LONG TIME BEHAVIOR OF AN AGE- AND LEAKY MEMORY-STRUCTURED NEURONAL POPULATION EQUATION*

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Abstract. We study the asymptotic stability of a two-dimensional mean-field equation, which takes the form of a nonlocal transport equation and generalizes the time-elapsed neuron network model by the inclusion of a leaky memory variable. This additional variable can represent a slow fatigue mechanism, such as spike-frequency adaptation or short-term synaptic depression. Even though two-dimensional models are known to have emergent behaviors, such as population bursts, which are not observed in standard one-dimensional models, we show that in the weak connectivity regime, two-dimensional models behave like one-dimensional models, i.e., they relax to a unique stationary state. The proof is based on an application of Harris's ergodic theorem and a perturbation argument, both adapted to the case of a multidimensional equation with delays.

Key words. long time behavior, nonlocal transport equation, mean-field equation, Doeblin's and Harris's theorems, piecewise deterministic Markov process, spiking neuron, spike-frequency adaptation, short-term synaptic plasticity

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1. Introduction. Multidimensional mean-field models in theoretical neuroscience are challenging to analyze [41, 48, 2, 32], but their study is a necessary step towards understanding how multiple timescales present at the single-neuron level [40, 45] affect the dynamics of large networks of neurons.

One-dimensional mean-field equations for populations of spiking neurons with deterministic drift and stochastic jumps have been a subject of mathematical studies since the works of Pakdaman, Perthame, and Salort [35, 36, 37], providing rigorous foundations for earlier works in theoretical neuroscience [50, 21, 18, 19]. These population equations correspond to the mean-field limit of large networks of interacting neurons [10, 15, 6]. However, they are derived from spiking neuron models that are of the "renewal" type (with the exception of [37]), which means that, while they capture the effect of neuronal refractoriness, they neglect slower neuronal timescales, such as those of spike-frequency adaptation and short-term synaptic plasticity.

To take into account slow neuronal timescales, state-of-the-art phenomenological spiking neuron models must be multidimensional [28, 45] or kernel-based [46, 38, 39] (see also [20, Chap. 6.4]). In the following, we consider a class of neuron models that characterize neuronal refractoriness by an "age" variable (the time elapsed since the last spike) and effects of spike-frequency adaptation or short-term synaptic plasticity by a "leaky memory" variable. For this class of neuron models, the mean-field limit is characterized by a multidimensional transport equation with a nonlocal boundary

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condition [43]. In this work, we study in the two-dimensional case the long time behavior of the solutions to the equation proposed in [43].

1.1. The age- and leaky memory-structured model. The population model we consider describes the evolution of a density ρ_t over the state-space $(a, m) \in \mathbb{R}_+ \times \mathbb{R}^*_+$, where a and m are the "age" and "leaky memory" variables of the neuron, and $\rho_t(a, m)$ represents the density of neurons in state (a, m) at time t.

The nonlinear evolution problem for the density ρ_t , for the initial datum u_0 , reads

(1.1a)
$$\partial_t \rho_t + \nabla \cdot (b\rho_t) = -f(a, m, \varepsilon x_t)\rho_t,$$

(1.1b)
$$\rho_t(0,m) = \mathbb{1}_{m > \gamma(0)} \left| (\gamma^{-1})'(m) \right| \int_0^\infty f(a,\gamma^{-1}(m),\varepsilon x_t) \rho_t(a,\gamma^{-1}(m)) da,$$

(1.1c)
$$x_t = \int_0^t \int_0^\infty \int_0^\infty h(t-s,a,m) f(a,m,\varepsilon x_s) \rho_s(a,m) dadm ds,$$

(1.1d) $\rho_0 = u_0.$

The dynamics of the model can be decomposed into three elements: (i) the behavior of neurons between spikes, (ii) the spike-triggered jumps, and (iii) the interaction between neurons, which we discuss in turn.

(i) Between spikes, neurons are transported along the vector field $b(a,m) = (1, -\lambda m)$, with $\lambda > 0$ ($\nabla \cdot$ denotes the divergence operator over the state-space).

(ii) Neurons spike at rate $f(a, m, \varepsilon x_t)$, where $f : \mathbb{R}_+ \times \mathbb{R}^+_+ \times \mathbb{R} \to \mathbb{R}_+$ is the "firing rate function" corresponding to the stochastic intensity of the spike generation process, and $\varepsilon \in \mathbb{R}$ is the connection strength. When a neuron spikes, its age a is reset to 0, and its leaky memory variable m jumps to $\gamma(m)$, where $\gamma : \mathbb{R}_+ \to \mathbb{R}^+_+$ is the "jump mapping" and is assumed to be a strictly increasing C^1 -diffeomorphism. As a consequence, the border condition (1.1b) has a simple interpretation: the density of neurons in state (0, m) at time t is equal to the marginal density of those neurons that have their leaky memory variable in state $\gamma^{-1}(m)$ and spike at time t. The indicator function $\mathbb{1}_{m > \gamma(0)}$ reflects the fact that m is always strictly positive, and the term $|(\gamma^{-1})'(m)|$ is necessary to guarantee the conservation of the total mass of neurons. Indeed, formally,

$$\partial_t \int \rho_t = \int \mathbb{1}_{m > \gamma(0)} \left| (\gamma^{-1})'(m) \right| \int_0^\infty f(a, \gamma^{-1}(m), \varepsilon x_t) \rho_t(a, \gamma^{-1}(m)) dadm - \int f(a, m, \varepsilon x_t) \rho_t = 0,$$

by a change of variable.

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(iii) Neurons interact through the "total postsynaptic potential" x_t , which integrates the past spiking activity of the population, filtered by the "interaction function" $h: \mathbb{R}_+ \times \mathbb{R}_+ \times \mathbb{R}^*_+ \to \mathbb{R}$, and which weighted by the connection strength $\varepsilon \in \mathbb{R}$, influences the firing rate f. If we write N(t) the population activity (the mean firing rate)

$$N(t) := \int_0^\infty \int_0^\infty f(a, m, \varepsilon x_t) \rho_t(a, m) da dm,$$

and if we take h independent of a and m, then x_t takes the form

$$x_t = \int_0^t h(t-s)N(s)ds,$$

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where h is now a simple delay kernel, as in [21, 18, 19, 35]. In our formulation, h in (1.1c) allows us to model more general interactions. For example, in subsection 1.2.2, we show that by choosing $h(t, a, m) = \hat{h}(t)(1 - m)$, we can include the effects of a classical short-term synaptic plasticity model [47].

1.2. Motivation. The model (1.1) extends the time-elapsed neuron network model [35] (see also [18, 19]) by the addition of a leaky memory variable which can accumulate over spikes (as opposed to the age variable which is reset to 0 at each spike) and hence introduces a slow timescale in the population dynamics. Such a slow timescale is typically used to account for some form of fatigue mechanism, which can act on the spiking activity (spike-frequency adaptation) or on synaptic transmission (short-term synaptic depression). Slow fatigue at the single-neuron level can lead to nontrivial emergent behaviors at the population level, such as population bursts [49, 22, 17] (see Figure 1), which have not been observed in the age- or voltagestructured models of [35] and [10] (but see [37]). Even though some population equations have been successfully used in the computational neuroscience literature to study emergent behaviors in networks of neurons with fatigue, these population equations were obtained at the cost of a timescale separation approximation [22, 17] or a "mixing" assumption [33, 44], making them inexact. In contrast, the model (1.1) is the exact mean-field limit [43] for spiking neuron models with spike-frequency adaptation or short-term synaptic depression, as we discuss now.

1.2.1. Spike-frequency adaptation. The recent spike history of a neuron can modulate its firing rate f, leading to spike-frequency adaptation [3]. If h is independent of a and m, and if $\gamma(m) = m + \hat{\Gamma}$ for a fixed $\hat{\Gamma} > 0$, (1.1) becomes

(1.2a)
$$\partial_t \rho_t + \nabla \cdot (b\rho_t) = -f(a, m, \varepsilon x_t)\rho_t,$$

(1.2b)
$$\rho_t(0,m) = \mathbb{1}_{m>\hat{\Gamma}} \int_0^\infty f(a,m-\hat{\Gamma},\varepsilon x_t)\rho_t(a,m-\hat{\Gamma})da,$$

(1.2c)
$$x_t = \int_0^t h(t-s) \int_0^\infty \int_0^\infty f(a,m,\varepsilon x_s) \rho_s(a,m) dadm ds,$$

(1.2d)
$$\rho_0 = u_0.$$

If $\eta : \mathbb{R}_+ \to \mathbb{R}$ is a bounded function such that $\lim_{a\to+\infty} \eta(a) = 0$ (η is the "refractory kernel" [20, sect. 9.3]), we can define f more explicitly as

(1.2e)
$$f(a, m, \varepsilon x_t) := \hat{f}(\eta(a) - m + \varepsilon x_t),$$

where $\hat{f} : \mathbb{R} \to \mathbb{R}_+$ is typically a nondecreasing function. Since *m* makes jumps of size $\hat{\Gamma} > 0$ at each spike and decays exponentially at rate λ between spikes, *m* accumulates over spikes, which decreases the firing rate *f* (1.2e), leading to spikefrequency adaptation [3]. More specifically, (1.2) is a population equation for adaptive SRM₀ (Spike Response Model) neurons [27, 20].

Populations of spiking neurons with spike-frequency adaptation exhibit selfsustained population bursts when the connectivity strength is sufficiently strong [49, 22, 17]. We call self-sustained population bursts a periodic pattern of activity characterized by an alternation between periods of low population activity and sequences of population spikes (short time intervals where almost all the neurons in the population fire). This definition is borrowed from the following definition of single-neuron bursting [26]: "When neuron activity alternates between a quiescent state and repetitive spiking, the neuron activity is said to be *bursting*." In Figure 1, we show simulations of



FIG. 1. Depending on the connectivity strength ε , a population of adaptive SRM_0 neurons can exhibit self-sustained bursts ($\varepsilon \gg 0$) or relaxation to a stationary state (small ε). We show simulations of a network of $5 \cdot 10^5$ adaptive SRM_0 neurons, approximating the mean-field limit (1.2), with identical parameters (except for ε) and identical initial conditions. The raster plots below the plots for the time-evolution of the total postsynaptic potential x_t represent the spikes of 100 randomly selected neurons.

(1.2) for two different connectivity strengths ε . For large ε , we observe self-sustained bursts, whereas for small ε , we observe relaxation to a stationary state. Note that the neurons considered here are not intrinsically bursting: if an adaptive SRM₀ neuron receives no input (or a constant input), it does not burst. Population bursts are therefore an emergent behavior of the mean-field model (1.2).

For comparison, in Appendix A we show similar simulations for the time-elapsed neuron network model [35], where, as expected, we only observe self-sustained oscillations or relaxation to a stationary state.

1.2.2. Short-term synaptic depression. The recent spike history of a presynaptic neuron can modulate the synaptic transmission, leading to short-term synaptic plasticity [51]. We will consider here the case of depressive synapses and use the model of [47] (with a change of variable for convenience). In this case, the state-space is $(a, m) \in \mathbb{R}_+ \times [0, 1[$. Taking f independent of m, and choosing h and γ of the form

 $h(t, a, m) := \tilde{h}(t)(1-m)$ and $\gamma(m) := 1 - v + vm$ for a fixed $v \in [0, 1[, (1.1) \text{ becomes}]$

(1.3a)
$$\partial_t \rho_t + \nabla \cdot (b\rho_t) = -f(a, \varepsilon x_t)\rho_t,$$

(1.3b)
$$\rho_t(0,m) = \mathbb{1}_{m > \gamma(0)} \frac{1}{\upsilon} \int_0^\infty f(a,\varepsilon x_t) \rho_t(a,\gamma^{-1}(m)) da,$$

(1.3c)
$$x_t = \int_0^t \hat{h}(t-s) \int_0^1 \int_0^\infty (1-m) f(a,\varepsilon x_s) \rho_s(a,m) dadm ds,$$

(1.3d)
$$\rho_0 = u_0.$$

Note that the term $\frac{1}{v}$ on the right-hand side of (1.3b) simply comes from the fact that $|(\gamma^{-1})'(m)| = \frac{1}{v}$ for all $m \in]0, 1[$. Here, at each spike, m makes strictly positive jumps whose size tends to 0 as m tends to 1 (since $\gamma(1) = 1$) and decays exponentially at rate λ between spikes. If m is close to 1, synaptic transmission is weak because of the factor (1 - m) in (1.3c).

As observed in [42], the stationary state of populations of neurons with shortterm synaptic plasticity can be described by a simple formula, which we rederive in subsection 4.3.

1.3. Assumptions and main results. The main result of this work is the exponential stability of (1.1) in the weak connectivity regime (Theorem 1.4)—or, more explicitly, there exists $\varepsilon^{**} > 0$ such that (1.1) is exponentially stable for all connectivity strength $\varepsilon \in] - \varepsilon^{**}, +\varepsilon^{**}[$. Before proving the exponential stability, we first establish the well-posedness of (1.1) in the appropriate function space (Theorem 1.2) and show that stationary solutions exist and are unique for sufficiently weak connectivity (Theorem 1.3).

Here, we study the weak solutions to (1.1) for an initial datum in $L^1_+ := L^1(\mathbb{R}_+ \times \mathbb{R}_+^*, \mathbb{R}_+)$ and write $L^1_+(\mathbb{R}_+^*) := L^1(\mathbb{R}_+^*, \mathbb{R}_+)$.

DEFINITION 1.1 (solutions). $(\rho, x) \in \mathcal{C}(\mathbb{R}_+, L^1_+) \times \mathcal{C}(\mathbb{R}_+)$ is a solution to (1.1), for the initial datum $u_0 \in L^1_+$, if (1.4a)

$$x_t = \int_0^t \int_0^\infty \int_0^\infty h(t - s, a, m) f(a, m, \varepsilon x_s) \rho_s(a, m) dadm ds \qquad \forall t \ge 0$$

and if for all $\varphi \in \mathcal{C}^{\infty}_{c}(\mathbb{R}_{+} \times \mathbb{R}_{+} \times \mathbb{R}_{+}^{*}),$

(1.4b)
$$0 = \int_0^\infty \int_0^\infty u_0(a,m)\varphi(0,a,m)dadm + \int_0^\infty \int_0^\infty \int_0^\infty \rho_t(a,m) \Big\{ [\partial_t + \partial_a - \lambda m \partial_m]\varphi + (\varphi(t,0,\gamma(m)) - \varphi(t,a,m))f(a,m,\varepsilon x_t) \Big\} dadmdt.$$

To prove the well-posedness of (1.1), we need some simple assumptions on the firing rate function f and the interaction function h.

ASSUMPTION 1. f is bounded and L_f -Lipschitz, i.e.,

$$|f(a,m,x) - f(a^*,m^*,x^*)| \le L_f(|a-a^*| + |m-m^*| + |x-x^*|),$$

and h is bounded and continuous.

Since we want to apply Harris's theorem, the well-posedness in L^1 (which is treated in [43]) is not enough, and we need the well-posedness in a weighted L^1 space (where the weight satisfies a Lyapunov condition [29]) with a global-in-time estimate in the weighted L^1 norm.

Using the weight function

$$w: \mathbb{R}_+ \times \mathbb{R}_+ \to [1, \infty[, (a, m) \mapsto 1 + m,$$

we define the function space

$$L^{1}_{+}(w) := \left\{ g \in L^{1}(\mathbb{R}_{+} \times \mathbb{R}^{*}_{+}, \mathbb{R}_{+}) \mid \|g\|_{L^{1}(w)} := \int_{0}^{\infty} \int_{0}^{\infty} g(a, m)w(a, m)dadm < \infty \right\}.$$

To obtain a global-in-time estimate in the $L^1_+(w)$ norm, we further require the jump sizes of γ to be bounded.

Assumption 2. There exists a bounded function $\Gamma : \mathbb{R}^*_+ \to \mathbb{R}^*_+$ such that for all $m \in \mathbb{R}^*_+, \ \gamma(m) = m + \Gamma(m).$

THEOREM 1.2 (well-posedness). Grant Assumption 1. For any initial datum $u_0 \in L^1_+$, there exists a unique weak solution (ρ, x) to (1.1). This solution satisfies

- (i) $(L^1$ -stability) $\|\rho_t\|_{L^1} = \|u_0\|_{L^1}$ for all t > 0.
- (ii) (Global bound in $L^1_+(w)$) If, in addition, Assumption 2 holds and $u_0 \in$ $L^1_+(w)$, then

(1.5)
$$\forall t > 0, \qquad \|\rho_t\|_{L^1(w)} \le \|u_0\|_{L^1(w)} e^{-\alpha t} + \frac{b}{\alpha} (1 - e^{-\alpha t}) \|u_0\|_{L^1}$$

for some constants $\alpha > 0$ and $b \in \mathbb{R}$.

In contrast to [43], the well-posedness proof presented here does not involve any probabilistic argument. The proof consists of two consecutive applications of Banach's fixed-point theorem, where a first fixed-point gives the unique solution to a linearized version of (1.1) which is then used in a second fixed-point treating the nonlinearity of (1.1).

