

### Efficiency-boosted semiconductor optical amplifiers via modedivision multiplexing

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## OPTICA

# Efficiency-boosted semiconductor optical amplifiers via mode-division multiplexing

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Semiconductor optical amplifiers (SOAs) are a fundamental building block for many photonic systems. However, their power inefficiency has been setting back operational cost reduction, circuit miniaturization, and the realization of more complex photonic functions such as large-scale switches and optical phased arrays. In this work, we demonstrate significant gain and efficiency enhancement using an extra degree of freedom of light—the mode space. This is done without changing the SOA's material design, and therefore high versatility and compatibility can be obtained. Light is multiplexed in different guided modes and reinjected into the same gain section twice without introducing resonance, doubling the interaction length in a broadband manner. Up to 87% higher gain and 300% higher wall-plug efficiency are obtained in a double-pass SOA compared to a conventional single-pass SOA, at the same operating current, in the wavelength range of 1560–1580 nm. © 2023 Optica Publishing Group under the terms of the Optica Open Access Publishing Agreement

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#### **1. INTRODUCTION**

The pursuit of high energy efficiency has fueled major advances in photonic microchips, as it promises denser integration, lower costs, and reduced environmental impact of advanced technologies [1–5]. However, the semiconductor optical amplifier (SOA), a crucial element in photonic integrated circuits (PICs), still suffers from being very power-hungry and inefficient, particularly when weak signals are to be amplified. For an SOA before saturation, the same electrical power is needed to maintain a certain optical gain, regardless of the input optical power. This is due to the fact that the strength of stimulated emission depends linearly on the local photon density itself [6]. Consequently, the injected carriers are not efficiently utilized for stimulated emission before the photon density is sufficiently high, resulting in an overall low wall-plug efficiency (WPE) and uneven depletion of carriers. Furthermore, when signals are low, high levels of amplified spontaneous emission (ASE) are generated, contributing to increased noise rather than signal amplification.

Several approaches have been taken to improve the SOA's energy efficiency, including increased optical confinement [7–9], deeper quantum wells (QWs) [10–12], better heat sinking [13–17], higher injection efficiency [18–20], etc. Slow light has also been employed for SOA efficiency enhancement [21], but only optical pumping has been demonstrated. To the best of

our knowledge, all these existing works focus on improving the SOA at the device level, so they are specific to a certain technology platform. They are all single-pass devices, so they still suffer from the problem of spatially inhomogeneous and inefficient carrier utilization.

Here, we propose a novel direction of SOA efficiency enhancement at the circuit level, leveraging on-chip mode-division multiplexing (MDM) [22–24]. In our approach, light gets reinjected back into the gain section, reusing the same gain material in both directions. Therefore, the interaction length is multiplied, and the available carriers are utilized in a more homogeneous way. Resonance is suppressed since each pass is spatially multiplexed in separate modes. Since it is a circuit-level method, it is generic and can be applied to many active–passive (A-P) platforms with distinct realizations of the components.

Efficiency enhancement using MDM has been demonstrated for linear devices such as thermo-optic [25] or electro-optic [26–28] modulators, but it has never been done for SOAs, which are highly nonlinear: the gain process is naturally exponential along the propagation direction, if not considering saturation. Therefore, compared to linear devices, a higher efficiency gain can be expected in SOAs with the same number of passes, or fewer passes may be needed to achieve a substantial efficiency boost. The challenge of integrating the MDM technique with active devices mainly lies in their distinct waveguide cross sections. Low-loss, compact, and broadband MDM (de)multiplexers normally require high-confinement rectangular or ridge waveguides for efficient evanescent coupling [22,29]. Therefore, they are normally realized on high-confinement photonic platforms such as Si-on-insulator (SOI), which uses a thick dielectric layer for index contrast. In contrast, typical SOAs are realized in III-V material and have their optical modes buried in micrometer-scale semiconductor layer stacks for efficient electrical pumping [30].

In this paper, the integration challenge is tackled by the InP-membrane-on-silicon (IMOS) technology [31,32]. Via wafer-to-wafer bonding creating a high-index contrast cross section similar to that of SOI, the InP passive circuitry is miniaturized into a nanophotonic membrane. SOAs can be conveniently integrated into the same membrane by leveraging epitaxial growth at the wafer scale, allowing for a scalable PIC solution. Leveraging the active-passive IMOS technology, we demonstrate here, for the first time, significant SOA efficiency enhancement through an MDM double-pass (DP) gain section. The design, fabrication, and characterization results are subsequently presented in the following sections.

