

High-Resolution Infrared Studies of Perdeutero-Spiropentane, C₅D₈

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Abstract

Perdeutero-spiropentane (C_5D_8) has been synthesized and infrared and Raman spectra are reported for the first time. Wavenumber assignments are made for most of the fundamental vibrational modes. Gas phase infrared spectra were recorded at a resolution (0.002 cm^{-1}) sufficient to resolve individual rovibrational lines and show evidence of strong Coriolis and/or Fermi resonance interactions for most bands. However a detailed rovibrational analysis of the fundamental ν_{15} (b_2) parallel band proved possible and a fit of more than 1600 lines yielded a band origin of $1053.84465(10)\text{ cm}^{-1}$ and ground state constants (in units of cm^{-1}): $B_0 = 0.1120700(9)$, $D_J = 1.51(3) \times 10^{-8}$, $D_{JK} = 3.42(15) \times 10^{-8}$. We note that the B_0 value is significantly less than a value of $B_a = 0.1140\text{ cm}^{-1}$ calculated using structural parameters from an earlier electron diffraction (ED) study, whereas one expects B_a to be lower than B_0 because of thermal averaging over higher vibrational levels. A similar discrepancy was noted in an earlier study of C_5H_8 [1]. The structural and spectroscopic results are in good accord with values computed at the anharmonic level using the B3LYP density functional method with a cc-pVTZ basis set.

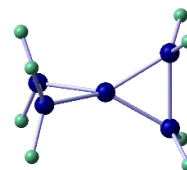
Introduction

In recent papers we have reported the results of high-resolution infrared studies of several interesting carbon ring compounds of high symmetry, including spiro[3.3]heptane [1,2], propellane [3-5] and bicyclo[2.2.1]heptane [6,7]. In each case, the molecules are of considerable interest to structural chemists because of their high ring strain and, in the case of propellane, because of the very unusual axial bond. All three molecules are symmetric tops and relatively small but none had previously been examined at a spectroscopic resolution sufficient to discern individual rovibrational transitions. This may be because none of these compounds is available from commercial sources, although synthesis methods for each have been published [8-10]. Using these methods we have prepared small amounts of each compound with sufficient purity to yield excellent high-resolution infrared spectra.

Part of the motivation for studying such molecules comes from an interest in the use of quantum calculations as an aid in the analysis of the complex vibration-rotation spectral patterns one often observes in high-resolution spectroscopy. Such complexity stems from changes in rotational constants with vibrational states and from Fermi resonance and Coriolis couplings between states. For small molecules, useful estimates of the rovibrational parameters accounting for these effects can be obtained from anharmonic potential energy surfaces calculated by quantum programs such as Gaussian [11]. Such parameters enable spectral simulations that are quite helpful in making initial quantum number assignments of rovibrational transitions in complex spectra. Moreover, once accurate molecular parameters are determined from the spectral analysis, their values serve as added tests of the method and basis sets chosen in the quantum calculations. In this manner, for example, it has been determined that the Gaussian

B3LYP density functional (DFT) method applied to a cc-pVTZ basis set gives quite good rovibrational parameters at a reasonable investment of computer time. In addition, the calculations give helpful information about relative infrared and Raman band intensities that can further aid in vibrational assignments.

Given in our earlier papers [1, 2] is a review of the synthesis and previous studies of spiro[2.2]pentane (spiropentane, C_5H_8):



The present work involved a study of the perdeutero form of spiropentane (C_5D_8), synthesized after the method reported by House et al. [12]. Infrared and Raman spectra were obtained and, aided by the DFT calculations, a confident assignment is made of most vibrational fundamentals. High resolution infrared spectra obtained for gaseous C_5D_8 were analyzed to give accurate rovibrational constants for the ground state and the upper state of the ν_{15} transition. Good agreement is found with corresponding parameters obtained from DFT calculations but we note that the experimental B_0 rotational constant for C_5D_8 is inconsistent with a value calculated from the reported electron diffraction structure [13]. A similar inconsistency was seen in our earlier study of C_5H_8 [1] and in fact served as partial motivation for the work presented here.

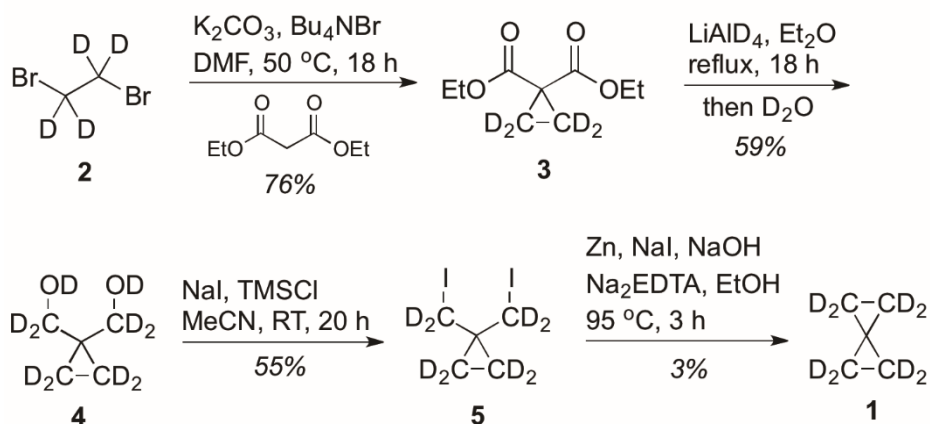
Experimental Details and Results

Synthesis

The perdeutero-spiropentane (**1**) was prepared using the method of House et al. [12], with some notable exceptions detailed in Appendix A. The compound was synthesized via four consecutive reactions starting with perdeutero-dibromo ethane (**2**), as shown in Figure 1. Alkylation of diethyl malonate with **2** gave cyclopropyl malonate **3**. Reduction with lithium aluminum deuteride effectively installed the remaining deuterium