The second step towards the exponential stability proof is the study of the existence and uniqueness of the stationary solutions to (1.1). For this step, we require the following assumption.

Assumption 3.

(i) There exist $\Delta_{abs} > 0$ and $\sigma > 0$ such that

$$f(a, m, x) \ge \sigma \qquad \forall (a, m, x) \in [\Delta_{abs}, +\infty[\times \mathbb{R}^*_+ \times \mathbb{R}.$$

- (ii) There exists $C_{\gamma} \in [0, 1]$ such that $C_{\gamma} \leq \gamma' \leq 1$. (iii) $\bar{h}(a, m) = \int_{0}^{\infty} h(t, a, m) dt$ is bounded.

The first point of Assumption 3 sets a lower bound on the firing rate function f for any $a > \Delta_{abs}$ while allowing neurons to have an absolute refractory period $\Delta_{abs} > 0$, i.e., a period of time following a spike during which f can be 0 (which is an important neurodynamical feature [20, sect. 1.1]). This assumption is also used in [5].

In the second point of Assumption 3, the lower bound $0 < C_{\gamma} \leq \gamma'$ guarantees that γ is strictly increasing, which reflects the idea that m is a "leaky memory" variable of the past neuronal activity. On the other hand, the upper bound $\gamma' \leq 1$, which can be rewritten in terms of the jump size function Γ as $\Gamma' \leq 0$, prevents the variable m

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from growing too fast and allows for a potential saturation of the memory, as in the example with short-term synaptic plasticity (1.3). The third point of Assumption 3 reflects the fact that a single spike has a finite impact on the neuron that receives it.

We emphasize that the two examples shown above—spike-frequency adaptation (1.2) and short-term synaptic depression (1.3)—satisfy Assumption 3.

THEOREM 1.3 (stationary solutions). Grant Assumptions 1–3.

- (i) There exists a stationary solution to (1.1).
- (ii) There exists $\varepsilon^* > 0$ such that for all $\varepsilon \in] \varepsilon^*, +\varepsilon^*[$, the stationary solution to (1.1) is unique.

Over the course of this work, we obtained the existence of the stationary solution by two different approaches. The first approach is based on the Doeblin–Harris method [23] and is similar to that of [5]. First, we show that when x_t is fixed and time-invariant in (1.1) (neurons are noninteracting), the system satisfies a Harris condition—this constitutes a key result of this work—and we can use Harris's theorem to get the stationary solution. Then, we use the Lipschitz continuity of the stationary solutions with respect to the fixed x to prove the existence of a stationary solution for arbitrary connectivity strengths ε . Finally, for ε small enough, we also get the uniqueness of the stationary solution by Banach's fixed-point theorem.

The second approach relies on the fact that the stationary solutions solve an integral equation, for which we can show that a solution exists by Schauder's fixed-point theorem. In the process, we get several estimates on the stationary solutions, namely that they are continuous, bounded, and exponentially decaying in m. However, this approach does not give uniqueness.

As mentioned above, the application of Harris's theorem requires us to consider solutions in the weighted space $L^1(w)$. However, in the case where the state-space of the leaky memory variable m is bounded, the situation is simpler: we can use Doeblin's theorem in L^1 . The following assumption guarantees that m stays in a bounded state-space.

Assumption 4. There exists G > 0 such that for all $m \in \mathbb{R}^*_+$, $\gamma(m) < G$.

Note that this assumption is satisfied in the example with short-term synaptic plasticity (1.3), with G = 1.

Finally, to study the exponential stability of (1.1), we need an exponential decay on h.

Assumption 5. There exists $\mathfrak{h}, C_h > 0$ such that $h(t, a, m) \leq C_h e^{-\mathfrak{h}t}$ for all $(t, a, m) \in \mathbb{R}_+ \times \mathbb{R}_+ \times \mathbb{R}_+^*$.

By a perturbation argument similar to that of [30], we obtain our main result.

THEOREM 1.4 (exponential stability in the weak connectivity regime). Grant Assumptions 1–3 and 5. For any W > 0, there exists $\varepsilon_W^{**} > 0$ such that for $\varepsilon \in$ $] - \varepsilon_W^{**}, + \varepsilon_W^{**}[$, there exist $C \ge 1$ and $c_W > 0$ such that for all initial data $u_0 \in L^1_+(w)$ with $\|u_0\|_{L^1} = 1$ and $\|u_0\|_{L^1(w)} \le W$, the solution (ρ, x) to (1.1) satisfies

(1.6)
$$\|\rho_t - \rho_\infty\|_{L^1(w)} + |x_t - x_\infty| \le C e^{-c_W t} \left(\|u_0 - \rho_\infty\|_{L^1(w)} + 1 \right) \quad \forall t \ge 0,$$

where $(\rho_{\infty}, x_{\infty})$ is the unique stationary solution given by Theorem 1.3(ii).

If, in addition, we grant Assumption 4, then there exists $\varepsilon^{**} > 0$ such that for all $\varepsilon \in] - \varepsilon^{**}, +\varepsilon^{**}[$, there exist $C' \geq 1$ and c > 0 such that for all initial data $u_0 \in L^1_+$

with $||u_0||_{L^1} = 1$,

(1.7)
$$\|\rho_t - \rho_\infty\|_{L^1} + |x_t - x_\infty| \le C' e^{-ct} \left(\|u_0 - \rho_\infty\|_{L^1} + 1 \right) \quad \forall t \ge 0.$$

From the neuronal modeling point of view, this result is not surprising: when the connection strength is weak enough, neurons do not synchronize, and the population activity converges to a stationary state. This was already proved for simpler one-dimensional models (see below), and the addition of a leaky memory variable carrying the effect of spike-frequency adaptation or short-term synaptic plasticity does not change this behavior.

1.4. Discussion of the methods. The asymptotic stability of the age-structured model of [35] in the weak connectivity regime has been studied using entropy methods (assuming that f is a step-function) [35, 36], spectral analysis of semigroups in Banach spaces [31, 30], and Doeblin's theorem [5]. For treatments of the strong connectivity regime, we refer the reader to [35, 36, 30].

The asymptotic stability of the closely related voltage-structured model of [10] in the weak connectivity regime has also been studied by Cormier, Tanré, and Veltz [8] using Laplace transform techniques. In addition, the same authors have analyzed the nonlinear stability of the stationary solutions [7] (see also [12]) and proved the existence of periodic solutions [9].

Doeblin's theorem has also been used in [14] in the case of the "threshold crossing" neuronal population equation of [34]. Note that closely related methods have been used by probabilists to study the ergodicity of single-neuron models [25, 13].

Our approach combines strategies from [31] and [5], even though [5] uses Doeblin's instead of Harris's theorem. On the one hand, our proof is based on the application of Harris's theorem for the linear problem, which simplifies the proof of [31]. On the other hand, we use an argument from [31] to deal with delay effects, which are not considered in [5]. Note that our model is two-dimensional (by the addition of the leaky memory variable), whereas the aforementioned works only considered one-dimensional models.

1.5. Plan of the paper. The proof of Theorem 1.2 (well-posedness) is presented in section 2. In section 3, we prove the exponential stability of (1.1) in the noninteracting case $\varepsilon = 0$ using Harris's or Doeblin's theorem. The proof of Theorem 1.3 (stationary solutions) is presented in section 4 which is divided into three parts: in the first part, we present a proof which uses the exponential stability of the noninteracting case; in the second part, we present an alternative proof for the existence of stationary solutions which does not involve the Doeblin–Harris method; and in the last part, we present a proof for the formula of [42] in the case of short-term synaptic plasticity (1.3). Finally, section 5 is dedicated to the proof of Theorem 1.4 (exponential stability in the weak connectivity regime).

2. Well-posedness. This section is dedicated to the proof of Theorem 1.2, which we decompose into several lemmas. First, we verify the a priori L^1 -stability of the solutions to (1.1), a technical result we use later in the proof. Then, we introduce a linearized version of (1.1) and show that it is well-posed by an application of Banach's fixed-point theorem. Another application of Banach's fixed-point theorem is used to treat the nonlinearity of (1.1) and concludes the proof of the well-posedness in L^1 . Finally, we prove the global bound in $L^1_+(w)$ (Theorem 1.2(ii)), which we will use to apply Harris's theorem in the next sections.

LEMMA 2.1 (a priori L^1 -stability). Grant Assumption 1. If (ρ, x) is a weak solution to (1.1) for the initial datum $u_0 \in L^1_+$, then

$$\|\rho_t\|_{L^1} = \|u_0\|_{L^1} \qquad \forall t > 0.$$

Proof. By a standard cut-off in time argument, we have that for all T > 0 and for all $\varphi \in C_c^{\infty}(\mathbb{R}_+ \times \mathbb{R}_+ \times \mathbb{R}_+^*)$,

$$\begin{split} \int_0^\infty \int_0^\infty \rho_T(a,m)\varphi(T,a,m)dadm &- \int_0^\infty \int_0^\infty u_0(a,m)\varphi(0,a,m)dadm \\ &= \int_0^T \int_0^\infty \int_0^\infty \rho_t(a,m) \Big\{ [\partial_t + \partial_a - \lambda m \partial_m]\varphi + (\varphi(t,0,\gamma(m)) \\ &- \varphi(t,a,m))f(a,m,\varepsilon x_t) \Big\} dadmdt \end{split}$$

Let χ be a function in $\mathcal{C}_c^{\infty}(\mathbb{R}_+ \times \mathbb{R}^*_+, \mathbb{R}_+)$ such that

$$\chi(a,m) = 1 \quad \forall \ a^2 + m^2 \le 1.$$

For all $n \in \mathbb{N}^*$, we write $\tilde{\varphi}^n \in \mathcal{C}^{\infty}(\mathbb{R}_+ \times \mathbb{R}_+ \times \mathbb{R}_+^*)$, the classical solution to the transport equation

(2.1a)
$$\partial_t \widetilde{\varphi}^n(t, a, m) + \partial_a \widetilde{\varphi}^n(t, a, m) - \lambda m \partial_m \widetilde{\varphi}^n(t, a, m) = 0,$$

(2.1b)
$$\widetilde{\varphi}^n(0, a, m) = \chi(a/n, m/n).$$

Because of the finite speed of propagation of the transport equation, for all n, there exists a function $\varphi^n \in \mathcal{C}^{\infty}_c(\mathbb{R}_+ \times \mathbb{R}_+ \times \mathbb{R}^*_+)$ such that $\varphi^n(t, a, m) = \tilde{\varphi}^n(t, a, m)$ for all $(t, a, m) \in [0, T] \times \mathbb{R}_+ \times \mathbb{R}^*$. Hence, for all $n \in \mathbb{N}^*$,

$$\begin{split} \int_0^\infty \int_0^\infty \varphi^n(T, a, m) \rho_T(a, m) dadm &- \int_0^\infty \int_0^\infty \varphi^n(0, a, m) u_0(a, m) dadm \\ &= \int_0^T \int_0^\infty \int_0^\infty \bigg\{ \partial_t \varphi^n + \partial_a \varphi^n - \lambda m \partial_m \varphi^n + \big(\varphi^n(t, 0, \gamma(m)) \\ &- \varphi^n(t, a, m)\big) f(a, m, \varepsilon x_t) \bigg\} \rho_t(a, m) dadm dt \end{split}$$

As φ_n is a solution to (2.1a) on time [0, T], we get

$$\int_0^\infty \int_0^\infty \varphi^n(T, a, m) \rho_T(a, m) dadm - \int_0^\infty \int_0^\infty \varphi^n(0, a, m) u_0(a, m) dadm$$
$$= \int_0^T \int_0^\infty \int_0^\infty \left\{ \left(\varphi^n(t, 0, \gamma(m)) - \varphi^n(t, a, m) \right) f(a, m, \varepsilon x_t) \right\} \rho_t(a, m) dadm dt$$

For all $(t, a, m) \in [0, T] \times \mathbb{R}_+ \times \mathbb{R}^*_+$, $\varphi^n(t, a, m) \xrightarrow[n \to \infty]{n \to \infty} 1$, since the initial datum tends to 1 as $n \to \infty$ (2.1b) and by finite speed of propagation. Thus, by dominated convergence, we get

(2.2)
$$\int_0^\infty \int_0^\infty \rho_T(a,m) dadm - \int_0^\infty \int_0^\infty u_0(a,m) dadm = 0.$$

Since ρ is nonnegative, this concludes the proof.

Lemma 2.1 will allow us to prove the well-posedness of (1.1) by means of fixedpoint arguments. Let us first introduce a linearized version of (1.1): for all $x \in \mathcal{C}(\mathbb{R}_+)$, we consider the linear evolution problem

(2.3a)
$$\partial_t \rho_t + \nabla \cdot (b\rho_t) = -f(a, m, \varepsilon x_t)\rho_t,$$

(2.3b)
$$\rho_t(0,m) = \mathbb{1}_{m > \gamma(0)} \left| (\gamma^{-1})'(m) \right| \int_0^\infty f(a,\gamma^{-1}(m),\varepsilon x_t) \rho_t(a,\gamma^{-1}(m)) da$$

(2.3c) $\rho_0 = u_0.$

We can see (2.3) as the Kolmogorov forward equation of a time-dependent Markov process. Indeed, we can rewrite (2.3a) and (2.3b) as

(2.4)
$$\partial \rho_t = \mathcal{L}_t \rho_t,$$

where, for all suitable test functions $\phi : \mathbb{R}_+ \times \mathbb{R}^*_+ \to \mathbb{R}$,

(2.5)
$$\mathcal{L}_t^*\phi(a,m) = b(a,m) \cdot \nabla \phi(a,m) + [\phi(0,\gamma(m)) - \phi(a,m)]f(a,m,\varepsilon x_t)$$

 \mathcal{L}_t^* is the time-dependent generator of a piecewise deterministic Markov process with degenerate jumps.

The linearized equation (2.3) will play a special role in the following sections and therefore deserves its own proposition.

PROPOSITION 2.2 (well-posedness of the linearized equation (2.3)). Grant Assumption 1. For any initial datum $u_0 \in L^1_+$ and any $x \in \mathcal{C}(\mathbb{R}_+)$, there exists a unique weak solution $\rho^x \in \mathcal{C}(\mathbb{R}_+, L^1_+)$ to (2.3). Furthermore, ρ^x satisfies the following:

(i) For all t > 0 and for all $m \in \mathbb{R}^*_+$,

$$\rho_t^x(0,m) = \mathbb{1}_{m > \gamma(0)} \left| (\gamma^{-1})'(m) \right| \int_0^\infty f(a,\gamma^{-1}(m),\varepsilon x) \rho_t^x(a,\gamma^{-1}(m)) da,$$

$$\rho_t^x(a,m) = \int u_0(a-t,e^{\lambda t}m) \exp\left(\lambda t - \int_0^t f(a-t+s,e^{\lambda(t-s)}m,\varepsilon x) ds\right) if a \ge t,$$

$$\int \rho_{t-a}^x(0, e^{\lambda a}m) \exp\left(\lambda a - \int_{t-a}^t f(a-t+s, e^{\lambda(t-s)}m, \varepsilon x) ds\right) if 0 < a < t.$$

(ii) For all t > 0 and for all $\phi \in \mathcal{C}_c^{\infty}(\mathbb{R}_+ \times \mathbb{R}_+^*)$,

(2.6)
$$\langle \rho_t^x, \phi \rangle = \langle u_0, \phi \rangle + \int_0^t \langle \rho_s^x, \mathcal{L}_s^* \phi \rangle ds$$

Proof. Fix $x \in \mathcal{C}(\mathbb{R}_+)$. For all $p \in \mathcal{C}(\mathbb{R}_+, L^1_+(\mathbb{R}^*_+))$ and $u_0 \in L^1_+$, we know from the standard theory of transport equations that there is a unique weak solution to

$$\partial_t \rho_t + \nabla \cdot (b\rho_t) = -f(a, m, \varepsilon x_t)\rho_t,$$

$$\rho_t(0, m) = p_t(m),$$

$$\rho_0 = u_0,$$

which we denote by $\rho^{x,p}$, and it is given by the representation formula,

$$\rho_t^{x,p}(a,m) := \begin{cases} u_0(a-t, e^{\lambda t}m) \exp\left(\lambda t - \int_0^t f(a-t+s, e^{\lambda(t-s)}m, \varepsilon x_s)ds\right) \text{ if } a \ge t, \\ p_{t-a}(e^{\lambda a}m) \exp\left(\lambda a - \int_{t-a}^t f(a-t+s, e^{\lambda(t-s)}m, \varepsilon x_s)ds\right) \text{ if } 0 < a < t \end{cases}$$

The solution $\rho^{x,p}$ is in $\mathcal{C}(\mathbb{R}_+, L^1)$ since

$$\forall t \in \mathbb{R}_+, \qquad \left\| \rho_t^{x,p} \right\|_{L^1} \le \|u_0\|_{L^1} + \int_0^t \|p_s\|_{L^1} \, ds$$

