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## Coupling Evidence From Lower Atmosphere to Mesosphere and Ionosphere Through Quasi 27-Day Oscillation

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2 **Coupling evidence from lower atmosphere to**  
3 **mesosphere and ionosphere through quasi**  
4 **27-day oscillation**  
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21       **Abstract.** Using meteor radar, radiosonde and digisonde observations and MERRA-2 reanalysis data  
22 from 12 August to 31 October 2006, we report a dynamical coupling from the tropical lower  
23 atmosphere to the mesosphere and ionospheric F2 region through a quasi 27-day intraseasonal  
24 oscillation (ISO). It is interesting that the quasi 27-day ISO is active in the troposphere and stratopause  
25 and mesopause regions, exhibiting a three-layer structure. In the mesosphere and lower thermosphere  
26 (MLT), the amplitude in the zonal wind increases from about  $4 \text{ ms}^{-1}$  at 90 km to  $15 \text{ ms}^{-1}$  at 100 km,  
27 which is different from previous observations that ISOs generally have the amplitude peak at about  
28 80-85 km, and then weakens with height. OLR and specific humidity data demonstrate that there is a  
29 quasi 27-day periodicity in convective activity in the tropics, which causes the ISO of the zonal wind  
30 and gravity wave (GW) activity in the troposphere. GW energy in the stratosphere also exhibits a sharp  
31 spectral speak at 27-day period, meaning that the convectively modulated GWs play a vital role in  
32 driving the oscillation in the MLT. The quasi 27-day variability arises clearly in the hmF2. Wavelet  
33 analysis shows that the dominant period and active time of the hmF2 oscillation are in good agreement  
34 with those in the zonal wind of the MLT and OLR rather than in the F10.7 and Kp index. Hence,  
35 tropical convective activity has an influence on the dynamics of the MLT and F2 region through  
36 modulated waves and ISOs.

37

## 38 **1. Introduction**

39       The circulation of the tropical stratosphere, mesosphere and lower thermosphere (MLT) is  
40 characterized by quasi-biennial (QBO), annual (AO), and semiannual (SAO) oscillations (Baldwin et al.,  
41 2001). In the tropics, convective activity is active due to intense solar radiation, which can excite  
42 planetary scale Kelvin and Rossby waves, and mid and small scale gravity waves (GWs). It is generally  
43 accepted that waves with different scales generated in the lower atmosphere play a vital role in driving

44 QBO and SAO at different heights through their upward propagation and interaction with the  
45 background flow (Lindzen, 1981; Dunkerton, 1997). Hence, oscillation and wave propagation and their  
46 intercoupling dominate the dynamical process in the tropical atmosphere.

47 Oscillation with seasonal time scale was noticed slightly later than the well-known QBO (Ebdon,  
48 1960; Reed et al., 1961). Madden and Julian (1971, 1972) discovered a 40 to 50-day oscillation of the  
49 zonal wind in the tropical troposphere, which is referred to as the Madden-Julian oscillation (MJO).  
50 Afterwards, as the advancement of atmospheric sounding, more periodic oscillations were found in the  
51 tropical zonal wind and temperature, thus intraseasonal oscillation (ISO) is often extended to be  
52 quasi-periodic variations of about 20-100 days (Ghil & Mo, 1991; Eckermann et al., 1997; Isoda et al.,  
53 2004). A number of observations reveal that ISOs are a prominent variability in the tropical atmosphere  
54 (Madden, 1986; Gutzler & Madden, 1989; Ghil & Mo, 1991; Hendon & Salby, 1994; Eckermann et al.,  
55 1997; Lieberman, 1998; Kumar et al., 2007; Deshpande et al., 2011; Niranjankumar et al., 2011;  
56 Guharay et al., 2017; Li et al., 2018). Until now, ISOs have attracted extensive attention because of not  
57 only their complex dynamics but also their important influences on monsoons (Lawrence & Webster,  
58 2001; Zhou & Chan, 2005; Karmakar & Krishnamurti, 2019), tropical cyclone (Huang et al., 2011;  
59 Yang et al., 2015), and cloud and precipitation (Benedict & Randall, 2007; Lau & Wu, 2010; Barnes &  
60 Houze Jr, 2013).

61 Theoretical, modeling and observational studies paid great efforts with regards to spatial structures,  
62 propagation patterns and generation mechanisms of ISOs. ISOs can propagate either eastward, showing  
63 Kelvin wave features, or westward, displaying Rossby wave characteristics, which is explained in terms  
64 of an equatorial convectively coupled Kelvin-Rossby couplet (Lau & Peng, 1990; Wang & Xie, 1996,  
65 1997). Interactions between eastward propagating ISOs and mean flow can also generate an unstable  
66 Rossby wave with horizontal scale of thousands of kilometers (Lau & Peng, 1990). Easterly vertical

67 shears (Jiang et al., 2004), air-sea interaction (Fu & Wang, 2004), eddy momentum transport from  
68 synoptic systems (Hsu & Li, 2011) could contribute to the northward propagation component of ISOs,  
69 and beta drift of synoptic motion may induce the northward propagation of ISOs (Boos & Kuang, 2010).  
70 Although tropical ISOs are mainly originated from convectively coupled wave dynamics involving  
71 Kelvin and Rossby waves through latent heat released in precipitation, atmospheric response to  
72 independent forcing, atmospheric instability, water vapor variation, nonlinear heating from sea surface  
73 temperature, multiscale interaction, and solar irradiance are proposed to be possible mechanisms of ISO  
74 generation (Madden & Julian, 1994; Zhang, 2005).

75 When ISO generated in the tropical lower atmosphere propagates upward, its amplitude often  
76 quickly decays above the tropopause (Madden & Julian, 1971; Rao et al., 2009; Guharay et al., 2014).  
77 However, some studies show that ISO may strengthen again in the upper stratosphere. Kumar and Jain  
78 (1994) proposed that ISO could propagate upward to the stratosphere from the tropical troposphere via  
79 leakage of its partial energy into the stratosphere. Ziemke and Stanford (1991) showed that strong ISO  
80 activity associated with tropical Rossby waves in the lower atmosphere of the Southern Hemisphere  
81 (SH) could propagate to the upper stratosphere by refracting out to mid latitudes and then refracting  
82 back to the equatorial stratosphere region, and similar phenomenon in the Northern Hemisphere (NH)  
83 was inferred based on rocketsonde and radiosonde observations (Nagpal et al., 1994). Niranjankumar et  
84 al. (2011) suggested that partial refraction and reflection in the troposphere are related with subtropical  
85 jet when investigating vertical and lateral propagation features of ISO. Numerical investigation  
86 demonstrated that Rossby waves forced by the induced heating on the equator could radiate poleward  
87 into the extratropical westerlies and vertically into the stratosphere in the two hemispheres (Salby et al.,  
88 1994).

89 In the MLT, ISOs with a wide period range of 20-60 days were reported in the zonal wind over

90 Christmas Island (2°N, 157°W) based on medium frequency (MF) radar observation (Eckermann &  
91 Vincent, 1994). The study indicated that it is unlikely that the ISOs propagate up to the MLT from the  
92 lower atmosphere although they were supposed to originate from in the lower atmosphere. A further  
93 investigation by Eckermann et al. (1997) found that ISO patterns in the MLT occur not only in zonal  
94 wind but also in GW variances and diurnal tidal amplitudes, thus they suggested that GWs and tides  
95 modulated by ISOs in tropical convection could induce in turn similar periodicities in the zonal flow of  
96 the MLT by transferring their momentum and energy into the mean flow as these waves propagate up to  
97 the MLT. By analyzing the correlations between ISO of the zonal wind and wave amplitudes in the  
98 equatorial MLT from three radar observations, Isoda et al. (2004) argued that nonmigrating tides are  
99 modulated at ISO period in the lower atmosphere, and by propagating to the MLT, the breaking and  
100 dissipating tides drive the variation of the zonal mean wind, but the contribution of GWs to the ISO in  
101 the MLT is unclear. There is a consistency of ISOs between convective activity in the lower atmosphere  
102 and the zonal wind in the MLT based on MF and meteor radar observations at low latitudes, implying  
103 similar ISO of the zonal wind in the MLT driven by modulated wave forcing (Kumar et al., 2007; Rao  
104 et al., 2009). Lower stratospheric GW potential energy derived from Constellation Observing System  
105 for Meteorology, Ionosphere, and Climate (COSMIC) radio occultation satellite measurement shows a  
106 strong correlation with the tropospheric zonal winds due to filtering of upward propagating GWs by the  
107 MJO wind (Moss et al., 2016). Middle stratospheric GW temperature variances observed by  
108 Atmospheric Infrared Sounder (AIRS) display an intraseasonal variability because the zonal wind  
109 around the tropopause likely controls GW propagation into the stratosphere by a critical level filtering,  
110 meaning that MJO can modulate the middle atmospheric circulation through regulating GWs by means  
111 of generation and propagation (Tsuchiya et al., 2016). However, using meteor radar data at mid and high  
112 latitudes, Pancheva et al. (2003) noticed an obvious ISO in the MLT, but there was not a corresponding

113 periodicity in atmospheric wave activities, including GWs, tides and quasi 2-day PWs. In addition,  
114 Huang et al. (2015) presented that an ISO with relatively short period of 27 days could penetrate  
115 tropopause region and weak westward wind field in the lower stratosphere to propagate upward into the  
116 MLT, while an ISO with long period of about 46 days does not so.

