



## H<sub>2</sub> detection mechanism in chemoresistive sensor based on low-cost synthesized WO<sub>3</sub> nanorods

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### ABSTRACT

Nanostructured WO<sub>3</sub> represents a promising material for fast and reliable molecular hydrogen detection through chemo-resistive effect. Here, an extended experimental investigation of WO<sub>3</sub>-H<sub>2</sub> interaction is presented and modeled. A powder of WO<sub>3</sub> nanorods (400 nm long, 50 nm large) is produced by hydrothermal technique and drop casted on Pt interdigitated electrode. H<sub>2</sub> sensing tests at different concentrations (2000–50,000 ppm) and temperatures (250–400 °C) are reported. Scanning Electron Microscopy (SEM), X-ray Diffraction analysis (XRD), and electrical measurements were performed. The response and recovery kinetics of H<sub>2</sub> sensing are carefully described by using a two-isotherms Langmuir model, and kinetics barriers for WO<sub>3</sub>-H<sub>2</sub> interaction are evaluated. Two microscopic processes lead to gas detection. A fast process (shorter than 4 s) is attributed to H<sub>2</sub> interaction with adsorbed oxygen at WO<sub>3</sub> nanorods surface. A slow process (20–1000 s), with activation energy of 0.46 eV, is attributed to oxygen vacancy generation in WO<sub>3</sub>. H intercalation in WO<sub>3</sub> is ruled out. The recovery of WO<sub>3</sub> after H<sub>2</sub> exposure is also modeled. The chemo-resistive effect leading to H<sub>2</sub> sensing by WO<sub>3</sub> is explained through the above processes, whose kinetic barriers have been quantified. These data open the route for the development of fast, sensitive, and low-temperature operating H<sub>2</sub> sensors based on WO<sub>3</sub>.

### 1. Introduction

In recent years the scientific community is devoting a large interest to hydrogen as a source of energy, since it is an excellent candidate to replace fossil fuels for many applications, such as chemical industry, power generation, medical treatment, and many others [1–3]. Soon, hydrogen production, transportation, and storage will become key issues for sustainable development. However, H<sub>2</sub> is a colorless, tasteless, and odorless gas which becomes inflammable and explosive when the concentration exceeds 4%. Therefore, fast and very sensitive sensors are necessary where H<sub>2</sub> is produced, stored, and transported to promptly find any leakage. Typically, fast H<sub>2</sub> sensors work at high temperatures (250–400 °C), increasing the explosion danger [4,5]. The progress in H<sub>2</sub> application cannot proceed without a parallel development of safe, fast, and reliable H<sub>2</sub> sensors. To this aim, a proper investigation is needed to unveil the interaction between H<sub>2</sub> and the sensing material leading to reliable detection.

One of the most promising materials for H<sub>2</sub> sensing is WO<sub>3</sub>, an *n*-type metal oxide semiconductor with an indirect bandgap ranging between

2.6 and 3.2 eV [6]. It has been widely studied for its reversible chromogenic ability under external stimuli such as electrical bias (electrochromism) or reducing ambient (gaschromism) [7], for its ability to detect many gases such H<sub>2</sub>S, NH<sub>3</sub>, NO<sub>2</sub>, and H<sub>2</sub> [8–11] and for its properties in electrochemical storage application [12,13]. Nanostructured WO<sub>3</sub> shows better performances thanks to a larger surface-to-volume ratio, higher specific surface area, and crystal faceting with marked surface reactivity [14–16]. Further, improvement of performances is obtained by surface decoration with noble mono-metallic, or bimetallic nanoparticles, such as Pt [17], Pd [18], Au [19] or proper combination of these [20,21] leading to synergistic effects [22].

The affinity of WO<sub>3</sub> for reducing gases makes it an excellent candidate for optical and chemoresistive H<sub>2</sub> sensor realization since low concentrations can be detected also below 250 °C. Zhang et al. [11] synthesized urchin-like hexagonal WO<sub>3</sub> and tested them at 250 °C for H<sub>2</sub> sensing at concentrations ranging between 10 and 80 ppm. Wu et al. [18] synthesized Pd loaded mesoporous WO<sub>3</sub> for H<sub>2</sub> sensing test. They compared bare WO<sub>3</sub>, mesoporous WO<sub>3</sub>, and Pd-loaded mesoporous WO<sub>3</sub>, showing that the latter is the best one in terms of response

