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β -diketonate Ancillary Ligands in Heteroleptic Iridium Complexes: a balance between synthetic advantages and photophysical troubles

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β-diketones are an important class of bidentate cyclometalating compounds, used in organometallic chemistry as ancillary ligands because of their wide commercial availability and the easy synthesis. They are employed to finely tune the electronic, spectroscopic and physical properties of metal complexes. Heteroleptic iridium complexes often advantage from the use of β-diketonates ligands, being their properties similar to those of the corresponding homoleptic triscyclometalated. Nevertheless, in some cases their use results in a complete quenching of the phosphorescence. Aiming to understand the origin of this drawback, we designed a suitable class of heteroleptic complexes and studied their thermal stability (DSC/TGA). We explored the effect of the ancillary ligand in a series of Ir(III) complexes bearing 2-phenylpyridine (*ppy*) as cyclometalated ligand and *acac* (acetylacetonate), *tta* (2-thienoyltrifluoroacetonate), *dtdk* (1,3-di(thiophen-2-yl)propane-1,3-dionate) and *BPhen* (4,7-Diphenyl-1,10-phenanthroline) as ancillary ligands. Through photochemical and electrochemical investigations, whose results agree and support our density functional theory calculations, we demonstrate that β-diketonate ligand with low triplet energy generates dark triplet excited states with negligible coupling to the ground state which indeed promote non-radiative relaxation through population of higher states.

archetypal of the phosphorescent emitters for OLEDs, with

impressive green emission efficiency, approaching unity. The frontier orbitals in **Ir3** are localized on the pyridyl π -orbitals in

the LUMO, whereas the phenyl group and the iridium d-

orbitals mostly contribute to the HOMO.³ Hence, through

careful engineering of a C^N ligand structure, the

photophysical and electrochemical properties of the complex

can be controlled to achieve emission from UV to near-IR.⁴⁻⁶

The replacement of one of the C^N moieties with another

cyclometalated "ancillary" ligand, yielding heteroleptic

complexes with the optimal features of the homoleptic

 $Ir(C^N)_3$, provides the tool to fine control the color emission by

selective stabilization/destabilization of the HOMO and LUMO

levels. When chromophoric ancillary ligands are employed,

additionally, the photophysics of the resulting molecule could

be characterized by interligand energy transfer, in particular

when the ligand triplets are close in energy.⁷ Typical examples

most popular and used. Many β-diketones are low cost,

commercially available and the synthesis of the corresponding

complexes is relatively easy. In addition, either their

introduction produces no change in the luminescence color or

they can induce a slight emission red-shift.^{6,8} Functionalized

1,3-β-diketones have been widely used in lanthanide

chemistry¹³ and in recent years they demonstrated their

potential as highly tunable conjugate scaffolds for the

synthesis of luminescent materials.^{14,15}

of ancillary ligands are β -diketone,⁸ picolinic acid,⁹ salicylimine^{10,11} and pyrazolyl-borate¹² derivatives, to mention a few; among them, β -diketonate (O^O) ligands stand as the

Introduction

Tris-cyclometalated iridium complexes $(Ir(C^N)_3, where C^N)$ is a monoanionic ligand) have drawn great attention in the last decades, in particular due their excellent performances as phosphorescent emitters in optoelectronic applications, such as organic light emitting diodes (OLEDs).^{1,2} In these complexes, the intense phosphorescence arises from a triplet ligand centered (³LC) excited state with the admixture of high-lying metal-to-ligand charge transfer state (MLCT). The strong spinorbit coupling (SOC) induced by the iridium center favors the radiative emission from the triplet state to the singlet ground state, formally forbidden. Moreover, SOC promotes a high rate of intersystem crossing (ISC) from singlet excited states to the emitting triplet, thus allowing the harvesting of both singlet and triplet electro-generated excitons in OLEDs.

Fac-tris(2-phenylpyridinato-N^C2) Iridium(III) (Ir3) is the

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In particular, upon coordination to B(III) Lewis acids like in diaroyl(methanato)boron difluoride systems, their photophysical behaviour have been investigated in both solution and solid state.^{13,16,17} Acetylacetone (Hacac) is the prototype of a non chromophoric β -diketonate ligand with the highest triplet energy.¹⁸ Alkyl, fluoroalkyl, aromatic and/or heteroaromatic groups have been introduced as substituents on the carbonyl functions. All these functional groups affect different properties of the final β -diketonate complex; for instance, branched alkyl groups serve to enhance solubility and volatility, while aromatic substituents enhance light absorption but also influence the position of the energy levels of the ligand (*i.e.* singlet and triplet states). β -diketonates with heteroaromatic electron rich groups, such as thiophenes, were also proposed to be useful in certain optoelectronic applications, such as OLEDs, possibly enhancing the charge trapping ability of a complex or to provide sites (e.g. thiophene position) for a further modification.

