

Toward biologically meaningful net carbon exchange estimates for tall, dense canopies: multi-level eddy covariance observations and canopy coupling regimes in a mature Douglas-fir forest in Oregon

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Abstract

We sought to improve net ecosystem exchange (NEE) estimates for a tall, dense, mature Douglas-Fir forest in the Oregon Coast range characterized by weak flows, systematic wind directional shear, and limited turbulent mixing throughout the diurnal period. We used eddy covariance (EC) observations at two levels and concurrent biological measurements of carbon and water fluxes collected over a period of 6 years (2006-2011) to develop and test a conceptual framework to i) reduce uncertainty by retaining more measurements for the computation of annual NEE sums, and ii) produce defensible and biologically meaningful estimates by accounting for the missing sub-canopy respiration. The framework assumes that a) the scalar exchange between vertical layers can be categorized into discrete canopy coupling regimes, and b) advection leads to a systematic loss of scalar from the observational volume

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that can indirectly be estimated and accounted for as sub-canopy respiration flux when canopy layers are decoupled.

Periods with a decoupled sub-canopy layer dominated and occupied 65 and 88 % of the day- and nighttime periods, respectively. Annual NEE derived from the new framework was estimated as $480 \text{ gC m}^{-2} \text{ yr}^{-1}$, which was reduced by $620 \text{ gC m}^{-2} \text{ yr}^{-1}$ compared to traditional estimates from single-level EC data filtered using a critical friction velocity. The reduced NEE was due to an enhanced ecosystem respiration (RE), while gross ecosystem productivity remained unchanged. Improved RE estimates agreed well with those from independent estimates based on soil, stem, and foliage respiration within 3 %. Risks and limitations of the new framework are discussed. We conclude that concurrent above- and sub-canopy EC observations are essential to measure a meaningful carbon balance in tall, dense forests since they do not lend themselves to traditional, standardized processing. The new framework may help to include more tall and dense forests in global carbon cycle synthesis and modeling efforts.

Keywords: Net ecosystem exchange, Ecosystem respiration, Advection, Eddy covariance, u-star correction, Canopy flow, Douglas-Fir, Turbulence

1. Introduction

2 The net ecosystem exchange NEE is the single most important parameter
3 describing the strength of the carbon sink or source of terrestrial ecosystems.
4 Its estimation has received much attention in the literature and a commonly
5 identified obstacle is the 'nighttime problem' when weak turbulent mixing
6 becomes limiting and the computed NEE from a simplified mass balance ap-

7 proach may not reflect ecosystem functioning (see Goulden et al., 1996; Aubi-
8 net et al., 2000; Baldocchi, 2003; Papale et al., 2006, and references therein).
9 The simplified mass balance approach defines NEE as the sum of the turbu-
10 lent vertical carbon dioxide or methane flux observed above the canopy and
11 the temporal change in storage term from profile observations. In contrast,
12 during the day when mixing is enhanced through stronger flows and signifi-
13 cant heat flux, estimates of NEE are typically assumed to reflect ecosystem
14 response to environmental drivers such as light, nutrients, and water inde-
15 pendent of the strength of the turbulent transport. Global NEE estimates
16 are modeled based on continental observational networks representing the
17 major biomes, but the selection of individual sites within the networks may
18 be biased toward short vegetation such as grass, open shrubland, and forest.
19 In these canopies, mixing of the scalar sinks and sources can sufficiently well
20 be estimated using variety of mixing indicators. Quantities that have been
21 proposed as a proxy for the turbulent mixing strength include the standard
22 deviation of the vertical velocity variance σ_w (e.g. Acevedo et al., 2009), the
23 friction velocity u_* (e.g. Goulden et al., 1996), their non-dimensional ratio
24 $\sigma_w u_*^{-1}$ termed the integral turbulence characteristic (e.g. Foken and Wichura,
25 1996; Thomas and Foken, 2002), the mean wind speed U (e.g. Suyker et al.,
26 2003), and a modified turbulent kinetic energy scale u_{TKE} (Wharton et al.,
27 2009). The indicators are typically evaluated from eddy covariance (EC)
28 measurements with a fixed perturbation time scale of 30 or 60 min taken
29 at a single level above the main canopy. However, dense canopies pose ad-
30 ditional challenges as they suffer from a night- and daytime problem, since
31 the dense crown space with the maximum plant area index (PAI) presents

32 a mechanical barrier and efficient momentum sink throughout the diurnal
33 cycle leading to a frequent, persistent decoupling of the sub-canopy from the
34 overstory and above-canopy layers (e.g. Thomas and Foken, 2007; Belcher
35 et al., 2008). In contrast, short vegetation and open forests are only tempo-
36 rally limited by turbulent mixing when surface heating and the mean flow are
37 weak, which typically occurs at night in the absence of significant synoptic
38 meso-scale forcing. As a result, forests with high-PAI, closed canopies are
39 often excluded from network syntheses for reasons of uncertain data quality
40 and insufficient number of observations, since many measurements need to
41 be discarded for the computation of seasonal and annual NEE because of the
42 systematic turbulence limitations. However, these ecosystems may be very
43 efficient carbon sinks as demonstrated by their high PAI, which can only
44 be sustained in high productivity ecosystems. Hence, we identify a signif-
45 icant observational, modeling, and interpretational problem when assessing
46 regional to global carbon balances and their dynamics without the inclusion
47 of tall and dense forests.

48 In this study we seek to improve NEE estimates for a very dense mature
49 Douglas-fir stand analyzing six years of concurrent EC flux observations at
50 two levels, above the main canopy crown and in the clear bole space below
51 the main canopy crown, with the following objectives:

- 52 • identify a simple and meaningful estimator for canopy mixing, coupling,
53 and decoupling that reflects characteristics of the mean flow and the
54 turbulent carbon, sensible and latent heat, and momentum fluxes,
- 55 • construct an alternative and practical theoretical framework for the
56 evaluation of multi-level EC observations to estimate NEE,

- 57 • increase the fraction of sub-daily NEE estimates that are retained for
58 the computation of seasonal and annual sums, which are assumed to
59 reflect ecosystem behavior,
- 60 • compute an improved carbon balance by applying the proposed frame-
61 work to the observations.

62 We do not expect this study to solve the problem of overestimating NEE
63 by systematically missing ecosystem respiration, but it may be an important
64 step toward producing defensible and biologically meaningful estimates for
65 dense canopy sites. The ultimate goal is to include these ecologically im-
66 portant sites into network syntheses and global estimates. We further aim
67 at demonstrating the utility of concurrent sub-canopy EC observations to
68 better understand turbulent mixing and other micrometeorological processes
69 in dense canopies. In the literature, only few studies exist that utilize sub-
70 canopy EC observations with the intention of either incorporating their flux
71 estimates into the carbon mass balance or partitioning net carbon fluxes into
72 its components (e.g. Misson et al., 2007; Falk et al., 2008). A number of
73 recent studies has focused on evaluating the advective terms directly using
74 sensor networks and include their flux contributions to the mass balance
75 (e.g. Feigenwinter et al., 2004; Staebler and Fitzjarrald, 2004; Aubinet et al.,
76 2005), while the success of these efforts has recently been questioned (Aubi-
77 net et al., 2010).

78

79 **2. Conceptual framework**

80 Figure 1 placed here.

81

82 Diagnosing the dynamics of the canopy flow and its turbulence is critical
83 to connecting the biologically active surfaces such as the foliage, stems, and
84 the soil with the micrometeorological sensors used to quantify the ecosystem
85 scalar exchange. The conceptual framework presented here is based on two
86 major assumptions: first, coupling between vertical layers of the soil-plant-
87 air continuum and thus the exchange of scalar flux varies depending on the
88 strength of the turbulent mixing, which can be diagnosed using multi-level
89 turbulence statistics. Secondly, both vertical and horizontal advection is the
90 main transport mechanism removing the scalar from the observational vol-
91 ume when layers are decoupled. The concept of vertical communication of
92 air across the canopy profile is based on the exchange regimes proposed by
93 Thomas and Foken (2007) for a tall, moderately dense spruce canopy with
94 a PAI of 5.2. Their work differentiates between different conceptual vertical
95 layers which together comprise most of the roughness sub-layer (Fig. 1): i)
96 the above-canopy layer between the top of the overstory and the upper EC
97 observation height, ii) the overstory where most of the PAI is concentrated,
98 iii) the sub-canopy layer or the clear bole space with minimum PAI, and iv)
99 the ground layer including the soil surface and understory including herbs
100 and shrubs often comprising a secondary maximum in the PAI profile. For
101 a typical EC setup located above the canopy at $zh_c^{-1} \approx 1.2$ to 1.8, where
102 z is the observation and h_c the mean canopy height, to measure the entire
103 ecosystem, i.e., integrate over all its vertically distributed scalar sources and

104 sinks, one must assure that turbulent eddies communicate across all three in-
105 terfaces between these four layers. In their approach horizontal heterogeneity
106 and transport were originally neglected, but were investigated later by Ser-
107 afimovich et al. (2011). Thomas and Foken (2007) proposed five exchange
108 regimes with increasing degree of vertical communication between the layers
109 using turbulence measurements at five observation levels to determine the
110 penetration depth and flux contribution of mixing-layer type eddies. They
111 demonstrated that the sub-canopy layer is often decoupled from the crown
112 and free roughness sub-layers even during the day when the above-canopy
113 flow and turbulent mixing are significant. A recent study by Foken et al.
114 (2012) confirmed the utility of this concept when investigating the dynamics
115 of ecosystem fluxes of volatile organic compounds, nitrous oxides, and ozone.
116 The authors concluded that the observed concentration profiles and fluxes
117 could only be explained when the cross-interface transport was diagnosed us-
118 ing the exchange regimes, while other simpler methods including the critical
119 u_* threshold approach (Goulden et al., 1996) failed.

120 Since extensive multi-level EC observations and complex post-processing are
121 impractical and probably cost-prohibitive for long-term ecosystem studies
122 targeting measuring NEE on annual and decadal time scales, we simplified
123 the method. The method proposed here diagnoses the cross-interface trans-
124 port based on only two observational heights located in the above- and sub-
125 canopy layers while retaining the concept of exchange regimes and layers.
126 We further added a concept originally proposed by Scanlon and Albertson
127 (2001) to use scalar-scalar cross-correlations between perturbations in carbon
128 dioxide and water vapor mixing ratios to partition scalar sinks and sources.

129 Instead of partitioning, we here invert the original idea and apply it to di-
130 agnose the communication of air across the canopy profile by relating the
131 scalar-scalar fingerprint to the turbulent mixing strength. The simplified
132 conceptual framework differentiates between three exchange regimes:

133

- 134 • Fully coupled canopy (C): Above-canopy EC fluxes fully integrate over
135 all scalar sinks and sources and are representative of the entire ecosys-
136 tem. This is the ideal state typically assumed for traditional EC anal-
137 ysis of NEE.
- 138 • Decoupled sub-canopy (Ds): Fluxes observed with the sub-canopy EC
139 system integrate the scalar sinks and sources of the ground and sub-
140 canopy layers, but are decoupled from those of the overstory and above-
141 canopy layers.
- 142 • Decoupled ground layer (Dg): Fluxes observed with the sub-canopy EC
143 system are not coupled to the soil and ground vegetation, i.e., do not
144 represent soil CO₂ efflux, photosynthesis and autotrophic respiration
145 of the ground layer. This regime was originally not included in the
146 concept by Thomas and Foken (2007) since transport across the soil-
147 air interface was neglected.