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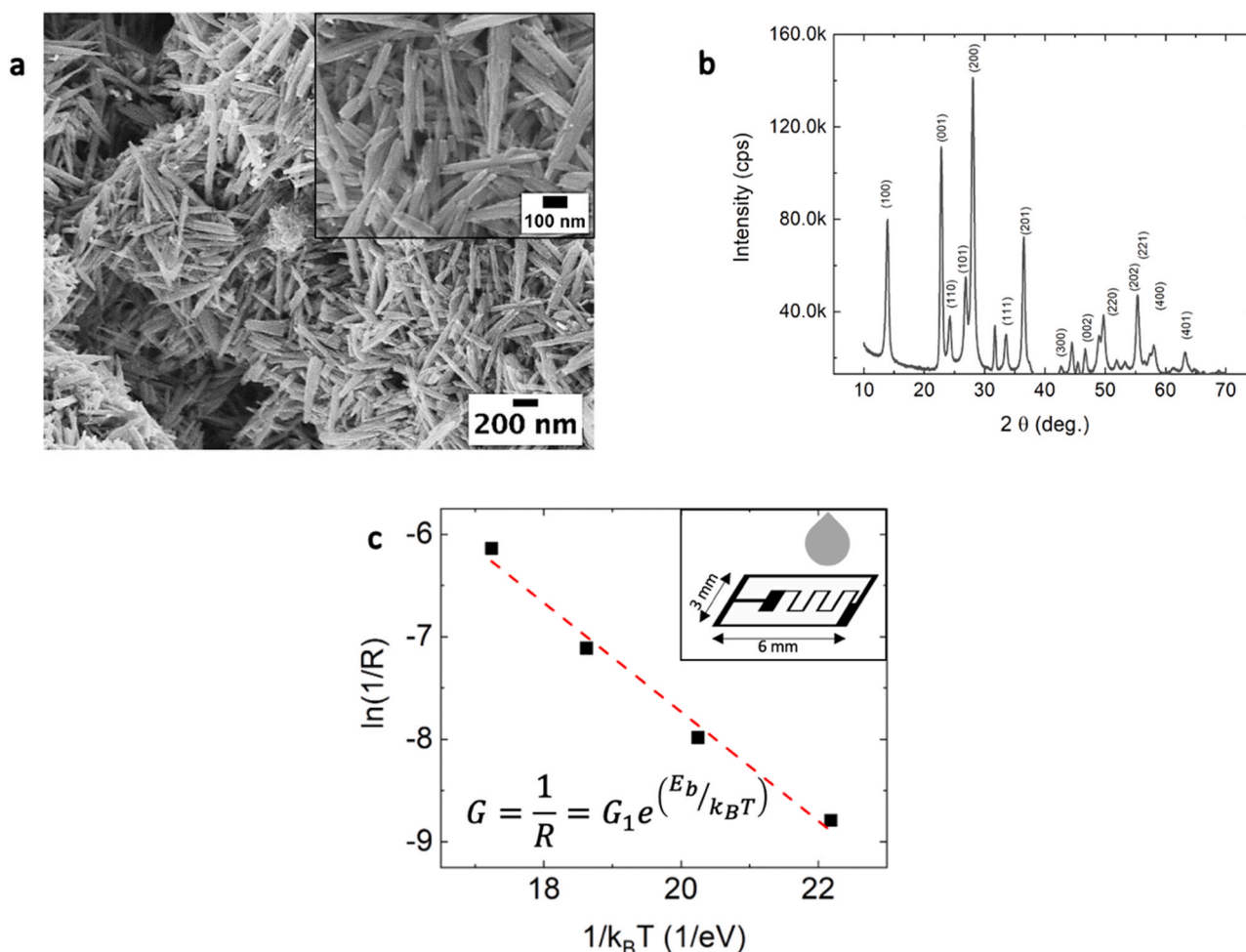
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**Fig. 1.** (a) low and (inset) high magnification SEM image of WO<sub>3</sub> nanorods deposited on the interdigitated substrate, (b) XRD pattern, and (c) Arrhenius plot of inverse of measured resistance (R); inset in (c): schematic of interdigitated substrate drop coated.

efficiency at room temperature at 5000 ppm of H<sub>2</sub>. Chang et al. [23] studied the gas sensing performances of Pt-loaded WO<sub>3</sub> thin-film at 200 °C under different concentrations of H<sub>2</sub> (between 1 and 10,000 ppm). Mattoni et al. [17] synthesized a single-crystal Pt-decorated WO<sub>3</sub> thin film and tested it at room temperature under sub-ppm concentration of hydrogen. In 2017, using a density functional theory approach, Tian et al. [24] proposed a model for the interaction between H<sub>2</sub> and hexagonal WO<sub>3</sub>, according to which the H<sub>2</sub> sensing mechanism proceeds through H<sub>2</sub>O molecules formation after interaction of H<sub>2</sub> with oxygen ions adsorbed onto noble metal nanoparticles, if any. Experimental works by Wu and Chang claimed Tian's model. On the other hand, Mattoni et al. proposed a different model based on adsorption of H<sub>2</sub> on Pt nanoparticles, H<sub>2</sub> dissociation and intercalation of H<sup>+</sup> ions inside the lattice of WO<sub>3</sub>. A comprehensive understanding of WO<sub>3</sub>-H<sub>2</sub> interaction is still missing, univocal experimental evidence of chemical reactions or energy barriers do not result from literature and a detailed description of temperature-activated kinetics is however required. A basic comprehension of actual processes underlying the WO<sub>3</sub>-H<sub>2</sub> interaction could promote a step forward in H<sub>2</sub> sensing by WO<sub>3</sub>.

In this paper, we synthesized by hydrothermal technique hexagonal WO<sub>3</sub> nanorods, and we studied the interaction with H<sub>2</sub> molecules at different temperatures and concentrations. The experimental kinetic curves of the material response and recovery were satisfactorily modeled by using two thermally activated processes leading to chemoresistive effect. A comprehensive description of the WO<sub>3</sub>-H<sub>2</sub> interaction is discussed.

## 2. Material and methods

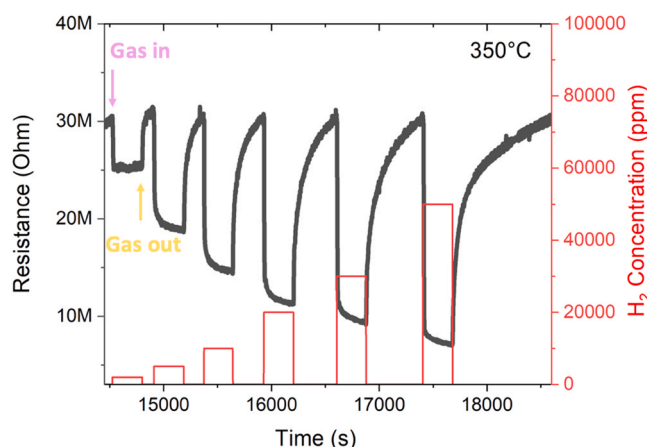
### 2.1. Synthesis

Hydrothermal synthesis of hexagonal WO<sub>3</sub> nanorods is employed by using sodium tungstate (Na<sub>2</sub>WO<sub>4</sub>), sodium chloride (NaCl), and HCl [25]. The experimental procedure is detailed here: 0.825 g of Na<sub>2</sub>WO<sub>4</sub> is added to 19 ml of deionized water. After stirring for 10 min, a certain quantity of HCl is added to solution to get a 2.4 pH. The solution is then put in a stainless-steel autoclave (25 ml). Before closing, 0.5 g of NaCl is added as a capping agent which leads to hexagonal nanorods morphology [26]. The sealed autoclave is put in an oven at 180 °C for 3 h. The autoclave is cooled down naturally, the obtained powder is separated (by centrifugation at 6000 rpm for 10 min) and washed with water and ethanol several times, finally dried on a hot plate for 1 h.

