

Effect of biological polymers on mobility and run-out distance of cohesive and non-cohesive sediment gravity flows

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1 Effect of biological polymers on mobility and run-out distance

² of cohesive and non-cohesive sediment gravity flows

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7 ABSTRACT

8 Lock-exchange experiments were carried out to investigate the effect of biologically cohesive 9 extracellular polymeric substances (EPS) on the mobility of sediment gravity flows laden with 10 physically cohesive clay, non-cohesive coarse silt and non-cohesive fine sand. The results reveal significant differences in the head velocity, run-out distance and deposit shape of these flows related 11 12 to differences in physical cohesion, particle size, and EPS content. These differences are captured in a three-way coupling model of turbulent forces, cohesive forces, and particle settling velocity. In 13 14 general, biological cohesion reduces flow mobility, demonstrated most clearly by a progressive 15 decrease in the run-out distance of the silt and clay flows, as the EPS concentration is increased. This 16 reduction in flow mobility is caused by the dominance of cohesive forces over turbulent forces, which 17 comprise turbulence attenuation and the bulk settling of a biologically cohesive gel in which EPS form a pervasive network of bonds between the sediment particles. However, sand-laden gravity flows 18 19 were found to behave in a markedly different way, in that the head velocity and run-out distance first 20 increase and then decrease, as the EPS concentration is increased. The increase in sand flow mobility 21 is inferred to be caused by a reduction in the settling velocity of the sand particles, as the EPS cause 22 an increase in flow viscosity at EPS concentrations that are sufficiently low to maintain turbulent flow. 23 Once the EPS concentration is high enough for turbulence attenuation, the sand flows start to agree

24 with the silt and clay flows in establishing a negative correlation between flow mobility and EPS 25 concentration caused by gelling. The experimental data also uncovered that deposits formed by EPSrich, turbulence-attenuated flows are shorter and thicker and have more abrupt terminations than 26 27 deposits formed by EPS-free or EPS-poor turbulent flows. The larger thickness of these deposits is 28 partly caused by the ability of EPS to retain water and form matrix-supported textures. Earlier work 29 has shown that EPS is common in many sedimentary environments, including those where sediment 30 transport takes place regularly by particulate density currents. Combined with the increasing rate at 31 which man-made structures, such as pylons and communication cables, appear in these 32 environments, we argue that there is a need to incorporate the results of this study in applied models that aim to mitigate damage to such structures by sediment gravity flows. 33

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35 Key words: Sediment gravity flows; Physical experiments; EPS; cohesion; clay; silt; sand

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37 **1. Introduction**

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39 Sediment gravity flows (SGFs) comprise mixtures of water and sediment driven by excess density and 40 gravity forces (Middleton and Hampton, 1973). Bottom-hugging SGFs originate from sediment-laden 41 river flows that plunge after entering seas and lakes, i.e., hyperpycnal flows, and from slope failures 42 caused by, for example, earthquakes, storm waves, slope oversteepening after rapid sedimentation, and fishing gear dragged across a loose substrate (Postma, 2011). SGFs are a global phenomenon, with 43 44 the capacity of transporting large volumes of sediment, carbon, nutrients, and pollutants (e.g. 45 microplastics) to the deep ocean (Postma, 2011; Baker et al., 2017; Heerema et al., 2019; Craig et al., 46 2020). These flows can cause serious damage to subaqueous communication cables and other deep-47 water engineering infrastructure (Inman et al., 1976; Talling et al., 2015), and they have been linked 48 to the formation of tsunami (Johnson et al., 2017). Moreover, three-dimensionally stacked deposits

49	of SGFs create submarine fans that are amongst the largest reservoirs of oil and gas on Earth (Reading
50	and Richards, 1994; Heerema et al., 2019). Despite recent successes in studying SGFs in modern lakes
51	and oceans (Talling et al., 2013; Zabala et al., 2017), and valuable understanding gained from
52	traditional outcrop and core studies of SGF deposits, laboratory simulations remain a valued method
53	for obtaining a physical understanding of the dynamics of SGFs and the style of their deposits (e.g.
54	Baker et al., 2017; Craig et al., 2020). Laboratory experiments are unique in allowing the parameters
55	that control the complex dynamics of SGFs to be studied in isolation. This includes SGFs that are
56	cohesive because of the presence of physically active clay particles (Baker et al., 2017) and biologically
57	active extracellular polymeric substances (EPS) (Craig et al., 2020). The effect of biological cohesion
58	on the mobility of clay, silt, and sand flows is the principal target of the present study.

59

60 **2. Background and rationale of study**

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62 2.1. Types of sediment gravity flow

63

64 SGFs have been classified using rheology and particle support mechanism (Middleton and Southard, 1984; Postma, 1986; Dasgupta, 2003; Haughton et al., 2009). In its simplest form, the rheological 65 properties of a SGF can be characterized by yield strength, the resistance to deformation, and dynamic 66 67 viscosity, the dependence of shear rate on applied shear stress (George and Qureshi, 2013; 68 Widyatmoko, 2016). Turbidity currents are relatively dilute SGFs that exhibit Newtonian rheological 69 behaviour (Middleton, 1993), i.e., without a yield strength and with a constant molecular viscosity 70 (Chereskin and Price, 2001). Examples of denser SGFs with non-Newtonian behaviour are subaqueous 71 debris flows, mud flows, and slides (Iverson, 1997). These flows may have a yield strength and a 72 viscosity that decreases or increases as the shear stress is increased (Allen, 2009; Sumner et al., 2009; 73 Manica, 2012). The main particle support mechanism in turbidity currents is turbulence (Middleton 74 and Hampton, 1973; Heerema et al., 2019). Particles in debris flows, mud flows, and slides are

75 supported mostly by matrix strength (Lowe, 1979; Dasgupta, 2003), provided by non-cohesive silt, 76 physically cohesive clay (Baker et al., 2017) and biologically cohesive EPS (Craig et al., 2020). Cohesive 77 forces tend to attenuate or fully suppress turbulent forces (Baker et al., 2017; Craig et al., 2020), often 78 leading to a reduction in SGF mobility. Here, mobility is defined as the ability of SGFs to transport 79 suspended load over a certain maximum distance, referred to as the run-out distance. Flow mobility 80 is also governed by other processes, such as sediment erosion and deposition, mixing with ambient 81 fluid, and particle settling velocity, as coarser particles are more likely to be deposited and less likely 82 to be eroded from the sediment bed than finer particles, especially if these particles are non-cohesive.