117 Quasi 27-day ISO arouses attention because this variability has the same period as solar rotation.  
118 Quasi 27-day periodicity is often observed in the zonal wind, temperature, and trace gases at different  
119 heights from the troposphere to the MLT (Fioletov, 2009; Huang et al., 2015; Hood, 2016; Guharay et  
120 al., 2017; Thiéblemont et al., 2018). Although trace gases have a clear response to solar 27-day cycle,  
121 the attribution of 27-day periodicity in the atmosphere to a solar cause is complicated by the fact that  
122 internal variability of the atmosphere can itself also generate quasi 27-day oscillation (Hoffmann & von  
123 Savigny, 2019). Both chemistry-climate model (CCM) and whole atmosphere community climate  
124 model (WACCM) results show that a 27-day variability in the wind and temperature is an inherent  
125 feature of the atmosphere, thus is not necessarily related to the solar rotational period since CCM and  
126 WACCM can output quasi 27-day variation even without solar rotational forcing (Schanz et al., 2016;  
127 Sukhodolov et al., 2017). Similar to the response of ozone in the upper stratosphere to solar rotation  
128 through physicochemical mechanisms, the presence of quasi 27-day oscillation in the ionospheric  
129 variability is a natural event as solar radiation is a major source of energy and ionization. The effects of  
130 solar rotation on ionospheric key parameters, such as total electron content (TEC), peak electron density  
131 of F2 layer (NmF2), maximum height of F2 layer (hmF2), and critical frequency of F2 layer (foF2), are  
132 extensively investigated by observational and modeling studies (Pancheva et al., 1991; Rich et al., 2003;  
133 Min et al., 2009; Xu et al., 2011; Coley & Heelis, 2012; Ma et al., 2012; Ren et al., 2018).

134 Nevertheless, many investigations show that long-period solar flux changes of 11-year solar cycle  
135 are the major source of ionospheric variability, and compared to long-term solar changes, 27-day solar



136 rotation and annual and semiannual variations associated solar zenith angle represent small variability  
137 (Forbes et al., 2000; Richards, 2001; Liu et al., 2006; Solomon et al., 2018). The ionospheric variability  
138 with periods between 2 and 30 days is mainly caused by geomagnetic activity and equally important  
139 "meteorological influences" transmitted from lower levels (Forbes et al., 2000, 2018; Rishbeth &  
140 Mendillo, 2001; Laštovička, 2006; Borries & Hoffmann, 2010). Mikhailov et al. (2004) analyzed  
141 ionospheric F2-layer disturbances not related to geomagnetic activity based on records over 26  
142 ionosonde stations. Ionospheric disturbances in quiet time of the Sun are not only rather frequent but  
143 their amplitude is comparable to the amplitude of moderate F2-layer storm effects, and these  
144 perturbations may partly be attributed to the impact from below. By using hmF2 data from digisonde  
145 measurement at Millstone Hill, Pancheva et al. (2002) demonstrated that the 27-day oscillations in  
146 hmF2 are generated by geomagnetic activity and by 27-day oscillation present in the zonal neutral wind  
147 of the MLT. Hence, the dynamics in the MLT play a significant role in ionospheric variability.

148 In this work, a quasi 27-day ISO activity in the tropics is reported. The oscillation exhibits a  
149 characteristic variation with height in the MLT, being different from previous observations, and leads to  
150 a corresponding change in the ionosphere. We attempt to reveal its generation in the lower atmosphere  
151 and the coupling from the troposphere to the ionosphere through this oscillation. Section 2 covers a  
152 brief explanation of the data that we utilized. In section 3, we provide the quasi 27-day variability  
153 throughout the troposphere to the thermosphere. Section 4 discuss the origination of the oscillation in  
154 different atmospheric layers, and a summary is presented in Section 5.

155

## 156 **2. Data**

### 157 **2.1 Meteor Radar Observation**

158 The horizontal wind from a meteor radar located in Kihei on Maui, Hawaii, at 20.75°N, 156°W, is

159 used in the study. The radar system is an all-sky interferometric meteor (SKiYMET) radar with an  
160 operating frequency of 40.92 MHz, which adopts a 3-element Yagi antenna pointing zenith to illuminate  
161 meteor trails. Five 3-element Yagi antennas oriented along two orthogonal baselines are used to receive  
162 echoes from the meteor trails. The hourly horizontal wind at 1 km height interval in the height range of  
163 80-100 km is estimated based on meteor trail position and Doppler shift. A technical description of  
164 SKiYMET radar can be found in the work of Hocking et al. (2001). The Maui meteor radar system and  
165 the wind computation are described in detail by Franke et al. (2005). The meteor radar observation has  
166 been applied to exploring dynamical processes in the MLT over Maui (Lu et al., 2011; Liu et al., 2013;  
167 Huang et al., 2013a, 2013b).

168 In the paper, the horizontal wind data for 81 days from 12 August to 31 October 2006 is utilized. In  
169 this period of our attention, the missing data decreases to a minimum value of 1.3% at 90 km from  
170 about 12.9% at 80 km and 15.6% at 100 km, thus the highly effective observation is suitable for the  
171 investigation of the atmospheric oscillation in the MLT region.

## 172 **2.2 Digisonde Observation**

173 We use the hmF2 to examine the ionospheric response to the atmospheric oscillation. The hmF2  
174 data is obtained from the digital ionosonde database (DIDBase) of University of Massachusetts Lowell  
175 (Galkin et al., 1999; Reinisch et al., 2004), which is accessed from the website at  
176 <http://www.digisonde.com>. The hmF2 is observed from the DPS-4 digisonde installed in Kwajalein  
177 Island (9.4°N, 167.4°E), with a temporal resolution of 5 min.

## 178 **2.3 Radiosonde Observation**

179 The United States radiosonde observations at three tropical stations distributed in the NH and SH are  
180 applied to the analysis of the oscillation activity in the troposphere and lower stratosphere. The data are  
181 archived and provided freely by the National Climatic Data Center (NCDC) of National Oceanic and

182 Atmospheric Administration (NOAA) through the Stratospheric Processes and Their Role in Climate  
183 (SPARC) Data Center at <ftp://ftp.ncdc.noaa.gov>. The three stations are situated at Hilo (19.72°N,  
184 155.07°W), Kauai (21.98°N, 159.35°W) on Hawaii, and at Pago Pago (14.33°S, 170.72°W) on Samoa.  
185 Observational sites are sparse in the Pacific, thus we choose the radiosonde and digisonde stations as  
186 close as possible to the meteor radar station.

187 Routine radiosondes are usually launched twice daily at 00:00 and 12:00 UT. As a balloon rises,  
188 atmospheric horizontal wind, temperature, pressure, and relative humidity are sensed by balloon-borne  
189 platform. The sampled heights depend on ascent rate of balloon, ranging from 10 to 100 m. For  
190 convenience, we interpolate the raw data linearly to a uniform interval of 50 m. The maximum height of  
191 a radiosonde observation is its balloon burst altitude. In the period that we focus on, about 90%, 88%  
192 and 90% of balloons reached 30 km in Hilo, Kauai and Pago Pago, respectively, but only about 63%,  
193 87% and 87% attained 31 km. Hence, we select the altitude of 30 km as the upper height limit of  
194 radiosonde observation in our analysis.

195 We follow the method proposed by Alduchov and Eskridge (1996) to derive specific humidity from  
196 temperature, pressure and relative humidity measured by radiosonde.

## 197 **2.4 Outgoing Longwave Radiation**

198 Specific humidity and outgoing longwave radiation (OLR) have often been used as proxies of  
199 convective activity over the tropical region (Arkin & Ardanuy, 1989). We expect to explore the origin of  
200 the tropical oscillation by combining specific humidity and OLR. Daily OLR data is obtained from the  
201 NOAA at the Website of <https://www.esrl.noaa.gov/psd>, with a 2.5° latitudinal and longitudinal  
202 resolution (Liebmann & Smith, 1996).

## 203 **2.5 Solar and Geomagnetic indices**

204 We also examined solar and geomagnetic forcing as potential controlling factors of the atmosphere.

205 Daily solar 10.7 cm radio flux (F10.7) and 3-h Kp index data provided from the National Centers for  
206 Environmental Information (NCEI) of NOAA at the Website of <https://www.ngdc.noaa.gov> are used as  
207 measures of solar and geomagnetic activities. The measurements were taken during the late declining  
208 phase of solar cycle 23.

## 209 **2.6 Reanalysis Data**

210 The Modern-Era Retrospective Analysis for Research and Applications (MERRA) reanalysis of the  
211 National Aeronautics and Space Administration (NASA) is an ideal candidate to examine the features of  
212 the oscillation in the zonal, meridional and vertical directions. The product of “inst6\_3d\_ana\_Nv” in the  
213 version 2 of MERRA (MERRA-2) is available through the NASA Goddard Earth Sciences Data and  
214 Information Services Center (GES DISC) online archive at <https://disc.gsfc.nasa.gov/datasets>. The  
215 reanalysis data is 6-hourly instantaneous analysis fields on a  $0.5^\circ \times 0.625^\circ$  latitude-by-longitude grid at  
216 72 model levels from ground up to 0.01 hPa (Gelaro et al., 2017).

217 All the data used in the present study are in the same period from 12 August to 31 October 2006,  
218 and 12 November 2006 is referred to as day 1.

219

## 220 **3. Quasi 27-Day Oscillation**

### 221 **3.1 Oscillation in MLT**

222 Figure 1 shows the daily averaged zonal and meridional winds measured by the meteor radar for 81  
223 days from 12 August to 31 October 2006. The mean zonal wind (positive eastward) has a maximum of  
224  $50.7 \text{ ms}^{-1}$  at 95 km on day 59 and a minimum of  $-67.0 \text{ ms}^{-1}$  at 100 km on day 24, while the mean  
225 meridional wind (positive northward) is between  $-59.3$  and  $40.0 \text{ ms}^{-1}$ . The different temporal scales of  
226 the perturbations can be noted in the wind field.