For the electrode realization, WO<sub>3</sub> nanorods powder is dissolved in 5 ml of deionized water and sonicated for 20 min. Several drops are then dipped on an interdigitated Pt electrode until all the fingers are fully coated.

### 2.2. Characterization

The morphological analyses were carried out by a scanning electron microscope (SEM) Gemini Field Emission SEM Carl Zeiss SUPRATM 25 (FEG-SEM, Carl Zeiss Microscopy GmbH, Jena, Germany). X-ray diffraction (XRD) patterns were acquired through a Bruker-AXSD50050-θ diffractometer, using a Göbel mirror to parallel the Cu Kα radiation



**Fig. 2.** Dynamic resistance variation of  $\text{WO}_3$ -based sensor (black curve) at  $350^\circ\text{C}$  under different  $\text{H}_2$  concentrations (red curve). Arrows indicate when  $\text{H}_2$  is fluxed in the chamber (named “gas in”) or it is stopped (named “gas out”). (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

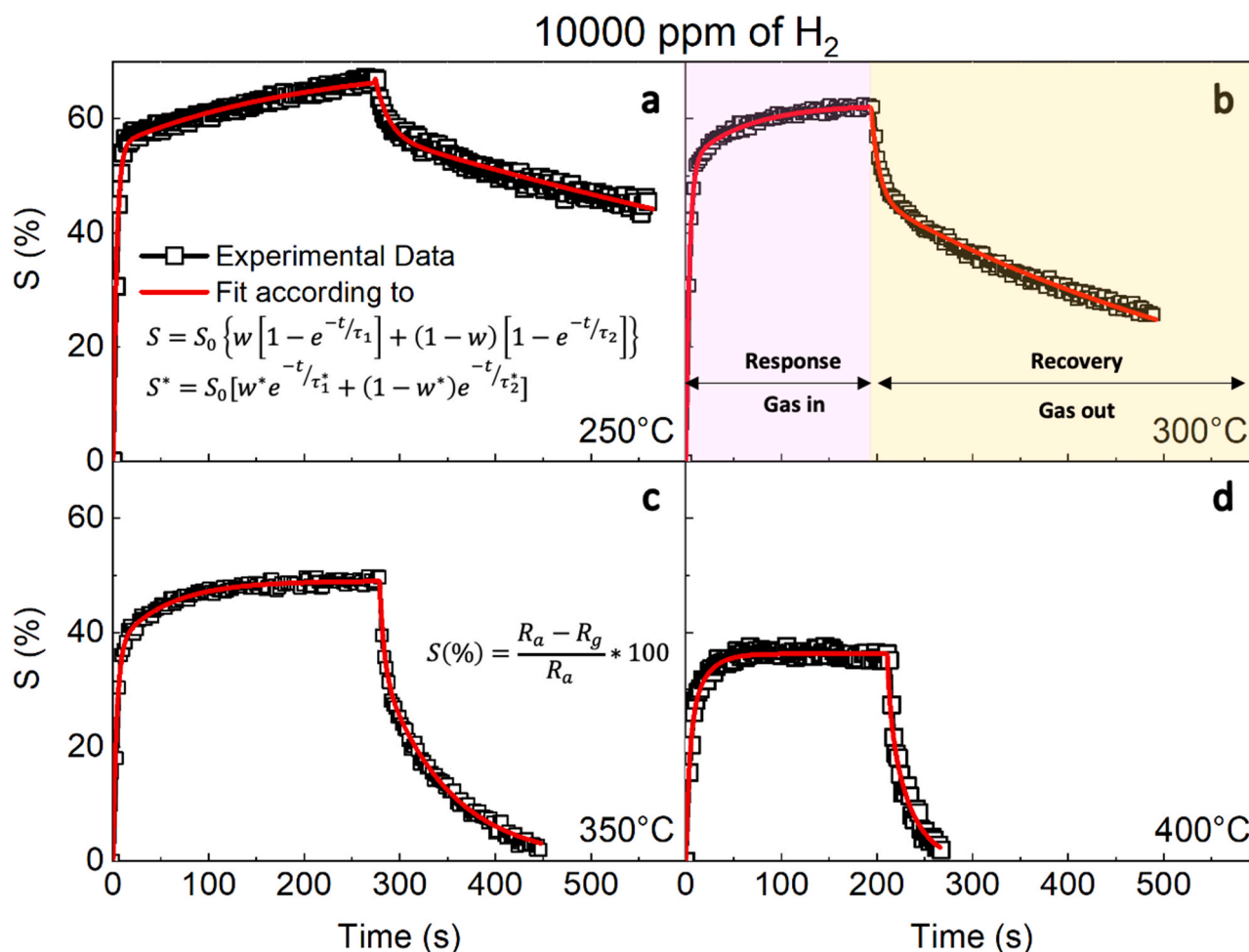
operating at 40 kV and 30 mA. The transmittance spectra were obtained by using a UV–VIS–NIR spectrophotometer Varian Cary 500. Gas sensing measurements were performed in a stainless-steel test chamber which allowed measurements in a controlled atmosphere. Gases coming from

certified bottles were used and diluted in synthetic dry air at a given concentration by using mass flow controllers. The sensor was heated from room temperature ( $\sim 20^\circ\text{C}$ ) up to  $400^\circ\text{C}$  under a dry air ( $\text{RH} < 3\%$ ) total stream of 100 sccm by using a dual-channel power supplier instrument Agilent E3632A to bias the built-in heater. The response to different gases was evaluated by recording the resistance at an applied voltage of 1.0 V through a Keithley 6487 picometer with a time sensitivity of about 2.5 s. A detailed description of the sensor platform and the setup for  $\text{H}_2$  sensing tests are reported in a previous paper [27].