We have

$$\left(\left. \mathbb{1}_{m > \gamma(0)} \left| (\gamma^{-1})'(m) \right| \int_0^\infty f(a, \gamma^{-1}(m), x) \rho_t^{x, p}(a, \gamma^{-1}(m)) da \right)_{(t, m) \in \mathbb{R}_+ \times \mathbb{R}_+^*} \\ \in \mathcal{C}(\mathbb{R}_+, L^1_+(\mathbb{R}_+^*))$$

since

$$\begin{aligned} \forall t \in \mathbb{R}_+, \qquad \int_{\gamma(0)}^{\infty} \left| (\gamma^{-1})'(m) \right| \int_0^{\infty} f(a, \gamma^{-1}(m), x) \rho_t^{x, p}(a, \gamma^{-1}(m)) dadm \\ \leq \|f\|_{\infty} \left\| \rho_t^{x, p} \right\|_{L^1} \leq \|f\|_{\infty} \left(\|u_0\|_{L^1} + \int_0^t \|p_s\|_{L^1} \, ds \right) \end{aligned}$$

Hence, we can define, for any T > 0, the operator Φ_T^x :

$$\begin{aligned} \mathcal{C}([0,T], L^{1}_{+}(\mathbb{R}^{*}_{+})) &\to \mathcal{C}([0,T], L^{1}_{+}(\mathbb{R}^{*}_{+})) \\ p &\mapsto \left(\mathbb{1}_{m > \gamma(0)} \left| (\gamma^{-1})'(m) \right| \int_{0}^{\infty} f(a, \gamma^{-1}(m), x) \rho_{t}^{x, p}(a, \gamma^{-1}(m)) da \right)_{(t,m) \in [0,T] \times \mathbb{R}^{*}_{+}} \end{aligned}$$

For any $p, q \in \mathcal{C}([0, T], L^1_+(\mathbb{R}^*_+)),$

$$\begin{split} \left\| \Phi_T^x(p) - \Phi_T^x(q) \right\|_{\mathcal{C}([0,T],L^1)} &\leq \|f\|_{\infty} \sup_{t \in [0,T]} \left\| \rho_t^{x,p} - \rho_t^{x,q} \right\|_{L^1} \\ &\leq \|f\|_{\infty} \int_0^T \|p_s - q_s\|_{L^1} \, ds \\ &\leq T \|f\|_{\infty} \|p - q\|_{\mathcal{C}([0,T],L^1)} \, . \end{split}$$

Therefore, if $0 < T < ||f||_{\infty}^{-1}$, Φ_T^x is a contraction. By Banach's fixed-point theorem, there exists a unique $\rho^x \in \mathcal{C}([0,T], L_+^1)$ solving (2.3). Since the choice of the contracting T does not depend on the initial datum, we can iterate the above argument on successive time intervals of length T and conclude that there exists a unique $\rho^x \in \mathcal{C}(\mathbb{R}_+, L_+^1)$ solving (2.3), for which formula (i) is satisfied. Then, (ii) follows from a standard cut-off-in-time argument.

Now, we can prove the existence and uniqueness of a solution to the nonlinear problem (1.1) by means of a second application of Banach's fixed-point theorem.

Proof of the well-posedness of (1.1) in L^1 . For any $x \in \mathcal{C}(\mathbb{R}_+)$, we take the ρ^x given by Proposition 2.2. We have

$$\left(\int_0^t \int_{\mathbb{R}_+ \times \mathbb{R}_+^*} h(t-s) f(\varepsilon x_s) \rho_s^x \, dadm ds\right)_{t \in \mathbb{R}_+} \in \mathcal{C}(\mathbb{R}_+)$$

since

$$\forall t \in \mathbb{R}_+, \qquad \left| \int_0^t \int_{\mathbb{R}_+ \times \mathbb{R}^*_+} h(t-s) f(\varepsilon x_s) \rho_s^x \, dadm ds \right| \le \|h\|_{\infty} \|f\|_{\infty} \int_0^t \|\rho_s^x\|_{L^1} \, ds \le \|h\|_{\infty} \|f\|_{\infty} \int_0^t \|h\|_{\infty} \|f\|_{\infty} \int_0^t \|\rho_s^x\|_{L^1} \, ds \le \|h\|_{\infty} \|f\|_{\infty} \|f\|_{\infty$$

Hence, for any T > 0, we can define the operator

$$\begin{split} \Psi_T : \mathcal{C}([0,T]) &\to \mathcal{C}([0,T]) \\ x \mapsto \left(\int_0^t \int_{\mathbb{R}_+ \times \mathbb{R}^*_+} h(t-s) f(\varepsilon x_s) \rho_s^x \, dadm ds \right)_{t \in [0,T]} \end{split}$$

For any $x, y \in \mathcal{C}([0, T])$, we have

$$\begin{split} \|\Psi_{T}(x) - \Psi_{T}(y)\|_{\mathcal{C}([0,T])} &\leq T \|h\|_{\infty} \sup_{t \in [0,T]} \int_{\mathbb{R}_{+} \times \mathbb{R}^{*}_{+}} |f(\varepsilon x_{t})\rho_{t}^{x} - f(\varepsilon y_{t})\rho_{t}^{y}| \, dadm \\ &\leq T \|h\|_{\infty} \sup_{t \in [0,T]} \left(\varepsilon L_{f} |x_{t} - y_{t}| \, \|\rho_{t}^{x}\|_{L^{1}} + \|f\|_{\infty} \|\rho_{t}^{x} - \rho_{t}^{y}\|_{L^{1}}\right). \end{split}$$

By Grönwall's lemma, $\|\rho_t^x\|_{L^1} \leq \|u_0\|_{L^1} \exp(\|f\|_\infty t)$ since

$$\forall t \in [0,T], \qquad \|\rho_t^x\|_{L^1} \le \|u_0\|_{L^1} + \|f\|_{\infty} \int_0^t \|\rho_s^x\|_{L^1} \, ds$$

On the other hand, we have, for all $t \in [0, T]$,

$$\begin{split} \left\| \rho_t^x - \rho_t^y \right\|_{L^1} &\leq \int_0^t \int_0^\infty \left| \rho_s^x(0,m) \exp\left(-\int_s^t f(u-s, e^{-\lambda(u-s)}\gamma(m), \varepsilon x_u) du \right) \right. \\ &\left. - \rho_s^y(0,m) \exp\left(-\int_s^t f(u-s, e^{-\lambda(u-s)}\gamma(m), \varepsilon y_u) du \right) \right| dm ds \\ &\leq \|f\|_\infty \int_0^t \|\rho_s^x - \rho_s^y\|_{L^1} ds + t\varepsilon \|f\|_\infty L_f \|x-y\|_{\mathcal{C}([0,T])} \int_0^t \|\rho_s^x\|_{L^1} ds. \end{split}$$

Hence, by Grönwall's lemma, for all $t \in [0, T]$,

$$\left\|\rho_t^x - \rho_t^y\right\|_{L^1} \le \varepsilon L_f \|u_0\|_{L^1} \frac{\left(\exp(\|f\|_{\infty} t) - 1\right)^2}{\|f\|_{\infty}} \|x - y\|_{\mathcal{C}([0,T])}.$$

Gathering the bounds, we get

$$\begin{aligned} \left\| \Psi_T(x) - \Psi_T(y) \right\|_{\mathcal{C}([0,T])} \\ &\leq T \varepsilon \|h\|_{\infty} L_f \|u_0\|_{L^1} \exp(\|f\|_{\infty} T) \left[1 + \exp(\|f\|_{\infty} T) \right] \|x - y\|_{\mathcal{C}([0,T])} \end{aligned}$$

For T small enough, Ψ_T is a contraction and, by Banach's fixed-point theorem, has a unique fixed-point. Thus, there exists a unique solution $(\rho, x) \in \mathcal{C}([0, T], L^1_+)$. Since, by Lemma 2.1, $\|\rho_T\|_{L^1} = \|u_0\|_{L^1}$, we can iterate this argument on successive time intervals of length T and conclude that there exists a unique solution in $\mathcal{C}(\mathbb{R}_+, L^1_+)$.

To conclude the proof of Theorem 1.2, it remains to show the estimate (1.5). Under Assumption 2, the weight function

$$w: \mathbb{R}_+ \times \mathbb{R}_+ \to [1, \infty[, \qquad (a, m) \mapsto 1 + m]$$

satisfies $w(a,m) \to \infty$ when $m \to \infty$ and the Lyapunov condition on m:

(2.7)
$$\exists \alpha > 0, b \ge 0$$
 such that $\mathcal{L}_t^* w \le -\alpha w + b.$

Indeed, for all $(t, a, m) \in \mathbb{R}_+ \times \mathbb{R}_+ \times \mathbb{R}_+^*$,

$$\mathcal{L}_{t}^{*}w(a,m) = -\lambda m + \Gamma(m)f(a,m,\varepsilon x_{t}) \leq -\lambda w(a,m) + \lambda + \|\Gamma\|_{\infty} \|f\|_{\infty}$$

Importantly, the constants α and b do not depend on x.

LEMMA 2.3 (global bound in $L^1_+(w)$). Grant Assumptions 1 and 2. If the initial datum u_0 is in $L^1_+(w)$, then $\rho_t \in L^1_+(w)$ for all $t \ge 0$. Moreover,

(2.8)
$$\forall t > 0, \qquad \|\rho_t\|_{L^1(w)} \le \|u_0\|_{L^1(w)} e^{-\alpha t} + \frac{b}{\alpha} (1 - e^{-\alpha t}) \|u_0\|_{L^1},$$

where the constants α and b are taken from the Lyapunov condition (2.7).

Proof. We divide the proof into two steps. First, we prove that the solution is stable in $L^1_+(w)$ with a weaker and time-dependent bound; then we use this first bound to apply the dominated convergence theorem and obtain (2.8) by Grönwall's lemma.

Step 1. Fix any T > 0. Let $\chi \in C_c^{\infty}(\mathbb{R}_+, \mathbb{R}_+)$ be a nonincreasing function such that $\chi(x) = 1$ if $0 \le x \le 1$ and $\chi(x) = 0$ if x > 2. For all $n \in \mathbb{N}^*$, let us write $\varphi_k(a)\chi_n(m) := \chi(a/k)\chi(m/n)$. We also consider $g_M(w)$ a smooth approximation of $w \land M$, such that $\|g'\|_{\infty} \le 1$ and $M \mathbb{1}_{w \ge M} \le g(w) \le M$. For all n, k, and M, $g_M(w)\chi_n\varphi_k \in \mathcal{C}_c^{\infty}(\mathbb{R}_+ \times \mathbb{R}_+, \mathbb{R}_+)$. Hence (by Proposition 2.2(ii)) the solution (ρ, x) to (1.1) satisfies

$$\forall n \in \mathbb{N}^*, \qquad \langle \rho_T, g_M(w) \chi_n \varphi_k \rangle = \langle u_0, g_M(w) \chi_n \varphi_k \rangle + \int_0^T \langle \rho_t, \mathcal{L}_x^*(g_M(w) \chi_n \varphi_k) \rangle dt,$$

where

$$\begin{aligned} \mathcal{L}_{x}^{*}(g_{M}(w)\chi_{n}\varphi_{k}) &= \partial_{a}(g_{M}(w)\chi_{n}\varphi_{k}) - \lambda m\partial_{m}(g_{M}(w)\chi_{n}\varphi_{k}) \\ &+ \left(g_{M}(w(\gamma(m)))\chi_{n}(\gamma(m))\varphi_{k}(0) - g_{M}(w)\chi_{n}\varphi_{k}\right)f \\ &= g_{M}(w)\chi_{n}\frac{1}{k}\chi'(a/k) - \lambda mg_{M}(w)\varphi_{k}\frac{1}{n}\chi'(m/n) - \lambda mg'_{M}(w)\chi_{n}\varphi_{k} \\ &+ \left(g_{M}(w(\gamma(m)))\chi_{n}(\gamma(m))\varphi_{k}(0) - g_{M}(w)\chi_{n}\varphi_{k}\right)f. \end{aligned}$$

From the L^1 -stability and the fact that both $g_M(w)\partial_m\chi_n$ and $g_M(w)\chi_n$ are bounded and have compact support, we can go to the limit in k by dominated convergence:

(2.9)
$$\langle \rho_T, g_M(w)\chi_n \rangle = \langle u_0, g_M(w)\chi_n \rangle$$

 $+ \int_0^T \left\langle \rho_t, -\lambda m g_M(w) \frac{1}{n} \chi'(m/n) - \lambda m g'_M(w)\chi_n \right\rangle dt$
 $+ \int_0^T \left\langle \rho_t, (g_M(w(\gamma(m)))\chi_n(\gamma(m)) - g_M(w)\chi_n)f \right\rangle dt.$

On the other hand, from the properties of χ and g_M , we have

$$\left|\lambda m g_M(w) \frac{1}{n} \chi'(m/n)\right| \le \lambda g_M(w) \frac{2n}{n} \|\chi'\|_{\infty} \le 2\lambda M \|\chi'\|_{\infty}$$

and

$$|\lambda m g'_M(w)\chi_n| \le \lambda g_M(w) \le \lambda M,$$

whence

$$\begin{aligned} \langle \rho_T, g_M(w)\chi_n \rangle &\leq \langle u_0, g_M(w)\chi_n \rangle \\ &+ \int_0^T \langle \rho_t, \lambda g_M(w) \|\chi'\|_\infty + \lambda g_M(w) \rangle dt \\ &+ \int_0^T \langle \rho_t, (g_M(w(\gamma(m)))\chi_n(\gamma(m)) - g_M(w)\chi_n)f \rangle dt, \end{aligned}$$

and we can take the limit in n by dominated convergence:

$$\begin{aligned} \langle \rho_T, g_M(w) \rangle &\leq \langle u_0, g_M(w) \rangle \\ &+ \int_0^T \langle \rho_t, \lambda g_M(w) \| \chi' \|_\infty + \lambda g_M(w) \rangle dt \\ &+ \int_0^T \langle \rho_t, (g_M(w(\gamma(m))) - g_M(w)) f \rangle dt. \end{aligned}$$

From the properties of γ , we get

$$w(0, \gamma(m)) \le w(0, m + \|\Gamma\|_{\infty}) \le (1 + \|\Gamma\|_{\infty})w(a, m)$$

and

$$g_M(w(0,\gamma(m))) \le (1 + \|\Gamma\|_{\infty})g_M(w(m))$$

This, together with the fact that f is bounded, shows that there exists a constant C, which does not depend on M, such that

$$\langle \rho_T, g_M(w) \rangle \leq \langle u_0, g_M(w) \rangle + C \int_0^T \langle \rho_t, g_M(w) \rangle dt,$$

and we can apply Grönwall's lemma to obtain

$$\langle \rho_T, g_M(w) \rangle \le \langle u_0, g_M(w) \rangle e^{Ct}.$$

Finally, it follows from Fatou lemma that $\rho_T \in L^1_+(w)$.

Step 2. To improve the previous estimate, we come back to (2.9) and use dominated convergence in n and M (domination being guaranteed by Step 1) to show

$$\langle \rho_T, w \rangle = \langle u_0, w \rangle + \int_0^T \langle \rho_t, \mathcal{L}_x^* w \rangle dt$$

By the Lyapunov condition (2.7),

$$\|\rho_T\|_{L^1(w)} \le \|u_0\|_{L^1(w)} - \alpha \int_0^T \|\rho_t\|_{L^1(w)} \, dt + Tb$$

and by Grönwall's lemma, we have, for all $t \in [0, T]$,

$$\|\rho_T\|_{L^1(w)} \le \|u_0\|_{L^1(w)} e^{-\alpha t} + \frac{b}{\alpha} (1 - e^{-\alpha t})$$

Since T can be chosen arbitrarily large, this achieves the proof.

Remark 2.4. Following the same steps as in the proof above, we can show that the bound (2.8) also holds for the linearized equation (2.3) and does not depend on x or the constants α and b.

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3. Exponential stability in the noninteracting case. If $x \in C(\mathbb{R}_+)$ in the linearized equation (2.3) is time-invariant, i.e., $x \equiv \tilde{x}$ for some $\tilde{x} \in \mathbb{R}$, then (2.3) can be seen as the dynamics of a noninteracting population of neurons. In this section, we prove the exponential stability in the noninteracting case using Harris's or Doeblin's theorem. This is the key result of this work and will allow us to prove the existence and uniqueness of the stationary solution to (1.1) (section 4) and the exponential convergence towards it (section 5).

For $\tilde{x} \in \mathbb{R}$, $u_0 \in L^1$, we denote by $\rho_t^{\tilde{x}}$ the unique solution to (2.3) for the initial datum u_0 and for $x \equiv \tilde{x}$ given by Proposition 2.2. We write, using the semigroup notation,

$$(3.1) S_t^{\tilde{x}} u_0 := \rho_t^{\tilde{x}} \forall t \ge 0$$

To show that the semigroup (3.1) is exponentially stable, we will use Harris's theorem in the general case or Doeblin's theorem if Assumption 4 is granted. The original theorems of Doeblin [11] and Harris [24] have since been refined and extended—see the well-known works of Meyn and Tweedie [29] and Hairer and Mattingly [23]. More recently, these theorems have been generalized to stochastic semigroups [16, 5, 1, 4]. Below, we give general statements of Doeblin's and Harris's theorems. For completeness, a short yet enlightening proof of Doeblin's theorem is presented in Appendix B. A proof of Harris's theorem can be found in the recent work of Cañizo and Mischler [4].

