Research Article

Radium in groundwater hosted in porous aquifers: estimation of retardation factor and recoil rate constant by using NAPLs

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Abstract

A new method for estimating retardation factor and recoil constant of radium isotopes in groundwater hosted in porous aquifers is described. The method is based on the evidence that alpha-recoiled radium ions, supplied by thorium parent atoms which occur in phases immersed in NAPL (Non-Aqueous Phase Liquids), are not adsorbed on solid phases. Experimental evidence is given that manganese dioxide, zeolite 4A, natural clay, monazite and weathered volcanic rock, all phases normally adsorbing radium from aqueous solutions, when immersed in NAPL adsorb negligible amounts of radium. This allows using experimental data on rock samples, representative of porous aquifers, for estimating Ra retardation factor and its alpha recoil constant in groundwater, without using Rn data as a comparison term. Unlike estimation of retardation factor between the "NAPL method" and the method based on comparison with radon depends on the different process of entry from aquifer rock into groundwater for radon and radium. Precise estimates of retardation factor and recoil constants of apply equations ruling the temporal evolution of radium isotopes in groundwater and to determine its age. Implications, useful for measuring the contamination age of soils by NAPL fluids, are described as well.

Keywords Retardation factor · Recoil constant · Radium isotopes · NAPL · Groundwater

1 Introduction

The goal of this paper is to provide new insights for determining groundwater ages using radium isotopes. When using radioactive tracers, the tracer age differs from the groundwater age by a factor, which is named tracer retardation factor. This factor, in case of conservative tracers, equals the unity. It should be noted that the groundwater age remains only an estimate of the mean residence time (MRT) of groundwater in aquifer [1]. In fact the shape of the age distribution in real-world groundwater is unknown. Thus, MRT is a weighted average over many idealized ages and the age of a single sample of ground-water can differ notably from the MRT of groundwater.

Ra isotopes are a quartet of radioactive isotopes with half lives extending from some days up to thousands of years (²²⁴Ra $t_{1/2}$ =3.66 days, ²²³Ra $t_{1/2}$ =11.4 days, ²²⁸Ra $t_{1/2}$ =5.75 years, and ²²⁶Ra $t_{1/2}$ =1600 years). They belong to the natural radioactive series of ²³⁸U (²²⁶Ra), ²³⁵U (²²³Ra) and ²³²Th (²²⁴Ra and ²²⁸Ra). The hydrogeochemistry of radium depends mainly on four factors [2]: (1) a low ionic potential that causes strong adsorption of Ra ions on the

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surface of aquifer rocks; (2) coprecipitation with barium in sulphate and carbonatic groundwaters; (3) mobilization from aquifer rocks by chemical processes like dissolution of Ra-hosting minerals or Ra desorption from aquifer rocks in saline groundwaters; (4) mobilization from aquifer rocks in groundwater by the physical processes of alpha-recoil.

The temporal evolution of Ra isotopes in groundwater hosted in porous aquifers, when secular equilibrium in host rocks is assumed and Ra precipitation rate can be neglected, is given by the following master equation, whose derivation is given in Appendix 1: [2–7]: product of the constant decay λ_i of ^{*i*}Ra and f_{α_i} the fraction of ^{*i*}Ra_r lost by α -recoil [11]:

$$\Gamma_i = \lambda_i \cdot f_\alpha \tag{3}$$

1.1.1 Alpha-recoil constant

Rn isotopes supply to groundwater, according Krishnawami [10], is considered as a measure of the supply of Ra isotopes as well. In this paper, this approach will be questioned. We will provide experimental evidence that

$${}^{i}\mathsf{Ra}_{w} = \left[\frac{\left(\Delta_{i} + \Gamma_{i}\right)}{\left(\lambda_{i} \cdot \mathsf{R}_{i}\right)}\right] \cdot \left(1 - \mathrm{e}^{-\lambda_{i} \cdot \mathsf{R}_{i} \cdot \mathsf{T}}\right) \cdot \left({}^{i}\mathsf{Ra}_{r} \cdot \frac{\rho_{r}}{\Phi}\right) + {}^{i}\mathsf{Ra}_{w(0)} \cdot \left(\mathrm{e}^{-\lambda_{i} \cdot \mathsf{R}_{i} \cdot \mathsf{T}}\right)$$
(1)

where, for a Ra isotope *i*, ${}^{i}Ra_{w}$ is its activity in water (Bq·L⁻¹) at time *T* (a) (i.e. groundwater age), ${}^{i}Ra_{w(0)}$ its initial activity in water (Bq·L⁻¹), ${}^{i}Ra_{r}$ the activity in the aquifer rock (Bq·kg⁻¹), R_i its retardation factor (dimensionless), λ_{i} (a⁻¹) its decay constant, Δ_{i} its dissolution rate constant (a⁻¹), Γ_{i} its α-recoil rate constant (a⁻¹) and finally Φ and ρ_{r} the porosity and the bulk density of the aquifer rock.

If ${}^{i}Ra_{w(0)}$ can be neglected, Eq. 1 reduces to:

$${}^{i}\mathrm{Ra}_{w} = \left[\frac{\left(\Delta_{i} + \Gamma_{i}\right)}{\left(\lambda_{i} \cdot \mathrm{R}_{i}\right)}\right] \cdot \left(1 - \mathrm{e}^{-\lambda_{i} \cdot \mathrm{R}_{i} \cdot T}\right) \cdot \left({}^{i}\mathrm{Ra}_{r} \cdot \frac{\rho_{r}}{\Phi}\right)$$
(2)

If information about $\Delta_{ir} \Gamma_{ir} R_{ir}^{i} R_{ar} \Phi$ and ρ_r is provided, the measurement of Ra isotopes in groundwater gives an estimate of the age T of groundwater. Some methods, either in the field or in laboratory, partially provide this information [2, 4, 5, 8]. Here we focus on a new experimental tool for determining Γ_{ir} and R_{jr} .

1.1 Retardation factor and α-recoil rate constant

In aquifers, and especially in those hosted in porous rocks, sorption reactions cause partitioning of reactive solute onto the solid phases composing the porous rocks. Sorption retards the transport of reactive solutes (as radium ions) making their average velocity over the travel path less than that of the groundwater flow, which is generally assumed equal to that of nonreactive/conservative solutes (as radon atoms). In the case of linear adsorption reactions [9], the ratio of nonreactive solute velocity to reactive solute velocity, which corresponds to the inverse of the ratio of their MRTs, is called retardation factor, R. Aquifer rock supplies Rn predominantly by α-recoil and this supply is orders of magnitude higher than the contribution from in situ Ra decay or rock dissolution [10]. The α -recoil rate constant Γ_i of ^{*i*}Ra, or recoil ejection rate of ^{*i*}Ra [5], which has the dimension of the inverse of a time, is defined as the

SN Applied Sciences A SPRINGER NATURE journal the modality of Ra supply to groundwater differs notably from that of Rn. Consequently, not only R_i but also the alpha recoil constant Γ_i of 'Ra in groundwater should be reconsidered, having been overestimated. In this frame, it is useful to consider the Rn emanation from solid materials into air for understanding the peculiar behaviour of Rn in comparison with Ra in groundwater [12]. A convincing analysis of the production of Rn is given by Maraziotis [13] and it can be applied to any other α -recoiled radionuclide like Ra as well. The Rn emanation coefficient, f, has three components: f_{r} , direct recoil; f_{p} , intraparticle pore recoil; f_{dr} solid-state diffusion. The direct recoil component, f_{rr} is the fraction of radon escaping from a particle due to alpha-recoil processes which follow the radium decays in the outer nanometric thickness of the particle. The intraparticle pore recoil component, f_{p} , is the fraction of radon that enters the nanopore framework (i.e.nanocrack framework) of the particle due to alpha-recoil processes from radium decays inside the particle. Finally, the solid-state diffusion component, $f_{d'}$ is the fraction of radon that exits a solid particle due to its diffusion within definite (crystalline) or non-definite (amorphous) lattice of the particle. The definition of these components is given by Eqs. 35–37 (see Appendix 2). For ²²⁰Rn (Fig. 1), due to its short halflife, the situation is simpler, since f_d is negligible at ambient temperature [14]. These components depend only on variables linked to the particle from what Rn emanates, so Eqs. 35-37 apply to Ra emanation from solids into water as well. Let us to consider now the Rn and Ra emanation into groundwater in saturated aquifers. The diffusive component f_d is considered null for Ra emanation (see Appendix 3). Thus, if nanopores were quite water-filled and Rn and Ra were produced only by α -recoil ($f_d = 0$), then the Rn/Ra ratio in groundwater should represent actually R_{Ra}, assuming steady state conditions and secular radioactive equilibrium (i.e. that exponential term in Eq. 1: $e^{-\lambda_i \cdot R_i \cdot T} \approx 0$). Nevertheless, if confined air occurs, it interrupts the water



Fig. 1 Main components (f_p =intraparticle pore recoil, f_r =direct recoil) of ²²⁰Rn emanation coefficient vs. particle radius, according the model of Maraziotis [13]. Bottom: ratio between the main components vs. particle radius. Assigned conditions (Eq. 35–37): k=0.30, r_n 30 nm, R_r=30 nm

continuity and α -recoiled Ra ions, unlike neutral Rn atoms, fastly attach to the walls of nanopores [15, 16] (Fig. 2a). In a porous medium, f_{α} is the sum of f_r and f_{p_r} provided that diffusion and implantation effects in adjacent grains can

be neglected. Because of the presence of air in nanopores, intraparticle pore recoil component in general can be neglected as well, then $f_a = f_r$. Although the presence of nanopores as intermediary channels for Rn transport has been guestioned [17, 18], other two factors advice against the use of Rn as reference element for α-recoil of Ra. The first one is the possible different distribution of Th and Ra isotopes in the particle grains of an aquifer [19, 20]. Let us consider, for example, the ²³⁸U radioactive series. Insoluble ²³⁰Th, parent of soluble ²²⁶Ra, could be mainly localized within particle grains of aguifer and within newly formed insoluble Th-phases occurring as coatings on grains [21]. Thus, ²²⁶Ra emanation depends mainly on size of particle grain or on thickness of the newly formed coatings, at least when they are in the micronic range [22]. Differently, soluble ²²⁶Ra, parent of ²²²Rn, due to its prolonged adsorption from groundwater, is localized, more frequently than²³⁰Th, on the surface of the particle grain (Fig. 2b). In this case, the dependence of ²²²Rn emanation on size of particle grain or on coating thickness still holds (due to the dependence of surface area of the grain) but it is different from that of ²²⁶Ra.

The second factor, called "back diffusion," predicts that a-recoils damage the ordered structures of crystalline phases [11, 23]. Diffusion in radiation damage is in general faster than in ordered structures. Since Rn solid diffusion coefficient (Fig. 2c) is several orders higher than that of Ra (see Appendix 3), a further initial heavy bias is introduced when Ra retardation factor is calculated by comparison with Rn emanation. Some models include the effect of embedding (implantation of recoiled atom in the solid)

Fig. 2 Different models of radon and radium production (applied to ²³⁸U series). A: ²²⁶Ra and ²²²Rn production in nanopores. B: ²²⁶Ra and ²²²Rn production by direct recoil: ²²⁶Ra is distributed differently from ²³⁰Th covering also the grain surface. C: "back diffusion" of ²²²Rn inside the damage



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inside a nanopore as well [22, 24]. Others studies take in account the implantation in adjacent grains, recognizing its importance for Rn emanation in water-saturated porous media [25, 26]. In general, when parent nuclide is distributed on grain surface, the fraction f_{α} of α -recoiled radionuclides, topping in pore space but not implanting in adjacent grains, is defined by:

$$f_{\alpha} = \frac{W}{2 \cdot R_F} \tag{4}$$

where *w* is the pore size between grains and R_F is the range of recoil (nm) in the fluid [26, 27]. Thus, in a watersaturated medium, when Ra is distributed at the surface of submicron particles, the relationship between Rn emanation and particle diameter is no longer inverse, due to implantation effects. Actually, any attempt of modelling the Rn and Ra release by α -recoil effects in an unitary scheme is complicated and ultimately unrealistic. In our opinion, it is necessary to resort to experimental tools for establishing the real Γ_i and R_i . This is important both for the challenge of utilizing Ra isotopes for groundwater dating [28–30] and for the problem of radionuclide transport in groundwater because of possible release of nuclear waste products into aquifers [10].