227 We carry out a Lomb-Scargle spectrum analysis (Scargle, 1982), with a 4-times oversampling, on

228 the daily averaged zonal and meridional winds to investigate their spectral components. Figure 2 shows  
229 the Lomb-Scargle spectral of the mean zonal and meridional winds. A confidence level of 95%  
230 corresponds to a spectral amplitude of  $7.4 \text{ ms}^{-1}$ . It is interesting that the spectral components in the  
231 zonal wind are very weak below 90 km, while above 90 km, a quasi 27-day oscillation is a predominant  
232 component, and strengthens gradually with height. Its spectral amplitude grows to  $15.2 \text{ ms}^{-1}$  at 100 km  
233 from  $4.2 \text{ ms}^{-1}$  at 90 km. Based on over five years of wind data acquired by a MF radar at Christmas  
234 Island, Eckermann et al. (1997) showed that the  $\sim 60$ -day ISO in the zonal wind of the MLT is relatively  
235 strong in winter-early spring (December-April). Its amplitude reaches a peak of about  $10\text{-}15 \text{ ms}^{-1}$  at  
236 80-85 km, and then clearly decreases with height. Similarly, by using a 20-year radar observation at  
237 four sites from  $70^\circ\text{N}$  to  $30^\circ\text{N}$ , the investigation indicated that the oscillation in the zonal wind with  
238 period between 20 and 40 days at mid and high latitudes is active in winter (November-March), with a  
239 strength reduction from about  $10 \text{ ms}^{-1}$  at 75-85 km to about  $5 \text{ ms}^{-1}$  near 100 km (Luo et al., 2001).  
240 Huang et al. (2015) presented a quasi 27-day oscillation event in December 2004 to March 2005 over  
241 Maui with large amplitudes throughout the MLT. Its amplitude has a maximum value of  $20.3 \text{ ms}^{-1}$  in the  
242 zonal wind at 89 km, and then monotonously drops to  $7.5 \text{ ms}^{-1}$  at 96 km. Hence, in the paper, we note  
243 that a relatively intense quasi 27-day oscillation occurs in August-October, and its evolution with height  
244 in the MLT is different from those in the previous studies (Eckermann et al., 1997; Luo et al., 2001;  
245 Huang et al., 2015). The quasi 27-day periodicity also arises in the meridional wind and shows a similar  
246 vertical variation but with a weak spectral magnitude. This is consistent with early observations that  
247 atmospheric oscillation occurs mainly in the zonal wind (Eckermann et al., 1997; Luo et al., 2001;  
248 Huang et al., 2015), thus we will concentrate on the zonal wind oscillation.

249 In order to determine the vertical propagation of the oscillation, a sinusoidal wave fitting under a  
250 27-day period is applied to the time series of the mean zonal wind at 90-100 km. The fitted amplitude

251 and phase are shown in Figure 3. Here, the phase is described by the time when the oscillation attains its  
252 maximum value (Huang et al., 2015). It can be seen from Figure 3 that the amplitude of the oscillation  
253 gradually increases from  $4.3 \text{ ms}^{-1}$  at 90 km to  $15.2 \text{ ms}^{-1}$  at 100 km, which is in good agreement with the  
254 spectral result in Figure 2. A linear fitting of the phase exhibits a slow downward phase progression in  
255 the MLT.

### 256 **3.2 Oscillation in Ionosphere**

257 It is generally recognized that the 27-day periodicity in the ionosphere can originate from the effects  
258 of the solar rotation, geomagnetic activity and zonal neutral wind in the MLT (Pancheva et al., 2002;  
259 Liu et al., 2006; Borries & Hoffmann, 2010; Ma et al., 2012). Here, we use the hmF2 data over the  
260 tropical Kwajalein Island to examine its variation during this time. A wavelet transform is performed on  
261 the hourly mean hmF2 time series. Morlet wavelet function which consists of a plane wave modulated  
262 by a Gaussian envelope is chosen as mother wavelet. Wavelet analysis decomposes a time series into a  
263 two-dimensional time-frequency domain, thus it can provide not only the dominant components but  
264 also the variation of these components with time.

265 Figure 4 shows the wavelet spectrum of the hourly mean hmF2 from 12 August to 31 October 2006.  
266 The perturbation periods are limited to 15-45 days to highlight the quasi 27-day periodicity. For  
267 comparison, the wavelet transform for the linearly interpolated zonal wind at 100 km above Maui is  
268 also presented in Figure 4. Figure 4 illustrates that the quasi 27-day oscillation takes place in the hmF2,  
269 and the onset of the hmF2 oscillation lags behind the oscillation in the zonal wind at 100 km. Similarly,  
270 Pancheva et al. (2002) found that the quasi 27-day variation in the hmF2 above Millstone Hill can be  
271 generated by the quasi 27-day oscillation in the zonal neutral wind of the MLT but with 5-6 day delay.  
272 Some observational and modeling studies indicated that most GWs and tides could propagate into the  
273 ionosphere directly from the MLT, and oscillations at periods of 2-30 days might propagate upward to

274 the F-region indirectly via the potential mechanism of modulation by upward propagating tides  
275 (Laštovička, 2006; Pancheva et al., 2006; Borries & Hoffmann, 2010; Forbes et al., 2018). Hence, the  
276 quasi 27-day oscillation can lead to the dynamical and electrodynamical coupling between the MLT and  
277 the F2-region in this indirect way.

### 278 **3.3 Oscillation in Troposphere and Stratosphere**

279 We use the radiosonde observations and MERRA-2 reanalysis data to examine the activity of the  
280 quasi 27-day oscillation in the troposphere and stratosphere. Figure 5 presents the Lomb-Scargle  
281 spectrum of the zonal wind measured by the radiosonde in Hilo and Kauai on Hawaii close to the Maui  
282 meteor radar. The spectral analysis in Figure 5 illustrates that the quasi 27-day periodicity appears  
283 clearly below about 17 km with a maximum magnitude of about  $7 \text{ ms}^{-1}$  around 13 km. However, it does  
284 not have a significant spectral peak above 17 km where the tropopause is, indicating that it is difficult  
285 for the oscillation to penetrate through the tropopause region into the lower stratosphere. This is in  
286 agreement with rapid attenuation of ISOs above the tropopause presented in many early observations  
287 (Madden & Julian, 1971; Ziemke & Stanford, 1991; Rao et al., 2009; Niranjankumar et al., 2011;  
288 Guharay et al., 2014).

289 In order to extend zonal wind information to higher altitudes, Figure 6 plots the Lomb-Scargle  
290 spectrum of the zonal wind above Maui from the MERRA-2 reanalysis. The right vertical axis marks  
291 the approximate heights of the pressure levels derived from logarithmic pressure-height formula. The  
292 spectral feature below 10 hPa ( $\sim 32$  km) in the reanalysis data is roughly consistent with that in the  
293 radiosonde observations. A quasi 27-day periodicity occurs below 108.7 hPa ( $\sim 16$  km) level with an  
294 intensity similar to the radiosonde measurements in Hilo and Kauai. It is interesting that the quasi  
295 27-day oscillation evidently appears between 3.3 and  $3.27 \times 10^{-2}$  hPa ( $\sim 40$ -72 km) levels with a spectral  
296 peak of  $11.2 \text{ ms}^{-1}$  at 0.48 hPa ( $\sim 53$  km) level, and has a narrower spectral width than in the lower

297 atmosphere. However, there is not a significant spectral component of quasi 27-day period at  
298  $3.27 \times 10^{-2}$ - $1.0 \times 10^{-2}$  (~72-80 km) levels, close to the lowest altitude of 80 km in the meteor radar  
299 observation. Based on the radar observation, this oscillation is markedly strengthened above 90 km  
300 again, as shown in Figure 2. Hence, the quasi 27-day oscillation exhibits a three-layer structure from the  
301 troposphere to the MLT over Maui.

302 To examine the meridional feature of the quasi 27-day oscillation, we make a sinusoidal wave fitting  
303 with a 27-day period on the zonal wind at 176.93 and 0.48 hPa (~12 and 53 km) levels along the  
304  $156.25^\circ\text{W}$  longitude based on the MERRA-2 reanalysis data for a same duration of 81 days. The two  
305 pressure levels are chosen because the quasi 27-day component at these levels is strong, as shown in  
306 Figure 6. Figure 7 depicts the evolutions of the fitted amplitude and phase with latitude. The phase  
307 varies slowly with latitude, which seems to show a tendency of the oscillation propagation from the SH  
308 to the NH. The amplitudes have several extreme values between  $25^\circ\text{S}$  and  $25^\circ\text{N}$ , especially the maximal  
309 values around  $20^\circ\text{S}$  and  $20^\circ\text{N}$  at 176.93 hPa. In that case, we present the Lomb-Scargle spectrum of the  
310 zonal wind from the radiosonde observation in Pago Pago of the SH, which is shown in Figure 5(c). It  
311 can be noted that the quasi 27-day oscillation in the troposphere is more prominent above Pago Pago  
312 than above Hilo and Kauai. Therefore, the oscillation is active at low latitudes in the both hemispheres.

313

#### 314 **4. Discussion**

315 As shown in Figures 5 and 6, the quasi 27-day oscillation can be identified in the troposphere, but  
316 does not appear clearly in the lower stratosphere. Eckermann and Vincent (1994) reported the ISOs with  
317 periods of 35-60 days in the zonal wind in the equatorial MLT based on MF radar observation, and  
318 argued that it is unlikely for these ISOs to propagate to the MLT from the lower atmosphere. Instead,  
319 they proposed that GWs and tides generated by tropical convection are modulated by the oscillations in



320 the lower atmosphere and then in turn drive similar periodicities in the MLT when propagating to the  
321 MLT and depositing their energy and momentum into the mean flow through instability (Eckermann &  
322 Vincent, 1994; Eckermann et al., 1997). According to this mechanism of wave coupling among the  
323 different atmospheric layer, we investigate the GW activity in the lower atmosphere based on the  
324 radiosonde observations in the tropic Hilo, Kauai and Pago Pago. The GW perturbations are analyzed  
325 separately at 1-9 km of the troposphere and at 20-28 km of the lower stratosphere, which is for three  
326 reasons: 1) avoiding the sharp variations of the zonal wind and temperature from the upper troposphere  
327 to the lower stratosphere; 2) the approximately constant buoyancy frequency in each chosen height  
328 ranges; and 3) examining the GW features in the convective source region of 1-9 km and in the  
329 propagation region of 20-28 km.