Let Ω denote a general state-space, and let $(S_t)_{t\geq 0}$ be a stochastic semigroup; i.e., for all $t \geq 0$, S_t is a mass and positivity preserving linear operator on $L^1(\Omega)$; S_0 is the identity operator; and for all $t, s \geq 0$, $S_t S_s = S_{s+t}$. We say that $\rho_{\infty} \in L^1_+(\Omega)$ is an *invariant probability measure* of the semigroup $(S_t)_{t\geq 0}$ if $\|\rho_{\infty}\|_{L^1} = 1$ and if, for all $t \geq 0$, $S_t \rho_{\infty} = \rho_{\infty}$.

THEOREM 3.1 (Doeblin). If there exist T > 0 and a nonzero $\nu \in L^1_+$ such that

(3.2)
$$S_T u_0 \ge \nu \|u_0\|_{L^1} \quad \forall u_0 \in L^1_+.$$

then there exists a unique invariant probability measure ρ_{∞} , and for all initial data $u_0 \in L^1_+$ with $\|u_0\|_{L^1} = 1$,

(3.3)
$$\|S_t u_0 - \rho_\infty\|_{L^1} \le K e^{-\mathfrak{a} t} \|u_0 - \rho_\infty\|_{L^1} \quad \forall t \ge 0,$$

with

$$K = \frac{1}{1 - \|\nu\|_{L^1}}; \quad \mathfrak{a} = -\frac{\log(1 - \|\nu\|_{L^1})}{T} > 0.$$

We call (3.2) the Doeblin minoration condition. Very loosely speaking, the Doeblin minoration condition is best suited for compact state-spaces (but see [16, 5] for examples on \mathbb{R}_+). In the case of unbounded state-spaces, Harris's theorem tells us that the Doeblin minoration condition can be relaxed to a more local form if there is a Lyapunov-type localization condition.

THEOREM 3.2 (Harris). Let $w : \Omega \to [1, +\infty)$ be a measurable weight function. If there exists T > 0 such that

(i) (operator Lyapunov condition) there exist $A \in [0, 1]$ and $B \ge 0$ such that

$$(3.4) ||S_T u_0||_{L^1(w)} \le A ||u_0||_{L^1(w)} + B ||u_0||_{L^1} \forall u_0 \in L^1_+(w);$$

(ii) (Harris minoration condition) there exist a nonzero $\nu \in L^1_+$ and R > 0 such that

(3.5)
$$S_T u_0 \ge \nu \int_{\mathfrak{C}} u_0 \quad \forall u_0 \in L^1_+, \quad with \ \mathfrak{C} = \{x \in \Omega \mid w(x) \le R\},\$$

then there exists a unique invariant probability measure ρ_{∞} such that $\rho_{\infty} \in L^{1}_{+}(w)$,¹ and there exist $K \geq 1$ and $\mathfrak{a} > 0$ such that for all initial data $u_0 \in L^1_+(w)$ with $||u_0||_{L^1} = 1,$

(3.6)
$$\|S_t u_0 - \rho_\infty\|_{L^1(w)} \le K e^{-\mathfrak{a}t} \|u_0 - \rho_\infty\|_{L^1(w)} \quad \forall t \ge 0.$$

For the model (1.1) considered in the present work, Assumption 3 will be necessary to show the Doeblin or Harris minoration conditions. We then have to distinguish two cases: either Assumption 2 holds, and we can use Harris's theorem since the Lyapunov condition (2.7) implies the operator Lyapunov condition (3.4) by Lemma 2.3 (the constants are then $A = e^{-\alpha T}$ and $B = \frac{b}{\alpha}$); or Assumption 4 holds, and we can simply use Doeblin's theorem. The main technical difficulty is to verify the minoration conditions, as the jumps of the process described by (2.5) are degenerate, and the model is two-dimensional. The rest of the section is devoted to the verification of the minoration condition.

LEMMA 3.3 (minoration condition). Grant Assumptions 1 and 3. Fix any $x \in \mathbb{R}$. For all R > 0, there exist T > 0 and a nonzero $\nu \in L^1_+$ such that

(3.7)
$$S_T^{\tilde{x}} u_0 \ge \nu \int_{\mathbb{R}_+ \times [0,R]} u_0 \, dadm \quad \forall u_0 \in L^1_+.$$

If, in addition, Assumption 4 holds, then there exist T > 0 and a nonzero $\nu \in L^{1}_{+}$ such that

(3.8)
$$S_T^{\tilde{x}} u_0 \ge \nu \|u_0\|_{L^1} \quad \forall u_0 \in L^1_+.$$

Proof. We proceed in two steps. First (Step 1), we choose a time T > 0 and a rectangle $[0, \overline{a}] \times [\underline{m}, \overline{m}] \subset \mathbb{R}_+ \times \mathbb{R}^*_+$ (with nonzero Lebesgue measure) and show that the density $S_T^{\tilde{x}} u_0 \in L^1$ has a lower bound on $[0, \bar{a}] \times [\underline{m}, \overline{m}]$ which depends on a Lebesgue integral in \mathbb{R}^2_+ involving u_0 . Then (Step 2), we perform a change of variable to express this lower bound in terms of $\int_{\mathbb{R}_+\times [0,R]} u_0 \, dadm$. The proof only relies on the expression of $S_t^{\tilde{x}} u_0$ given by the method of characteristics (see Proposition 2.2), and this allows treating a typically probabilistic question—the Doeblin/Harris minoration condition—from a transport point of view. This is possible because $S_t^{\tilde{x}}$ is the stochastic (mass-conservative) semigroup of a piecewise deterministic Markov process.

The constants Δ_{abs} , σ , and C_{γ} are taken from Assumption 3. **Step 1.** Fix R > 0. Since $\gamma(e^{-\lambda \Delta_{\text{abs}}}\gamma(0)) > \gamma(0)$ and $\gamma(e^{-\lambda t}\gamma(e^{-\lambda \Delta_{\text{abs}}}R)) \to \gamma(0)$ as $t \to \infty$, there exist $\bar{a} > 0$ and $T > \bar{a} + \Delta_{abs}$ such that

$$(3.9) \qquad \underline{m} =: \gamma(e^{-\lambda(T-\bar{a}-\Delta_{\rm abs})}\gamma(e^{-\lambda\Delta_{\rm abs}}R)) < e^{-\lambda\bar{a}}\gamma(e^{-\lambda\Delta_{\rm abs}}\gamma(0)) =: \overline{m}.$$

Equation (3.9) has the following heuristic interpretation: if we see $S_t^{\tilde{x}}$ as the stochastic semigroup of the piecewise deterministic Markov process defined by the generator

¹Note that Harris's theorem does not exclude the existence of an invariant probability measure with infinite $L^1(w)$ norm.

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(2.5), for any initial point $(a_0, m_0) \in \mathbb{R}_+ \times [0, R]$ and any landing point $(a, m) \in [0, \bar{a}] \times [\underline{m}, \overline{m}]$ at time T, there is a "possible" trajectory going from (a_0, m_0) to (a, m), with exactly two jumps (spikes). Since the trajectories of the process are determined by the jump times, we will exploit the fact that these "possible" trajectories correspond to jump times with strictly positive probability density. Below, we take a transport point of view on this probabilistic argument.

For all $(a, m) \in [0, \overline{a}] \times [\underline{m}, \overline{m}]$,

$$\begin{split} & (S_T^{\tilde{x}}u_0)(a,m) \\ & \geq \mathbbm{1}_{\{a < T\}}(S_{T-a}^{\tilde{x}}u_0)(0,e^{\lambda a}m)\exp\left(\lambda a - \int_{T-a}^T f(a-T+s,e^{\lambda(T-s)}m,\tilde{x})ds\right) \\ & \geq \mathbbm{1}_{\{a < T\}}e^{-\|f\|_{\infty}T}e^{\lambda a}(S_{T-a}^{\tilde{x}}u_0)(0,e^{\lambda a}m) \\ & \geq \mathbbm{1}_{\{a < T\}}e^{-\|f\|_{\infty}T}\sigma e^{\lambda a}\left|(\gamma^{-1})'(e^{\lambda a}m)\right|\int_{\Delta_{\mathrm{abs}}}^{\infty} (S_{T-a}^{\tilde{x}}u_0)(a',\gamma^{-1}(e^{\lambda a}m))da' \\ & = \mathbbm{1}_{\{a < T\}}e^{-\|f\|_{\infty}T}\sigma\left|\frac{d}{dm}\gamma^{-1}(e^{\lambda a}m)\right|\int_{\Delta_{\mathrm{abs}}}^{\infty} \underbrace{(S_{T-a}^{\tilde{x}}u_0)(a',\gamma^{-1}(e^{\lambda a}m))}_{(\star)}da'. \end{split}$$

Above, we went back in time to the last jump time T-a. Let us note that $\gamma^{-1}(e^{\lambda a}m) \ge \gamma^{-1}(e^{\lambda a}\underline{m}) > 0$. We can therefore define

$$a_{a,m}^* := \frac{1}{\lambda} \left(\log \gamma(0) - \log \gamma^{-1}(e^{\lambda a}m) \right).$$

Note that $a_{a,m}^*$ satisfies $\gamma^{-1}(e^{\lambda a_{a,m}^*}\gamma^{-1}(e^{\lambda a}m)) = 0$. In other words, $a_{a,m}^*$ is the minimal time between the last and second-to-last jumps for a trajectory landing at (a,m) at time T. We can easily verify that, by our choice of $\{T, \overline{a}, \underline{m}, \overline{m}\}, \Delta_{abs} \leq a_{a,m}^* < T - a - \Delta_{abs}$. This guarantees that it is possible to make two jumps in [0, T] and land at (a,m) at time T while respecting the absolute refractoriness of the neuron (i.e., there needs to be a time interval $\geq \Delta_{abs}$ between jumps). This allows us to go further back in time to the second-to-last jump as follows: for all $a' \in [a_{a,m}^*, T - a - \Delta_{abs}]$,

$$(\star) \geq \mathbb{1}_{\{a' < T-a\}} e^{-\|f\|_{\infty} T} \sigma \left| (\gamma^{-1})' (e^{\lambda a'} \gamma^{-1} (e^{\lambda a} m)) \right| e^{\lambda a'} \\ \times \int_{\Delta_{\mathrm{abs}}}^{\infty} \underbrace{(S_{T-a-a'}^{\tilde{x}} u_0)(a'', \gamma^{-1} (e^{\lambda a'} \gamma^{-1} (e^{\lambda a} m)))}_{(\star \star)} da''.$$

Then, we can go further back to time 0 to get u_0 :

$$(\star\star) \ge \mathbb{1}_{\{a'' \ge T - a - a'\}} e^{-\|f\|_{\infty} T} e^{\lambda(T - a - a')} \\ \times u_0(a'' - (T - a - a'), e^{\lambda(T - a - a')} \gamma^{-1}(e^{\lambda a'} \gamma^{-1}(e^{\lambda a} m))).$$

Putting all the lower bounds together, we get

$$(S_T^{\tilde{x}}u_0)(a,m) \ge \mathbb{1}_{\{a < T\}} e^{-3||f||_{\infty}T} \sigma^2 \int_{a_{a,m}^*}^{T-a-\Delta_{abs}} \int_{T-a-a'}^{\infty} \left| \frac{d}{dm} e^{\lambda(T-a-a')} \gamma^{-1}(e^{\lambda a'} \gamma^{-1}(e^{\lambda a}m)) \right| u_0(a'' - (T-a-a'), e^{\lambda(T-a-a')} \gamma^{-1}(e^{\lambda a'} \gamma^{-1}(e^{\lambda a}m))) da'' da'$$

Since $\gamma' \leq 1$ (Assumption 3),

$$\left|\frac{d}{dm}e^{\lambda(T-a-a')}\gamma^{-1}(e^{\lambda a'}\gamma^{-1}(e^{\lambda a}m))\right| \ge e^{\lambda T}.$$

Thus,

(3.10)
$$(S_T^{\tilde{x}}u_0)(a,m) \ge \mathbb{1}_{\{a < T\}} e^{(\lambda-3\|f\|_{\infty})T} \sigma^2 \times \int_{a_{a,m}^*}^{T-a-\Delta_{abs}} \int_0^\infty u_0(a_0, e^{\lambda(T-a-a')}\gamma^{-1}(e^{\lambda a'}\gamma^{-1}(e^{\lambda a}m))) da_0 da'.$$

We have obtained that on $[0, \bar{a}] \times [\underline{m}, \overline{m}]$, the density $(S_T^{\tilde{x}} u_0)$ is lower bounded by a constant depending on a Lebesgue integral on \mathbb{R}^2_+ involving u_0 .

Step 2. Now, we want to express the lower bound (3.10) in terms of

$$\int_{\mathbb{R}_+\times]0,R]} u_0 \, dadm$$

by a change of variable. Let us define the function $\psi_{a,m}^T$:

$$\psi_{a,m}^T : [a_{a,m}^*, T - a - \Delta_{\text{abs}}] \to \mathbb{R}_+, \quad a' \mapsto e^{\lambda(T - a - a')} \gamma^{-1}(e^{\lambda a'} \gamma^{-1}(e^{\lambda a}m)).$$

We verify that $(\psi_{a,m}^T)' > 0$ as follows: for all $a' \in [a_{a,m}^*, T-a]$,

$$(3.11)$$

$$(\psi_{a,m}^{T})'(a')$$

$$= \lambda e^{\lambda(T-a-a')} \bigg\{ (\gamma^{-1})' (e^{\lambda a'} \gamma^{-1}(e^{\lambda a}m)) e^{\lambda a'} \gamma^{-1}(e^{\lambda a}m) - \gamma^{-1}(e^{\lambda a'} \gamma^{-1}(e^{\lambda a}m)) \bigg\}.$$

As $\Gamma > 0$ and $\gamma' \leq 1$ (Assumption 3), we have

$$\begin{aligned} (\psi_{a,m}^T)'(a') &> \lambda e^{\lambda(T-a-a')} \left\{ (\gamma^{-1})'(e^{\lambda a'}\gamma^{-1}(e^{\lambda a}m))e^{\lambda a'}\gamma^{-1}(e^{\lambda a}m) - e^{\lambda a'}\gamma^{-1}(e^{\lambda a}m) \right\} \\ &= \lambda e^{\lambda(T-a)}\gamma^{-1}(e^{\lambda a}m) \left\{ \underbrace{(\gamma^{-1})'(e^{\lambda a'}\gamma^{-1}(e^{\lambda a}m))}_{\geq 1} - 1 \right\} \geq 0. \end{aligned}$$

Therefore, $\psi_{a,m}^T$ is a strictly increasing C^1 -diffeomorphism from $[a_{a,m}^*, T - a - \Delta_{abs}]$ to $[\psi_{a,m}^T(a_{a,m}^*), \psi_{a,m}^T(T - a - \Delta_{abs})]$. We can now rewrite (3.10) as

$$(S_T^{\tilde{x}}u_0)(a,m) \ge e^{(\lambda-3\|f\|_{\infty})T} \sigma^2 \int_{a_{a,m}^*}^{T-a-\Delta_{abs}} \int_0^\infty u_0(a_0,\psi_{a,m}^T(a')) da_0 da'$$

= $e^{(\lambda-3\|f\|_{\infty})T} \sigma^2 \int_{\psi_{a,m}^T(a_{a,m}^*)}^{\psi_{a,m}^T(T-a-\Delta_{abs})} \int_0^\infty u_0(a_0,m_0) \left| ((\psi_{a,m}^T)^{-1})'(m_0) \right| da_0 dm_0$

Going back to (3.11) and using the fact that there exists C_{γ} such that $C_{\gamma} \leq \gamma' \leq 1$ (Assumption 3), we have, for all $a' \in [a_{a,m}^*, T - a - \Delta_{abs}]$,

$$(\psi_{a,m}^T)'(a') \le \lambda e^{\lambda(T-a-a')} C_{\gamma}^{-1} e^{\lambda a'} \gamma^{-1}(e^{\lambda a}m) \le \lambda e^{\lambda T} C_{\gamma}^{-1} \overline{m}$$

Hence,

$$(S_T^{\tilde{x}}u_0)(a,m) \ge \frac{e^{-3\|f\|_{\infty}T}\sigma^2 C_{\gamma}}{\lambda\overline{m}} \int_{\psi_{a,m}^T(a_{a,m}^*)}^{\psi_{a,m}^T(T-a-\Delta_{\mathrm{abs}})} \int_0^{\infty} u_0(a_0,m_0) da_0 dm_0.$$

In addition, by our choice of $\{T, \overline{a}, \underline{m}, \overline{m}\}$, we have

$$\begin{split} \psi_{a,m}^{T}(a_{a,m}^{*}) &= 0, \\ \psi_{a,m}^{T}(T - a - \Delta_{\text{abs}}) &= e^{\lambda \Delta_{\text{abs}}} \gamma^{-1}(e^{\lambda (T - a - \Delta_{\text{abs}})} \gamma^{-1}(e^{\lambda a}m)) > R \end{split}$$

Therefore,

$$(S_T^{\tilde{x}}u_0)(a,m) \ge \frac{e^{-3\|f\|_{\infty}T}\sigma^2 C_{\gamma}}{\lambda \overline{m}} \int_0^R \int_0^{\infty} u_0(a_0,m_0) da_0 dm_0.$$

Since we have assumed that $(a, m) \in [0, \overline{a}] \times [\underline{m}, \overline{m}]$, this concludes the proof of (3.7).