in diol **4**. Activation as the diiodide gave **5**. The yields were comparable to those of House et al., except for the last reaction, for which we utilized Zn metal for the halide extraction and ring-closure step. Our yield for this reduction step was disappointingly low (<3%), much less than a value of 20% reported by House et al. using sodium metal in place of zinc. This poor yield was partially due to the small scale of our reaction amounts, which amplified the losses in the various workup and transfer steps done to purify the final sample after the compound was formed. However the final sample was quite pure, with only about 10% C₅D₇H isotopic impurity, and the amount of C₅D₈ proved sufficient for the high resolution measurements reported here.

Fig. 1. Reaction sequence for the synthesis of perdeutero-spiropentane (C₅D₈) (**1**).



Spectroscopy

In their synthesis paper for C₅D₈, House et al. indicated that infrared and Raman spectra would be reported in a subsequent paper, however we have been unable to find any such published report. Accordingly, we show in Figure 2 our infrared spectra for gaseous C₅D₈, along with Raman spectra for the liquid phase. The latter were recorded at about 6 cm⁻¹ resolution using a Thermo DXR SmartRaman spectrometer with a 150 mW 780 nm excitation source.

The high-resolution infrared spectra were obtained with the Bruker IFS 125HR Fourier transform spectrometer located at the Pacific Northwest National Laboratory.[†] In recording the FTIR spectra, a Globar light source, KBr beamsplitter, and a MCT detector were used. A 20 cm long gas cell with KBr windows and temperature-stabilized to 25.0 °C was filled to about 750 Pa (5.68 torr) of sample and 4160 interferograms were coadded. Instrument resolution, which is defined as $0.9/(\text{maximum optical path difference})$, was set to 0.002 cm^{-1} . The averaged interferograms were transformed using a Cooley-Tukey FFT algorithm with four times zero filling, Mertz phase correction, phase resolution set to 1.0 cm^{-1} , and Boxcar apodization. Spectra were saved from 500 to 4000 cm^{-1} . The transition center wavenumber locations were determined by using a second derivative peak-picking algorithm with a nine point Savitsky-Golay smoothing function. Wavenumber calibration of the various regions was accomplished with an OCS gas reference sample, with wavenumber standards taken from NIST reference tables [14]. The wavenumber corrections were about 0.0002 cm^{-1} and the absolute wavenumber transition center uncertainty is estimated at $\pm 0.00015\text{ cm}^{-1}$.

[†] Certain commercial equipment, instruments, and materials are identified in the paper to adequately specify the experimental procedure. Such identification does not imply recommendations or endorsements by the Pacific Northwest National Laboratory, nor does it imply that the materials or equipment identified are necessarily the best available for the purpose.

Table 1

Experimental and calculated fundamental frequencies (cm^{-1}) for spiropentane (C_5D_8)

Mode		Exp. ^a		Theory (B3LYP/cc-pVTZ)			Raman int. ^b
		IR	Raman	anharmonic	harmonic	IR int. ^b	
a_1	ν_1	...	2206	2202	2265	...	209.0
	ν_2	...	1163	1164	1191	...	12.9
	ν_3	...	918	925	947	...	20.2
	ν_4	...	798	807	824	...	9.5
	ν_5	...	530	520	533	...	11.2
a_2	ν_6	2294	2374
	ν_7	892	908
	ν_8	609	616
b_1	ν_9	...	2312	2300	2378	...	94.6
	ν_{10}	...		939	954	...	4.3
	ν_{11}	...		733	749	...	0.5
	ν_{12}	...		236	237	...	0.6
b_2	ν_{13}	2190		2189	2269	17.7	4.8
	ν_{14}			1485	1529	0.1	0.1
	ν_{15}	1054		1051	1079	1.3	0.6
	ν_{16}	847	840	845	870	4.6	10.7
	ν_{17}	724		725	735	16.4	1.3
e	ν_{18}	2321	2327	2300	2380	26.7	60.0
	ν_{19}	2190	2172	2174	2250	23.8	38.1
	ν_{20}	1081		1074	1101	10.1	0.2
	ν_{21}	1032	1026	1029	1050	1.6	10.4
	ν_{22}		861	874	891	0.3	5.4
	ν_{23}		719	719	731	6.8	15.3
	ν_{24}		554	560	563	0.5	2.7
	ν_{25}			264	263	0.3	0.2

^aInfrared values are for the gas phase, Raman values for the liquid phase.^bInfrared and Raman intensities are harmonic values.

