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Mars upper atmospheric responses to the 10 September 2017 solar flare: A global, time-dependent simulation

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Key Points:
Ionospheric perturbation follows the flare in time and is concentrated mostly below 110 km altitude.
Neutral atmospheric percent changes increase with altitude and is important above 150 km altitude.
It takes the neutral atmosphere 2.5 hours to reach the peak and 10 more hours to generally recover.

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23 Abstract

We report the first global, time-dependent simulation of the Mars upper atmospheric 24 responses to a realistic solar flare event, an X8.2 eruption on 10 September 2017. The 25 Mars Global Ionosphere-Thermosphere Model runs with realistically specified flare irradi-26 ance, giving results in reasonably good agreement with the Mars Atmosphere and Volatile 27 EvolutioN spacecraft measurements. It is found that the ionized and neutral regimes of the 28 upper atmosphere are significantly disturbed by the flare but react differently. The iono-29 spheric electron density enhancement is concentrated below ~110 km altitude due to en-30 hanced solar X-rays, closely following the time evolution of the flare. The neutral atmo-31 spheric perturbation increases with altitude and is important above ~ 150 km altitude, in 32 association with atmospheric upwelling driven by solar EUV heating. It takes ~ 2.5 hours 33 past the flare peak to reach the maximum disturbance, and then additional ~10 hours to 34 generally settle down to pre-flare levels. 35

1 Introduction

Solar flares represent an important type of space weather event, in which a tremen-37 dous amount of energy is released into the heliosphere in the form of radiation bursts and 38 hence imposes significant disturbances upon planetary atmospheres. With dramatic pertur-39 bations on solar irradiance, solar flares offer an invaluable opportunity to test our under-40 standing and constrain first-principles modeling of how solar ionizing and heating fluxes 41 dissipate and redistribute the energy in atmospheric and ionospheric systems. An accu-42 rate description of upper atmospheric processes is critical not only for understanding the 43 higher-altitude plasma environment and atmospheric loss by solar wind stripping, but also 44 for the safety of current and future Mars orbital platforms. 45

46 While there have been numerous studies on the effectiveness of solar flares at Mars,

⁴⁷ nearly all of them focus on ionospheric responses [*Gurnett et al.*, 2005; *Nielsen et al.*,

48 2006; *Mendillio et al.*, 2006; *Haider et al.*, 2009; *Mahajan et al.*, 2009; *Lollo et al.*, 2012;

Haider et al., 2012; Fallows et al., 2015] and little is known about the thermospheric im-

⁵⁰ pact of solar flares [e.g., *Thiemann et al.*, 2015]. Historically, the main challenge in the

study of the Mars upper atmosphere has been the lack of systematic and comprehensive

neutral species observations except for limited knowledge derived from sparse aerobraking

activities [e.g., *Bougher et al.*, 2000, and references therein]. Moreover, there has been a

lack of solar irradiance measurements at the Mars' orbit until the NASA Mars Atmosphere 54 and Volatile EvolutioN (MAVEN) mission [Jakosky et al., 2015], which for the first time 55 carries both solar EUV and neutral particle detectors, suitable for solving the cause-and-56 effect connection between the Sun and Mars. Different from previous unpublished confer-57 ence presentations performing generic model runs for solar flares, in this study we make 58 the first numerical attempt to quantify global perturbations of the Martian upper atmo-59 sphere in response to a real solar flare event using realistic flare irradiance, and to make 60 direct model-data comparisons for the flare effects. 61

