



Remarks to the paper "On Lie Derivatives in Areal Spaces"

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Remarks to the paper "On Lie Derivatives in Areal Spaces"

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Abstract

In the previous paper³⁾, the author extended the notion of Lie derivative to the areal space of general type, by the aid of theories^{1),2)} mainly.

There, we treated p as m-ple element p_a^i in the fundamental function F(x, p) of the areal space $A_a^{(m)}$.

However, the theory of the areal space was essentially started from the treatise p as m-dimensional "area"-element $p^{t[m]}$, so it is desirable that the theory of Lie derivative of $A_n^{(m)}$ is rewritten from this point of view. In this paper, we try to rewrite in the above-mentioned way.

More interesting results will be developed in the forth-coming paper.

1. By the reviewer⁴⁾, the summary of the previous paper is as follows: The author considers an infinitesimal transformation of type

$$(1. 1) \bar{x}^i = x^i + \xi^i(x) dt$$

which maps a point x of a surface $V_m: x^i = x^i(u^a)^{**}$ to a point \bar{x} of a surface $\bar{V}_m: \bar{x}^i = \bar{x}^i(u^a)$. Under this transformation the m-ple element $p_a^i \equiv \partial x^i/\partial u^a$ is transformed to

$$(1. 2) \bar{p}_{\alpha}^{i} = p_{\alpha}^{i} + \xi_{i,j}^{i} p_{\alpha}^{j} dt \quad \text{with} \quad \xi_{i,j}^{i} \equiv \partial \xi^{i} / \partial x^{j}.$$

When a geometric object $\Omega(x, p)$ is transformed to $\bar{\Omega}(\bar{x}, \bar{p})$ by (1.1), the Lie derivative of Ω with respect to ξ is defined as

(1.3)
$$\mathbf{\pounds}_{\varepsilon} Q = \lim_{dt \to 0} \left\{ \bar{\Omega}(\bar{x}, \bar{p}) - \Omega(\bar{x}, \bar{p}) \right\} / dt.$$

We call a transformation (1.1) satisfying $\pounds_{\epsilon}F = 0$ an areal motion, because it does not change the area $S = \int \cdots \int F du^1 \cdots du^m$ of an m-dimensional surface in the space.

The main results of the paper are as follows.

- (A). In order that the space admits an areal motion, it is necessary and sufficient that the Lie derivative of the metric m-tensor $g_{i[m],j[m]}$ vanishes.
- (B). If the vector ξ^i in (1, 1) is transversal to p^i_α with respect to F, then the transformation (1, 1) is an areal motion.

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^{**)} Latin indices run over $1, 2, \dots, n$; Greek indices over $1, 2, \dots, m$ $(1 \le m < n)$ in section 1 and over 1, 2 in other sections.

(C). If the space $A_n^{(m)}$ admits an areal motion, then $\mathfrak{L}_{\varepsilon}C_{i,l}^{i\alpha}=0$.

If the space $A_n^{(m)}$ is of submetric class, the metric tensor g_{ij} can be introduced. Then we have:

- (D). A motion is an areal motion.
- (E). When the areal space $A_n^{(m)}$ admits a motion (1.1), then $\mathfrak{L}_{\varepsilon}\tau_{jk}^{*i}=0$ and $\mathfrak{L}_{\varepsilon}C_{j,l}^{i,a}=0$.
- 2. In this section, we take in the areal space $A_n^{(2)}$ in place of $A_n^{(m)}$ for convenience. Under the transformations (1.1) and (1.2), the bivector p^{ij} is transformed as

$$\begin{split} \overline{p}^{ij} &= 2 \, \overline{p}^i_{11} \, \overline{p}^j_{21} = 2 (p^i_{11} + \xi^i_{,|k|} p^k_{11} dt) (p^j_{21} + \xi^j_{,|k|} p^l_{21} dt) \\ &= p^{ij} + 2 \xi^{[i}_{,h} p^{[h|j]} dt + \xi^i_{,k} \xi^j_{,l} p^{ki} dt^2, \end{split}$$

therefore, the variations of x^i and p^{ij} are represented as follows;

(2. 1)
$$\delta x^i \equiv \bar{x}^i - x^i = \xi^i dt$$

$$(2. 2) \delta p^{ij} \equiv \overline{p}^{ij} - p^{ij} = 2\xi_{,h}^{[i]} p^{[h]j]} dt + \xi_{,k}^{i} \xi_{,h}^{j} p^{kl} dt^{2}.$$

