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5	3D Multi-source Model of Elastic Volcanic Ground Deformation
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7	Antonio G. Camacho ^a , José Fernández ^{a,*} , Sergey V. Samsonov ^b ,
8	Kristy F. Tiampo ^c , and Mimmo Palano ^d
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10	^a Institute of Geosciences (CSIC-UCM), C/ Doctor Severo Ochoa, 7, Facultad de Medicina
11	(Edificio entrepabellones 7 y 8, 4 ^a planta), Ciudad Universitaria, 28040, Madrid, Spain. (e-
12	mail: antonio_camacho@mat.ucm.es, jft@mat.ucm.es)
13	^b Canada Centre for Mapping and Earth Observation, Natural Resources Canada, 560
14	Rochester Street, ON K1A 0E4, Ottawa, Canada. (e-mail: sergey.samsonov@canada.ca)
15	^c CIRES and Geological Sciences. 216 UCB, University of Colorado, Boulder, CO 80309,
16	USA. (e-mail: kristy.tiampo@colorado.edu)
17	^d Istituto Nazionale di Geofisica e Vulcanologia, Osservatorio Etneo - Sezione di Catania,
18	Piazza Roma 2, 95125 Catania, Italy (e-mail: mimmo.palano@ingv.it)
19	
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21	*Corresponding author: jft@mat.ucm.es, +34-913944632 (JF)
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26 ABSTRACT: Developments in Interferometric Synthetic Aperture Radar (InSAR) and GNSS 27 (Global Navigation Satellite System) during the past decades have promoted significant 28 advances in geosciences, providing high-resolution ground deformation data with dense 29 spatio-temporal coverage. This large dataset can be exploited to produce accurate assessments 30 of the primary processes occurring in geologically active areas. We present a new, original 31 methodology to carry out a multi-source inversion of ground deformation data to better 32 understand the subsurface causative processes. A nonlinear approach permits the 33 determination of location, size and three-dimensional configuration, without any a priori 34 assumption as to the number, nature or shape of the potential sources. The proposed method 35 identifies a combination of pressure bodies and different types of dislocation sources (dip-36 slip, strike-slip and tensile) that represent magmatic sources and other processes such as 37 earthquakes, landslides or groundwater-induced subsidence through the aggregation of 38 elemental cells. This approach has the following features: (1) simultaneous inversion of the 39 deformation components and/or line-of-sight (LOS) data; (2) simultaneous determination of 40 diverse structures such as pressure bodies or dislocation sources, representing local and 41 regional effects; (3) a fully 3D context; and (4) no initial hypothesis about the number, 42 geometry or types of the causative sources is necessary. This methodology is applied to Mt. 43 Etna (Southern Italy). We analyze the ground deformation field derived from a large InSAR 44 dataset acquired during the January 2009 - June 2013 time period. The application of the 45 inversion approach models several interesting buried structures as well as processes related to the volcano magmatic plumbing system, local subsidence within the Valle del Bove and 46 47 seaward motion of eastern flank of the volcano.

49 1. Introduction.

50 Recent technical developments in geodesy have resulted in significant advances in 51 volcanology (Fernández et al., 2017; and references therein). For example, Global Navigation 52 Satellite System (GNSS) produced sub-centimeter precision in positioning while the 53 development of Advanced Differential Interferometric Synthetic Aperture Radar (A-DInSAR) 54 techniques have resulted in the estimation of 1D to 3D deformation field with dense spatio-55 temporal coverage. Therefore, high resolution, high precision measurements of ground 56 deformation with extensive coverage are available to explore complex models of ground 57 deformation in volcanic areas.

In this context, surface displacements are inverted to infer valuable constraints on the active magmatic sources (e.g., Rymer and Williams-Jones, 2000; Fernández et al., 2001; Dzurisin, 2007; Cannavò et al., 2015). Surface deformation is a direct consequence of the dynamics of volcanic plumbing systems, and reflect the shape of magma intrusions, the volume of intruding/arising magma, and the emplacement mechanisms. Normally, regular geometries (point sources, disks, prolate or oblate spheroids, etc.) are assumed at the initial stages (Lisowski, 2007) and the resulting inversion is carried out in a linear context.

Surface deformation also has been inverted in order to provide insight into the geometry and slip of buried seismic dislocations. The initial geometry of the buried dislocation is generally assumed based on prior information obtained from various sources such as local geology, fault mapping, and earthquake focal mechanisms. Again, the inversion is generally conducted in a linear framework (Segall, 2010; Pascal et al., 2014).

Camacho et al. (2011a) developed an original methodology aimed at the determination of the 3D geometry and the location of the causative bodies by inverting ground deformations and gravity changes due to pressure and/or mass anomalies embedded into an elastic medium.
Such a fully nonlinear inversion has led to interesting results in volcanic environments, where ground deformations are related to over-pressured magmatic bodies (Camacho et al., 2011a;
Samsonov et al, 2014; Cannavò et al, 2015; Camacho et al., 2018; Camacho and Fernández,
2019).

Most volcanically active regions are characterized by complicated patterns of ground deformation resulting from multiple natural (e.g., inflation, deflation, dike intrusion, active faulting, flank instability and landslides) and anthropogenic sources (Fernández et al., 2005, 2017; Tiampo et al., 2013; Samsonov et al., 2014). For example, Mt. Etna volcano is characterized by short-term inflation/deflation episodes related to the magmatic dynamics of its plumbing system, by a near-continuous seaward motion of its eastern flank (Palano, 2016) and by regional tectonic processes (Palano et al., 2012).

An extension of the former successful nonlinear approach, which only estimated elastic deformation due to pressure sources applicable to specific volcanic areas, is required for more general geophysical active regions, where more varied types of deformation sources are present.

We present a new inversion process that extends the previous methodology by including dislocation sources as given by Okada (1985), in order to obtain a more general inversion method that estimates non-subjective models of the observed deformation process within an almost entirely automatic framework.

Here, we describe this new approach, some simulation cases, and its application to actual ground deformation at Mt. Etna estimated from advanced *A-DInSAR* data. A second test case, the interpretation of the co-seismic deformation for the 2014 earthquake in Napa Valley (California) (Polcari et al., 2017), is presented in the Supplementary Material. The results allow us to evaluate the power of the methodology for 3D multi-source modelling of volcanic deformation data.