⁶² 2 The 10 September 2017 Solar Flare Irradiance at Mars

On 10 September 2017, one of the most powerful solar flares in the recent decade 63 erupted from the solar active region AR2673 and impacted Mars. The activities from 64 AR2673 also include an eruption of a fast and wide coronal mass ejection (see *Lee et al.* 65 [2018] for an overview). The X8.2-class solar flare eruption manifests itself in dramatic 66 enhancement over a broad wavelength range including X-ray and extreme ultraviolet (EUV). 67 It has been found by terrestrial solar flare studies that thermospheric responses are more dependent on time-integrated energy inputs than on peak irradiance fluxes [e.g., Pawlowski 69 and Ridley, 2008, 2011]. Therefore, to yield a reasonable assessment of the flare effec-70 tiveness in the Martian upper atmosphere, we need not only a detailed description of the 71 flare irradiance spectra but also their evolution with time during the event. There is also 72 a need for extrapolating direct solar irradiance measurements by the MAVEN EUV Mon-73 itor (EUVM) within three discrete finite-wavelength channels (0.1-7 nm, 17-22 nm, and 74 121-122 nm, see *Eparvier et al.* [2015]) to a broad radiation range that is of importance 75 to atmospheric absorption. Because of an especially high solar corona temperature associ-76 ated with the flare, we adopt a physics-based spectral irradiance model for the wavelength 77 range of 0.1-36 nm except for 30.5 nm, which uses flare plasma temperature measure-78 ments made from Earth and soft X-ray irradiance measurements made by EUVM. The 79 EUVM 121.6 nm channel is used to estimate the 30.5 nm irradiance, and direct flare spec-80 tral measurement made from Earth by SDO EVE are used from 36-106 nm. The routine 81 estimates of FISM-M [Chamberlin et al., 2007, 2008; Thiemann et al., 2017] are used 82 above 106 nm. A detailed description of this composite irradiance spectrum has been 83 given by *Thiemann et al.* [2018], in which flare irradiance observations at Earth and pho-84 toelectron observations at Mars indicate that the spectra used here are an improvement 85

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over the EUVM Level 3 (L3) spectra. It is thus speculated that the error/uncertainty of the
 spectral irradiance model for this study is better than that of the L3 model, whose upper
 limit is about 40% [*Thiemann et al.*, 2017].

Figure 1 shows the calculated solar irradiance spectra (in 1-nm wavelength resolu-89 tion) and their evolution with time during the flare event. The transient nature of flares 90 is well demonstrated in Figure 1a: the photon fluxes had an abrupt rise within ~ 15 min 91 before reaching the peaks and then gradually recovered and largely dropped back to the 92 pre-flare level about 4 hours later. For the comparison purpose, we select three time points 93 on September 10 15:00 UT, 16:15 UT, and 17:42 UT as representatives of pre-flare, peak-94 flare, and post-flare conditions, respectively. It is well known that solar flares have differ-95 ent time scales in onset and decay characteristics at different wavelengths [e.g., Fletcher 96 et al., 2011, and references therein], which is also seen in Figure 1a in the X-ray and EUV 97 irradiance changes with time. The choice of the flare peak at 16:15 UT is thus some-98 what arbitrary, which, nevertheless, is adequate to help characterize the time scales in 99 association with atmospheric perturbations. Our results, which will be shown later, illus-100 trate that the time scale in the responses of the neutral regime of the upper atmosphere 101 is much longer than that in the flare spectral variability. In addition, the post-flare time 102 of 17:42 UT is selected, because it is the time when MAVEN reached periapsis of ~ 155 103 km altitude. Note that the orbital period of the spacecraft is about 4.5 hours, which means 104 MAVEN missed the chance to closely observe the upper atmospheric responses during the 105 peak of the flare event. This, on the other hand, underscores the importance and irreplace-106 ability of global modeling in a time-evolving manner, like in the present study. The brief 107 bite outs within 1-10 nm wavelengths at a time cadence of the MAVEN orbital period are 108 not real but caused by the instrument effects of EUVM, which either pointed away from 109 the Sun or happened to not open its aperture. These radiation bite outs have an insignifi-110 cant effect because of being well outside of the flare event. 111