If a contravariant vector $X(x^h, p^{ij})$ is transformed to $\overline{X}(\overline{x}^h, \overline{p}^{ij})$ by (1.1), then

(2.3)
$$dX^{i} = X^{i}_{,h} \delta x^{h} + X^{i}_{;kl} \delta p^{kl}$$

$$= X^{i}_{,h} \xi^{h} dt + X^{i}_{;kl} \{ 2 \xi^{lk}_{,h} p^{|h|l|} dt + \xi^{k}_{,r} \xi^{l}_{,s} p^{rs} dt^{2} \}.$$

Now, in the other hand, we interpret that (1.1) is an infinitesimal coordinate transformation, then

(2.4)
$$\frac{\partial \bar{x}^i}{\partial x^j} = \delta^i_j + \xi^i_{,j} dt, \qquad \frac{\partial x^i}{\partial \bar{x}^j} = \delta^i_j - \xi^i_{,j} dt, \cdots$$

neglecting higher order terms with respect to dt.

If the contravariant vector X^i is transformed by the coordinate transformation (1.1), then

(2.5)
$$\overline{X}^{i} = \frac{\partial \overline{x}^{i}}{\partial x^{j}} X^{j} = (\delta^{i}_{j} + \xi^{i}_{,j} dt) X^{j} = X^{i} + \xi^{i}_{,j} X^{j} dt ,$$

$$dX^{i} \equiv \overline{X}^{i} - X^{i} = \xi^{i}_{,j} X^{j} dt .$$

Substituting (2.3) and (2.5) into the definition (1.3), we have

$$\mathfrak{L}_{\xi} X^{i} \! = \! X^{i}_{,h} \xi^{h} \! + \! 2 X^{i}_{;kl} \xi^{\lceil h}_{,h} p^{\lceil h \rceil l \rceil} \! - \! \xi^{s}_{,h} X^{h} \, .$$

This is the new definition of Lie derivative of contravariant vector X^i with respect to ξ . For a covariant vector $Y(x^h, p^{ij})$ and for a tensor $T(x^h, p^{ij})$ of 1–1 type, we have analogously that

$$\mathfrak{L}_{\varepsilon}Y_{i} = Y_{i,h}\xi^{h} + 2y_{s;kl}\xi^{[h]}_{h}p^{[h]l]} + \xi^{h}_{s}Y_{h},$$

(2.8)
$$\mathbf{\pounds}_{\xi} T_{j}^{i} = T_{j,h}^{i} \xi^{h} + 2 T_{j;hl}^{i} \xi_{,h}^{[h]} p^{[h|l]} - \xi_{,h}^{i} T_{j}^{h} + \xi_{,j}^{h} T_{h}^{i}.$$

These expressions (2.6), (2.7) and (2.8) are new definition of Lie derivative in $A_n^{(2)}$ with use of p^{ij} .

3. In the previous paper, the author defined the Lie derivative of the contravariant vector $X(x^i, p^{ij})$ in the form;

$$\mathfrak{L}_{\varepsilon}X^{i} = X^{i}_{,h}\xi^{h} + X^{i\alpha}_{;\ell}\xi^{l}_{,h}p^{h}_{\alpha} - \xi^{l}_{,h}X^{h}.$$

The term $X_{:i}^{i\alpha}(\hat{\xi}^{l}_{h}p_{\alpha}^{h})$ is rewritten as follows;

$$\begin{split} X^{ia}_{;l}\xi^{l}_{,h}p^{h}_{\alpha} &= X^{i}_{;jk}p^{jk\alpha}_{;l}\xi^{l}_{,h}p^{h}_{\alpha} \\ &= X^{i}_{;jk}(\delta^{j}_{l}\delta^{n}_{l}p^{k}_{2} + p^{j}_{l}\delta^{k}_{l}\delta^{n}_{2} - \delta^{j}_{l}\delta^{n}_{2}p^{k}_{1} - p^{j}_{2}\delta^{k}_{2}\delta^{n}_{1})\,\xi^{l}_{,h}p^{h}_{\alpha} \\ &= X^{i}_{;jk}\Big\{\xi^{j}_{,h}(p^{h}_{1}p^{k}_{2} - p^{h}_{2}p^{k}_{1}) - \xi^{k}_{,h}(p^{h}_{1}p^{j}_{2} - p^{h}_{2}p^{j}_{1})\Big\} = 2X^{i}_{;jk}\,\xi^{l,j}_{,h}p^{|h|k|}_{1}\,. \end{split}$$

Hence, we can conclude as follows:

Lemma 1. The new definition (2.6) of Lie derivative is coinside with the definition (3.1) of the previous paper.