In the case where Assumption 4 also holds, the proof of (3.8) is similar except that we can simply take $R = +\infty$ and $\underline{m} =: \gamma(e^{-\lambda(T-\bar{a}-\Delta_{abs})}G) < e^{-\lambda\bar{a}}\gamma(e^{-\lambda\Delta_{abs}}\gamma(0)) =: \overline{m}$.

In summary, by Harris's theorem, we have the following.

THEOREM 3.4. Grant Assumptions 1–3. For all $\tilde{x} \in \mathbb{R}$, there exists a unique $\rho_{\infty}^{\tilde{x}} \in L_{+}^{1}(w)$ with $\|\rho_{\infty}^{\tilde{x}}\|_{L^{1}} = 1$ such that $S_{t}^{\tilde{x}}\rho_{\infty}^{\tilde{x}} = \rho_{\infty}^{\tilde{x}}$ for all $t \geq 0$, and there exist $K \geq 1$ and $\mathfrak{a} > 0$ such that for all initial data $u_{0} \in L_{+}^{1}(w)$ with $\|u_{0}\|_{L^{1}} = 1$,

(3.12)
$$\left\| S_t^{\tilde{x}} u_0 - \rho_{\infty}^{\tilde{x}} \right\|_{L^1(w)} \le K e^{-\mathfrak{a}t} \left\| u_0 - \rho_{\infty}^{\tilde{x}} \right\|_{L^1(w)} \quad \forall t \ge 0.$$

Furthermore, by Lemma 2.3, we have that $\|\rho_{\infty}^{\tilde{x}}\|_{L^{1}(w)} \leq \frac{b}{\alpha}$, where the constants α and b are taken from the Lyapunov condition (2.7).

If Assumption 2 is replaced by Assumption 4, we can simply apply Doeblin's theorem.

THEOREM 3.5. Grant Assumptions 1, 3, and 4. For all $\tilde{x} \in \mathbb{R}$, there exists a unique $\rho_{\infty}^{\tilde{x}} \in L_{+}^{1}$ with $\|\rho_{\infty}^{\tilde{x}}\|_{L^{1}} = 1$ such that $S_{t}^{\tilde{x}}\rho_{\infty}^{\tilde{x}} = \rho_{\infty}^{\tilde{x}}$ for all $t \geq 0$, and there exist $K \geq 1$ and $\mathfrak{a} > 0$ such that for all initial data $u_{0} \in L_{+}^{1}$ with $\|u_{0}\|_{L^{1}} = 1$,

(3.13)
$$\left\|S_t^{\tilde{x}}u_0 - \rho_{\infty}^{\tilde{x}}\right\|_{L^1} \le Ke^{-\mathfrak{a}t} \left\|u_0 - \rho_{\infty}^{\tilde{x}}\right\|_{L^1} \quad \forall t \ge 0.$$

Note that both theorems imply the following.

COROLLARY 3.6. Grant Assumptions 1–3 (or Assumptions 1, 3, and 4). For all $\tilde{x} \in \mathbb{R}$, there exists a unique $\rho_{\infty}^{\tilde{x}} \in L^{1}_{+}(w)$ (or $\in L^{1}_{+}$) with $\|\rho_{\infty}^{\tilde{x}}\|_{L^{1}} = 1$ solving

(3.14a)
$$\partial_a \rho_{\infty}^{\tilde{x}}(a,m) - \lambda \partial_m(m\rho_{\infty}^{\tilde{x}}(a,m)) = -f(a,m,\tilde{x})\rho_{\infty}^{\tilde{x}}(a,m),$$

(3.14b)
$$\rho_{\infty}^{\tilde{x}}(0,m) = \mathbb{1}_{m > \gamma(0)} \left| (\gamma^{-1})'(m) \right| \int_{0}^{\infty} f(a,\gamma^{-1}(m),\tilde{x}) \rho_{\infty}^{\tilde{x}}(a,\gamma^{-1}(m)) da$$

in the weak sense. Furthermore, we have that $\rho_{\infty}^{\tilde{x}} \in \mathcal{C}(\mathbb{R}_+, L^1_+(\mathbb{R}^*_+)) \cap L^{\infty}(\mathbb{R}_+, L^1_+(\mathbb{R}^*_+)).$

4. Stationary solutions for arbitrary connectivity strength. In this section, we study the stationary solutions to (1.1), namely the solution to

(4.1a)
$$\partial_a \rho_\infty(a,m) - \lambda \partial_m(m\rho_\infty(a,m)) = -f(a,m,\varepsilon x_\infty)\rho_\infty(a,m),$$

(4.1b) $\rho_\infty(0,m) = \mathbbm{1}_{m>\gamma(0)} \left| (\gamma^{-1})'(m) \right| \int_0^\infty f(a,\gamma^{-1}(m),\varepsilon x_\infty)\rho_\infty(a,\gamma^{-1}(m))da,$
(4.1c) $x = \int_0^\infty \int_0^\infty \bar{h}(a,m)f(a,m,\varepsilon x_\infty)\rho_\infty(a,m)dadm$

(4.1c) $x_{\infty} = \int_{0} \int_{0} \bar{h}(a,m)f(a,m,\varepsilon x_{\infty})\rho_{\infty}(a,m)dadm.$

DEFINITION 4.1. $(\rho_{\infty}, x_{\infty}) \in L^1_+(w) \cap \mathcal{C}(\mathbb{R}_+, L^1_+(\mathbb{R}^*_+)) \cap L^{\infty}(\mathbb{R}_+, L^1_+(\mathbb{R}^*_+)) \times \mathbb{R}_+$ is a stationary solution to (1.1) if $\|\rho_{\infty}\|_{L^1} = 1$ and if it solves (4.1) in the weak sense.

4.1. Existence and uniqueness using the Doeblin–Harris method. We present two Lipschitz continuity results, which will allow us to prove the existence (and the uniqueness when ε is small) of stationary solutions. The following lemma plays the same role as Theorem 4.5 in [5].

LEMMA 4.2 (Lipschitz continuity at finite T). Grant Assumptions 1–3. For all initial data $u_0 \in L^1_+(w)$ and for all T > 0, there exists a constant $C_{T,||u_0||_{L^1(w)}} > 0$ such that

(4.2)
$$\forall \widetilde{x_1}, \widetilde{x_2} \in \mathbb{R}, \qquad \left\| S_T^{\widetilde{x_1}} u_0 - S_T^{\widetilde{x_2}} u_0 \right\|_{L^1(w)} \le C_{T, \|u_0\|_{L^1(w)}} |\widetilde{x_1} - \widetilde{x_2}|.$$

Proof. For all t > 0,

...

$$\begin{split} S_t^{\widetilde{x_1}} u_0 - S_t^{\widetilde{x_2}} u_0 \Big\|_{L^1(w)} \\ &= \int_0^\infty \int_t^\infty \Big| u_0(a-t, e^{\lambda t}m) \exp\left(\lambda t - \int_0^t f(a-t+s, e^{\lambda(t-s)}m, \widetilde{x_1}) ds\right) \\ &- u_0(a-t, e^{\lambda t}m) \exp\left(\lambda t - \int_0^t f(a-t+s, e^{\lambda(t-s)}m, \widetilde{x_2}) ds\right) \Big| w(a, m) dadm \\ &+ \int_0^\infty \int_0^t \Big| \rho_{t-a}^{\widetilde{x_1}}(0, e^{\lambda a}m) \exp\left(\lambda a - \int_{t-a}^t f(a-t+s, e^{\lambda(t-s)}m, \widetilde{x_1}) ds\right) \\ &- \rho_{t-a}^{\widetilde{x_2}}(0, e^{\lambda a}m) \exp\left(\lambda a - \int_{t-a}^t f(a-t+s, e^{\lambda(t-s)}m, \widetilde{x_2}) ds\right) \Big| w(a, m) dadm \\ &=: Q_1 + Q_2. \end{split}$$

ī

$$\begin{split} Q_1 &= \int_0^\infty \int_0^\infty u_0(a,m) \left| \exp\left(-\int_0^t f(a+s,e^{-\lambda s}m,\widetilde{x_1})ds\right) \right. \\ &- \exp\left(-\int_0^t f(a+s,e^{-\lambda s}m,\widetilde{x_2})ds\right) \right| w(a+t,e^{-\lambda t}m) dadm \\ &\leq \int_0^\infty \int_0^\infty u_0(a,m) \left(\int_0^t \left| f(a+s,e^{-\lambda s}m,\widetilde{x_1}) \right. \\ &\left. -f(a+s,e^{-\lambda s}m,\widetilde{x_2}) \right| ds\right) w(a+t,e^{-\lambda t}m) dadm \\ &\leq tL_f |\widetilde{x_1} - \widetilde{x_2}| \int_0^\infty \int_0^\infty u_0(a,m) w(a+t,e^{-\lambda t}m) dadm \\ &\leq tL_f ||u_0||_{L^1(w)} ||\widetilde{x_1} - \widetilde{x_2}|, \end{split}$$

where, in the last inequality, we used

(4.3)
$$w(a+t, e^{-\lambda t}m) \le w(a, m) \quad \forall a \ge 0, m \ge 0.$$

$$Q_{2} = \int_{0}^{\infty} \int_{0}^{t} \left| \rho_{t-a}^{\widetilde{x_{1}}}(0,m) \exp\left(-\int_{t-a}^{t} f(a-t+s,e^{\lambda(t-s-a)}m,\widetilde{x_{1}})ds\right) - \rho_{t-a}^{\widetilde{x_{2}}}(0,m) \exp\left(-\int_{t-a}^{t} f(a-t+s,e^{\lambda(t-s-a)}m,\widetilde{x_{2}})ds\right) \right| w(a,e^{-\lambda a}m) dadm.$$

By changes of variables,

$$\begin{split} Q_2 &= \int_0^\infty \int_0^t \left| \rho_s^{\widetilde{x_1}}(0,m) \exp\left(-\int_0^{t-s} f(u,e^{-\lambda u}m,\widetilde{x_1})du \right) \right. \\ &\quad - \rho_s^{\widetilde{x_2}}(0,m) \exp\left(-\int_0^{t-s} f(u,e^{-\lambda u}m,\widetilde{x_2})du \right) \left| w(t-s,e^{-\lambda(t-s)}m)dsdm \right. \\ &\leq \int_0^\infty \int_0^t \rho_s^{\widetilde{x_1}}(0,m) \left| \exp\left(-\int_0^{t-s} f(u,e^{-\lambda u}m,\widetilde{x_1})du \right) \right. \\ &\quad - \exp\left(-\int_0^{t-s} f(u,e^{-\lambda u}m,\widetilde{x_2})du \right) \left| w(t-s,e^{-\lambda(t-s)}m)dsdm \right. \\ &\quad + \int_0^\infty \int_0^t \left| \rho_s^{\widetilde{x_1}}(0,m) - \rho_s^{\widetilde{x_2}}(0,m) \right| w(t-s,e^{-\lambda(t-s)}m)dsdm \\ &\quad =: Q_{2,1} + Q_{2,2}. \end{split}$$

$$\begin{aligned} Q_{2,1} &\leq t \|f\|_{\infty} L_{f} |\widetilde{x_{1}} - \widetilde{x_{2}}| \int_{0}^{t} \int_{0}^{\infty} \int_{0}^{\infty} \left| (\gamma^{-1})'(m) \right| \rho_{s}^{\widetilde{x_{1}}}(a, \gamma^{-1}(m)) w(t, m) dadmds \\ &\leq t \|f\|_{\infty} L_{f} |\widetilde{x_{1}} - \widetilde{x_{2}}| \int_{0}^{t} \int_{0}^{\infty} \int_{0}^{\infty} \rho_{s}^{\widetilde{x_{1}}}(a, m) w(t, m + \|\Gamma\|_{\infty}) dadmds \\ &\leq t (1 + \|\Gamma\|_{\infty}) \|f\|_{\infty} L_{f} |\widetilde{x_{1}} - \widetilde{x_{2}}| \int_{0}^{t} \left\| \rho_{s}^{\widetilde{x_{1}}} \right\|_{L^{1}(w)} ds, \end{aligned}$$

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where, in the last inequality, we used

(4.4)
$$w(t, m + \|\Gamma\|_{\infty}) = 1 + m + \|\Gamma\|_{\infty} \le (1 + \|\Gamma\|_{\infty})w(a, m) \quad \forall a \ge 0, m \ge 0.$$

By Lemma 2.3,

$$Q_{2,1} \le t^2 (1 + \|\Gamma\|_{\infty}) \|f\|_{\infty} L_f \left(\|u_0\|_{L^1(w)} + \frac{b}{\alpha} \right) |\widetilde{x_1} - \widetilde{x_2}|$$

$$\begin{aligned} Q_{2,2} &\leq \|f\|_{\infty} \int_{0}^{t} \int_{0}^{\infty} \int_{0}^{\infty} \left| (\gamma^{-1})'(m) \right| \left| \rho_{s}^{\widetilde{x_{1}}}(a,\gamma^{-1}(m)) - \rho_{s}^{\widetilde{x_{2}}}(a,\gamma^{-1}(m)) \right| w(t,m) dadmds \\ &\leq \|f\|_{\infty} \int_{0}^{t} \int_{0}^{\infty} \int_{0}^{\infty} \left| \rho_{s}^{\widetilde{x_{1}}}(a,m) - \rho_{s}^{\widetilde{x_{2}}}(a,m) \right| w(t,m+\|\Gamma\|_{\infty}) dadmds \\ &\leq (1+\|\Gamma\|_{\infty}) \|f\|_{\infty} \int_{0}^{t} \left\| S_{s}^{\widetilde{x_{1}}}u_{0} - S_{s}^{\widetilde{x_{2}}}u_{0} \right\|_{L^{1}(w)} ds, \end{aligned}$$

where again, in the last inequality, we used (4.4). Fix T > 0. Gathering the bounds for Q_1 , $Q_{2,1}$, and $Q_{2,2}$ we see that there exist constants C > 0 and $C'_{T,||u_0||_{L^1(w)}} > 0$ such that, for all $t \in [0, T]$,

$$\left\|S_t^{\widetilde{x_1}}u_0 - S_t^{\widetilde{x_2}}u_0\right\|_{L^1(w)} \le C \int_0^t \left\|S_s^{\widetilde{x_1}}u_0 - S_s^{\widetilde{x_2}}u_0\right\|_{L^1(w)} ds + tC'_{T,\|u_0\|_{L^1(w)}}|\widetilde{x_1} - \widetilde{x_2}|.$$

By Grönwall's lemma, for all $t \in [0, T]$,

(4.5)
$$\left\| S_t^{\widetilde{x_1}} u_0 - S_t^{\widetilde{x_2}} u_0 \right\|_{L^1(w)} \le \frac{C'_{T, \|u_0\|_{L^1(w)}} |\widetilde{x_1} - \widetilde{x_2}|}{C} \left(\exp(Ct) - 1 \right).$$

Since (4.5) holds for all $t \in [0, T]$, this achieves the proof.

LEMMA 4.3 (Lipschitz continuity at $T = \infty$). Grant Assumptions 1–3. Writing $\rho_{\infty}^{\tilde{x}} \in L^{1}_{+}(w)$, the invariant probability measure given by Theorem 3.2 for any $\tilde{x} \in \mathbb{R}$, we see that the function

$$\Upsilon: \mathbb{R}_+ \to \mathbb{R}_+, \qquad \Upsilon(x) = \int_0^\infty \int_0^\infty \bar{h}(a,m) f(a,m,\varepsilon x) \rho_\infty^{\varepsilon x}(a,m) dadm$$

is Lipschitz, and there exists C > 0 such that

$$\forall x_1, x_2 \in \mathbb{R}_+, \qquad |\Upsilon(x_1) - \Upsilon(x_2)| \le |\varepsilon| C |x_1 - x_2|.$$

Proof. Since f is Lipschitz in x, we have, for any $x_1, x_2 \in \mathbb{R}_+$,

$$\begin{aligned} |\Upsilon(x_1) - \Upsilon(x_2)| &\leq \left\| \bar{h} \right\|_{\infty} \left\{ \|f\|_{\infty} \|\rho_{\infty}^{\varepsilon x_1} - \rho_{\infty}^{\varepsilon x_2}\|_{L^1} + L_f |\varepsilon| |x_1 - x_2| \right\} \\ &\leq \left\| \bar{h} \right\|_{\infty} \left\{ \|f\|_{\infty} \|\rho_{\infty}^{\varepsilon x_1} - \rho_{\infty}^{\varepsilon x_2}\|_{L^1(w)} + L_f |\varepsilon| |x_1 - x_2| \right\}, \end{aligned}$$

from where we only need to bound the first term on the right-hand side. We can use Theorem 3.2 and Lemma 4.2 as follows: for any $T \in \mathbb{R}_+$,

$$\begin{aligned} \|\rho_{\infty}^{\varepsilon x_{1}} - \rho_{\infty}^{\varepsilon x_{2}}\|_{L^{1}(w)} &= \|S_{T}^{\varepsilon x_{1}}\rho_{\infty}^{\varepsilon x_{1}} - S_{T}^{\varepsilon x_{1}}\rho_{\infty}^{\varepsilon x_{2}} + S_{T}^{\varepsilon x_{1}}\rho_{\infty}^{\varepsilon x_{2}} - S_{T}^{\varepsilon x_{2}}\rho_{\infty}^{\varepsilon x_{2}}\|_{L^{1}(w)} \\ &\leq Ke^{-\mathfrak{a}T}\|\rho_{\infty}^{\varepsilon x_{1}} - \rho_{\infty}^{\varepsilon x_{2}}\|_{L^{1}(w)} + C_{T}|\varepsilon||x_{1} - x_{2}|, \end{aligned}$$

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where K and a are the exponential stability constants of Theorem 3.2. Choosing T such that $Ke^{-\mathfrak{a}T} = 1/2$, we get

$$\|\rho_{\infty}^{\varepsilon x_1} - \rho_{\infty}^{\varepsilon x_2}\|_{L^1(w)} \le 2C_T |\varepsilon| |x_1 - x_2|.$$

Gathering the bounds concludes the proof.