330 We follow the analysis technique of Allen and Vincent (1995) and Vincent and Alexander (2000) to  
331 derive the GW perturbations from the profiles of radiosonde sounding. Assuming that the observed  
332 zonal wind, meridional wind and temperature  $[u, v, T]$  mainly consist of the background  $[\bar{u}, \bar{v}, \bar{T}]$   
333 and GW perturbations  $[u', v', T']$ , we estimate the background  $[\bar{u}, \bar{v}, \bar{T}]$  by fitting a second-order  
334 polynomial to the vertical profiles of  $[u, v, T]$  in the chosen height interval. This second-order  
335 polynomial fitted background was widely used in GW analysis from radiosonde observations (Allen &  
336 Vincent, 1995; Vincent & Alexander 2000; Zhang & Yi, 2007; Huang et al., 2018), which can reduce  
337 the effect of long vertical wavelength waves. The total fluctuation quantities are obtained from the  
338 observed profiles by removing the fitted second-order polynomial. In order to remove fluctuations due  
339 to smaller-scale effects, such as measurement error, drag variation of balloon, noise introduced by the  
340 interpolation process (Zhang et al, 2012), we apply a high-pass filter to extract the GW components  
341 from the total fluctuation. The cutoff wavelength of the filter is selected to be 0.5 km since extensive  
342 observations show the GW vertical wavelengths are in general more than 0.5 km. The total energy per

343 unit mass ( $E$ ) is used as a measurement for GW activity, which is written as follows (Allen & Vincent,  
 344 1995; Vincent & Alexander 2000),

$$345 \quad E = \frac{1}{2} \left[ \overline{u'^2} + \overline{v'^2} + \frac{g^2 \widehat{T'^2}}{N^2} \right] \quad (1)$$

346 where  $\widehat{T'} = \frac{T'}{T}$  is the normalized temperature perturbation;  $g$  is the acceleration due to gravity;  $N$  is  
 347 the buoyancy frequency; and the overbar means an unweighted average over height. The vertical wind  
 348 perturbation of GWs is neglected in Eq. (1) because there is no vertical wind in radiosonde observation,  
 349 and the vertical wind perturbation of GWs is much smaller than their horizontal wind perturbation.

350 We derive the total energy of GWs from the observed profiles of radiosonde in Hilo, Kauai and  
 351 Pago Pago, and then calculate the daily averaged energy in Pago Pago of the SH and between the two  
 352 station of the NH, respectively, for the sake of reducing the impact of chance events. Figure 8 shows the  
 353 Lomb-Scargle spectra of daily mean GW energies at 1-9 km of the troposphere and 20-28 km of the  
 354 lower stratosphere in the NH and SH. One can note from Figure 8 that a quasi 27-day oscillation of GW  
 355 energy occurs in both the troposphere and the lower stratosphere. The quasi 27-day periodicity has a  
 356 stronger spectral intensity in the lower stratosphere than in the troposphere because the GW amplitude  
 357 increases with height owing to the exponentially decreasing atmospheric density. Importantly, although  
 358 the quasi 27-day oscillation does not evidently appear in the lower stratospheric wind field, its spectral  
 359 peak in the GW energy is much sharper in the lower stratosphere than in the troposphere, which may be  
 360 due to the filtering effect by the quasi 27-day variation in the tropospheric wind field (Moss et al., 2016;  
 361 Tsuchiya et al., 2016), as shown in Figure 5. Therefore, the GW energy investigation seems to support  
 362 the coupling mechanism that GWs modulated by oscillation in the lower atmosphere induce the similar  
 363 periodicity in the MLT through their energy and momentum transport (Eckermann & Vincent, 1994;  
 364 Eckermann et al., 1997).

365 As for the ISO enhancement in the upper stratosphere again, there are two possible mechanisms  
366 proposed in previous studies. Kumar and Jain (1994) argued that ISO propagates directly upward to the  
367 stratosphere from the tropical troposphere through partial energy transmission, and Huang et al. (2015)  
368 confirmed the ISO penetration into the stratosphere based on the radiosonde observation. Ziemke and  
369 Stanford (1991) suggested that upward propagating ISO is refracted to mid latitudes, and then refracted  
370 back into the tropical upper stratosphere. The phase in the stratosphere shown in Figure 7(a) seems to  
371 imply that the oscillation propagates towards the NH low latitudes from the NH and SH mid latitudes.  
372 In addition, the quasi 27-day oscillation may be partially attributable to GWs. As the GWs modulated at  
373 a 27-day period propagate to the upper stratosphere, they may deposit their partial momentum and  
374 energy into the background flow by means of instability and background wind filtering. The GWs with  
375 their remaining energy can further propagate to higher levels and further induce the quasi 27-day  
376 variability in the mesopause region. Still, the contribution of tidal waves cannot be ruled out since we  
377 cannot examine the tidal activity in the lower atmosphere due to a coarse temporal resolution of 12 h in  
378 the radiosonde measurement.

379 Earlier studies discussed possible mechanisms of ISO generation (Zhang, 2005). It is generally  
380 recognized that ISOs in the tropics originate mainly from convective activity. OLR is often used as a  
381 proxy of convection activity. We choose the OLR data at the locations of (20°N, 157.5°W) and (15°S,  
382 170°W), close to Maui and Pago Pago, to examine the convective activity, and then carry out a wavelet  
383 transform on the OLR. Figure 9 presents their wavelet spectra. It shows that there is a quasi 27-day  
384 variation in the OLR over the two tropical sites. We also calculate the specific humidity from the  
385 radiosonde observation in Pago Pago. Figure 10 shows the Lomb-Scargle spectrum of the specific  
386 humidity from the radiosonde measurement and the reanalysis data above Pago Pago. The spectra of  
387 specific humidity between the observation and reanalysis data are consistent with each other. The quasi

388 27-day periodicity is the dominant component in the water vapor variation, and the spectral analysis  
389 above 7 km (not presented here) indicates that the dominant 27-day component extends all the way to  
390 the tropopause region. Hence, the periodicity in the OLR and water vapor demonstrates that there is a  
391 quasi 27-day variability in the convective activity in the tropics. Since convections are a main source of  
392 tropical waves and ISOs (Fritts & Alexander, 2003; Zhang, 2005), their variability can lead to the quasi  
393 27-day ISO not only in the zonal wind but also in the GW activity in the tropical troposphere. When the  
394 GWs generated in the convection source region propagate upward, they are further modulated by the  
395 quasi 27-day oscillation of the zonal wind, resulting in a more prominent 27-day periodicity of the GW  
396 activity in the stratosphere than in the troposphere, as shown in Figure 8.

397 The ionospheric activity is complex because it is affected not only by solar and geomagnetic  
398 activities, but also by neutral-plasma interaction, thus the oscillation in the hmF2 may have several  
399 possible origins (Forbes et al., 2000; Pancheva et al., 2002). We investigate the relationship of the hmF2  
400 oscillation with solar radiation and geomagnetic activity to reveal its main origin. Here, F10.7 and Kp  
401 index are chosen as proxies for solar radiation and geomagnetic activity, respectively. Figure 11 depicts  
402 that the wavelet spectra of the F10.7 and Kp index in the duration of our attention. One can see from  
403 Figure 11 that the F10.7 has an oscillation with a dominant period less than 27 days. The F10.7  
404 oscillation is weak because the year 2006 is in the late declining phase of solar cycle 23, and the weak  
405 oscillation arises mainly before day 40. The geomagnetic activity exhibits a significant quasi 27-day  
406 variation between about day 10 and 36. The hmF2 oscillation shown in Figure 4(a) is robust after about  
407 day 30. By comparing Figure 4(a) with Figures 4(b), 9 and 11, it can clearly be noted that both the  
408 dominant period and active time of the hmF2 oscillation are in better correlation with the variabilities in  
409 the wind field and OLR rather than in the F10.7 and geomagnetic activity. Hence, the hmF2 periodicity  
410 originates mainly from the oscillation in the convection, waves and wind field, which demonstrates that

411 the quasi 27-day variability in the ionosphere can also come from the natural wind and "meteorological  
412 influences" from below (Forbes et al., 2000; Pancheva et al., 2002). In other words, the variability in the  
413 tropical convective activity drives the ISO in the zonal wind and modulates the wave activity in the  
414 lower atmosphere, which leads to a coupling from the lower atmosphere to the MLT and F2 region  
415 through wave propagation and wave-flow interactions (Forbes et al., 2000, 2018).

416

## 417 **5 Summary**

418 Combining the meteor radar, radiosonde and digisonde observations and the MERRA-2 reanalysis  
419 data for 81 days from 12 August to 31 October 2006, we study a quasi 27-day ISO event in the tropical  
420 zonal wind from the troposphere to MLT, and the corresponding hmF2 variability generated from the  
421 oscillation.

422 The radiosonde observations at the three tropical stations distributed in the NH and SH show that  
423 the quasi 27-day oscillation originates from the lower atmosphere, and has an amplitude of about  $7 \text{ ms}^{-1}$   
424 in the zonal wind in the upper troposphere. Above the tropopause, the oscillation attenuates rapidly,  
425 which is in agreement with most previous observations. The reanalysis data indicates that the quasi  
426 27-day oscillation is strengthened in the upper stratosphere, and attains a magnitude of about  $11 \text{ ms}^{-1}$ .  
427 The ISO enhancement in the stratopause region is also reported in early studies. Whereas, the  
428 oscillation decays quickly again from about 70 km. The meteor radar measurement illustrates that the  
429 oscillation increases again from about 90 km, and reaches an intensity of about  $15 \text{ ms}^{-1}$  at 100 km.  
430 Hence, the quasi 27-day oscillation exhibits an interesting three-layer structure from the troposphere to  
431 the MLT. The oscillation in the MLT is different from the previous observations in which the ISO in the  
432 zonal wind generally has the amplitude peak at about 80-85 km, and then decreases with height  
433 (Eckermann et al., 1997; Luo et al., 2001; Huang et al., 2015). In the period that we focus on, the

434 digisonde observation shows that the quasi 27-day variability arises clearly in the hmF2. This implies  
435 that a coupling from the lower atmosphere to the MLT and F2 region occurs through the quasi 27-day  
436 oscillation.

437 As proxies of convection activity, the OLR and specific humidity observations demonstrate that  
438 there is a quasi 27-day periodicity in the convective activity in the tropics. The convection is a main  
439 excitation source of tropical ISOs and waves, thus the quasi 27-day variability takes place not only in  
440 the zonal wind in the troposphere but also in the convectively modulated GW activity. As the GWs  
441 propagate upward, they can be modulated by the zonal wind oscillation in the upper troposphere. Hence,  
442 the quasi 27-day variability in the GW activity has a sharper spectrum in the lower stratosphere relative  
443 to that in the troposphere, even though the zonal wind in the lower stratosphere does not display a quasi  
444 27-day ISO. The GWs may deposit their partial energy and momentum into the background flow in the  
445 stratopause region through the instability and background wind filtering, which might contribute partly  
446 to driving the quasi 27-day variability there, despite that the quasi 27-day ISO may be refracted back to  
447 the tropical stratopause region from midlatitudes. After partial energy loss, the GWs can further  
448 propagate upward to higher altitudes and induce the quasi 27-day ISO in the mesopause region through  
449 their momentum and energy deposition. Therefore, GWs play an important role in the connection of the  
450 quasi 27-day oscillation between the lower atmosphere and the MLT, as proposed by Eckermann and  
451 Vincent (1994) and Eckermann et al. (1997).

