between 100 and 10 000 h, as evaluated by proximity histograms. The miscibility gap composition was calculated using the TCFE7 database in Thermo-Calc.

At%	Phase	Cr	Ni	Mn	Si	Cu	Mo
2101-100 h	α	17.34 \pm 0.29	1.51 \pm 0.02	3.99 \pm 0.03	1.77 \pm 0.02	0.27 \pm 0.01	0.19 \pm 0.01
	α'	44.56 \pm 0.72	1.13 \pm 0.03	5.36 \pm 0.06	2.35 \pm 0.04	0.17 \pm 0.01	0.19 \pm 0.01
2101-1000 h	α	12.79 \pm 0.10	1.67 \pm 0.02	3.67 \pm 0.03	1.54 \pm 0.02	0.38 \pm 0.01	0.18 \pm 0.01
	α'	48.25 \pm 0.44	1.01 \pm 0.14	6.52 \pm 0.12	2.92 \pm 0.28	0.12 \pm 0.02	0.21 \pm 0.04
2101-10 000 h	α	11.19 \pm 0.06	1.62 \pm 0.02	3.67 \pm 0.04	1.59 \pm 0.02	0.36 \pm 0.01	0.20 \pm 0.01
	α'	70.25 \pm 0.19	0.31 \pm 0.02	6.36 \pm 0.10	2.64 \pm 0.07	0.03 \pm 0.01	0.26 \pm 0.02
2101 miscibility gap composition by TCFE 7	α	9.58	1.64	1.78	2.47	NA	NA
	α'	70.27	0.07	14.78	0.03	NA	NA
2205-100 h	α	15.50 \pm 0.85	5.78 \pm 0.13	1.73 \pm 0.03	1.08 \pm 0.06	0.75 \pm 0.05	1.83 \pm 0.07
	α'	40.63 \pm 1.72	3.79 \pm 0.11	2.05 \pm 0.09	1.25 \pm 0.07	0.25 \pm 0.03	2.51 \pm 0.09
2205-1000 h	α	13.20 \pm 0.12	8.83 \pm 0.21	3.22 \pm 0.13	1.65 \pm 0.09	1.17 \pm 0.08	1.65 \pm 0.09
	α'	51.39 \pm 0.49	2.37 \pm 0.10	1.75 \pm 0.08	1.11 \pm 0.07	0.03 \pm 0.01	3.35 \pm 0.11
2205-10 000 h	α	9.01 \pm 0.18	7.64 \pm 0.16	2.62 \pm 0.10	1.72 \pm 0.08	0.82 \pm 0.05	1.21 \pm 0.07
	α'	64.15 \pm 0.33	1.51 \pm 0.16	1.43 \pm 0.13	0.78 \pm 0.10	0.05 \pm 0.03	2.92 \pm 0.20
2205 miscibility gap composition by TCFE 7	α	7.11	5.28	0.56	1.35	NA	NA
	α'	80.85	0.06	6.14	0.02	NA	NA

annealing. Previous studies indicate that significant embrittlement is caused by a larger lattice mismatch due to the increase of Fe–Cr phase decomposition [21], shear modulus variation [21] or alteration of the moving dislocation morphologies [22]. However, as the aging time goes beyond 1 000 h, the hardness of the less decomposed ferrite (2205 alloy) is much higher than that in the 2101 alloy (table 1), which is in contrast with previous predictions. Table 3 presents the composition of the α and α' phases after aging between 100 and 10 000 h, as evaluated by proximity histograms. Different partitioning of elements, such as Ni and Cu, will lead to a different kinetics and morphology of secondary phase precipitation, which potentially provides an additional hardening effect in ferrite. The extent of clustering inside these two alloys has been statistically quantified.