148 Table 1 placed here.

149

150 We emphasize that the biological processes of respiration and assimila-
151 tion are continuous in time. However, when ground and sub-canopy layers
152 are decoupled (Dg, Ds regimes), their signals and flux contributions are not

153 properly captured by the micrometeorological instrumentation and simplified
 154 mass balance approach due to limited turbulent transport and mixing. As in
 155 most NEE studies, the simplified mass balance approach deployed here does
 156 not directly measure the advective transport and neglects the pressure trans-
 157 port and the horizontal flux divergence terms. For a fully coupled canopy (C),
 158 advection is assumed to be negligible and the ecosystem-scale net ecosystem
 159 exchange is the defined as the sum of vertical turbulent flux and temporal
 160 change in storage term, i.e.,

$$NEE = FCO_{2,top} + \Delta S. \quad (1)$$

161 The terms are defined as

$$FCO_2 = \overline{\rho w' \chi_c'}, \quad \Delta S = \int \left(\frac{d\chi_c}{dt} \right) dz, \quad (2)$$

162 with the subscript 'top' denoting the above-canopy EC level, w is the vertical
 163 wind component, χ_c the mixing ratio of carbon dioxide, ρ air density, and
 164 the overbar represents the temporal average and the prime the perturbation
 165 thereof. However, when decoupling limits the vertical exchange across the
 166 interface between the sub-canopy and canopy layers (Ds regime), we assume
 167 that advection becomes important and can indirectly be estimated and ac-
 168 counted for by adding the sub-canopy respiration flux Re_{sub} to the simplified
 169 carbon mass balance, which leads to

$$NEE = FCO_{2,top} + Adv + \Delta S \approx FCO_{2,top} + Re_{sub} + \Delta S. \quad (3)$$

170 Adv is the sum of vertical and horizontal advection, and the subscript 'sub'
171 denotes the sub-canopy EC level. The validity of this assumption will be dis-
172 cussed in Section 4.6. Re_{sub} was quantified following the method of Thomas
173 et al. (2008), which applies a relaxed eddy accumulation approach to the
174 sub-canopy EC observations to sample episodic respiration pulses used to
175 construct flux estimates. A brief overview of this method is given in Section
176 3.2. Note that $Re_{sub} \neq FCO_{2,sub}$, the latter being the net carbon exchange
177 observed by the sub-canopy EC system. Although this approach was origi-
178 nally proposed for daytime conditions only, the method is equally applicable
179 to nighttime conditions due to the persistence of the scalar-scalar correla-
180 tions between carbon dioxide and water vapor in the sub-canopy.

181 In traditional single-level EC NEE studies only data for a fully coupled
182 canopy (C), typically diagnosed by exceeding a critical threshold of the fric-
183 tion velocity, can be used to compute a biologically meaningful NEE. In con-
184 trast, the improved method proposed here allows for computation of NEE
185 also when the turbulent exchange of air between the sub-canopy and the
186 crown is significantly inhibited (C + Ds). Both methods fail to produce a
187 meaningful NEE for conditions when the sub-canopy is decoupled from the
188 ground layer (Dg). A summary of the applied mass balances to compute
189 the improved NEE estimates depending on the light and canopy exchange
190 regimes is given in Table 1. Note that for our site the photosynthetic uptake
191 of the understory GEP_{sub} was neglected for NEE computations because of the
192 very low photosynthetic photon flux densities measured in the sub-canopy
193 that reach maximum values of $< 50 \mu\text{mol m}^{-2} \text{s}^{-1}$ during summer months
194 (Thomas, 2011).

195 3. Methods & Materials

196 3.1. Measuring the turbulent fluxes and change in storage term

197 The turbulent fluxes of momentum τ and its related generalized friction
198 velocity $u_* = \sqrt{\frac{\tau}{\rho}} = \sqrt[4]{\overline{u'w'^2} + \overline{v'w'^2}}$, carbon dioxide FCO_2 , and water vapor
199 $\lambda E = \rho \overline{w'q'}$ were computed as covariances from high-frequency measurements
200 sampled at 20 Hz. The symbols u , v , w stand for the along, cross, and
201 vertical wind velocities respectively, and q is the water vapor mixing ratio.
202 Wind vector components were transformed prior to computing fluxes and
203 flow statistics using a double-rotation matrix consisting of median rotation
204 angles computed for wind direction sectors that contain an equal number of
205 data over one year. The philosophy behind this rotation is similar to that
206 of the planar-fit rotation method (Wilczak et al., 2001), but the streamlines
207 are not forced through a plane and the statistical uncertainty of the rotation
208 angles across all wind direction sectors is identical. At our site, the mean
209 streamlines do not fall on to a plane because of the mountainous terrain.
210 The perturbation time scale to compute excursions from mean quantities for
211 the above-canopy system was chosen as 30 min, while 6 min were used for
212 the sub-canopy observations based on co-spectral analysis to minimize the
213 influence of motions other than turbulence introducing non-stationarities. It
214 also enables one to relate the turbulent fluxes and flow statistics to properties
215 of the generally weak sub-canopy mean flow (see Section 3.4). Subsequently,
216 five 6-min sub-canopy fluxes were averaged into one 30-min estimate to min-
217 imize the random error and to match the temporal resolution between both
218 EC systems. The upper EC system was located above the canopy at 37.4
219 m above ground level (agl), i.e., at $zh_c^{-1} = 1.4$, and consisted of a sonic

220 anemometer (Model CSAT3, Campbell Sci., Logan, UT, USA) in combina-
221 tion with open- (Model Li-7500, Licor Env., Lincoln, NE, USA) and closed-
222 path (Model Li-7000, Licor Env., Lincoln, NE, USA) infra-red gas analyzers
223 recorded by a data logger (Model CR5000, Campbell Sci., Logan, UT, USA).
224 From Jan 2007 until Dec 2008, both infra-red gas analyzers were run in paral-
225 lel continuously for cross-validation of carbon dioxide and water vapor fluxes
226 showing only small deviations from a 1:1 relationship. Starting Jan 2009, the
227 closed-path gas analyzer was chosen to be the primary instrument for flux
228 measurements because of the frequent and abundant winter precipitation,
229 while the open-path analyzer was brought online for episodic comparison
230 during the drier summer each year. The closed-path analyzer was housed in
231 a custom-designed, insulated, and temperature-controlled aluminum case at
232 the top of the tower. The sample gas was drawn via an electric brushless
233 pump (Model 6025SE, Thomas, Sheboygan, WI, USA) through a 7 m-long
234 insulated, heated, and pressure regulated (Model 640, MKS Instr., Andover,
235 MA, USA) stainless steel line of 6.35 mm (1/4 in) outer diameter equipped
236 with two 1 μm Gelman filters yielding a mean flow of 5 to 8 l min^{-1} depend-
237 ing on pump supply voltage. The flow rate was monitored, recorded, and in
238 combination with the analyzer status flag used to quality-filter the data. The
239 lower, sub-canopy system was located at 4 m agl, i.e., at $zh_c^{-1} = 0.14$, and
240 consisted of a sonic anemometer in combination with an open-path analyzer
241 recorded by a data logger identical to those used for the upper EC system.
242 Fluxes computed from open-path analyzer measurements were subject to a
243 post-hoc density correction (Webb et al., 1980). Aspirated air temperature
244 T , relative humidity rH , net radiation R_{net} , and barometric pressure p , were

245 continuously measured at both observations levels, while short- and longwave
246 radiation components and the photosynthetic photon flux density (Models
247 CNR-1 and PARlite, Kipp&Zonen, Delft, The Netherlands) were observed
248 only above the canopy at the tower top.

249 The change in storage term ΔS (see Eq. 2) was estimated by measuring the
250 temporal dynamics of the carbon dioxide mixing ratio across a vertical profile
251 sampled at 7 heights at 0.5, 1.7, 4.0, 9.0, 18.0, 27.0, 37.9 m agl several times
252 within a 30-min period using an infra-red gas analyzer (Model Li-820, Licor
253 Env., Lincoln, NE, USA) operating at 1 Hz. This custom-made CO₂ profile
254 system followed the design by Stephens et al. (2006) with major modifica-
255 tions to minimize fluctuations in pressure, air temperature, and water vapor
256 (Gockede et al., 2010). Calibration and quality control of the raw data was
257 based on four NOAA-ESRL gas standards sampled at regular intervals.

258 *3.2. Computing the sub-canopy respiration flux*

259 When the above- and sub-canopy layers are decoupled (Ds), the sub-
260 canopy respiration flux is an important quantity in the carbon mass balance
261 accounting for the advective loss (see Eq. 3 and Table 1). Re_{sub} was com-
262 puted using the modified relaxed eddy accumulation method of Thomas et al.
263 (2008) applied to the sub-canopy EC observations. The sub-canopy respira-
264 tion includes the contributions of the root, litter, and microbial respiration
265 and respiratory signals of the coarse and fine woody detritus and under-
266 story vegetation. For each 6-min interval, the method identifies individual
267 data points in the continuous EC time series that represent excursions from
268 similarity-theory predictions of the correlation structure between perturba-
269 tions of carbon dioxide and water vapor mixing ratios that are associated

270 with updrafts. These excursions can then collectively be used to construct
 271 a flux estimate. The underlying assumption is that updrafts ($w' > 0$) carry-
 272 ing the sub-canopy respiration signal have a different, positively correlated
 273 scalar-scalar fingerprint with $\chi_c' > 0$ and $q' > 0$. In contrast, motions carry-
 274 ing the fingerprint of the canopy crown show a negative correlation between
 275 χ_c' and q' . To include only those excursions into the computation of Re_{sub}
 276 that carry a strong sub-canopy respiration signal, a hyperbolic hole of the size
 277 $H_{Re}=0.5$ is applied to the correlation diagram constructed from plotting χ_c'
 278 versus q' . The excluded data falling into the area of $H_{Re} < 0.5$ may contami-
 279 nate the resulting sub-canopy respiration fluxes. Thomas et al. (2008) found
 280 a hole size of $H_{Re} = 0.25$ to be adequate using data from a range of different
 281 sites including the one investigated here for both above- and sub-canopy EC
 282 observations. A comparison of different hole sizes for our data yielded that
 283 Re_{sub} was only slightly sensitive to the choice of H_{Re} . Re_{sub} estimates for
 284 individual 6-min periods were only accepted and processed further when the
 285 fraction of excursions from similarity theory predictions exceeded 8 % of the
 286 total data (τ_{Re} , for details see Fig. 7 in Thomas et al., 2008). This criterion
 287 corresponds to $r_{c,q} \geq 0.1$ for the threshold of the transition from the decou-
 288 pled ground layer (Dg) to decoupled sub-canopy (Ds) regimes discussed later
 289 (see Fig. 5a). As mentioned in the previous section, flux estimates from a
 290 maximum of 5 subsequent 6-min periods were then combined into one 30-min
 291 average. If none of the 6-min periods within a 30-min interval fulfilled this
 292 condition, NEE could not be computed from observations according to Eq.
 293 3, but was modeled. One must note that $Re_{sub} \neq FCO_{2,sub}$ for the decoupled
 294 exchange regime since the turbulent eddies transporting the CO_2 molecules

295 released by the biologically active surfaces into the sub-canopy air do not
296 exchange with the layers aloft, but transport the scalar passively in up- and
297 downdrafts with their stochastic motion (see center panel of Fig. 4 in Zeeman
298 et al., 2012). As a result, the sub-canopy net carbon dioxide flux $FCO_{2,sub}$
299 underestimates the sub-canopy respiration flux assumed to be advected.