In a recent work,¹⁹ we compared a series of NIR emissive heteroleptic complexes having β-diketonate ancillary ligands conjugation: with increasing acac. tta (2-(1,3-di(thiophen-2thienoyltrifluoroacetonate) and dtdk yl)propane-1,3-dionate). Although all the compounds have comparable photophysical and electrochemical properties, the phosphorescence efficiency was lower in the case of the tta derivative and that was attributed to a low-lying CT quenching state that can be thermally accessible from the lowest emissive state. Even worst could be the situation of some known compounds, like the Ir(ppy)2dbm (dbm, 1,3diphenylpropane-1,3-dione) where a complete quenching of the luminescence is observed.⁸

In this work we explore a series of complexes of formula $Ir(ppy)_{2}L$ (Chart 1) and carrying phenyl-pyridine (*ppy*) as the cyclometalated ligand, where L is a monoanionic or neutral bidentate ligand: *ppy*, giving the homoleptic **Ir3**; *acac*, *tta*, *dtdk* as β -diketonate ligands of different conjugation size giving **Ir2acac**, **Ir2tta** and **Ir2dtdk**, and *BPhen* (4,7-Diphenyl-1,10-phenanthroline) giving the cationic complex **Ir2BPhen**. *tta* and *dtdk* were selected according to their triplet energies

compared to that of the *ppy* counterpart. We observed that the emission color and efficiency drastically change along this series: while **Ir3** and **Ir2acac** are highly efficient green emitters, a complete quenching of the emission occurs in **Ir2tta** and **Ir2dtdk**; **Ir2BPhen**, which exhibits optical and electronic features similar to the latter two complexes, displays a quite efficient orange emission. We will compare the electrochemical and optical properties of the complexes and, under the guidance of the theoretical modelling, we will address the origin of the quenched emission in *dtdk* and *tta* derivatives.

Results

Redox Properties

Electrochemical studies were performed to identify the nature and localization of the redox centers and gain informations on the frontier electronic levels of the complexes. Selected results, collected by cyclic voltammetry on a glassy carbon (GC) electrode in DMF solution with a 0.1 M TBAP supporting electrolyte, are reported in Table 1, while Figure 1 displays a synopsis of the CV patterns recorded at 0.1 V s⁻¹. The data of **Ir3** are taken from the literature and show reversibility both in the oxidation and reduction.⁵ Hacac and Hppy ligands have redox potentials that fall outside the solvent limits and thus would not be considered.

In all the β -diketonate complexes, a single, chemically and electrochemically reversible oxidation wave is detected above 0.4 V vs Fc⁺/Fc. It is attributed, in analogy to previous data,^{3,4} to a predominantly metal-centered process involving the cyclometalated phenyl-pyridine moiety,^{20,21} as also evidenced by the DFT calculations. Similarly, to our previous study,¹⁹ we observe a progressive shift of the complex oxidations (*i.e.* an HOMO level stabilization) towards more positive potentials following the order *acac* \leq *dtdk* < *tta*. Although it should be considered that the negative charge on the *tta* and *dtdk* could be stabilized by their larger π framework compared to the parent *acac*, thus resulting less available, the slight HOMO

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stabilization (20 mV) occurring in the **Ir2dtdk**, compared to **Ir2acac**, could be more likely be attributed to the π -back-donation from the Ir(III) center to the *dtdk* ligand (again due to its extended conjugation).²² Instead, the stronger positive shift (150 mV) of **Ir2tta** results from the combined effects of π -back-donation (in analogy with **Ir2dtdk**) and the electron-attracting effect of the CF₃ group on the *tta* moiety.

Ir2BPhen is the only ionic complex of the group. In this case, the oxidation potential undergoes a further positive shift accounting for the net positive charge centered on the iridium, that favors the π -donation from the N atoms of *BPhen* to Ir(III) while weakening the π -back-donation from the metal to *BPhen*. The complex oxidation, close to the solvent limit, is irreversible, although Dragonetti *et al.* reported the same process to be more reversible in acetonitrile solution and a return peak can be perceived increasing the scan rates.²¹

Within the solvent limits, no clear reduction waves are perceivable in the case of Ir2acac although a quite broad signal, on the solvent edge with onset around -2.28 V is observable. This would point to a peak close to -2.6 V as reported in previous works⁸ and attributable to a one-electron reduction processes localized on the phenyl-pyridine ligand. Ir2dtdk and Ir2tta show a clear electrochemical and chemical reversible reduction at -2.16 and -1.99 V, respectively. The reduction cannot be reasonably attributed to a ppy based orbital (as in Ir3 and Ir2acac) since the positive shift of 0.5 V cannot be explained with the effect of the ancillary ligand. Therefore, given that the free β -diketone are characterized by a reduction wave between -1.88 and -1.55 V (Table 1), the reduction process in Ir2dtdk and Ir2tta can be safely attributed to orbitals localized on the β -diketonate themselves. It is interesting to note that the reversibility is largely improved compared to the free ligands²³ and that upon coordination to the iridium center the reduction potential of dtdk undergoes a smaller negative shift (0.32 V) compared to the 0.48 V observed in tta. The different shifts provide a further inside on the donation and back-donation processes and kindly match with the interpretation that tta ligand favors a stronger back donation from the metal center.