3.4. G-phase precipitation behavior inside ferrite

The influence of G-phase on thermal embrittlement is still unclear [5, 23, 24]. In Fe–Cr alloys, the appearance of G phase is accompanied by Fe–Cr phase separation, which provides enough flux of constitutional atoms such as Ni, Mn, and Si. Previous APT works by Pareige *et al* [24] suggest that the coarsening of G phase has no influence on the microhardness of duplex stainless steel after 30 000 h ageing at 350 °C. However, the α – α' phase evolves concurrently with the G phase [25], and thus, the influence of the α – α' phase separation on the microhardness cannot be completely deconvoluted from the G-phase evolution. G-phase precipitates were also found in other steels such as irradiated reactor pressure vessel steels [26]. It was documented that sufficient G-phase-precipitate coarsening can reduce the hardness after thermal aging at 454 °C for 168 h [27], indicating that a fine distribution of G-phase contributes an additional hardening effect. Li *et al* [11] has published an

**Figure 6.** A summary of amplitude determination by LBM and proxigram analysis methods for 2101 alloys and 2205 alloys aged at 427 °C.

opposite view, in which they annealed the Fe–Cr steel sample above the miscibility gap, which removes the α/α' phase separation, but retains the G-phase precipitation within the matrix. The hardness of the annealed sample containing the G-phase was also reduced significantly. Contradictory observations like these have motivated a comparison of the precipitation behavior of the nanoscale G-phase with the material properties and α – α' phase separation.

For small precipitates, APT uniquely provides a detailed characterization, including size, composition, number density, and volume density inside the 3D volume. Figure 7 displays the various stages during the nucleation and growth of Cu–Ni rich clusters as a function of aging time for the 2101 alloy. The 1 h aged sample exhibits a homogeneous distribution of