452 The 27-day variability in the ionosphere could originate from the periodicity of solar rotation,  
453 geomagnetic activity and zonal neutral wind or meteorological influences from below. The wavelet  
454 spectra of the F10.7 and Kp index indicate that in the time period we studied, the effects of solar  
455 rotation and geomagnetic activity are weak, and their quasi 27-day period arises before day 40. The  
456 hmF2 oscillation is robust after day 30, and its dominant period and active duration are in good

457 agreement with those in the wind field of the MLT and OLR rather than in the F10.7 and Kp index. This  
458 demonstrates that the hmF2 periodicity originates mainly from the oscillation in the wind field  
459 associated with the convective activity. Therefore, the tropical convective activity can influence the  
460 dynamics of the MLT and F2 region through the propagation and interaction of these convectively  
461 modulated waves and ISOs.

462

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749           **Caption:**

750   **Figure 1.** Daily mean (a) zonal and (b) meridional winds from Maui meteor radar observation for 81  
751 days from 12 August to 31 October 2006.

752   **Figure 2.** Lomb-Scargle spectra of (a) zonal and (b) meridional winds observed by Maui meteor radar.  
753 The dashed vertical line corresponds to a period of 27 days.

754   **Figure 3.** (a) Amplitude and (b) phase of quasi 27-day oscillation in zonal wind from Maui meteor radar  
755 observation. The asterisk denotes the value derived from the sinusoidal wave fitting, and the dashed line  
756 in Panel (b) denotes the linearly fitted phase.

757   **Figure 4.** Wavelet spectra of (a) ionospheric hmF2 in Kwajalein and (b) zonal wind at 100 km in Maui.

758   **Figure 5.** Lomb-Scargle spectrum of zonal wind from radiosonde observations in (a) Hilo, (b) Kauai  
759 and (c) Pago Pago. The dashed vertical line corresponds to a period of 27 days.

760   **Figure 6.** Lomb-Scargle spectrum of zonal wind over Maui obtained from MERRA-2 reanalysis data.  
761 The dashed vertical line corresponds to a period of 27 days.

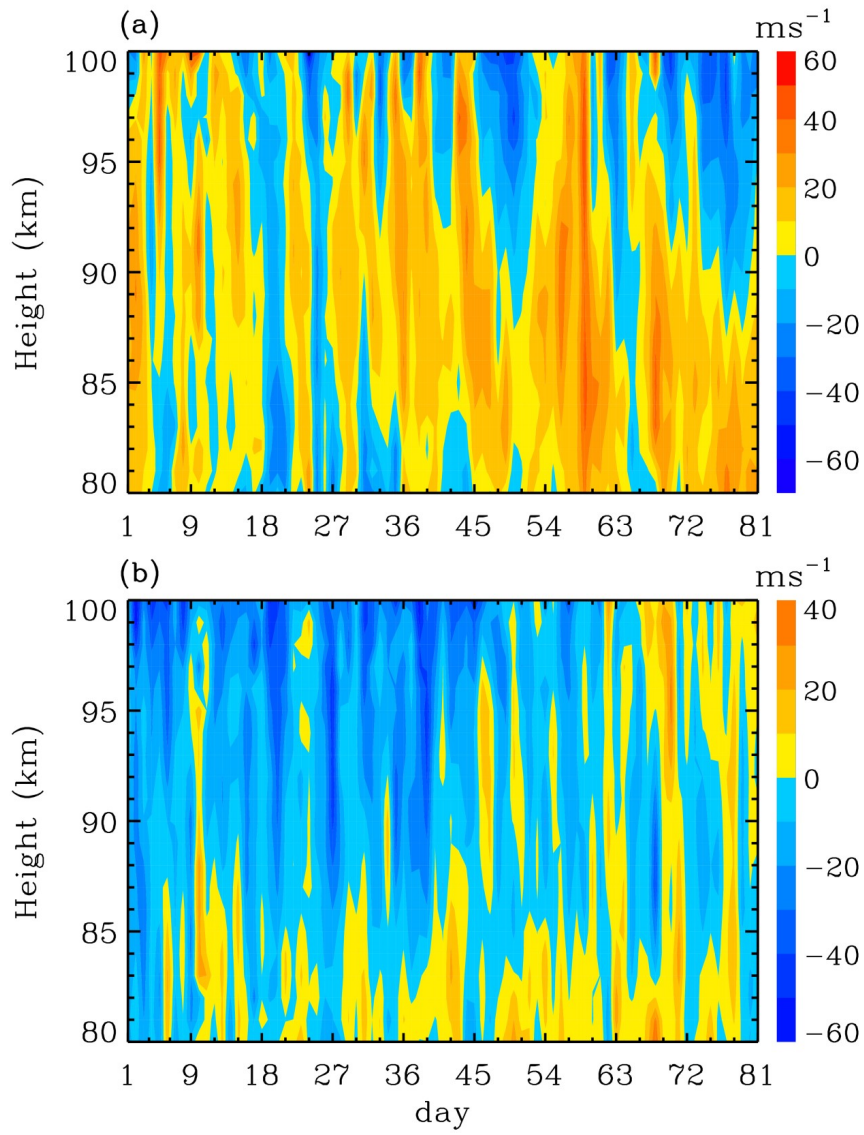
762   **Figure 7.** Latitudinal evolutions of fitted amplitude and phase in zonal wind at (a) 0.48 hPa and (b)  
763 176.93 hPa levels from MERRA-2 reanalysis data. The asterisk denotes the fitted value.

764   **Figure 8.** Lomb-Scargle spectra of daily mean GW energies at (a, c) 1-9 km and (b, d) 20-28 km  
765 derived from radiosonde observations. Panels (a, b) represent the spectra of GW energy averaged  
766 between Hilo and Kauai, and Panels (c, d) represent the spectra of averaged GW energy in Pago Pago.  
767 The dashed vertical line corresponds to a period of 27 days.

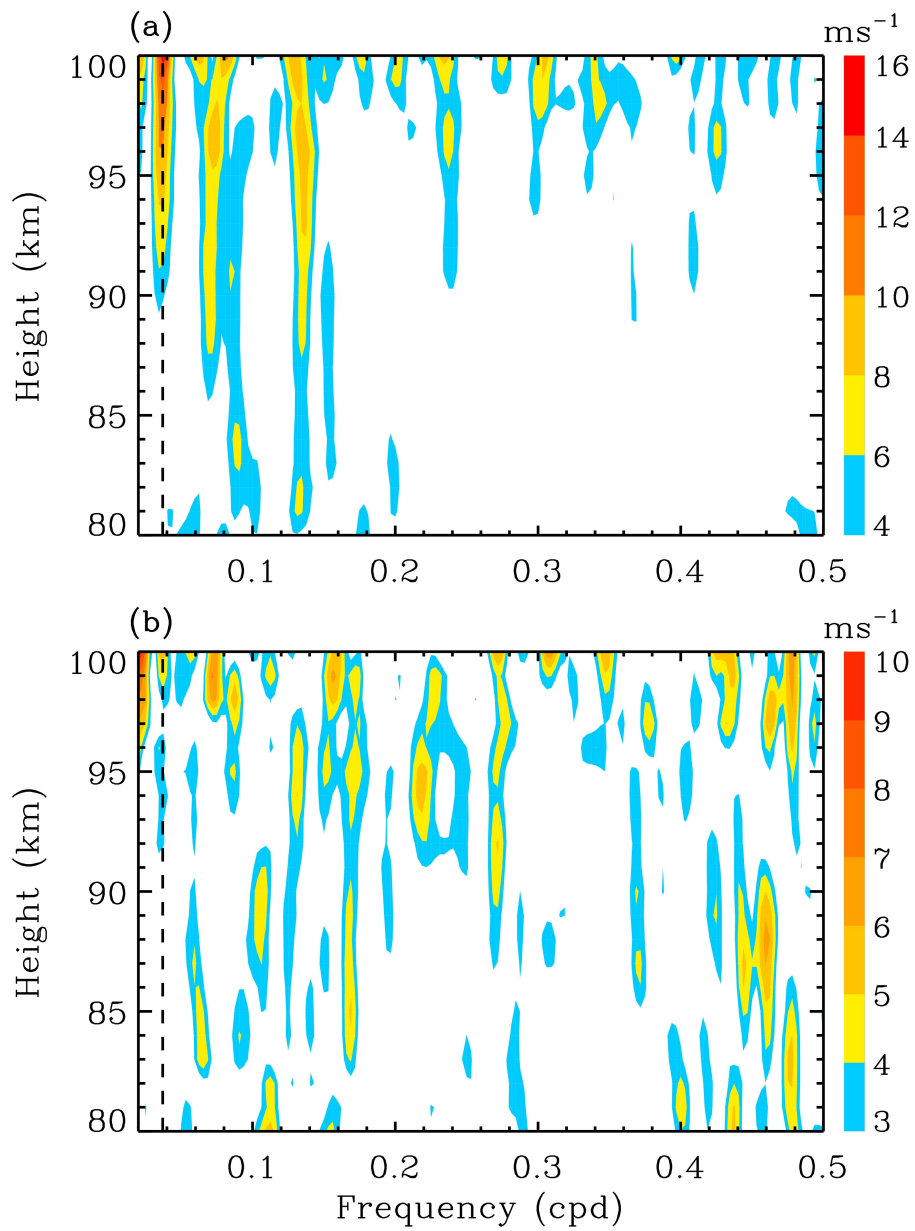
768   **Figure 9.** Wavelet spectrum of OLR at (20°N, 157.5°W) and (15°S, 170°W), close to Maui and Pago  
769 Pago, respectively.