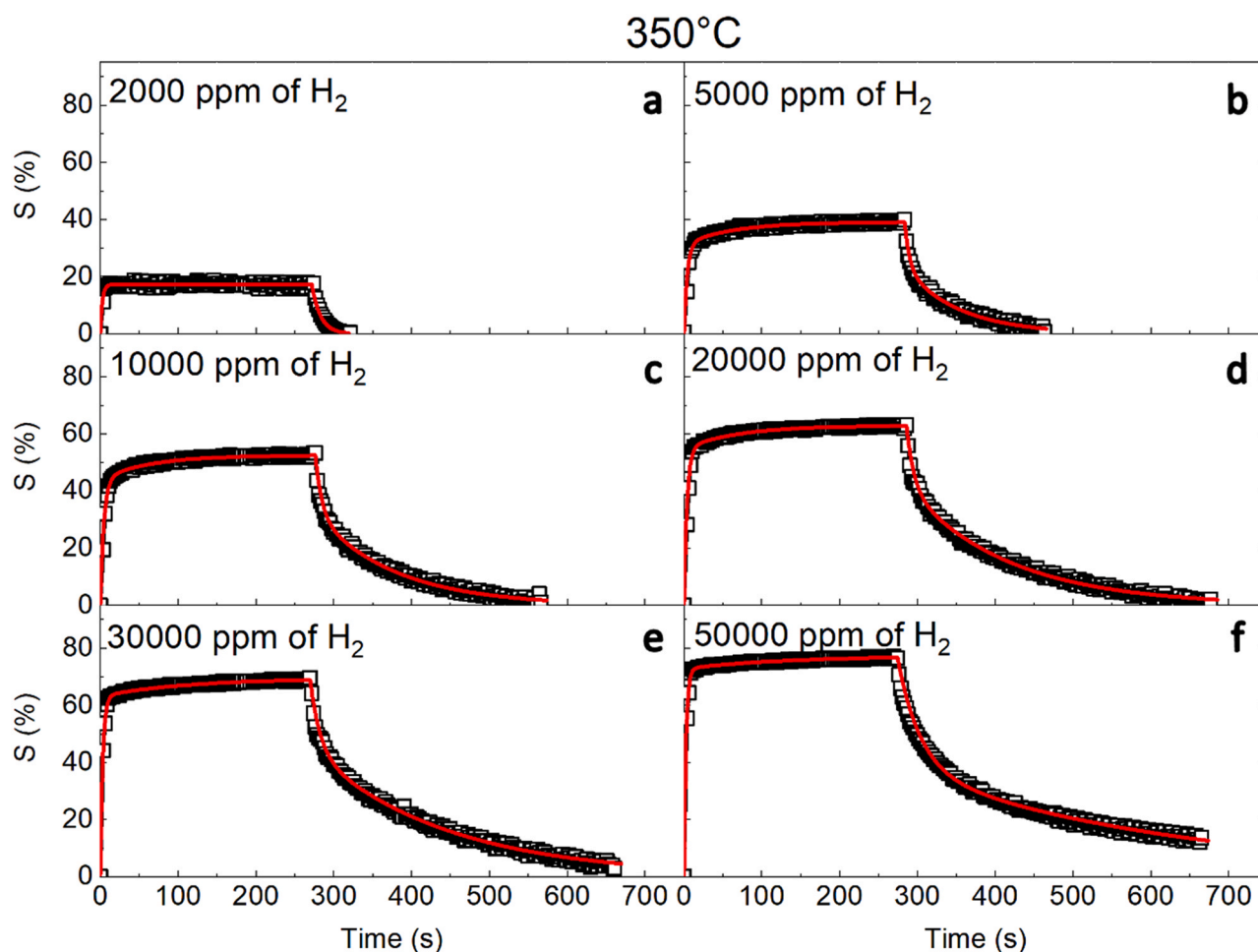
### 3. Results and discussion

#### 3.1. Structural and electrical properties of $\text{WO}_3$ nanorods

Fig. 1a shows low and high (inset) magnification SEM images of obtained powder on interdigitated electrode, confirming the nanorods morphology. An aspect ratio of about 8 is determined by SEM images showed in Fig. S1 (nanorods average length and diameter of 400 and 50 nm, respectively). All interdigitated fingers were covered to maximize the response during gas sensing measurements. Some drops of  $\text{WO}_3$  nanorods solution (powder well dissolved in deionized water) were drop casted on a corner glass and XRD analyses were performed to identify the crystal structure of our nanostructures. From the powder diffraction pattern shown in Fig. 1b, the typical peaks of hexagonal  $\text{WO}_3$  (PDF card 00 075 2187) can be observed, with the corresponding Miller Indexes reported. This result confirms the hexagonal structure of  $\text{WO}_3$  nanorods.



**Fig. 3.** Dynamic responses of  $\text{WO}_3$ -based sensor (squared symbols) and fit (red lines) obtained at (a)  $250^\circ\text{C}$ , (b)  $300^\circ\text{C}$ , (d)  $350^\circ\text{C}$ , and (d)  $400^\circ\text{C}$  under 10,000 ppm of  $\text{H}_2$  as a function of elapsed time starting when  $\text{H}_2$  is fluxed in chamber. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)



**Fig. 4.** Dynamic responses of  $\text{WO}_3$ -based sensor (squared symbols) and fit (red lines) obtained at  $350^\circ\text{C}$  under (a) 2000 ppm, (b) 5000 ppm, (c) 10,000 ppm, (d) 20,000 ppm, (e) 30,000 ppm and (f) 50,000 ppm fluxes of  $\text{H}_2$  as a function of elapsed time. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

We investigated the electrical resistance  $R$  in dry air at different temperatures  $T$  (Fig. 1c). By increasing the temperature from  $250^\circ\text{C}$  to  $400^\circ\text{C}$ , the electrical conductance  $G (=1/R)$  of  $\text{WO}_3$  nanorods increases by almost three orders of magnitude, following an exponential behavior with an activation energy  $E_b$  of 0.53 eV. In an  $n$ -type semiconductor, as the  $\text{WO}_3$ , the conductivity depends on electron concentration in conduction band and on electron mobility. Disregarding the thermal dependence of electron mobility, the measured activation energy for conductance could account for the energy position (0.53 eV) below the conduction band of the intrinsic donor-like defects of undoped  $\text{WO}_3$ , in agreement with Kalanur et al. [28].