Figures 1b and 1c show that the flare spectral intensity has the most pronounced variability at short wavelengths, particularly <20 nm. The short-wavelength end of the spectrum undergoes rapid changes in both rising and decay phases. The time sequence in Figure 1a shows that at 16:15 UT, the total solar fluxes integrated over 0-10 nm, 10-20 nm, and 20-100 nm are enhanced by a factor of 8.68, 2.90, 1.23, respectively, in comparison with the pre-flare level at 15:00 UT. The respective irradiance enhancement factors significantly dropped to 2.92, 1.39, 1.13 at 17:42 UT, and further to 1.63, 1.08, 1.06 at

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20:00 UT. This indicates two main characteristics of the solar flare development: short du ration (~4 hours for this case) and wavelength-dependent variability (greater changes at
 the shorter wavelengths).

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3 Numerical Simulation of Upper Atmospheric Effects

The Mars Global Ionosphere-Thermosphere Model (MGITM) [Bougher et al., 2015a,b] 123 is adopted to investigate the solar flare impact on the Martian upper atmosphere. MGITM 124 combines the terrestrial GITM framework of Ridley et al. [2006] with Mars fundamental 125 physical parameters, ion-neutral chemistry, and key radiative processes to capture the ba-126 sic observed features of the thermal, compositional, and dynamical structure of the Mars 127 atmosphere from the ground to \sim 300 km altitude. MGITM solves for the bulk horizontal 128 neutral winds, while in the vertical direction, the momentum equation is solved for each 129 of the major species. Key neutral species include CO₂, CO, O, N₂, O₂, N(⁴S), N(²D), NO, 130 Ar, and He. Key ion species include O^+ , O^+_2 , CO^+_2 , N^+_2 , and NO^+ . An important feature 131 of MGITM distinct from conventional general circulation models is the use of altitude 132 grids instead of pressure grids. The altitude-based system allows for the relaxation of the 133 hydrostatic equilibrium assumption and enables the model to capture sound and gravity 134 waves in vertical and horizontal directions. In the present study, MGITM runs at a high 135 resolution of 2.5° longitude by 2.5° latitude by 2.5-km altitude (~0.25 scale height). The 136 time resolution of the model is about a few seconds (which is dynamically adjusted), al-137 though we output the model results every 5 minutes during the flare time period. The lo-138 calized crustal magnetic field, which adds complexity to the near-Mars space environment 139 [e.g., Fang et al., 2015, 2017], is neglected. In this work, we focus more on the flare im-140 pact from a system perspective than small-scale or regional disturbances. 141

In order to reasonably describe the Martian thermospheric and ionospheric state 142 changes during the space weather event, we start the MGITM run ~ 60 Martian solar days 143 prior to the flare onset, assuming constant solar irradiance inputs at a pre-event level of 144 2017-09-03/00:00 (>7 days before the X-flare). The purpose of the preconditioning run is 145 to spin up the global dynamics to achieve a pseudo steady state before the flare. MGITM 146 then runs using time-varying, realistically configured solar inputs (at 1-minute time ca-147 dence, as seen in Figure 1) in the next 9 days from 09-03/00:00 till 09-12/00:00. Note that 148 several relatively weak M-class solar flares happened during September 8-9 prior to the 149 examined X-class flare. Figures 2a-2e present the abundance altitude profiles of five key 150

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151	neutral species (CO ₂ , O, CO, N ₂ , and Ar) retrieved from the MGITM results along three
152	MAVEN periapsis passages. These spacecraft tracks span the pre-flare, near-post-flare,
153	and far-post-flare phases of the event, with periapsis passage times of 09-10/08:49, 09-
154	10/17:42, and 09-11/02:34, respectively. Figures 2f-2j show the percentage changes in the
155	neutral densities along the near-post-flare and far-post-flare periapsis passages relative to
156	the pre-flare values at the same altitudes. The in-situ neutral measurements for comparison
157	are from the MAVEN Neutral Gas and Ion Mass Spectrometer (NGIMS) [Mahaffy et al.,
158	2014, 2015; Benna et al., 2015]. Here we use only inbound segments to exclude potential
159	contamination on the instrument. Complementary discussions of the MAVEN observations
160	of the Martian upper atmosphere and ionosphere during this event have been given by <i>El</i> -
161	rod et al. [2018] and Thiemann et al. [2018], respectively.