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THEOREM 4.4 (stationary solutions). Grant Assumptions 1–3. We have the following:

- (i) There exists a stationary solution to (1.1).
- (ii) There exists $\varepsilon^* > 0$ such that for all $\varepsilon \in] \varepsilon^*, +\varepsilon^*[$, the stationary solution to (1.1) is unique.

Proof. For all $\tilde{x} \in \mathbb{R}$, let us write $\rho_{\infty}^{\tilde{x}} \in L_{+}^{1}(w)$ the unique invariant measure given by Theorem 3.2, and let us also take the function Υ from Lemma 4.3. By Corollary 3.6, $(\rho_{\infty}, x_{\infty}) \in L_{+}^{1}(w) \cap \mathcal{C}(\mathbb{R}_{+}, L_{+}^{1}(\mathbb{R}_{+}^{*})) \cap L^{\infty}(\mathbb{R}_{+}, L_{+}^{1}(\mathbb{R}_{+}^{*})) \times \mathbb{R}_{+}$ is a weak solution to (4.1) if and only if $\rho_{\infty} = \rho_{\infty}^{\varepsilon_{\infty}}$ and x_{∞} is a fixed-point of Υ . Hence, the study of the existence and uniqueness of stationary solutions is reduced to the study of the existence and uniqueness of the fixed-point of Υ .

Since for all $x \in \mathbb{R}_+$, $\|\rho_{\infty}^{\varepsilon x}\|_{L^1} = 1$, we have that for all $x \in \mathbb{R}_+$, $\Upsilon(x) \leq \|\bar{h}\|_{\infty} \|f\|_{\infty}$. Therefore, the set $[0, \|\bar{h}\|_{\infty} \|f\|_{\infty}]$ (which is compact and convex) is stable by Υ . Then, the continuity of Υ guarantees the existence of a fixed-point, which proves (i).

To obtain (ii), we observe that the Lipschitz constant of Υ is $|\varepsilon|C$: if we take $|\varepsilon| < \varepsilon^* := C^{-1}$, then Υ is a contraction, and we can apply Banach's fixed-point theorem to conclude the proof.

4.2. Alternative proof for the existence using Schauder's fixed-point theorem. We include here an alternative proof for the existence of a stationary solution, which is interesting for two reasons: on the one hand, it does not rely on the Doeblin–Harris method, and on the other hand, it provides some estimates on the stationary solutions.

For any $(\tilde{u}, \tilde{x}) \in L^1_+(]\gamma(0), +\infty[) \times \mathbb{R}$, consider the transport equation

$$\partial_a \varrho(a,m) - \lambda \partial_m (m \varrho(a,m)) = -f(a,m,\tilde{x})\varrho(a,m),$$

$$\rho(0,m) = \tilde{u}(m).$$

It has a unique weak solution $\rho_{\infty}^{\tilde{u},\tilde{x}} \in \mathcal{C}(\mathbb{R}_+, L^1_+(\mathbb{R}^*_+)) \cap L^{\infty}(\mathbb{R}_+, L^1_+(\mathbb{R}^*_+))$ given by the method of characteristics, i.e., for all $(a, m) \in \mathbb{R}_+ \times \mathbb{R}^*_+$,

(4.6)
$$\rho_{\infty}^{\tilde{u},\tilde{x}}(a,m) = \tilde{u}(e^{\lambda a}m) \exp\left(\lambda a - \int_{0}^{a} f(s,e^{\lambda(a-s)}m,\tilde{x})ds\right).$$

We can now define the operator $\Phi := (\Phi_1, \Phi_2)$ on $L^1_+(]\gamma(0), +\infty[) \times \mathbb{R}$, where, for all $(\tilde{u}, \tilde{x}) \in L^1_+(]\gamma(0), +\infty[) \times \mathbb{R}$,

(4.7a)
$$\Phi_{1}(\tilde{u}, \tilde{x})(m) := \mathbb{1}_{m > \gamma(0)} \left| (\gamma^{-1})'(m) \right| \int_{0}^{\infty} f(a, \gamma^{-1}(m), \tilde{x}) \rho_{\infty}^{(\tilde{u}, \tilde{x})}(a, \gamma^{-1}(m)) da,$$

(4.7b)
$$\Phi_{2}(\tilde{u}, \tilde{x}) := \int_{0}^{\infty} \int_{0}^{\infty} \bar{h}(a, m) f(a, m, \tilde{x}) \rho_{\infty}^{\tilde{u}, \tilde{x}}(a, m) da dm.$$

A stationary solution $(\rho_{\infty}, x_{\infty})$ is a fixed-point of Φ , and vice versa. Therefore, we have the following a priori estimates.

LEMMA 4.5. Grant Assumptions 1 and 3. There exists $\theta \in]0,1[$ such that for all $(\tilde{u},\tilde{x}) \in L^1_+(]\gamma(0),+\infty[) \times \mathbb{R}$, we have the following:

- (i) $\left\| \Phi_1(\tilde{u}, \tilde{x}) \right\|_{L^1} = \left\| \tilde{u} \right\|_{L^1}.$
- (ii) For all $m \in \mathbb{R}^+_+$, $|\Phi_1(\tilde{u}, \tilde{x})(m)| \leq \mathbb{1}_{m > \gamma(0)} \frac{\|f\|_{\infty}}{\lambda \gamma^{-1}(m)} \|\tilde{u}\|_{L^1}$.

(iii)

$$\int_0^\infty \Phi_1(\tilde{u}, \tilde{x})(m) m dm \le \max\left(\int_0^\infty \tilde{u}(m) m dm, \frac{\gamma(0)}{1-\theta} \|\tilde{u}\|_{L^1}\right)$$

(iv) For all $\beta \in \left]0, \frac{\min(f)}{\lambda}\right[$,

$$\int_{\gamma(0)}^{\infty} \frac{\Phi_1(\tilde{u}, \tilde{x})(m)}{\gamma^{-1}(m)^{\beta}} dm \le \frac{\|f\|_{\infty}}{\lambda\gamma(0)^{\beta}} \left(\frac{\min(f)}{\lambda} - \beta\right) \|\tilde{u}\|_{L^1}.$$

(v) $\Phi_2(\tilde{u}, \tilde{x}) \le \|\bar{h}\|_{\infty} \|\tilde{u}\|_{L^1}$.

Proof. (i) By changes of variables on m,

$$\begin{split} \left\| \Phi_1(\tilde{u},\tilde{x}) \right\|_{L^1} &= \int_0^\infty \int_0^\infty f(a,m,\tilde{x}) \tilde{u}(e^{\lambda a}m) \exp\left(\lambda a - \int_0^a f(s,e^{\lambda(a-s)}m,\tilde{x})ds\right) dadm \\ &= \int_0^\infty \tilde{u}(m) \underbrace{\int_0^\infty f(a,e^{-\lambda a}m,\tilde{x}) \exp\left(-\int_0^a f(s,e^{-\lambda s}m,\tilde{x})ds\right) da}_{=1 \quad \text{(by Assumption 3(i))}} dadm. \end{split}$$

(ii)

$$\begin{split} |\Phi_{1}(\tilde{u},\tilde{x})(m)| &\leq \mathbb{1}_{m > \gamma(0)} \|f\|_{\infty} \int_{0}^{\infty} \tilde{u}(e^{\lambda a} \gamma^{-1}(m)) \exp(\lambda a) da \\ &= \mathbb{1}_{m > \gamma(0)} \frac{\|f\|_{\infty}}{\lambda \gamma^{-1}(m)} \int_{0}^{\infty} \tilde{u}(e^{\lambda a} \gamma^{-1}(m)) \gamma^{-1}(m) \lambda \exp(\lambda a) da \\ &= \mathbb{1}_{m > \gamma(0)} \frac{\|f\|_{\infty}}{\lambda \gamma^{-1}(m)} \underbrace{\int_{\gamma^{-1}(m)}^{\infty} \tilde{u}(y) dy}_{\leq \|\tilde{u}\|_{L^{1}}}, \end{split}$$

where, for the last equality, we used the change of variable $y = e^{\lambda a} \gamma^{-1}(m)$.

(iii) Performing the same change of variable as for (i) and using the fact that $\gamma(m) \leq \gamma(0) + m$ for all $m \in \mathbb{R}_+$ (since $\gamma' \leq 1$), we have

$$\begin{split} &\int_{0}^{\infty} \Phi_{1}(\tilde{u},\tilde{x})(m)mdm \\ &= \int_{0}^{\infty} \tilde{u}(m) \int_{0}^{\infty} \gamma(e^{-\lambda a}m) f(a,e^{-\lambda a}m,\tilde{x}) \exp\left(-\int_{0}^{a} f(s,e^{-\lambda s}m,\tilde{x})ds\right) dadm \\ &\leq \int_{0}^{\infty} \tilde{u}(m)m \underbrace{\int_{0}^{\infty} e^{-\lambda a} f(a,e^{-\lambda a}m,\tilde{x}) \exp\left(-\int_{0}^{a} f(s,e^{-\lambda s}m,\tilde{x})ds\right) da}_{=:\vartheta(m)} dm + \gamma(0) \|\tilde{u}\|_{L^{1}} \end{split}$$

There exists $\theta \in]0,1[$ such that for all $m \in \mathbb{R}^*_+$, $\vartheta(m) < 1$, as we show in the

following. Fix $\epsilon > 0$.

Therefore,

$$\int_0^\infty \Phi_1(\tilde{u}, \tilde{x})(m) m dm \le \theta \int_0^\infty \tilde{u}(m) m dm + \gamma(0) \|\tilde{u}\|_{L^1}.$$

 $\leq 1 - (1 - e^{-\lambda \epsilon}) \exp(-\|f\|_{\infty} \epsilon) =: \theta < 1.$

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 $+\int_{\epsilon}^{\infty} e^{-\lambda\epsilon} f(a, e^{-\lambda a}m, \tilde{x}) \exp\left(-\int_{0}^{a} f(s, e^{-\lambda s}m, \tilde{x}) ds\right) da$

To see that

$$\int_0^\infty \Phi_1(\tilde{u}, \tilde{x})(m) m dm \le \max\left(\int_0^\infty \tilde{u}(m) m dm, \frac{\gamma(0)}{1-\theta} \|\tilde{u}\|_{L^1}\right),$$

we can distinguish three cases: if $\int_0^\infty \tilde{u}(m)mdm = \infty$, the inequality is trivial; if $\frac{\gamma(0)}{1-\theta} \|\tilde{u}\|_{L^1} \leq \int_0^\infty \tilde{u}(m)mdm < +\infty$, then

$$\begin{split} \int_0^\infty \Phi_1(\tilde{u}, \tilde{x})(m) m dm &\leq \int_0^\infty \tilde{u}(m) m dm - (1-\theta) \int_0^\infty \tilde{u}(m) m dm + \gamma(0) \|\tilde{u}\|_L \\ &\leq \int_0^\infty \tilde{u}(m) m dm; \end{split}$$

and, finally, if $\int_0^\infty \tilde{u}(m)mdm < \frac{\gamma(0)}{1-\theta} \|\tilde{u}\|_{L^1}$, then

$$\int_0^\infty \Phi_1(\tilde{u}, \tilde{x})(m) m dm \le \theta \frac{\gamma(0)}{1-\theta} \|\tilde{u}\|_{L^1} + \gamma(0) \|\tilde{u}\|_{L^1} = \frac{\gamma(0)}{1-\theta} \|\tilde{u}\|_{L^1} \,.$$

(iv)

$$\begin{split} &\int_{\gamma(0)}^{\infty} \frac{\Phi_1(\tilde{u},\tilde{x})(m)}{\gamma^{-1}(m)^{\beta}} dm \\ &= \int_0^{\infty} \int_0^{\infty} \frac{1}{m} f(a,m,\tilde{x}) \tilde{u}(e^{\lambda a}m) \exp\left(\lambda a - \int_0^a f(s,e^{\lambda(a-s)}m,\tilde{x})\right) dadm \\ &\leq \|f\|_{\infty} \int_0^{\infty} \int_0^{\infty} \frac{1}{m^{\beta}} \tilde{u}(e^{\lambda a}m) \exp\left(\lambda a - \min(f)a\right) dadm. \end{split}$$

Making the change of variable $y = e^{\lambda a} m$, we get

$$\begin{split} &= \|f\|_{\infty} \int_{0}^{\infty} \int_{m}^{\infty} \frac{1}{\lambda m^{1+\beta}} \tilde{u}(y) \exp\left(-\min(f) \frac{1}{\lambda} \ln\left(\frac{y}{m}\right)\right) dy dm \\ &= \frac{\|f\|_{\infty}}{\lambda} \int_{0}^{\infty} \int_{m}^{\infty} m^{\min(f)/\lambda - 1 - \beta} \tilde{u}(y) y^{-\min(f)/\lambda} dy dm, \end{split}$$

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and, using Fubini's theorem and the fact that $\min(f)/\lambda - \beta > 0$, we get

$$\begin{split} &= \frac{\|f\|_{\infty}}{\lambda} \int_0^{\infty} \tilde{u}(y) y^{-\min(f)/\lambda} \underbrace{\int_0^y m^{\min(f)/\lambda - 1 - \beta} dm}_{=\frac{y^{\min(f)/\lambda - \beta}}{\min(f)/\lambda - \beta}} dy \\ &= \frac{\|f\|_{\infty}}{\lambda} \left(\frac{\min(f)}{\lambda} - \beta\right) \int_0^{\infty} \tilde{u}(y) y^{-\beta} dy. \end{split}$$

Finally, it is easy to check that $\int_0^\infty \tilde{u}(y)y^{-\beta}dy \leq \gamma(0)^{-\beta} \|\tilde{u}\|_{L^1}$.

(v) Use (4.7b) and see the proof of (i).

By these estimates, we see that there exists $\beta, C_1, \ldots, C_4 > 0$ such that the set $\mathscr{C} \times B \subset L^1(]\gamma(0), +\infty[) \times \mathbb{R}$, where

$$\begin{split} \mathscr{C} &:= \bigg\{ u \in L^1_+(]\gamma(0), +\infty[) \ \bigg| \ \|u\|_{L^1} \le 1; \\ u &\le \frac{C_1}{\gamma^{-1}(\cdot)} \ \text{a.e.}; \int_0^\infty u(m) m dm \le C_2; \ \int_{\gamma(0)}^\infty \frac{u(m)}{\gamma^{-1}(m)^\beta} dm \le C_3 \bigg\}, \end{split}$$

and $B := [-C_4, +C_4]$ is stable by the operator Φ .

In order to apply Schauder's fixed-point theorem, we will need the next lemma.

LEMMA 4.6. Grant Assumptions 1 and 3. C is convex, closed, and compact for the weak topology $\sigma(L^1, L^\infty)$.

Proof. It is easy to verify that \mathscr{C} is convex. Since \mathscr{C} is convex, it suffices to show that it is strongly closed to show that it is weakly closed. Let u_n be a sequence of elements of \mathscr{C} which converge strongly to $u \in L^1(]\gamma(0), +\infty[)$. By the strong convergence, $||u||_{L^1} \leq 1$. We can extract a subsequence u_{n_k} such that u_{n_k} converges to u a.e. Taking the pointwise limit, we have that $u \leq \frac{C_1}{\gamma^{-1}(\cdot)}$ a.e. Furthermore, by Fatou lemma,

$$\int_{\gamma(0)}^{\infty} u(m)mdm \le \liminf_{k \to +\infty} \int_{\gamma(0)}^{\infty} u_{n_k}(m)mdm \le C_2$$

and

$$\int_{\gamma(0)}^{\infty} \frac{u(m)}{\gamma^{-1}(m)^{\beta}} dm \le \liminf_{k \to +\infty} \int_{\gamma(0)}^{\infty} \frac{u_{n_k}(m)}{\gamma^{-1}(m)^{\beta}} dm \le C_3$$

Hence, \mathscr{C} is strongly closed.