300 3.3. Computing alternative ecosystem respiration

301 In order to verify the ecosystem respiration estimates computed from the
302 new framework, they were compared against a completely independent es-
303 timate (see Sections 4.4 and 4.5). This alternative estimate of ecosystem
304 respiration RE_{alt} was constructed from chamber-based respiration measure-
305 ments and models of foliage and stem tissues, as well as total soil respiration.
306 Foliage respiration Re_{fol} was measured on a 8 branches from four individ-
307 ual trees from mid to upper canopy positions at the end of summer (DOY
308 243, 2011) using a portable photosynthesis system (Model LiCor 6400, Licor
309 Env., Lincoln, NE, USA). Branches were cut prior to sunrise, cut ends were
310 placed in water and returned to the lab where measurements were made at
311 a standard temperature (25 °C) between 06:00 and 09:00 local time. Foliage
312 samples were dried at 40 °C and analyzed for carbon and nitrogen concen-
313 trations (Model LECO CNS 2000 analyzer, Leco Corp., St. Joseph, MI,
314 USA). The resultant respiration values were standardized to 20 °C and com-
315 pared to published values from meta-analysis work where reference values of
316 mass based respiration (R_{20} , respiration values at 20 °C) were a function of
317 nitrogen (N) concentration (Reich et al., 2008). Our foliar respiration per
318 unit N values were slightly lower and required a correction factor of 1.13.
319 Additionally, a seasonal time course of N was used to scale the R_{20} values

320 for temporal N allocation patterns (Matson et al., 1994). The seasonally
321 dynamic R20 was then combined with the acclimation response of Q10 to
322 temperature history (Atkin and Tjoelker, 2003; Tjoelker et al., 2001). The
323 combined model was driven by sub-canopy air temperature while the base
324 rate was modulated by seasonal N and the Q10 was dynamic in response to
325 the 4 day running mean of sub-canopy air temperature. These mass-based
326 estimates of foliage respiration were then scaled to ground area using a leaf
327 mass per area of 78.12 gC m^{-2} and a site level PAI of $9.4 \text{ m}^2 \text{ m}^{-2}$. Stem res-
328 piration Re_{st} on a sapwood volume basis and concurrent stem temperature at
329 0.05 m depth was measured on 8 trees (see methodology in Law et al., 1999)
330 over multiple temperature conditions with 3 trees monitored continuously
331 for detailed temperature responses over different 3 day periods to provide
332 a large temperature range (temperature range of 11 to 19 °C). The aver-
333 aged temperature responses were adjusted using the mean values from the
334 remaining trees. This temperature response was then applied to continuous
335 records of soil temperature at 0.08 m depth, which closely matched point
336 measurements of stem temperature, and scaled to the site using year specific
337 estimates of sapwood volume. Soil respiration Re_{soil} was measured every 4
338 hours continuously on six chambers (methodology and system construction
339 can be found in Irvine and Law, 2002). Hours between measurements were
340 linearly interpolated while gaps larger than 2 measurement cycles were filled
341 using a temperature and soil moisture model which performed well at this site
342 and adequately captured the variation and timing of diurnal trends (Martin
343 et al., 2012). Temporal coverage of available 4 hour periods was $> 65 \%$.
344 All modeling and gap filling was performed on individual chambers; these

345 records were then averaged and scaled to the site using extensive spatial
346 data from 3 to 5 periodic spatial surveys per year on 20 separate locations
347 (see details in Thomas et al., 2009; Vickers et al., 2012). Soil temperature was
348 recorded hourly using 3 arrays of thermocouples at 0.02, 0.04, 0.08, 0.16, 0.32
349 and 0.64 m depth. Soil volumetric water content was measured using probes
350 integrating over the upper 0.3 m (Model CS615, Campbell Sci., Logan, UT,
351 USA).

352 3.4. Architecture and mean flow characteristics of the research site

353 Figure 2 placed here.

354 The dataset was collected in a 36 year-old mature Douglas-Fir (*Pseudotsuga*
355 *menziesii*, *Mirb.*) forest located in the coast range of western Oregon, USA
356 (AmeriFlux site US-Fir, 44.646° N latitude, 123.551° W longitude, 310 m el-
357 evation) surrounded by mountainous terrain with a flat saddle located to the
358 northeast of the tower at a distance of approximately 600 m (see Figure 1 in
359 Thomas, 2011). Observations described in this study were collected between
360 May 01, 2006 and Nov 03, 2011, i.e., over a period of approximately 5.5 years.
361 The vertical structure of the vegetation canopy with a mean height $h_c = 28$
362 m consists of a sparse understory composed mainly of Salal (*Gaultheria shal-*
363 *lon*, *Pursh*) with a plant height of up to 0.8 m agl and the main tree crown
364 space extending from 15 to 28 m agl separated by a clear bole space. The
365 canopy is very dense with a plant area index (PAI) of 9.4 m² m⁻² optically
366 measured in 2004 (Model LAI2000, Licor Env., Lincoln, NE, USA).

367 The vertical structure of the canopy and the high PAI lead to generally very
368 weak flows in the sub-canopy characterized by a median horizontal wind
369 speed $U \approx 0.5 \text{ m s}^{-1}$ and a 98 %-percentile of 1 m s^{-1} irrespective of light

370 regime (Fig. 2b). In contrast, above-canopy wind speeds are significantly
371 stronger with a median of 1.6 and 1.2 m s⁻¹ during day and night respec-
372 tively (Fig. 2a). In comparison with other sites, both above- and sub-canopy
373 speeds are generally weak because of the orographic sheltering provided by
374 the surrounding ridges. All wind speed distributions have a positive skewness
375 indicating a preference for weaker winds. Above the canopy, the distribution
376 of wind directions is strongly bimodal with preferred flows from west-south-
377 west (250°) and east (90°), which is representative of the regional flow. At
378 night, easterly drainages are relatively more common compared to daytime
379 at the expense of westerly flows (Fig. 2c). The comparison of wind direc-
380 tions between both levels indicates a systematic wind directional shear at the
381 site. For above-canopy flows centered around south-west and north-east the
382 wind directional shear reaches a maximum of 150° indicating almost opposite
383 flow directions. The existence of wind directional shear is insensitive to the
384 light regime, but its north-east maximum broadens at night to span a sec-
385 tor between north-west and north-east. Although these wind directions are
386 less common and the wind directional shear for the two main wind sectors
387 is typically $\Delta\phi \leq 40^\circ$, the directional shear in conjunction with significant
388 stratification generated additional mixing and impacted the classification into
389 exchange regimes as discussed in Section 4.2.

390 Previous studies investigating processes and mechanisms contributing to ver-
391 tical coupling and decoupling at this sites yielded two main results: Thomas
392 et al. (2008) reported that the attempt to estimate the sub-canopy daytime
393 respiration from above-canopy EC measurements failed as the dense over-
394 story acts as a mechanical barrier, which leads to a loss of the carbon dioxide

395 - water vapor fingerprint of air originating in the sub-canopy and ground lay-
396 ers. A follow-up study by Zeeman et al. (2012) found that larger-scale mo-
397 tions including mixing-layer type coherent structures and convective plumes
398 are not the main flow mode responsible for transporting scalars through
399 the dense overstory, but concluded that small-scale stochastic turbulence
400 was the main transport path. The authors confirmed that the scalar-scalar
401 cross-correlations observed above the canopy showed a weak dependence on
402 ground scalar sources and sinks. In comparison with other sites, both stud-
403 ies demonstrated that the method to partition the sub-canopy respiration
404 flux from NEE measurements succeeded in more open, sparse canopies, as
405 their architecture allows larger motions to penetrate deeply into the canopy
406 forcing the exchange with the sub-canopy and ground layers.

407 **4. Results & Discussion**

408 *4.1. Comparing possible mixing indicators*

409 Figure 3 placed here.

410 In order to assess the strength of the turbulent mixing that serves as an esti-
411 mator for the degree of coupling across the canopy profile, one needs to find
412 a suitable proxy. We here focus on quantities that are commonly evaluated
413 during EC post-processing. The comparison of friction velocity and the stan-
414 dard deviation of the vertical velocity variance between the above- and sub-
415 canopy levels yielded similarities, but also important differences (Fig. 3a, c):
416 At night when winds are weak, the turbulence strength in the sub-canopy is
417 largely independent of that above the canopy showing an asymptotic behav-
418 ior toward a constant value, but scales almost linearly after the above-canopy

419 flow crossed the threshold of $u_* = 0.25 \text{ m s}^{-1}$ and $\sigma_w 0.45 \text{ m s}^{-1}$ respectively.
420 During the day when winds are stronger, the behavior is similar, but the
421 thresholds for the transition from the asymptotic to the linear response are
422 smaller at 0.12 and 0.20 m s^{-1} respectively. The scatter on both abscissa and
423 ordinate is significant, particularly for the asymptotic, independent regime.
424 The independence of the sub-canopy turbulence strength from that aloft
425 indicates either different, independently acting mechanisms generating or
426 destroying turbulence, or that the above- and sub-canopy flows are largely
427 decoupled, or a combination of both. In case of the former, the above-canopy
428 flow and its wind speed shear at the canopy top are not the main source of
429 turbulence in the sub-canopy as the penetration of eddies deep into the trunk
430 space is limited by the dense overstory. This interpretation is in agreement
431 with the findings of Zeeman et al. (2012) who investigated the penetration
432 depth of larger-scale including mixing-layer type structures and convective
433 plumes at the same site. If the above- and sub-canopy flows are decoupled,
434 the existence and passage of sub-meso motions on scales larger than those of
435 turbulence could be responsible for the generation of sub-canopy turbulence
436 (Thomas, 2011). Sub-meso motions are always present in the atmosphere,
437 but tend to dominate dispersion and turbulent diffusion only when the mean
438 flow is very weak (Mahrt et al., 2009), such as in canopy flows. Note that
439 for a given turbulence strength above the canopy, the sub-canopy equivalent
440 is larger at night than during the day. An explanation for this observation is
441 the temporally reversed stability regime between the above- and sub-canopy
442 layers typical for dense forests: during the day when the radiative heating
443 of the overstory leads to free convection and a maximum temperature lo-