Moreover, we have the following:

Lemma 2. The Lie derivative defined by (2.6) and (2.7) satisfies the Leibnitz' rule, that is,

(3. 2)
$$\mathbf{\pounds}_{\varepsilon}(X^{i}Y_{i}) = (\mathbf{\pounds}_{\varepsilon}X^{i}) Y_{i} + X^{i}(\mathbf{\pounds}_{\varepsilon}Y_{i}).$$

Proof) On account of (2.8), we can see

$$\begin{split} \pounds_{\varepsilon}(X^{\varepsilon}Y_{i}) &= (X^{\varepsilon}Y_{i})_{,h} + 2(X^{\varepsilon}Y_{i})_{;kl} \xi_{,h}^{\mathbb{E}} p^{|h|l_{-}^{1}} - \xi_{,h}^{\varepsilon}(X^{h}Y_{i}) + \xi_{,\varepsilon}^{h}(X^{\varepsilon}Y_{h}) \\ &= (X_{,h}^{\varepsilon} \xi^{h} + 2X_{;kl}^{\varepsilon} \xi_{,h}^{\mathbb{E}} p^{|h|l_{-}^{1}} - \xi_{,h}^{\varepsilon}X^{h}) Y_{\varepsilon} \\ &\quad + X^{\varepsilon} (Y_{i:k} \xi^{k} + 2Y_{i:kl} \xi_{,h}^{\mathbb{E}} p^{|h|l_{-}^{1}} + \xi_{,t}^{h}Y_{h}) \,. \end{split}$$

On the end of this section, we consider two infinitesimal transformation;

$$\bar{x}^i = x^i + \hat{\xi}^i(x) dt$$
, $\bar{x}^i = x^i + \eta^i(x) dt$.

If we operate $\mathfrak{L}_{\varepsilon}$ and \mathfrak{L}_{η} successively, then, after somewhat complicated calculations,

$$\begin{split} \mathbf{\pounds}_{\boldsymbol{\eta}} \mathbf{\pounds}_{\boldsymbol{\xi}} X^i - \mathbf{\pounds}_{\boldsymbol{\xi}} \mathbf{\pounds}_{\boldsymbol{\eta}} X^i &= X^i_{,h} (\boldsymbol{\xi}^h_{,j} \cdot \boldsymbol{\eta}^j - \boldsymbol{\xi} \boldsymbol{\eta}^h_{,j}) \\ &+ 2 X^i_{,rs} (\boldsymbol{\xi}^{[r}_{,h,j} \boldsymbol{\eta}^{[j]} + \boldsymbol{\xi}^{[r}_{,j} \cdot \boldsymbol{\eta}^{[j]}_{,h} - \boldsymbol{\xi}^{[j]}_{,h} \boldsymbol{\eta}^{[r}_{,j} - \boldsymbol{\xi}^{[j]} \boldsymbol{\eta}^{[r}_{,h,j}) \, \boldsymbol{\mathcal{P}}^{[h|s]} \\ &- (\boldsymbol{\xi}^i_{,h,j} \boldsymbol{\eta}^j + \boldsymbol{\xi}^i_{,j} \boldsymbol{\eta}^j_{,h} - \boldsymbol{\xi}^j_{,h} \boldsymbol{\eta}^i_{,j} - \boldsymbol{\xi}^j \boldsymbol{\eta}^i_{,h,j}) \, \boldsymbol{X}^h \;, \end{split}$$

hence,

$$(3.3) \qquad \mathbf{\pounds}_{r}\mathbf{\pounds}_{\bar{r}}X^{i} - \mathbf{\pounds}_{r}\mathbf{\pounds}_{\bar{r}}X^{i} = \mathbf{\pounds}_{r}X^{i},$$

(3.4)
$$\pounds_{\zeta} X^{i} = X^{i}_{,h} \zeta^{h} + 2X^{i}_{;rs} \zeta^{[r}_{,h} p^{|h|s]} - \zeta^{i}_{,h} X^{h},$$

where we put

$$\zeta^i \equiv \hat{\xi}^i_{,j} \eta^j - \hat{\xi}^j \eta^i_{,j} = \mathbf{\pounds}_{\scriptscriptstyle{\eta}} \xi^i = - \mathbf{\pounds}_{\scriptscriptstyle{\xi}} \eta^i$$

These facts tell us the following:

Theorem 1. If (1, 1) belongs to a transformation group, that is, ξ^i , η^i , ζ^i , \cdots are elements of an r-parameter group of transformation, then \mathfrak{L} 's in (2, 6) are r infinitesimal operators of an r-parameter group of transformations and (3, 3) with (3, 4) holds good.