770   **Figure 10.** Lomb-Scargle spectrum of specific humidity derived from (a) radiosonde observation and (a)  
771 reanalysis data over Pago Pago. The dashed vertical lines correspond to a period of 27 days.

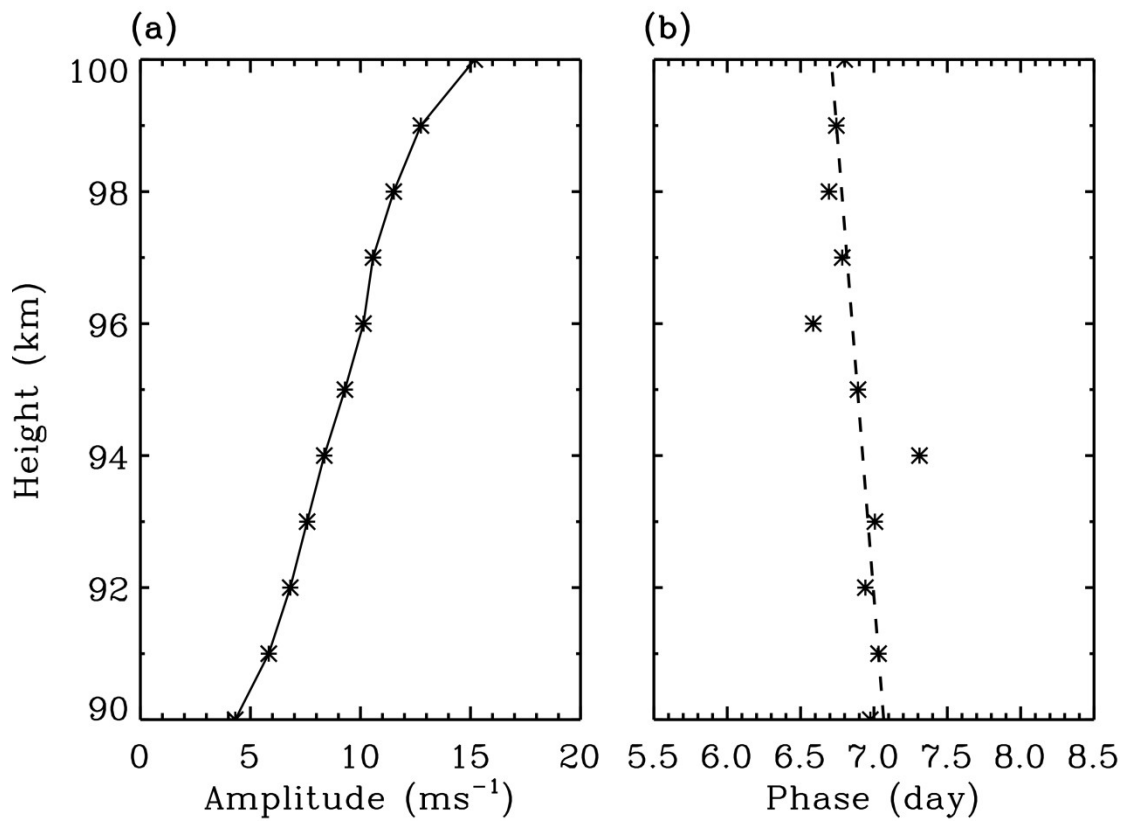
772 **Figure 11.** Wavelet spectra of (a) F10.7 and (b) Kp index.



773  
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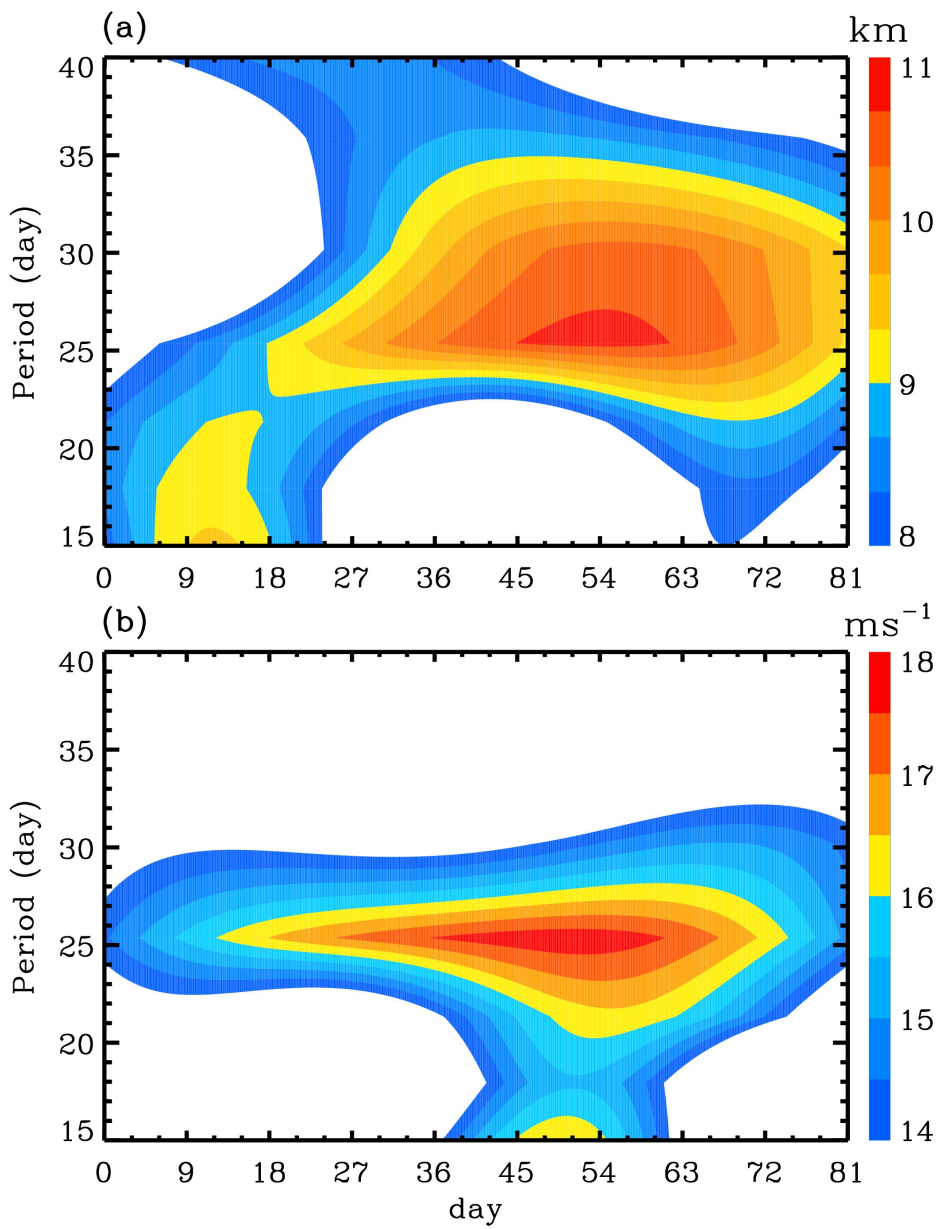


776  
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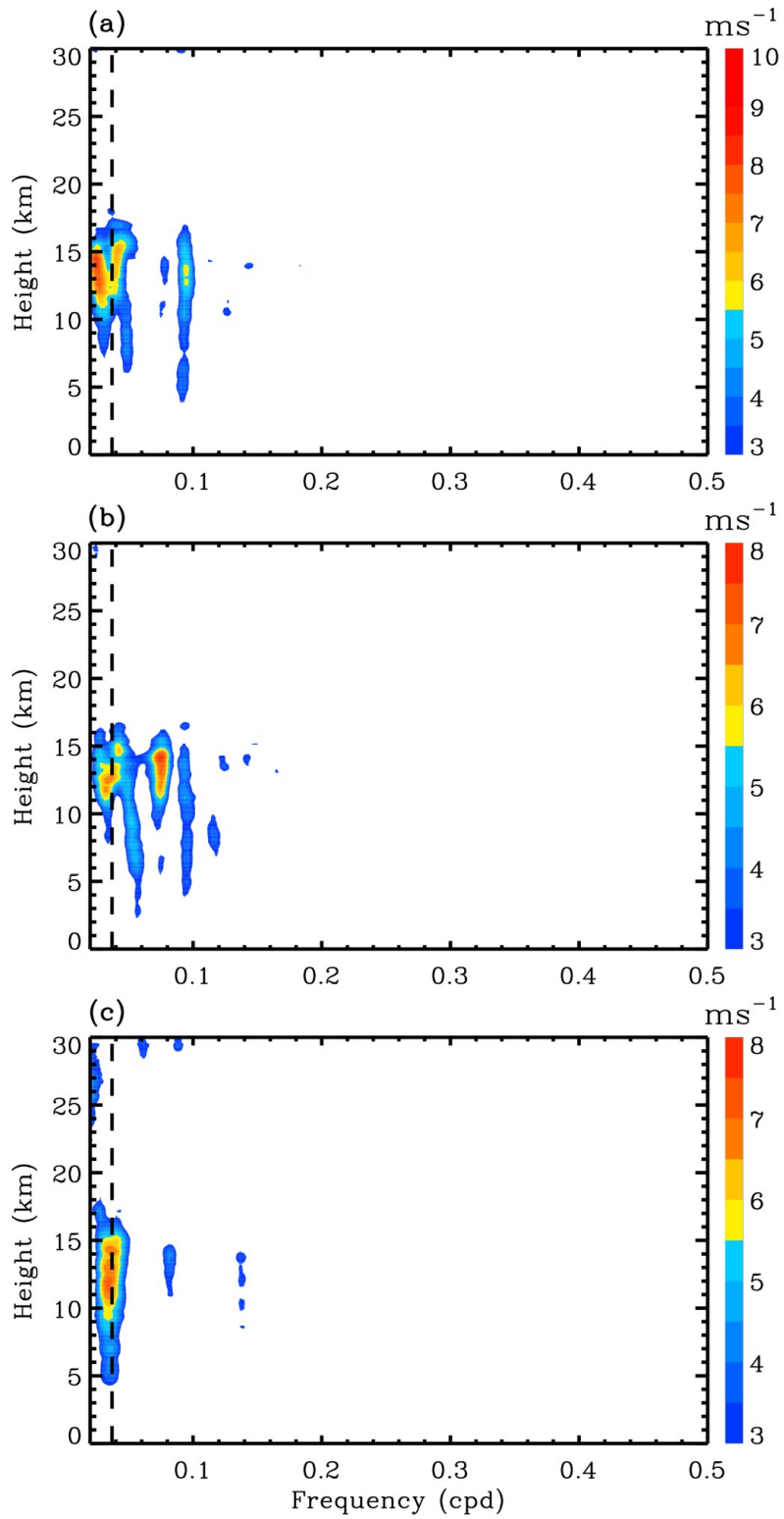
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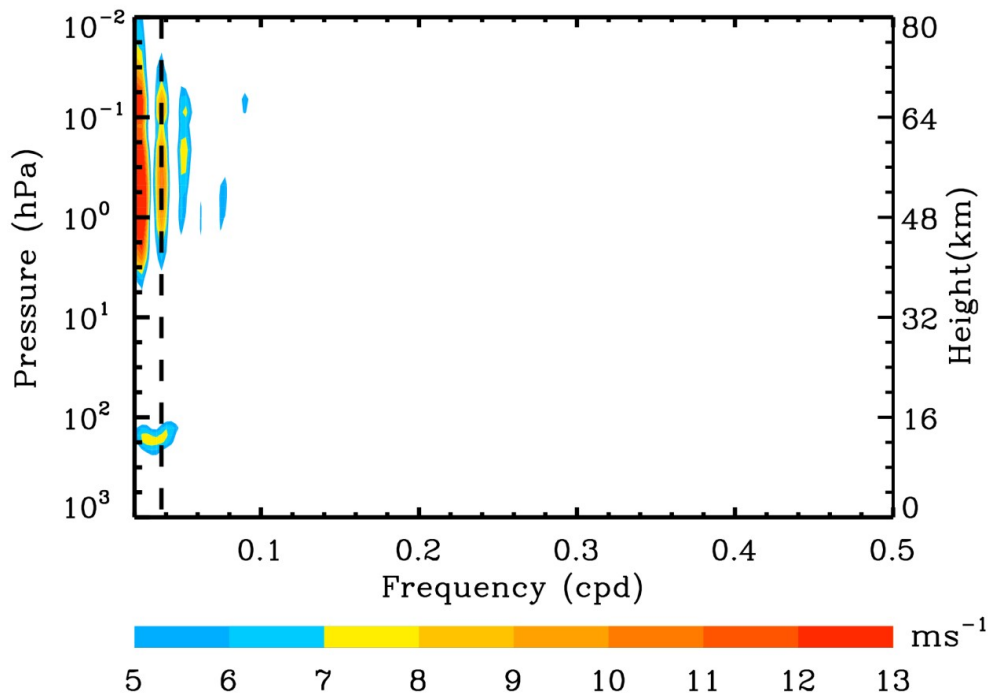
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785