99

100 2. Inversion Methodology.

101 Camacho et al. (2011a) presented an original methodology for simultaneous inversion of 102 displacement determined using terrestrial and/or space techniques and gravity changes, 103 adapted from a previous methodology for gravity inversion (Camacho et al, 2007 and 2011b). Assuming simple isotropic elastic conditions, the approach determines a general geometrical 104 105 configuration of pressurized and/or density sources corresponding to prescribed values of 106 anomalous density and pressure. These sources are described as an aggregate of pressure and 107 density point sources, and they fit the entire dataset within some regularity conditions. In this 108 methodology, the representation of single sources as the sum of elementary solutions 109 representing 3D irregular geometries, as is typically done for dislocation sources representing 110 faults (Segall, 2010). For pressure sources, this is applied by assuming that the model is linear 111 in the pressure perturbation, with an assumed constant value of pressure change, and the 112 media is assumed to be isotropic, allowing for superposition (Geerstma and Van Opstal, 1973; 113 Brown et al., 2014; Fernández et al., 2018). In a mathematical appropriate way, pressure and 114 mass sources can be combined together (Rundle, 1982; Fernández and Rundle, 1994). The 115 approach works in a step-by-step growth process that constructs very general geometrical 116 configurations (Camacho et al., 2007; 2011a, b).

This approach provided useful results for volcanic areas when deformations come from magmatic sources considered as a combination of pressure and mass variations, if displacement and gravity change data are available; or just pressure sources if only displacement data exist. Nevertheless, for many volcanic regions, observed deformations often are caused by additional phenomena not related to pressurization. These include fault dislocations, sliding and subsidence phenomena that cannot be satisfactorily modelled with the former approach. Therefore, here we propose an improvement of the original inversion methodology which incorporates these new sources, allowing us to obtain a general model of all the observed deformation composed of multiple simultaneous and combined 3D sources.

In this new approach, superposition is still allowable for modeling single sources, as in
Camacho et al. (2011a). For combination of different sources of the same or different nature,
we apply the results of Pascal et al. (2014).

The medium is divided into a 3D partition of elemental cells. The aggregation of elemental sources and the superposition of their contribution forms the geometry of the extended causative bodies. One key aspect is to select some simple expressions for cell contribution in order to fit thousands of data points by the superposition of thousands of cells in a short time, thus allowing the methodology to be used for real time monitoring during unrest (Cannavò et al., 2015; Camacho and Fernández, 2019).

135 **2.1 Elementary sources. Direct formulae.**

We consider a point P(X, Y, Z) located on the surface of a semi-infinite elastic medium where an elemental source is located at (x, y, z). For the surface deformation due to elemental dislocation sources we use the expressions by Okada (1985), and for the elemental pressure sources we use the expressions by Geertsma and Van Opstal (1973).

140 **2.1.1.** Surface deformation due to shear and tensile elemental dislocations.

141 Displacements u_x , u_y , u_z at *P* produced by a buried dislocations point source located at 142 (x,y,z) in an elastic half-space are given by (Okada, 1985):

143 (a) for strike-slip:

144
$$\begin{pmatrix} u_x \\ u_y \\ u_z \end{pmatrix} = -\frac{U_1 \Delta S}{2 \pi} \left[\frac{3 \, dx \, q}{R^5} \begin{pmatrix} dx \\ dy \\ dz \end{pmatrix} + \sin \delta \begin{pmatrix} I_1 \\ I_2 \\ I_4 \end{pmatrix} \right]$$
(1)

(b) for dip-slip: 145

146
$$\begin{pmatrix} u_x \\ u_y \\ u_z \end{pmatrix} = -\frac{U_2 \Delta S}{2 \pi} \left[\frac{3 p q}{R^5} \begin{pmatrix} dx \\ dy \\ dz \end{pmatrix} - \sin \delta \cos \delta \begin{pmatrix} I_3 \\ I_1 \\ I_5 \end{pmatrix} \right]$$
(2)

(c) for tensile: 147

148
$$\begin{pmatrix} u_x \\ u_y \\ u_z \end{pmatrix} = -\frac{U_3 \Delta S}{2 \pi} \left[\frac{3 q^2}{R^5} \begin{pmatrix} dx \\ dy \\ dz \end{pmatrix} - \sin^2 \delta \begin{pmatrix} I_3 \\ I_1 \\ I_5 \end{pmatrix} \right]$$
(3)

149 where:

150	δ : dip angle of the fault plane,
151	α : azimuth angle

152
$$\begin{pmatrix} dx \\ dy \end{pmatrix} = \begin{pmatrix} \sin \alpha & \cos \alpha \\ -\cos \alpha & \sin \alpha \end{pmatrix} \begin{pmatrix} X - x \\ Y - y \end{pmatrix},$$

$$dz = Z - z,$$

$$I_{1} = \frac{\mu}{\lambda + \mu} dy \left[\frac{1}{R(R + dz)^{2}} - dx^{2} \frac{3R + dz}{R^{3}(R + dz)^{3}} \right],$$

$$I_{2} = \frac{\mu}{\lambda + \mu} dx \left[\frac{1}{R(R + dz)^{2}} - dy^{2} \frac{3R + dz}{R^{3}(R + dz)^{3}} \right],$$

$$I_{3} = \frac{\mu}{\lambda + \mu} dx \left[\frac{1}{R^{3}} \right] - I_{2},$$

$$I_{4} = \frac{\mu}{\lambda + \mu} \left[-dx \, dy \frac{2R + dz}{R^{3}(R + dz)^{2}} \right],$$

$$I_{5} = \frac{\mu}{\lambda + \mu} \left[\frac{1}{R(R + dz)} - dx^{2} \frac{2R + dz}{R^{3}(R + dz)^{2}} \right],$$

$$\binom{p}{q} = \left(\frac{\sin \delta}{-\cos \delta} \frac{\cos \delta}{\sin \delta} \right) \binom{dz}{dy},$$

$$\frac{p}{4}$$

$$R^{2} = (X - x)^{2} + (Y - y)^{2} + (Z - z)^{2} = dx^{2} + dy^{2} + dz^{2} = dx^{2} + q^{2} + p^{2}.$$

154 2.1.2. Surface deformation due to a pressure elemental prismatic body.