83 The change from turbulent to turbulence-suppressed SGF depends on the flow velocity and the 84 suspended sediment type and concentration, with high-velocity, low-viscosity flows more likely to be 85 turbulent than low-velocity, high-viscosity flows. It is therefore not possible to define a single 86 threshold sediment concentration for changes between turbidity current and mud or debris flow 87 (Baker et al., 2017; Heerema et al., 2019). A range of transitional flows with both turbulent and laminar 88 behaviour has been defined, including high-density turbidity currents (Middleton and Hampton, 1973; 89 Lowe, 1982; Baker et al., 2017), lower and upper transitional plug flows (Baas et al., 2009, 2011), 90 hybrid flows (Haughton et al., 2009), and top-transitional plug flows (Hermidas et al., 2018; Craig et 91 al., 2020). Particularly relevant to this study are the high-density turbidity currents and top-transitional 92 plug flows, which have a lower region of attenuated turbulence separated from a fully turbulent upper 93 region where mixing with the ambient fluid takes place. Such flows are intermediate between fully 94 turbulent, low-density turbidity currents and laminar mud or debris flows (called plug flows by 95 Hermidas et al. [2018] and Craig et al. [2020]).

96

97 2.2. Sediment type and physical and biological cohesion

99 In the laboratory experiments with clay-laden and silt-laden SGFs of Baker et al. (2017), cohesion 100 started to affect the flow mobility at a volumetric concentration of 10%. Above 10%, the head 101 velocities of the flows started to diverge; the silt flows retained a higher velocity than the clay flows. 102 This threshold concentration will vary with flow velocity, but this difference in flow mobility signifies 103 a stronger network of particle bonds in the clay flows and therefore a stronger turbulence attenuation 104 (e.g. Kuenen, 1965; Felix et al., 2005). Frictional forces started to reduce the mobility, and also the 105 run-out distance, of the non-cohesive silt flows at concentrations that were at least three times 106 greater than those of the clay flows (Baker et al., 2017; their figure 10).

107 More recent research by Craig et al. (2020) showed that biological cohesion can have a similar effect 108 on the mobility of SGFs as physical cohesion, but at concentrations that are several orders of 109 magnitude lower. Biological cohesion is caused by EPS secreted by microorganisms, such as diatoms 110 and bacteria, mainly for protection, communication and interaction, carbon storage, nutrient entrapment, and aggregation (Wingender et al., 1999; Wolfaardt et al., 1999; Sandhya and Ali, 2015; 111 112 Wang et al., 2015; Costa, 2018). These polymers form a gel-like, three-dimensional structure that 113 stabilises the microbial aggregates via different physicochemical mechanisms, including dispersion 114 forces, electrostatic interactions, and hydrogen bonds (Flemming et al., 2000). The microorganisms 115 and their EPS are found in many depositional environments, from rivers and estuaries to hypersaline 116 lakes and deep-sea hydrothermal vents (Decho and Gutierrez, 2017). EPS also act as a cohesive binder 117 between the organisms and the surrounding sediment particles (Chenu, 1995; Wolfaardt et al., 1999), 118 forming biofilms on the sediment surface that are resistant to erosion by overriding flows (Malarkey 119 et al., 2015). Moreover, EPS induce more pervasive cohesion in muddy and sandy substrates below 120 biofilms (e.g. Malarkey et al., 2015; Hope et al., 2020) and assist in flocculation and gelling of clay in 121 suspension flows (Flemming and Wingender, 2010; Gerbersdorf and Wieprecht, 2014). Craig et al. 122 (2020) found that the mobility of SGFs carrying 15%, 22% and 23% kaolinite clay was reduced 123 significantly by adding between 0.052% and 0.265% by weight of EPS, matching concentrations 124 measured by Craig et al. (2020) in surficial deep-water sediment offshore New Zealand.

125

126 2.3. Research aims

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128 The present study aims: (a) to determine the effect of non-cohesive silt and sand on the mobility of 129 SGFs with and without EPS; (b) to compare the effect of EPS on flow mobility between physically 130 cohesive and non-cohesive SGFs. Craig et al. (2020) used mixtures of clay and EPS, but SGFs also carry 131 non-cohesive silt and sand (Kuenen, 1951; Britter and Simpson, 1978; Parker et al., 1987; Middleton 132 and Neal, 1989; Baas et al., 2005). Like clay-laden SGFs, silt and sand flows may contain EPS supplied 133 from the source area or eroded from the basin floor during transport. We hypothesise that EPS change 134 the flow properties from Newtonian turbidity currents to non-Newtonian transitional flows and mud 135 or debris flows with attenuated turbulence, which in turn reduces the flow velocity and run-out 136 distance, and modifies the properties of their deposits, as for the mixed clay-EPS flows of Craig et al. 137 (2020). Testing this hypothesis is essential for mitigating damage to subaqueous engineering 138 infrastructure, forecasting the dispersal of nutrients and pollutants in the deep sea, and predicting the 139 three-dimensional architecture of submarine fans.