Figure 1. Cyclic voltammograms of the complexes recorded versus Fc^*/Fc in DMF at 298 K under a N₂ atmosphere (scan rate = 100 mV s⁻¹).

Finally, **Ir2BPhen** displays two clear reversible reductions at -1.67 and -2.27 V vs $Fc^+|Fc$. The less negative process is ascribed to orbitals localized onto the *BPhen* ligand as already discussed in the literature;²¹ the second reductive event is presently difficult to assign. In fact, the uncoordinated *BPhen* ligand is characterized by two *quasi*-reversible and nearly merging reductions (see Fig. SI8).²⁴ Once it is cyclometalated to iridium, it acts as neutral N donor compound, therefore its electron poor character allows for easier reduction compared to the uncoordinated system. Hence, a one electron reduction in **Ir2BPhen** removes the degeneracy. Therefore, the peak at -2.27 V may belong either to a second reduction on the ancillary *BPhen* ligand or it can be ascribed to a *ppy* based process (considering that, being the oxidation at 0.93 V, this potential is consistent with an optical band gap of 2.5 eV).

DFT/TDDFT Calculations

The molecular orbitals (MOs) obtained at the DFT-B3LYP level are reported in Figures 2 and SI9. The visual inspection shows that the HOMO are localized mainly on the *ppy* moiety. The LUMO of **Ir3** and **Ir2acac** are also localized on the *ppy*, whereas the LUMO of **Ir2dtdk** and **Ir2tta** are on the ancillary ligand, with a modest contribution from the metal d-states. Interestingly, **Ir2BPhen** has a two-fold degenerate LUMO, localized on the *BPhen* moiety. Once, the LUMO get occupied by one electron, the degeneracy is lifted because of Coulomb repulsion and the once-degenerate companion MO, is pushed higher in energy.

Table 1. Electrochemical data of the con	nplexe
--	--------

	Eox	E _{ox onset}	E [E ⁰] ()()	E ()()	$E_{g \max}^{0}$	E _{g,onset}	номо	LUMO
	$[E_{ox}^{o}]$ (V)	(V)	Ered [Ered] (V)	E _{red,onset} (V)	(eV)	(eV)	(eV)	(eV)
Ir3 ⁵	//	+0.31	//	-2.70	//	3.01	-5.11	-2.10
lr2acac	+0.43' [+0.40]	+0.32	//	-2.28	//	2.60	-5.12	-2.52
lr2dtdk	+0.46' [+0.42]	+0.35	-2.20' [-2.16]	-2.09	2.58	2.44	-5.15	-2.71
lr2tta	+0.58 ^r [+0.55]	+0.47	-2.03' [-1.99]	-1.92	2.54	2.39	-5.27	-2.88
Ir2BPhen	+0.93	+0.75	-1.70 [′] [-1.67], - 2.30 [′] [-2.27]	-1.60	2.63	2.35	-5.55	-3.20
Hdtdk	+1.06	+0.85	-1.88 ⁹ , -2.17 ⁹	-1.73	//	2.58	-5.65	-3.07
Htta	//	+1.04	-1.55 ['] , -2.09 [']	-1.41	//	2.45	-5.84	-3.39
BPhen	11	//	-2.40 ^q , -2.53 ^q	-2.17	//	//	//	-2.63

 $E^{o}=1/2(E_{pa}+E_{pc});\ DMF/TBAP\ (0.1\ M),\ vs\ Fc^{+}|\ Fc\ at\ scan\ rate\ 100\ mVs^{-1};\ ^{r}\ denotes\ reversible;\ ^{i}\ denotes\ irreversible\ and\ ^{q}\ denotes\ quasi-reversible\ peaks.\ E_{g}^{0}_{,max}=(E_{ox}^{0}^{o}-E_{red}^{o}).\ E_{g,onset}=(E_{ox,onset}-E_{red,onset}).\ The\ HOMO\ energy\ level\ was\ calculated\ using\ the\ equation\ E_{HOMO}\ (eV)=-(E_{red,onset}+4.8)\ and\ LUMO\ energy\ level\ was\ calculated\ using\ the\ equation\ E_{LUMO}\ (eV)=-(E_{red,onset}+4.8).\ The\ bandgap\ was\ determined\ from\ the\ equation\ E_{g}\ (eV)=-(E_{HOMO}-E_{LUMO}).\ The\ onset\ potentials\ of\ oxidation\ and\ reduction\ are\ determined\ from\ intersection\ of\ the\ tangents\ drawn\ at\ the\ rising\ current\ side\ of\ the\ peak\ and\ the\ extrapolated\ baseline\ charging\ current\ of\ the\ CV.$

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Figure 2. Frontier molecular orbitals calculated at the B3LYP level.

Note that we did not include a compensating anion in the calculation of **Ir2BPhen** and, because of the net charge, all MO energy levels are lower in energy with respect to the others systems. Thus, excluding **Ir2BPhen**, the energy of the HOMOs and of the LUMOs correlate fairly well with the oxidation and reduction peaks of the CV experiments.