The model-data comparison from pre-flare to post-flare in Figures 2a-2e shows that 162 MGITM generally captures the basic structures of the upper atmospheric density pro-163 files along all the three examined MAVEN orbits. The model results agree reasonably 164 well with the data for CO₂, CO, and Ar, while significant model deviation is found, in-165 cluding underestimation of the abundances for O (particularly below ~180 km altitude) 166 and for N₂. The detailed examination of the atmospheric density perturbations in per-167 cent, as presented in Figures 2f-2j for both the model and data, illustrates a dramatic den-168 sity enhancement in all the key neutral species during the flare and then a general recov-169 ery along the far-post-flare orbit. The MAVEN data indicate that the densities along the 170 near-post-flare orbit (in red) increase more with increasing altitude, from by up to about 171 50% at altitudes lower than \sim 190 km to by a factor of 3 or more at higher altitudes. The 172 model captures the increasing trend with altitude, while the great enhancement ampli-173 tude above the exobase (which is typically located at around 200 km altitude) is missed 174 by the model. This is partly because the model is subject to more limitations in physics 175 as neutral species gradually change from a fluid-like behavior in the thermosphere to-176 ward a ballistic motion across the exobase. Along the far-post-flare orbit (in blue), the 177 model accurately reproduces the slight decrease in the thermospheric concentrations but 178 misses the reversed change in the exosphere. In addition, the wave-like structures in the 179 observations are not accounted for in the model run. Nevertheless, the comparisons as 180 seen in Figure 2 show that our simulation reasonably captures the neutral density enhance-181 ment during the flare and the subsequent recovery, on both spatial and temporal scales. It 182 should be pointed out that no ad hoc tuning or adjustment has been made to the MGITM 183

model for this specific event, except for the solar irradiance specification as described be-184 fore. Considering the complexity and challenging nature of modeling a global system in 185 a time-evolving fashion, the agreement as seen in Figure 2 is remarkable and underscores 186 the usefulness of the model in understanding of the Martian upper atmospheric behavior 187 of the first order [Bougher et al., 2015b]. While the model-data discrepancy indicates an 188 opportunity to identify potential processes that could be improved or considered in future 189 work (see Bougher et al. [2015a] for discussions of MGITM simplification and empirical 190 approximations), the numerical study that we report here represents one of the best model-191 ing capabilities that are currently available to the Mars upper atmospheric community. 192

The direct orbit-to-orbit comparison is straightforward but does not necessarily rep-193 resent the true atmospheric perturbations solely due to the space weather event. The Mars 194 system is dynamic in nature and is seldom in a steady state even under quiescent solar 195 conditions. Large orbit-to-orbit variability has been reported in the Martian upper atmo-196 sphere [Bougher et al., 2015b, 2017; Zurek et al., 2017]. The changes as seen from orbit 197 to orbit implicitly result from many variability sources other than the flare, including, for 198 example, longitudinal variations of atmospheric heating due to largely inhomogeneous dis-199 tributions of thermal inertia and albedo [e.g., Putzig et al., 2005]. The wide longitudinal 200 span among the orbits due to planetary rotation contributes in part to the changes shown 201 in Figure 2. To add to the complexity, the MAVEN orbital projection in the Mars-centered 202 Solar Orbital (MSO) coordinate system is also not fixed but precesses with time. In or-203 der to reliably retrieve the thermospheric perturbations only due to the 10 September 2017 204 flare, we run a benchmark case for the non-flare scenario, similar to the approach taken 205 by the terrestrial study of Pawlowski and Ridley [2008]. The non-flare case runs under the 206 identical conditions over the same time frame as used in the flare case except that the so-207 lar irradiance starting from 09-10/15:00 is held constant at the minimum post-flare level 208 during 2017-09-11. A comparison of these two time-varying cases enables us to quantify 209 the net effects that the flare has on the upper atmosphere and their time evolution. 210