140

141 **3. Methods**

142

Eleven laboratory experiments were conducted using a lock-exchange tank, 5 m long, 0.2 m wide and 0.5 m deep, in the Hydrodynamic Laboratory, School of Ocean Sciences, Bangor University (Fig. 1). The 0.31-m long reservoir behind the lock gate was filled with 0.3 m of a sediment–water or sediment– EPS–water mixture, while the channel downstream of the lock gate was filled with water only to the same depth. All experiments used natural seawater (density, $\rho = 1027$ kg m⁻³; salinity, s = 35 PSU) sourced from the Menai Strait, a tidal strait next to the Hydrodynamic Laboratory. Three different sediment types were used: cohesive kaolin clay (median size, $D_{50} = 0.0091$ mm); non-cohesive, very well-sorted, silt-sized, spherical glass beads ($D_{50} = 0.050$ mm); and non-cohesive, very well-sorted, sand-sized glass beads ($D_{50} = 0.213$ mm). All experiments used a volumetric sediment concentration of 15% and weight concentrations of EPS of 0%, 0.15%, and 0.3%. The sand experiments used additional EPS concentrations of 0.05% and 0.2%. Xanthan gum was used to represent natural EPS (cf., Malarkey et al., 2015; Craig et al., 2020). These sediment and xanthan gum concentrations were informed by the SGF experiments of Baker et al. (2017) and Craig et al. (2020).

156 Each starting suspension in the reservoir was prepared in the same way to account for any time-157 dependent behaviour, following procedures described by Baker et al. (2017) and Craig et al. (2020). 158 First, the xanthan gum and sediment were mixed in dry form in a concrete mixer for 10 minutes. The 159 seawater was then added to the dry material and mixed for another 10 minutes. Subsequently, the 160 wet slurry was decanted into a large bucket and mixed for another 3 minutes with a handheld mixer 161 to break up any remaining clumps of sediment. The slurry was then left to rest for 60 minutes, before 162 being mixed a third time for 3 minutes (except for the silt-water mixture which required 6 minutes of 163 mixing). Thereafter, a 180-ml subsample was taken from the slurry (except for 0.2% EPS sand) for 164 subsequent particle settling velocity analysis and the slurry was added to the reservoir in the lock-165 exchange tank and mixed for 30 seconds with the handheld mixer, before lifting the lock gate and thus 166 generating a bottom-hugging SGF. The moving head of each SGF was recorded using a high-definition 167 video camera, with the aim to document changes in head shape, internal flow structure, and head 168 velocity, following procedures described by Baker et al. (2017). Deposit thicknesses were recorded 169 along the sidewall of the tank for the SGFs that halted before reaching the end of the tank. Replicates 170 of three sand experiments showed that head velocities and run-out distances of the SGFs were 171 reproducible (Table 1).

The 180-ml subsamples were used to determine particle settling velocities of the 15% pure sediment
and mixed EPS-sediment suspensions. This involved timing the falling interface between clear water

and the settling suspension in the sampling container, after fully homogenising the starting
suspension, following the procedure described by Baas et al. (2022, their figure 1).

176 Head velocities of the SGFs were calculated from the video footage at a spatial resolution of 0.1 m along the tank using Microsoft Movie Maker. With the exception of the clay flows laden with 0.3% 177 EPS, each SGF could be subdivided into a period of constant head velocity followed by a period of 178 179 waning head velocity. The period of constant head velocity was characterised by the mean pre-180 deceleration head velocity between 23% and 64% of either the run-out distance (red segment of the 181 blue curve in Fig. 2) or 4.6 m for flows that reflected off the end of the tank (Fig. 2). For most SGFs that 182 stopped before reaching the end of the tank — referred to as the main flow below — a dilute cloud 183 bypassed the deposit of the main flow. This cloud, referred to as the bypassing flow below, continued 184 to move down the tank at a slower rate than the main flow. The velocity of these bypassing flows was 185 measured using the same method as for the main flow (orange curve in Fig. 2).

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187	4.	Resu	lts
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189 4.1. Settling velocity experiments

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191 *4.1.1. EPS-free sediment*

The EPS-free 15% kaolin clay suspension had a settling velocity at 0.0025 mm s⁻¹ (Fig. 3). This was two
orders of magnitude lower than the settling velocity of the EPS-free 15% coarse silt suspension. As
expected, the EPS-free 15% fine sand suspension had the highest settling velocity of 4.7 mm s⁻¹ (Fig. 3).
3).

196

197 4.1.2. EPS-sediment mixtures

198 All EPS-laden suspensions had a lower settling velocity than the equivalent EPS-free suspensions, and 199 the settling velocity decreased, as the EPS concentration was increased, for all sediment types (Fig. 3). 200 This decrease stretched over five orders of magnitude for the non-cohesive silt and sand suspensions, 201 but the decrease was confined to a factor of only 4.5 for the cohesive clay suspensions between 0% 202 and 0.30% EPS. Moreover, Fig. 3 shows that the greatest decrease in settling velocity for the silt 203 suspensions was between 0% and 0.15% EPS, whereas the decrease in the settling velocity for the 204 sand suspensions was relatively small between 0% and 0.05% EPS. At 0.30% EPS, the settling velocities 205 were within one order of magnitude for the three sediment types (Fig. 3). Hence, the particle settling 206 was less dependent on sediment type and size at 0.30% EPS than at lower EPS concentrations.

207

208 4.2. Visual observations and flow type classification

209

Five out of the eleven experimental runs conducted, i.e., all EPS-free flows and flows Sa-0.05 and Sa-0.15 (Table 1), were characterised by fully turbulent behaviour. The head of these flows had a pointed, semi-elliptical shape with a prominent nose in a vertical section parallel to the flow direction (Fig. 4A; Table 1). These flow properties match the low-density turbidity currents of Baker et al. (2017; their table 3) and the turbidity currents s.s. of Craig et al. (2020; their table 2).