It should be noted however, that the molecular orbital energies are an approximation to the CV oxidation/reduction levels.²⁵ If the Koopman's condition would be fulfilled exactly by the exchange-correlation potentials, and in absence of orbital relaxation, they two would correspond exactly. It has been shown that hybrid-XC functional can fulfil the Koopman's condition quite well. The total energy as a function of the number of electrons is nearly piece-wise linear (the small deviation from linearity is the Hubbard energy, i.e. the charging/discharging energy in CV). This Hubbard energy is due to the residual electron self-interaction and is of different magnitude between occupied (HOMO) and unoccupied (LUMO) states.²⁶

The calculated absorption spectra, at the TDDFT-B3LYP level, are reported in Figure SI11 and in Tables from SI2 to SI11. The energy of the transitions and their oscillator strength are in overall good agreement with the experimental spectra.

Table 2. Calculated phosphorescence emission energies

	ΔSCF(0-0)	∆SCF(vertical)	TDA
lr3	2.5262 eV (486.9 nm)	2.2714 eV (545.8 nm)	2.2580 eV (549.1 nm)
lr2acac	2.4378 eV (508.6 nm)	2.2500 eV (551.0 nm)	2.1700 eV (571.3 nm)
lr2dtdk	2.0545 eV (603.5 nm)	1.6178 eV (766.4 nm)	1.4290 eV (867.6 nm)
lr2tta	2.0285 eV (611.2 nm)	1.4802 eV (837.6 nm)	1.2981 eV (955.1 nm)
Ir2BPhen	2.3438 eV (552.6 nm)	2.0953 eV (591.7 nm)	1.7459 eV (710.1 nm)

In Table 2 we report the calculated emission wavelengths, according to three different methods. As shown in previous publications, ^{19,27} the Δ SCF method slightly overestimates the wavelength in these systems, but is more accurate than TDDFT (in the Tamm-Dancoff Approximation). As the emission shape is temperature dependent in experiments, the Δ SCF(0-0) and Δ SCF(vert.) results provide in practice, upper and lower bounds for the emission. This is true for **Ir3**, **Ir2acac** and **Ir2BPhen**. However, the calculated emission wavelength of **Ir2dtdk** and **Ir2tta** strongly overestimates the experimental emission peak. This could mean that either the lowest T1 is not emissive, or emission arises from a different state, competitive with T1.

Photophysical Properties

The absorption spectra of the complexes in CH₂Cl₂ are reported in Figure 3. All the compounds show intense bands below 320 nm, with absorptivity exceeding 2.5x10⁴ L mol⁻¹ cm⁻¹ ¹ (Table 3). While **Ir3** displays two distinct peaks at 243 nm and 282 nm, in the heteroleptic complexes the transition slightly red-shifts along the series: Ir2acac (260 nm), Ir2tta (266 nm), Ir2dtdk (268 nm) and Ir2BPhen (272 nm). These bands in the UV region match the free ligand absorptions (See Figure SI 6 in ESI) and are attributed to the spin allowed π - π * transition (¹LC) of the organic moieties. In all the complexes, the weaker absorption at lower energy (320 - 440 nm, ε < 23000 L mol⁻¹ cm⁻¹) accounts primarily for the spin allowed d- π^* transition (¹MLCT states), although contributions from ¹LC states of the β-diketonate ancillary ligands are present in this spectral region for Ir2tta and Ir2dtdk. Finally, the broad and unstructured absorption, in the range 440 - 550 nm, with extinction coefficient below 4000 L mol⁻¹ cm⁻¹, accounts for

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Figure 3. UV/Vis absorption spectra in CH₂Cl₂ at 298 K of complexes Ir3 (black inverted triangles), Ir2acac (red squares), Ir2tta (green triangles), Ir2dtdk (blue circles) and Ir2BPhen (purple stars). SO–T1 absorption bands in the inset.

transitions of triplet character. Overall, giving attention to the absorption onset there is a progressive shift towards low energies on going from Ir3 > Ir2acac > Ir2Bphen and Ir2tta \approx Ir2dtdk

Photoluminescence experiments were performed in degassed solution of CH₂Cl₂ at room temperature and in glassy matrix of 2-MeTHF at 77 K (Figures 4 and SI7, Tables 3 and SI1). Ir3 and Ir2acac have been already extensively studied. Both are known to be highly phosphorescent green emitters (from a predominant ³MLCT state involving the occupied Ir-5d and unoccupied *ppy* π^* orbitals).³ At room temperature the two complexes emit at 515 nm and 520 nm, respectively, (Figure 3) with phosphorescence quantum yield and lifetime of 95% and 1.5 µs for Ir3, 62% and 1.6 µs for Ir2acac. The heteroleptic structure of the Ir2acac and the different geometric orientation of the ppy ligands (Chart 1) result in an increase of $k_{\rm nr}$ by one order of magnitude with respect to the homoleptic complex. The modification of the β -diketonate ligand by the introduction of one thiophene ring and one electron-drawing CF₃ group, as in Ir2tta, or the by bis-substitution with thiophenes, as in Ir2dtdk, result in an almost completely quenched room temperature phosphorescence (QY < 0.01). The emission maintains the unresolved broad feature of the parent Ir2acac complex although slightly blue shifted at 519 nm and 515 nm for Ir2tta and Ir2dtdk, respectively. Despite the low QY, the excited state lifetimes are barely reduced compared to Ir2acac, being 1.3 µs and 1.5 µs, respectively.