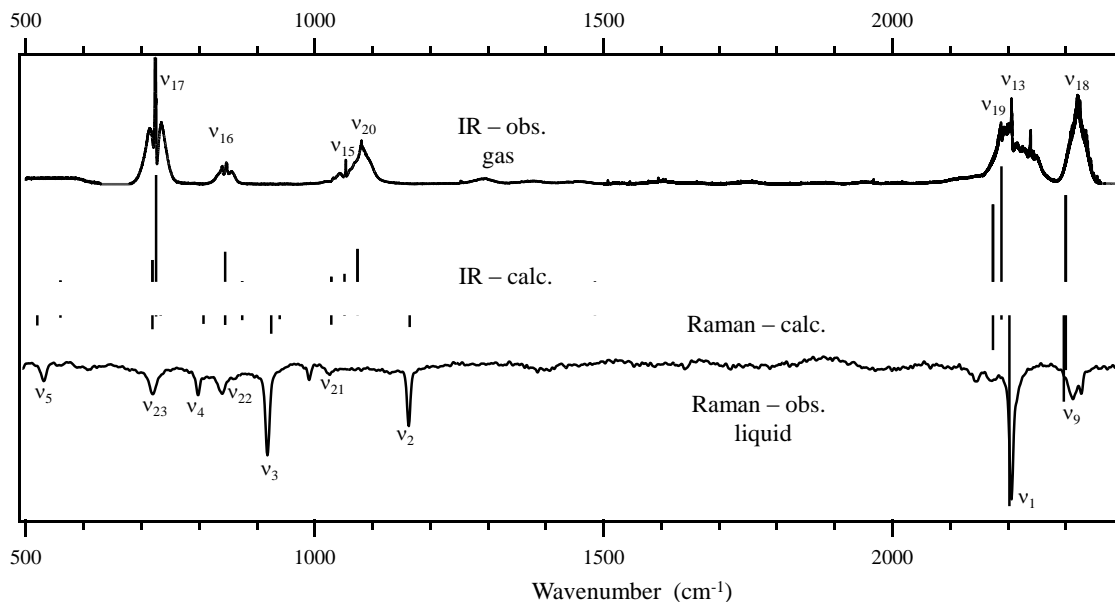


Fig. 2. Infrared and Raman spectra of perdeutero-spiropentane (C_5D_8). For display purposes, intense features due to CO_2 and H_2O were deleted from the spectra. Assignments of most of the fundamental modes are indicated in the spectrum. Vertical sticks indicate the positions calculated at the anharmonic level using Gaussian (B3LYP/cc-pVTZ), with stick heights representing relative intensities computed at the harmonic level.

Vibrational Assignments

Spiropentane has the symmetry of the D_{2d} point group and the representation of the $3N-6$ normal modes of vibration is

$$\Gamma_{\text{vib}} = 5a_1(\text{R}) + 3a_2 + 4b_1(\text{R}) + 5b_2(\text{IR, R}) + 8e(\text{IR, R}).$$

The infrared and Raman activities for the fundamental transitions are shown in parentheses. Table 1 lists fundamental frequencies assigned in this work, along with the harmonic and anharmonic values calculated by us using Gaussian 09 (B3LYP method with a cc-pVTZ basis). The harmonic predictions of relative infrared and Raman intensities are also shown for all fundamentals and are indicated as stick heights in Fig. 2.

The five a_1 modes were prominent in the Raman spectrum and match closely the calculated anharmonic values. The a_2 modes are unassigned since they are inactive in both

infrared absorption and Raman scattering processes. Of the four Raman active b_1 modes, only ν_9 is assigned, a CD₂ antisymmetric stretch at 2312 cm⁻¹. The five b_2 fundamental bands are parallel transitions with sharp Q -branches that are relatively obvious in the infrared spectrum for four modes. A detailed analysis of the ν_{15} b_2 mode, described below, gave a very accurate value of the rovibrational parameters for this fundamental. The fifth b_2 fundamental, a CD₂ bending mode, ν_{14} , is predicted to have very low infrared and Raman intensity and could not be detected.

The band centers of the eight perpendicular e bands are less well-defined but tentative assignments are made for seven of these, based on comparisons with the frequencies and intensities predicted by the Gaussian calculations. The normal modes obtained in these calculations predict the symmetric CD₂ stretching modes ($\nu_1, \nu_{13}, \nu_{19}$) to be about 2200 cm⁻¹ while the antisymmetric stretches (ν_6, ν_9, ν_{18}) are higher, near 2300 cm⁻¹. Because of band overlaps and Coriolis couplings between b_2 and e states, most gas phase band centers have an uncertainty of about 2 to 5 cm⁻¹ and a similar uncertainty applies to the Raman values, which are for the liquid phase. Overall, the agreement with frequencies and intensities predicted by the DFT calculations is considered quite good.

Rovibrational Band Analyses

Rovibrational level expressions

Spiropentane is a prolate symmetric top molecule with principal moments of inertia $I_A < I_B = I_C$. The rotational constants are then given as $A > B$ in which $A = h/(8\pi^2cI_A)$ and likewise for B , where h is Planck's constant and c is the speed of light in vacuum. The physical constants are chosen such that the rotational constants are expressed in units of cm⁻¹.