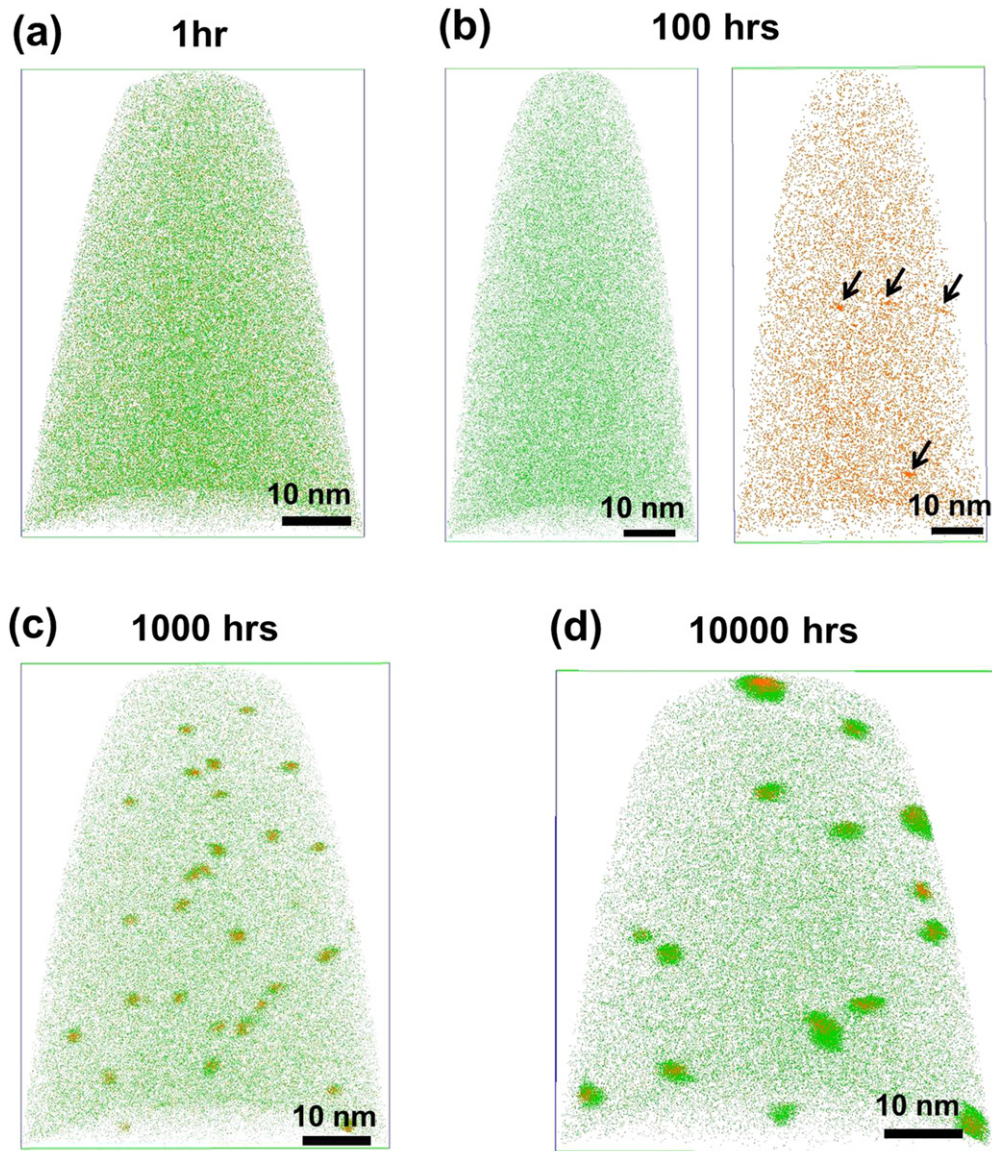


Figure 7. APT elemental maps of Cu (orange) and Ni (green) atoms of the 2101 alloy aged at 427 °C for different times: (a) 1 h; (b) 100 h; (c) 1000 h, and (d) 10 000 h. The arrows in (b) mark the positions of initial Cu clustering.

Ni and Cu atoms. As the aging time proceeds to 100 h, very small and largely spaced Cu rich clusters appear inside the tip; however, the Ni distribution is still rather homogeneous (figure 7(b)). It has been suggested that local strain can induce Cu clustering in Fe–Cr alloys [28]. Hence, the Cu clusters are mostly found close to the α – α' phase domains possessing high misfit strain. From figure 7(c), it is observed that Ni also starts to cluster and the number density of Ni–Cu rich clusters is at its maximum at 1 000 h, where the Ni atoms aggregate around the Cu clusters, forming a core–shell precipitate structure. After 10 000 h, the number of Cu–Ni-rich clusters decreases while the particle sizes start to increase, indicating that there is cohesion of the particles. Overall, the evolution of Cu–Ni clustering occurs in three stages: stage I corresponds to the onset of Cu clustering; stage II relates to the Cu and Ni co-clustering with a Cu-rich core and Ni-rich shell structure; stage III manifests as the disappearance of small precipitates

and the further growth of large precipitate via the Ostwald ripening mechanism [29].

Figure 8 shows 1 nm thick atom maps of the center of selective clustering for the first two stages to further clarify the precipitation growth behavior. The major elements, i.e., Fe and Cr, are not shown for clarity. In stage I (figure 8(a)), Cu atoms start to cluster in advance of all other non-major elements, i.e. Ni, Mn, and Si. A 2 at% Cu isosurface borders the cluster from the matrix, which is used to create a proximity histogram, shown in figure 8(c). The Cu concentration reaches 30 at% within the precipitate, while the concentration of other elements remains consistent with that in the matrix, further elucidating stage I precipitate formation. For stage II, it is evident from figure 8(b) that Cu, Ni, Mn, and Si atoms are also enriched around the precipitates such that the precipitates exhibit a core–shell elemental distribution. In figure 8(d), the proxigram analyses in terms of 10 at% Cu isosurfaces reveal that the solute