#### 2. RESULTS

#### A. Circuit Design

The generic circuit architecture of the DP-SOA is shown in Fig. 1(a). The light path is as follows: input fiber grating coupler (FGC, TE<sub>0</sub>); MDM (de)multiplexer (TE<sub>0</sub> passing through); SOA (TE<sub>0</sub> gain); MDM (de)multiplexer (TE<sub>0</sub> passing through); U-bend (TE<sub>0</sub>); MDM (de)multiplexer (TE<sub>0</sub> converted to TE<sub>1</sub>); SOA (TE<sub>1</sub> gain); MDM (de)multiplexer (TE<sub>1</sub> converted back to TE<sub>0</sub>); output FGC (TE<sub>0</sub>). Although a high-order mode is used inside the circuit, fiber interfacing is done using only the fundamental mode. As seen in Fig. 1(b), by breaking apart the U-bend and connecting to FGCs, an MDM  $2 \times 2$  SOA is obtained, so the TE<sub>0</sub> and TE<sub>1</sub> gain channels can be independently accessed and characterized. This single-pass TE<sub>0</sub> channel can be used as a benchmark to quantify the gain performance of the DP-SOA.

All photonic components on the IMOS platform are based on a thin InP-based epilayer adhesively bonded on Si. The SOA gain section is based on a regrowth-free twin-guide scheme [31,32]. The active functional layers including n- and p-InP claddings, InGaAs metal contacting layer, and the core are epitaxially grown on top of the 300 nm thick i-InP layer, in which the passive waveguide and other photonic circuitries are fabricated. The core is made of InGaAsP with a bandgap corresponding to the 1250 nm



**Fig. 1.** Schematics of the MDM-SOA circuits (not to scale). (a) Top view of the DP-SOA. (b) Top view of the MDM  $2 \times 2$  SOA. (c) Cross-sectional structure of the passive waveguide and SOA. The color scheme is the same as in (a). (d) 3D illustration of the SOA active–passive taper, looking from the substrate side. (e) Zoom-in schematics of the MDM (de)multiplexer. (f) Mode coupling between TE<sub>0</sub> and TE<sub>1</sub> through the MDM (de)multiplexer in (e).

wavelength (Q1.25) and contains a compressively strained 4multi-QW (MQW). Halfway fabrication, the whole layer stack is flipped and buried into the benzo-cyclobutene (BCB) bonding layer, creating the cross section shown in Fig. 1 (c). The whole wafer is also planarized and buried in polyimide (PI). The Q1.25 layer immediately under the i-InP layer is a wet etch-stop layer, which ensures an accurate etch depth for the SOA ridge and tapers without damaging the passive waveguide (see Section 1 in Supplement 1). The width of the SOA is designed to be 2.2  $\mu$ m, supporting both TE<sub>0</sub> and TE<sub>1</sub> modes in the core region. Computed modal fields are shown in Supplement 1, Fig. S2. The confinement factors in the 4-MQW are calculated to be 4.31% and 4.29% for TE<sub>0</sub> and TE<sub>1</sub>, respectively.

A-P coupling between the SOA and passive waveguide is done via a vertical evanescent taper, as seen in Fig. 1(d). In this taper, by separately tapering down the thick p-InP cladding to eliminate the vertical high-order cladding modes, unwanted mode conversions are suppressed (see Section 3 in Supplement 1 for details). The p-cladding, core, and n-cladding taper lengths are 10  $\mu$ m, 35  $\mu$ m, and 20  $\mu$ m, respectively. The p-cladding and core taper both start at the gain section, and therefore the total taper length is 55  $\mu$ m. The tip widths of the tapers are set to 200 nm for fabrication robustness [33]. According to eigenmode expansion (EME) simulations, a high 90% transmission is achieved for both TE<sub>0</sub> and TE<sub>1</sub> modes in this taper, if not considering material absorption.

Figure 1(e) shows the zoom-in schematics of the dual-core adiabatic tapers as an MDM (de)multiplexer. This MDM (de)multiplexer contains a multimode bus connecting directly to the SOA, and a TE<sub>0</sub> access waveguide for adiabatic coupling from/to the TE<sub>1</sub> mode in the bus waveguide. The gap between waveguides is designed to be 100 nm, allowing for simultaneous good manufacturability and a small footprint (50  $\mu$ m coupling length). Detailed structural parameters and performance of the MDM (de)multiplexer can be found in Section 4 of Supplement 1. Figure 1(f) shows the light propagation through this MDM

(de)multiplexer when  $TE_0$  light of 1550 nm wavelength is launched in the access waveguide. This design is demonstrated to have high fabrication tolerance on waveguide width ( $\pm 25$  nm), low insertion loss (1 dB), and broad optical bandwidth (90 nm) [22].

#### **B.** Fabrication

The fabrication is done on both sides of the InP membrane (both before and after bonding to Si; see Section 1 in Supplement 1 for details) to realize the twin-guide S-shaped cross-sectional structure. After bonding, the distance between the InGaAs layer and Si layer is around 2  $\mu$ m, as seen in Fig. 1(c). This BCB bonding layer planarizes surface topology, which enables high fabrication tolerance. The passive circuitries are fabricated after bonding from the back side of the original InP wafer and are therefore on top of the BCB layer. A microscope image of the fabricated DP-SOA is shown in Fig. 2(a). The zoom-in scanning electron microscope (SEM) images of the A-P taper before bonding and the coupling region of the MDM (de)multiplexer are shown as insets of the figure. Since e-beam lithography (EBL) is used for pattern definition, high overlay accuracy (~20 nm) is achieved for all mask layers.

#### C. Enhanced Amplified Spontaneous Emission

The DP-SOA and a nearby  $2 \times 2$  SOA with the same gain section dimensions ( $2.2 \times 500 \mu m$ ) are measured to characterize the gain enhancement. The wafer stage is kept at 10°C, and cleaved single-mode fibers (SMFs) are used for optical interfacing with the FGCs at an emission angle of 80° with respect to the waveguides. First, the LIV curves are obtained using a current source and broadband optical power meter for both channels of the  $2 \times 2$ SOA and DP-SOA. As seen in Fig. 2(b), owing to the uniform wafer-scale process, the IV curves of the  $2 \times 2$  SOA and DP-SOA overlap well with each other, and they both have a series resistance of  $\sim 7 \Omega$ . As can be seen, due to an extra pass in the gain section,



**Fig. 2.** ASE enhancement of the DP-SOA. (a) Microscope image of the fabricated DP-SOA (scale bar: 100  $\mu$ m). The insets are zoom-in SEM images. Top left: SOA active–passive taper before bonding (scale bar: 10  $\mu$ m). Here the i-InP waveguide is not yet fabricated. Top right: MDM (de)multiplexer fabricated after bonding. (b) LIV curves of the DP-SOA and reference 2 × 2 SOA of the same length. The received optical power for each SOA channel is labeled in the figure. Inset: Zoom-in of the IV curves around 1 kA cm<sup>-2</sup>. The difference in bias at 1 kA cm<sup>-2</sup> is 0.002 V. (c) ASE spectra of the DP-SOA and reference 2 × 2 SOA of the same length at a current density of 4.1 kA cm<sup>-2</sup>.

 $2 \times$  higher ASE power is obtained at the same current density for the DP-SOA compared to the  $TE_0$  channel of the 2 × 2 SOA (i.e., single pass). Since the TE<sub>0</sub> gain channel does not involve mode coupling over the MDM (de)multiplexer [22], the excess loss is negligible (Supplement 1, Fig. S6). Therefore, this channel can be used to benchmark the performance of the DP-SOA. The  $TE_1$  ASE is lower than that of the  $TE_0$ , which is due to the  $TE_1$ channel having: (1) extra loss from the MDM (de)multiplexer (Supplement 1, Fig. S6). 2) a slightly lower confinement factor and higher sidewall scattering loss in the gain section. Optical power saturation occurs at around 5 kA cm<sup>-2</sup>, and it decreases beyond this point. ASE spectra collected by an optical spectrum analyzer (OSA) at the current density of 4.1 kA cm<sup>-2</sup> for the DP-SOA and  $2 \times 2$  SOA are shown in Fig. 2(c). Near-smooth spectra (<0.5 dB fluctuation) are obtained for both TE<sub>0</sub> and TE<sub>1</sub> channels, indicating the high mode conversion efficiency and low cross talk in the MDM (de)multiplexers. More significant fluctuations can be observed in the spectra of the DP-SOA. These fluctuations come from the residual optical feedback:  $\sim 4\%$  reflection of the SMF facet,  $\sim 1\%$  refection of the grating couplers, and  $\sim 1\%$  effective reflection from the MDM (de)multiplexers.

#### **D. Enhanced Optical Gain and Power Efficiency**

Since MDM SOA circuits involve more than one mode, measurement of optical gain is done by directly comparing the output and input optical powers. To measure the small-signal gain, an external tunable laser is used as the input, whose peak wavelength is tuned from 1520 nm to 1600 nm with a step of 2 nm. The on-chip input power after the FGC is set at -13.5 dBm. The input polarization is set for maximum optical transmission under a constant current injection. Since the FGCs and MQWs are designed for TE, and TM gain is far less than TE gain, a TE input polarization is assumed. Then, at various injection current densities, optical spectra are collected by an OSA for each input–output channel combination, for both the 2  $\times$  2 SOA and DP-SOA. Then, the peak powers of the spectra are extracted. In this way, the broad ASE power can be excluded from the gain measurement. Nearby (<100  $\mu$ m distance) reference waveguides are also measured at each corresponding wavelength point using the same configuration to calibrate out the FGC loss. The passive propagation loss in the access waveguides between FGCs and MDM (de)multiplexers is compensated for, and therefore zero-net-loss levels are obtained.

Figures 3(a)-3(c) show the received optical spectra at the current density of  $4.1 \text{ kA cm}^{-2}$ , with the blue dots indicating zero-net-loss levels, for the 2 × 2 SOA with input through the TE<sub>0</sub> port, input in the TE<sub>1</sub> port, and the DP-SOA, respectively. As can be seen, positive optical gain is achieved in the wavelength range of 1540–1600 nm for all three cases, and there are no signs of lasing of the SOAs themselves. As seen in Figs. 3(a) and 3(b), for the 2 × 2 SOA, the modal cross talk (TE<sub>0</sub>-TE<sub>1</sub> and TE<sub>1</sub>-TE<sub>0</sub> transmission) levels are -16.2 dB to 28.6 dB and -16.1 dB to -27.8 dB in the wavelength range of 1520–1570 nm for TE<sub>0</sub> and TE<sub>1</sub> input, respectively. This indicates that independently controllable and accessible gain channels are achieved. Generally, the cross talk increases as wavelength decreases, which could be due to stronger absorption and re-emission [34,35] in the gain section at shorter wavelengths.

Figure 3(d) shows the gain values by comparing the measured peak powers and zero-net-loss levels. As can be seen, the gain of the DP-SOA matches well with the summation of values obtained from the separate gain channels. At shorter wavelengths where the material gain is negative, the DP loss also corresponds well to the total loss in TE<sub>0</sub> and TE<sub>1</sub> channels. This indicates that light successfully travels through the gain section twice, and the zero-net-loss levels are correctly set. The 3 dB gain bandwidth measured for the DP-SOA is around 50 nm, confirming its broadband operation. The gain difference measured between TE<sub>0</sub> and TE<sub>1</sub> modes is



**Fig. 3.** Gain enhancement of the DP-SOA. (a), (b) Optical spectra of the  $2 \times 2$  SOA obtained at current density of 4.1 kA cm<sup>-2</sup>, using a tunable laser with -13.5 dBm on-chip input power. The spectra of different input wavelengths are overlapped, and the input and output channels are labeled on the bottom right and top right of the figures, respectively. Blue dots indicate zero-net-loss levels measured by reference waveguides. (c) Optical spectra of the DP-SOA at the same measurement conditions as in (a) and (b). (d) Optical transmission through the SOAs, obtained by subtracting the zero-net-loss levels from the peak powers in (a)–(c). Dots are actual data points, while solid curves are Wiener-filtered to eliminate the fluctuations from FGC and fiber reflections.

1.5–2.8 dB, which is partly due to the wavelength-dependent loss of the (de)multiplexers (Supplement 1, Fig. S6).

Compared to single-pass in TE<sub>0</sub>, significant gain enhancement is achieved in the wavelength range of 1538-1600 nm for the DP-SOA. In the wavelength range of 1560-1580 nm, over 158% (58% more) optical gain is achieved, without increasing the operating current. The WPE (WPE =  $P_{optical}/P_{electrical}$ ) can be calculated using the on-chip input and output optical powers  $(P_{\text{optical}} = P_{\text{out}} - P_{\text{in}})$  and total electrical power  $(P_{\text{electrical}} = \text{VI})$ . In this 20 nm bandwidth, a > 180% extra WPE is achieved. At 1570 nm in particular, 11.6 dB gain is achieved for the DP-SOA, while the gain measured for single-pass TE<sub>0</sub> is 6.2 dB. Therefore, the net gain enhancement is 5.4 dB, or 87% more than the original value, for the same operating current. The WPE for single-pass and DP are 0.3% and 1.2% respectively. A 300% extra WPE is achieved. At shorter wavelengths of 1520-1536 nm, >100% absorption enhancement can be observed, which partly attributes to the extra loss from the MDM couplers.

Figures 4(a)–4(c) show the gain as a function of injection current density, at various wavelengths. As can be seen, the DP-SOA and single-pass  $TE_0$  channel have similar zero-net-loss current densities, but beyond zero-net-loss, the DP gain increases faster and consistently outperforms the single-pass  $TE_0$  gain. From another perspective, the same gain can be achieved with a lower injection current density, or less energy consumption, for the DP-SOA.

#### E. Saturation Effects

Saturation effects of the DP-SOA are measured by inserting sequentially an erbium-doped fiber amplifier (EDFA) and a variable optical attenuator (VOA) between the tunable laser and the input fiber. The measurement schematics can be seen in Supplement 1, Fig. S7. The optical power before the VOA is set to 17 dBm (tunable laser + EDFA), while the power at the input

fiber is controlled by the VOA. Optical spectra are recorded at the output for each VOA setting, and the optical gain values are obtained by the peak powers subtracting zero-net-loss levels, the same as mentioned above. To calculate the on-chip input power, reference waveguides are measured to calibrate out the FGC loss as mentioned above.

Figure 5(a) shows the on-chip gain as a function of the on-chip input power for the DP-SOA at a current density of 4.1 kA cm<sup>-2</sup> and 1570 nm wavelength. The same result for the  $TE_0$  channel of the  $2 \times 2$  SOA is also shown as a benchmark. As can be seen, compared to the benchmark, the DP-SOA consistently offers a higher gain of up to 5 dBm on-chip input power. At higher input power, saturation occurs. The input saturation power  $P_{\text{sat}}$  of the DP-SOA is -1.8 dBm, which is lower than that of the single-pass TE<sub>0</sub> (6.8 dBm). This is as expected since the DP-SOA has effectively a doubled interaction length, which significantly enhances light-matter interaction, and therefore, carriers are utilized more exhaustively and deplete faster. To further increase the saturation power, the optical confinement can be lowered, but at the cost of optical gain and WPE. In Fig. 5(a), optical gain increases slightly with input power before it goes down, which could be due to the beginning and end of the SOA being not well pumped and acting as saturable absorbers [36]. This can be solved by increasing the pad thickness and making the SOA taper passive.

Figure 5(b) shows the logarithmic plots of the WPE for both the DP-SOA and single-pass  $TE_0$ . As can be seen, since the extra gain is obtained without needing more current injection, up to 300% more WPE is obtained for the DP-SOA. The advantage shrinks as the gain becomes saturated, but the DP-SOA remains 100% more efficient until the on-chip input power reaches -5 dBm. At high input powers, the WPE goes down, and this is anticipated to result from increased nonlinear absorption.



**Fig. 4.** (a), (b) Optical gain as functions of the injection current densities at various wavelengths, calculated by subtracting the zero-net-loss powers from the obtained peak powers. Solid lines and dashed lines are transmissions from the signal channels ( $TE_0$ - $TE_0$  and  $TE_1$ - $TE_1$ ) and cross talk channels ( $TE_0$ - $TE_1$  and  $TE_1$ - $TE_0$ ), respectively. (c) Optical transmission of the DP-SOA obtained using the same method as in (a) and (b).



**Fig. 5.** Saturation effects of the DP-SOA. (a) On-chip gain as functions of on-chip input power at a wavelength of 1570 nm, measured with the  $TE_0$  channel of the 2 × 2 SOA and DP-SOA. The input saturation points are indicated by empty squares in the figure. (b) WPE as functions of on-chip input power, calculated for the TE<sub>0</sub> channel of the 2 × 2 SOA and DP-SOA. The inset shows the extra amount of WPE obtained with the DP-SOA.

#### 3. DISCUSSION

In this study, a maximum 6.2 dB single-pass gain is achieved at  $4.1 \text{ kA cm}^{-2}$ . This single-pass gain achieved is comparable to what has been realized on III-V-on-Si transfer-printing platforms with alignment-tolerant taper designs [37], if scaled proportionally in length. It is lower compared to those on native InP substrates with a similar geometric design and the same MQW, in which ~10 dB optical gain can be achieved [38]. With specific optimizations on the material and geometric dimensions, gain values of more than 20 dB may be achieved [39]. However, as our work targets gain boost at the circuit level, its effectiveness is independent of the individual SOA component used. We expect a significant gain and efficiency boost if our method is combined with high-gain SOA with optimized component-level design.

The optical gain difference between the single-pass case in this work and a conventional SOA with a similar component-level design and the same MQW, but on the native substrate [38], is attributed to the following: the loss in the taper, thermal isolation due to the bonding layer, and the effectively deep-etched design. The passive loss in a single taper is simulated to be around 0.5 dB (Section 3 in Supplement 1), while the absorption is estimated to be 0.8 dB, assuming a 5000 dB cm<sup>-1</sup> material loss in a single QW. Therefore, the coupling loss for one taper is estimated to be around 1.3 dB. This taper loss is comparable to heterogeneously integrated InP/Si SOAs, which have similar A-P tapers [40]. This supports our estimation as the refractive indices of Si and InP are close to each other, meaning similar taper dimensions. Based on this estimation, an 8.8 dB TE<sub>0</sub> single-pass gain in the active region can be achieved without this coupling loss. The remaining  $\sim 1 \text{ dB}$ difference is attributed to worse thermal dissipation and higher sidewall recombination due to the more exposed active section. It should be noted that neither of these is a fundamental limitation: the taper can be made passive through QW intermixing [41] or ion implantation, or even be completely eliminated using butt-joint regrowth [42,43]; the thermal performance can be improved by thermal shunting to the silicon substrate [14], or using highthermal-conductivity substrates [16]. The sidewall recombination velocity can be greatly reduced by wet chemical treatments [18].

Excluding grating couplers, which is not a fundamental limiting factor, the maximum possible gain of the DP-SOA is determined by the cross talk of the MDM (de)multiplexer, since cross talk manifests itself as reflections into the same mode at the U-bend. In this work, -20.3 dB cross talk is obtained, but this may be further reduced to -33 dB by altering the waveguide design [44]. This

level of reflection is comparable to that in bulk InP SOAs with butt-joint A-P interfaces [45].

In conclusion, we demonstrate that gain and efficiency can be significantly enhanced by reinjecting the amplified light for a second pass through the SOA gain section. Each pass in the SOA is in separate transversal modes, so resonance is suppressed and broadband operation is achieved. Light also passes through the gain section from opposite directions, therefore homogeneously utilizing the available carriers. Compared to the single-pass case, the ASE power is improved by >100%, and 58%–87% more small-signal gain is obtained at the same injection current in the wavelength range of 1560–1580 nm; 300%–100% more WPE is achieved for on-chip input powers up to -5 dBm.

Since the method demonstrated in this paper works at the circuit level, it can be exploited in combination with other efficiency enhancement methods to achieve even better performance. It also allows gain and efficiency enhancement without changing the SOA layer stack itself, providing more compatibility and versatility in circuit design. Although demonstrated on the InP membrane platform, in principle, the concept is applicable to a wide range of A-P integration platforms including hybrid, heterogeneous, or butt-joint technologies, such as III-V/Si platforms with Si or  $SiN_x$  waveguides. With the increasing demand for building block performance in densely integrated PICs, this work may be used in various applications such as optical beam steering, neuromorphic photonics, photonic matrix computation, photonic switching, etc. The DP-SOA may also be used to realize novel laser types, and the design principle can be leveraged to boost the extinction ratio of electro-absorption modulators without increasing the length, so modulation speed will not be compromised. It is also worth noting that this is the first time that MDM is realized in active photonic circuits on-chip, which opens up new opportunities in device and circuit design, such as multi-wavelength lasers or channel multi-casting.

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**Data availability.** Data underlying the results are available upon reasonable request.

Supplemental document. See Supplement 1 for supporting content.

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