444 cated near the PAI maximum, the sub-canopy is slightly stably stratified
445 acting to suppress turbulence; at night when progressive radiative cooling of
446 the crown leads to increasing static stability in the above-canopy layer, the
447 sub-canopy layer becomes isothermal or slightly unstable acting to enhance
448 turbulence. One feature worth noting is the increase of sub-canopy turbu-
449 lence strength for very weak above-canopy turbulence, which was associated
450 with the generation of additional mixing by wind directional shear, which
451 will be discussed in the next section. Wind speeds in the above- and sub-
452 canopy layers are related by a weakly non-linear relationship depending on
453 the light regime (Fig. 3b). As for the turbulence, the nighttime sub-canopy
454 flow tends to be stronger for a given above-canopy wind speed exceeding \geq
455 1.5 m s^{-1} than during the daylight hours. During the day, both above-canopy
456 u_* and σ_w respond non-linearly to an increase in wind speed up until ≈ 1.5
457 m s^{-1} , after which the relationship becomes strictly linear (Fig. 3d). The
458 initial maximum observed around 1 m s^{-1} can be explained by the gener-
459 ation of turbulence through buoyancy resulting from the radiative heating
460 of the crown, which is important when flows are weak. The linear response
461 for stronger flows indicates the shear-driven turbulence. For the subcanopy,
462 the dependence of the turbulence strength on the mean sub-canopy flow was
463 much weaker (not shown here), accentuating the importance of generating
464 mechanisms for turbulence other than the commonly discussed wind speed
465 shear. At night, both u_* and σ_w scale linearly with wind speed and approach
466 a constant minimum value for $U \rightarrow 0 \text{ m s}^{-1}$. The identical behavior was
467 observed by Mahrt et al. (2012) for calm near-surface flows above short grass
468 during stable conditions at three different sites in structured, mountainous

469 terrain. Both ratios $\sigma_w u_*^{-1}$ and $\sigma_w U^{-1}$ showed a maximum for very weak
470 flows while asymptotically approaching a constant value for stronger flows
471 (Fig. 3d).

472 For the purpose of this study which aims at improving NEE estimates through
473 determining meaningful exchange regimes based on multi-level flow statistics,
474 the most suitable proxies are either the friction velocity or the standard de-
475 viation of the vertical velocity variance. Both quantities showed an obvious
476 transition between regimes that may be linked to vertical coupling and decou-
477 pling processes as discussed below. We prefer the vertical velocity variance
478 σ_w as the primary indicator of mixing in this study for the following reasons:
479 i) it relates the turbulence strength directly to properties of the mean tur-
480 bulent flow and not to its turbulent momentum transport as in the case of
481 u_* ; ultimately, the turbulent eddies and their associated vertical velocity per-
482 turbations lead to diffusion and mixing, and not the momentum transport;
483 ii) it appears to be more sensitive to alternative generating mechanisms for
484 turbulence including sub-meso motions and wind directional shear, which are
485 important in dense canopies; and iii) both day- and nighttime values con-
486 verge on to the same linear relationship and show less scatter for stronger
487 mixing facilitating the selection of universal thresholds presented in the next
488 section. Our findings are in agreement with Acevedo et al. (2009) who also
489 preferred σ_w over u_* as the mixing indicator to filter nighttime CO₂ fluxes for
490 several forest sites in Brazil. With their choice of σ_w , thresholds for limiting
491 vertical mixing were easier to identify and nighttime NEE estimates were
492 increased in magnitude.

493 *4.2. Enhancing mixing through wind directional shear*

494 Figure 4 placed here.

495 The systematic, substantial directional shear found in the wind climatology
 496 (Fig. 2d) led to an increase in mixing strength at night when combined with
 497 significant above-canopy flows (see Fig. 4a). For weaker flows, the enhanced
 498 mixing generated by the wind directional shear is expected to lead to conver-
 499 gence of both flows resulting in $|\Delta\phi_{top-sub}| \rightarrow 0^\circ$. Since wind speed shear, i.e.,
 500 the vertical differences of horizontal wind speeds across the canopy profile,
 501 is commonly thought of as the main mechanism inducing flow instabilities
 502 and creating overturning eddies to generate turbulence at night, the wind
 503 directional shear is typically not included in the dynamic stability analysis
 504 for canopy flows (e.g. Paw U et al., 1992; Raupach et al., 1996). However,
 505 the gradient Richardson number, defined as the ratio of buoyancy to shear
 506 terms in the turbulent kinetic energy equation (e.g. Stull, 2000)

$$Ri_{gr} = \frac{g}{\theta_v} \Delta\theta_v \Delta z (\Delta X^2 + \Delta Y^2)^{-1}, \quad (4)$$

507 theoretically includes the effect of directional shear through the differences
 508 in zonal ΔX and meridional ΔY wind components; g is the gravitational
 509 acceleration on earth and θ_v the potential virtual air temperature. Evalu-
 510 ating Ri_{gr} yielded that the sub-canopy mixing strength $\sigma_{w,sub}$ was enhanced
 511 only for $U_{top} \geq 1 \text{ m s}^{-1}$ in combination with strong stability $Ri_{gr} > 1$ and
 512 $|\Delta\phi_{top-sub}| > 70^\circ$ (Fig. 4b). This enhancement was 50 to 80 % compared
 513 to periods when either $|\Delta\phi_{top-sub}| < 50^\circ$ or $0.2 \leq Ri_{gr} \leq 4$ (Fig. 4b, Eq.
 514 4). This additional mixing impacted the found exchange regime and thus
 515 the computation of NEE. For moderately strong winds $2 < U_{top} \leq 3 \text{ m s}^{-1}$

516 the additional mixing led to a fully mixed canopy, albeit the strong stability
517 $Ri_{gr} > 10$ (Fig. 4b). This observation is, however, counter-intuitive given
518 the definition of Ri_{gr} and requires some discussion. Since the directional
519 shear contributes to the denominator in Eq. 4, one would expect Ri_{gr} to
520 decrease with increasing $|\Delta\phi|$. Since the opposite behavior is observed, we
521 conclude that the gradient Richardson number insufficiently reflects the addi-
522 tional mixing induced by the wind directional shear. In contrast, Burns et al.
523 (2011) evaluated flow and scalar statistics in a moderately dense coniferous
524 forest in mountainous terrain and found that binning by the Richardson num-
525 ber was useful to explain vertical scalar mixing. However, the authors used
526 the bulk formulation of the Richardson number Ri_b that is sensitive to wind
527 speed shear only. At their site, the mean wind directional shear vanished
528 for $Ri_b > 0.1$, while its variance increased significantly possibly indicating
529 meandering of the flow.

530 Note that in our study the presence of wind directional shear between above-
531 and sub-canopy flows is not synonymous with vertical decoupling. On the
532 contrary, it leads to enhanced mixing and a fully coupled canopy under con-
533 ditions of significant stratification. Note that while mixing is enhanced, it is
534 still weak compared to unstable conditions.

535 *4.3. Defining mixing thresholds for the exchange regimes*

536 Figure 5 placed here.

537 After identifying a suitable proxy for the turbulent mixing strength, the next
538 step is to connect its behavior to the signals of the biological carbon sinks
539 and sources. The goal is to define biologically meaningful thresholds for the
540 exchange regimes proposed earlier. To that end, we use the same concept

541 applied earlier when quantifying the sub-canopy respiration flux (see Section
542 3.2) to distinguish between turbulent eddies that carry the fingerprint of the
543 ground layer from that of the canopy layer. One recalls that following this
544 concept, eddies carrying the fingerprint of the ground layer are expected to
545 show a positive correlation between χ_c' and q' . In contrast, eddies carrying
546 the fingerprint of the overstory are expected to show a negative correlation
547 between χ_c' and q' .

548 The correlation coefficient between perturbations of carbon dioxide and wa-
549 ter vapor perturbations $r_{c,q}$ evaluated from the sub-canopy EC data showed
550 a distinct behavior with the turbulent mixing strength irrespective of light
551 regime and seasonality (Fig. 5a): for very weak mixing $\sigma_{w,sub} \leq 0.03 \text{ m s}^{-1}$
552 the correlation is close to zero and erratic with frequent changes in sign. For
553 $0.03 < \sigma_{w,sub} \leq 0.10 \text{ m s}^{-1}$ the correlation becomes systematically positive
554 with $r_{c,q} \geq 0.1$ and increases with increasing mixing strength. If the mixing
555 strength increases beyond $\sigma_{w,sub} \geq 0.10 \text{ m s}^{-1}$, the correlation systematically
556 declines again and ultimately becomes negative for $\sigma_{w,sub} > 0.2 \text{ m s}^{-1}$. We
557 here define the first threshold at $\sigma_{w,sub} = 0.03 \text{ m s}^{-1}$ as the transition from the
558 decoupled ground layer (Dg) to the decoupled sub-canopy layer (Ds) regime.
559 Beyond this threshold the sub-canopy EC observations clearly capture the
560 scalar-scalar fingerprint of the sources in the ground layer. Note that the
561 magnitude of the correlation is sensitive to sub-canopy air temperature as
562 the magnitude of the sub-canopy fluxes of respiration and evapotranspiration
563 is seasonally dependent. Strongest correlations were found for the greatest
564 temperatures when fluxes are maximum. The second transition at $\sigma_{w,sub} =$
565 0.10 m s^{-1} reflects the transition from the decoupled sub-canopy (Ds) to

566 the fully coupled canopy (C) regime. As eddies start to communicate across
567 the entire canopy profile, the growing influence of the scalar-scalar finger-
568 print of the overstory with $r_{c,q} < 0$ leads to a progressive decorrelation of the
569 ground layer fingerprint with $r_{c,q} > 0$. Note that correlating $r_{c,q}$ to the sub-
570 canopy friction velocity showed the same overall behavior (not shown here),
571 but yielded different transitions at $u_{*,sub} = 0.02 \text{ m s}^{-1}$ from Dg to Ds, and at
572 $u_{*,sub} = 0.10 \text{ m s}^{-1}$ from Ds to C (Table 1). Our interpretation of the con-
573 nection between increasing sub-canopy mixing strength and enhanced mass
574 transport exchanging scalars across the entire canopy is different from that
575 based on analyzing the momentum exchange in tall, but low PAI forests (e.g.
576 Shaw and Zhang, 1992; Dupont et al., 2012). Based on observations and
577 large eddy simulations in stands with $\text{PAI} \leq 2.5$, these studies concluded
578 that much of the sub-canopy turbulence is generated by pressure transport
579 from aloft. While pressure diffusion may be an important mechanism gen-
580 erating mixing in the sub-canopy, these studies did not directly evaluate the
581 air and scalar exchange which may follow perturbations in the pressure field.
582 Flow visualizations using machine-generated fog at our site confirmed our
583 interpretation of mass exchange between the ground, sub-canopy, overstory,
584 and above-canopy layers that is associated with significant sub-canopy mix-
585 ing. The validity of our interpretation to define meaningful exchange regimes
586 based on $\sigma_{w,sub}$ is further discussed in Section 4.6.
587 For above-canopy observations, the concept of $r_{c,q}$ changing sign is not appli-
588 cable since carbon dioxide and water vapor perturbations are systematically
589 negatively correlated. We here use the traditional nighttime plot of $(\text{FCO}_{2,top}$
590 $+ \Delta S)$ versus turbulent mixing strength to define a threshold for coupling be-

591 tween the above-canopy EC system and the tree crown. At night, ($\text{FCO}_{2,top} +$
592 ΔS) becomes independent of turbulent mixing strength for $\sigma_{w,top} \geq 0.45 \text{ m s}^{-1}$
593 independent of season (Fig. 5b). The friction velocity equivalent was $u_{*,top} \geq$
594 0.35 m s^{-1} (not shown here). We here select this threshold as the transition
595 from the decoupled sub-canopy layer (Ds) to the fully coupled canopy regime
596 (C). For daytime conditions, we assume that above-canopy turbulent mixing
597 is not limiting as long as sub-canopy mixing is sufficiently strong.
598 Comparing the probability of occurrence across all exchange regimes, one
599 finds that the decoupled sub-canopy is the most abundant regime occupying
600 approximately 65 and 88 % of the observations during the day and night,
601 respectively, while the fully coupled canopy occurred 34 and 11 % and the
602 decoupled ground layer $< 1\%$ of the time, respectively (Table 1).