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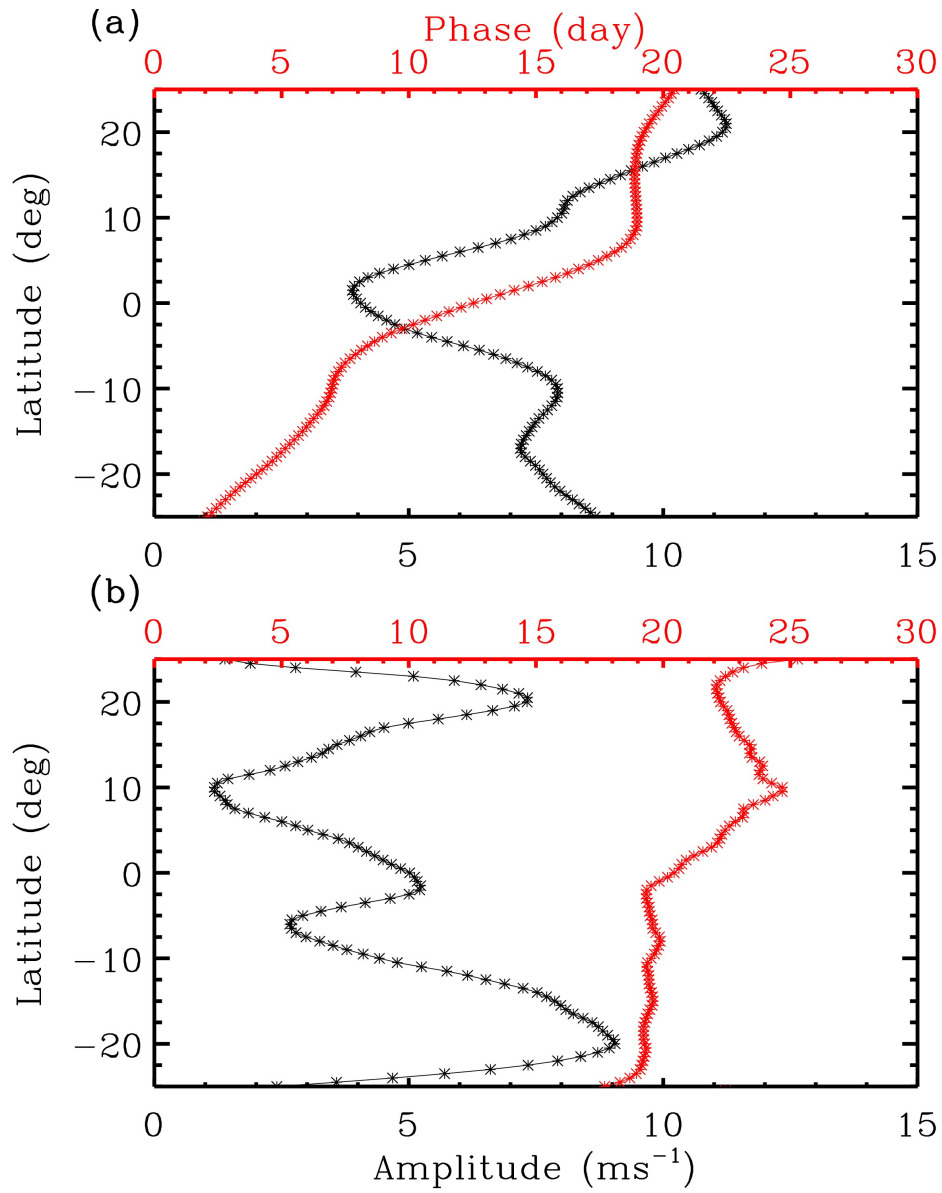
787 and (c) Pago Pago. The dashed vertical line corresponds to a period of 27 days.



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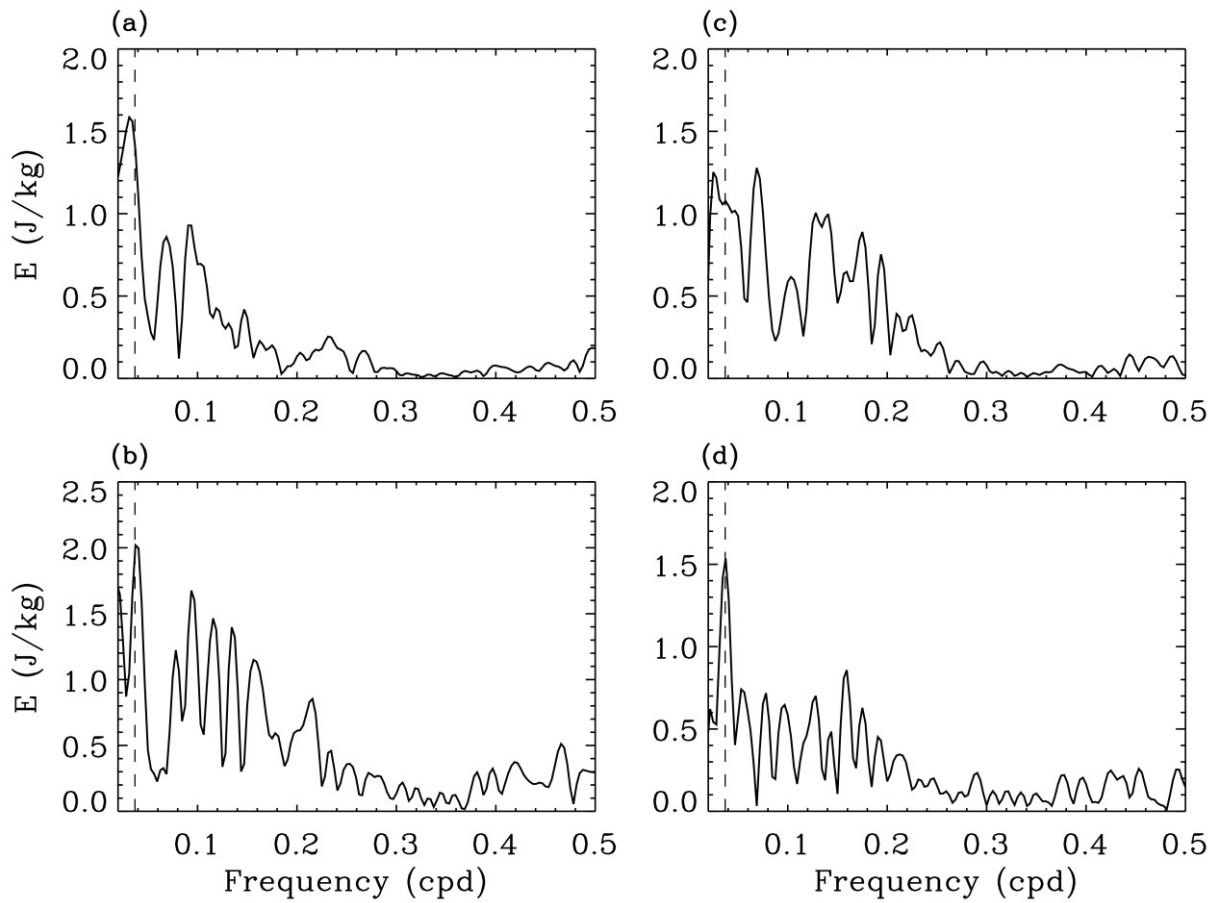
790 The dashed vertical line corresponds to a period of 27 days.