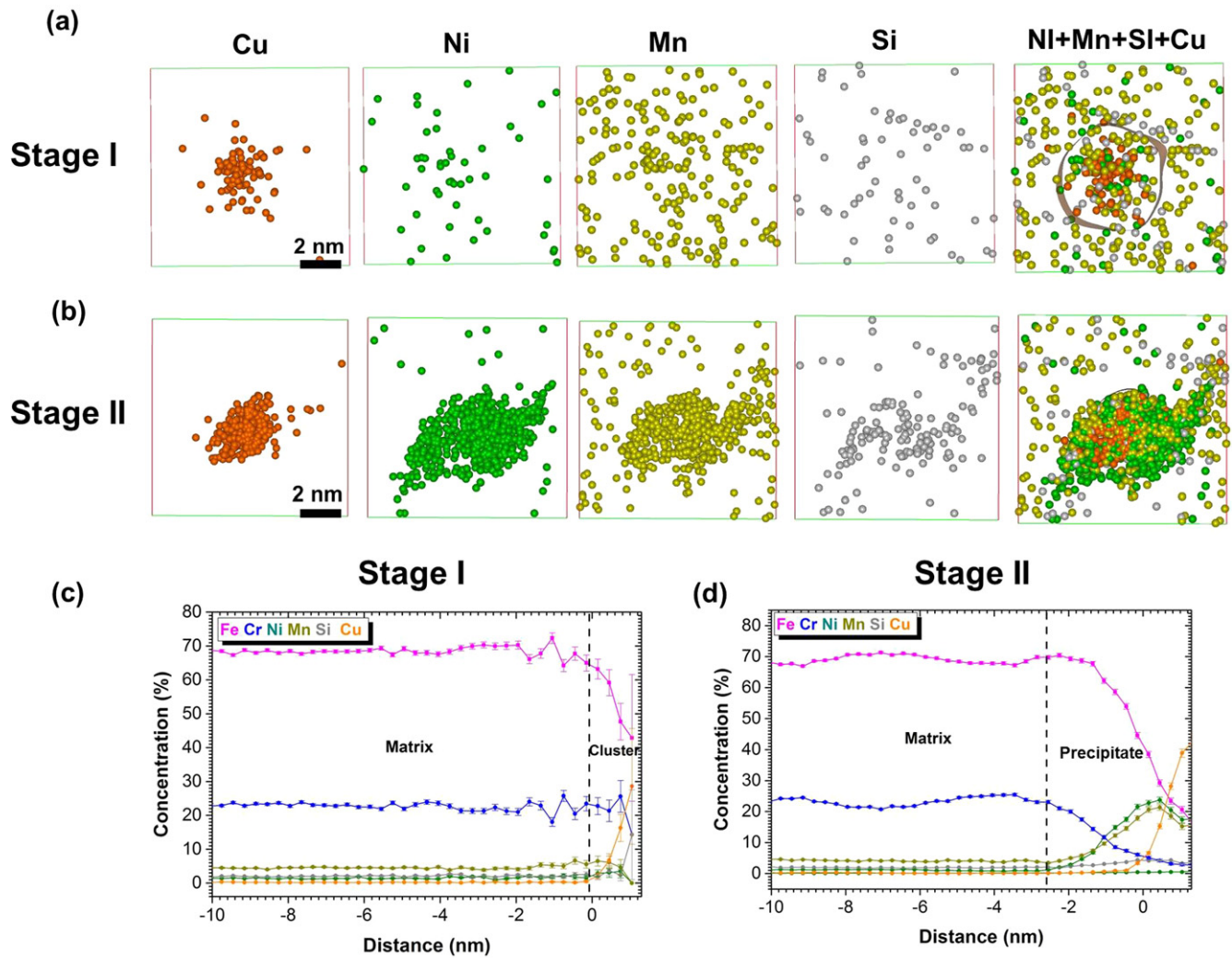


Figure 8. 1 nm thick maps showing two stages of forming the Ni–Mn–Si–Cu precipitates (a), (b) in the 2101 alloy and proximity histograms of 2 at% Cu isosurfaces at two different stages (c), (d). The dashed line in (c) and (d) mark the potential interfaces between the cluster (precipitate) and the matrix.

levels reach 24.5 ± 1.8 at% Ni, 20.4 ± 1.4 at% Mn, and 4.0 ± 0.5 at% Si at the precipitate–matrix interface.