The spectral features suggest that the two complexes still emit from a ³MLCT state centered on *ppy* ligand, although radiative rates ($k_r = 7-8 \times 10^3 \text{ s}^{-1}$) are two order of magnitude smaller than that of **Ir2acac**, while k_{nr} constants only slightly increase.

The cationic complex **Ir2BPhen** has a complete different behavior with emission maximum at 588 nm, quantum yield of 77% and a lifetime of 1.1 μ s. It is considerably red-shifted with respect to all the other complexes. In agreement with the reported literature, its emission can be assigned to a ³MLLCT (metal-ligand-to-ligand-charge transfer) emitting state from

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the *ppy* ligand to the *BPhen* one. The broad band (FWHM = 3100 cm^{-1}) and the large Stokes shift confirm a greater spatial separation of the involved states.

In glassy matrix at 77 K, **Ir3** and **Ir2acac** exhibit a well resolved vibronic structure with maxima at 494 and 506 nm, respectively. The emission blue-shift with respect to room temperature, i.e. rigidochromic effect, accounts for 825 and 532 cm⁻¹, respectively. This effect is due the stabilization of the emitting state upon solvent reorganization at room temperature, whereas it is hampered in rigid matrix at 77 K, and it is directly related to the degree charge redistribution between excited and ground states.¹⁹ Complex **Ir2BPhen** is quite emblematic since it undergoes over 2400 cm⁻¹ emission blue-shift at 77 K with a complete change in spectral shape becoming structured and equivalent to a *ppy* based process.

At low temperature **Ir2dtdk** and **Ir2tta**, still show very low emission intensity with peculiar behaviors. Both complexes display instead a new broad and unstructured low energy band at around 630 nm with microsecond decay regime. Moreover, **Ir2tta** displays a residual high energy emission at 497 nm, absent in the case of **Ir2dtdk**.

Discussion

All complexes were synthesised in good overall yields by means of standard procedures as described in the ESI.

The iridium complexes can be used as emitters in a simple multi-layer OLED architecture (+)ITO/HTL/EML/ETL/LIF/Al(-), where HTL (hole transport layer) could be 4,4'-bis[N-(naphthyl)-N-phenyl-amino]biphenyl (NPD), ETL (electron transport layer) could be 2,9-dimethyl-4,7-diphenyl-1,10-phenanthroline (BCP) and the EML (emitting layer) could be a blend of the iridium complex doped into a host matrix.



Figure 4. Normalized emission spectra in degassed CH₂Cl₂ solution at 298 K (black line); excitation (blue line) and emission (red line) spectra in glassy 2-MeTHF at 77 K, for the complexes.

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Table 3. Photophysical data of the complexes.

	, , ,		298 К ^[b]					77 K ^[c]	
	A_{abs} (nm) (ε (10 ³ L mol ⁻¹ cm ⁻¹)) ^{<i>loj</i>}	λ _{em} (nm)	τ (μs)	$arPsi_{ extsf{L}}$	$k_{\rm r}$ (10 ⁵ s ⁻¹)	$k_{\rm nr}$ (10 ⁵ s ⁻¹)	$\lambda_{\scriptscriptstyle em}$ (nm)	τ (μs)	
lr3	243 (54.8), 282 (53.4), 340 [sh] (11.1), 378 (14.1), 405 (9.52), 455 (3.26), 485(1.45)	515	1.5	0.95	6.17	0.325	494	4.5	
lr2acac	260 (54.5), 306 (23.0) [sh], 338 (12.6), 366 (8.16), 405 (5.26), 457 (457), 492 (1.33)	520	1.6	0.62	3.76	2.30	506	5.0	
Ir2dtdk	268 (39.8), 344 (22.8), 358 (20.7) [sh], 401 (11.4), 465 (3.16)	515	1.5	< 0.01	≈0.066	≈6.6	600, 630	-	
Ir2tta	266 (41.7), 303 (26.0), 331 (20.0), 363 (14.1), 450 (3.47)	519	1.3	< 0.01	≈0.077	≈7.5	493, 645	-	
lr2BPhen	272 (57.7), 288 (51.0), 338 (17.9), 391 (10.2), 472 (1.55)	588	1.1	0.77	6.91	2.02	513	5.7	

^[a] Absorption maxima and molar extinction coefficients in CH₂Cl₂ at 298 K. ^[b] Phosphorescence maxima, lifetimes and quantum yields in degassed CH₂Cl₂ at 298 K. Radiative and non-radiative rate constants are calculated using the equations $k_r = \Phi_L/\tau$ and $k_{nr} = 1/\tau - k_r$, on the assumption $\Phi_{ISC} = 1$. ^[c] Phosphorescence maxima and lifetimes in 2-MeTHF at 77 K. k_r and k_{nr} for **Ir2tta** and **Ir2dtdk** are indicative given the too large variability in accessing Φ_L .