The simplest method which still provides a good overall estimate of the spatial subsidence distribution for compacting reservoirs of arbitrary 3D shape and change in reservoir pressure is based on the lineal elastic theory of nuclei of strain in the half-space (Geertsma and Van Opstal, 1973). Assuming linearity of the stress-strain relation and isotropy of the material, the displacements u_x , u_y , u_z at a surface point *P* due to a buried small prismatic source with overpressure Δp and sides Δx , Δy , Δz , located at (*x*,*y*,*z*) in an elastic half-space can be determined as:

162
$$\begin{pmatrix} u_{x} \\ u_{y} \\ u_{z} \end{pmatrix} = \Delta p \; \frac{1-\nu}{\mu} \; \frac{3}{4\pi} \int_{x-\Delta x/2}^{x+\Delta x/2} \int_{y-\Delta y/2}^{y+\Delta y/2} \int_{z-\Delta z/2}^{z+\Delta z/2} \binom{X-\xi}{Y-\eta} \frac{d\xi \; d\eta \; d\zeta}{\left((X-\xi)^{2}+(Y-\eta)^{2}+(Z-\zeta)^{2}\right)^{3/2}}$$
(4)

163 where v is the Poissson's ratio and μ is the shear modulus.

Assuming that displacements u_x , u_y , u_z at the surface happen to be almost directly proportional to the thickness Δz of the reservoir, the volume integrations for a parallelepiped cell of sides Δx , Δy , Δz and overpressure Δp in equations (4) can be simplified to integration in the horizontal plane only given rise to (Geertsma and Van Opstal, 1973):

168
$$\begin{pmatrix} u_x \\ u_y \\ u_z \end{pmatrix} = \Delta p \; \frac{1-\nu}{\mu} \frac{3}{4\pi} I \; \Delta z$$
 (5)

169 where:

170
$$I = I_i \left(X - x + \frac{\Delta x}{2}, Y - y + \frac{\Delta y}{2}, Z - z \right) - I_i \left(X - x + \frac{\Delta x}{2}, Y - y - \frac{\Delta y}{2}, Z - z \right) -$$

171
$$I_i\left(X - x - \frac{\Delta x}{2}, Y - y + \frac{\Delta y}{2}, Z - z\right) + I_i(X - x - \frac{\Delta x}{2}, Y - y - \frac{\Delta y}{2}, Z - z),$$

and integrals I_i for displacements along *i*-directions are given by:

173
$$I_{z}(p,q,r) = \frac{1}{2} \frac{p}{|p|} \frac{q}{|q|} \left\{ \arcsin \frac{p^{2}q^{2} - r^{2}(p^{2} + q^{2} + r^{2})}{(p^{2} + r^{2})(q^{2} + r^{2})} + \frac{\pi}{2} \right\},$$

174
$$I_x(p,q,r) = \operatorname{arcsinh} \frac{p}{\sqrt{q^2 + r^2}}$$

175
$$I_y(p,q,r) = \operatorname{arcsinh} \frac{q}{\sqrt{p^2 + r^2}}$$

Equations (1)-(3), (5) provide the surface displacement due to elemental cells for pressure and dislocations. The total effect of a single anomalous structure described as an aggregation of *m* small parallelepiped cells is obtained as an addition (discrete integration) of the partial effects (Geertsma and Van Opstal, 1973; Okada, 1985).

The topography of volcanoes can have an important effect on deformation changes (Supplementary Material). We take this effect into account by incorporating the varyingelevation analytical solution approach (Williams and Wadge, 1998) into the equations and code. This direct formulation is used to carry out the inverse approach and to determine the pressure and dislocation 3D source structures responsible of the observed deformation.

185 2.2. Inversion methodology

The perturbing 3D sources are described as an aggregate of elemental sources that fits the entire dataset within some regularity conditions. The approach works in a step-by-step growth process (Camacho et al., 2007; 2011b) constructing very general geometrical configurations.

189 The observation equations are:

190

$$dr = dr^c + \varepsilon \tag{6}$$

191 where dr, dr^c represent the vector of observed and calculated 3D displacements, and ε the 192 residual values coming from inaccuracies in the observations and from insufficient model fit.

In Camacho et al. (2011a), the surface deformations, dr^c , due to a buried over pressure structure are calculated aggregating the effects for several Mogi point sources (Masterlark, 195 2007). In the present paper dr^c corresponds to the addition of the pressure sources and the 196 Okada's dislocation sources (strike-slip, dip-slip and tensile). Moreover, we substitute the 197 simple point source calculus (Masterlark, 2007) by the more accurate calculus by Geertsma 198 and Van Opstal (1973) for 3D pressure structures.

199 2.2.1. Model description.

200 General geometrical single structures will be described by aggregation of elementary sources filled with causative perturbations (pressure, and strike-slip, dip-slip and tensile 201 202 dislocations). We consider a partition of the medium into a dense 3D grid of m small cells 203 located in (x_i, y_i, z_i) and with small volumes $\Delta V_i = \Delta x_i \cdot \Delta y_i \cdot \Delta z_i$ and small dislocation 204 surfaces ΔSi , i=1,...,m. The data spatial resolution conditions the smaller cell size. Each small 205 cell effect can be modeled by the effect of an elementary source located in its geometric 206 center. We carry out the partitioning by means of small rectangular prisms on horizontal 207 layers, looking for a similar average quadratic deformation effect of each cell upon the whole 208 data set. Then, we calculate the deformation effects, $dr_i (dX_{i}, dY_{i}, dZ_i)$ (see Figure 1), in the n 209 surface (not necessarily gridded) points, $P_i(X_i, Y_i, Z_i)$ (j=1, ..., n) by accumulation of the effects 210 of the filled cells (for $i \in \text{set } \Phi_P$ of pressured, $i \in \text{set } \Phi_S$ of strike dislocation cells, $i \in \text{set } \Phi_D$ 211 for dip-slip and thrust dislocation cells, and $i \in \text{set } \Phi_T$ for tensile dislocation cells):

212
$$dr^{c} = \sum_{\Phi P} \begin{pmatrix} u_{x} \\ u_{y} \\ u_{z} \end{pmatrix} + \sum_{(\Phi S, \Phi D, \Phi T)} \begin{pmatrix} u_{x} \\ u_{y} \\ u_{z} \end{pmatrix}$$
(7)

213
$$dr_j^c = \sum_{i \in \Phi_P} \Delta V_i \Delta \rho_i f_p(r_{ij}) + \sum_{i \in \Phi_S} \Delta S_i \Delta \sigma_i f_S(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_D} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f_D(r_{ij}, \alpha_i, \delta_i) + \sum_{i \in \Phi_P} \Delta S_i \Delta \sigma_i f$$

214
$$\sum_{i \in \Phi_T} \Delta S_i \Delta \sigma_i f_T(r_{ij}, \alpha_i, \delta_i)$$

216 with u_x , u_y , u_z given throw equations (1), (2), (3) and (5).

