# The Maximum Clique Problem in Spinorial Form 

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#### Abstract

We present a formulation of the Maximum Clique problem of a graph as a geometrical problem of null vectors in complex space. This problem is then translated in spinorial language.


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## THE MAXIMUM CLIQUE PROBLEM

Given a graph of size $n$ and adjacency matrix $A$, a clique is a subgraph with pairwise adjacent vertices and the Maximum Clique (MC) problem asks for the size $\omega(A)$ of the largest clique. It is an old, well studied, NPcomplete problem and there are reviews with hundreds references [1].

Any symmetric matrix like $A$ may be expressed in the form

$$
\begin{equation*}
A=B^{\prime} B=B B=B^{2} \tag{1}
\end{equation*}
$$

where $B$ is a complex, symmetric matrix that we can think as formed by $n$ complex vectors $\mathbf{z}_{i} \in \mathbb{C}^{n}$ called null vectors since they have zero length: $\mathbf{z}_{j}^{2}=a_{i j}=0$.

It is well known [1] that the MC problem of $A$ is identical to the problem of the maximum independent set of the conjugate graph $\bar{A}=J-\mathbb{1}-A$ and a formulation with appealing properties is

$$
\begin{equation*}
\max _{\mathbf{x} \in\left\{\{0,1\}^{n}: \mathbf{x}^{\prime} \bar{A} \mathbf{x}=0\right\}} \mathbf{x}^{\prime} \mathbf{x}=\omega(A) . \tag{2}
\end{equation*}
$$

This problem has the following geometrical interpretation: the null vectors $\overline{\mathbf{z}}$ span the space $V \subseteq \mathbb{C}^{n}$ and any couple of linearly independent vectors $\overline{\mathbf{z}}_{j}$ and $\overline{\mathbf{z}}_{k}$ span a two-dimensional space contained in $V$. If $\overline{\mathbf{z}}_{j}^{\prime} \overline{\mathbf{z}}_{k}=\bar{a}_{j k}=0$ it is easy to verify that this space has the property that all of its elements are null vectors and are all mutually orthogonal: this space is called a Totally Null Plane (TNP). If $\bar{A}$ contains at least one nondiagonal zero element, then $V$ contains at least one two dimensional TNP.

The solution of the MC problem provides the largest subset $\overline{\mathbf{z}}_{j_{1}}, \overline{\mathbf{z}}_{j_{2}}, \ldots, \overline{\mathbf{z}}_{j_{k}}$ of $\overline{\mathbf{z}}$ 's which define a TNP contained in $V$; [2].

## SPINORS, WITT AND FOCK BASIS

Cartan was the first to show [3] that the geometry of null vectors can be treated elegantly with spinors; we
will follow this road. A spinor $\Phi$ is a vector belonging to the spaces $S$ of endomorphism of $C l(2 n)=E n d S$ and is defined by the Cartan's equation:

$$
\begin{equation*}
v \Phi=\left(\sum_{j=1}^{2 n} v_{j} \gamma_{j}\right) \Phi=0 \tag{3}
\end{equation*}
$$

where $v \Phi$ is a Clifford product $v \Phi=v \_\Phi+v \wedge \Phi$.
Let us define the null, or Witt, basis of $C l(2 n)$ :

$$
\begin{equation*}
p_{j}=\frac{1}{2}\left(\gamma_{2 j-1}+i \gamma_{2 j}\right) \quad \text { and } \quad q_{j}=\frac{1}{2}\left(\gamma_{2 j-1}-i \gamma_{2 j}\right) \tag{4}
\end{equation*}
$$

for $j=1,2, \ldots, n$ with the properties

$$
\left[p_{j}, p_{k}\right]_{+}=\left[q_{j}, q_{k}\right]_{+}=0 \quad \text { and } \quad\left[p_{j}, q_{k}\right]_{+}=\delta_{j k} \mathbb{1} .
$$

With this basis $\mathbb{C}^{2 n}$ is easily seen as the direct sum of two maximal TNP $P$ and $Q$ spanned by null vectors $\left\{\mathbf{p}_{j}\right\}$ and $\left\{\mathbf{q}_{j}\right\}$ respectively:

$$
\mathbb{C}^{2 n}=P \oplus Q
$$

since $P \cap Q=\emptyset$ each vector $\mathbf{v} \in \mathbb{C}^{2 n}$ may be expressed in the form $\mathbf{v}=\sum_{i=1}^{n}\left(\alpha_{i} \mathbf{p}_{i}+\beta_{i} \mathbf{q}_{i}\right)$ with $\alpha_{i}, \beta_{i} \in \mathbb{C}$.

A spinor $\Phi \in S$, defined by Cartan equation (3), may be represented by Minimal Left Ideals (MLI) of $\mathrm{Cl}(2 n)$ [4]. Consider the $2^{n}$ MLI that form the Fock basis in spinor space [5]

$$
\begin{align*}
& \omega_{0}=p_{1} p_{2} \ldots p_{n} \\
& \omega_{1}=q_{1} \omega_{0}, \quad \omega_{2}=q_{2} \omega_{0}, \ldots, \quad \omega_{2^{n-1}}=q_{n} \omega_{0} \\
& \omega_{3}=q_{1} q_{2} \omega_{0}, \quad \omega_{5}=q_{1} q_{3} \omega_{0}, \quad \ldots  \tag{5}\\
& \ldots \ldots \\
& \omega_{2^{n}-1}=q_{1} q_{2} \ldots q_{n} \omega_{0}
\end{align*}
$$

When we write the Cartan equation (3) in the basis, defined in (4) and (5), we get

$$
\begin{equation*}
v \Phi=\left(\sum_{i=1}^{n} \alpha_{i} p_{i}+\beta_{i} q_{i}\right)\left(\sum_{s=0}^{2^{n}-1} \xi_{s} \omega_{s}\right)=0 \tag{6}
\end{equation*}
$$

and this equation can be read in two ways depending on wether $v$ or $\Phi$ plays the role of the unknown. In [6] we have proved the
Proposition 1 Given $v:=\operatorname{Span}\left(x_{1}, \ldots, x_{k}\right)$, there exists a spinor $\Phi$, satisfying the Cartan equation ( 6 ) if, and only if, v is a TNP.