To show that \mathscr{C} is weakly compact, we will show the following:

- (a) $\sup_{u \in \mathscr{C}} ||u||_{L^1} < \infty.$
- (b) For all $\epsilon > 0$, there exists R > 0 such that $\int_{R}^{\infty} u(m) dm < \epsilon$ for all $u \in \mathscr{C}$.
- (c) \mathscr{C} is equi-integrable, i.e., for all $\epsilon > 0$, there exists $\delta > 0$ such that for all Borel set $A \subset \mathbb{R}_+$ with $|A| \leq \delta$ and for all $u \in \mathscr{C}$, $\int_A u(m) dm \leq \epsilon$.

Then use the Dunford–Pettis theorem. (a) is clearly verified. (b) is also verified since for all R > 0, $\int_{R}^{\infty} u(m)dm \leq \frac{1}{R} \int_{0}^{\infty} u(m)mdm \leq \frac{C_2}{R}$. To show (c), let us first observe that for all $\delta_1 > 0$,

$$\int_{\gamma(0)}^{\gamma(0)+\delta_1} u(m)dm \le \gamma^{-1}(\gamma(0)+\delta_1)^{\beta} \int_{\gamma(0)}^{\infty} \frac{u(m)}{\gamma^{-1}(m)^{\beta}}dm \le \gamma^{-1}(\gamma(0)+\delta_1)^{\beta}C_3.$$

For any $\epsilon > 0$, let us choose $\delta_1 > 0$ such that $\gamma^{-1}(\gamma(0) + \delta_1)^{\beta}C_3 \leq \frac{\epsilon}{2}$. Then, for all Borel set $A \subset \mathbb{R}_+$ with $|A| \leq \delta$,

$$\int_{A} u(m) dm \leq \int_{\gamma(0)}^{\gamma(0)+\delta_{1}} u(m) dm + \int_{A \setminus [0,\gamma(0)+\delta_{1}]} u(m) dm \leq \frac{\epsilon}{2} + \delta \frac{C_{1}}{\gamma^{-1}(\gamma(0)+\delta_{1})}.$$

Hence, we can choose $\delta = \min(\delta_1, \frac{\epsilon \gamma^{-1}(\gamma(0) + \delta_1)}{2C_1})$, and (c) is verified. By the Dunford–Pettis theorem, \mathscr{C} is weakly relatively compact. Finally, since \mathscr{C} is weakly closed, \mathscr{C} is weakly compact.

We can now give an alternative proof of the existence of stationary solutions to (1.1) for arbitrary connectivity strength ε .

Proof of Theorem 1.3(i). We verify that the operator Φ is weakly continuous: for any sequence $(u_n, x_n) \to (u, x)$ in $\mathscr{C} \times \mathbb{R}$ and for any $\varphi \in L^{\infty}(\mathbb{R}_+)$,

$$\left| \int (\Phi_1(u_n, x_n) - \Phi_1(u, x)) \varphi(m) dm \right| \le Q_1^n + Q_2^n + Q_3^n,$$

where

$$\begin{split} Q_1^n &:= \left| \int_0^\infty \int_0^\infty (u_n(e^{\lambda a}m) - u(e^{\lambda a}m))\varphi(\gamma(m))e^{\lambda a}f(a,m,x)e^{-\int_0^a f(\tau,e^{\lambda(a-\tau)}m,x)d\tau}dadm \right|,\\ Q_2^n &:= \|\varphi\|_\infty \int_0^\infty \int_0^\infty u_n(e^{\lambda a}m)e^{\lambda a}|f(a,m,x) - f(a,m,x_n)|e^{-\int_0^a f(\tau,e^{\lambda(a-\tau)}m,x)d\tau}dadm,\\ Q_3^n &:= \|\varphi\|_\infty \int_0^\infty \int_0^\infty u_n(e^{\lambda a}m)e^{\lambda a}f(a,m,x_n) \left| e^{-\int_0^a f(\tau,e^{\lambda(a-\tau)}m,x)d\tau} - e^{-\int_0^a f(\tau,e^{\lambda(a-\tau)}m,x_n)d\tau} \right| dadm. \end{split}$$

Making the change of variable $ydy = e^{\lambda a}mdm$ in Q_1 , we get

$$Q_1^n = \left| \int_0^\infty (u_n(y) - u(y)) \int_0^\infty \varphi(\gamma(ye^{-\lambda a})) f(a, ye^{-\lambda a}, x) e^{-\int_0^a f(\tau, e^{-\lambda \tau}y, x)} dadm \right|.$$

Since u_n converges to u in $\sigma(L^1, L^\infty)$ and

$$\begin{split} \int_0^\infty \varphi(\gamma(ye^{-\lambda a})) f(a, ye^{-\lambda a}, x) e^{-\int_0^a f(\tau, e^{-\lambda \tau} y, x)d\tau} da \\ \leq \|\varphi\|_\infty \int_0^\infty f(a, ye^{-\lambda a}, x) e^{-\int_0^a f(\tau, e^{-\lambda \tau} y, x)d\tau} da = \|\varphi\|_\infty \,, \end{split}$$

 Q_1^n converges to 0. On the other hand, since f is bounded and Lipschitz, $Q_2^n, Q_3^n \leq ||u_n||_{L^1}C|x_n-x| \leq C|x_n-x|$. Therefore, Φ_1 is a continuous operator with respect to the weak topology $\sigma(L^1, L^\infty)$.

The continuity of Φ_2 is shown analogously, taking $\varphi = h$ (*h* is bounded).

Since \mathscr{C} is stable by Φ , convex, and weakly compact (Lemma 4.6), we can apply Schauder's fixed-point theorem to obtain the existence of a fixed-point, which gives the existence of a stationary solution.

COROLLARY 4.7. Grant Assumptions 1 and 3. If f is of class C^k , then u(m) is a function of class C^k for all $m > \gamma(0)$. Consequently, the stationary solutions of (1.1) are of class C^k .

Proof. If (u, \tilde{x}) is a fixed-point of Φ , then

(4.8)
$$u(m) = \mathbb{1}_{m > \gamma(0)} \left| (\gamma^{-1})'(m) \right| \int_0^\infty f(a, \gamma^{-1}(m), \tilde{x}) u(e^{\lambda a} \gamma^{-1}(m)) \times \exp\left(\lambda a - \int_0^a f(s, e^{\lambda(a-s)} \gamma^{-1}(m), \tilde{x}) ds\right) da.$$

Making the change of variable $y = e^{\lambda a} \gamma^{-1}(m)$ in a, as in estimate (ii) of Lemma 4.5, we obtain

(4.9)
$$u(m) = \mathbb{1}_{m > \gamma(0)} \frac{\left| (\gamma^{-1})'(m) \right|}{\lambda \gamma^{-1}(m)} \int_{\gamma^{-1}(m)}^{\infty} f(g(y,m), y, \tilde{x}) u(y)$$
$$\times \exp\left(-\int_{0}^{g(y,m)} f(s, e^{s}y, \tilde{x}) ds \right) dy,$$

where $g(y,m) = \ln \frac{y}{\lambda(\gamma^{-1}(m))}$. We conclude with a bootstrap argument: if u is L^1 , then the right-hand side of (4.9) is a continuous function of m, meaning that u is continuous. But if u is continuous, then the right-hand side is of class C^1 , etc.

4.3. Formula in the case of short-term synaptic depression. In general, there is no explicit formula for the invariant probability measure solving (3.14). However, in the case of short-term synaptic depression (1.3), we can derive an explicit expression for the total postsynaptic potential

$$X(\tilde{x}) := \int_0^\infty \hat{h}(t) \int_0^1 \int_0^\infty (1-m) f(a,\tilde{x}) \rho_\infty^{\tilde{x}}(a,m) dadm dt$$

for any $\tilde{x} \in \mathbb{R}$. This fact has been reported in the theoretical neuroscience literature [42]; we provide here a rigorous and analytic justification for it.

For all $\tilde{x} \in \mathbb{R}$, let us introduce the quantities

$$I^{\tilde{x}} := \int_0^\infty a f(a, \tilde{x}) \exp\left(-\int_0^a f(s, \tilde{x}) ds\right) da = \int_0^\infty \exp\left(-\int_0^a f(s, \tilde{x}) ds\right) da,$$
$$P^{\tilde{x}}(\lambda) := \int_0^\infty e^{-\lambda a} f(a, \tilde{x}) \exp\left(-\int_0^a f(s, \tilde{x}) ds\right) da.$$

The value $I^{\tilde{x}}$ can be interpreted as the mean interspike interval of a neuron receiving a constant input \tilde{x} . The value $P^{\tilde{x}}(\lambda)$ can be seen as the Laplace transform of the interspike interval distribution of that neuron evaluated in λ .

PROPOSITION 4.8. Grant Assumptions 1 and 3. For all $\tilde{x} \in \mathbb{R}$,

$$X(\tilde{x}) = \int_0^\infty \hat{h}(t) dt \, \frac{1}{I^{\tilde{x}}} \left\{ \frac{1 - P^{\tilde{x}}(\lambda)}{1 - v P^{\tilde{x}}(\lambda)} \right\}.$$

Proof. Using the method of characteristics (i.e., combining (4.6) and (3.14b)), we

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have

$$\begin{split} 1 &= \int_0^1 \int_0^\infty \rho_\infty^{\tilde{x}}(a,m) dadm \\ &= \int_0^1 \int_0^\infty \mathbbm{1}_{e^{\lambda a}m < 1} \rho_\infty^{\tilde{x}}(0, e^{\lambda a}m) \exp\left(\lambda a - \int_0^a f(s, \tilde{x}) ds\right) dadm \\ &= \int_0^1 \int_0^\infty \rho_\infty^{\tilde{x}}(0,m) \exp\left(-\int_0^a f(s, \tilde{x}) ds\right) dadm = I^{\tilde{x}} \int_0^1 \rho_\infty^{\tilde{x}}(0,m) dm \end{split}$$

Therefore,

$$\int_{0}^{1} \int_{0}^{\infty} f(a, \tilde{x}) \rho_{\infty}^{\tilde{x}}(a, m) dadm = \int_{0}^{1} \rho_{\infty}^{\tilde{x}}(0, m) dm = \frac{1}{I^{\tilde{x}}}.$$

On the other hand,

$$\begin{split} \int_0^1 \int_0^\infty mf(a,\tilde{x})\rho_\infty^{\tilde{x}}(a,m)dadm \\ &= \int_0^1 \int_0^\infty \mathbbm{1}_{e^{\lambda a}m < 1} mf(a,\tilde{x})\rho_\infty^{\tilde{x}}(0,e^{\lambda a}m) \exp\left(\lambda a - \int_0^a f(s,\tilde{x})ds\right) dadm \\ &= \int_0^1 \int_0^\infty e^{-\lambda a} mf(a,\tilde{x})\rho_\infty^{\tilde{x}}(0,m) \exp\left(-\int_0^a f(s,\tilde{x})ds\right) dadm \\ &= P^{\tilde{x}}(\lambda) \int_0^1 m\rho_\infty^{\tilde{x}}(0,m) dm \end{split}$$

and

$$\begin{split} \int_0^1 m\rho_\infty^{\tilde{x}}(0,m)dm &= \int_0^1 m \mathbbm{1}_{m>1-\upsilon} \frac{1}{\upsilon} \int_0^\infty f(a,\tilde{x})\rho_\infty^{\tilde{x}} \left(a,1-\frac{1-m}{\upsilon}\right) dadm \\ &= \int_0^1 (1-\upsilon+\upsilon m) \int_0^\infty f(a,\tilde{x})\rho_\infty^{\tilde{x}}(a,m) dadm \\ &= \frac{1-\upsilon}{I^{\tilde{x}}} + \upsilon P^{\tilde{x}}(\lambda) \int_0^1 m\rho_\infty^{\tilde{x}}(0,m) dm. \end{split}$$

Therefore,

$$\int_0^1 m \rho_\infty^{\tilde{x}}(0,m) dm = \frac{1-\upsilon}{I^{\tilde{x}}(1-\upsilon P^{\tilde{x}}(\lambda))}$$

 $\quad \text{and} \quad$

$$\int_0^1 \int_0^\infty mf(a,\tilde{x})\rho_\infty^{\tilde{x}}(a,m)dadm = \frac{P^{\tilde{x}}(\lambda)(1-\upsilon)}{I^{\tilde{x}}(1-\upsilon P^{\tilde{x}}(\lambda))}.$$

Finally, we have

$$\begin{split} X(\tilde{x}) &= \int_{0}^{\infty} \hat{h}(t) dt \left\{ \int_{0}^{1} \int_{0}^{\infty} f(a, \tilde{x}) \rho_{\infty}^{\tilde{x}}(a, m) da dm - \int_{0}^{1} \int_{0}^{\infty} m f(a, \tilde{x}) \rho_{\infty}^{\tilde{x}}(a, m) da dm \right\} \\ &= \int_{0}^{\infty} \hat{h}(t) dt \; \frac{1}{I^{\tilde{x}}} \left\{ 1 - \frac{P^{\tilde{x}}(\lambda)(1-v)}{(1-vP^{\tilde{x}}(\lambda))} \right\} = \int_{0}^{\infty} \hat{h}(t) dt \; \frac{1}{I^{\tilde{x}}} \left\{ \frac{1-P^{\tilde{x}}(\lambda)}{1-vP^{\tilde{x}}(\lambda)} \right\}. \quad \Box$$

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5. Exponential stability in the weak connectivity regime. To study the long time behavior (1.1) in the weak connectivity regime, we perturb the noninteracting case (3.1), taking $\tilde{x} = \varepsilon x_{\infty}$, where x_{∞} is given by the unique stationary solution to (1.1) when $\varepsilon \in] -\varepsilon^*, +\varepsilon^* [(\varepsilon^* \text{ is taken from Theorem 1.3(ii)}).$ In this section, we keep the small ε fixed, and we work under Assumptions 1–3 and 5. We roughly follow the same line of argument as in [30, sect. 5].

For convenience, we first rewrite (1.1) in a more formal and compact form,

(5.1a)
$$\partial_t \rho_t = -\partial_a \rho_t + \lambda \partial_m (m\rho_t) - f(\varepsilon x_t) \rho_t + \delta_0^a (\gamma \circ \Pi)_* (f(\varepsilon x_t) \rho_t),$$

(5.1b)
$$x_t = \int_0^t \int h(t-s) f(\varepsilon x_s) \rho_s \, dadm \, ds,$$

(5.1c) $\rho_0 = u_0,$

where δ_0^a indicates that (singular) mass enters in $a = 0,^2 \Pi : (a, m) \mapsto m$ is the projection on m, and * denotes the pushforward measure. To write (5.1) as an evolution equation, we introduce an auxiliary transport equation on $\mathbb{R}_+ \times \mathbb{R}_+ \times \mathbb{R}_+^*$,

$$\partial_t \zeta_t = -\partial_s \zeta_t + \delta_0^s f(\varepsilon x_t) \rho_t,$$

$$\zeta_0 = 0.$$

which solution is given by the method of characteristics,

$$\zeta_t(s) = \mathbb{1}_{s < t} f(\varepsilon x_{t-s}) \rho_{t-s} \qquad \forall (t,s) \in \mathbb{R}^*_+ \times \mathbb{R}_+.$$

Using the auxiliary equation, (5.1) is equivalent to

(5.2a)

$$\begin{aligned} \partial_t(\rho_t, \zeta_t) &= \left(-\partial_a \rho_t + \lambda \partial_m(m\rho_t) - f(\varepsilon x_t) \rho_t + \delta_0^a (\gamma \circ \Pi)_* \left(f(\varepsilon x_t) \rho_t \right), -\partial_s \zeta_t + \delta_0^s f(\varepsilon x_t) \rho_t \right), \\ \end{aligned}$$
(5.2b)
$$(\rho_0, \zeta_0) &= (u_0, 0), \end{aligned}$$

where $x_t := \int_0^\infty \int h(s)\zeta_t(s) \, dadm \, ds$. By Theorem 1.3, for all $\varepsilon \in] - \varepsilon^*, +\varepsilon^*[$, there exists a unique stationary solution $(\rho_{\infty}, x_{\infty})$, and we have

(5.3)
$$-\partial_a \rho_\infty + \lambda \partial_m (m\rho_\infty) - f(\varepsilon x_\infty) \rho_\infty + \delta_0^a (\gamma \circ \Pi)_* \left(f(\varepsilon x_\infty) \rho_\infty \right) = 0.$$

Now, we write (5.2) as the sum of a linear equation and a perturbation,

(5.4a)
$$\partial_t(\rho_t, \zeta_t) = \Lambda(\rho_t, \zeta_t) + (Z_t^{(1)}, Z_t^{(2)})$$

(5.4b)
$$(\rho_0, \zeta_0) = (u_0, 0)$$

$$\partial_t \rho_t = -\partial_a \rho_t + \lambda \partial_m (m\rho_t) - f(\varepsilon x_t) \rho_t + \delta_{0=a} \left\{ (\gamma \circ \Pi)_* \left(f(\varepsilon x_t) \rho_t \right) - \rho_t(0, \cdot) \right\}.$$