Figure 3 describes the net flare effects in the dayside upper atmosphere. Figures 3ah give the percentage changes by subtracting the non-flare case from the flare case and then dividing the difference by the non-flare case. The examined parameters in panels ah correspond to electron density, neutral temperature, neutral pressure, CO_2 , O, CO, N_2 densities, and O/CO₂ density ratio, respectively. The altitude profiles for comparison are obtained by averaging over the entire dayside for solar zenith angle (SZA) less than 90°,

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using corresponding horizontal areas as weights. A prominent feature as seen in Figure 3 is that from a system perspective, the Martian ionosphere and neutral atmosphere on average undergo significant increase in density and temperature and apparent decrease in the mixing ratio of O relative to CO_2 in response to the solar irradiance enhancement during the flare. It takes the upper atmosphere more than 12 hours past the flare peak to generally settle down to the pre-flare level. In what follows, we discuss in detail how the Mars system is disturbed.

One response difference between the upper atmospheric neutral and ionized regimes 224 is on their temporal development: they both react instantaneously but with distinctly dif-225 ferent time scales. The ionospheric density increase, which is the most pronounced below 226 110 km altitude, is closely in line with the increase in X-ray photon fluxes and thus the 227 resulting photoionization. The short reaction time of the ionosphere is due to fast pho-228 tochemical reactions. This is also seen in the negligible time delay between brief iono-229 spheric depletions (after ~21:55 UT and ~23:35 UT) and artificial solar shortwave radi-230 ation bite-outs (as discussed in Figure 1a). Since these instrument effects hardly impact 231 the atmosphere, we didn't make corrections but instead find them useful as a diagnostic 232 of the ionospheric response. As a comparison, the atmospheric disturbances gradually 233 develop following the flare onset and reach the highest level approximately at 18:45 UT, 234 about 2.5 hours after the flare peak. The significantly slower response is because of the 235 time needed for neutrals to accumulate, dissipate, and redistribute the absorbed solar en-236 ergy. Similar findings have been found in terrestrial flare-impact studies [e.g., Liu et al., 237 2007; Pawlowski and Ridley, 2008], showing that there is no apparent one-to-one corre-238 spondence between solar inputs and upper atmospheric states. Instead, the integral of solar 239 radiation over a time history is more important than instantaneous irradiance. This poses 240 the difficulty of attributing neutral perturbations to solar irradiance at a specific time point. 241

The other difference between the ionospheric and atmospheric responses is on the 242 perturbation domain and magnitude. Our results suggest that the ionospheric electron den-243 sity may increase substantially by up to an order of magnitude in this flare event, mostly 244 concentrated at low altitudes of ~55-105 km (with the maximum percentage increase at 245 \sim 70 km). Note that the electron concentration in this region (where photoionization is 246 from solar X-rays) is orders of magnitude lower than that in the main ionospheric layer 247 (which is typically above 120 km with photoionization mainly from solar EUV). Figure 248 3 shows that the main ionospheric density enhancement is indeed moderate: up to 25%249