The energy (in units of cm^{-1}) of a given vibrational state, ν , is given by the term value expression

$$E_\nu = G(\nu, l) + F_\nu(J, K, l) \quad (1)$$

where $G(\nu, l)$ is the vibrational term, and $F_\nu(J, K, l)$ is the rotational term, which may be divided into two parts

$$F_\nu(J, K, l) = F_\nu(J, K) + F_{\nu l}(J, k, l). \quad (2)$$

The contracted subscript ν is used to represent both quantum number and mode number of a vibrational state. For the ground state $\nu_1 = \nu_2 = \nu_3 = \dots = 0$. The zero energy is defined as the $J = K = 0$ level of the ground state *i.e.*, $\nu_0 = G(\nu, l) - G(0, 0)$. For this and other states with vibrational angular momentum quantum number $l = 0$, the second term of Equation 2 is dropped and $F_\nu(J, K)$ is given by

$$\begin{aligned} F_\nu(J, K) = & B_\nu J(J+1) + (A_\nu - B_\nu)K^2 - D_{\nu J} J^2(J+1)^2 - D_{\nu JK} J(J+1)K^2 - D_{\nu K} K^4 \\ & + H_{\nu J} J^3(J+1)^3 + H_{\nu JK} J^2(J+1)^2 K^2 + H_{\nu KJ} J(J+1)K^4 + H_{\nu K} K^6 \\ & \pm \delta_{2K} \Delta_{\nu 2} [J(J+1)][J(J+1)-2]. \end{aligned} \quad (3)$$

In the above expressions J is the total angular momentum quantum number (exclusive of nuclear spin), K is the quantum number of the projection of the vector J onto the principal C_2 symmetry axis, k is that of the signed projection of J onto the symmetry axis, so that $K = |k|$. The D and H terms are quartic and sextic corrections for centrifugal distortion and the last $\Delta_{\nu 2}$ term in Eq. (3) accounts for the centrifugal rotational K -type splitting of the $k = 2$ levels into a B_1, B_2 pair for the A_1 ground vibrational state or, for a B_2 vibrational state, into a A_1, A_2 pair. The pattern of levels predicted in the case of a parallel band such as ν_{15} is shown in Figure 3. The nuclear spin weights g_{ns} indicated in this figure for the ground levels are 1701 for K -even levels of vibronic symmetry A_1 or B_1 , and 1620 for K -

even levels of vibronic symmetry A_2 or B_2 . For D_{2d} molecules, the allowed infrared transitions between rovibrational levels are $A_1 \leftrightarrow B_1$, $A_2 \leftrightarrow B_2$, and $E \leftrightarrow E$ [15].