### 3.2. $\text{H}_2$ sensing results

Experimental  $\text{H}_2$  sensing curves are obtained by measuring resistance variations as a function of time while a mixture of  $\text{H}_2$  and dry air (response phase) or dry air only (recovery phase) are fluxed in the chamber. Resistance variations measured at different temperatures under dry synthetic air and  $350^\circ\text{C}$  under different  $\text{H}_2$  concentrations are reported in Fig. S2. To investigate the  $\text{WO}_3$ - $\text{H}_2$  interaction, we tested different concentrations (2000–50,000 ppm) and different temperatures ( $250$ – $400^\circ\text{C}$ ). Fig. 2 shows resistance variations (black curve) at  $350^\circ\text{C}$  under different  $\text{H}_2$  concentrations (red curve), confirming a strong dependence of resistance on  $\text{H}_2$  concentration. Resistance decreases when  $\text{H}_2$  is fluxed in the test chamber (gas in arrow), while increases when  $\text{H}_2$  flux is stopped (gas out arrow).

From experimental curves of resistance variation, a sensor response  $S$  (%) can be defined. Since  $\text{WO}_3$  is an  $n$ -type semiconductor and  $\text{H}_2$  is a reducing gas,  $S$  (%) is determined as follows:

$$S(\%) = (R_a - R_g) / R_a * 100 \quad (3.2.1)$$

in which  $R_g$  and  $R_a$  are the measured resistances in presence and absence of  $\text{H}_2$  gas, respectively. To understand  $\text{WO}_3$ - $\text{H}_2$  interaction,  $S$  (%) is defined for each tested temperature and concentration. Fig. 3 shows sensor response curves obtained at each temperature ( $250$ – $400^\circ\text{C}$ ) as a function of elapsed time during the “response” (pink background in panel b) and “recovery” (yellow background in the panel b) phase, respectively.

In each case of Fig. 3, the response curve increases rapidly when  $\text{H}_2$  is fluxed into the chamber, while it tends to saturate after about 20 s. When  $\text{H}_2$  flux is stopped, the recovery mechanism starts, and the curves decrease in a temperature dependent way.

$\text{H}_2$  concentration tests are carried out at  $350^\circ\text{C}$  and Fig. 4 shows response and recovery curves obtained under (a) 2000 ppm, (b) 5000 ppm, (c) 10,000 ppm, (d) 20,000 ppm, (e) 30,000 ppm, and (f) 50,000 ppm fluxes of  $\text{H}_2$ .

Response curves increase rapidly when  $\text{H}_2$  gas is fluxed in chamber. After about 20 s the curves tend to saturation and when  $\text{H}_2$  flux is stopped the recovery mechanism starts, with the curves decreasing rapidly at first but then slowly tending to the initial value.

About the stability, different pulses of 10,000 ppm of  $\text{H}_2$  were fluxed into the chamber at  $350^\circ\text{C}$ , as shown in Fig. S3(a). The sensor response



**Table 1**

Results of the fitting procedure to response and recovery transients at various temperatures and at fixed concentration of H<sub>2</sub> (10,000 ppm) based on Eqs. (3.3.1) and (3.3.2).

Temperature (°C)	S <sub>0</sub> (%)	τ <sub>1</sub> (s)	τ <sub>2</sub> (s)	w (%)	τ* <sub>1</sub> (s)	τ* <sub>2</sub> (s)	w* (%)
250	70	4	200	78	14	1137	15
300	63	4	61	83	8	483	29
350	49	4	54	76	5	71	32
400	36	4	16	60	3	23	25

**Table 2**

Results of the fitting procedure to response and recovery transients at various concentrations of H<sub>2</sub> at a temperature of 350 °C, based on Eqs. (3.3.1) and (3.3.2).

Concentration 10 <sup>3</sup> (ppm)	S <sub>0</sub> (%)	τ <sub>1</sub> (s)	τ <sub>2</sub> (s)	w (%)	τ* <sub>1</sub> (s)	τ* <sub>2</sub> (s)	w* (%)
2	17	3	–	100	11	–	100
5	39	4	63	80	6	71	42
10	53	5	65	84	8	102	40
20	63	4	70	87	10	132	34
30	69	3	111	91	12	174	36
50	77	3	126	94	29	359	50

shows a slight drift after 4 pulses, demonstrating a fair stability and a good reactivation, since the resistance is restored at the initial condition. Moreover, we tested WO<sub>3</sub> sensor selectivity under pulses of different gases (10,000 ppm of H<sub>2</sub>, 500 ppm of NO and 50 ppm of CO) at temperatures ranging between 250 ° and 400 °C as shown in Fig. S3(b). The H<sub>2</sub> response is higher than the NO response, especially at higher temperatures, while no response is obtained for CO whatever the temperature is, confirming a certain selectivity of the WO<sub>3</sub> sensor towards H<sub>2</sub> as

previously reported in the literature [7]. The calibration curve at 350 °C is reported in Fig. S3(c). The LoD (Limit of Detection), calculating by considering S/N = 3, is 1076 ppm [29].