603 *4.4. Comparison of chamber-based soil and sub-canopy respiration*

604 Figure 6 placed here.

605 Since the respiration components captured by the sub-canopy respiration
606 Re_{sub} and the soil respiration chambers Re_{soil} are largely identical, one can
607 investigate the differences between these fluxes with regard to turbulent mix-
608 ing strength relevant for the computation of NEE. Both quantities include
609 contributions by the roots, litter, and microbes as well as fine woody detritus,
610 while Re_{soil} additionally captures the respiration from coarse woody debris
611 and the understory vegetation. The fraction of the latter two components to
612 total ecosystem respiration, however, is typically small, particularly at our
613 site, which is characterized by low light levels and few coarse woody debris
614 on the ground.

615 For the decoupled sub-canopy regime (Ds), Re_{sub} was sensitive to the turbu-

616 lent mixing strength, while Re_{soil} was independent of $\sigma_{w,sub}$ (Fig. 6). Transi-
617 tioning to a fully coupled canopy (C), Re_{sub} also became independent and ap-
618 proached values that were approximately 10 to 20 % lower compared to Re_{soil}
619 for a given temperature class. Note that the boundaries of the temperature
620 classes are equivalent since they were determined by fitting a quadratic re-
621 gression of air versus soil temperature. The sub-canopy respiration observed
622 for the highest temperatures was an exception since it failed to converge to
623 a value independent of σ_w . Separating fluxes in this broad temperature class
624 yielded that the sustained increase was caused by a small group of extreme
625 values observed for very high temperatures $T > 26$ °C. Differences between
626 the flow conditions in the sampling volume of each technique can explain the
627 differences for the Ds regime: diffusion of CO₂ molecules out of the soil and
628 transport through the sampling path of the sub-canopy EC system is limited
629 by the weak turbulent mixing for $\sigma_{w,sub} \leq 0.10$ m s⁻¹. The sub-canopy trans-
630 port in the dense canopy therefore remains limiting to the flux, despite the
631 steepness of the vertical CO₂ gradient of several hundred ppm m⁻¹. In con-
632 trast, the transport inside the soil respiration chamber is different since both
633 molecular diffusion and turbulent mixing are artificially enhanced by actively
634 pumping air. CO₂ concentration gradients at the soil-air interface are simi-
635 lar for both techniques. We emphasize that both respiration estimates Re_{sub}
636 and Re_{soil} are biologically meaningful, while they show systematic differences
637 arising from the mixing strength in their sampling volumes.

638 4.5. Improved carbon balance

639 Figure 7 placed here. Table 2 placed here.

640 We now proceed and compute the improved carbon balance using the pro-

641 posed framework with the exchange regimes and equations summarized in
642 Table 1. Accounting for the advective losses to total ecosystem respiration
643 for the decoupled sub-canopy regime (Ds) led to a close agreement between
644 nighttime NEE \equiv RE and the independently computed alternative ecosystem
645 respiration RE_{alt} over a wide temperature range (Fig. 7). The interquartile
646 range of RE_{alt} closely matched that of NEE for a given temperature bin,
647 while the scatter in NEE was significantly larger. One must recall that stem
648 and foliage respiration, which on average contributed 18 % and 22 % respec-
649 tively to RE_{alt} , were modeled using air temperature as the only predictor
650 (see Section 3.3). Hence, their estimates are expected to closely follow the
651 temperature sensitivity of the model acting to reducing the scatter in RE_{alt} .
652 In addition, the soil CO_2 efflux, which on average contributed the remaining
653 60 % to RE_{alt} , represent a spatial averages across six different chambers also
654 acting to reduce the scatter. In contrast, the nighttime respiration computed
655 as $(FCO_{2,top} + \Delta S)$ derived using a 'traditional' filtering and gap-filling ap-
656 proach was significantly smaller by 36 % compared to both improved RE
657 from the new framework and RE_{alt} . The traditional approach was based
658 solely on single-level above-canopy EC fluxes with nighttime values filtered
659 by a friction velocity criterion of $u_* \geq 0.35 \text{ m s}^{-1}$ determined from plotting
660 $(FCO_{2,top} + \Delta S)$ against u_* for several temperature bins (Goulden et al.,
661 1996). Systematic discrepancies between RE and RE_{alt} were found only for
662 the highest temperatures $T_{sub} \geq 23 \text{ }^\circ C$ occurring during late summer. Dur-
663 ing this season, the trees increasingly experience soil moisture limitations in
664 the late afternoon hours (Fig. 7), which leads to a reduction of photosyn-
665 thesis and autotrophic respiration, and thus ecosystem respiration captured

666 by the EC fluxes. The soil moisture limitations may further dampen het-
667 erotrophic soil respiration, which enhances this effect. In contrast, stem and
668 foliage respiration were modeled using a temperature-dependent model that
669 lacks moisture limitations, while Re_{soil} showed a reduction during the dry,
670 late summer (not shown here). These discrepancies however, had a negligi-
671 ble effect on the uncertainty of the carbon balance estimates since very few
672 observations fell into the high temperature classes (see grey solid line in Fig.
673 7).

674 The annual improved NEE was significantly reduced compared to that de-
675 rived from the 'traditional approach' yielding on average $622 \text{ gC m}^{-2} \text{ yr}^{-1}$ less
676 carbon uptake (Tab. 2). The 6-year average of improved NEE was -479 gC
677 $\text{m}^{-2} \text{ yr}^{-1}$ compared to $-1101 \text{ gC m}^{-2} \text{ yr}^{-1}$ from the traditional approach. The
678 difference between the two methods was attributed to an enhanced ecosys-
679 tem respiration, while gross ecosystem productivity, computed as the resid-
680 ual from $\text{GEP}=\text{NEE}-\text{RE}$, remained almost unchanged. The GEP estimates
681 agreed well with those predicted by Waring et al. (2008) for the same site
682 using the process-based forest growth model 3-PG driven by observed me-
683 teorologic data. As expected, improving NEE was primarily accomplished
684 by increasing nighttime NEE by including the advective losses estimated
685 through the sub-canopy respiration (Eq. 3). The improved annual NEE
686 agreed well estimates for a Douglas-Fir forest of similar age, stocking and
687 climate (Chen et al., 2009; Jassal et al., 2010, 2007; Krishnan et al., 2009),
688 and fit more favorably into global meta-analyses given the mean temperature
689 and precipitation at the site (Luyssaert et al., 2007). Note that we did not
690 apply an 'ecological plausibility filtering' to force $\text{RE}>0$ and $\text{GEP}<0 \mu\text{mol}$

691 $\text{m}^{-2} \text{s}^{-1}$. However, the fractions for these 'ecologically implausible' periods
692 were 0.7 % and 2 % over the entire observational period, respectively. The
693 positive GEP values were mainly associated with hours shortly after dawn
694 and before dusk and may thus be biologically plausible, while the negative
695 nighttime RE values are expected to result from random variability associ-
696 ated with turbulent fluxes. Given their small abundance, these anomalies
697 have a negligible impact on annual NEE estimates. The fraction of observed
698 nighttime NEE data retained to constrain the Arrhenius-type model used for
699 gap-filling and to model daytime respiration was doubled compared to the
700 traditional approach (35 % versus 17 %, respectively, Tab. 2). In contrast,
701 the fractions of observed NEE during the day remained unchanged.

702 *4.6. Limitations and risks of the approach*

703 The improved NEE approach bears the risk of double-counting an un-
704 known fraction of CO_2 molecules by adding the missing, advected sub-canopy
705 respiration to NEE for the decoupled sub-canopy regime (Ds). In other
706 words, our assumption of a 'hermetical' decoupling between the sub-canopy
707 and above-canopy layers may not be correct in all instances. However, several
708 observations support our claim that the diagnosed decoupling is meaningful
709 and that a significant fraction of total ecosystem respiration is not captured
710 by $(\text{FCO}_{2,top} + \Delta\text{S})$:

- 711 • The identified thresholds of $\sigma_{w,top} \geq 0.45 \text{ m s}^{-1}$ and $\sigma_{w,sub} \geq 0.10 \text{ m s}^{-1}$
712 for transition from the Ds to C regime agree well with the transi-
713 tion from the asymptotic to the linear response regime in the mixing
714 strength (see Fig. 3c, and Section 4.1). This agreement suggests a co-
715 hesion between the physical transport and the biological signal of CO_2

716 sinks and sources in the canopy. A recent review of canopy flows in
717 mountainous terrain by Belcher et al. (2008) also supports the assump-
718 tion of decoupled canopy and sub-canopy layers from the flow aloft.
719 Their explanation is based on fundamental differences in momentum
720 and heat transport from the air to the canopy elements, in spite of
721 significant above-canopy flows.

722 • For $\sigma_{w,top} \leq 0.45 \text{ m s}^{-1}$, the traditional nighttime ($\text{FCO}_{2,top} + \Delta S$)
723 indicated a dependence on turbulence mixing strength as well as on
724 temperature (Fig. 5b). Since nighttime ecosystem respiration is pri-
725 marily driven by temperature except for the late summer season where
726 soil moisture may become limiting, this observation is evidence that the
727 above-canopy EC flux does not integrate over all sinks and sources sup-
728 porting the argument of Goulden et al. (1996). This underestimation is
729 systematic for low turbulence conditions that are intrinsically related
730 to the dense, tall overstory of the Douglas-fir trees and exacerbated by
731 the sheltering topography.

732 • The variability of the sub-canopy flow and temperature fields and the
733 resultant horizontal advection of sensible heat at the site was inves-
734 tigated by Thomas (2011). Both the vector and scalar field showed
735 systematic variability that was dependent upon the time scale, the loca-
736 tion within the domain, and the total size of the domain. Furthermore,
737 horizontal advection of heat was significant and systematic in sign and
738 magnitude for horizontal scales exceeding 180 m. We therefore consider
739 our assumption of a systematic loss in ecosystem respiration through

740 advection to be plausible. However, the lack of direct evidence for the
741 scalar similarity between the advective transport of carbon dioxide and
742 sensible heat, and the absence of direct estimates of sub-canopy CO₂
743 advection such as in Staebler and Fitzjarrald (2004) remains a weak-
744 ness of the investigation presented here.