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near 210 km altitude. For the neutral upper atmosphere, its perturbations are concen-250 trated at high altitudes (mostly above 150 km), and the percentage increase grows with 251 increasing altitude. Within the MGITM spatial domain of <300 km altitude, the maximum 252 flare-induced changes in the dayside-averaged properties are 7% for the neutral tempera-253 ture, 46% for the thermal pressure, 122%, 34%, 73%, and 66% for the densities of CO_2 , 254 O, CO, and N_2 , respectively. Due to the different increase in O and CO₂, their density 255 ratio is reduced by up to -40% in the event. The high-altitude concentration of the at-256 mospheric effects can be explained by the fact that solar EUV heating dominates at high 257 altitudes and quickly drops below ~160 km [e.g., Bougher and Dickinson, 1988]. The pre-258 dicted perturbation amplitudes are consistent with the enhancement of EUV inputs (see 259 Figure 1). However, the real impact in the exosphere (above 200 km) would probably have 260 been greater, where an underestimation of the model is implied by Figure 2. Moreover, 261 because MGITM uses a single temperature to approximate the bulk behavior of atmo-262 spheric species, the actual heating effect on some species could be greater than our pre-263 diction here [Elrod et al., 2018]. 264

In Figure 3i, we assess the upper atmospheric movement during the flare event by 265 evaluating the altitude change (in units of km) of fixed pressure levels between the MGITM 266 non-flare and flare cases. The pressure levels of 10^{-8} Pa, 10^{-5} Pa, and 10^{-2} Pa are lo-267 cated near the altitudes of 260 km, 135 km, and 86 km, respectively, at 09-10/15:00 in 268 the non-flare case. Given that the pressure is a proxy of the atmospheric column mass, 269 Figure 3i illustrates that the solar flare results in a significant upwelling in the dayside 270 Martian atmosphere. At the time of the atmospheric disturbance peak (18:45 UT), the 271 vertical expansion ranges from ~ 1 km near 135 km altitude to ~ 10 km near 260 km al-272 titude. The upper atmospheric upwelling is consistent with the increase of the neutral 273 species abundances at high altitudes (Figures 3d-3g) and also explains the ionospheric 274 density enhancement there (Figure 3a). The ionospheric intensification at low altitudes 275 (<110 km) is caused by the enhanced solar ionizing fluxes in the flare event, specifically 276 in hard and soft X-ray wavelengths. The ionospheric density increase at high altitudes 277 (>150 km), however, needs a careful examination. Its increase during the main flare burst 278 directly results from the irradiance enhancement in the EUV range. On the other hand, 279 the remarkable increase, which lasts >8 hours with the maximum amplitude reached hours 280 after the flare peak, indicates an indirect effect. Because a photochemical equilibrium ap-281 proximation is taken for the ionosphere in MGITM, the high-altitude ionospheric enhance-282

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ment during the flare recovery phase must be caused by the atmospheric expansion, which brings more neutral species to high altitudes and leads to more local solar ionizing energy absorption. It is realized that the calculated ionospheric results as presented here are subject to model limitations due to the neglect of transport effects (whose importance starts to increase above ~180 km altitude). This study focuses more on the understanding of neutral disturbances, and a more accurate modeling of the ionosphere could be included in a future work using a magnetohydrodynamic approach.

Figure 4 shows the horizontal distributions of the flare-induced atmospheric pertur-290 bations at 251.25 km altitude, as a function of MSO latitude and local time. We select 291 four representative time points to examine the percentage differences between the MGITM 292 non-flare and flare cases: 2017-09-10/16:15 (approximately flare peak), 2017-09-10/18:45 293 (approximately atmospheric perturbation peak), 2017-09-11/00:00 and 2017-09-11/05:00 294 (in the recovery tail, ~ 8 hours and ~ 13 hours after the flare peak, respectively). These 295 horizontal variations provide supplemental information to the dayside-averaged altitude 296 profile examination as conducted in Figure 3. It is illustrated that the upper atmospheric 297 disturbances start and accumulate on the Sun-facing side in response to the flare impact, 298 and at the same time propagate and diffuse into the nightside. The dayside perturbations 299 demonstrate a general SZA dependence, although a dawn-dusk asymmetry exists with the 300 maximum percentage increase in the morning sector. In the late recovery phase, while the 301 dayside disturbances have mostly subsided, some residual changes are seen on the night-302 side. These results underscore the complexity of the upper atmospheric responses to solar 303 flares, on both temporal and spatial variations. 304