### 3.3. H<sub>2</sub> sensing kinetics

To determine the kinetics of the H<sub>2</sub> sensing, all experimental response curves were fitted by considering the Langmuir absorption theory with two isotherms (two processes are assumed to play independent roles). According to Langmuir, sensor response during the “gas in” phase, S, and during the “gas out” phase, S\*, can be fitted as follows [29]:

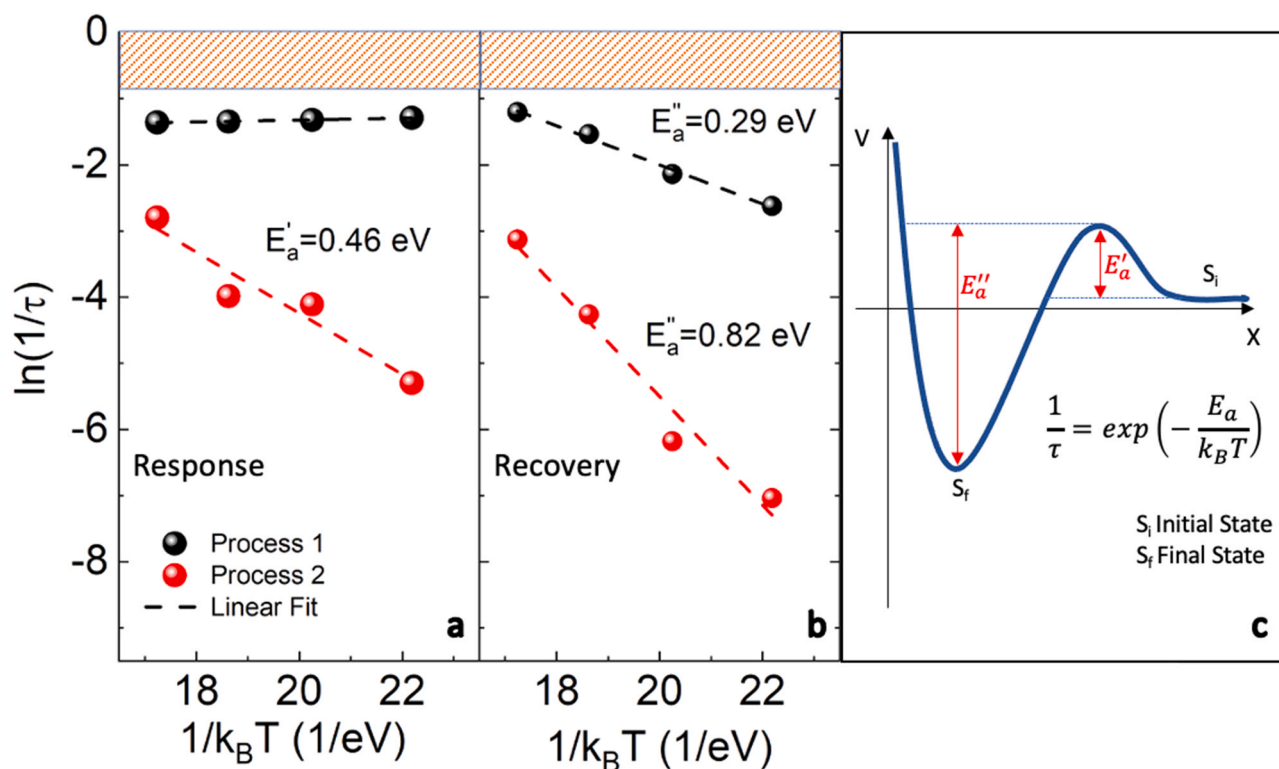
$$S = S_0 \{ w [1 - \exp(-t/\tau_1)] + (1-w) [1 - \exp(-t/\tau_2)] \} \quad (3.3.1)$$

$$S^* = S_0 [w^* \exp(-t/\tau_1^*) + (1-w^*) \exp(-t/\tau_2^*)] \quad (3.3.2)$$

in which S<sub>0</sub> is the maximum value of resistance change (the same value was fixed for response and recovery phases), τ<sub>n</sub> and τ<sub>n</sub><sup>\*</sup> (n stands for 1 or 2) are lifetimes associated to the two isotherms during response and recovery phases, respectively; w and w\* are weights attributed to process 1 during response and recovery phases, respectively.

The fitting curves (red lines) shown in Fig. 3 are in excellent agreement with experimental data (black squares), supporting the hypothesis of a fast (P1) and a slow (P2) process, independent and concomitant. The fitting parameters are listed in Table 1.

The maximum response value (S<sub>0</sub>) changes with temperature: it is maximum (70 %) at 250 °C, as expected given the very high resistance at this temperature. Lifetimes and weights are reported in Fig. S4 as a function of temperature. P1 is characterized by times of few seconds (roughly the time resolution of experimental set-up), both in response and in recovery phases and regardless of temperature, while P2 is characterized by long times (10–200 s) in response and (20–1100 s) in recovery phases. In both cases, τ<sub>2</sub> and τ<sub>2</sub><sup>\*</sup> are clearly dependent on



**Fig. 5.** Arrhenius plot of inverse of lifetime τ of P1 (black sphere) and P2 (red sphere) during (a) response and (b) recovery phase. Dotted lines are linear fit to each set of data. The evidenced region indicates the time resolution of experimental set-up. (c) schematic representation of energy barriers for the sensing processes. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

temperature, since they decrease when temperature increases. Concerning the weight,  $w$ , P1 is predominant ( $w$  around 60–84%) during the response phase while it is minority during the recovery phase. Such a result could suggest different mechanisms occurring during response and recovery phases.

We performed an analogous analysis when tests were carried out at fixed temperature by varying  $H_2$  concentration. For all these cases shown in Fig. 4, except at 2000 ppm, Eqs. (3.3.1) and (3.3.2) can be used to fit experimental curves by obtaining good agreement. In the case of the lowest concentration (Fig. 4(a)) only one process (P1) is identified both during the response and the recovery phase, and so,  $w$  is maximum. Table 2 reports all parameters obtained from fitted curves. The maximum response values depend on concentration since the highest value (77%) is obtained at a higher concentration (50,000 ppm). Moreover, lifetimes and weights are reported in Fig. S5 as a function of concentration.

Both in response and recovery phases, P1 is characterized by short lifetimes, regardless of concentration, while P2 has increasing lifetimes with increasing concentration. P1 is predominant in the response phase since  $w$  is around 80–95%, but not during the recovery phase. Also, in this case, this can be due to different mechanisms occurring during response and recovery phases. Fig. S4(a) shows the  $w$  trend as a function of  $H_2$  concentration in cases where 2 processes are identified, revealing an exponentially increasing trend up to saturation. It should be noted that, within the Langmuir theory, the surface coverage  $\theta_m$  by the adsorbed gas molecules [30] is expected to increase similarly. This evidence suggests that process P1 is related to a surface coverage by  $H_2$ .