745
746 In the literature, there exists indirect observational and modeling evidence
747 for the importance of advection to the carbon dioxide budget: Sun et al.
748 (2006) used a large eddy simulation method to model the transport of car-
749 bon dioxide for a similar tall forest in sloped mountainous terrain. They
750 found that advection caused by drainage flows and a weak land breeze was a
751 significant term in the scalar budget, which resulted in an underestimation
752 of nighttime respiration of about 20 %. Despite the similarities in canopy ar-
753 chitecture and geometry of the terrain between their study area and our site,
754 their nighttime NEE was dominated by the vertical turbulent flux, which
755 disagrees with our findings. However, their transport was dominated by the
756 portion contributed by the sub-grid scale parameterizations, so a direct com-
757 parison cannot be done. Van Gorsel et al. (2007) and a follow-up study by van
758 Gorsel et al. (2009) found indirect evidence for advective loss of carbon diox-
759 ide due to drainage flows. The authors investigated the temporal dynamics
760 of $(F_{CO_2} + \Delta S)$ over the diurnal period showing a clear maximum during
761 the early evening, whose magnitude agreed well with independent ecosys-
762 tem respiration estimates from chambers. With progression of the night,
763 the difference between the two respiration estimates increased indicating the
764 growing importance of advective transport and loss of CO₂ in the simplified

765 carbon budget.

766 One must keep in mind that the transitions between all three exchange
767 regimes are rather gradual, in spite of our definition of numerically exact
768 thresholds. The significant variability in relationships between correlation
769 coefficients, ecosystem carbon and water vapor fluxes, and turbulent mix-
770 ing proxies will inevitably lead to a false classification of exchange regimes
771 for individual periods. However, we expect the determined thresholds to be
772 suitable for the purpose of constructing seasonal and annual carbon budgets.
773 Errors resulting from miss-classification are assumed to be random and thus
774 to cancel out over a sufficiently large sample size.

775 Double-counting of CO₂ molecules and the uncertainty associated with the
776 determination of thresholds for the exchange regimes remain to be realistic
777 limitations to the improved NEE approach. Its estimates may therefore serve
778 as a lower limit for the strength of the terrestrial carbon sink resulting from
779 a possibly overestimated ecosystem respiration. Considering that the well-
780 documented nighttime limitation to the traditional EC approach results in a
781 systematic underestimation of ecosystem respiration and an overestimation
782 of the carbon sink, the framework proposed here may make an important
783 contribution to the discussion of uncertainty in NEE estimates for tall, dense
784 forests.

785 **5. Conclusions**

786 We arrive at the following conclusions with respect to our objectives listed
787 in the introduction:

- 788 • The standard deviation of the vertical velocity variance σ_w was the

789 most suitable proxy for the turbulent mixing strength in this dense
790 canopy. It is preferable over alternatives including the friction velocity
791 since it is a direct measure of the turbulent kinetic energy rather than
792 its momentum flux. It was more sensitive to alternative mechanisms
793 generating turbulence creating additional mixing. Alternative mech-
794 anisms include wind directional shear in combination with significant
795 above-canopy flows and strong stability, as well as sub-meso structures
796 present in the sub-canopy (Thomas, 2011). The standard deviation of
797 the vertical velocity variance is routinely calculated, which makes it
798 readily available to the user.

799 • The proposed framework of exchange regimes was found to be mean-
800 ingful with respect to both biological scalar signals and the physical
801 transport of the turbulent eddies. In this dense canopy, it is imper-
802 ative to have direct measurements of the sub-canopy mixing strength
803 to estimate the vertical coupling. We therefore recommend including
804 continuous, concurrent sub-canopy EC observations to the standard
805 measurement protocol in addition to the traditional above-canopy flux
806 measurements in dense, tall forests.

807 • The improved NEE method led to a significant enhancement of the
808 fraction of observations used to parametrize the analytical gap-filling
809 models and to compute the annual carbon budget. For nighttime data,
810 this fraction was doubled compared to the traditional approach based
811 on above-canopy EC fluxes at a single level filtered by a u_* criterion.
812 It may therefore help to add more dense forests to continental and

813 global scale carbon cycle studies by correcting the bias toward sites
814 with shorter and more open canopies characterized by stronger flows
815 and more vigorous turbulent mixing.

816 • The annual NEE computed from the improved method was significantly
817 lower than that of the traditional approach by $\approx 620 \text{ gC m}^{-2} \text{ yr}^{-1}$.
818 The improved estimated strength of the carbon sink for this forest
819 compared better to estimates reported for similar ecosystems. Based
820 on the discussion of risks and limitations of the method, this estimate
821 can be regarded as a lower limit.

822 • Improving upon the traditional approach used to compute NEE from
823 single-level EC observations was not difficult. The improved method
824 does not require additional instrumentation or measurements assuming
825 that EC observations from at least two levels exist. It can readily be
826 applied to historic data. The thresholds for the transitions between
827 exchange regimes are expected to be site-specific, while their underly-
828 ing concepts and correlations can be transferred. Since NEE estimates
829 from the improved and traditional approaches yielded large discrepan-
830 cies, we conclude that dense and tall forests do not lend themselves to
831 standardized processing, which has been proposed for various continen-
832 tal networks. The benefit of having an additional, lower boundary for
833 NEE reflecting an enhanced ecosystem respiration may, however, out-
834 weigh the additional effort required to determine the exchange regimes
835 and to compute sub-canopy respiration.

836 **6. Acknowledgements**

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Table 1: Mass balance equations for net ecosystem exchange (NEE) comprised of the vertical turbulent CO₂ flux (FCO₂), the temporal change in storage term (ΔS), net photosynthesis (GEP), and subcanopy respiration (Re_{sub} , determined following Thomas et al. (2008) applied to the sub-canopy EC observations, see Section 3.2) depending on observation height within the above-canopy layer (*top*) or sub-canopy layer (*sub*) for the three different exchange regimes described in Section 2: Decoupled ground layer (Dg), decoupled sub-canopy (Ds), and fully coupled canopy (C). Also shown are the determined thresholds of the standard deviation of the vertical velocity variance, $\sigma_w = \sqrt{\overline{w'w'}}$, where w' is the vertical velocity perturbation.

Exchange regime	Mass balance		Criteria & Thresholds		Fraction [#]	
	Day	Night	Day	Night	Day	Night
Dg	NEE not measurable	NEE not measurable	$\sigma_{w,sub} < 0.3 \text{ m s}^{-1}$	$\sigma_{w,sub} < 0.3 \text{ m s}^{-1}$	< 1 %	< 1 %
Ds	NEE = FCO _{2,top} + Re _{sub} + (GEP _{sub}) + ΔS	NEE = FCO _{2,top} + Re _{sub} + ΔS	$0.03 \leq \sigma_{w,sub} < 0.10 \text{ m s}^{-1}$	$0.03 \leq \sigma_{w,sub} < 0.10 \text{ m s}^{-1}$	65 %	88 %
C	NEE = FCO _{2,top} + ΔS	NEE = FCO _{2,top} + ΔS	$\sigma_{w,sub} \geq 0.10 \text{ m s}^{-1}$	$\sigma_{w,sub} \geq 0.10 \text{ m s}^{-1}$ & $\sigma_{w,top} \geq 0.45 \text{ m s}^{-1}$	34 %	11 %

[#]: averaged over entire observational period 2006 through 2011

Table 2: Annual sums of carbon balance components for the observational period 2006 through 2011: improved net ecosystem exchange (NEE), ecosystem respiration (RE), and gross ecosystem productivity (GEP) computed using the equations listed in Table 1 from both above- and sub-canopy EC data. For comparison, estimates are also provided for a traditional approach to compute components solely based on the above-canopy EC observations with nighttime fluxes filtered by a critical threshold of $u_{*,top} \geq 0.35 \text{ m s}^{-1}$ (NEE_{trad} , RE_{trad} , and GEP_{trad}). Statistics are provided for percentage (p) of observed (obs) and modeled (mod) data for day (d) and night (n) conditions. Also listed are the annual sums of alternative ecosystem respiration (RE_{alt}) based on scaled soil CO_2 efflux chambers (Re_{soil}), and modeled stem (Re_{st}) and foliage respiration (Re_{fol}) described in Section 3.3. Gaps in EC data were filled using Arrhenius type and Michaelis-Menten type models fitted to the night- and daytime data, respectively, for each year (see Ruppert et al., 2006; Thomas et al., 2009, for details). $NEE < 0$ indicates a net carbon sink.

Year	Improved approach					Traditional approach					Difference=			RE_{alt} [gC m ⁻² yr ⁻¹]
	Two heights; mixing based on σ_w			p (obs/mod ⁺)		Single height; mixing based on u_*			p (obs/mod ⁺)		Improved-Traditional			
	NEE	RE	GEP	d&n	n	NEE_{trad}	RE_{trad}	GEP_{trad}	d&n	n	Δ NEE	Δ RE	Δ GEP	
[gC m ⁻² yr ⁻¹]			[%]		[gC m ⁻² yr ⁻¹]			[%]		[gC m ⁻² yr ⁻¹]				
2006*	-694	1972	-2666	37/63	31/69	-1348	1254	-2603	39/61	15/85	655	718	63	2086
2007	-563	2009	-2572	52/48	48/52	-1258	1233	-2492	48/52	20/80	695	776	-80	2076
2008	-495	2092	-2588	48/52	46/54	-1121	1557	-2678	49/51	21/79	626	535	91	1994
2009	-305	2291	-2596	40/60	40/60	-882	1689	-2571	41/59	18/81	577	602	-25	1976
2010	-336	2231	-2567	33/67	32/68	-926	1599	-2526	41/59	20/80	590	632	-42	1964
2011	-482	1847	-2329	16/84	15/85	-1072	1275	-2347	21/79	9/91	590	572	18	2008

*: annual sum was computed over period 01-May-2006 through 01-May-2007; model coefficients determined from observations in 2007 were used for gap-filling.