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4 Summary and Discussion

In this study we use the MGITM model to perform a global, time-dependent nu-306 merical simulation of the Mars upper atmospheric and ionospheric responses to the X8.2-307 class solar flare eruption during 10 September 2017. The flare irradiance for driving the 308 model, covering a broad wavelength range of 0-190 nm at 1-minute time cadence, is spec-309 ified by a spectral irradiance model using both in-situ MAVEN EUVM measurements and 310 Earth measurements for improved accuracy. By comparing two time-dependent runs for 311 the non-flare and flare scenarios, we find that the solar flare results in instantaneous inten-312 sification in the dayside ionospheric electron density, most pronounced at altitudes lower 313 than ~ 110 km due to the dominance of the flare enhancement at the short-wavelength end 314

of the spectrum. There is a close correlation between the changes of electron densities 315 and solar ionizing fluxes in both perturbation magnitude and in time scale. In contrast, 316 the solar flare effectiveness in the neutral atmosphere proceeds through accumulation and 317 redistribution processes on the Sun-facing side, with the maximum perturbations reached 318 about 2.5 hours after the flare peak. Our model results predict a remarkable increase in 319 neutral species abundances: by up to 122%, 73%, 66%, and 34% for CO_2 , CO, N_2 , and 320 O, respectively. The neutral atmospheric disturbance is primarily concentrated at altitudes 321 higher than ~ 150 km, generally increasing its amplitude with rising altitude. In accor-322 dance with the flare-induced atmospheric upwelling due to solar EUV heating (ranging 323 from an upward movement of ~ 1 km at 135 km altitude to ~ 10 km at 260 km), the high-324 altitude ionosphere during the recovery phase of the flare is subject to a moderate increase 325 of up to 25% at ~ 210 km altitude through the photoionization increase. It is also shown 326 that the dayside atmospheric disturbance propagates and diffuses into the nightside. It 327 takes the Mars system more than 12 hours in total to generally recover to pre-flare levels. 328

The MGITM results have been compared with MAVEN in-situ measurements along 329 spacecraft periapsis passages. While the comparison with the MAVEN data suggests that 330 the model may have underestimated the solar flare impact at high altitudes, the general 331 model-data agreement is satisfactory. The atmospheric density perturbations are reason-332 ably captured during the flare and the subsequent recovery, on both spatial and temporal 333 scales. There are two noteworthy advantages of the modeling approach to satellite obser-334 vations. First, not limited to the investigation of the atmospheric time sequence during the 335 flare event, our numerical study enables retrieval of net flare effects. By subtracting the 336 MGITM results of the non-flare (pseudo) case from those of the flare (realistic) case, we 337 effectively minimize the impact of the current modeling challenge in replicating all the 338 details of satellite-observed atmospheric states. Furthermore, we mitigate the interference 339 from other variability sources that are implicitly included in orbit-to-orbit changes, such 340 as longitudinal effects. Our results reflect our best understanding of the Mars system's 341 response solely to the solar flare, which stems from our current understanding of upper at-342 mospheric physical processes that are included in the model. The general validity of the 343 model has been confirmed [Bougher et al., 2015a,b]. Second, the flare disturbance is as-344 sessed in a spatially global and temporally continuous manner. As a comparison, in-situ 345 data have very limited spatial and temporal coverages. This work represents the first nu-346 merical attempt to realistically simulate the Mars upper atmospheric responses to a real 347

solar flare event and to make direct model-data comparisons for the resulting perturbations. It is illustrated that the neutral regime is not exempt from the influence by space
weather events, including solar flares (this work) and interplanetary coronal mass ejections
[*Fang et al.*, 2013]. It is of great science interest to explore in the future whether and how
flare-induced perturbations in the upper atmosphere and ionosphere could propagate upward to the magnetosphere through coupling processes, particularly during stronger solar
flares.