(8)

Volumes ΔV_i , surfaces ΔS_i and intensity factors $\Delta \rho_i$ (pressure, MPa) and $\Delta \sigma_i$ (dislocation, cm) appear as linear factors in the observation equations (6)-(8), allowing for simple cell aggregation, but the other model parameters (orientation angles α and δ , and sets Φ_P , Φ_S , Φ_D , and Φ_T of filled cells) are nonlinear, necessitating a non-linear inversion approach.

221 2.2.2. Misfit conditions.

Assuming a Gaussian uncertainty given by a covariance matrix Q_D for displacement data, a minimization condition for observation residuals ε , as $\varepsilon^T Q_D^{-1} \varepsilon = \min$, leads to the maximum likelihood solution. For a simplified treatment, Q_D is considered as a diagonal matrix of estimated variances corresponding to the displacement data.

During inversion of geophysical data, problems of singularity and instability for the solution can arise due to inadequate data coverage (normally the number of data points is smaller than the number of unknowns), inaccuracy of the data, and intrinsic ambiguity of the design problem. In this case, they can occur if we assume that positive and negative anomalous pressure/dislocations can be contemporaneously present in the model. A process to avoid instabilities is to consider additional minimization or smoothing conditions for the norm of the solution model as:

233

$$\mathbf{m}^T \mathbf{Q}_M^{-1} \mathbf{m} = \min \;, \tag{9}$$

where the vector \boldsymbol{m} is constituted by the values $\Delta \rho_i$ and $\Delta \sigma_i$ (i=1,...,m) for the filled cells of the model (sets Φ_P , Φ_S , Φ_D , and Φ_T) and \boldsymbol{Q}_M is a suitable covariance matrix corresponding to the physical configuration of cells and data points. This matrix provides a balanced model, avoiding very shallow solutions. We propose a normalizing diagonal matrix \boldsymbol{Q}_M with elements q_i (*i*=1,..., *m*) given for volumes ΔV_i and distances r_{ij} as

239
$$q_i = \frac{\Delta V_i}{n} \sum_{j=1}^n \frac{|z_j - Z_i|}{r_{ij}^3},$$
 (10)

that takes into account the average effect of the *i*-th cell upon all data points.

Condition (9) is a stabilizing term for control on the entire pressure, and dislocations of the structures (Farquharson and Oldenbourg, 1998; Bertete-Aguirre et al., 2002). Weighting by matrix Q_M prevents the occurrence of very large fictitious values of pressure/dislocations, resulting from a model that is poorly determined model due, e.g., to coupling of some positive and negative sources, peripheral sources, etc.

Finally, a mixed minimization equation

247
$$S(\mathbf{m}) = \boldsymbol{\varepsilon}^T \mathbf{Q}_D^{-1} \boldsymbol{\varepsilon} + \gamma \, \mathbf{m}^T \mathbf{Q}_M^{-1} \mathbf{m} = \min \,. \tag{11}$$

is adopted for the constraining equation (6) for residuals and for model magnitude. γ is a factor that provides a balance between fitness and smoothness of the model. Low γ values produce very good data fit but often result in extended and/or irregular models. Conversely, high γ values can produce concentrated and smooth models but with a poorer data fit. The optimal choice is determined by an autocorrelation analysis of the residual values, the value producing a null (planar) autocorrelation distribution (Moritz, 1980; Camacho et al., 2007).

254 **2.2.3** Exploration approach for solving the system.

255 The model system (6)-(8) must be satisfied within the minimization constraint condition 256 (11). It constitutes a nonlinear optimization problem with respect to the geometrical properties 257 (orientation angles α and δ for the dislocations, and sets Φ_P , Φ_S , Φ_D , and Φ_T of filled cells).

258 Considering the very large number of degrees of freedom necessary to describe the 259 pressure and dislocation sources, a general exploratory inversion approach simultaneously 260 applied to the aggregation of thousands of small cells filled with anomalous values would be 261 ineffective. A necessary reduction of the model space is obtained by limiting the possible 262 orientation angles (dip δ and azimuth α), considering only certain orientations. We limit 263 values of α from 0° to 180°, and of δ from 0° to 90°, with step 10°, resulting in 190 possible 264 orientations for each elemental dislocation. After tests on simulated and real data we have 265 concluded that this offers enough detail for most practical applications. Any arbitrary 266 dislocation direction can be fit by a combination of these basic directions.

The model space to be explored is composed by: (1) *m* possible cells to be filled; (2) four primary source possibilities (pressure change, strike, dip or tensile dislocation) for each cell; (3) positive or negative value for each pressure/dislocation cell; and (4) 190 possible orientations for dislocation elements. As previously pointed, coefficients ΔV_i , ΔS_i , $\Delta \rho_i$ and $\Delta \sigma_i$ appear in linear mode and they are solved by a scaled, linear fit.

272 Despite the reduction in angular options, a general exploration of the extensive model domain that considers all possible combinations of thousands of cells, and angles, signs, and 273 274 source natures, would be inefficient. An alternative approach is to build the anomalous 3D 275 structures by means of a step-by-step growth process. The key idea is to substitute a unique 276 global exploratory approach by successive explorations. For each step of the growth process, 277 that exploration allows for selection of only one new optimal cell (and additional parameters) (Camacho et al., 2007; 2011b). This approach explores a model domain clearly smaller at 278 279 every step, composed only of the "empty" cells.