215 Flows Cl-0.15, Si-0.3 and Sa-0.3 consisted of a dense lower layer, in which turbulence attenuation was 216 clearly visible, and a more dilute, lighter-coloured and fully turbulent, upper layer, characterised by 217 mixing with the ambient water (Fig. 4B; Table 1). In combination with hydroplaning at the base of the 218 head of these flows (Fig. 4B), the presence of linear coherent fluid entrainment structures (sensu Baker 219 et al., 2017; Figs. 4C, D) and a more rounded flow front than the low-density turbidity currents, these 220 flows resemble the high-density turbidity currents of Baker et al. (2017; their table 3) and the top-221 transitional plug flows of Hermidas et al. (2018) and Craig et al. (2020; their table 2). High-density 222 turbidity currents Cl-0.15, Si-0.3 and Sa-0.3 also included a dilute turbulent flow released from the

upper part of the main flow after it stopped moving (Figs. 4E, F). These bypassing flows moved slowlydown the tank.

Flows Si-0.15 and Sa-0.2 (Table 1) started as fully turbulent low-density turbidity current but they changed to stratified high-density turbidity current in the final phase of movement by vertical settling of sediment particles into the lower part of the flow. Flow Sa-0.2 exhibited coherent fluid entrainment structures, but fewer than in the high-density turbidity currents.

Finally, flow CI-0.3 slid out of the reservoir as a coherent mass of sediment; it was devoid of any internal turbulence (Figs. 4G, H; Table 1). This viscous flow, which showed little mixing at the top and lacked hydroplaning, is classified as a slide (Baker et al., 2017 [their table 3]; Craig et al., 2020 [their table 2]). Immediately after the slide had stopped moving, a dilute turbidity current formed at the top of the slide. This current bypassed the front of the slide and then moved slowly towards the downstream end of the tank.

235 The flow types defined in this paper are plotted as a function of EPS concentration and sediment type 236 in Fig. 5. This flow-type phase diagram reveals that the boundaries between low-density turbidity 237 currents, transitional low-density to high-density turbidity currents, high-density turbidity currents, 238 and slides move to higher EPS concentrations, as the sediment type changes from clay via coarse silt 239 to fine sand. Fig. 5 also shows a stability field for mud flows, in the case of suspended clay, or debris 240 flows, in the case of suspended silt or sand (sensu Baker et al., 2017), equivalent to the plug flows of 241 Hermidas et al. (2018) and Craig et al. (2020; their table 2). These flow types were not observed in the 242 present experiments, but they are expected in clay flows that contain between c. 0.2% and 0.25% EPS.

243

244 4.3. Flow velocity and run-out distance

245

246 4.3.1. EPS-free control experiments

The EPS-free clay and silt flows behaved in a similar manner hydrodynamically. Both flows had a mean pre-deceleration head velocity of 0.35 m s⁻¹ down to a distance, *x*, of 3 m along the tank (Figs. 6A, 7). The head velocity of these flows then gradually decreased to c. 0.2 m s⁻¹, before the flows reflected off the end of the tank (Figs. 6A, B, 8).

In contrast, the EPS-free sand flow decelerated earlier than flows Cl-0 and Si-0. After maintaining a mean pre-deceleration head velocity of 0.34 m s⁻¹ for 1.2 m, flow Sa-0 waned quickly, reaching a runout distance of 2.5 m (Figs. 6C, 7, 8). A weak bypassing flow formed after the main flow had stopped, but it lost forward momentum quickly and ran out at x = 3 m (Fig. 6C).

255

256 4.3.2. Mixed sediment–EPS experiments

257 The head velocity profiles and run-out distances of the EPS-laden flows varied with sediment type and 258 EPS concentration (Figs. 6–8). Except for slide Cl-0.3, which started to decelerate immediately after 259 leaving the reservoir, all flows were similar to the EPS-free control flows in that the head velocity was 260 constant for a certain distance along the tank, before the flow started to decelerate exponentially. 261 The differences in mean pre-deceleration head velocity were small, ranging from 0.32 m s⁻¹ to 0.37 m 262 s⁻¹ across all sediment types and EPS concentrations, again excluding slide Cl-0.3 (Fig. 7; Table 1). Yet, 263 Fig. 7 reveals that the mean pre-deceleration head velocity in the mixed clay–EPS and silt–EPS flows 264 decreased slightly, as the EPS concentration was increased and the flows changed from low-density 265 to high-density turbidity current in the silt runs and from low-density turbidity current to slide in the 266 clay runs. In contrast, the mean pre-deceleration head velocity in the mixed sand-EPS flows increased 267 from 0% to 0.05% EPS and decreased from 0.15% to 0.30% EPS (Fig. 7). These trends are mimicked by 268 changes in the location at which the flows started to decelerate, and, particularly, by changes in the 269 run-out distance of the flows (cf. Figs. 7, 8). The run-out distance of the mixed clay–EPS and silt–EPS 270 flows decreased, as the EPS concentration was increased, and the run-out distance in the mixed sand-271 EPS flows reached a maximum value of 4.4 m around 0.15% EPS (Fig. 8). Fig. 8 also shows that the runout distance of the clay flows decreased more rapidly with increasing EPS concentration than the runout distance of the silt flows and that the run-out distances of flows Si-0.20 and Si-0.30 merge with
those of flows Sa-0.20 and Sa-0.30 (Fig. 8). This merging of run-out distances thus appears to be
confined to the turbulence-attenuated turbidity currents carrying non-cohesive coarse silt and fine
sand.

277

- 278 4.4. Deposit properties
- 279

The shape and thickness distribution of the deposits formed by the experimental flows varied with 280 281 sediment type and EPS concentration, and therefore with flow type. The deposit of slide Cl-0.30 was 282 short and thick; it dipped steeply and almost uniformly from the point of entry into the tank (Fig. 9A). 283 The deposits of the low-density turbidity currents and the transitional low-density to high-density 284 turbidity currents tended to be wedge-shaped, with a gradual termination near the location of run-285 out (e.g. flows Sa-O and Sa-O.20; Fig 9B). The high-density turbidity current deposits had a more abrupt 286 termination, which is most apparent in the deposit of flow CI-0.15 (Fig. 9A). The EPS-free and low-EPS 287 silt and sand flows formed firm deposits, whereas adding large amounts of EPS to these flows resulted in soft, water-rich and gel-like beds. Such EPS-rich beds were significantly thicker than the EPS-free 288 289 and EPS-poor beds (e.g. Fig. 9B).