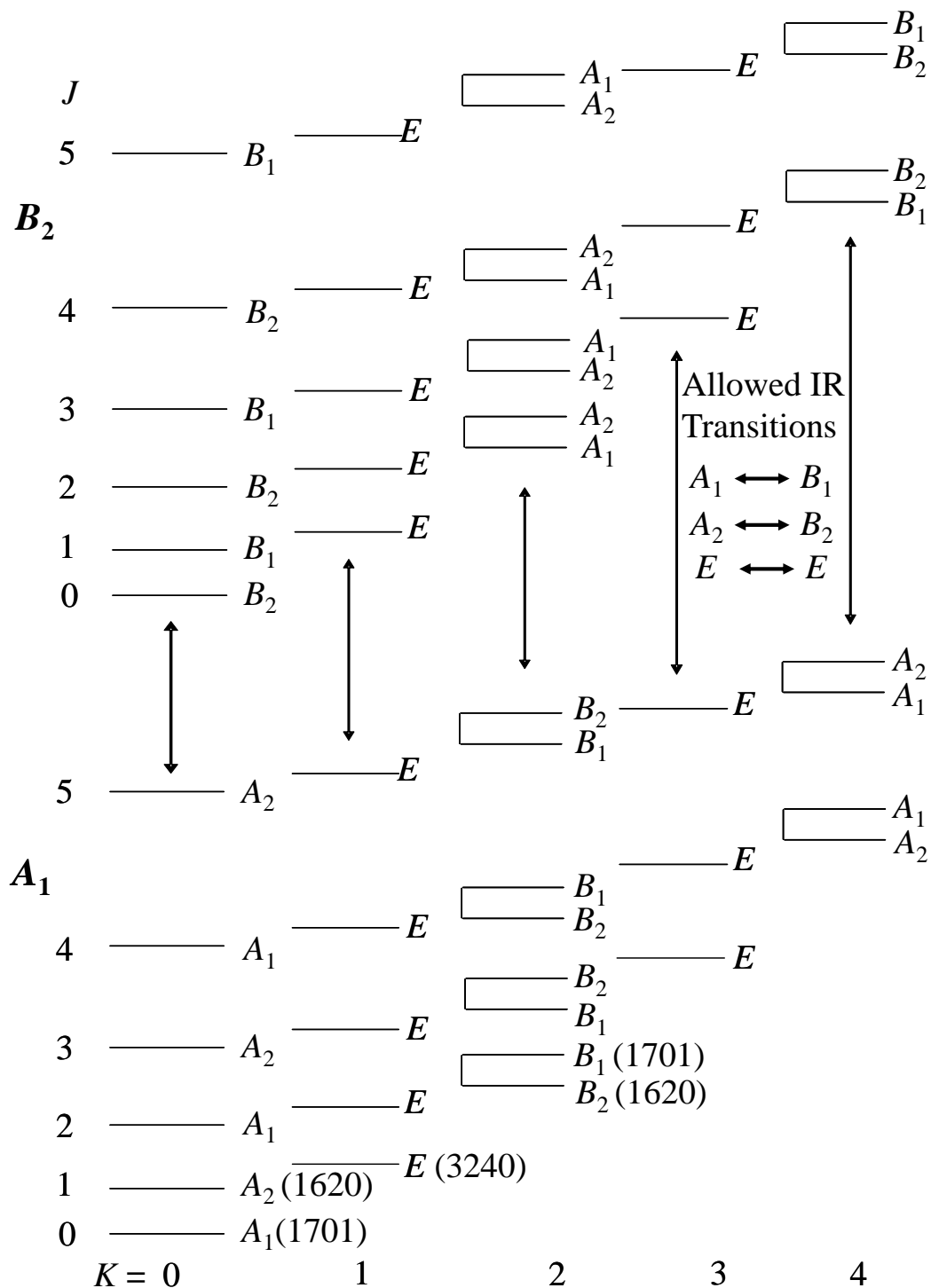


Figure 3. Diagram for the allowed transitions between A_1 and B_2 vibrational states for the ν_{15} band of perdeutero-spiropentane, C_5D_8 .

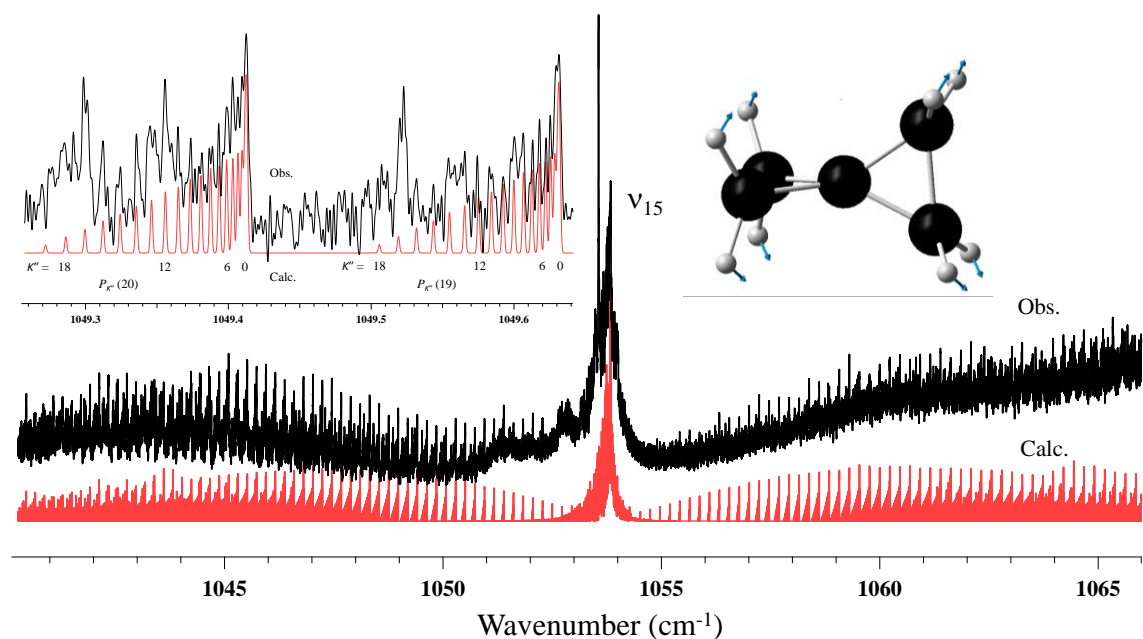


Figure 4. ν_{15} (b_2) parallel band of perdeutero-spiropentane. The inset to the left shows the K'' -structure of the $J'' = 20$ and 19 P -lines. Some of the extra features, as well as the sharp Q -feature on the left side of the Q -branch are attributed to hot bands, such as $(\nu_{25} + \nu_{15}) - \nu_{25}$. The rising baseline in the R -branch is due to overlap with a nearby ν_{20} (e) perpendicular band at 1081 cm^{-1} .

The ν_{15} 1054 cm^{-1} parallel band

The ν_{15} band, seen in Figure 4, shows the characteristic P - Q - R shape of a parallel band and the right hand edge of the Q -branch defines the band origin, making the J assignment of the P - and R -branch lines relatively straight-forward. The left inset of Figure 4 shows the individual resolved K -transitions of the $J'' = 19$ and 20 K -stacks of the P -branch recorded in the present work. The nuclear spin statistical weights g_{ns} in the ground vibrational state are shown in Figure 3 and, for the special case $K'' = 0$, the J -even lines have a weight of 1701 while the J -odd lines are weaker, with weight 1620. In contrast to the case for C_5H_8 , where the even/odd weights are $136/120 = 1.13$, for the deuterated form this even/odd ratio is 1.05 and the effect is too small to be discerned in the intensity variation with J of the unresolved $K'' = 0, 1$ features of the P - and R -branch of Figure 4. Similarly, no evidence is seen for any splitting of the $K'' = 2$ and higher K -

even levels of C_5D_8 , even though a small B_1, B_2 splitting was detected for the $K'' = 2$ transitions of C_5H_8 [1,2].