Further, we assume that pressure values $(\Delta \rho)$ and dislocation amplitude values $(\Delta \sigma)$ will be the same over the entire model. These will be expressed as proportional to some basic fixed values $\Delta \rho_o$ and $\Delta \sigma_o$, $\Delta \rho = f \times \Delta \rho_o$ and $\Delta \sigma = f \times \Delta \sigma_o$, f > 0 being a scale factor. $\Delta \rho_o$ and $\Delta \sigma_o$ are arbitrary small fixed values with a fixed ratio $\Delta \rho_o / \Delta \sigma_o$ so that the average effect upon the data of an arbitrary cell with dislocation $\Delta \sigma_o$ will be similar to the one of a pressure arbitrary cell with pressure $\Delta \rho_o$. 286 Considering these conditions, we implement the step-by-step growth process. For the *k*-th 287 step of the growth process, *k* cells have been filled with the prescribed anomalous values for 288 pressure $\Delta \rho_o$ and dislocation amplitude $\Delta \sigma_o$, giving rise to modeled values dr^c from the model 289 equations, which now include a scale factor. For the (k+1)-th step, we fill a new cell fitting 290 the system,

$$dr = f_{k+1} dr^c + \varepsilon$$
(12)

292
$$\boldsymbol{\varepsilon}^{T}\boldsymbol{Q}_{\mathrm{D}}^{-1}\boldsymbol{\varepsilon} + \gamma f_{k+1}^{2}\boldsymbol{m}^{T}\boldsymbol{Q}_{M}^{-1}\boldsymbol{m} = \min, \qquad (13)$$

where $0 < f_{k+1} < f_k$ is a scale factor that fits the modeled deformation field for the provisional, not fully developed, model and the observed deformations. We calculate the value $e^2 = \varepsilon^T Q_D^{-1} \varepsilon + \gamma f_{k+1}^2 m^T Q_M^{-1} m$ for the empty cells according to a general exploratory approach with random selection. We choose as optimal cell to be filled for the (k+1)-th step that *j*-th cell giving:

298
$$e_j^2 = \min$$
. (14)

Throughout the process both f and e^2 decrease. Note that considering the scale factor fmodifies the process from a unique general exploration of the extensive model domain, which would be inefficient, to a much more affordable task: the exploration of aggregation possibilities for a new cell, in a step-by step growth process. This is the primary feature of the inversion approach.

The process continues until: (1) *f* reaches a prescribed small value according to a defined criterion based on previous trials and inspection of the resulting model; or (2) aggregation of a new cell does not produce smaller values of *f* and e^2 . Case (2) produce the larger model, with smaller values for $\Delta \rho$ and $\Delta \sigma$. One potential definition of the stopping criteria could be a prescribed ratio between successive e^2 values. At the final step, we arrive at a 3D model virtually automatically. That model is the aggregation of some filled elementary cells: (1) pressure elementary sources filled with the prescribed anomalous values; and (2) dislocation elementary sources with the appropriate orientation and magnitude. Together, they fit the observed displacement within some error margin and appropriate set of model bounds.

A final test on the validity of the inversion results is done by comparing their geographical
distribution and distances between differences sources, as in Pascal et al. (2014).

Additional details about the practical implementation of the inversion approach aredescribed in Section B of the Supplementary Material.

318 3. Synthetic test cases.

319 To demonstrate the efficiency of this inversion process, we consider a simulated example 320 described in Figure 2, composed of four different deformation sources: a vertical ellipsoid 321 with homogeneous negative pressure (-3 MPa located at 2.5 km depth below the surface and 322 with semi-axes of 2 km and 1.4 km) (Figure 2a); a sub-horizontal strike-slip fault (azimuth 323 65° and dip angle 20° from the horizontal, length 5 km and width 3 km, located at 1.5 km 324 depth) with 12 cm dislocation (Figure 2b); a nearly vertical dip-slip fault (azimuth 30° and tilt 325 angle 20° from the vertical, 4 km vertical side and 7 km horizontal side, mean depth 2.5 km 326 below the surface) with 9 cm dislocation (Figure 2c); and a tensile fault (azimuth 20°, tilt 327 angle 5°, dimensions 2 km and 4 km, mean depth 2 km) with 10 cm of opening (Figure 2d). 328 Above these buried anomalous structures, a planar distribution grid of 800 data points is 329 delineated, with a grid size of 400 m and total diameter of 12 km (Figure 3). The anomalous 330 pressure body, which is sensitive to the diameter of the survey area, occupies a central 331 position below the survey area. The fault structures are located in the borders of the survey 332 area. In this case, we employ a magnitude of 6 MPa (for pressure) and 9 cm (for all 333 dislocation kinds). Figure 3 shows the (a) Up, (b) EW and (c) NS components of the simulated displacement vectors (u_x, u_y, u_z) at the 800 surface points (X_i, Y_i, Z_i) . The average amplitudes of these 800 data values are 2.1 cm, 1.2 cm, and 1.4 cm respectively.

336 For the simulated data (Figure 3), we apply the inversion approach without any a priori 337 assumptions about the 3D structure of the sources. First, we determine a complete 3D 338 partition of the subsurface volume into several thousands of cells with mean side 170 m 339 (Figure 1). The primary decision required concerns the γ parameter. It is selected, after several 340 trials, as that larger value producing a (nearly) null autocorrelation distribution of the final 341 residues for the three components. A secondary assessment is made for the growth stopping criteria. Here we employ a standard threshold value for the ratio e_k^2/e_{k-1}^2 between successive 342 values of the misfit parameter, given as a default value in our software. Once these are 343 344 selected, the 3D model for the deformation sources is obtained automatically. This resulting 345 model is composed of a large aggregation (thousands) of elementary (pressure and 346 dislocations) cells. Figure 3 (right) shows the fit between the simulated (orange) and modelled 347 (blue) data. That fit is quite good, about 0.01 cm for all three components.

Figure 4 shows a flat view from the top of the obtained 3D model defined by aggregation of elemental deformation sources. They reproduce the simulated data (Figure 3) very well and fit the original simulated, pressurized ellipsoid and faults, represented by dashes lines in Figure 4, reasonably well, given that the inversion fit is unconstrained, and involves several simultaneous possibilities for the active structures without specific a priori hypothesis about the number, nature or shape of sources.