1.1.2 Retardation factor

The retardation factor is connected to the dimensionless adsorption distribution coefficient K_d^* by:

$$R = K_d^* + 1 \tag{5}$$

 K_d^* is related to the dimensional adsorption distribution coefficient K_d (L·kg⁻¹) by [33]:

$$\mathsf{K}_d \cdot \frac{\rho_r}{\Phi} = \mathsf{K}_d^*. \tag{6}$$

In fresh waters, radium ion (Ra²⁺) approximates a conservative behaviour in fracturated aquifers as carbonatic or karst aguifers [28, 29] or in some volcanic aguifers where Ra-adsorbing mineral phases are scarce [31]. Generally, in fresh waters hosted in common fracturated crystalline rocks and in porous rock, Ra is strongly adsorbed [32], displaying an increasing retardation factor ($R_{Ra} > > 1$) with decreasing grain size [3]. Ra desorption and consequent decrease in retardation factor has been observed under conditions of: (i) increasing salinity due to processes of cation competition and displacement [33, 34], (ii) decreasing mineral surface charge and (iii) increasing stability of inorganic complexes [35]. In coastal aquifers and even in marine environment Ra adsorption is still appreciable with the exception of hypersaline environments where radium can be considered conservative $(R_{Ra} \sim 1)$ [7].

SN Applied Sciences A SPRINGER NATURE journal According to Krishnawami, at steady-state conditions and when the term $e^{-\lambda_i \cdot R_i \cdot T}$ in Eq. 1 trends to 0 (i.e. $\lambda_i T \approx 5$), retardation factors of ²²⁶Ra, ²²⁸Ra and ²²⁴Ra are equal to the ²²²Rn/²²⁶Ra, ²²⁰Rn/²²⁸Ra and ²²⁰Rn/²²⁴Ra activity ratios in groundwater, respectively [10, 36]. In terms of kinetics, the general equation for R_iof radioactive solutes is defined by [5, 10]:

$$R_{i} = 1 + \left[\frac{k_{s}}{\left(k_{d} + \lambda_{i}\right)}\right]$$
(7)

where k_s and k_d are the sorption and desorption rate constants (t^{-1}) of the radioactive solute *i*, respectively. This equation trends to the known equation for not radioactive solutes when λ_i trends to zero or when $\lambda_i < k_d$:

$$R_i = 1 + \left(\frac{k_s}{k_d}\right) = 1 + K_d^*$$
(8)

Returning to Eq. 2, for Ra isotopes of short mean life (²²³Ra, ²²⁴Ra) both dissolution constant and exponential term approximate to zero in few days [2]. Therefore, in the case of ²²⁴Ra in groundwater, it is realistic to admit that:

$$^{224} \operatorname{Ra}_{w} = \left(\frac{\Gamma_{224}}{\lambda_{224} \cdot \operatorname{R}_{\operatorname{Ra} - 224}}\right) \cdot \left(^{224} \operatorname{Ra}_{r} \cdot \frac{\rho_{r}}{\Phi}\right)$$
(9)

Combining Eq. 9 with Eq. 3:

$$\mathsf{R}_{\mathsf{Ra}-224} = \left(\frac{f_{a224}}{^{224}\mathsf{Ra}_{w}}\right) \cdot \left(^{224}\mathsf{Ra}_{r} \cdot \frac{\rho_{r}}{\Phi}\right) \tag{10}$$

Equations 7–10 are valid for groundwater in moderateto-low salinity conditions and near neutral pH.

Equation 10 is alternative to Eq. 7 for estimating R_{Ra-224}. As for R_{Ra-228} and R_{Ra-226}, being $\lambda_i < < k_d$, which is generally the case for ²²⁸Ra and ²²⁶Ra in groundwaters, [10, 29, 37]:

$$R_{Ra-228} \approx R_{Ra-226} \approx 1 + \left[\frac{k_s}{k_d}\right] \neq R_{Ra-224} = 1 + \left[\frac{k_s}{(k_d + \lambda_{224})}\right]$$
(11)

2 Estimation of retardation factor and recoil rate constant by "NAPL method"

Equations like Eqs. 7 and 10 are valid when the system is at a steady state and at radioactive equilibrium relatively to a particular ^{*i*}Ra. In this case, the sum of the production and desorption rate equals the sum of decay and adsorption rate [5]. For shallow aquifers, this condition is generally satisfied for Ra isotopes of shortest life (²²³Ra, ²²⁴Ra). Usually the Ra production rate in water is monitored by Rn which, however, has an emanation efficiency different from Ra.

This difficulty can be avoided, in case of porous aguifers, by using the "NAPL method." It consists in immersing two representative samples of aquifer rock in water and in a "non-aqueous phase liquid" (NAPL) and, after a comparison of measured Ra activities, in calculating the real f_{a} . This approach assumes that adsorption of inorganic ions on mineral grains in NAPLs is negligible. The most complete study on this topic is the one of Wainipee and coworkers [38]. They have shown as NAPLs coating reduces As (V) adsorption by decreasing the available surface area of clay minerals. In fact, polar organic compounds (as negatively charged carboxylic, naphthenic acids and positively charged amines, cationic surfactants, pyridines) contained in NAPLs, easily adsorb onto the mineral surfaces. Clay edges and asperities represent pinning points for NAPLs, being negatively or positively charged. Polar components of NAPLs with opposite charges, adsorbed onto the surface of the aquifer rock, provide surfaces of low interaction energy [39, 40]. In this way, saturation of adsorption sites and low mobility of the adsorbed organic polar compounds resemble what occurs in hypersaline aguifers where R_i trends to unity [6]. For a porous rock immersed in a non-polar liquid like NAPL, dissolution does not occur and, if adsorption is negligible, then R_i=1. In such a system, after a time of 20 days (required for reaching radioactive equilibrium between ²²⁴Ra and ²²⁸Th), Eq. 10 can be rewritten:

$$1 = \left(\frac{f_{a224}}{^{224}\mathsf{Ra}_{\mathsf{NAPL}}}\right) \cdot \left(^{224}\mathsf{Ra}_{r} \cdot \frac{\rho_{r}}{\Phi}\right)$$
(12)

where $^{224}\text{Ra}_{\text{NAPL}}$ (Bq·L $^{-1}$) is the ^{224}Ra activity measured in NAPL; therefore:

$$^{224}\text{Ra}_{\text{NAPL}} = f_{\alpha 224} \cdot \left({}^{224}\text{Ra}_r \cdot \frac{\rho_r}{\Phi} \right)$$
(13)

Experimentally, $f_{\alpha 224}$ is found by:

$$f_{\alpha 224} = \frac{^{224}\text{Ra}_{\text{NAPL}} \cdot \text{L of NAPL in immersion test}}{^{224}\text{Ra}_{\text{r}} \cdot \text{kg of rock in immersion test}}$$
(14)

Then f_{a224} through Eq. 3, gives the value of Γ_{224} .

By an analogue immersion test of 20 days, using water instead of NAPL, ²²⁴Ra_w can be measured for values of ρ_r and Φ equal to those of NAPL immersion. Thus, if all ²²⁴Ra which recoils from solid to fluid during the immersion test, is desorbed from solid surface to NAPL (*i.e.* The dimensionless distribution coefficient of ²²⁴Ra between solid and NAPL is equal to 0), R_{Ra-224} can be easily determined by substituting Eq. 13 in Eq. 10:

$$R_{Ra-224} = \frac{{}^{224}Ra_{NAPL}}{{}^{224}Ra_{w}}$$
(15)

In this paper, the experimental way to estimate R_{Ra-224} and Γ_{224} by using Eqs. 13–15 will be called "NAPL method". Figure 3 shows the standard procedure for obtaining R_i and Γ_i of Ra isotopes from representative samples of porous aquifer rocks. The essential for reproducibility of the method is that the duration of the immersion and the ratio between the weight of sample and the volume of liquid be the same in the two immersion tests (NAPL and water). Although R_{Ra-224} and Γ_{224} can be estimated also theoretically, their determination in laboratory improves the application of Ra isotopes to groundwater dating.

Now, how to find the retardation factor for ²²⁸Ra and ²²⁶Ra? Theoretically, the R₂₂₈/R₂₂₄ ratio varies between 1 and the $\lambda_{224}/\lambda_{228}$ ratio (~ 273) [10]. In the real world, this range could be narrower, between 1 and 3, as shown in Appendix 4.

Let us consider now the values of R_{Ra-228} and R_{Ra-224} calculated from values of k_d and k_s available in the literature and obtained by the Krishnawami method (25 groundwaters of 12 different aquifers, Fig. 4). These data can be used to estimate the R_{Ra-228}/R_{Ra-224} ratio. When plotted together, the best fitting through origin gives a straight line whose slope, equal to the R_{Ra-228}/R_{Ra-224} ratio, is 1.75 ± 0.21. R_{Ra-228} and R_{Ra-226} therefore are estimated to about 1.75 times R_{Ra-224}. It should seem strange that data derived from a questioned method obtain a valid relationship between the R_is of radium. However, if absolute values of R_is depend strictly on the used method, it is not so for their ratio. Any R_iin fact is equal to the ratio of the production rate (monitored by Rn) and the expected content of the Ra isotope i at secular equilibrium [41]. Therefore, when R_is are divided each other, eventual biases of the production rates, measured using Rn, cancel together.

Alternatively, the NAPL method can be also applied directly to ²²⁸Ra, by immersing two samples representative of aquifer rock in NAPL and in water for a time sufficient to detect α -recoiled ²²⁸Ra (some months to 2 years, depending on analytical sensitivity). In this case ²²⁸Ra does not achieve the radioactive equilibrium conditions and R_{Ra -228} is easily obtained by solving (e.g. graphically) the following nonlinear equation, where ²²⁸Ra_{NAPL}, R_{Ra -228}, λ_{228} and *T* are known:

$$\frac{^{228}\text{Ra}_{\text{NAPL}}}{^{228}\text{Ra}_{w}} = \frac{\text{R}_{\text{Ra}-228} \cdot (1 - e^{-\lambda_{228} \cdot T})}{(1 - e^{-\lambda_{228} \cdot \text{R}_{\text{Ra}-228} \cdot T})}$$
(16)



Fig. 3 Standard procedure for determining $R_{Ra\text{-}224}$ and Γ_{224} of radium by NAPL method

3 Study site and application to groundwater dating by Ra isotopes: Red Pozzolane aquifer (Alban Hills Volcano, Rome)

and by dividing Eqs. 17–18 each other and remembering that $R_{Ra-228} = R_{Ra-226'}$ it follows:

The equations describing the temporal evolution of ²²⁸Ra_w and ²²⁶Ra_w (Eq. 1 for *i* = 226,228) can be expressed as:

$$\left[{}^{228}\mathsf{Ra}_{w} - \left({}^{228}\mathsf{Ra}_{w(0)} \cdot e^{-\lambda_{228} \cdot \mathsf{R}_{\mathsf{Ra}-228} \cdot T}\right)\right] = \left[\frac{\left(\Delta_{228} + \Gamma_{228}\right)}{\left(\lambda_{228} \cdot \mathsf{R}_{\mathsf{Ra}-228}\right)}\right] \cdot \left(1 - e^{-\lambda_{228} \cdot \mathsf{R}_{\mathsf{Ra}-228} \cdot T}\right) \cdot \left({}^{228}\mathsf{Ra}_{r} \cdot \frac{\rho_{r}}{\Phi}\right)$$
(17)

$$\left[{}^{226}\mathsf{Ra}_{w} - \left({}^{226}\mathsf{Ra}_{w(0)} \cdot e^{-\lambda_{226} \cdot \mathsf{R}_{Ra-226} \cdot \tau}\right)\right] = \left[\frac{\left(\Delta_{226} + \Gamma_{226}\right)}{\left(\lambda_{226} \cdot \mathsf{R}_{Ra-226}\right)}\right] \cdot \left(1 - e^{-\lambda_{226} \cdot \mathsf{R}_{Ra-226} \cdot \tau}\right) \cdot \left({}^{226}\mathsf{Ra}_{r} \cdot \frac{\rho_{r}}{\Phi}\right)$$
(18)

$\left[\frac{{}^{228}\text{Ra}_w - \left({}^{228}\text{Ra}_{w(0)} \cdot e^{-\lambda_{228} \cdot \text{R}_{\text{Ra}-228} \cdot T}\right)}{{}^{226}\text{Ra}_w - \left({}^{226}\text{Ra}_{w(0)} \cdot e^{-\lambda_{226} \cdot \text{R}_{\text{Ra}-228} \cdot T}\right)}\right] \cdot \frac{{}^{226}\text{Ra}_r}{{}^{228}\text{Ra}_r}$	$= \left[\frac{\left(\Delta_{228} + \Gamma_{228}\right)}{\left(\Delta_{226} + \Gamma_{226}\right)}\right] \cdot \frac{\lambda_{226}}{\lambda_{228}} \cdot \frac{\left(1 - e^{-\lambda_{228} \cdot R_{\text{Ra}-228} \cdot T}\right)}{\left(1 - e^{-\lambda_{226} \cdot R_{\text{Ra}-228} \cdot T}\right)}$	(19)
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Fig. 4 Radium retardation factors from literature: R_{Ra-228} and R_{Ra-224} in groundwater samples from 12 different aquifers. Samples from aquifers 1–7: data in [4]; from aquifers 8–10: data in [10]; from aquifers 11–12: data in [36, 37]

Equation 19 has the great advantage of not being dependent on Φ and ρ_r of the aquifer rock. Assuming $\Delta_{228} = \Delta_{226}$ (congruent dissolution), three typical cases for Eqs. 17–19 are developed (Eqs. 41–52, Appendix 5).