The Langmuir adsorption model with two isotherms very well describes our experimental data, pointing out that the  $WO_3$ - $H_2$  interaction proceeds through at least two ways. Indeed, the thermal dependence of  $\tau$  is highly significant as it will unveil the energy barrier for concurring processes. Fig. 5 reports the inverse of lifetimes for the two processes, both in response (a) and in recovery (b) phases in an Arrhenius plot, by clearly showing linear behaviors for the four cases. It should be noted that in the response phase the fast process is too close to the time resolution which hinders us to resolve high frequency (orange box in Fig. 5) so that the linear trend can be meaningful only for this process. The inverse of a lifetime gives the probability rate, which is thermally activated, as follows:

$$1/\tau = \exp(-E_a/k_B T) \quad (3.2.3)$$

where  $k_B$  is Boltzmann constant,  $T$  is temperature and  $E_a$  is activation energy ( $E_a$  for response phase and  $E'_a$  for recovery phase). The physical meaning of these activation energies is depicted in Fig. 5c, as kinetic energy barriers in the interaction potential  $V$  between the  $WO_3$  surface and the approaching  $H_2$  molecule. By analyzing the meaningful Arrhenius plots reported in Fig. 5 for P1 (black spheres) and P2 (red spheres) involved during (a) response phase and (b) recovery phase, we can extract  $E_a$  values for each involved process (dotted lines). The result of such an exercise gives  $E_a$  of 461 meV in the response phase for P2; 295 meV and 821 meV in the recovery phase for P1 and P2, respectively. The experimental determination of these activation energies is a key point to reach a solid description of the  $WO_3$ - $H_2$  interaction leading to  $H_2$  sensing.

### 3.4. Modeling of $WO_3$ - $H_2$ interaction

It is well assessed that H intercalation occurs in hexagonal  $WO_3$ , especially during the electrochromic process [31].  $H_2$  intercalation causes internal changes of  $WO_3$  structure also reflected by a significant change of optical transmittance. Mattoni et al., proposed that  $WO_3$ - $H_2$  interaction leads to  $H_2$  intercalation in  $WO_3$  crystals [17], while for other models such a process is not required. It is essential now to prove if  $H_2$  intercalation occurs or not in our samples. To this aim, several drops of  $WO_3$  nanorods solution were dropped on corning glass substrates,

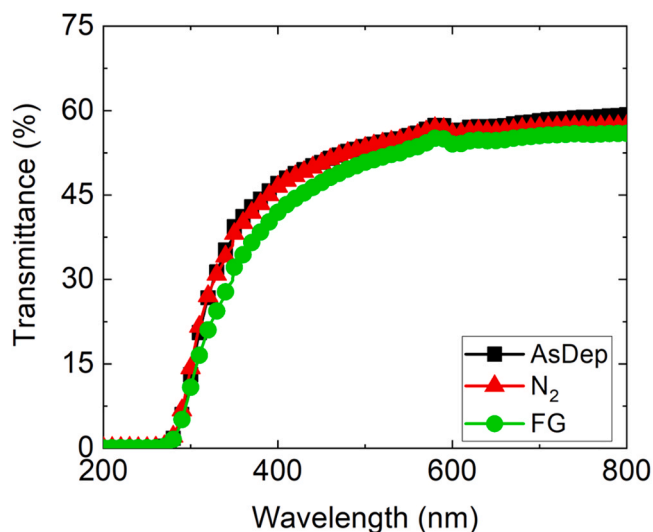
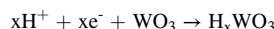


Fig. 6. Transmittance spectra of AsDep (black line + squares),  $N_2$  (red line + triangles), and FG (green line + circles) samples. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

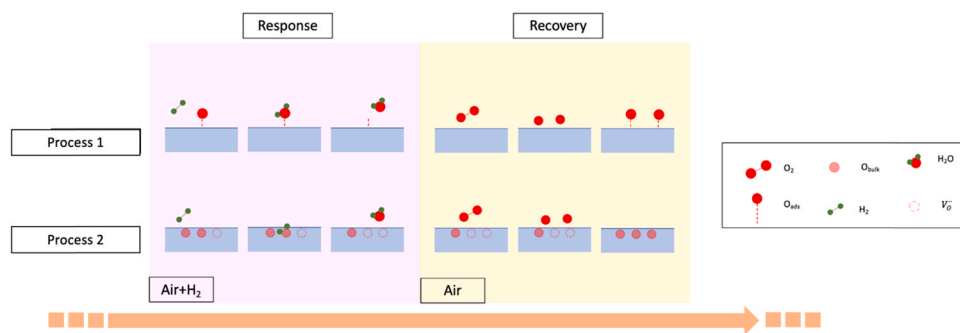
realizing a very thin film. Two  $WO_3$  coated glasses were annealed at 350 °C for 30 min in forming gas (FG, gas mixture of  $N_2:H_2 = 95:5$ ) or in  $N_2$ . Annealing in FG simulates the exposure of  $WO_3$  nanorods to  $H_2$  gas during the sensing test at 350 °C. Transmittance spectra acquired just after annealing of  $WO_3$  coated glasses are shown in Fig. 6. The annealing in inert ambient does not significantly modify the transmittance curve of the sample. FG annealing induces a barely appreciable reduction of transmittance, far away from the expected chromism induced by intercalating  $H_2$ . Moreover, the  $H_2$  intercalation typically produces an enhancement in measured resistance according to Mirzaei et al. [32] as following reaction consequence:



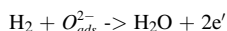
On the contrary, after  $H_2$  exposure, our samples reveal a marked decreased resistance at any temperature and any  $H_2$  concentration. Such evidence, together with transmittance spectra, tell us that during  $H_2$  exposure H intercalation in  $WO_3$  nanorods, if any, is not a key process and can be ruled out in modeling the  $WO_3$ - $H_2$  interaction.