354 Considering these results, we outline some observations on the operation of this 355 methodology:

(1) The SE dipping fault structure appears well-constrained, almost entirely composed of
 small dipping elements (yellow in Figures 4 and 5), whose aggregation describes an
 extended body with geometry and location similar to the original body. As expected,

359 considering the regularity conditions, the top of the structure is quite precise, but the360 bottom appears rounded and more diffuse.

- 361 (2) The pressure ellipsoidal structure also is well-characterized, composed largely of an
 362 aggregation of pressure cells (dark blue in Figures 4 and 5) and with geometry and
 363 location similar to the original body.
- (3) The NW sub-horizontal strike structure also is identified and modelled in the inversion
 approach, composed of an aggregation mostly strike elements (green in Figures 4 and
 5). Nevertheless, the sub-horizontal character of the original structure provided some
 difficulties, and we observe a large number of dipping cells whose effect, for subhorizontal structures and limited values of orientation angles, could be close to those
 of strike cells. The geometry also is less precise than for the SE nearly vertical
 structure.
- 371 (4) The tensile structure is the least faithful model here (Figures 4 and 5), composed
 372 primarily by tensile cells but with some distortions.

A well-understood drawback to and unconstrained inversion is that there is a known ambiguity regarding the true values of the magnitude of the sources (MPa for pressure and cm for displacement). The same deformation values can be reproduced with a high magnitude or intensity at deeper depth as with smaller, more shallow structures. Here this issue is related to the selection of the stopping point for the growth of the model.

In the Supplementary Material, Section C, we show additional synthetic cases. First, we show the inversion results for the former source bodies, but as isolated structures, and an isolated spherical source. These isolated studies offer better results that the former combined modelling. Second, we present different simulation studies for combinations of spherical pressurized bodies. Third, we repeat the previous synthetic case combining different structures, but adding a high level of synthetic Gaussian noise to the data (about a 33%). Results show the efficiency in noise filtering, but they also show some deterioration of the model due the noise effects. All this material provides an evaluation of the method's efficacy.

386 An additional consideration is that of the relative confidence corresponding to the fitted 387 model. First, a global confidence of the model comes, as previously pointed out, from the 388 study of the autocorrelation of the residuals and the choice of the value of the smoothing 389 parameter γ . Another interesting approach to the model confidence comes from a study of the 390 sensitivity of the data pixels to the different areas of the 3D model. Indeed, the model 391 characteristics (nature of the source, magnitude, orientation angles, sign) are identified with 392 varying clarity depending on the location of the cells within the subsurface volume and the 393 orientation of the dislocation sources. Cells located in very deep or peripheral areas with 394 respect to the pixels provide a smaller sensitivity and relative confidence (Supplementary 395 Material, Section E).

396 4. Mt. Etna application case.

Mt. Etna (Figure 6) constitutes an excellent test case for applying the inversion methodology detailed above. The volcano was characterized, over the past decade, by persistent volcanic activity as well as a continuous seaward motion of its eastern flank (Palano, 2016). In addition, the large number of SAR images over the region provides a high quality dataset of ground deformation at the scale of the entire volcano. We perform an application of the inversion methodology without a priori assumptions on the numbers, type and 3D geometry of the causative sources.

404 4.1. Deformation data.

405 To study ground deformation at Mt. Etna we collected 38 ascending and 59 descending
406 RADARSAT-2 Standard-3 (S3) images spanning the January 2009 - June 2013 period (see

Figure 7, and Table S1, Supplementary Material, Section D). Each SAR dataset was 407 408 processed independently with the GAMMA software (Wegmuller and Werner, 1997). A 409 single master for each set was selected and remaining images were re-sampled into the master 410 geometry. The spatially averaged interferograms were computed and the topographic phase 411 was removed using the 30 m resolution Advanced Spaceborne Thermal Emission and 412 Reflection Radiometer (ASTER) Digital Elevation Model. Differential interferograms were 413 filtered using adaptive filtering with a filtering function based on local fringe spectrum 414 (Goldstein and Werner, 1998) and unwrapped using the minimum cost flow algorithm 415 (Costantini, 1998). The residual orbital ramp was corrected by applying a baseline refinement 416 algorithm implemented in GAMMA software. For this, the area experiencing large ground 417 deformation was masked out and baseline parameters were re-estimated from the 418 measurements of interferometric phase and topographic height. Minor interpolation of each 419 interferogram was performed in order to improve the spatial coverage reduced by 420 decorrelation. Then, 494 ascending and 298 descending interferograms were geocoded and 421 resampled to a common lat/long grid with the uniform spatial sampling (Table 1). The 422 advanced Multidimensional Small BAseline Subset (MSBAS) method (Samsonov and 423 d'Oreye, 2012; Samsonov et al., 2014) was employed to produce horizontal and vertical time 424 series of ground deformation.

Several inflation/deflation episodes occurred during the 2009-2013 period. However, the GNSS time series show a clear long-term trend, similar to the InSAR average deformation rates. Therefore, as a first order approximation, use of the average deformation rates to describe the long-term trend is justified. Working directly with ascending and descending *LOS* displacements would avoid some uncertainties but it would make interpretation of results significantly more tedious. Two-dimensional deformation rates produced by MSBAS can be easily understood by any user, independently of their knowledge on InSAR. From the total dataset (approximately 451000 pixels) for both *Up* and *EW* components, we extract a reduced subset as input for the inversion approach. We use pixels which verify three conditions: located within 20 km distance from the Mt. Etna summit, with a distance between consecutive pixels of 800 m, and with a coherence value higher than 0,6. The result of this selection is a dataset of 1613 pixels (Figure 7). At this reduced size the inversion method runs faster and the main features of the resulting model will be nearly the same that for the total dataset (see Supplementary Material).

439 4.2 Inverse modelling.

We apply the inversion approach as in the simulated examples, without any assumptions about the nature of the active elements, although we only use two displacement components (*Up* and *EW*) obtained from the InSAR data. First, we consider a 3D partition of the subsurface volume into several thousands of small cells, with an average side of 500 m. As in the synthetic test case, the only one decisions to take before to carry out the inversion involve the selection of the value for the smoothing γ parameter (we apply a trial study about correlation of residuals) and a value (our usual default value) for determining the growth end.