The 2205 alloy exhibits a different precipitation behavior of the Ni–Mn–Si–Cu particles. At the onset of precipitation, initial Cu clustering is not observed before the other elements cluster. Instead, Cu and Ni atoms tend to cluster simultaneously and distribute homogeneously inside the precipitates. Figure 9 shows an example of such a feature. After 10 000 h annealing, no core–shell structure was detected throughout the APT tip. As shown in the proxigram analysis in figure 9(b), the selected precipitates display 37.2 ± 4.2 at% Ni, 19.1 ± 3.3 at% Mn, 6.4 ± 1.6 at% Si, and 5.2 ± 1.2 at% Cu inside the clusters. The absence of the core–shell elemental distribution is similar to the G phases previously observed in SAF 2507 steels [25].

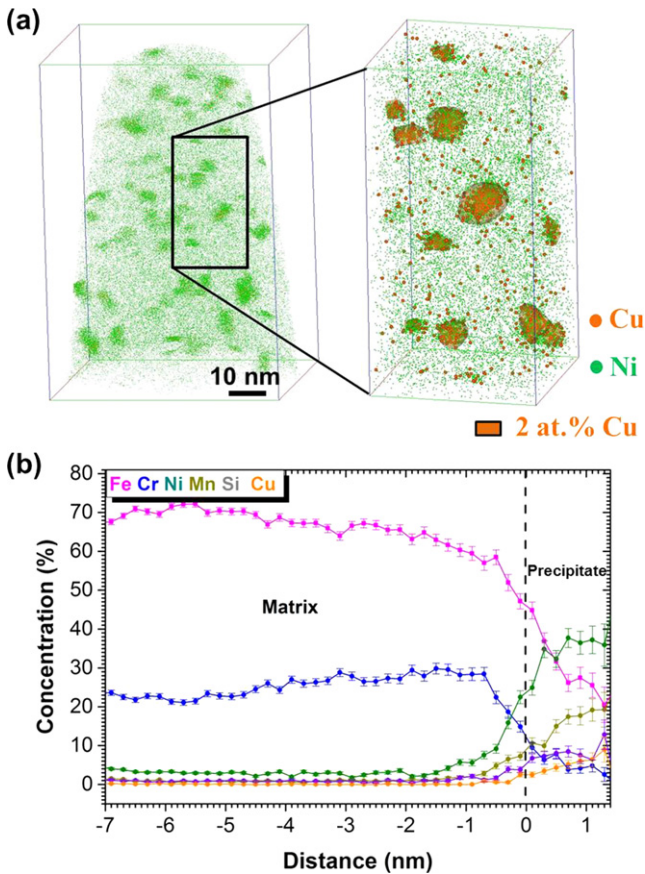
The discrepancies of precipitates in the 2205 alloy from that in the 2101 alloy are mainly due to the higher Ni content or higher Ni/Cu ratio within the ferrite. Table 3 provides the time dependent elemental partitioning inside ferrite. Also, the miscibility gap compositions of both alloys are evaluated by

Thermo-Calc using only Fe, Cr, Ni, Mn, and Si elements. The calculations predict a significant partitioning of Mn in α' phase and Ni in α phase. For the 2101 alloy, there are abrupt Ni concentration changes from 1.1 to 0.3 at% in α' phases, indicating that the diffusion of Ni atoms from the α' phases determines the kinetics of G-phase precipitation. For the 2205 alloy, much higher Ni concentrations are shown in both α' and α phases. The average Ni content inside the ferrite, as shown in table 2, is 1.13 ± 0.01 at% for the 2101 alloy and 4.22 ± 0.02 at% for the 2205 alloy. In addition, the atom maps of Ni and Cu for the 2205 alloy display a larger particle density of precipitates than the 2101 alloy for the same condition. Much more precipitates form during the decomposition process in 2205 than the 2101 alloy, resulting in a higher precipitation hardening effect.

APT unambiguously provides the quantitative statistics for such precipitates, and thus allows for the quantification of the influence of particles on hardening. The parameters used for bordering the precipitates were identical for all the datasets (delocalization, 1.5 nm; voxel size, 0.7 nm; Cu cluster

Table 4. Statistics and the estimated hardness increase for Ni–Mn–Si–Cu clusters in ferrite of the 2101 and 2205 alloys aged at 427 °C.