To model the  $WO_3$ - $H_2$  interaction, we now need to discuss the P1 and P2 occurring in response and recovery phase. Firstly, it must be specified that during the response phase the interaction will involve  $H_2$ , while during the recovery phase, only residual  $H_2$ , if any, can be considered. Experimental results give us much information about active processes. The evidence can be summarized as follows: (i) P1 is faster than P2, both in response and in recovery phases.  $\tau_1$  is around 4 s, regardless of temperature and concentration, while  $\tau_1^*$  slightly depends on temperature and concentration; (ii) P1 is more probable than P2, both in response and recovery phases, and at any temperature and any concentration; (iii) the  $H_2$  concentration dependence of  $w$  suggests a surface process for P1; (iv) Arrhenius plots hint that P1 is a very fast process in response phase, while it is thermally activated ( $E''_a = 295$  meV) at recovery phase when the temperature is ranging between 250 °C and 400 °C. P2 needs high activation energy both in response (461 meV) and recovery phases (821 meV) in the same temperature range; (v) at very low  $H_2$  concentration (below 5000 ppm), only one process is active.

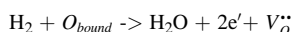
Considering the above discussion on experimental evidence, and assuming that  $WO_3$  nanorods are typically adsorbed with oxygen at surface, the following model is proposed. In response phase, P1 can be described as a *surface process* where a spontaneous interaction occurs between  $H_2$  molecules and oxygen ions adsorbed on the  $WO_3$  surface leading to a water molecule as follows:



**Fig. 7.** Schematic model of surface processes (P1) and bulk processes (P2) involved during the response (pink evidenced area) and the recovery (yellow evidenced area) phase. H<sub>2</sub> interacts with adsorbed oxygen and oxygen directly bound to W to produce a water molecule in response phase, while O<sub>2</sub> can be adsorbed on WO<sub>3</sub> surface or can occupy oxygen vacancy for restoring the initial condition. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

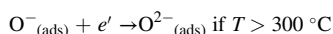
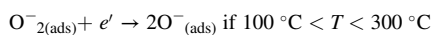
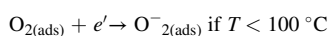
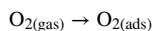


In addition, once water evaporated and oxygen adsorbed on the surface has been reduced, H<sub>2</sub> interacts with the bare WO<sub>3</sub> surface through a *bulk process* involving oxygen atoms bound to surface W atoms [33]. This interaction can lead to P2 in response phase when a water molecule is produced leaving an oxygen vacancy V<sub>O</sub><sup>••</sup>, as follows:

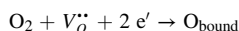


During response phase, both P2 and P1 produce a decrease in measured resistance, according to experimental data.

A different scenario occurs in the recovery phase, where the most probable interaction occurs between O<sub>2</sub> present in dry synthetic air and WO<sub>3</sub> surface. P1 can arise from interaction with oxygen adsorbed on WO<sub>3</sub> surfaces, depending on working temperature as follows [23]:



P2 can be described as the interaction among O<sub>2</sub> and oxygen vacancies, leading to defect recombination as follows:



Processes occurring in the recovery phase let WO<sub>3</sub> resistance increase, according to experimental data. At very low H<sub>2</sub> concentration, only surface processes occur, both in response and recovery phases. The H<sub>2</sub>-WO<sub>3</sub> interaction, under exposure to high H<sub>2</sub> concentration, could produce a so high density of oxygen vacancies at surface that a diffusion towards bulk can occur. Such a scenario could explain the long response and recovery times at high H<sub>2</sub> concentration (Table 2).

Fig. 7 shows a schematic representation of the above-proposed model during response and recovery phases (pink and yellow evidenced area, respectively), where the distinction between *surface* and *bulk process* (P1 and P2 respectively) is clear.

#### 4. Conclusion

We investigated WO<sub>3</sub>-H<sub>2</sub> interaction by testing WO<sub>3</sub> nanorods exposed to H<sub>2</sub> (2000–50,000 ppm) fluxes at different temperatures (250–400 °C). The Langmuir adsorption theory was successfully used to model experimental data of sensing kinetics, evidencing that two processes occur both in response and in recovery phases, regardless of the temperature, for concentrations larger than 5000 ppm. The fast process (lifetimes of few seconds) is attributed to H<sub>2</sub> interaction with adsorbed O at WO<sub>3</sub> surface, in the response phase, and to O adsorption (barrier of 0.29 eV), in the recovery phase. Since a high concentration of H<sub>2</sub> consumes all adsorbed oxygen leaving WO<sub>3</sub> uncovered, the slow process

(lifetimes of hundreds of seconds) is attributed to generation (barrier of 0.46 eV) and recombination (barrier of 0.82 eV) of oxygen vacancies in WO<sub>3</sub> nanorods. No H intercalation in WO<sub>3</sub> bulk is observed. At low H<sub>2</sub> concentrations, only fast processes are evidenced, which are linked to surface processes. Our data and modeling show that WO<sub>3</sub> nanorods are a promising material for H<sub>2</sub> sensing and, to improve the sensitivity at low hydrogen concentrations, increasing the exposed surface of WO<sub>3</sub> nanorods is beneficial for fast detection.

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#### CRediT authorship contribution statement

**G. Mineo, S. Mirabella, E. Bruno:** conceptualization; **G. Mineo, K. Moulae, G. Neri, S. Mirabella, E. Bruno:** investigation; **G. Mineo, K. Moulae, G. Neri, S. Mirabella, E. Bruno:** methodology; **G. Neri, S. Mirabella, E. Bruno:** supervision; **G. Mineo:** writing—original draft; **G. Mineo, G. Neri, S. Mirabella, E. Bruno:** writing—review and editing. All authors have read and agreed to the published version of the manuscript.

#### Declaration of Competing Interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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#### Appendix A. Supporting information

Supplementary data associated with this article can be found in the online version at [doi:10.1016/j.snb.2021.130704](https://doi.org/10.1016/j.snb.2021.130704).

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