447 Figure 7 shows the data fit for both components (Up and EW) corresponding to the 448 selected ground deformation data for the considered period (orange dots, Figure 7). The 449 remaining residuals do not contain significant spatial autocorrelation, and their standard 450 deviations are approximately 0.1 cm/yr. It is interesting to observe that the proposed inversion 451 approach allows for the identification of data noise and outlier values, while modelling the 452 deformation signal (blue dots, Figure 7). The inversion approach helps to separate 453 perturbations in the deformation data, such as inexact orientation or regional effects. In fact, 454 in such cases, the model will introduce fictitious sources, located in very shallow or very 455 peripheral locations, limiting distortion effects on the real sources.

The resulting 3D model is composed of approximately 12,000 cells filled with one of the deformation elemental patterns: negative or positive pressure and/or dislocation (with values of about 0.5 MPa and 2 cm) for the available directions, plus several thousands of empty cells (see Figures 8-10).

460 *4.3 Discussion*.

The final cell aggregation appears as a rather complex model (Figure 8). However, by isolating subareas and dynamic components of this combined model, the resulting structures identify several interesting features and support several conclusions about the active sources below Mt. Etna volcano.

One important caution is that the input data correspond only to the *Up* and *EW* displacement components, neglecting the *NS* component. But the aim of this paper is not to carry out a complete analysis of Mt. Etna sources. We use this case study in order to show an example of the efficiency and robustness of the method. A detailed discussion about the inversion results would be the objective of another paper. Below we briefly discuss the main sources of deformation inferred from our InSAR data modelling.

471 4.3.1 Plumbing system.

472 In Figure 9, we show some isolated source structures (pressure and tensile cells) that may 473 be related to the plumbing system of Mt. Etna. These appears to be composed of two echelon 474 pressurized reservoirs located at depths of approximately 3000 and 11000 m below sea level, 475 bsl, and a shallower SSW elongated dike structure at a mean depth of 1500 km bsl (Figures 8-476 10). These plumbing structures are located below the western slope of the volcano edifice. 477 The deeper reservoirs are located progressively more SSW, suggesting an ancient location of 478 the eruptive system. Their overall shape and position correspond to the crustal volume where 479 a number of inflating/deflating sources, feeding the volcanic activity during 2009-2013, have

480 been inferred by GNSS-based models (e.g. Patanè et al., 2013; Spampinato et al., 2015;
481 Cannata et al., 2015).

482 Curiously, the shallower structure connected with the plumbing system in this model is a 483 tensile elongated structure (purple color in Figures 8-10), located at approximately 1500 m 484 bsl, that seems to extend almost into the volcano summit. Considering that it is located in an 485 area that is sensitive to this modelling method (Supplementary Material, Section E), we infer 486 that it corresponds to dike structures, separate from the deeper reservoir structures that 487 appears as pressurized cells. Such a structure aligns with the so called "West Rift", a zone of 488 weakness on the western flank where numerous monogenetic pyroclastic cones are aligned 489 along 240-260°N (e.g. Mazzarini and Armienti, 2001).

490 **4.3.2** Pernicana fault and sliding system.

491 Our model also suggests a complex pattern of deformation on the eastern flank of the 492 volcano. In the aggregation model shown in Figures 8 and 10, the main source components 493 are cells for strike- and dip slip dislocations. There is a shallower strike system close to the 494 Pernicana fault (Figure 6), which shows a tilted geometry (see Figure 10e) and a deeper sub-495 horizontal central striking system at a depth about 4 -5 km (Figure 10c). There is also a 496 dipping system in three parts (Figure 10): (a) the shallow header of the downward sliding, 497 both close to the summit (Figure 10a), and inside Valle del Bove (Figure 9b), (b) an intense 498 downward dipping region at 4 km depth (Figure 10c), and (c) a third dipping zone (5 km 499 depth) that corresponds to the thrusting final section of the sliding system.

The strike cells largely correspond to sub-horizontal sliding, and the dip cells determine the dipping pattern (normal in the header and Valle del Bove, and thrusting in the last half). We observe that sub-horizontal dislocations dip and strike sources are combined, similar to the synthetic case. 504 This geometry is rather different from that proposed in the literature, resulting from 505 geophysical-geochemical and magnetotelluric data (e.g. Siniscalchi et al., 2012) and geodetic 506 inversion models (e.g., Palano, 2016 and references therein). The seaward motion of eastern 507 flank of the volcano occurs along a shallow sliding surface bounded by the North Rift -508 Pernicana fault system and the South Rift - Mascalucia - Tremestieri - San Gregorio -509 Acitrezza fault system, northward and southward, respectively (e.g. Palano, 2016). Since no a 510 priori constraints have been adopted during the inversion, the south-dipping planar surface 511 resulting from the inversion probably represents an "average source" of the sub-horizontal 512 sliding surface and the $\sim 60^{\circ}$ S-dipping Pernicana fault system. However, where the modelled 513 planar dislocation intersects the volcano surface corresponds to the Pernicana fault system, 514 capturing the boundary between the undeformed sector (northward of the fault) and the 515 unstable region of the eastern flank of the volcano.

The localized subsidence structure below Valle del Bove, represented by dip cells and a depressurized body, may be related to: (i) the cooling and compaction of the lava flows that in the last decade accumulated on the western side of Valle del Bove (e.g. De Beni et al., 2015), and/or (ii) a process of relaxation of the substrate in response to loading produced by deposited lavas (e.g. Briole et al., 1997).

521 5. Conclusions.

We have presented a new inversion methodology for modeling geodetic displacement data in active volcanic areas which permits simultaneous inversion of the several components of surface deformations and allows for a global fit of the data. Non-planar and non-gridded data can be employed in this approach.

526 The method allows for objective modelling of diverse causative structures as pressure 527 bodies, and general dislocations (strike-slip, dip-slip and tensile). Well-known analytical

528 expressions from Okada (1985), for elemental dislocation sources, and Geertsma and Van 529 Opstal (1973), for pressured small prisms, are used for direct calculation. They assume a 530 semi-infinite elastic medium, characterized by some values of the elastic parameters. The 531 assumptions of linear elasticity and isotropy allows for the final modeling by superposition of 532 effects for elemental components (prisms and dislocations) form the obtained aggregated 533 geometry.