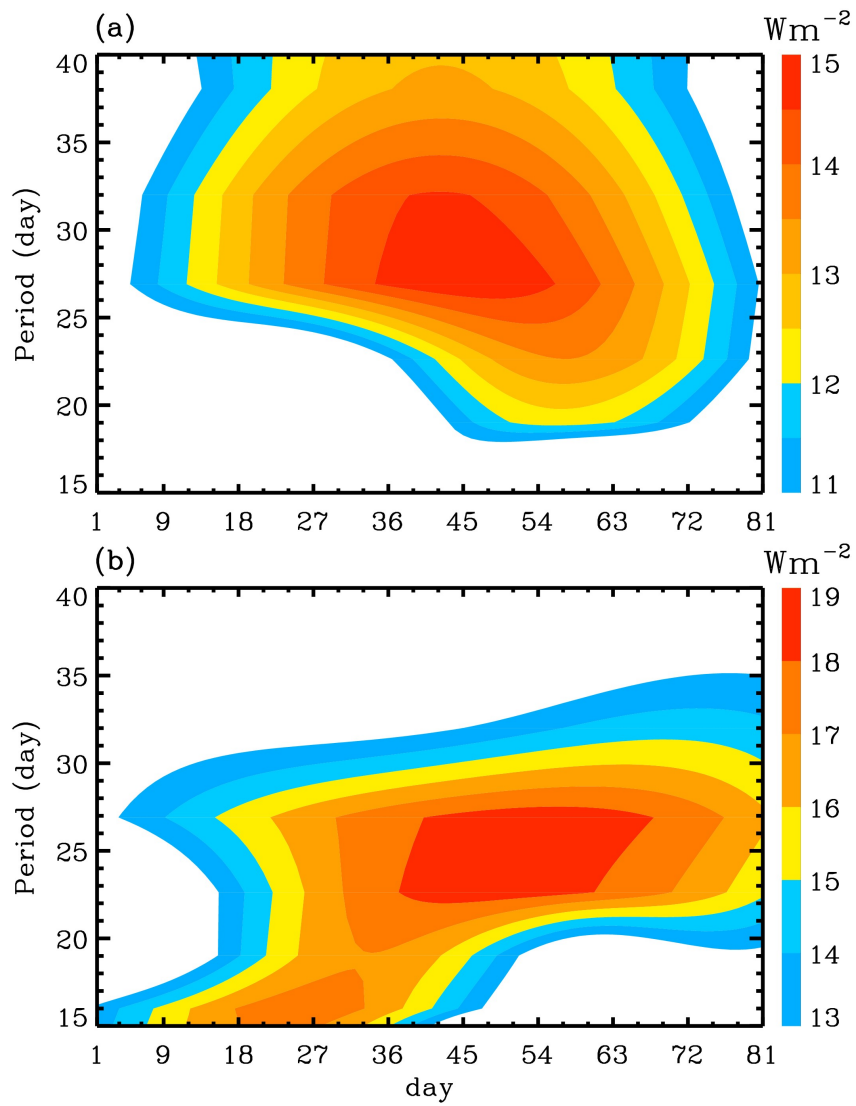
791

792 **Figure 7.** Latitudinal evolutions of fitted amplitude and phase in zonal wind at (a) 0.48 hPa and (b)

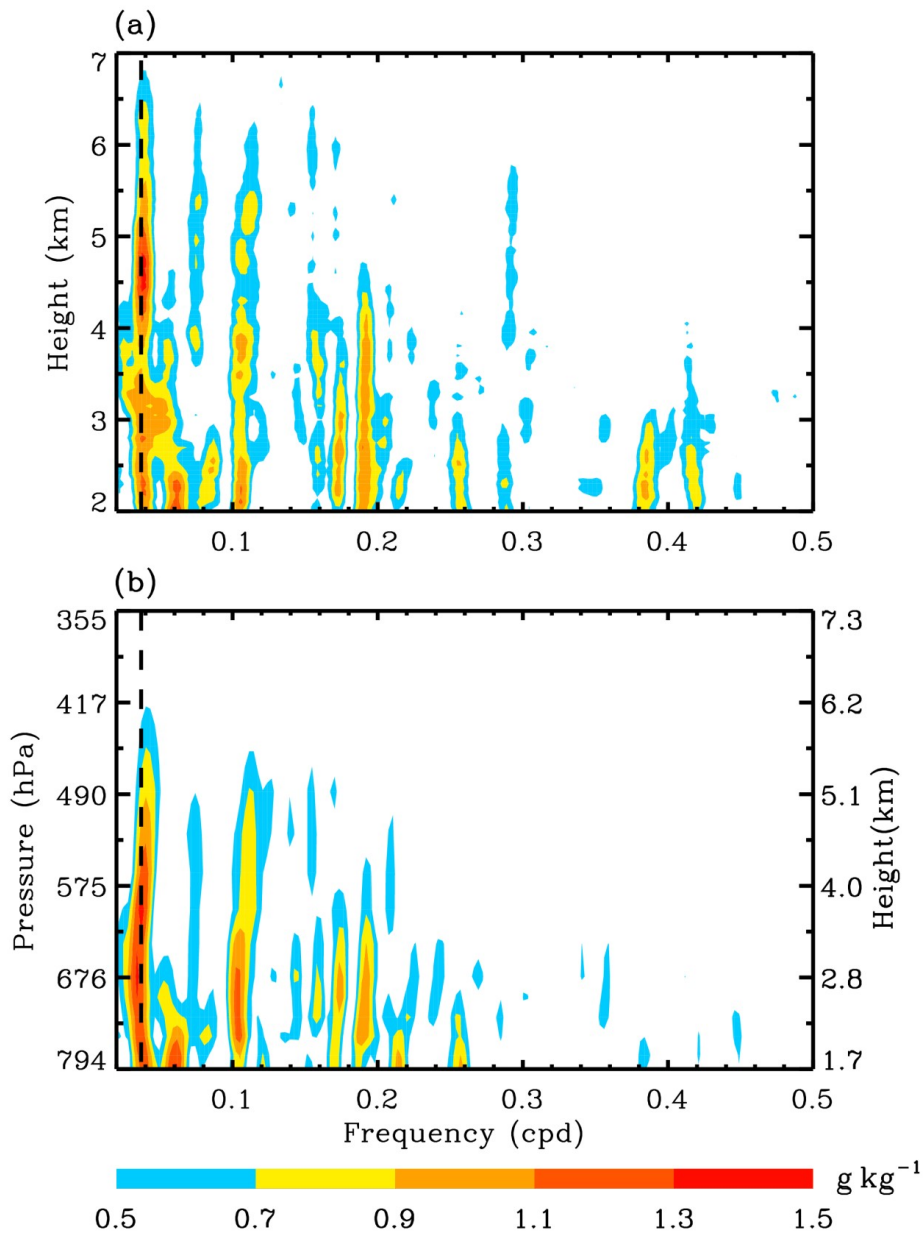
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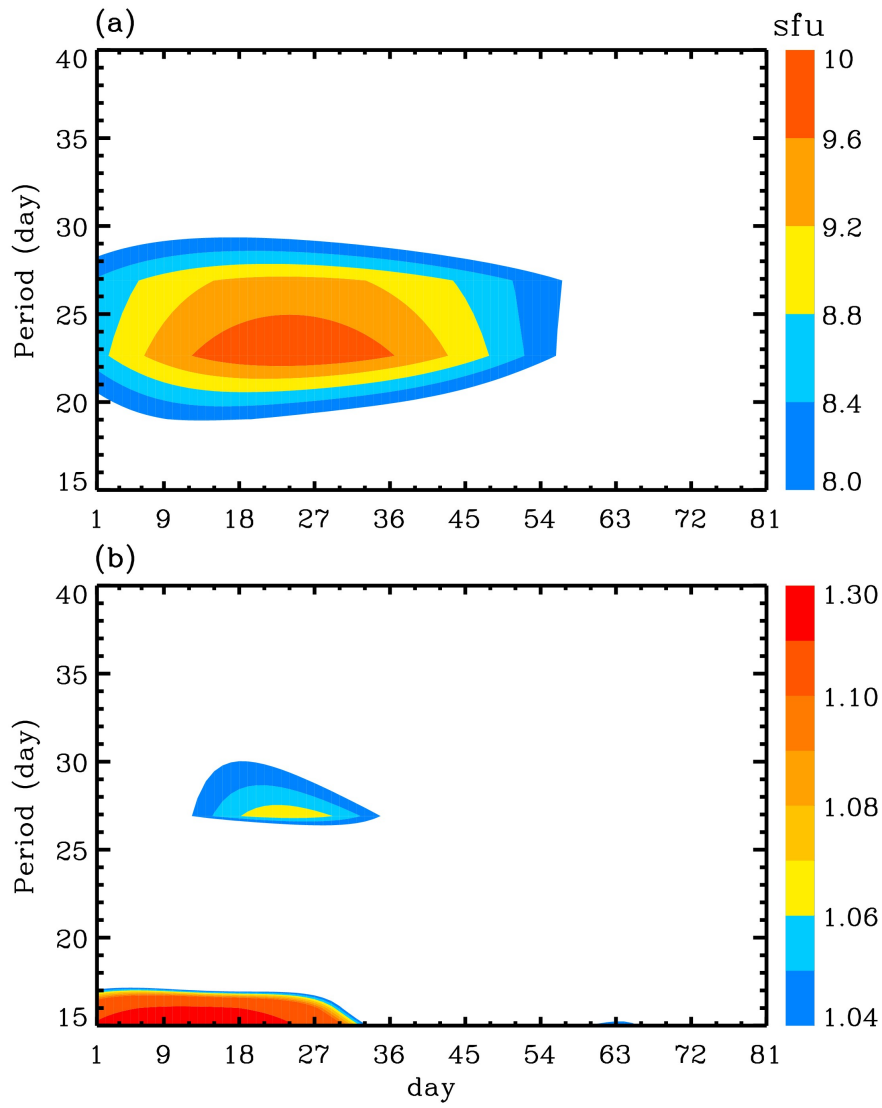
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805

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