4. A. Kawaguchi and Y. Katsurada¹⁾ defined a line-metric connection in the areal space $A_n^{(2)}$ in the form;

$$(4.1) DX^i = X^i_{/h} dx^h + X^i_{/hl} \omega^{hl},$$

where

$$(4. 2) X_{h}^{i} = X_{h}^{i} - X_{ikl}^{i} B_{h}^{kl} + \Gamma_{jh}^{*i} X^{j},$$

$$(4.3) X_{\ell\hbar}^i \equiv FX_{,kl}^i + FC_{j,kl}^i X^j,$$

and

$$B_h^{kl} \equiv 4p_{\uparrow 1}^{\lceil k} B_{z \downarrow h}^{l \rceil}, \qquad B_{\alpha h}^{l} \equiv \Gamma_{jk}^{*l} p_{\alpha}^{j}.$$

Sice the transformation vector ξ^i depends only on position x, so

(4.4)
$$\xi_{/h}^{i} = \xi_{,h}^{i} - \Gamma_{jh}^{*i} \xi^{j}.$$

From (4.2) and (4.4), the expression of defintion (2.6) is rewritten such that

$$\begin{split} \mathbf{\pounds}_{\varepsilon} X^{i} &= (X^{i}_{/h} + X^{i}_{;hl} B^{kl}_{h} - \Gamma^{*i}_{jh} X^{j}) \, \xi^{h} \\ &+ 2 X^{i}_{;kl} (\xi^{i}_{/h} - \Gamma^{*i}_{jh} \xi^{[j]}) \, p^{[h|l]} - (\xi^{i}_{/h} - \Gamma^{*i}_{jh} \xi^{j}) \, X^{h} \, , \end{split}$$

and by means of

$$(4.5) B_h^{kl} = 4p_{1h}^{\lceil k} B_{2h}^{l \rceil} = p_1^k B_{2h}^l - p_1^l B_{2h}^k + p_2^l B_{1h}^k - p_2^k B_{1h}^l = 2\Gamma_{hh}^{*lk} p_1^{|h|l \rceil},$$

we obtain finally

$$\mathfrak{L}_{\varepsilon}X^{i} = X^{i}_{lh}\xi^{h} + 2X^{i}_{:hl}\xi^{h}_{lh}\xi^{h} - \xi^{i}_{lh}X^{h}.$$

Now, we apply the Lie derivative (4.5) to the fundamental function

$$\mathfrak{L}_{\varepsilon} F = F_{/\hbar} - F_{:rs} B_{\hbar}^{rs}.$$

If we recall the relation $F_{rs}=2F_{rs}^{1}p_{s}^{2}$ and (4.5), we can easily see that

$$F_{rs}B_h^{rs}=F_{rl}^{\alpha}B_{ah}^l$$
.

Accordingly, we can say that the expression of the Lie derivative (4.7) of the fundamental functions is coinside with that of the previous paper.

Between the metric bitensor $g_{ij,kl}$ and F, there is a relation

$$g_{ij,kl}p^{ij}p^{kl}=4F^2.$$

Differentiating both sides of this relation by p^{ij} , we have

$$g_{ij,kl;rs}p^{ij}p^{kl} + g_{ij,kl}\delta_{rs}^{ij}p^{kl} + g_{ij,kl}\delta_{rs}^{ij}\rho^{kl} + g_{ij,kl}p^{ij}\delta_{rs}^{kl} = 8FF_{:rs}$$

and making use of $g_{ij,kl;rs}p^{kl}=0$, $g_{ij,kl}p^{kl}=2G_{ij}$ and $g_{ij,kl}p^{ij}=g_{kl,ij}p^{ij}=2G_{kl}$, and putting $F_{/h}=0$, finally we have

(4.8)
$$\mathfrak{L}_{\xi} F = \frac{1}{2F} G_{rs} \xi_{/h}^{rr} p^{|h|s|}.$$

In view of (4.8), we can conclude the following:

Theorem 2. If the transversal bivector G_{ij} vanishes, then the space $A_n^{(2)}$ admits an areal motion.

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