534 The approach works in a fully 3D context, although it employs, for faster operation, 535 elementary dislocation sources limited to a discrete set of orientations. A free 3D geometry of 536 the causative structures is described by aggregation of small elemental cells. There are not 537 additional a priori requirements on the geometry and types of the causative sources. The 538 method is able to automatically determine the number, nature and 3D geometry of the 539 causative source structures, and supports different type of deformation data, such as 540 GNSS/GPS, InSAR (horizontal and vertical components, or ascending and/or descending LOS 541 data), leveling data, and others. The inversion process constitutes an interesting tool for 542 integrating simultaneously terrestrial and spatial data, providing mapped models which incorporate all the available data. 543

544 This new methodology allows for a nearly automatic approach that takes advantage of 545 the large and precise datasets coming from ground-based deformation and advanced DInSAR 546 techniques and carries out an exhaustive inversion of ground deformation data to better 547 understand the subsurface causative structures and elastic processes, without preconceived 548 hypotheses. It can be applied on large regional scales to model tectonic plate movements and 549 subduction, volcanic activity and, on more local scales, to model deformation from landslides, 550 volcanic eruptions, and anthropogenic subsidence due to mining and extraction of oil, gas, or 551 groundwater. Additionally, this new inversion methodology can be used to invert coseismic 552 geodetic deformation data, as detailed in Section F, Supplementary Material.

In particular, for the InSAR data of Mt. Etna 2009-2013, the application of this methodology resulted in a model for several subsurface sources corresponding to the plumbing system, the subsidence within Valle del Bove and the seaward motion of the eastern flank of the volcano.

557 Several precautions should be noted. First, as for other geophysical inversions, the 558 problem has an intrinsic ambiguity. It is solved by use of regularity conditions. Solutions must 559 be interpreted carefully as informative models constrained by limitations in data and 560 smoothing constraints, particularly when applied, as here, within a range of potential 561 causative sources. Second, confidence in the solutions is not uniform. Peripheral or very deep 562 elements will be relatively less valuable (Supplementary Material, Section E). Third, in some 563 cases this approach allows for the separation of perturbing effects (noise, outliers, etc.) in the 564 deformation data.

Finally, there are some potential limitations on the validity of the results depending on the combination and sizes of the detected sources. The resulting combination of 3D sources, nature, geometries and relative distances should be examined for inconsistent results, as described in Pascal et al. (2014).

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Figures



- Figure 1. Partition of the subsurface volume below the survey into a 3D grid of thousands of
- small right prisms. Blue triangles correspond to data points (terrestrial stations or pixels) P_i with coordinates (x_i, y_i, z_i) and observed deformation vector dr_i .





Figure 2. Synthetic source structure composed by: (a) a vertical ellipsoid with a
decreasing pressure (blue), (b) a sub-horizontal strike slip fault (green), (c) a nearly vertical
dip slip fault (yellow), and (d) a nearly vertical tensile fault (brown). See text for details on
the sources characteristics.



Figure 3. Left panels show the simulated deformation values corresponding to the active
structures shown in figure 2 for the 800 data points used for inversion. (a) Up component. (b)
EW component. (c) NS component. Right panels show the data fit corresponding to the
inverse model for each component. Observed data are plotted in orange, modelled in blue.





Figure 4. Horizontal view from the top of the resulting source structures, described as
aggregation of different elemental source cells, and obtained by application of the inverse
approach. Dots indicate data sites, and discontinue lines the location of the synthetic bodies.
Arrows show displacement patterns for sources.





Figure 5. Inverse 3D source model as aggregation of elemental cells. (a) Horizontal sections

- *at 1500 m depth;* (b) EW vertical profile across the strike structure (green cells); (c) EW
- *vertical profile across the low pressure structure (blue cells) and dip-slip structure (yellow*
- *cells*). *Modelling magnitudes are 0.5 MPa and 1.5 cm.*







Figure 7. Observed displacement rate (cm/year): (a) Vertical and (b) EW component for the
 1613 pixels selected from the total dataset (Figure 6). Comparison between observed and
 modelled values: (c) Vertical and (d) EW component. Final residuals corresponding to local
 effects: (e) Vertical and (f) EW component.





Easting UTM (m)

2 km

-Dip

+Dip +Tens. -Tens.

773 denotes the Permicana Fault, Valle del Bove limit and coast line (see Figure 6).





Figure 9. Perspective 3D view of those source elements in Figure 6 corresponding to the
plumbing system: increasing pressure cells (large deep reservoir with mean depth ~11 km bsl,
SW of Etna, and shallow small reservoir with mean depth 3 km bsl. NW of Etna and along the
elongation of the deep reservoir), and expanding tensile cells (at levels 1 km and 2 km bsl,
with elongated pattern SW-NW). Black lines correspond to Etna summit, Pernicana Fault and
Valle del Bove limit (see Figure 6).



Figure 10. Some horizontal and vertical sections of the tensile structure (purple), strike-slip
structure (green), dip-slip structure (yellow) and high pressure (red) from the inverse model.
(a),(b),(c) and (d): Horizontal sections at depths 500, 1000, 4000 and 9000 m bsl. (e) SN
vertical section across the strike structure. (f) and (g): WE vertical sections. (h) Vertical
section with azimuth 100°. Green lines correspond to Etna summit, Pernicana Fault and Valle
del Bove. Location of vertical sections (e) to (h) are indicated by lines in panel (c).









Figure 4 Click here to download Figure: Figure4.pdf















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Declaration of interests

 \boxtimes The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

□The authors declare the following financial interests/personal relationships which may be considered as potential competing interests:

CRediT author statement

Antonio G. Camacho: Conceptualization, Methodology, Software, Validation, Formal Analysis, Investigation, Writing-Original Draft, Writing-Review & Editing, Funding acquisition. José Fernández: Conceptualization, Methodology, Software, Validation, Formal Analysis, Investigation, Writing-Original Draft, Writing-Review & Editing, Project administration, Funding acquisition. Sergey V. Samsonov: Formal analysis, Investigation, Data curation, Writing-Review & Editing. Kristy F. Tiampo: Formal Analysis, Investigation, Data curation, Writing-Review & Editing. Mimmo Palano: Formal Analysis, Investigation, Writing-Review & Editing.