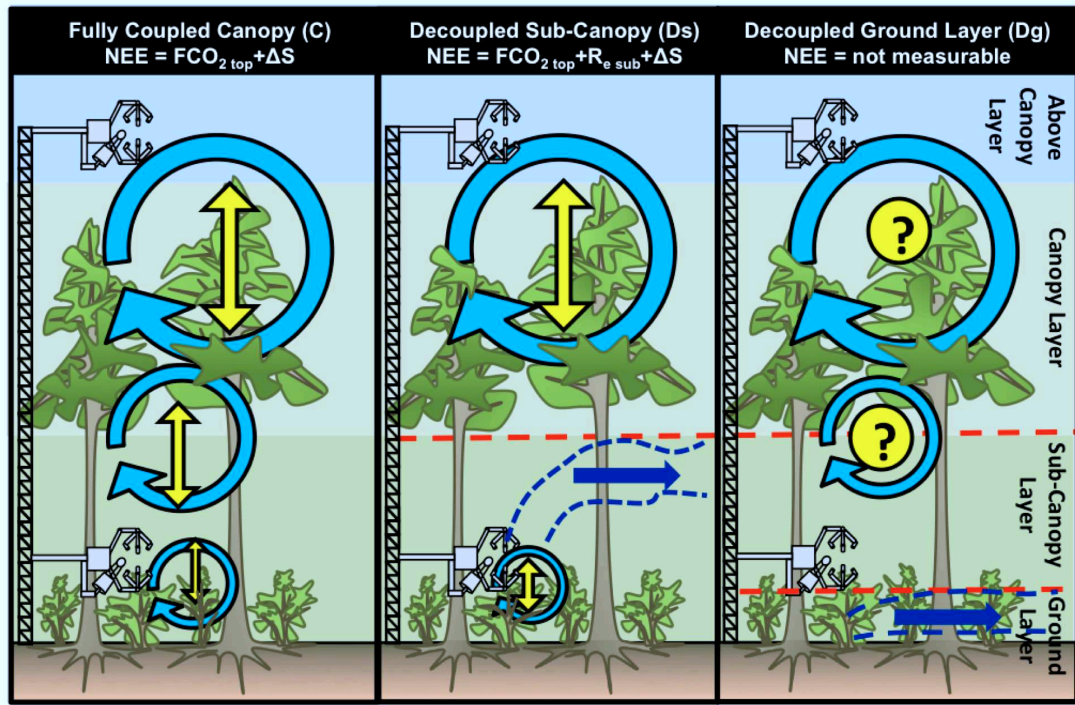


Figure 1: Illustration of the conceptual framework used to improve estimates of net ecosystem exchange (NEE) in tall, dense canopies. The vertical canopy coupling decreases from left to right from the fully coupled canopy (C), over the decoupled sub-canopy (Ds), to the decoupled ground layer (Dg) regimes. For the fully coupled canopy, it is assumed that the above-canopy eddy covariance (EC) flux ($FCO_{2,top}$) integrates the turbulent exchange of all carbon sinks and sources. For the decoupled sub-canopy regime, the above-canopy EC flux only integrates over sinks and sources located in the overstory, while the sub-canopy EC observations can be used to indirectly determine the advective loss through accounting for the sub-canopy respiration (Re_{sub}). For the decoupled ground layer regime, turbulent mixing is very weak and advection is assumed to be the dominant term in the carbon balance, which is not captured by any of the EC systems. ΔS is the temporal change of the storage term. Horizontal red dashed lines indicate decoupling between layers, dark blue arrows symbolize advection, and yellow vertical arrows turbulent exchange. Note that the light-blue circles symbolize turbulent structures in general, and not specific flow modes or length scales. See Section 2 for additional explanation.

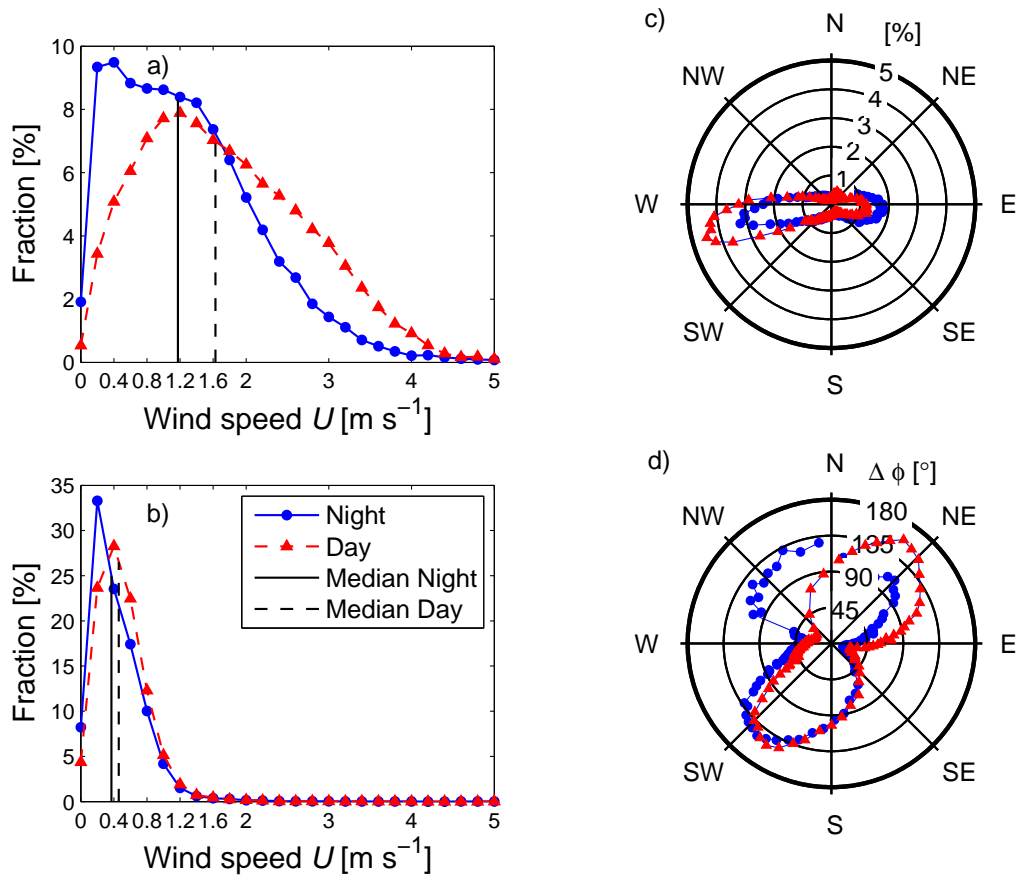


Figure 2: Characteristics of the mean flow at the dense, mature Douglas-Fir site observed between Mar 2006 and Nov 2011 as a function of the light regime: Probability density functions of the a) above-canopy (38.4 m agl) and b) sub-canopy (4 m agl) horizontal wind speeds, and c) above-canopy wind directions. Subplot d) shows the wind direction difference between above- and sub-canopy flow, $\Delta\phi$, i.e., the dimensional wind directional shear, as a function of the above-canopy wind direction.

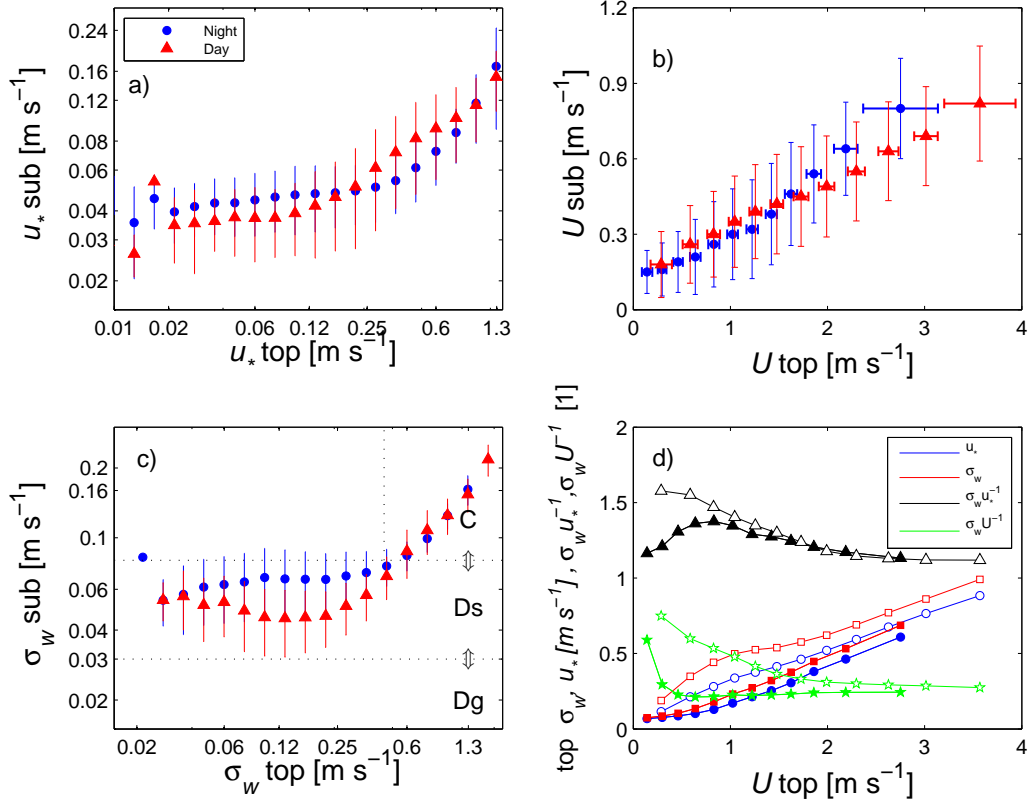


Figure 3: Comparison of various mixing indicators between the above-canopy (*top*) and sub-canopy (*sub*) levels: Friction velocity u_* , horizontal wind speed U , standard deviation of the vertical velocity variance σ_w , integral turbulence characteristic of the vertical velocity $\sigma_w u_*^{-1}$, and integral turbulence intensity $\sigma_w U^{-1}$. In subplot d), symbols are for u_* (circles), σ_w (squares), $\sigma_w u_*^{-1}$ (triangles), and $\sigma_w U^{-1}$ (pentagrams); filled symbols stand for night, and open symbols for day. Data were binned, bars represent one standard deviation of data within each bin. Subplot c) also contains the thresholds (dotted lines) used to define the three proposed exchange regimes: fully coupled canopy (C), decoupled sub-canopy (Ds), and decoupled ground layer (Dg), see Section 2 for details.