In the left inset of Figure 4, it is clear that there are extra features in the spectrum and these are attributed to hot bands that originate from the low frequency ν_{12} (b_1) and ν_{25} (e) levels, calculated to occur at 236 cm^{-1} and 264 cm^{-1} , respectively. The anharmonic frequencies calculated for these hot band transitions predict Q -branches shifted from ν_{15} by $+0.07$ and -0.27 cm^{-1} , respectively. The latter shift is in reasonable accord with a measured shift of -0.33 cm^{-1} for the sharp feature seen in Fig. 2 on the left side of the fundamental Q -branch; hence it is assigned as $(\nu_{25} + \nu_{15}) - \nu_{25}$. A weaker, broader feature also occurs $+0.10\text{ cm}^{-1}$ above the ν_{15} fundamental and this is assigned as $(\nu_{12} + \nu_{15}) - \nu_{12}$.

In addition to added features due to contributions from hot bands, the rovibrational spectrum of C_5D_8 is generally condensed compared to that of C_5H_8 , due to the smaller rotational constants of the former. Moreover, the density of nearby vibrational levels is greater for C_5D_8 , resulting in a higher probability of Fermi and Coriolis interactions with the upper states of the fundamental transitions. The latter made difficult, for example, the analysis of the apparent P - Q - R structure of the ν_{17} (b_2) perpendicular band predicted near 725 cm^{-1} , since this mode can Coriolis couple to the nearby ν_{11} (b_1) and ν_{23} (e) modes predicted at 733 and 719 cm^{-1} , respectively. Evidence for Coriolis perturbation of line positions and intensities were also seen for the intense, overlapping ν_{16} (b_2) and ν_{22} (e) fundamentals, predicted at 845 and 874 cm^{-1} , respectively. A further complication here is a Fermi resonance interaction with the $\nu_8(a_2)+\nu_{12}(b_1)$ combination state of b_2 symmetry which is predicted at 845 cm^{-1} .

Surprisingly, the ν_{15} (b_2) parallel band at 1051 cm^{-1} proved to be relatively unperturbed, despite the presence of nearby ν_{21} and ν_{22} (e) fundamentals predicted at

higher (1074 cm^{-1}) and lower (1029 cm^{-1}) wavenumber values, respectively. We note that the “push” on the b_2 rovibrational levels by these two e states is in opposing directions so this compensating effect may account for fact that over 1600 transitions could be confidently assigned and fitted for ν_{15} without including any Coriolis interactions. This set did not include any of the Q -transitions nor the $K'' = 0$ or 1 lines, since none of these were resolved. Despite this good fit, to avoid any possible errors due to unaccounted mixing of the upper levels of the ν_{15} transitions, ground state constants were deduced from rotational level separations obtained from $R(J''-1, K'') - P(J''+1, K'')$ differences in the ν_{15} transitions. 458 such differences were obtained and fitted, yielding the three ground state parameters B_0 , D_J , and D_{JK} shown in Table 2. We note that these ground state parameters are the same, within the listed uncertainties, as those obtained when these parameters were varied in fitting the full set of the ν_{15} transitions. The resultant rovibrational parameters for the upper state of the ν_{15} fundamental band are presented in Table 3 and all transitions used in both fits are provided as supplementary data for this article.

Table 2Rotational parameters (in cm^{-1}) for the ground state of normal and perdeutero-spiropentane

Parameter ^a	C ₅ H ₈ [Ref. 1]		C ₅ D ₈	
	Experimental	Theoretical ^b	Experimental	Theoretical ^b
A_0	[0.35059] ^c	0.35059	[0.24891]	0.24891
B_0	0.1394741(2)	0.13877	0.1120700(9)	0.11145
$D_J \times 10^8$	2.462(2)	2.42	1.505 (27)	1.35
$D_{JK} \times 10^8$	8.66(4)	8.59	3.42(15)	3.51
$D_K \times 10^8$	[5.21]	5.21	[2.09]	2.09
$H_J \times 10^{13}$	[0.0]	0.053	[0.0]	0.029
$H_{JK} \times 10^{13}$	[0.0]	-5.45	[0.0]	-2.45
$H_{KJ} \times 10^{13}$	[0.0]	8.99	[0.0]	4.29
$H_K \times 10^{13}$	[0.0]	-3.90	[0.0]	-1.82
$\Delta_2 \times 10^9$	1.172(16)	1.22	[0.0]	0.71
$J_{max} =$	96		52	
$K_{max} =$	31		37	
# differences	5612		458	
rms dev. (cm^{-1})	0.00037		0.00052	

^aGiven in parentheses here and in subsequent tables are uncertainties in the last digits (twice the standard deviation = type B, k=2, uncertainties as defined in Ref. [16].)

^bB3LYP/cc-pVTZ calculation using Gaussian 09 with Anharm/Vibrot option.

^cValues given in square brackets were held fixed.

Table 3

 ν_{15} rovibrational parameters (cm^{-1}) for perdeutero-spiropentane

Parameter ^a	Experimental	Theoretical ^b
ω_0		1078.81
ν_0	1053.84465(10)	1051.43
$\Delta A \times 10^3$	-0.2533(6)	-0.234
$\Delta B \times 10^3$	0.1336(3)	0.165
$\Delta D_J \times 10^8$	-0.19(2)	
$\Delta D_{JK} \times 10^8$	5.17(6)	
$\Delta D_K \times 10^8$	-4.13(6)	
$\Delta H_J \times 10^{12}$	-0.65(4)	
$\Delta H_{JK} \times 10^{12}$	10.53(18)	
$\Delta H_{KJ} \times 10^{12}$	-12.26(24)	
$\Delta H_K \times 10^{12}$	[0]	
$J_{max} =$	60	
$K_{max} =$	49	
# transitions	1633	
rms dev. (cm^{-1})	0.00050	

^a $\Delta B = B' - B''$, $\Delta C = C' - C''$, etc.^b B3LYP/cc-pVTZ calculation using Gaussian 09 with Anharm/Vibrot option

Discussion

Ground state rotational constants and molecular structure

Table 2 offers comparison of the experimental and theoretical ground state rovibrational constants for normal and perdeutero forms of spiropentane. For both isotopomers, level differences in the ground state were fitted to the symmetric top relation in Equation 3 and it was found that only the constants B_0 , D_J , and D_{JK} were needed to obtain excellent fits. No improvement was achieved by including H sextic distortion constants so these were fixed at zero values. Of course, the K -dependent A_0 , D_K , and H_K parameters cannot normally be determined from the spectra of symmetric tops but estimates of these parameters can be computed with the Gaussian 09 program [11] and these are also listed in the table. All parameters observed or calculated for C_5D_8 are significantly smaller than their C_5H_8 counterparts. In the case of the $K'' = 2$ centrifugal splitting constant Δ_2'' , the reduction is 42% and no splitting in any of the $K'' = 2$ transitions was resolved for C_5D_8 .

The Gaussian calculations of the quartic distortion constants do not contain any anharmonic contribution and the D_J and D_{JK} values correspond to values for the equilibrium structure. The anharmonic corrections are expected to be small however, and the experimental ground state D_J and D_{JK} values agree with the theoretical ones in sign and in magnitude to within about 2%, except for D_J of C_5D_8 , where the difference is 12%. The uncertainties in the C_5D_8 parameters are much larger than for those for C_5H_8 , which is not surprising since the latter data set was much larger and included differences from several fundamental bands, with a much greater range of J values.

The B_0 values from the C_5H_8 and C_5D_8 analyses are very well-determined, with uncertainties in the seventh decimal place. For both isotopomers, the experimental values

0.1394741(1) and 0.1120700(9) cm^{-1} , respectively, fall between the predicted theoretical values of B_e and B_0 and agree with the latter within 0.6%. We note that the experimental value of B_0 for C_5H_8 is significantly smaller (1.68%) than a value of $B_a=0.1418 \text{ cm}^{-1}$ calculated from the electron diffraction (ED) thermal average (r_a) structure (given in Table 2 of Ref. 1). Similarly, for C_5D_8 , the experimental value 0.1120700(9) cm^{-1} is comparably smaller (1.70%) than a value of 0.1140 cm^{-1} calculated from the electron diffraction structure. Since the latter includes contributions from all thermally populated states, it would be expected that the B_a value deduced from the ED structure would actually be *lower*, not higher, than B_0 . Several authors [17-19] have raised questions about the electron diffraction results, particularly about the CH_2 angle, 118.4(9), which is about 4 degrees larger than the Gaussian prediction of 114.4 degrees, and is also larger than a value of 115.0(7) degrees deduced from NMR measurements [18]. As discussed in our earlier paper [1], we suspect a scaling error in the electron diffraction measurements, but this remains to be confirmed.

The ν_{15} vibration

Table 3 lists the rovibrational parameters for the excited state of ν_{15} and it can be seen that the frequency of the band center, 1053.84465(10) cm^{-1} is in good accord (0.2 %) with the theoretical anharmonic frequency (1051.4 cm^{-1}). The calculated harmonic frequency (1078.8 cm^{-1}) is slightly higher, indicating that the anharmonicity of this vibrational mode is modest (2.4%). According to the Gaussian harmonic normal mode calculations, the ν_{15} mode is a mixture of ring stretching and CH_2 bending motions, with the latter emphasized in the inset of Figure 3, which is a snapshot of the animation of the mode given by the GaussView program.

The Anharmon/Vibrot options of the Gaussian program also permit calculation of the anharmonic contribution to the vibration-rotation constants for each mode i , $\alpha_{Ai} = A_0'' - A_i'$ and $\alpha_{Bi} = -\Delta B_i$. As seen in Table 3, the experimentally-determined ΔA and ΔB constants for spiropentane agree in sign with the theoretical values but differ in magnitude by 8% and -19%, respectively. The ΔB % difference is somewhat larger than typical values of less than 10% that we have seen for propellane and bicyclopentane and other cases where Fermi and/or Coriolis mixing with nearby states is absent. It is likely that unaccounted x,y Coriolis mixing of ν_{15} with the nearby ν_{21} and ν_{22} (e) states is responsible, since this can change the ΔB value (but not the ΔA value).

Summary

This work presents the first infrared and Raman spectra of perdeuterio-spiropentane and analysis of the high resolution infrared spectrum of gaseous C_5D_8 yielded accurate rovibrational parameters for the ground and excited states of the ν_{15} fundamental mode. The B_0 value determined for the ground state is consistent with our earlier speculation [1] that the structure obtained from electron diffraction measurements is not accurate. Comparison of the experimental and theoretical parameters for the ground state and the upper state levels of the ν_{15} parallel band is generally favorable and illustrates the value of quantum calculations as a complement in the analysis of complex molecular spectra. Evidence was seen of significant l -doubling and Fermi resonance/Coriolis interactions for some of the other bands of C_5D_8 and efforts to analyze these are underway.

Appendix A. Synthesis of Perdeutero-Spiropentane (C₅D₈)

Reaction 2 to 3. To a solution of diethyl malonate (15.3 mL, 100.2 mmol) in DMF (20.0 mL) were added potassium carbonate (34.6 g, 250.5 mmol), 1,2-dibromoethane-d₄ (Alfa Aesar, 99% D), **2** (25.0 g, 130.3 mmol), and tetrabutyl ammonium bromide (1.62 g, 5.01 mmol). The mixture was allowed to stir at room temperature for 12 hours, then was heated to 50°C and stirred for another 6 hours. The mixture was diluted with ethyl acetate and was washed with water and brine and dried over MgSO₄. The resulting mixture was purified by liquid chromatography (silica gel column) and the solvents removed under reduced pressure to yield 15.9 g of deuterated 1,1-dicarbethoxycyclopropane **3** (75.5%).

Reaction 3 to 4. Diester **3** (14.0 g, 74.9 mmol) was dissolved in 50 mL ether and added over 1 hour at 0°C to a solution of LiAlD₄ (Alfa, 98%D) (4.5 g, 98.9 mmol) in 200 mL ether. The reaction was then heated to reflux and allowed to stir for 18 hours. The reaction was quenched by adding sequentially 4.2 mL D₂O, 4.2 mL 15% NaOH solution in D₂O, then 12.6 mL D₂O. The mixture was filtered and dried over MgSO₄. The compound was purified by liquid chromatography and the solvents removed under reduced pressure to yield 4.96 g of deuterated 1,1-bis-(hydroxymethyl)cyclopropane **4** (58.9%).

Reaction 4 to 5. To a solution of **4** (4.90 g, 48.7 mmol) in 95 mL acetonitrile were added NaI (21.15 g, 141.1 mmol) and tetramethylsilyl chloride (17.9 mL, 141.1 mmol) under Ar(g). The mixture was allowed to stir at room temperature for 20 hours and was diluted with equal volume water and DCM. The organic layer was washed with 10% Na₂S₂O₃ (aq) and brine and dried over MgSO₄. The mixture was filtered and the solvents evaporated under reduced pressure. The crude sample was diluted with hexanes and passed through a silica plug to give 7.48 g of pure deuterated 1,1-bis-(iodomethyl)cyclopropane **5** (54.8%).

Reaction 5 to 1. For convenience, this reaction step employed the same apparatus used in synthesizing ordinary spiropentane from pentaerythrityl tetrabromide, as described in reference 1. The system consisted of a 500 mL round bottom flask attached to a water condenser and two successive traps, the first at -78°C, the second at 77 K. The system was flushed with He(g) which was flowed slowly through the system for the entire experiment to carry the product out of the reaction flask. Na₂ EDTA – 2H₂O (28.52 g, 76.6 mmol), NaOH (8.59 g, 215 mmol), NaI (0.75 g, 5.01 mmol), 150 mL H₂O, and 35 mL ethanol were added to the flask and the mixture was stirred while heating to 95°C. Zn(s) (6.50, 99.4 mmol) was then added to the flask and over 1 hour **5** (7.48 g, 22.6 mmol) was added to the stirring solution. The mixture was allowed to stir for 2.5 hours. A small amount of bromine was then added to the spiropentane in the second trap to remove known side products 2-methyl-1-butene and methylene cyclobutane. The overall yield of perdeutero-spiropentane **1** was quite low (<3%) but, from infrared spectra, the isotopic purity was high ~90% and the sample quantity was sufficient for the high resolution spectroscopic measurements. We note that a higher yield (20%) in this last reaction step has been reported by House and Lord [12] using Na metal to remove the iodine atoms and close the second ring in **1**, so their method for this last step is recommended.

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Supplementary Data

Supplementary data for this article are available on ScienceDirect (www.sciencedirect.com) and as part of the Ohio State University Molecular Spectroscopy Archives (http://msa.lib.ohio-state.edu/jmsa_hp.htm).

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