290

291 **5. Discussion**

292

293 5.1. Controlling factors on the mobility of the biologically cohesive flows

295 The amount of EPS added to the flows was two orders of magnitude lower than the sediment 296 concentration, i.e., 0–0.3% EPS versus 15% sediment. The density difference with the ambient water, 297 which drives the mobility of sediment gravity flows (Middleton and Hampton, 1973; Kneller and 298 Buckee, 2000), was therefore nearly constant and unable to explain the observed variations in flow 299 mobility and run-out distance. Other factors that controlled the flow mobility and run-out distance in 300 the experiments were: (1) drag-induced turbulent forces; (2) cohesive forces of physical and biological 301 origin, caused by kaolin clay and EPS, respectively; (3) particle settling velocity; and (4) hydroplaning 302 (e.g. Middleton and Hampton, 1973; Lowe, 1982; Mohrig et al., 1998; Baker et al., 2017; Craig et al., 303 2020). Upward-directed turbulent forces tend to promote flow mobility by keeping particles in suspension and therefore maintaining the density difference with the ambient water (Middleton and 304 305 Hampton, 1973). However, turbulent forces are counteracted by cohesive forces and particle settling, 306 which both tend to hinder flow mobility and promote deposition. The settling velocity of single 307 particles increases with increasing particle size, but it decreases with increasing particle concentration, 308 because of the hindered settling effect (e.g. Richardson and Zaki, 1954; Baas et al., 2022). However, 309 the effect of hindered settling on flow mobility is expected to have been small in the experiments, 310 because all flows carried the same volumetric concentration of sediment particles. Hydroplaning 311 reduces the drag at the base of dense SGFs and it therefore promotes flow mobility and run-out 312 distance (Mohrig et al., 1998). Below, the combined effect of turbulent forces, cohesive forces, particle 313 settling velocity, and hydroplaning is used to explain the differences in mobility between the 314 experimental flows.

315

316 5.2. Comparison of flow mobility: EPS-free control flows

317

All the EPS-free SGFs behaved as fully turbulent, low-density turbidity currents (*sensu* Baker et al. 2017). However, the mean pre-deceleration head velocity of flow Sa-0 was slightly lower than that of

320 flow Cl-0 and Si-0 (0.34 m s⁻¹ and 0.35 m s⁻¹, respectively; Table 1) and flow Sa-0 started to decelerate 321 c. 1.5 m earlier than the finer-grained flows (Fig. 6). These differences are inferred to reflect the higher 322 single particle settling velocity of the fine sand (Fig. 3), which worked against the upward-directed 323 particle support by turbulence and caused the flow Sa-0 to have a significantly shorter run-out 324 distance than flows Cl-0 and Si-0 (Fig. 8). The role of physical cohesion in flow Cl-0 was small, 325 considering that the non-cohesive silt flow and the cohesive clay flow had almost identical head 326 velocity profiles (Fig. 6). This agrees with Baker et al. (2017), who found that their 15% fine silt and 327 kaolin flows had similar mean head velocities.

328

329 5.3. Comparison of flow mobility: Biologically cohesive flows

330

331 Our hypothesis that biologically cohesive EPS reduce the flow mobility, i.e., the run-out distance (Fig. 332 8), is supported for the silt and clay flows. This reduction is also expressed by a lower mean pre-333 decelerating head velocity (Fig. 7). The biologically cohesive forces induced by the EPS attenuated the 334 turbulence and changed the low-density turbidity currents to high-density turbidity currents (sensu 335 Lowe, 1982), and eventually to slides (Mohrig and Marr, 2003) at the maximum EPS concentration in 336 the clay flows (Table 1). These changes in flow type were accompanied by a reduction in the run-out 337 distance through a change in the settling behaviour from single particle settling and hindered settling 338 to bulk, 'en-masse', settling of gel-like substances, in which the EPS form a volume-filling network of 339 bonds between the clay or silt particles (Craig et al., 2020). Flow transformation along the tank from 340 low-density to high-density turbidity current, i.e., from fully turbulent flow to top transitional plug 341 flow (Hermidas et al., 2018; Craig et al., 2020), signifies a downflow shift in the force balance from 342 dominantly turbulent to dominantly cohesive during the deceleration phase, resulting in bulk settling.

343 In contrast to the similarities in the mobility of the pure silt and clay flows, described above, the clay-344 EPS flows started to decelerate earlier and had shorter run-out distances than the silt–EPS flows (Fig. 345 6). These differences in the influence of biological cohesion may have two possible causes. Firstly, as 346 for clay concentration, EPS concentration has an exponential relationship with flow viscosity and yield 347 strength (e.g. Wan, 1982). Therefore, adding EPS to an already weakly cohesive clay flow may lead to 348 a larger reduction in flow mobility than adding the same amount of EPS to a non-cohesive silt flow, 349 even if this difference in physical cohesion is not reflected in a significant difference in the head 350 velocity profile between these EPS-free flows (Figs. 6A, B). This interpretation assumes that sediment 351 particles and EPS act independently in changing the cohesion of the flows. This assumption is not 352 necessarily valid, because different types of sediment and EPS may interact in yet unknown ways in 353 gravity flows. Further work is needed to explore such interactive processes. Secondly, the total particle 354 surface area in clay flows is larger than in silt flows, because of the larger number of particles for a 355 given concentration, which might increase the ease with which to establish particle bonds by EPS in 356 the clay flows and thus a stronger network of particle bonds (Craig et al., 2020).