The R_{Ba-228} , R_{Ba-226} and the Γ_{228} , Γ_{226} found by the NAPL method will be used to estimate the residence time of groundwater hosted in Red Pozzolane (RP) aquifer of the volcanic region of Alban Hills, Rome (Fig. 5). The RP hydrogeological complex consists of massive and chaotic pyroclastic with mean thickness of 20 m. It is characterized by a medium to medium-high permeability due to a high porosity locally diminished by zeolitization. The complex constitutes the main aquifer of Rome, particularly in the area on the left bank of Tiber river, where it is spread with continuous and relevant thicknesses [42]. Scoria clasts in the RP flow deposits are scarcely porphyritic containing millimetre-sized fresh leucite with subordinate clinopyroxene. Primary leucite is often replaced with analcime [43]. The groundmass is largely formed by glass turned to zeolites and/or halloysite and by leucite microcrysts with star-like habit [44]. The abundance of leucite required a study of Ra adsorption onto this mineral. As application of



Fig. 5 Geological section and hydrogeological setting of the Red Pozzolane Aquifer (Alban Hills, Rome)

the NAPL method, we describe an attempt of measuring the age of groundwater from a piezometer well and from a mineral spring located in proximity of the peripheral part of the PR hydrogeological complex. Initially, however, some experimental evidence of the validity of the assumptions of the NAPL method will be provided.

4 Materials, methods and experimental procedures

Tests on different materials, normally adsorbing Ra from aqueous solutions, were carried out to validate the assumption of negligible effects of Ra adsorption at the interface between NAPL and mineral grains. Manganese dioxide, zeolite 4A and natural clays were selected for Ra adsorption experiments using aqueous and organic solutions enriched in radium isotopes. Monazitic sand and a natural Th-enriched weathered volcanic rock were selected for testing the reliability of assuming retardation factors equal to unity for Ra and even for Th ions when aquifer rock is immersed in NAPLS for a long time. Finally Ra adsorption experiments on leucite, carried out in IGAGlab in the past and never published, are presented as an aid for discussing experimental data obtained for groundwater of RP aquifer.

The size distribution of samples was obtained in wet mode by a laser particle size analyser, model Sympatec Helos, equipped with He–Ne laser source (Supplementary material).

XRPD analysis was carried out with a fully automated Bruker AXS D8 advance diffractometer, operating in reflection mode with θ - θ geometry, equipped with high-resolution energy-dispersive 1-D Linxeye XE detector. Measurement parameters used were: Cu Ka, 40 kV, 30 mA, scan angle (2θ) = 5°-80°; step width (2θ) = 0.020°.

Determination of Ra isotopes in solids and liquids was carried out by high-resolution γ-spectrometry with an EG&G Ortec solid-state photon detector (HPGe) and Coaxial detector system, following procedures and standardization methods described in a previous paper [2]. Samples were sealed in Rn-tight containers (to avoid Rn escaping). Ra isotopes were measured from the y-emission of their daughters after waiting the conditions of radioactive equilibrium. In particular 226 Ra, 228 Ra and 224 Ra were determined from γ -emission of 214 Pb, 228 Ac and 212 Pb, respectively. Radioactive equilibrium between ²²⁶Ra and ²¹⁴Pb was reached in 20 days, while between ²²⁸Ra and ²²⁸Ac in 1 day and between ²²⁴Ra and ²¹²Pb in 2 days. ²²⁴Ra was determined after correction for decay occurring between separation time and initial counting and during y-counting. In some experiments, ²¹²Pb was also measured directly for determining the adsorption of lead.

 ^{220}Rn was measured following the method named "Rad H_2O " using a continuous Rn Gas Monitor model Rad 7, Durridge Company, equipped with a solid-state ion-implanted planar silicon detector [45].

Kerosene (Alfa Aesar) and pure n-hexane (Merck) were chosen as apolar fluids representative of light NAPLs. They display similar apolarity but only kerosene, due to its complex nature, contains also organic polar compounds (e.g. naphtenic acids). They differ notably by volatility and by cost as well. Both were previously checked by γ -spectrometry, for excluding the detectable presence of Ra isotopes.

Th-Welsbach mantles [46] were immersed for 20 days in 50 ml of kerosene obtaining a ²²⁴Ra-enriched kerosene phase (K20D) and for 2 years in 50 ml of kerosene obtaining a ²²⁸Ra-, ²²⁴Ra-enriched kerosene phase (K2Y). From aqueous Ra solution, obtained by chromatographic separation of Ra from standard Th solution (Spectrascan Analytical Standards), ²²⁸Ra was extracted by synergistic complexation with TTA (Thenoyltrifluoroacetone) and TOPO (Trioctylphosphine oxide) in n-hexane [47] producing an organic phase (RaC) containing complexed radium. K20D, K2Y and RC were used in experiments after measurement of their Ra activity by γ-spectrometry. Table 1 resumes the plan of the experiments carried out in this work.

4.1 Ra adsorption on MnO_2 in kerosene and recoil constant of ²²⁴Ra recoiled from ²²⁸Th adsorbed on MnO_2

MnO₂ was selected for its known capacity of adsorbing bivalent cations in aqueous solutions [48]. Thus, it can be useful to compare Ra adsorption in water and in NAPL. The adsorption kinetics in water is fast: after 30 min about 90% of Pb and 50% of Ra are removed by adsorption on MnO₂ [48–50]. The < 0.125 mm fraction of manganese dioxide black powder (Carlo Erba, BET surface area: 4 m²g⁻¹), MDBP, was selected for Ra adsorption experiments. Particle bulk density ρ_r and porosity Φ of MDBP were evaluated according standard procedures [51].

4.1.1 Experiments MDBP-1 K and MDBP-1H

This experiment was carried out to test the adsorption of bivalent simple ions like Ra^{2+} and Pb^{2+} when MnO_2 is immersed in NAPL (kerosene and n-hexane). 10 ml of K2Y were added to 0.1 g of MDBP (particle /liquid ratio = 0.01 kg·L⁻¹). After agitation for 30 min, the organic phase (MDBP-1 K) was separated by centrifugation and suddenly counted for ²¹²Pb while ²²⁸Ra was detected after 1 day.10 ml of n-hexane were added to 0.1 ml of liquid resulting from evaporation of 10 ml of K2Y and counted for ²²⁸Ra and ²¹²Pb. Then this (²²⁸Ra,²¹²Pb)-enriched

SAMPLE	EXPERIMENT	CONTACT TIME	ANALYSED PHASE	DETECTED ISOTOPES	MAIN RESULTS AND DERIVED PARAMETERS
MDBP-1 K	MDBP + K2Y	30 min	Ker	²¹² Pb, ²²⁸ Ra	moder. Pb ads./ moder. Ra ads
MDBP-1H	MDBP + K2Y + Hex	30 min	Hex+(Ker)	²¹² Pb, ²²⁸ Ra	neglig. Pb ads./ neglig. Ra ads
MDBP-2	MDBP + RaC + Ker	30 min	Ker + Hex	²²⁸ Ra-complex	high Ra adsorption
MDBP-3a	Th-MDBP + Ker	20 days	Ker	²²⁴ Ra	f_{aa224} - Γ_{Ra} (NAPL method)
MDBP-3b	Th-MDBP+Water	5 min	Water	²²⁰ Rn	$f_{a220} - \Gamma_{Rn-220}$
Z4A-1	Z4A + K20D	2 days	Ker	²²⁴ Ra	negligible Ra adsorption
Z4A-2	Z4A + RaC + Ker	2 days	Ker + Hex	²²⁸ Ra-complex	moderate Ra adsorption
Z4A-3	Z4A + ²²⁸ Ra-Water	2 days	Water	²²⁸ Ra	high Ra adsorption
S8N1-1 K	S8N1+K20D	1 days	Ker	²²⁴ Ra	negligible Ra adsorption
S8N1-3 K	S8N1+K20D	3 days	Ker	²²⁴ Ra	negligible Ra adsorption
S8N1-3 W	S8N1+ ²²⁸ Ra-Water	3 days	Water	²²⁸ Ra	high Ra adsorption
MON-1 K	monazite + Ker	2 days	Ker	²²⁴ Ra	f_{a224} - Γ_{Ba} (NAPL method)
MON-1 K	monazite + Ker	27 months	Ker	²²⁸ Th, ²²⁸ Ra	Treal≡Tmeasured(<i>i.e.</i> noadsorption)
CMB-1 K	CM + Ker	2 years	Ker	²²⁸ Ra	R _{Ra-228} (NAPL method)
CMB-1 W	CM+Water	2 years	Water	²²⁸ Ra	Treal \equiv Tmeasured (<i>i.e.</i> no adsorption R_{Ra-228} (NAPL method)
CM-1aK	CM-1 + Ker	20 days	Ker	²²⁴ Ra	R _{Ra-224} (NAPL method)
CM-1bW	CM-1+Water	20 days	Water	²²⁴ Ra	R _{Ra-224} (NAPL method)
CM-1cW	CM-1+Water	5 min	Water	²²⁰ Rn	f ₂₂₀ -Γ _{Rn-220}
RP-2 K	RP-2 + Ker	20 days	Ker	²²⁴ Ra	R _{Ra-224} (NAPL method)
RP-2 W	RP-2+Water	20 days	Water	²²⁴ Ra	R _{Ra-224} (NAPL method)
RP-2R	RP-2+Water	5 min	Water	²²⁰ Rn	f ₂₂₀ -Γ _{Rn-220}

Ker = Kerosene; moder. = moderate; neglig. = negligible

hexane was added to 0.1 g of MDBP (particle /liquid ratio = 0.0099 kg \cdot L⁻¹) and stirred for 30 min. Finally the organic phase (MDBP-1H) was separated and counted for ²¹²Pb. ²²⁸Ra was.

detected after 1 day.

Table 1 Experiment plan

4.1.2 Experiment MDBP-2

The experiment is aimed to highlight the eventual different adsorption on MnO_2 when Ra occurs as organic complex (RaC: TOPO-TTA Ra complex). RaC organic phase was mixed with kerosene (1:2 volume ratio) and counted after 1 day for detecting initial ²²⁸Ra. Then 15 ml of this mixture were added to 0.1 g of MDBP (particle /liquid ratio = 0.0067 \cdot 10⁻² kg ·L⁻¹) and, after 30 min of agitation, the organic phase (MDBP-2) was separated and counted for ²²⁸Ra after 1 day.

4.1.3 Experiment MDPB-3

The following experiment was used to test whether the radon emanation of ²²⁰Rn, f_{220} , approximates the alpha recoil fraction of ²²⁴Ra, f_{a-224} , measured by NAPL method, when parent radionuclides (²²⁴Ra and ²²⁸Th, respectively)

are distributed in the same way exclusively as adsorbed atoms on particle grains. 10 ml of a standard Th solution in nitric acid (1000 ppm of ²³²Th) were slowly dried on 16.0 g of MDBP by continuous stirring at T < 80 °C. The HNO₃ treatment greatly increases the surface area [52]. In these conditions, Th nitrate (density = 2.8 $q \cdot cm^{-3}$) should cover the MnO₂ grains (specific surface = $4 \cdot 10^4$ cm² g⁻¹) with a thickness < 0.2 nm. This is a trascurable fraction of the R_c in Th nitrate (R_c = 30 nm). In this way ²²⁸Th occurs only as adsorbed ion, allowing ²²⁴Ra-recoil exclusively from the surface of grains. At dryness, 8 g of Th-enriched MDBP were immersed in 30 ml of kerosene (particle /liguid ratio = 0.27 kg \cdot L⁻¹) up to 20 days to reach radioactive equilibrium between surfacial ²²⁸Th and recoiled ²²⁴Ra. Then kerosene (MDBP-3a) was separated by centrifugation and filtration and counted after 2 days for ²²⁴Ra. Another subsample of Th-enriched MDBP (8 g) was immersed in 30 ml of water (particle /liquid ratio = 0.27 kg \cdot L⁻¹) and ²²⁰Rn activity in water (MDBP-3b) was measured.

4.2 Ra adsorption on zeolite 4A

The importance of these experiments relies on the evidence that in many volcanic aquifers, like RP aquifer, zeolites are important Ra-sinking mineral phases. We have used zeolites of A family, known as LTA (Linde Type A), which are extensively used as standard Ra adsorbers [53, 54]. Zeolite 4A (Bis Italia s.r.l.) was used in these experiments. In zeolite 4A, negative charges, formed in the tetrahedral framework, are compensated by exchangeable Na⁺.

The first experiment (Z4A-1) tests Ra adsorption (as simple ion) on zeolite immersed in NAPL. The second experiments (Z4A-2) tests adsorption on zeolite in NAPL when Ra occurs as organic complex (RaC: TOPO-TTA Ra complex). A third experiment (Z4A-3) was carried out to compare Ra adsorption (as simple ion) on zeolite when immersed in water.

Ten ml of K20D were added to 2 g of zeolite 4A (particle /liquid ratio = 0.2 kg \cdot L⁻¹) and stirred for 2 days. Then kerosene (Z4A-1) was separated and after 2 days was counted for ²²⁴Ra.

5 ml of RaC were mixed with 20 ml of kerosene and counted after 1 day for detecting ²²⁸Ra. Then the resulting organic phase was added to 2 g of zeolite 4A (particle /liquid ratio = 0.08 kg \cdot L⁻¹). After 2 days of stirring, the organic phase (Z4A-2) was separated and, after 1 day, counted for ²²⁸Ra.

20 ml of water solution (pH 5.5), containing a known amount of 228 Ra, obtained by chemical separation from Th standard solution, was added to 2 g of zeolite 4A (particle /liquid ratio=0.1 kg·L⁻¹) and stirred for 2 days. Then water (Z4A-3) was separated and, after 1 day, measured for 228 Ra.

4.3 Ra adsorption on clay

The pliocenic clay of Dunarobba fossil forest [55, 56], containing numerous clay minerals as montmorillonite, saponite, mixed layered chlorite-montmorillonite, illite, chlorite and kaolinite, was chosen for adsorption experi-

liquid ratio = 0.03 kg \cdot L⁻¹). (ml/g = 33.3), were stirred for 1 and 3 days, respectively. Then kerosene phases were separated by centrifugation and counted for ²²⁴Ra after 2 days.

25 ml of an aqueous solution (pH 7.0) containing a known amount of 228 Ra (3.8·10³ Bq·L⁻¹), obtained by chemical separation from Th standard solution, was added to 0.75 g of S8N1 (particle /liquid ratio = 0.03 kg·L⁻¹) and stirred for 3 days. Then water (S8N1-3 W) was separated by centrifugation and, after 1 day, 228 Ra activity was measured.

4.4 Monazite in contact with NAPL

This experiment was carried out to test the validity of assuming R_{Ra} and R_{Th} in NAPL (kerosene) equal to 1, by measuring the ²²⁸Th/²²⁸Ra age of contact between monazite and NAPL.

10 ml of kerosene was added to 35 g of pure monazitic sand (particle /liquid ratio = 0.29 kg L⁻¹). After 20 days of stirring, kerosene (MON⁻¹ K) was separated and after 2 days ²²⁴Ra activity was measured for determining f_{a224} . After the γ -counting, kerosene was newly put in contact with the monazitic sand up to 27 months. Then kerosene was separated and, after 2 weeks, ²²⁸Ra and ²²⁸Th (the latter one through ²²⁴Ra in secular equilibrium with ²²⁸Th) were measured. Monazite does not dissolve in kerosene, thus the ²²⁸Th/²²⁸Ra ratio in kerosene depends only on α -recoil phenomena. The temporal variation of this radioactive pair is described by the following equation when retardation factor of Th and Ra are both equal to 1 [58]:

$$\frac{{}^{228}\text{Th}}{{}^{228}\text{Ra}} = \frac{1 - e^{-\lambda_{\text{Th}}} - 228 \cdot T} - \left[\left(\frac{\lambda_{\text{Th}}}{\lambda_{\text{Th}}} - 228} \frac{\lambda_{\text{Th}}}{\lambda_{\text{Th}}} \right) \cdot \left(e^{-\lambda_{\text{Ra}}} - 228 \cdot T} - e^{-\lambda_{\text{Th}}} - 228 \cdot T} \right) \right]}{1 - e^{-\lambda_{\text{Ra}}} - 228 \cdot T}$$
(20)

Instead, considering explicitly retardation factors of Th and Ra, Eq. 20 becomes:

$$\frac{^{228}\text{Th}}{^{228}\text{Ra}} = \frac{R_{\text{Ra}}}{R_{\text{Th}}} \cdot \frac{1 - e^{-\lambda_{\text{Th}-228} \cdot \mathcal{T} \cdot R_{\text{Th}}} - \left[\left(\frac{\lambda_{\text{Th}-228} \cdot R_{\text{Th}}}{\lambda_{\text{Th}-228} \cdot R_{\text{Th}} - \lambda_{\text{Ra}-228} \cdot \mathcal{T} \cdot R_{\text{Ra}}} \right) \cdot \left(e^{-\lambda_{\text{Ra}-228} \cdot \mathcal{T} \cdot R_{\text{Ra}}} - e^{-\lambda_{\text{Th}-228} \cdot \mathcal{T} \cdot R_{\text{Th}}} \right) \right]}{1 - e^{-\lambda_{\text{Ra}-228} \cdot \mathcal{T} \cdot R_{\text{Ra}}}}$$
(21)

ments on clay. The sample S8N1 (60–120 mesh fraction), with a low content of organic matter (3%), was selected for its negligible 224 Ra activity (0.01 Bq g⁻¹) in secular equilibrium with 228 Th [57].

Samples S8N1-1 K and S8N1-3 K, prepared in the same way, by adding 25 ml of K20D to 0.75 g of S8N1 (particle /

At secular equilibrium, Eq. 21 predicts that $^{228}Th/^{228}Ra=R_{Ra}/\,R_{Th}$ [59].

It's clear that only if both R_{Ra} and R_{Th} in kerosene approximate to 1, the age of MON-1 K, calculated by Eq. 20, corresponds to the true one (27 months).

4.5 Application of the NAPL method for determining Ra retardation factor and recoil constant for a Th-enriched weathered volcanic rock

A Th-enriched, highly porous pyroclastic rock, composed by diatomite beds mixed with deeply altered leucititictephrite in the guarry of Casale Morticini (CM) near Montefiascone (Vulsini volcano, Latium, Italy) was selected. In this area, subduction-related metasomatic enrichment of incompatible elements in the mantle source, coupled with magma differentiation within the upper crust, gave rise to U, Th and K-enriched melts [60]. Alteration by CO₂ and H₂S-rich gaseous exhalation resulted in abundant amorphous silica with scarce halloysite 10 Å and allophane [61]. The XRD spectrum of the sample (see Supplementary material) reveals hydrothermal quartz and jarosite among main minerals, as in similar degassing areas of Latium [62]. A broad halo between 20°-30° (20) indicates that the bulk of material is amorphous. The ²³²Th activity of CM, measured by high resolution y-spectrometry, is close to 1,980 Bg kg⁻¹ (500 ppm Th), about five times the average of Vulsini volcanic products [60].

4.5.1 Experiments CMB-1 K and CMB-1 W

This experiment tests the capacity of NAPL method to find directly R_{Ra-228} and the validity of assuming R_{Ra} and R_{Th} in kerosene equal to 1 by measuring the resulting ²²⁸Th/²²⁸Ra age of contact between volcanic rock and NAPL. CM was dryed at 105 °C and splitted in two portions of 50 g. They were saturated and rinsed with 225 ml of kerosene (CMB-1 K) and distilled water (CMB-1 W), respectively, using flasks of 250 cc (particle /liquid ratio = 0.222 kg $\cdot L^{-1}$). After 2 years, the liquid phases were separated and their ²²⁸Ra activities determined. The ²²⁸Th activity of sample CMB-1K was also measured.

4.5.2 Experiments CM-1aK, CM-1bW, CM-1cW

These experiments have the objective to compare the NAPL method with the Rn method in determining Ra retardation factor. NAPL method on ²²⁴Ra was carried out on subsample (CM-1), obtained by sieving a pulverized sample of CM and selecting the < 125 μ fraction (120 ASTM mesh). This for working with a sample having a more defined granulometry and also higher in thorium, which generally enriches in fine fraction. The fraction displayed a measured ²³²Th activity of 2,900 Bq kg⁻¹ (732 ppm). Particle bulk density, porosity and distribution size of CM-1 (see Supplementary material) were determined and 18 g of CM-1 were immersed in 41 ml of kerosene (particle /liquid ratio = 0.44 kg ·L⁻¹) for 20 days. Successively kerosene was

separated (CM-1aK) and after 2 days the ²²⁴Ra_{NAPL} activity was measured. A portion of CM-1 (18 g) was immersed in 41 ml of distilled water (particle /liquid ratio = 0.44 kg ·L⁻¹) for 20 days. Then water was separated (CM-1bW) and detected for ²²⁴Ra_w after 2 days. Finally, 18 g of CM-1 were immersed in 250 ml of distilled water (CM-1cW, particle / liquid ratio = 0.072 kg ·L⁻¹) and the ²²⁰Rn emanation coefficient (f_{220}) in water was measured.

4.6 Ra adsorption onto leucite

²²⁶Ra adsorption on leucite at ambient temperature was studied. A solution of known ²²⁶Ra activity, deriving from extraction and chemical separation of uraninite (Black Hills), was used for adsorption experiments. Five samples, each one composed by several leucite crystals from Pozzolanelle Unit of Alban Hills volcanic district [63] were selected on the basis of their size. The composite total surface area of each sample was determined. They were immersed in different volumes of an aqueous ²²⁶Ra-solution of known activity at pH 7.0 (T = 298°K) and stirred for 1 week. Then aqueous solutions were separated and kept in sealed containers for 20 days before of determining the desorbed ²²⁶Ra activity.

4.7 Study site: Ra isotopes in groundwater from the Red Pozzolane aquifer (Alban Hills)

40 L of groundwater (RP-1) of the aquifer hosted in the RP hydrological unity were extracted from a piezometric well (Fig. 5) located near the Arco di Travertino (Appia road, Rome). Groundwater was sampled at a depth of 12 m from ground level and about 1.5 m below the local water table.

20 g of Ba(NO3)₂ and a calculated excess of H₂SO₄ were added to RP-1 for precipitating (Ba,Ra) sulphate. The efficiency of precipitation was obtained by the weight ratio of recovered BaSO₄ to the added amount of Ba(NO3)₂ [64]. Dry powder from precipitation was measured after 2 days for detecting ²²⁴Ra and ²²⁸Ra. The powder was then sealed in an air-tight container and, after 20 days, ²²⁶Ra was measured.

192 g of RP (RP-2) were measured for Ra isotopes by γ -spectrometry. Then two portions of RP-2 were treated by NAPL method using kerosene (RP-2 K) and water (RP-2 W) (ml /g = 0.36). A third aliquot of RP-2 was immersed in a 250-ml flask filled with distilled water (RP-2R) and the ²²⁰Rn emanation in water (f_{220}) was measured. Since the high hydrolytic reactivity of RP formation is well known [65], its dissolution at the first stage of contact with water was studied experimentally. RP was reacted with Millipore water (ml/g = 2.5) at 1 atm in a constant-temperature water bath at 298 °K for a total time of 45 days. The mixture was agitated magnetically at intervals of 1 h (during the first

Table 2 Adso	rption experimer	nts								
SAMPLE	Bulk density	Porosity	<u>A ads(Bq)</u> A des (Bq)	Liquid (L) particle (kg) real	Liquid (L) particle (kg) experimental	K_d (L·kg ⁻¹) n=0= K_d_{exp} n=0.5= K_d_{real}	$\begin{array}{c} K_d^* \\ n=0=K_{dexp}^* \\ n=0.5=K_{dreal}^* \end{array}$	$\begin{array}{l} R_{H20} \\ n = 0 = R_{H20 \ exp} \\ n = 0.5 = R_{H20 \ real} \end{array}$	$\begin{array}{l} R_{NAPL} \\ n=0=R_{NAPL\ exp} \\ n=0.5=R_{NAPL\ real} \end{array}$	Radionuclide
MDBP-1 K	0.48	06.0	0. 33± <i>0.01</i>	1.88	100	33.0±1.0 4.50±0.13	17.5±0.5 2.39±0.07		18. 5 ± 0.5 3.39 ± 0.07	²¹² Pb simple ion
MDBP-1H	0.48	06.0	0.03± <i>0.02</i>	1.88	101	3.0±2.0 0.41±0.27	1.6±1.1 0.22±0.15		2.6±1.1 1.22±0.15	²¹² Pb simple ion
MDBP-1 K	0.48	06.0	0.14± <i>0.0</i> 3	1.88	100	14.0± <i>3.0</i> 1.92±0.41	7.4±1.6 1.02±0.21		8.4 ± 1.6 2.02 ± 0.21	²²⁸ Ra simple ion
MDBP-1H	0.48	06.0	0.03± <i>0.02</i>	1.88	101	3.0±2.0 0.41±0.27	1.6±1.1 0.22±0.15		2.6±1.1 1.22±0.15	²²⁸ Ra simple ion
MDBP-2	0.48	06.0	1.18±0.07	1.88	149	175.8±10.4 19.7±1.2	93.5 ± 5.5 10.5 ± 0.6		94.5 ± 5.5 11.5 ± 0.6	²²⁸ Ra ion complex
MnO ₂ litera- ture values			1-10 *			100-1000 **				²²⁶ Ra simple ion
Z4A-1	0.70	0.63	0.02± <i>0.03</i>	06.0	Ŋ	0.10 ± 0.15 0.04 ± 0.06	0.11 ± 0.17 0.04 ± 0.07		1.11 ± 0.17 1.04 ± 0.07	²²⁴ Ra simple ion
Z4A-2	0.70	0.63	0.32± <i>0.02</i>	06.0	12.5	4.0±0.3 1.07±0.07	4.4±0.3 1.19±0.07		5.4 ± 0.3 2.19 ± 0.07	²²⁸ Ra ion complex
Z4A-3	0.70	0.63	4.5 ±0.7	06.0	10	45.0±7.0 13.5±2.1	50.0± <i>8.0</i> 15.0±2.3	51.0±8.0 16.0±2.3		²²⁸ Ra simple ion
S8N1-1 K	1.16	0.43	0.03± <i>0.02</i>	0.37	33.3	0.99± <i>0.66</i> 0.10± <i>0.07</i>	2.67 ± 1.78 0.27 ± 0.18		3.67 ± 1.78 1.27 ± 0.18	²²⁴ Ra simple ion
S8N1-3 K	1.16	0.43	0.03± <i>0.03</i>	0.37	33.3	0.99± <i>0.99</i> 0.10±0.10	2.67 ± 2.67 0.27 ± 0.27		3.67 ± 3.67 1.27 ± 0.27	²²⁴ Ra simple ion
S8N1-3 W	1.16	0.43	2.34±0.23	0.37	33.3	77.9±7.7 8.21±0.82	210.5 ± <i>21.0</i> 22.2 ± 2.2	211.5± <i>21.0</i> 23.2±2.2		²²⁸ Ra simple ion
(*) Literature	/alues from refer	ences: [48–50]; (**) Literature	values from 1	reference: [33]					

Table 7 Ade

SN Applied Sciences A SPRINGER NATURE journal day) and successively of 24 h. The extent of initial dissolution of the rock was monitored by a pH meter (pH340-A, WTW, Germany) and an electrical conductivity meter (LF340, WTW, Germany). The probes were kept in continuous contact with solution and values recorded periodically. Electrical conductivity was converted, after calibration, in terms of TDS of solution and expressed as TDSL of pore water according the formula:

TDSL = TDS of solution $(250/100) \cdot (\rho_d/\Phi)$.

5 Results and discussion

Results of adsorption experiments are reported in Table 2. The latter ones are expressed in terms of the dimensional and dimensionless adsorption distribution coefficients, $K_{d exp}$ (L·kg⁻¹) and $K_{d exp}^*$, respectively. $K_{d exp}$ was calculated by the following equation:

$$K_{dexp} = \frac{\text{Activity adsorbed(Bq)}}{\text{Activity desorbed (Bq)}} \cdot \left(\frac{\text{Liquid (L)}}{\text{particle (kg)}}\right)_{exp} \quad (22)$$

where $\left(\frac{\text{Liquid}(L)}{\text{particle}(kg)}\right)_{\text{exp}}$ is the ratio between liquid (water or NAPL) to solid during the experiments given in Sect. 4.1, whereas K_{dexp} is calculated by Eq. 6. The $\left(\frac{\text{Liquid}(L)}{\text{particle}(kg)}\right)_{exp}$ ratio is widely recognized as a critical component in determining retardation factors of Ra as evidenced by many authors [33, 66, 67] that have discussed the sediment–water ratio artifact or particle concentration effect. The particle concentration effect predicts that experimental K_d differs from real K_d when $\left(\frac{\text{Liquid}(L)}{\text{particle}(kg)}\right)_{exp}$ ratio differs from real $\left(\frac{\text{Liquid}(L)}{\text{particle}(kg)}\right)_{real}$ ratio (equal to the real $\frac{\Phi}{\rho_r}$ ratio) following a kind of scaling law [68, 69] that can be expressed as:

$$K_{d \text{ real}} = \left[\frac{\left(\frac{\text{Liquid}}{\text{Particle}}\right)_{\text{real}}}{\left(\frac{\text{Liquid}}{\text{Particle}}\right)_{\text{exp}}} \right]^{n} \cdot K_{d \text{ exp}}$$
(23)

where $n \ge 0$ is an empirical parameter called the particle effect index which measures the degree of the particle concentration effect. In Table 3, $K_{d real}$ and $K_{d real}^*$ values and real retardation factors calculated applying Eq. 5 are given for n = 0.5. The value of n = 0.5 fits very well the Ra adsorption experiments on sands provided by Beck and Cochran [33]. This value is close to the particle effect index found for Zn ²⁺on manganite (0.41) [69]. Such a value of particle effect index (0.5) corresponds to a dependence of adsorption on the square root of the liquid to particle ratio, and it is assumed in this work as the most reliable value for correcting distribution coefficients (expressed as $K_{d real}^*$) and retardation factors (expressed as R_{real}) for the particle concentration effect. Results of radium analitical determinations in groundwater and rock samples together with the values of f_{a224} and R_{Ra} obtained by applying the NAPL method are reported in Table 3.

5.1 Ra adsorption on MnO_2 in kerosene and recoil constant of 224 Ra recoiled from 228 Th adsorbed on MnO_2

5.1.1 Experiments MDBP-1K and MDBP-1H

In MDBP-1K and MDBP1-H, the ratio of adsorbed ²¹²Pb to desorbed ²¹²Pb in kerosene and n-hexane was equal to 0.33 ± 0.01 and 0.03 ± 0.02 , respectively. K^*_{dreal} and $R_{NAPL real}$ of ²¹²Pb vary between 2.39±0.07 (kerosene) to 0.22 ± 0.15 (n-hexane) and between 3.39 ± 007 (kerosene) to 1.22 ± 0.15 (n-hexane), respectively. The ratio of adsorbed ²²⁸Ra to desorbed ²²⁸Ra in NAPLs was equal to 0.14 ± 0.03 in MDBP-1K and 0.03 ± 0.02 in MDBP-1H. The values of K^*_{dreal} vary between 1.02 ± 0.21 (kerosene) to 0.22 ± 0.15 (n-hexane) whereas $R_{NAPL real}$ varies between 2.02 ± 0.21 (kerosene) to 1.22 ± 0.15 (n-hexane), respectively. All measured values (Table 2) appear much lower than those displayed by Pb and Ra in experiments of adsorption on MnO₂ carried out in water for comparable temporal range and "mass adsorbent /volume solution" ratio [33, 48-50]. The notable difference of adsorption of ²¹²Pb and ²²⁸Ra ions between NAPL (kerosene or hexane) and water suggests that some mechanism hinders the adsorption of cations in apolar liquids allowing to achieve R_{NAPL real} values close to unity(like in n-hexane). Two processes can contribute to this effect. The first is the adsorption of organic polar compounds onto MnO₂. This process could partially work for kerosene but not for pure hexane. The second is the phenomenon of electrostriction in apolar liquids [70]. In these fluids, a spherical positive impurity ion extends notably its electrostrictive influence into the liquid environment. The inclusion of a charge in a non-polar medium with a low permittivity has a range effect on any other present charge much longer than in an aqueous medium [71]. According to the Bjerrum theory, ions might exist separately only if their size exceeds the Bjerrum radius [72]. Bjerrum radius is of about 0.7 nm in water and of 28 nm in non-polar liquids of low permittivity (\approx 1.8) like kerosene or hexane. Electrostriction causes the formation of a cluster of non-polar species around the ion. It may be expected that the ion exists in a polarized atmosphere of neighbour molecules as a relatively large polarized clump that moves through the liquid [73]. The adjacent layers around the ion are layers of high viscosities which approach the bulk viscosity of the liquid only further away from the ionic core [70]. These layers "hinder" the adsorption of simple ions as Ra on adsorption sites. The

Table 3	Radiu	n analy	tical determinat	cions and resu	lts of N∕	VPL method a	application								
SAM- PLE	Bulk density kg·dm ⁻³	Poros- ity of rock	²²⁴ Ra, Bq·kg ⁻¹	²²⁸ Ra, Bq·kg ⁻¹	²²⁶ Ra _r Bq·kg ⁻¹	f_220	f _{a224} (1) NAPL method	f _{r 224} (2)	Meas- ured P. size (D) µm (3)	Calcu- lated P. Size (D) µm (4)	²²⁴ Ra Bq-L ⁻¹	Bq.L ⁻¹ E	²²⁶ Ra Bq·L ⁻¹	R _{Ra-224} R _{Ra-228} (5) (6) NAPL method	R _{1a-228}
CM-1aK CM-1bW	0.50	0.75 0.75	2,900±16 2,900±16				0.0073±0.0009	0.0064	7.0 7.0	6.2±0.7 6.2±0.7	14.1±1.7 6.4±1.1			2.2 ± 0.4	
CM-1 cW	0.50	0.75	2,900± <i>16</i>			0.050±0.005			7.0	6.2±0.7					
CMB-1K			1,980± <i>15</i>	1,980±15								1.40± <i>0.12</i> (10)		4.5 ± 1.2	0
CMB-1W			1,980±15	1,980± <i>15</i>								0.31±0.08			
MDBP-3	a 0.48	06.0	148,750±2,975				0.024±0.002		14.5		952±69				
MDBP-3	o 0.48	06.0	148,750±2,975			0.029±0.003			14.5						
MON-1K	2.51	0.45	331,000± <i>740</i>	331,000± <i>740</i>			$0.000037 \pm *0.00002$ $0.000023 \pm **0.000001$	0.000030	1,00	1,288	43.0±2.4	14.62± <i>0.82</i> (11)			
RP-1	1.35	0.5									0.0089 ± 0.0009	0.0084 ± 0.0005 (0.0600± <i>0.0029</i>		130– 260 (7)
RP-2			299±9	299±9	172±7										
RP-2K RP-2W	1.35	0.5					0.0016±0.0001				1.30± <i>0.11</i> 0.009± <i>0.001</i>			144±20 253±2	9 114 ± 10 (8) 143 ± 13 (9)
RP-2R	1.35	0.5				0.033± 0.003									
(1) calc methou measur errors a	ulated d (Eqs.1 ed afte ire given	from Eq 13–15); r 21 mc 1 at 1-σ	 13 *or Eq. 14 (6) calculated france onths; in the sar level 	**; (2) calcula om: R _{Ra-228} = 1 ne sample, af	ted fron .75.R _{Ra-2} ter 21 n	ר Eq. A2 and 24; (7) calcula 10nths: ²²⁸ Th	measured average ited from Eqs. 28–26 = 0.36 Bq. L ⁻¹ (11) v	particle si 9; (8) calcu value mea:	ze; (3) c lated frc sured af	alculatec om Eq. 3(ter 27 m	l by using Eq.) and $\Phi_c = \Phi =$ onths; in the :	24; (4) calculate 0.5; (9) calculat same sample, af	ed from Eq. 25 ed from Eq. 3(fter 27 month	5; (5) measurec 0 and	l by NAPL (10) value 3q L ⁻¹ . All

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SN Applied Sciences A Springer Nature journat resilient occurrence of Ra adsorption in MDBP-1 K is likely due to the exceptional specific surface of MnO₂ that even in extreme saline solutions, where the retardation factor of Ra is close to 1 (as in Dead Sea water), is still capable to adsorb radium [34] with an efficiency up to 15%.

5.1.2 Experiments MDBP-2

In sample MDBP-2, ratio of adsorbed ²²⁸Ra to desorbed ²²⁸Ra in kerosene is equal to 1.18. K^*_{dreal} and $R_{NAPL real}$ of ²²⁸Ra are 10.5±0.6 and 11.5±0.6. These values (Table 2) show that Ra, when occurs in apolar fluids as a large complexed ion, is moderately adsorbed on MnO₂. Therefore, electrostriction is important as adsorption-hindering mechanism only for simple positive Ra ions being less efficient when Ra occurs as large sized complexed ion.

5.1.3 Experiments MDPB-3

Results of this experiment are given on Table 3. The mean diameter of MDBP particles is ~ 14.5 μ m, when calculated from mode (7 μ m) and median (12 μ m) of the size (diameter) distribution. For asymmetric distribution, the three statistics are connected by the empirical relation [74]:

$$Mean = \frac{\left[(3 \cdot Median) - Mode\right]}{2}.$$
 (24)

In MDBP-3 experiment, initial Ra occurs exclusively as adsorbed ion on the external surface of the sample. In this case, the model of Maraziotis (Eqs.35-37) does not apply, the theoretical fraction from α-recoil being given by Eq. 4. The f_{a224} of MDBP-3a, obtained by the NAPL method (Eq. 14), is equal to $24 \pm 2 \cdot 10^{-3}$. The mean of 12 measurements of ²²⁰Rn emanation coefficient, f_{220} . (MDBP-3b) is similar: $29 \pm 3 \cdot 10^{-3}$. In case of R_{*Ra*224} = 1, in this particular experiment, f_{a224} and f_{220} are expected to be quite similar (apart a difference of 5% due to their Q_{av} the α -decay energy of their parent atoms). Their difference is of 20%, corresponding to a R_{Ra224} of 1.21 ± 0.16 which is very close, taking in account the associated error, to the assumed value of $R_{Ra224} = 1$ in NAPL. The low absolute values of f_{a224} and f_{220} (compared with the theoretical value of 0.50 in absence of embedding effects) are likely due to embedding phenomena [26, 27]. The consistence of the embedding model can be tested by applying Eq. 4 and by comparing the obtained f_{a224} with f_{220} . By considering the recoil energy of ²²⁴Ra and its R_F in kerosene (similar to polybutadiene [75]), it obtains a pore size (w) of 5.2 nm. A comparable pore size (5.3 nm) is calculated by considering the recoil energy of 220 Rn and its R_E in water. This pore size matches that observed in synthetic and commercial MnO₂ [76, 77]. This experiment shows that only if Ra is exclusively

located in grains as adsorbed ion, it enters the liquid phase by the same mechanism of radon (direct recoil) and the two methods, NAPL and radon comparison, coincide.

5.2 Ra adsorption on zeolite 4A

In Z4A-1, the ratio of adsorbed ²²⁴Ra to desorbed ²²⁴Ra, K^*_{dreal} and $R_{NAPL real}$ in kerosene are equal to 0.02 ± 0.03 , 0.04 ± 0.07 and 1.04 ± 0.07 , respectively. These results (Table 2) confirm the assumption that in NAPL the retardation factor of Ra simple ions is close to unity. Differently in Z4A-2 (Ra complex ion in kerosene) and especially in Z4A-3 (Ra simple ion in water), the real Ra retardation factors are higher $(2.19 \pm 0.07 \text{ and } 16.0 \pm 2.3, \text{ respectively})$. The formation of clusters of molecules of apolar fluid around simple Ra ions, induced by electrostriction, explains this behaviour. They prevent Ra adsorption on elite in kerosene especially when the electrostrictive effect of the Ra charge is not lowered by formation of Ra-complexes with large organic molecules. As large complexed ion (Z4A-2), instead, Ra is adsorbed significantly on zeolite (like on MDBP) in apolar liquids. TTA-TOPO-Ra complexes, together to the crown ether-Ra complexes, are the only ones known to be soluble in apolar liquids but both are not present in common NAPLs. Thus, the suggestion of using common apolar liquids like kerosene to determine R_{Ra} is strengthen.

5.3 Ra adsorption on clay

As expected, in S8N1-3 W the real Ra retardation factor in water is high (23.2 ± 2.2) (Table 3). Differently, in S8N1-1 K and S8N1-3 K the R_{NAPL real} are close to unity $(1.27 \pm 0.18, 1.27 \pm 0.27, respectively)$. Clay minerals have the highest capacity of adsorbing NAPLs between common minerals [78]. By comparing Ra in NAPL (S8N1-1 K, S8N-3 K) and in water (S8N1-3 W), it can be concluded that NAPL molecules, wetting clay surfaces, hinder the development of double layer reducing the adsorption of inorganic ions as yet suggested [79]. Clay minerals are ubiquitous in soil and Ra isotopes seem to be unaffected by adsorption when clay is immersed in NAPL; therefore, the determination of the contact time between NAPLs and particle grains of sediments or soils is possible, as the following two sections will show.

5.4 Monazite in contact with NAPL

The surface of monazite grains is charged negatively and adsorbs positive ions as Ce, La and Ca [80]. In particular, adsorption of divalent elements like Ca, Ba and Ra, is high when they occur as Ca(Ra,Ba)OH⁺, via hydrogen bonding with phosphate-oxygen active sites negatively charged on the monazite surface [81, 82]. The f_{a224} of MON-1 K (NAPL



Fig. 6 Evolution of ²²⁸Th/²²⁸Ra activity ratio in kerosene according to Eq. 21 for different couples of Th and Ra retardation factors. Full black circle: ²²⁸Th/²²⁸Ra activity ratio in kerosene after 27 months of contact with monazite

method, Table 3) calculated by Eq. 13 ($3.7 \cdot 10^{-5}$) or Eq. 14 ($2.3 \cdot 10^{-5}$), is comparable to the f_r ($3.0 \cdot 10^{-5}$) resulting from Eq. 35, therefore α -recoiled Ra is not adsorbed onto monazite in apolar fluid.

The ²²⁸Th/²²⁸Ra activity ratio in kerosene, measured after 27 months of contact between kerosene and monazite, is equal to 0.34 ± 0.4 corresponding to a time of contact of 28 ± 3.5 months.

Figure 6 shows the good matching between ²²⁸Th/²²⁸Ra activity ratio measured in kerosene and its theoretical value after 27 months of contact between kerosene and monazite. This is a clear evidence, as mentioned before, that retardation factors of Ra and Th are both close to 1 in kerosene. The temporal evolution of the ²²⁸Th/²²⁸Ra activity ratio is sensible to the values of the retardation factors. Thus, this correspondence is a strong argument in favour of inhibition of the adsorption of inorganic positive ions on charged sites of mineral surfaces in contact with apolar fluids. The results suggest that this radioactive pair can be useful for dating NAPL pollution in sediments or soil which contain always some ²³²Th that generates ²²⁸Ra recoiled atoms into NAPL in contact with them [83].

5.5 Application of the NAPL method for determining retardation factor and recoil constant of Ra for a Th-enriched weathered volcanic rock

The ²²⁸Th/²²⁸Ra activity ratio in CMB-1 K, measured after 21 months of contact between kerosene and monazite, is equal to 0.26 ± 0.4 corresponding to a time of contact of 21 ± 4 months (Table 3). As before, this results confirms that retardation factors of Ra and Th are both close to 1 in kerosene and the possibility to use this radioactive pair to estimate the time of contact between NAPL and soil.

The mean particle diameter of CM-1 was calculated (Eq. 29) from mode (40 µm) and median (18 µm) obtaining a value of 7.0 µm. The f_{a224} of the CM-1aK is equal to $7.3 \pm 0.9 \cdot 10^{-3}$ (Table 3). Remembering that the intraparticle pore recoil component of Ra can be neglected then $f_a = f_r$ and combining Eq. 13 and Eq. 35 it obtains:

$$f_{r\,224} = f_{\alpha 224} = \frac{^{224} \text{Ra}_{\text{NAPL}}}{^{224} \text{Ra}_r} \cdot \frac{\Phi}{\rho_r} = \frac{3\text{R}_c}{4\alpha} - \frac{1}{16} \cdot \left(\frac{\text{R}_c}{\alpha}\right)^3 \quad (25)$$

From Eq. 25, knowing f_{a224} , measured by NAPL method, and R_c in the particle, the average particle diameter (2*a*) has been estimated (6.2 µm) resulting comparable with the value found by laser granulometry. By using the mean diameter of 7.0 μ m, f_{r224} of CM-1 is equal to 6.4·10⁻³, a value fitting the measured one (f_{a224} of CM-1aK=7.3·10⁻³). The intraparticle pore recoil component ($f_{p224} = f_{a224}$ ---- $f_{r224} = 0.9 \cdot 10^{-3}$) is clearly trascurable. This is not the case for 220 Rn. The average of 12 measurements of f_{220} (CM-1cW) is in fact equal to $50 \pm 5 \cdot 10^{-3}$. Therefore by subtracting the sum of f_{r220} and f_{d220} to $f_{220,}$ a relevant ²²⁰Rn intraparticle pore recoil component, $f_{p220} = 43 \cdot 10^{-3}$, is obtained. Thus the "Rn method" should give a R_{Ra-224} seven times higher than that obtained by "NAPL method". This difference depends on f_n : it is negligible for ²²⁴Ra but not for ²²⁰Rn. Nanobubbles of air occur commonly in solid-state nanopores immersed in liquids [84]. Liquids in nanopores of silicates, moreover, do not form a compact structure of connected molecules but rather a collection of small clusters forming plugs separated by air bubbles [85]. Thus, any nanometric air bubble along the path produces a fast electrostatic attachment of Ra ions on the walls of the pore, reducing notably f_{n224} . Differently, Rn atoms can continue to diffuse in the pore. The R_{Ra-224}, obtained by dividing the 224 Ra content of CM-1aK to that of CM-1bW, and the R_{Ra-228}, obtained by dividing the ²²⁸Ra content of CMB-1 K to that of CMB-1 W, are equal to 2.2 ± 0.4 and 4.5 ± 1.2 , respectively. The ratio $R_{Ra-228/} R_{Ra-224} = 2.04 \pm 0.64$ is therefore comparable with the mean value of 1.75 found from literature data. Finally, the relatively low R_{Ra}s are likely due to

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Fig. 7 Variation of $^{226}\rm{Ra}$ retardation factor due to leucite versus S. The $^{226}\rm{Ra}$ retardation factor is calculated from experimental values of K_A and S

the complete obliteration of leucite (a Ra-adsorber phase causing high retardation factors) by weathering processes.

5.6 Ra adsorption onto leucite

Figure 7 reports ²²⁶Ra adsorption data on leucite. The TSA (total surface area)/V_S (volume solution) ratio is plotted vs. the R_{Ra-226} calculated adopting the Sugita-Gillham formalism [86]. According to it, R_{Ra-226} depends on the product of the distribution coefficient K_A defined on an area basis (μm^{-1}) and the TSA/V_S ratio:

$$R_{Ra-226} = 1 + \left(\frac{K_A \cdot TSA}{V_S}\right) = 1 + \left(K_A \cdot S\right)$$
(26)

The TSA/V_s ratio equals S, the extent of the rock surface in contact with unit volume of groundwater (μm^{-1}). The distribution coefficient between adsorbed phase and solution, K_A, defined on an area basis, is equal to:

$$K_{A} = \frac{^{226}\text{Ra} \text{ activity per unit surface area}}{^{226}\text{Ra activity per unit volume of solution}}.$$
 (27)

The R_{Ra-226} due to adsorption on leucite, obtained by linear best-fitting through the plotted points, is:

$$R_{Ra-226} = 1 + \left(\frac{2.88 \cdot 10^3 \cdot TSA}{V_S}\right)$$
(28)



Fig. 8 Initial dissolution of Red Pozzolane in distilled water at pH 6.0–6.5: Recorded TDS values over time. Dotted line signs the limit value of 475 mg L^{-1}

In order to find the ratio TSA/V_s, which equals S, the Carman–Kozeny equation for porous strata and spherical particle grains was used [87]:

$$S(\mu m^{-1}) = \left\{ \frac{\left[\frac{\Phi^{3}}{(1-\Phi)^{2}}\right]}{k \cdot 5} \right\}^{0.5}$$
(29)

where k is the mean permeability coefficient.

By knowing that $\Phi = 0.50$ and $k = 5 \cdot 10^{-1} \mu m^2$ in RP [88, 89], a value for S of $0.45 \cdot m^{-1}$ is obtained. This value, multiplied for the modal fraction of leucite in RP (0.1–0.2), gives a R_{Ra-226} due to leucite in the range of 130–260. The implications of this result are discussed in the following section.

5.7 Study site: Ra isotopes in groundwater from the Red Pozzolane aquifer (Alban Hills)

Results of analytical determination of radium isotopes in RP groundwater are given in Table 3. The Ra isotopic composition of RP-1 is equal to $8.9 \pm 0.9 \cdot 10^{-3}$ Bq ·L⁻¹ for ²²⁴Ra, $8.4 \pm 0.5 \cdot 10^{-3}$ Bq ·L⁻¹ for ²²⁸Ra and $60.0 \pm 2.9 \cdot 10^{-3}$ Bq ·L⁻¹ for ²²⁶Ra. The f_{a224} of RP-2 K, obtained by the NAPL method, results equal to $1.6 \pm 0.1 \cdot 10^{-3}$. The R_{Ra-224}, calculated from RP-2 K and RP-2 W, by using the NAPL method (Table 3), is close to 144 ± 20 . The average of 12 measurements of the ²²⁰Rn emanation coefficient, f_{220} , carried out on RP-2R, results equal to $33 \pm 3 \cdot 10^{-3}$. Figure 8, where TDS is plotted

vs. T^{0.5}, shows the variation of initial dissolution of RP in groundwater. At the end of the initial dissolution (45 days), TDS trends to a stationary value of 487 mg·l⁻¹, close (about 85%) to that measured in the field (572 mg·l⁻¹). The increase of pH, at the end of experiment, is < 0.5 pH units (pH _{initial} = 6).²²⁶Ra_{w(0)} and ²²⁸Ra _{w(0)}, then, are calculated according to:

 226 Ra_{w (0)} (Bq ·L⁻¹) = TDS stationary (kg· L⁻¹)·²²⁶Ra_r (Bq·kg⁻¹) = 0.487 ·10⁻³ ·172 = 0.084 Bq ·L⁻¹,

 228 Ra_{w(0)} (Bq ·L⁻¹) = TDS stationary (kg· L⁻¹)·²²⁸Ra_r (Bq·k g⁻¹) = 0.487 · 10⁻³·299 = 0.146 Bq ·L⁻¹.

A first estimation of R_{Ra-228} and R_{Ra-226} was done by applying Eq. 45 to the value of R_{Ra-224} , found by NAPL method (see Appendix 5). It gives a value of 253 ± 29 , in the range of the values estimated for the only leucite which plays, in this aquifer, an important role in Ra adsorption. In RP, the contribute of the initial dissolution of aquifer rock to Ra activity cannot be neglected, as yet others evidenced [3]. Thus, Eqs. 50–52 (see Appendix 5) were applied. Since the first estimate of R_{Ra-228} is very high, exponential terms can be considered negligible and Eq. 50 becomes:

$$^{228}\mathsf{Ra}_{w} = \left[\frac{f_{\alpha 224} \cdot 0.74}{\mathsf{R}_{Ra-228}}\right] \cdot \left(^{228}\mathsf{Ra}_{r} \cdot \frac{\rho_{r}}{\Phi}\right)$$
(30)

The last equation, known f_{a224} (from NAPL method) and the other variables (Φ , ρ_r , ²²⁸Ra_r) of the aquifer, gives a second estimate of $R_{Ra-228} = R_{Ra-226} = 114 \pm 10$. This second estimate, lower than the value of R_{Ba-224} (144 ± 20) found with the NAPL method, seems to contradict the previous statement that the R_{Ra-228}/R_{Ra-224} ratio is ≥ 1 . Indeed up to now it has been implicitly assumed that porosity in porous aguifers corresponds to the pore connectivity. However this is not true especially for vesiculated volcanic aquifer rocks. In this case Φ in Eq. 36 should be substituted by the pore connectivity Φ_c which is the product of Φ and connectivity C [90]. The variable vesicularity of Red Pozzolane [42], from poorly to moderate (0-20%), suggests that pore connectivity could reach the value of 0.40. In this case, the contradiction should be solved because Eq. 36 yields $R_{Ra-228} = R_{Ra-226} = 143 \pm 13$. Anyway, after averaging the two extreme values (114 and 253), we assume 184 ± 15 as the more reliable value for $R_{Ra-228} = R_{Ra-226}$. According to Eq. 52, the age of groundwater sampled at the Arco di Travertino site results equal to 5.5 ± 0.4 years. This is the age of the groundwater at the sampled depth which, however, differs from the mean residence time, MRT. A critical problem in MRT determination, in fact, is the dispersion of groundwater age, i.e. the relationship between MRT and the depth of aquifer below the water table. Age increases with the depth of sampling respect to the level of water table, and the values can be spread by a factor 50 [91]. It is



Fig. 9 Mean residence time of Alban Hills groundwater at the Arco di Travertino site and at Egeria mineral water site, calculated from Eqs.32–33

however possible to use an empirical equation to correlate the recharge rate *R* to the groundwater age [92]:

$$R = \left\{ \ln \left[\frac{H}{(H-z)} \right] \right\} \cdot \left(\frac{H \cdot \Phi}{T} \right)$$
(31)

where *R* is the recharge rate (m y^{-1}), *T* the groundwater age, *H* the height of the water column and *z* is the distance from the sampling point and the upper level of the water table. Equation 31 can be expressed as well as:

$$\frac{z}{H} = 1 - e^{-\frac{RT}{H\Phi}}$$
(32)

When this last equation is rendered graphically, the MRT is given by the intersection point of z/H function with the line corresponding to $z/H = 1 - e^{-1} = 0.632$ (Fig. 9). In the case of the piezometric well at Arco di Travertino (Fig. 5) H = 23, z = 1.5, $\Phi = 0.5$ and $T = 5.5 \pm 0.4$ y, therefore $R = 0.141 \pm 0.011$ m y⁻¹. This value matches that estimated by others (0.152 m·y⁻¹) for the same hydrogeological basin [89].

In this exponential age distribution, the mean residence time (MTR) is given by [92]:

$$MTR = H \cdot \Phi/R. \tag{33}$$

Applying this equation to the site of Arco di Travertino, a MTR of 82 ± 6 years is obtained. Of course, establishing the exact average residence time of groundwater of the entire Alban Hills aquifer is beyond the scope of this work.

This example was made to emphasize that the estimated local MRT, even for a single well, is realistic when using the NAPL method. Differently, by using the ²²⁰Rn/²²⁴Ra ratio for estimating the retardation factor of ²²⁴Ra, a value of 3000 is obtained, ~ 21 times $(f_{220}/f_{a224} ~ 21)$ higher than that measured by NAPL method and, consequently, an unrealistic low MRT (3.9 years). The goodness of the model (Eq. 52 and Eq. 32) can be verified comparing these results to the age and MRT of the Egeria mineral water, springing at about 1 km at SE of the Arco di Travertino site, inside the same RP aguifer. At Egeria site, the presence of a CO_2 bubble plume has likely destructed the age stratification of groundwater [93] (Fig. 5). Therefore, in this site, the groundwater age T coincides with MRT. From available hydrogeological maps and geological profiles [42, 94] the thickness of the aquifer is, at Egeria site, of 27 m. By assuming the calculated R of 0.141 \pm 0.011 m y⁻¹, a MRT of 96 ± 8 years is obtained from Eq. 33 (Fig. 9). It is interesting to compare this predicted value with the one resulting by applying Eq. 50-52 (see Appendix 5) to Ra in groundwater for this site. Available data on ²²⁶Ra and ²²⁸Ra for the Egeria mineral water are equal to $3.56 \cdot 10^{-3}$ and $10.2 \cdot 10^{-3}$ Bg $\cdot L^{-1}$, respectively [95]. Assuming the same $R_{\text{Ra-}226}$ (184 \pm 15) and 226 Ra_{w(0)} (0.084 Bq ·L⁻¹) found for the Arco di Travertino site, it obtains a value of 84±7 years, quite comparable with the predicted one.

6 Conclusions

Adsorption experiments on Ra-adsorbing materials (manganese dioxide, zeolite, clay,) immersed in NAPL and measurements of the time of contact NAPL /monazite and NAPL/weathered volcanic rock by ²²⁸Th/²²⁸Ra activity ratio have been carried out. The experiments suggest that when these materials are immersed in non-aqueous phase liquids (NAPL) adsorption on particle grains is strongly reduced (as in the case of manganese dioxide) or absent (as in the case of the other tested materials). This is evident especially when distribution coefficient is corrected for the particle concentration effect. A method based on this observation, the NAPL method, has been described and applied to estimate recoil rate constant and retardation factor of Ra isotopes. When retardation factor and recoil rate constant of Ra isotopes are estimated by the NAPL method, the obtained values result always lower than those found using radon as comparing term. Indeed radon is introduced in solution by a more complex way than radium. Radium is introduced essentially by direct α-recoil, whereas radon by further mechanisms as back diffusion. The NAPL method is based on the theory of electrostriction which predicts that, in nonpolar solutions, simple ions are surrounded by non-polar molecules. These molecules are feebly polarized by the electrical field of ions occurring in non-polar fluids. The polarization power of these ions, in non-polar fluids, is about 30 times stronger than in polar fluids like water. At the interface between charged surface of adsorbent and NAPL, the same electrostrictive phenomena allow also a full saturation of adsorption sites by polar and non-polar molecules of NAPL. The result of the two effects is a negligible amount of adsorbed a-recoiled nuclides onto the surface of particle grains and an increase in the conservative behaviour of these radionuclides in NAPL in comparison with water. By knowing retardation factor and recoil constant of radium through NAPL method, the equations describing the evolution of radium isotopes in aguifers can be solved in many cases. An example is the determination of residence time in upper shallow aquifer of Alban Hills mainly hosted in Red Pozzolane units. An interesting implication of the results of this paper is also the theoretical possibility to date the NAPL contamination time of soils, from medium (> 10 ppm) to high content of thorium, by measuring the evolution of the ²²⁸Th/²²⁸Ra ratio in NAPLs [83].

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Availability of data and material The main data set of experiments is given in the paper in Table 1. Gamma spectra of analysed sample are stored and saved at IGAG -CNR Environmental Geochemistry Laboratory.

Compliance with ethical standards

Conflict of interest No competing interests among authors or between CNR and Università Roma TRE are involved in the preparation of this research paper.

Appendix 1

Derivation of the master equation

According Kiro et al. [6, 7], by assuming one-dimensional steady-state flow coupled with negligible hydrodinamic dispersion and negligible co-precipitation, the solution of the differential equation describing the temporal variation of Ra activity in groundwater along a single flow path x with constant velocity v_x and constant K_d^* is

$${}^{i}\text{Ra}_{w} = \left[\frac{P}{\lambda_{i}\cdot\left(K_{d}^{*}+1\right)}\right]\cdot\left(1-e^{-\lambda_{i}\cdot\left(K_{d}^{*}+1\right)\cdot\left(\frac{x}{v_{x}}\right)}\right) + {}^{i}\text{Ra}_{w(0)}\cdot e^{-\lambda_{i}\cdot\left(K_{d}^{*}+1\right)\cdot\left(\frac{x}{v_{x}}\right)}$$

According Kiro et al. [7], P is the production rate of ${}^{i}Ra_{w}$ activity (Bq·t⁻¹) from the aquifer solids by recoil and dissolution and it is given by:

 $\mathsf{P} = \left(\Delta_i + \Gamma_i\right) \cdot {}^i \mathsf{Th}_r \cdot \frac{\rho_r}{\Phi}.$

Firstly, in case of radioactive equilibrium in the aquifer rock, ${}^{i}Th_{r} = {}^{i}Ra_{r}$; then *ii*) $\left(\frac{x}{v_{x}}\right)$, the ratio between the length of the linear radionuclide path and its velocity, is equal to T; finally, *iii*) $\left(K_{d}^{*}+1\right) = R_{i}$.