It is suggested that the processes that shape the Mars upper atmosphere during and 355 after a solar flare are similar to those processes at Earth. Terrestrial studies have shown 356 that solar flares result in atmospheric expansion and thermospheric density increases [e.g., 357 Pawlowski and Ridley, 2008; Qian et al., 2011] and that the atmosphere slowly returns to 358 the pre-flare state after dissipating the absorbed solar flare energy [Pawlowski and Ridley, 359 2011]. Despite the similarities, at Mars there are differences that play a role in modifying 360 how its upper atmosphere responds to a flare event. For example, Pawlowski and Ridley 361 [2008] simulated the response of the terrestrial upper atmosphere to a stronger X17 flare 362 but found much weaker responses (in terms of percent changes) than what we present here 363 for the relatively weaker X8.2 flare at Mars. At a first glance, this is not straightforward 364 because solar forcing at Mars may be thought to play a less significant role in driving 365 thermospheric disturbances due to the longer distance to the Sun [Bougher et al., 2015a]. 366 Nevertheless, the thermospheric response is driven not only by the absorption of solar X-367 ray and EUV photons, but also by the efficiency of energy redistribution and dissipation. 368 The dominant energy loss mechanisms at Mars (i.e., thermal conduction and CO_2 cooling) 369 turn out to be less effective at removing the excess energy than at Earth (where O and NO 370 cooling are important). To further investigate the differences that the heating and cooling 371 processes play at their respective planets, it would be helpful to conduct a comparative 372 study for a same solar flare event. Such an investigation is the topic of future work. 373

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Figure 1. The calculated solar irradiance and variation with time during the 10 September 2017 solar flare event. Panel (a) shows the irradiance integrated within various wavelength ranges and the time evolution during the event. Panel (b) compares the detailed spectra at three time points as marked in the top panel, which are representative of pre-flare (black), peak-flare (red), and post-flare (green) conditions, respectively. Panel (c) shows the percentage increases of the spectral intensity at the peak- and post-flare phases relative to the pre-flare condition.



Figure 2. Comparison of the MGITM calculated CO_2 , O, CO, N₂, and Ar neutral densities with MAVEN NGIMS in-situ measurements along MAVEN pre-flare (green), near-post-flare (red), and far-post-flare (blue) orbits during the 10 September 2017 solar flare event. Figures 2a-2e present the neutral species abundances, and Figures 2f-2j present the percentage differences along the two post-flare orbits relative to the pre-flare orbit. The model results and MAVEN data are indicated by solid lines and open circles, respectively.



Figure 3. MGITM average dayside upper atmospheric perturbations, beginning from 2017-09-504 10/15:00:00, ~1 hour prior to the flare onset. Here are shown the time-varying percentage changes of the 505 dayside-averaged altitude profiles (SZA $< 90^{\circ}$) in the flare case compared with the non-flare case for (a) elec-506 tron density, (b) neutral temperature, (c) thermal pressure, (d) CO₂ density, (e) O density, (f) CO density, (g) 507 N_2 density, and (h) number density ratio of O to CO_2 . Figure 3i shows the altitude difference in units of km 508 between the pressure levels in the two cases. Note that the order of pressure on the vertical axis of Figure 3i 509 has been reversed to make altitude increase from the bottom to the top of the panel. In all the panels, we use 510 green-red colors to denote positive changes and use blue for negative changes. 511



Figure 4. The top row shows the MGITM-calculated horizontal distributions of (from left to right) neutral temperature, thermal pressure, CO₂ and O number densities at 251.25 km altitude prior to the flare onset at 2017-09-10/15:00. The results are shown in MSO latitude and local time, with the subsolar point located in the panel center. The subsequent four rows show the percentage differences between the non-flare case and the flare case at four representative time points: 2017-09-10/16:15, 2017-09-10/18:45, 2017-09-11/00:00, and 2017-09-11/05:00, respectively.