 $^{{}^{2}\}delta_{0}^{a}$ should not be confused with the Dirac distribution $\delta_{0=a}$. Using $\delta_{0=a}$, by integration by parts of weak solutions, (5.1a) should read

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where

$$\begin{split} \Lambda(\rho_t, \zeta_t) \\ &:= \left(-\partial_a \rho_t + \lambda \partial_m(m\rho_t) - f(\varepsilon x_\infty)\rho_t + \delta_0^a(\gamma \circ \Pi)_* \left(f(\varepsilon x_\infty)\rho_t \right), -\partial_s \zeta_t + \delta_0^s f(\varepsilon x_\infty)\rho_t \right), \\ Z_t^{(1)} &:= \left[f(\varepsilon x_\infty) - f(\varepsilon x_t) \right] \rho_t + \delta_0^a(\gamma \circ \Pi)_* (\left[f(\varepsilon x_t) - f(\varepsilon x_\infty) \right] \rho_t), \\ Z_t^{(2)} &:= \delta_0^s [f(\varepsilon x_t) - f(\varepsilon x_\infty)] \rho_t. \end{split}$$

Let us put $\zeta_{\infty}(s) := f(\varepsilon x_{\infty})\rho_{\infty}$ for all $s \in \mathbb{R}_+$. Then, using (5.3), by the linearity of the operator Λ and writing $\bar{\rho}_t := \rho_t - \rho_{\infty}$ and $\bar{\zeta}_t := \zeta_t - \zeta_{\infty}$, we get

(5.5a)
$$\partial_t(\bar{\rho}_t, \bar{\zeta}_t) = \Lambda(\bar{\rho}_t, \bar{\zeta}_t) + (Z_t^{(1)}, Z_t^{(2)}),$$

(5.5b)
$$(\bar{\rho}_0, \zeta_0) = (u_0 - \rho_\infty, -\zeta_\infty)$$

Writing $(S_t^{\Lambda})_{t \in \mathbb{R}_+}$ the semigroup associated with the operator Λ , we have, by Duhamel's formula,

(5.6)
$$(\bar{\rho}_t, \bar{\zeta}_t) = S_t^{\Lambda}(\bar{\rho}_0, \bar{\zeta}_0) + \int_0^t S_{t-s}^{\Lambda}(Z_s^{(1)}, Z_s^{(2)}) ds \quad \forall t \ge 0.$$

Let us define the weighted space

$$L^1_+(\mu) := \left\{ \zeta \in L^1(\mathbb{R}_+ \times \mathbb{R}_+ \times \mathbb{R}_+^*, \mathbb{R}_+) \ \Big| \ \int_0^\infty \left\| \zeta(s) \right\|_{L^1} \|h\|_\infty \ e^{-\mathfrak{h} s} ds < \infty \right\}.$$

Note that, for all $t \ge 0$,

$$\begin{aligned} |x_t - x_{\infty}| &= \left| \int_0^{\infty} \int h(s)\zeta_t(s) \, dadm \, ds - \int_0^{\infty} \int h(s)\zeta_{\infty}(s) \, dadm \, ds \right| \\ &\leq \int_0^{\infty} \|h\|_{\infty} \, e^{-\mathfrak{h}s} \left\| \zeta_t(s) - \zeta_{\infty}(s) \right\|_{L^1} ds = \left\| \bar{\zeta}_t \right\|_{L^1(\mu)}. \end{aligned}$$

Also, we have, for all $t \ge 0$,

(5.7a)
$$\begin{aligned} \left\| Z_{t}^{(1)} \right\|_{L^{1}} &\leq |\varepsilon| 2L_{f} \|\rho_{t}\|_{L^{1}} |x_{t} - x_{\infty}| \leq |\varepsilon| 2L_{f} \|\bar{\zeta}_{t}\|_{L^{1}(\mu)}, \\ (5.7b) \quad \left\| Z_{t}^{(1)} \right\|_{L^{1}(w)} &\leq |\varepsilon| 2L_{f} \|\rho_{t}\|_{L^{1}(w)} |x_{t} - x_{\infty}| \leq |\varepsilon| 2L_{f} \left(\|u_{0}\|_{L^{1}(w) + \frac{b}{\alpha}} \right) \|\bar{\zeta}_{t}\|_{L^{1}(\mu)}, \\ (5.7c) \quad \left\| Z_{t}^{(2)} \right\|_{L^{1}(\mu)} &\leq |\varepsilon| \|h\|_{\infty} L_{f} \|\rho_{t}\|_{L^{1}} |x_{t} - x_{\infty}| \leq |\varepsilon| \|h\|_{\infty} L_{f} \|\bar{\zeta}_{t}\|_{L^{1}(\mu)}, \end{aligned}$$

where we have used Theorem 1.2(ii) in the first line and Lemma 2.1 in the second.

LEMMA 5.1. Grant 1–3 and 5 and take $(\bar{\rho}_0, \bar{\zeta}_0)$ as in (5.5). There exist $K_1 \geq 1$ and $\mathfrak{a}_1 > 0$ such that, for all initial data $u_0 \in L^1_+(w)$ with $\|u_0\|_{L^1} = 1$,

(5.8)
$$\left\| S_t^{\Lambda}(\bar{\rho}_0, \bar{\zeta}_0) \right\|_{L^1(w) \times L^1(\mu)} \le K_1 e^{-\mathfrak{a}_1 t} \left\| (\bar{\rho}_0, \bar{\zeta}_0) \right\|_{L^1(w) \times L^1(\mu)} \quad \forall t \ge 0.$$

If in addition we grant Assumption 4, then there exist $K_2 \ge 1$ and $\mathfrak{a}_2 > 0$ such that, for all initial data $u_0 \in L^1_+$ with $||u_0||_{L^1} = 1$,

(5.9)
$$\left\| S_t^{\Lambda}(\bar{\rho}_0, \bar{\zeta}_0) \right\|_{L^1 \times L^1(\mu)} \le K_2 e^{-\mathfrak{a}_2 t} \left\| (\bar{\rho}_0, \bar{\zeta}_0) \right\|_{L^1 \times L^1(\mu)} \quad \forall t \ge 0.$$

Proof. We write $(S_t^{\Lambda}(\bar{\rho}_0, \bar{\zeta}_0)^{(1)}, S_t^{\Lambda}(\bar{\rho}_0, \bar{\zeta}_0)^{(2)}) := S_t^{\Lambda}(\bar{\rho}_0, \bar{\zeta}_0)$ the first and second components of $S_t^{\Lambda}(\bar{\rho}_0, \bar{\zeta}_0)$.

By Theorem 3.2, there exist $K \ge 0$ and $\mathfrak{a} > 0$ such that

$$\left\| S_t^{\Lambda}(\bar{\rho}_0, \bar{\zeta}_0)^{(1)} \right\|_{L^1(w)} \le K e^{-\mathfrak{a}t} \|\bar{\rho}_0\|_{L^1(w)} \qquad \forall t \ge 0$$

Then,

$$\begin{split} \left\| S_{t}^{\Lambda}(\bar{\rho}_{0},\bar{\zeta}_{0})^{(2)} \right\|_{L^{1}(\mu)} &= \int_{0}^{t} \left\| f(\varepsilon x_{\infty}) S_{t-s}^{\Lambda}(\bar{\rho}_{0},\bar{\zeta}_{0})^{(1)} \right\|_{L^{1}} C_{h} e^{-\mathfrak{h}s} ds \\ &+ \int_{t}^{\infty} \left\| \bar{\zeta}_{0}(s) \right\|_{L^{1}} C_{h} e^{-\mathfrak{h}s} ds \\ &\leq C_{h} \left\{ \| f \|_{\infty} K \int_{0}^{t} e^{-\mathfrak{a}(t-s)} e^{-\mathfrak{h}s} ds \| \bar{\rho}_{0} \|_{L^{1}(w)} + e^{-\mathfrak{h}t} \| \bar{\zeta}_{0} \|_{L^{1}(\mu)} \right\} \end{split}$$

Gathering the bounds on the two components and observing that the function $t \mapsto \int_0^t e^{-\mathfrak{a}(t-s)}e^{-\mathfrak{h}s}ds$ decays exponentially, we conclude that there exist $K_1 \geq 1$ and $\mathfrak{a}_1 > 0$ such that (5.8) holds.

For (5.9), we use Theorem 3.1 and follow the same argument. \Box

We can now prove our main result.

Proof of Theorem 1.4. By Duhamel's formula (5.6), (5.8) in Lemma 5.1, and the bounds (5.7), for all $t \ge 0$,

$$\begin{split} &\|(\bar{\rho}_{t},\bar{\zeta}_{t})\|_{L^{1}(w)\times L^{1}(\mu)} \\ &\leq \left\|S_{t}^{\Lambda}(\bar{\rho}_{0},\bar{\zeta}_{0})\right\|_{L^{1}(w)\times L^{1}(\mu)} + \int_{0}^{t} \left\|S_{t-s}^{\Lambda}(Z_{s}^{(1)},Z_{s}^{(2)})\right\|_{L^{1}(w)\times L^{1}(\mu)} ds \\ &\leq K_{1}e^{-\mathfrak{a}_{1}t}\|(\bar{\rho}_{0},\bar{\zeta}_{0})\|_{L^{1}(w)\times L^{1}(\mu)} + K_{1}\int_{0}^{t}e^{-\mathfrak{a}_{1}(t-s)}\left\|(Z_{s}^{(1)},Z_{s}^{(2)})\right\|_{L^{1}(w)\times L^{1}(\mu)} ds \\ &\leq K_{1}e^{-\mathfrak{a}_{1}t}\|(\bar{\rho}_{0},\bar{\zeta}_{0})\|_{L^{1}(w)\times L^{1}(\mu)} + |\varepsilon|\tilde{C}_{W}\int_{0}^{t}e^{-\mathfrak{a}_{1}(t-s)}\|(\bar{\rho}_{s},\bar{\zeta}_{s})\|_{L^{1}(w)\times L^{1}(\mu)} ds \\ &=:\mathcal{Q}(t), \end{split}$$

where \tilde{C}_K is a constant depending on W. We have, for all $t \ge 0$,

$$\begin{aligned} \frac{d}{dt}\mathcal{Q}(t) &= -\mathfrak{a}_{1}\mathcal{Q}(t) + |\varepsilon|\tilde{C}_{W} \| (\bar{\rho}_{t}, \bar{\zeta}_{t}) \|_{L^{1}(w) \times L^{1}(\mu)} \\ &\leq \left(-\mathfrak{a}_{1} + |\varepsilon|\tilde{C}_{W} \right) \mathcal{Q}(t), \end{aligned}$$

whence, by Grönwall's lemma,

$$\forall t \ge 0, \qquad \mathcal{Q}(t) \le K_1 \left\| (\bar{\rho}_0, \bar{\zeta}_0) \right\|_{L^1(w) \times L^1(\mu)} \exp\left(\left(-\mathfrak{a}_1 + |\varepsilon| \tilde{C}_W \right) t \right).$$

For all $t \ge 0$, we have

$$\|\rho_t - \rho_\infty\|_{L^1(w)} + |x_t - x_\infty| \le \|(\bar{\rho}_t, \bar{\zeta}_t)\|_{L^1(w) \times L^1(\mu)} \le \mathcal{Q}(t)$$

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and

$$\left\|\bar{\zeta}_{0}\right\|_{L^{1}(\mu)} \leq \int_{0}^{\infty} \left\|f(\varepsilon x_{\infty})\rho_{\infty}\right\|_{L^{1}} C_{h} e^{-\mathfrak{h}s} ds \leq \frac{\|f\|_{\infty} C_{h}}{\mathfrak{h}}.$$

Therefore, choosing $\varepsilon_W^{**} := \frac{\mathfrak{a}_1}{\tilde{C}_W} \wedge \varepsilon^*$, we easily see that there exist $C \ge 1$ and $c_W > 0$ such that (1.6) holds.

For (1.7), we use (5.9) instead of (5.8) and follow the same argument.

Appendix A. Time-elapsed neuron network model. Here, we compare simulations of (1.2) with simulations of the time-elapsed neuron network model [35]. If the firing rate function f does not depend on m and if we put

(A.1)
$$f(a, \varepsilon x_t) := \hat{f}(\eta(a) + \varepsilon x_t),$$

then (1.2) reduces to the time-elapsed neuron network model

(A.2a)
$$\partial_t \rho_t(a) + \partial_a \rho_t(a) = -f(a, \varepsilon x_t)\rho_t(a),$$

(A.2b)
$$\rho_t(0) = \int_0^\infty f(a, \varepsilon x_t) \rho_t(a) da,$$

(A.2c)
$$x_t = \int_0^t h(t-s) \int_0^\infty \int_0^\infty f(a, \varepsilon x_s) \rho_s(a) dads,$$

(A.2d)
$$\rho_0(a) = u_0(a).$$

Equation (A.2) is the population equation for nonadaptive SRM_0 neurons (or agedependent nonlinear Hawkes processes) [6]. As reported previously, (A.2) exhibits self-sustained oscillations for large ε or relaxation to a stationary state for small ε (see Figure 2). Note that in the special case where h is a Dirac delta distribution ("instantaneous transmission"), (A.2) can exhibit elaborate periodic patterns [36], but these patterns do not fulfill our definition of self-sustained population bursts.

Appendix B. Proof of Doeblin's theorem (Theorem 3.1). We follow [4]. We first observe that for any $t \ge 0$, S_t is nonexpansive, i.e.,

(B.1)
$$||S_t\mu||_{L^1} \le ||\mu||_{L^1} \quad \forall \mu \in L^1.$$

Indeed, writing μ_+ and μ_- the positive and negative parts of μ , respectively $(\mu = \mu_+ - \mu_- \text{ and } |\mu| = \mu_+ + \mu_-),$

$$|S_t\mu| \le |S_t\mu_+| + |S_t\mu_-| = S_t\mu_+ + S_t\mu_- = S_t|\mu|,$$

and we get (B.1) by integrating both sides. The result can then be shown in two steps.

Step 1. The Doeblin minorization condition (3.2) implies

(B.2)
$$\|S_T\mu\|_{L^1} \le (1 - \|\nu\|_{L^1}) \|\mu\|_{L^1} \quad \forall \mu \in L^1 : \int \mu = 0.$$

Indeed, the Doeblin minoration condition (3.2) and the fact that $\|\mu_+\|_{L^1} = \|\mu_-\|_{L^1} = \|\mu\|_{L^1/2}$ imply

$$|S_T\mu| \le |S_T\mu_+ - \|\mu_+\|_{L^1}\nu| + |S_T\mu_- - \|\mu_-\|_{L^1}\nu| = S_T|\mu| - \|\mu\|_{L^1}\nu,$$



FIG. 2. Same as Figure 1 but for the time-elapsed neuron network model (A.2). Simulations of a network of $5 \cdot 10^5$ nonadaptive SRM_0 neurons, approximating (A.2), with identical parameters (except for ε) and identical initial conditions. Neuronal parameters are the same as in Figure 1, except that f is replaced by (A.1). The ε have also been adapted.

and we get (B.2) by integrating both sides.

Then, by the estimate (B.2), there exists a unique $\rho_{\infty} \in L^1_+$ with $\|\rho_{\infty}\|_{L^1} = 1$ such that $S_T \rho_{\infty} = \rho_{\infty}$. The existence is obtained by taking a $\mu_0 \in L^1$ with $\|\mu_0\|_{L^1} = 1$ and defining $\mu_k = S_T \mu_{k-1}$ for all $k \ge 1$. The estimate (B.2) implies that $\{\mu_k\}_{k\ge 1}$ is a Cauchy sequence, and passing to the limit, we get that $\rho_{\infty} := \lim_{k\to\infty} \mu_k$ satisfies $S_T \rho_{\infty} = \rho_{\infty}$. If $\rho_{\infty} \in L^1_+$ with $\|\rho_{\infty}\|_{L^1} = 1$ also satisfies $S_T \rho_{\infty} = \rho_{\infty}$, taking $\mu = \rho_{\infty} - \rho_{\infty}$ in (B.2) implies $\rho_{\infty} = \rho_{\infty}$, whence the uniqueness.

Step 2. By the semigroup property, for any t > 0 we have

$$S_t \rho_\infty = S_t S_T \rho_\infty = S_T S_t \rho_\infty,$$

and from the uniqueness of the fixed-point of S_T , we get that $S_t \rho_{\infty} = \rho_{\infty}$. Hence, ρ_{∞} is the unique invariant probability measure.

The general estimate (3.3) is obtained by taking $\mu = u_0 - \rho_{\infty}$ and writing $t = \lfloor \frac{t}{T} \rfloor T + r_1$ with $0 \le r_1 < T$. Indeed, using the semigroup property and nonexpansivity, we get

$$\|S_t\mu\|_{L^1} = \|S_{r_1}S_{\lfloor \frac{t}{T} \rfloor}\mu\|_{L^1} \le \|S_{\lfloor \frac{t}{T} \rfloor T}\mu\|_{L^1} \le (1 - \|\nu\|_{L^1})^{\lfloor \frac{t}{T} \rfloor}\|\mu\|_{L^1},$$

which implies (3.3) for the given K and \mathfrak{a} .

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