By considering both the energy of the frontier levels and triplet states of the dopant, 4,4'-N,N'-dicarbazol-biphenyl (CBP) could be suitable as host material since its HOMO/LUMO levels comprise the ones of the emitters, thus allowing for an efficient charge trapping. Moreover, the CBP triplet energy is higher than the those of the emitters, as necessary to guarantee the triplet harvesting.²⁸ The requirements are fulfilled for all the investigated compounds, despite the negligible emitting efficiencies of **Ir2dtdk** and **Ir2tta** make them not suitable as OLED emitters.

In order to assess the possible use of these compounds as guest emitters in vacuum deposited OLEDs, thermal studies were performed by DSC and TGA methods, under nitrogen atmosphere. Good stability is found for all the complexes, except **Ir2BPhen** that is the only one showing an exothermic degradation process. **Ir3** results as the most stable compound up to 444°C (melting, T_m) with less than 5-8 % weight loss (*w.l.*) below 350°C; **Ir2acac** shows T_m =330°C with less than 5% *w.l.* below 290°C; **Ir2tta** shares similar stability as **Ir2acac** with T_m =327°C and less than 5% *w.l.* below 325°C. The new **Ir2dtdk** shows improved stability with T_m = 355°C and less than 5% *w.l.* below 345°C. The cationic complex **Ir2BPhen** instead displays an exothermic decomposition process starting at ca. 320°C.

Analysing the frontier electronic levels of the iridium complexes and their ligands from both the DFT calculations and the electrochemical studies (in full agreement), it is evident that in all the complexes the HOMO level is characterized by a strong contribution from the metal center and the cyclometalated phenyl group of *ppy* ligand. HOMOs experience a progressive energy stabilization in the sequence *acac*, *dtdk*, *tta*; **Ir2Bphen** stands slightly out of this series

simply for its formal cationic nature leading to over 0.4 eV shift of the HOMO down in energy compared to the other complexes. The observed irreversibility might be related to the experimental conditions (reversibility is restored at high scan rate) since previous investigations indicate the oxidation of Ir2Bphen reversible.²¹ The LUMO levels in Ir3 and Ir2acac, as widely discussed in the literature, are localized on the cyclometalated C^N ligand and mainly on the pyridine portion. Contrary, in the other systems, the LUMO level shifts onto the ancillary ligands (either the β -diketonate or the phenanthroline one). DFT investigation safely confirms this picture and indicates that the LUMOs+1 are localized on ppy moiety. Ir2BPhen displays an interesting behaviour: similar to the free BPhen ligand, it has two nearly degenerate levels, while the LUMO+2 shows a ppy character, although this situation may change upon one electron reduction (ppy on LUMO+1). In Ir2BPhen, upon cyclometalation of BPhen, its LUMO drops in energy (with respect to free *BPhen*) accounting for the π -bonding to the Ir center; the two complexes Ir2tta and Ir2dtdk, display a reverse behaviour since first reduction occurs at more negative potentials compared to the free ligands.

Considering the absorption spectra of the free ligands (see Figure SI6), their singlet excited state energy decreases along the series $Hppy \approx Hacac > BPhen > Htta > Hdtdk$. Taken from low temperature measurements, their triplet energy ranks Hacac as the higher T state at 3.14 eV, followed by the triplet of Hppy (2.88 eV),²⁹ Ir(ppy)₂ fragments (2.51 eV), *BPhen* (2.48 eV),³⁰ Htta (2.47 eV)³¹ and Hdtdk (2.33 eV)³¹.

Calculations highlight that the transition to the S1 state is optically inactive in all systems except for Ir3 and Ir2acac and

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in fact no feature of their MLLCT transition into the ancillary ligand are evident in the absorption spectra (UV-Vis shows spectra rather similar one to the other); in turn excitation spectra at 77K reveals long-wavelength signal in the case of **Ir2tta** and **Ir2dtdk** with offset around 585 and 570 nm and 520 nm for **Ir2BPhen**). Moreover, the corresponding T1 state, localized on the ancillary ligand, has a non-emissive character in the case of **Ir2tta** and **Ir2dtdk**; their weak phosphorescence at low energy is observed only at 77 K. TDDFT calculation predicts a triplet emission energy for **Ir2tta** and **Ir2dtdk** in good agreement with their red phosphorescence at 77 K, although the green emission at 298 K is not explained.