## MC PROBLEM IN SPINOR LANGUAGE

We are now ready to give a spinorial formulation of the MC problem and start by introducing new vectors $\overline{\mathbf{z}}_{i}$ in the Witt basis of $C l(2 n)$ :

$$
\overline{\mathbf{z}}_{i}=\mathbf{p}_{i}+\sum_{j=1}^{n} \bar{a}_{i j} \mathbf{q}_{j} \quad i=1,2, \ldots, n
$$

These vectors have the properties (immediate to prove):

- belong to $\mathbb{C}^{2 n}$ and are linearly independent;
- are null, i.e. $\overline{\mathbf{z}}_{i}^{\prime} \overline{\mathbf{z}}_{i}=0$ since $\bar{a}_{i i}=0$;
- satisfy $\overline{\mathbf{z}}_{i}^{\prime} \overline{\mathbf{z}}_{j}=\bar{a}_{i j}$.

To fully exploit the spinorial formulation we will consider the $\overline{\mathbf{z}}_{i}$ vectors as representative of the subspace they induce i.e. $\operatorname{Span}\left(\mathbf{p}_{i}, \bar{a}_{i 1} \mathbf{q}_{1}, \ldots, \bar{a}_{i n} \mathbf{q}_{n}\right)$ of dimension $\sum_{j=1}^{n} \bar{a}_{i j}+1$. We do this introducing in the definition arbitrary coefficients $\alpha$

$$
\begin{equation*}
\overline{\mathbf{z}}_{i}=\alpha_{i} \mathbf{p}_{i}+\sum_{j=1}^{n} \bar{a}_{i j} \alpha_{j} \mathbf{q}_{j} \quad i=1,2, \ldots, n \tag{7}
\end{equation*}
$$

The equation $\mathbf{x}^{\prime} \bar{A} \mathbf{x}=(\bar{B} \mathbf{x})^{2}=0$, representing the constraints of the MC problem in (2), may be linearized formulating the problem in spinorial form:

$$
\begin{equation*}
\bar{B} \mathbf{x} \Phi=\left(\sum_{i=1}^{n} x_{i} \overline{\mathbf{z}}_{i}\right) \Phi=0 \tag{8}
\end{equation*}
$$

or with $\overline{\mathbf{z}}_{i}$ from (7)

$$
\begin{equation*}
\left[\sum_{i=1}^{n} x_{i}\left(\alpha_{i} \mathbf{p}_{i}+\sum_{j=1}^{n} \alpha_{j} \bar{a}_{i j} \mathbf{q}_{j}\right)\right] \Phi=0 \tag{9}
\end{equation*}
$$

of the form (6). In this equation, in general, $x_{i}$ must be interpreted as complex variables, restricted to values in $\{0,1\}$ in formulation (2) of the MC problem .

We thus have a set of $n$ vectors $\overline{\mathbf{z}}_{i}$ defining an $n$ dimensional subspace of $\mathbb{C}^{2 n}$ and we will look for the spinors $\Phi$ that satisfy (8).

We study solutions of (9) with an example: let us suppose that $\bar{a}_{12}=0$, this means that $\overline{\mathbf{z}}_{1}$ and $\overline{\mathbf{z}}_{2}$ form a TNP and, setting $x_{3}=x_{4}=\ldots=x_{n}=0$, with (7) we get

$$
\left(x_{1} \overline{\mathbf{z}}_{1}+x_{2} \overline{\mathbf{z}}_{2}\right) \Phi\left(\overline{\mathbf{z}}_{1} \overline{\mathbf{z}}_{2}\right)=0
$$

where $\Phi\left(\overline{\mathbf{z}}_{1} \overline{\mathbf{z}}_{2}\right)$ represent the TNP $\operatorname{Span}\left(\overline{\mathbf{z}}_{1}, \overline{\mathbf{z}}_{2}\right)$. This example shows that it is simple to get particular solutions to (9) the real problem being to find the set of all solutions, for which in [6] we proved the following

Proposition 2 The set of nonzero spinors that solve the Cartan equation (9) is isomorphic to the set of cliques of A.

We can now reformulate our initial MC problem (2): it will correspond to that solution of (9) with the maximum intersection with $P$, i.e.

$$
\omega(A)=\max _{\Phi:\left(\sum_{i=1}^{n} x_{i} \overline{\bar{z}}_{i}\right) \Phi=0} \operatorname{dim}(P \cap M(\Phi))
$$

where $M(\Phi)$ represent the TNP corresponding to $\Phi$.

## CONCLUSIONS

We have thus shown that:

- the MC problem can be formulated in spinorial language;
- the problem of finding all possible spinor solutions of (9) is NP-complete since, given the set of all solutions, one gets also the solution of the MC problem.
- With respect to the MC formulation (2) we remark two main differences:
- the demanding restriction $\mathbf{x} \in\{0,1\}^{n}$ can be relaxed since all solutions of (9) necessarily have binary $x_{i}$.
- The quadratic constraint $\mathbf{x}^{\prime} \bar{A} \mathbf{x}=0$ of (2) is linearized here to $\bar{B} \mathbf{x} \Phi=0$.


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