Sample ID	Volume fraction (%)	Number density (#/m ³)	Average precipitate radius (nm)	Estimated hardness increase (HV)
2101 100 h	0.01	1.06×10^{22}	1.33	3.5
2205 100 h	0.72	1.25×10^{24}	1.30	28.0
2101 1000 h	0.34	1.01×10^{23}	1.90	32.2
2205 1000 h	0.94	3.51×10^{23}	1.68	49.0
2101 10000 h	0.56	5.74×10^{22}	2.60	45.0
2205 10000 h	1.74	5.32×10^{23}	2.04	75.5

**Figure 9.** The Cu and Ni atom maps of the 2205 alloy aged at 427 °C for 10 000 h (a) and a proxigram analysis of selected isosurfaces of 2 at% Cu (b).

isoconcentration surface value 2 at%; Ni isoconcentration value, 7 at%). The average radius was calculated by the sphere-equivalent volume approach. Table 4 summarizes the average radius, number density, and volume fraction of the G-phase as a function of aging time. Upon isothermal aging at 427 °C, the average radius of the Ni–Mn–Si–Cu clusters in both alloys increases continuously from ~1.3 to >2 nm. Note that the number density of the G-phase precipitates in the 2205 alloy is one order of magnitude higher than that in the 2101 alloy. According to the classic Ashby–Orowan equation [1], the stress increase contributing from the precipitation hardening is described as:

$$\sigma_p = \frac{0.538 \cdot Gb \cdot f_v^{\frac{1}{2}}}{d} \ln\left(\frac{d}{2b}\right), \quad (3)$$

where σ_p is the precipitation induced stress increase, G is the shear stress of the steel (8.16×10^4 MPa), b is the burgers vector, f_v is the volume fraction of particles, and d is the diameter of the particles [1]. The stress increase can be easily converted into a hardness increase using the empirical relation: $H_v \approx 0.3\sigma_p$ [30]. The calculated hardness value increases are also shown in table 4. With a low G-phase volume fraction (~0.01%), no noticeable influence on the hardening is predicted, as the estimated hardness increase is only 3.5 HV. Overall, the stress increase due to precipitate formation for the 2205 alloy is much higher than that for the 2101 alloy because of the higher volume fraction of particles for each aging time. Therefore, additional hardening from the G-phase should be considered, although the G-phase formation is regarded as a byproduct of Fe–Cr phase decomposition.

4. Conclusions

Two series of Fe–Cr alloys, 2101 and 2205, were selected for studying the phase separation of ferrite after time-dependent, intermediate-temperature annealing by APT. By characterizing the wavelength and amplitude of Fe-rich (α)/Cr-rich (α') phase decomposition, and G-phase precipitation, the main findings are as follows:

1. The expanded embrittlement effect from G-phase precipitation should be considered when evaluating the thermal embrittlement of DSS. The Ashby–Orowan stress analysis indicates that the increased hardness in 2205 compared to 2101 alloys arises from the larger number density of G-phase precipitates.
2. Although the hardness increase for the 2205 alloy is higher than the 2101 alloy after long time aging, a less enhanced phase separation behavior occurs for the 2205 alloy, as reflected by the smaller λ and ΔCr values.
3. The LBM method can quantify the ΔCr value more accurately than the proxigram approach for evaluating the amplitude of phase separation in the initial stage because the proxigram approach suffers from bordering random statistical variations in addition to the α/α' interfaces, which artificially increases the ΔCr value. However, LBM underestimates the ΔCr value from a highly decomposed sample. Instead, the proxigram approach would be more accurate to assess the exact ΔCr when the samples are aged beyond 1000 h.

4. The local chemistry and evolution procedure depends largely on the alloying elements within the ferrite. In the 2101 alloy with a Ni/Cu atomic ratio as low as six, the Cu-rich clusters serve as the nucleation sites for the G-phase precipitation, in which a Cu-rich core, Ni–Mn–Si-rich shell structure forms. For the 2205 alloy, with a Ni/Cu ratio larger than 17.5 within the ferrite, the G-phase starts to form with the Cu and Ni atoms enriching simultaneously without the appearance of a core–shell distribution.

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