In order to give an explanation of the emission in **Ir2dtdk** and **Ir2tta**, we extracted the Mulliken spin density on the Ir ion in the T1 state (Table SI12) and we examined the changes of the geometry between the S0 and the lowest T1 state (Table SI13). The spin density shows that the metal d-states contributes nearly the same in all heteroleptic compounds. This contribution is slightly smaller in **Ir2dtdk**, with respect to the other four molecules. In **Ir3**, there is charge donation from the ligands into the minority spin channel of the Ir ion, and as a consequence its spin density is reduced. Our findings point towards the fact that spin orbit (SO) effects are of the same magnitude in all the compounds and do not explain the difference in the phosphorescence colors.

Next, we extracted from our DFT calculations, the distance between Ir and its neighboring atoms, in the S0 and T1 states. Interestingly, we found that on going from the S0 to the T1 there is an average contraction of the first shell of atoms surrounding the metal. On average, this contraction is quite small (-0.005 Å) for Ir3 and Ir2acac, and larger (-0.021, -0.025 Å) for Ir2dtdk and Ir2tta, respectively. The case of Ir2BPhen is intermediate, with an average contraction of -0.015 Å. A large geometrical rearrangement reflects a large difference between the S0 and the T1 potential energy surfaces (PES). As a result, in Ir2dtdk and Ir2tta, non-radiative processes might dominate over the photoemission from the lowest triplet state.

Therefore, our DFT and TDDFT results suggest the following scenarios: the lowest singlet of **Ir2dtdk** and **Ir2tta** is of MLLCT character with a very small oscillator strength. Moreover, the singlet- triplet splitting is very small (~0.03 eV) and comparable with thermal energy at room temperature. Considering that the T1 state is not well vibronically coupled to the ground state, it is conceivable that the triplet is depopulated in favor of other higher T states by reverse internal conversion (rIC) or S states by reverse ISC; states which in turn emits to the S0, suggesting a possible thermally activated process (TAP).

On the contrary, the singlet-triplet splitting of **Ir3** and **Ir2acac** is one order of magnitude larger and these compounds show conventional phosphorescence. The case of **Ir2BPhen** is intermediate and, since the *BPhen* moiety is more rigid, it is conceivable that the T1 could have two competing pathways for emission: via vibronic coupling to the S0 and by TAP.

Conclusion

β-diketonate ancillary ligands are used in the chemistry of heteroleptic iridium complexes as a tool to finely tune the electronic, spectroscopic and physical properties of the resulting compounds. Nevertheless, in some cases their introduction may ultimately lead to a partial or complete quench of the emission. For this reason, we investigated a series of heteroleptic complexes carrying β-diketonate ancillary ligands with progressively lower triplet energy. By studying their electrochemical behavior, we determined the nature and localization of the frontier energy revealing that although the HOMO energy is sensitive to the ancillary ligands structure, it always locates on the Ir center and the phenyl group of ppy. On the contrary, the LUMO level shifts from the ppy ligand to the ancillary one in those complexes with quenched emission. Photophysics reveals a quite peculiar emissive behavior with a strong color change between room temperature and 77 K. Spectroscopy, matched to the electrochemical studies, agrees and grants support to the theoretical framework developed. TDDFT calculations revealed a complex situation when the β -diketonate ligand represent the lowest triplet state of the complex. The corresponding HOMO-LUMO transition from the ground state has negligible oscillator strength because of the poor wavefunction overlap. Moreover, the strong charge transfer character of the transition shrinks the singlet-triplet splitting to the order of the thermal energy at 298 K, promoting the back energy transfer to the higher states, as the main deactivation channel for the non-emissive triplet. This phenomenon looks to be more general than the example herein described and applies also to literature cases, 7,32,33,34,35 and it should be considered in the situation where one of the frontier orbital (i.e. LUMO), fully shifts onto the ancillary ligand.

Experimental

Supporting Information.

Reagents and solvents were purchased from Aldrich and TCI and used without further purification. All solvents, unless otherwise stated are degassed under argon prior the use. The ligand H*dtdk* was previously synthesized by some of us,³¹ while H*tta* and H*dpm*, commercially available, were used as received. NMR were recorded on Bruker Advance 400 MHz and the residual signal of the deuterated solvent was used as internal standard (e.g. 5.32 ppm for ¹H and 53.84 ppm for ¹³C of CD₂Cl₂). The chemical shifts are given in ppm and coupling constants in Hz. Ir3³⁶, Ir2acac³⁷, Ir2tta²³ and Ir2BPhenPF₆²¹ were synthesized according to known synthetic procedures. Ir2dtdk was synthesized in this work and details are given in

Complex **Ir2dtdk**: (400 MHz, CD_2CI_2) δ = 8.62 (d, *J*=5.8 Hz, 2H), 7.90 (d, *J*=8.1Hz, 2H), 7.75 (t, *J*=7.9Hz, 2H), 7.62 (d, *J*=7.8Hz, 2H), 7.60 (d, *J*=3.8Hz, 2H), 7.43 (d, *J*=4.7Hz, 2H), 7.13 (t, *J*=7.8Hz, 2H), 7.02 (dd, *J*=4.7Hz, 1H), 7.01 (d, *J*=3.8Hz, 1H), 6.88 (t, *J*=7.4Hz, 2H), 6.74 (t, *J*=7.4, 2H), 6.47 (s, 1H), 6.29 (t, *J*=7.6Hz, 2H).