357 The sand flows are more complex than our hypothesis advocates, because the run-out distance, in 358 particular, in the EPS-free sand flow was shorter than in the sand flows with low EPS concentrations. 359 This can be explained by a three-way coupling between turbulence, cohesion, and settling velocity. 360 EPS increase the viscosity of the water (Craig et al., 2020), which decreases the settling velocity of 361 sand particles and allows the sand to be kept in suspension for longer. In turn, this promotes the run-362 out distance. However, this is valid only for flows in which the EPS concentration is low, so that the 363 flow remains turbulent. At high EPS concentrations, the sand flow becomes transitional or laminar and 364 behaves in a similar way to the silt and clay flows, despite the even lower settling velocity of sand in 365 these EPS-rich flows. In other words, particle size is less important in controlling flow mobility in EPS-366 rich high-density turbidity currents, mud or debris flows, and slides. This is supported by the results of 367 the settling experiments, which show that the clay, silt, and sand flows had more similar settling 368 velocities at 0.3% EPS than at lower EPS concentrations (Fig. 3) and that the run-out distances of the

silt and sand flows merged at 0.2% and 0.3% EPS (Fig. 8). As in the silt-laden and clay-laden highdensity turbidity currents and slides, bulk settling of a cohesive gel with pervasively bonded sand particles explains why the EPS-rich sand-laden high-density turbidity currents became less mobile and had a shorter run-out distance (Fig. 8).

373 Hydroplaning was confined to high-density turbidity currents in the experiments. This phenomenon 374 occurs when the weight per unit area of material, here sediment and EPS, in the head of a flow is 375 exceeded by the dynamic pressure generated by the fluid just below the head (Mohrig et al., 1998; 376 Mohrig et al., 1999). Moreover, it is essential to prevent mixing of the overridden water with the flow 377 above, thus the permeability of the base of the flow has to be sufficiently low (Talling, 2013). This was 378 achieved fully in all the high-density turbidity currents, because of the high physical and biological 379 cohesive strength of these flows. Hydroplaning, however, did not occur in the low-density turbidity 380 currents, because the turbulence caused immediate mixing of water getting underneath the head of 381 these flows (cf., Baker et al., 2017). The base of the mixed clay-EPS slide probably was sufficiently 382 impermeable, but it did not allow hydroplaning either, because its low flow velocity prevented 383 ambient water from being forced underneath the head of the slide. Hydroplaning in the high-density 384 turbidity currents should have reduced the drag with the bed, but the shorter run-out distance of 385 these flows compared to the low-density turbidity currents suggest that the effect of increased 386 cohesion outweighed that of hydroplaning.

387

388 5.4. Comparison of deposits

389

The deposits of all the low-density turbidity currents that stopped before reaching the end of the lockexchange tank were thin and wedge-shaped (Fig. 9), which is typical for turbulent flows that comprise progressive single-particle settling and hindered settling (Baker et al., 2017). This includes the low-EPS

393 sand flows that were more mobile than the EPS-free sand flow, thus confirming that the EPS 394 concentration in these flows was sufficiently low to prevent the flow from turning into a cohesive gel. 395 The high-density turbidity currents and the slide, on the other hand, produced relatively thick deposits 396 with a more abrupt termination, typical for bulk settling of cohesive gels (Baker et al., 2017; Fig. 9). 397 The unusually large thickness of the EPS-rich deposits (Fig. 9) is testament to the ability of the EPS to 398 retain water in the cohesive gel and prevent the deposit from attaining a grain-supported texture 399 (Craig et al., 2020). As in physically cohesive clay beds (Mehta, 2013), such fluid-mud like deposits may 400 take many months to consolidate to firm deposits.

401

402 5.5. Implications for natural biologically cohesive flows and deposits

403

404 The present laboratory experiments show that EPS have a significant influence on the mobility and 405 run-out distance of cohesive and non-cohesive SGFs and on the shape of their deposits. As such, our 406 experiments extend the SGF experiments with mixed clay–EPS of Craig et al. (2020) to mixed silt–EPS 407 and sand-EPS SGFs, and provide evidence that particle size, in addition to turbulence and cohesion, 408 needs to be considered in predicting the behaviour of biologically cohesive SGFs (Fig. 10). The inferred 409 three-way coupling between particle settling velocity, turbulent forces and cohesive forces dictates 410 that EPS can increase the flow mobility, if EPS increase the flow viscosity whilst keeping the flow 411 turbulent, and decrease the flow mobility, if EPS attenuate the turbulence and change the flow to a 412 cohesive gel (Fig. 10). Within the parameter space of the experiments, mobility enhancement is 413 characteristic of sandy low-density turbidity currents, and mobility reduction is typical of high-density 414 currents, slides, and probably also mud flows and debris flows (Fig. 10).

These fundamental physical outcomes should also be relevant to natural flows. However, scaling of the experimental flows to natural prototypes is not possible at present, because standard scaling

417 methods do not allow for the inclusion of physical and biological cohesion in SGFs. Moreover, natural 418 SGFs are more complex than the experimental flows simulated here; for example, single-particle size 419 flows moving across a smooth, horizontal bed are an exception, rather than rule, in nature. In regard 420 to flow velocity, the experimental SGFs are suitable analogues for hyperpycnal flows at river mouths 421 and weak, single-surge, submarine SGFs triggered by earthquakes (Talling et al., 2013), sustained SGFs 422 with a frontal high-density basal layer (Zabala et al., 2017), and natural SGFs that have decelerated to 423 a similar velocity as in the experiments. It should also be mentioned that the experiments used a single 424 type of clay and EPS. The rheological properties of clay and EPS vary with chemical composition. For 425 example, kaolin clay has a lower cation exchange capacity, a key parameter for describing cohesive 426 properties (Yong et al., 2012), than illite and montmorillonite, which are also common in prototype 427 flows and deposits. Under given hydrodynamic conditions and clay concentration, SGFs laden with 428 illite or montmorillonite should therefore lose mobility at lower EPS concentrations than kaolinite-429 laden flows, and thus be more prone to change from low-density turbidity current to high-density 430 turbidity current, mud flows and slide.