Therefore, Eq. 1 is easily obtained:

$${}^{i}\mathsf{Ra}_{w} = \left[\frac{\left(\Delta_{i} + \Gamma_{i}\right)}{\left(\lambda_{i} \cdot \mathsf{R}_{i}\right)}\right] \cdot \left(1 - e^{-\lambda_{i} \cdot \mathsf{R}_{i} \cdot \mathsf{T}}\right) \cdot \left({}^{i}\mathsf{Ra}_{r} \cdot \frac{\rho_{r}}{\Phi}\right) + {}^{i}\mathsf{Ra}_{w(0)} \cdot \left(e^{-\lambda_{i} \cdot \mathsf{R}_{i} \cdot \mathsf{T}}\right)$$

Appendix 2

Components of Rn and Ra emanation

They are defined by [13]:

$$f_r = \frac{3R_c}{4\alpha} - \frac{1}{16} \cdot \left(\frac{R_c}{\alpha}\right)^3 \tag{35}$$

$$f_p = \frac{R_c \cdot k(1+k)}{2r_p} - \frac{1}{16} \cdot \left(\frac{R_c}{\alpha}\right)^3$$
(36)

$$f_{\rm d} = \frac{3}{\alpha} \cdot \left(1 - f_{\rm r} - f_{\rm p}\right) \cdot \sqrt{\frac{\mathsf{D}}{\lambda}} \tag{37}$$

where R_c is the recoil range in the particle (nm), α the particle radius (nm) and k its porosity, r_p the radius of the nanopores in the particle (nm), D the intraparticle diffusion coefficient (nm² s⁻¹) and λ the Rn decay constant (s⁻¹).

Appendix 3

Diffusion coefficient of Rn and Ra in minerals

No data are available in the literature for diffusion coefficient of radium in minerals at ambient temperature (298°K). By using data on barium [96], which can be considered the stable analogue of radium for the same ionic charge and similar ionic radius, the Ra diffusion coefficient in feldspar (a common rock-forming mineral), at ambient temperature, should be close to 10^{-80} m^{2 s-1}. Diffusion

SN Applied Sciences A Springer Nature journal coefficient of Rn in feldspar at the same temperature is tenth orders higher, about 10^{-31} m^{2 s-1} [12].

(34)

Appendix 4

Limits of variation of the R_{Ra-228}/R_{Ra-224} ratio

A recent paper [35] assumes $R_{228}/R_{224} = 1$ and desorption rate constant $k_d > \lambda_{224}$ (0.19 day⁻¹). Other authors, more

realistically, found that k_d values are not less than 0.5- λ_{224} [4, 97, 98]. Thus, by considering that $\lambda_{228} < < k_d$, from Eq. 11:

$$\left(\frac{R_{Ra-228}-1}{R_{Ra-224}-1}\right) = \frac{\left(k_d + \lambda_{224}\right)}{k_d} = 1 + \frac{\lambda_{224}}{k_d}$$
(38)

By constraining the value of k_d to $\ge 0.5 \lambda_{224}$, then:

$$3 \ge \left(\frac{R_{Ra-228} - 1}{R_{Ra-224} - 1}\right) \ge 1$$
(39)

or:

$$3 - \frac{2}{R_{Ra-224}} \ge \frac{R_{Ra-228}}{R_{Ra-224}} \ge 1$$
(40)

 R_{Ra-224} is by definition ≥ 1 , then the R_{Ra-228}/R_{Ra-224} ratio must vary in a narrow range, between 1 and 3.

Appendix 5

Assuming dissolution rate constants of ²²⁸Ra and ²²⁶Ra be equal (congruent dissolution), three typical different cases for Eqs.17–19 are possible:

(1) initial activity of ²²⁸Ra and ²²⁶Ra can be neglected / α -recoil prevails over dissolution, (2) initial activity of ²²⁸Ra and ²²⁶Ra can be neglected/ dissolution prevails over α -recoil 3) initial activity of ²²⁸Ra and ²²⁶Ra cannot be neglected/ α -recoil prevails over dissolution.

Initial activity of ²²⁸Ra and ²²⁶Ra can be neglected / α -recoil prevails over dissolution ($\Gamma >> \Delta$)

 $^{228}\text{Ra}_{w(0)}$ and $^{226}\text{Ra}_{w(0)}\!=\!0;\;\Delta_{2208}\;\Gamma_{228}\!=\!\Gamma_{228}$ and $\Delta_{228}\;\Gamma_{226}\!=\!\Gamma_{226}.$

Remembering that $\Gamma_{228} = f_{a228} \cdot \lambda_{228}$ and $\Gamma_{226} = f_{a226} \cdot \lambda_{226}$ Eq. 19 becomes:

$$\frac{\left(\frac{22^{8}Ra_{w}}{22^{6}Ra_{w}}\right)}{\left(\frac{22^{8}Ra_{r}}{2^{26}Ra_{r}}\right)} = \frac{f_{\alpha 228}}{f_{\alpha 226}} \cdot \frac{\left(1 - e^{-\lambda_{228} \cdot R_{Ra-228} \cdot T}\right)}{\left(1 - e^{-\lambda_{226} \cdot R_{Ra-228} \cdot T}\right)}$$
(41)

The f_a of a Ra isotope *i* depends on its recoil range in the particle, ${}^{i}R_{c}$ (see Eq. 35–36) which in turn can be calculated by the following equation [99]:

$${}^{i}R_{c} = (M_{1} + M_{2}) \cdot (M_{2} \cdot E_{M1} \cdot K) \cdot \frac{\left(Z_{1}^{2/3} + Z_{2}^{2/3}\right)^{1/2}}{(M_{1} \cdot Z_{1} \cdot Z_{2})} \quad (42)$$

where E_{M1} is the initial recoil energy of a nucleus of mass M_1 and atomic number Z_1 , moving in an absorber of average mass M_2 and average atomic number Z_2 , K is a constant value (6.02), ${}^{i}R_{c}$ is the recoil range in the absorber given in nm and E_{M1} is given in keV.

In general $f_{\alpha i} \setminus f_{\alpha i'} =^{i} R_{c} \setminus {}^{i'}R_{c} \approx E_{Mi}/E_{Mi'} \approx Q_{\alpha i} \setminus Q_{\alpha i'}$ (43)

where Q_a is the α -decay energy of thorium parent.

Since $Q_{aTh-228}/Q_{a-Th-230} = 1.15$ Eq. 41 can be written as:

$$\frac{\left(\frac{228}{26}Ra_{w}}{\left(\frac{228}{26}Ra_{r}}\right)}{\left(\frac{228}{226}Ra_{r}}\right)} = 1.15 \cdot \frac{\left(1 - e^{-\lambda_{228} \cdot R_{Ra-228} \cdot T}\right)}{\left(1 - e^{-\lambda_{226} \cdot R_{Ra-228} \cdot T}\right)}$$
(44)

$$R_{Ra-228} = R_{Ra-226} = 1.75 \cdot R_{Ra-224}$$
(45)

Initial activity of 228 Ra and 226 Ra can be neglected/ dissolution prevails over α -recoil ($\Delta >> \Gamma$)

 228 Ra_{*w*(0)} and 226 Ra_{*w*(0)} = 0; $\Delta_{228} = \Gamma_{228} = \Delta_{228}$ and Δ_{226} $\Gamma_{226} = \Delta_{228}$.

Equation (17) becomes:

. . . .

. . . .

$$\frac{\left(\frac{22^{8}\mathsf{Ra}_{w}}{22^{6}\mathsf{Ra}_{w}}\right)}{\left(\frac{22^{8}\mathsf{Ra}_{r}}{22^{6}\mathsf{Ra}_{r}}\right)} = \frac{\lambda_{226} \cdot \Delta_{228}}{\lambda_{228} \cdot \Delta_{226}} \cdot \frac{\left(1 - e^{-\lambda_{228} \cdot \mathsf{R}_{Ra-228} \cdot \mathsf{T}}\right)}{\left(1 - e^{-\lambda_{226} \cdot \mathsf{R}_{Ra-228} \cdot \mathsf{T}}\right)}$$
(46)

We can assume to fall back into this case only if the left member of Eq. 46 is $\geq (\lambda_{226}/\lambda_{228})$.

In case of congruent dissolution, $\Delta_{228} = \Delta_{226}$, therefore:

$$\frac{\left(\frac{228}{26}R_{a_{w}}}{\left(\frac{228}{226}R_{a_{r}}}\right)} = \frac{\lambda_{226}}{\lambda_{228}} \cdot \frac{\left(1 - e^{-\lambda_{228} \cdot R_{Ra-228} \cdot T}\right)}{\left(1 - e^{-\lambda_{226} \cdot R_{Ra-228} \cdot T}\right)}$$
(47)

If $\Delta_{228} \neq \Delta_{226}$ dissolution rates can be however estimated by literature values [2]. Also in this case, the equation can be solved for T, measuring R by NAPL method and applying Eq. 46

Initial activity of 228 Ra and 226 Ra cannot be neglected – α -recoil prevails over dissolution (Γ >> Δ)

If initial activity of ²²⁶Ra and ²²⁸Ra cannot be neglected and $\Delta_{228}+\Gamma_{228}=\Gamma_{228}$; $\Delta_{226}+\Gamma_{226}=\Gamma_{226}$, then (Eq. 1):

$${}^{228}\text{Ra}_{w} = \left\lfloor \frac{(\Gamma_{228})}{(\lambda_{228} \cdot R_{Ra-228})} \right\rfloor \cdot \left(1 - e^{-\lambda_{228} \cdot R_{Ra-228} \cdot T}\right) \cdot \left({}^{228}\text{Ra}_{r} \cdot \frac{\rho_{r}}{\Phi}\right) + {}^{228}\text{Ra}_{w(0)} \cdot \left(e^{-\lambda_{228} \cdot R_{Ra-228} \cdot T}\right)$$
(48)

$${}^{226}\text{Ra}_{w} = \left[\frac{\left(\Gamma_{226}\right)}{\left(\lambda_{226} \cdot \text{R}_{Ra-226}\right)}\right] \cdot \left(1 - e^{-\lambda_{226} \cdot \text{R}_{Ra-226} \cdot \text{T}}\right) \cdot \left({}^{226}\text{Ra}_{r} \cdot \frac{\rho_{r}}{\Phi}\right) + {}^{226}\text{Ra}_{w(0)} \cdot \left(e^{-\lambda_{226} \cdot \text{R}_{Ra-226} \cdot \text{T}}\right)$$
(49)

We can assume to fall within this case only if the left member of Eq. 44 is \geq 1. The equation can be solved for T, measuring R_{Ra-224} by NAPL method and remembering that:

Since $\Gamma_{228}/\lambda_{228} = f_{a228}$; $\Gamma_{226}/\lambda_{226} = f_{a226}$ and $f_{a228} = f_{a224} \cdot \frac{228}{R_c}/2^{224}R_c = f_{a224} \cdot 0.74$. $f_{a226} = f_{a224} \cdot \frac{226}{R_c}/2^{224}R_c = f_{a224} \cdot 0.86$. where R_c is the recoil range in the particle, then:

$${}^{228}\text{Ra}_{w} = \left[\frac{f_{\alpha 224} \cdot 0.74}{R_{Ra-228}}\right] \cdot \left(1 - e^{-\lambda_{228} \cdot R_{Ra-228} \cdot T}\right) \cdot \left({}^{228}\text{Ra}_{r} \cdot \frac{\rho_{r}}{\Phi}\right) + {}^{228}\text{Ra}_{w(0)} \cdot \left(e^{-\lambda_{228} \cdot R_{Ra-228} \cdot T}\right)$$
(50)

$${}^{226}\text{Ra}_{w} = \left[\frac{f_{\alpha 224} \cdot 0.86}{R_{Ra-226}}\right] \cdot \left(1 - e^{-\lambda_{226} \cdot R_{Ra-226} \cdot T}\right) \cdot \left({}^{226}\text{Ra}_{r} \cdot \frac{\rho_{r}}{\Phi}\right) + {}^{226}\text{Ra}_{w(0)} \cdot \left(e^{-\lambda_{226} \cdot R_{Ra-226} \cdot T}\right)$$
(51)

By manipulating Eqs. 50–51:

$$\left[\frac{2^{28}\text{Ra}_{w} - \left(2^{28}\text{Ra}_{w(0)} \cdot e^{-\lambda_{228} \cdot \text{R}_{Ra-228} \cdot \text{T}}\right)}{2^{26}\text{Ra}_{w} - \left(2^{26}\text{Ra}_{w(0)} \cdot e^{-\lambda_{226} \cdot \text{R}_{Ra-228} \cdot \text{T}}\right)}\right] = \frac{0.74}{0.86} \cdot \frac{2^{28}\text{Ra}_{r}}{2^{26}\text{Ra}_{r}} \cdot \frac{\left(1 - e^{-\lambda_{228} \cdot \text{R}_{Ra-228} \cdot \text{T}}\right)}{\left(1 - e^{-\lambda_{226} \cdot \text{R}_{Ra-228} \cdot \text{T}}\right)}$$
(52)

By knowing $^{228}Ra_{w(0)}$, $^{226}Ra_{w(0)}$ and measuring R_{Ra} by NAPL method, Eq. 52 can be solved for T.

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