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¹³C-NMR (100 MHz, CD₂Cl₂) δ = 172.9, 168.1, 147.7, 147.4, 145.7 (Cq), 148.8, 137.9, 133.6, 129.8, 129.4, 128.5, 127.6, 124.4, 122.5, 121.4, 119.1, 94.5 (CH) $LC-MS = 735.8 [M]^{+}, M^{2}(735.8) = 500.9 [M-dtdk]^{+}$ Anal. Calc. for C₃₃H₂₃IrN₂O₂S₂ : C, 53.86; H, 3.15; Ir, 26.12; N, 3.81; O, 4.35; S, 8.71; Found: C, 53.05; H, 3.3; N, 3.63.

Electrochemical measurements were performed using potentiostat-galvanostat AUTOLAB PGSTAT of EcoChemie (Utrecht, The Netherlands), equipped with software GPES version 4.9 for the processing of experimental data. Cyclic voltammetry patterns were recorded on a wide range of scan rates on a glassy carbon (GC) working electrode (Metrohm, 3 mm diameter), using a platinum counter electrode and an aqueous saturated calomel electrode (SCE) inserted in a separated compartment filled with the working medium and ending with a porous frit to avoid significant leakage of water and chloride anions. The experiments were carried out in dimethylformamide (DMF) at 3×10^{-3} M concentration of the complex in the presence of tetrabutylammonium perchlorate salt (0.1 M) as supporting electrolyte, removing oxygen by nitrogen bubbling. The potentials were subsequently referred to the potential of the Fc⁺/Fc couple recorded in the same conditions. Ohmic drop correction was applied by the positive feedback method.

UV/Vis absorption spectra were obtained on Agilent 8453 spectrophotometer in 1 cm path length quartz cell in 2x10⁻⁵ M CH₂Cl₂ solution. Photoluminescence quantum yields were measured with a C11347 Quantaurus - QY Absolute Photoluminescence Quantum Yield Spectrometer (Hamamatsu Photonics U.K), equipped with a 150 W Xenon lamp, an integrating sphere and a multi-channel detector.

Steady state emission and excitation spectra and photoluminescence lifetimes were obtained with a FLS 980 spectrofluorimeter (Edinburgh Instrument Ltd.). Continuous excitation for the steady state measurements was provided by a 450 W Xenon arc lamp. Photoluminescence lifetime measurements, determined by TCSPC (time-correlated singlephoton counting) method, were performed using an Edinburgh Picosecond Pulsed Diode Lasers EPL-445 and EPL-375 (Edinburgh Instrument Ltd.), with central wavelength 442.2 nm and 374.0 respectively, repetition rates 50 µs or 20 µs. Photoluminescence experiments at room temperature were carried out in nitrogen degassed 2x10⁻⁵ M CH₂Cl₂ solution. Measurements at 77 K were performed in 2-MeTHF frozen matrix.

DSC were measured in open vessel mode under nitrogen flow on a Mettler Toledo STARe SYSTEM DSC 500 and TGA analysis were performed under dry nitrogen flow on a Mettler Toledo STARe SYSTEM TGA/DSC 3+.

DFT and TDDFT simulations have been performed with the Gaussian package (Gaussian09 revision D.01).³⁸ We used the B3LYP hybrid exchange-correlation functional, which, in combination with TDDFT, has been shown to reproduce very well the adsorption spectra of Ir-complexes.^{39,40} We used the 6-31G(d,p) basis set for the light atoms and the aug-cc-pVDZ+ECP (Effective Core Potential) for iridium, which

includes spin orbit at the scalar-relativistic level.⁴¹ Thus, the effect of spin orbit on the valence states are not treated explicitly in our calculations. To obtain the optical absorption spectra up to 250-300 nm, we included 100 singlets and 100 triplets.

The triplet emission energies have been calculated with three different approaches, the first two based on the Δ SCF method, the third by TDDFT: (1) as the energy difference between T1 and the ground state SO, at their relaxed geometries. This is good approximation of the 0-0 emission peak at low temperature. (2) as the energy difference between the T1 and the S0 at the T1 relaxed geometry (vertical transition). (3) as the TDDFT vertical excitation at the T1 geometry in the Tamm-Dancoff approximation (TDA), to avoid the issues with the triplet instability.⁴² We did not include solvent effects because, as reported in literature, they have a minor effect on the absorption/emission wavelengths for this class of dyes. We also did not calculate account vibrational progressions, but rather we have applied a gaussian broadening of 0.2 eV to the calculated transitions.

Conflicts of interest

There are no conflicts to declare.

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How the triplet energy of β -diketonate ancillary ligands in Ir(III) complexes affects the phosphorescence emission: photochemical and electrochemical investigations and DFT calculations shed light on the dark triplet excited states.



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