431 The maximum EPS concentration of 0.3% used in the experiments was informed by the maximum EPS 432 concentration found in the seabed offshore New Zealand (Craig et al., 2020). Unless higher EPS 433 concentrations can be established elsewhere, $\leq 0.3\%$ EPS contained in faster flows should have a 434 smaller effect on flow mobility, since turbulent forces are positively correlated with flow velocity. 435 Some support is provided by the experiments of Craig et al. (2020), since their SGFs laden with 15% kaolin clay were faster than in the present study (0.42 m s⁻¹ versus 0.35 m s⁻¹), and the reduction in 436 437 run-out distance for EPS concentrations up to 0.265% was less than in the present experiments. The 438 exception may be the increased mobility of sandy SGFs with low concentrations of EPS, as the 439 behaviour of these flows is more dependent on the particle settling velocity than on turbulence 440 attenuation.

The shape of deposits of natural SGFs are expected to mimic that of the experimental SGFs. EPS generally cause SGF deposits to become shorter and thicker, except for deposits formed by highmobility SGF laden with sand and low concentrations of EPS. These differences in shape can be expressed at first order by a deposit steepness parameter, *S*:

445
$$S = \arctan\left(\frac{h_{max}}{X_R}\right),\tag{1}$$

446 where h_{max} is the maximum deposit thickness, X_R is the run-out distance, and S is given in degrees. The 447 deposit steepness parameter is plotted against EPS concentration in Fig. 11 for the deposits of all flows 448 that stopped before reaching the end of the tank. Short and thick deposits are expressed by relatively 449 high S values, whereas long and thin deposits have low S values. Fig. 11 shows that EPS-laden clay 450 flows form steeper deposits than sand and silt flows, because of their lower mobility induced by 451 combined physical and biological cohesion. Natural deposits of EPS-free sand flows are expected to 452 have higher S values than deposits of low-EPS sand flows, here for 0.05–0.2% EPS, because of their 453 low viscosity and high sand settling velocity. At high EPS concentrations, here for 0.2% and 0.3% EPS, 454 the deposit steepness increases with increasing EPS concentration for SGFs laden with clay or sand 455 (Fig. 11), and presumably also for silt-laden flows, because of reduced flow mobility caused by greater 456 biological cohesion. The steepness of deposits of low-EPS sand flows, here for 0.05-0.2% EPS, is 457 proportional to EPS concentration on a logarithmic scale (Fig. 11). Despite the observed increase in 458 mobility of these flows, which would cause the deposits to become longer and thinner and thus S to 459 become lower, anomalously high increases in deposit thickness lead the increase in the steepness 460 parameter, because of the above-mentioned development of soft, water-rich and gel-like beds. 461 However, such beds may be unstable and bed consolidation by water expulsion may lead to a rapid 462 decrease in deposit thickness, and therefore deposit steepness, under natural conditions. For natural 463 hydrodynamic conditions comparable to those simulated in the experiments, the S values of deposits 464 formed by high-EPS sand-laden and silt-laden SGFs, here for 0.3% EPS, are expected to be similar (Fig.

465 11). This supports the above interpretation that the size of non-cohesive particles becomes less466 important as the EPS concentration increases.

467 Because SGFs in nature can have run-out distances of hundreds to even thousands of kilometres, the results of this study should be included in the mitigation of possible damage to subaqueous 468 469 infrastructure that is in the path of SGFs. If the sediment in the source area of these SGFs, as well as 470 the sediment eroded from the bed during transport, is dominated by silt or clay, the presence of EPS 471 in these flows can be expected to reduce the risk of damage to, for example, subaqueous pylons and 472 communication cables. Although turbulent sand-rich SGFs generally are less mobile than turbulent 473 silt-rich and clay-rich SGFs, and therefore, less likely to reach subaqueous infrastructure from the same 474 source area, the enhanced mobility of sand flows carrying small amount of EPS may lead to a false 475 sense of security, because such flows may travel further than expected.

476

477 **6. Conclusions**

478

479 The laboratory experiments presented in this paper reveal that the mobility of biologically cohesive 480 sediment gravity flows and the shape of their deposits depend on the type and size of suspended 481 particles, i.e., physically cohesive clay versus non-cohesive silt and sand, and the concentration of 482 extracellular polymeric substances. Upon the addition of EPS at concentrations typical of deep-marine 483 environments, the induced biological cohesion causes fine-grained EPS-free flows to become 484 turbulence-attenuated, leading to a reduction in flow mobility. This reduction in mobility is reflected 485 in the progressive shortening of the run-out distance of the flows, as the EPS concentration is 486 increased. In contrast to clay and silt flows, the mobility of sand flows is affected not only by the 487 balance between turbulent and cohesive forces, but also by the high settling velocity of the sand 488 particles. The high settling velocity causes a low flow mobility under EPS-free conditions, reflected in 489 a short run-out distance. However, the presence of EPS in sand flows induces a considerable increase

490	in run-out distance by increasing the flow viscosity, provided that the EPS concentration is sufficiently
491	low to prevent turbulence attenuation in these flows. Above a certain threshold EPS concentration,
492	sand flows become less mobile and behave in a similar way to silt flows and clay flows, because the
493	formation of a biologically cohesive gel renders the settling velocity of subordinate importance. Since
494	EPS has been shown to be pervasive in depositional environments where sediment gravity flows are
495	expected to occur and where engineering infrastructure is being constructed at an increasing rate,
496	these results need to be incorporated in models that aim to mitigate damage to such infrastructure.
497	
498	All experimental data are available on FigShare via DOI 10.6084/m9.figshare.19960151 and URL
499	https://figshare.com/articles/dataset/Data_tables_for_Figs_3_6_9_docx/19960151.
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677 Figure captions

Table 1. Summary of experimental data. LDTC = low-density turbidity current; HDTC = high-density
 turbidity current, LDTC>HDTC = low-density changing to high-density turbidity current near location
 of run-out.