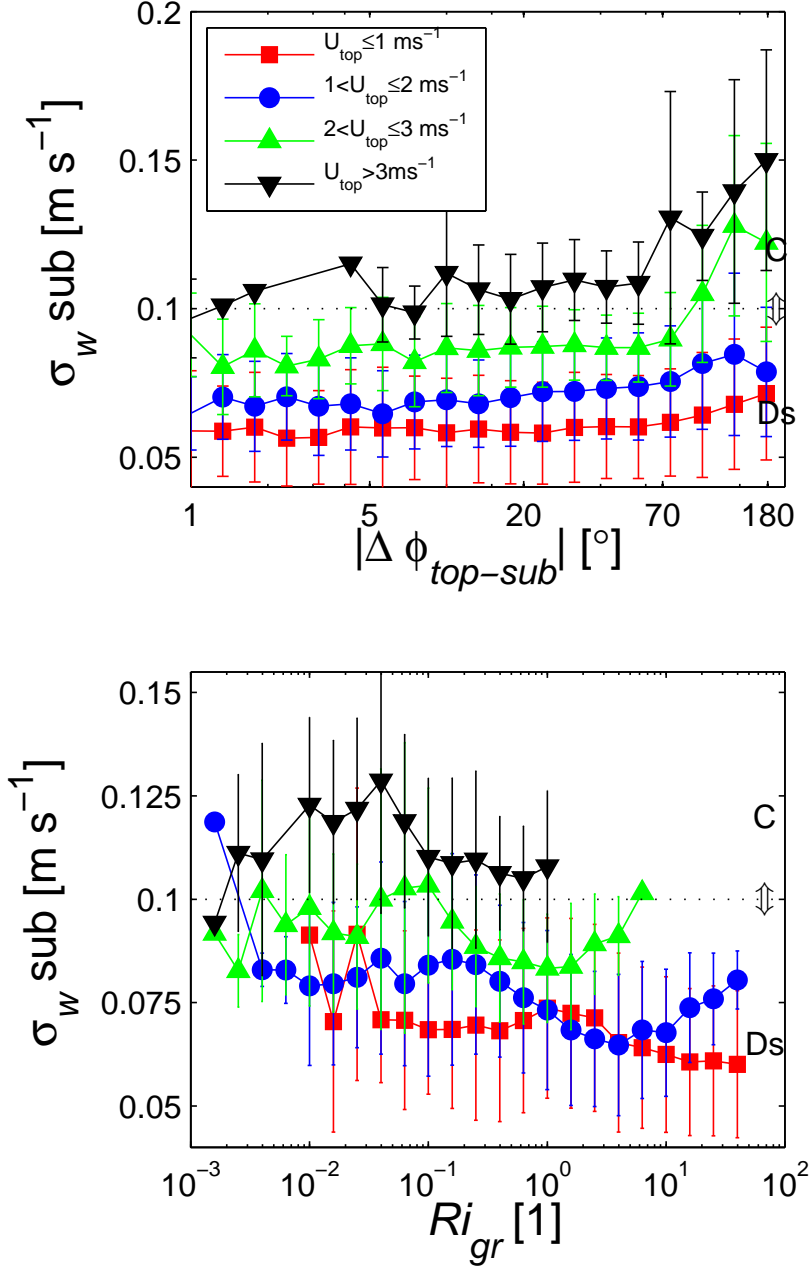


Figure 4: Enhancement of turbulent mixing through wind directional shear for different above-canopy horizontal wind speed classes (U_{top}): a) standard deviation of the vertical velocity variance σ_w as a function of the wind direction difference between above- and sub-canopy flows $|\Delta\phi_{top-sub}|$, and b) σ_w as a function of the canopy gradient Richardson number Ri_{gr} . The threshold of $\sigma_{w,sub} = 0.10$ m s⁻¹ used to delineate the exchange regimes decoupled sub-canopy (Ds) and fully coupled canopy (C) is also marked (dotted line).

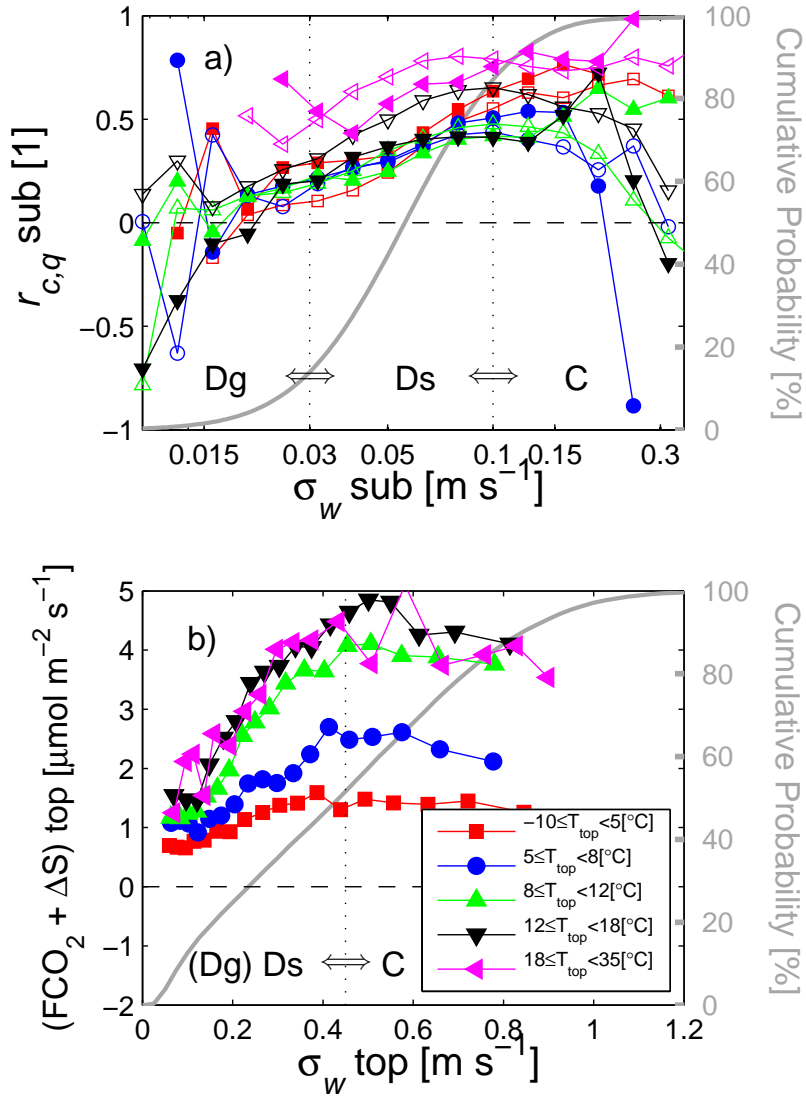


Figure 5: Determination of thresholds for the three exchange regimes: Fully coupled canopy (C), decoupled sub-canopy (Ds), decoupled ground layer (Dg); a) Correlation coefficient between perturbations of carbon dioxide and water vapor mixing ratios $r_{c,q}$ as a function of the turbulent mixing strength in the sub-canopy (sub) σ_w , and b) Nighttime net ecosystem exchange (NEE) as a function of the turbulent mixing strength above the canopy (top) σ_w for different air temperature classes (T_{top}). Each temperature class contains approximately 8000 30-min values. Also shown are the cumulative probability densities (grey lines) of σ_w combined for both night- and daytime data. Open symbols depict daytime, while closed symbols are for nighttime data. Error bars were omitted for clarity of presentation, but scatter was significant.

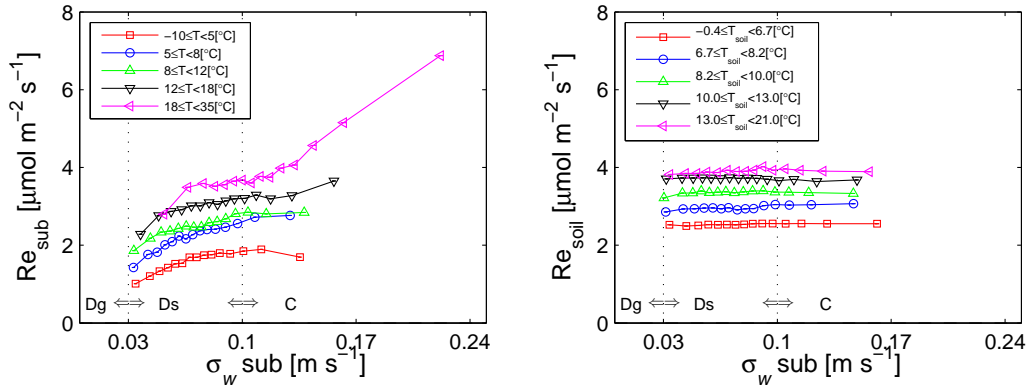


Figure 6: Comparing a) sub-canopy respiration Re_{sub} determined from the subcanopy EC data following Thomas et al. (2008), and b) chamber-based Soil CO₂ efflux Re_{soil} against the subcanopy turbulence mixing strength $\sigma_{w,sub}$ for classes of sub-canopy air T and soil temperature T_{soil} measured near the surface. Vertical dotted lines indicate the thresholds used to delineate the exchange regimes: fully coupled canopy (C), decoupled sub-canopy (Ds), and decoupled ground layer (Dg). Error bars were omitted for clarity of presentation, but scatter was significant.

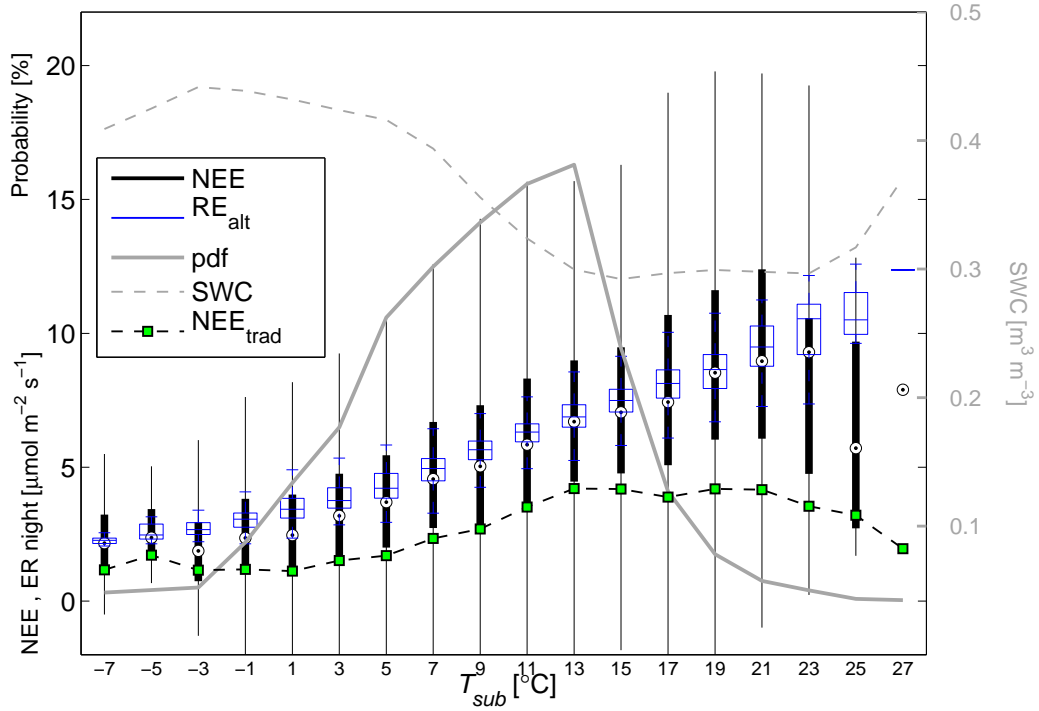


Figure 7: Temperature sensitivity of the improved estimates of nighttime net ecosystem exchange (NEE, for equations see Table 1), alternative ecosystem respiration (RE_{alt}) computed from scaled soil CO_2 efflux chambers and modeled stem and foliage respiration (see Section 3.3), and the 'traditional' single-level net ecosystem exchange using a critical u_* threshold ($NEE_{trad} = FCO_{2,top} + \Delta S$, see Section 4.5) for the observational period 2006 through 2011. NEE and RE_{alt} are presented in boxplots, in which the box is bounded by the 25 % and 75 % percentiles and the median is represented by a circle or horizontal line, respectively; length of the whiskers corresponds to 2.7σ or 99 % of the data in each bin. Also shown is the probability density function (pdf) of data across the sub-canopy air temperature bins T_{sub} of 2 K width, and the volumetric soil water content SWC (y-axis on the right).