Fig. 1. Top: Lock exchange tank. Bottom: Schematic diagram of the laboratory set up from the side.
HD = high-definition.

Fig. 2. Procedure used to calculate the mean pre-deceleration head velocity. The red curve delimits the 23% and 64% locations along the flow path, on which the mean pre-deceleration head velocity calculations are based, where 0% is the gate location and 100% is the run-out distance of the flow or the end of the tank for flows that reflected off the end wall.

Fig. 3. Particle settling velocity against EPS concentration for the kaolin clay, coarse silt and fine sandused in the lock-exchange experiments.

689 Fig. 4. Typical examples of the experimental sediment gravity flows. A) Head of low-density turbidity 690 current Si-O, characterised by pervasive turbulence, a pointed flow front, and lack of hydroplaning 691 (arrow). B) Head of high-density turbidity current Cl-0.15, characterised by pronounced hydroplaning 692 at the base (arrow) and a dense, turbulence-attenuated layer overlain by a more dilute, turbulent 693 layer; note the distinct density interface (dashed line). C, D) examples of coherent fluid entrainment 694 structures (arrows) in flows Si-0.15 and Si-0.3, respectively. E, F) examples of bypassing flow in high-695 density turbidity currents Cl-0.15 and Si-0.3. G, H) Slide Cl-0.3 after moving for 0.12 m and 0.74 m in 696 the tank, respectively. Note the lack of hydroplaning, the viscous flow character, and the minor mixing 697 with the ambient fluid. Flow is from left to right in all pictures. Scale at bottom of images is in centimetres. 698

Fig. 5. Flow-type phase diagram for different EPS concentrations and sediment types. LDTC = low density turbidity current; HDTC = high-density turbidity current, LDTC>HDTC = low-density changing

701 to high-density turbidity current near location of run-out; MF/DF = mud or debris flow. Dashed lines 702 denote estimated boundaries between flow types for the clay flows. The question mark refers to a 703 mud and debris flow phase that was observed in previous work (e.g. Baker et al., 2017; Craig et al., 704 2020), but that was not captured within the limited resolution of EPS concentrations in the present 705 experiments. Note that 18.5%, 26%, and 28% clay, instead of EPS, would have to be added to the clay 706 flows to cross the boundaries between LDTC, HDTC, MF, and slide, respectively (Baker et al., 2017). 707 These clay concentrations are two orders of magnitude higher than the boundary EPS concentrations, 708 demonstrating the strong cohesive properties of EPS.

Fig. 6. Head velocity against distance along tank for all flows laden with (A) kaolin clay, (B) coarse silt,
and (C) fine sand, with and without EPS. The head velocities of the bypassing part of flows are given
in orange.

Fig. 7. Mean pre-deceleration head velocity and standard deviation of the mean (vertical lines) against
EPS concentration for kaolin clay, coarse silt, and fine sand. LDTC = low-density turbidity current; HDTC
a high-density turbidity current.

Fig. 8. Run-out distance against EPS concentration for kaolin clay, coarse silt and fine sand flows.
Dashed blue line and green data point at 0% EPS denote minimum run-out distances, limited by the
4.6-m length of the lock-exchange tank. Note that the plotted run-out distances exclude the bypassing
part of the flows.

Fig. 9. Deposit thickness trends for all flows that did not reflect off the end of the lock-exchange tank.
A) Kaolin clay and coarse silt deposits. B) Fine sand deposits.

Fig. 10. Conceptual model summarising the effect of key parameters on flow mobility (top) and the dependence of type and run-out distance of sediment gravity flows on these parameters, informed by the laboratory experiments (bottom). The column labelled 'Dominant process' provides the dominant physical or biological controls on run-out distance.

- **Fig. 11.** Deposit steepness against EPS concentration for all flows that did not reflect off the end of
- the lock-exchange tank.

























<u>±</u>

- **Table 1**. Summary of experimental data. LDTC = low-density turbidity current; HDTC = high-density
- 2 turbidity current, LDTC>HDTC = low-density changing to high-density turbidity current near location
- 3 of run-out.

Experimental run	Sediment type	EPS concentration (%)	Pre-deceleration velocity (m s ⁻¹)	Froude number*	Run-out distance (m)	Flow type
CI-0	Clay	0	0.351	0.66	>4.6**	LDTC
Cl-0.15	Clay	0.15	0.329	0.62	3.2	HDTC
Cl-0.3	Clay	0.3	0.226	0.82	0.9	Slide
Si-0	Coarse silt	0	0.353	0.70	> 4.6**	LDTC
Si-0.15	Coarse silt	0.15	0.346	0.65	> 4.6**	LDTC> HDTC
Si-0.3	Coarse silt	0.3	0.338	0.64	4.1	HDTC
Sa-0	Fine sand	0	0.340	0.75	2.5 (2.6)***	LDTC
Sa-0.05	Fine sand	0.05	0.368	0.71	3.8	LDTC
Sa-0.15	Fine sand	0.15	0.362	0.71	4.4 (4.6)***	LDTC
Sa-0.2	Fine sand	0.2	0.357	0.74	4.4	LDTC> HDTC
Sa-0.3	Fine sand	0.3	0.324	0.66	4 (4)***	HDTC

4 * Densimetric Froude number was calculated from flow thickness at x = 1.5 m (except for Cl-03: x =

5 0.5 m), mean pre-deceleration head velocity, and initial flow density.

6 ** Flow reached the end of lock-exchange tank.

7 *** Run-out distances between brackets are based on replicate experiments.

Declaration of interests

⊠The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

□The authors declare the following financial interests/personal relationships which may be considered as potential competing interests: