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Euclid preparation: XXVIII. Modelling of the weak lensing angular power spectrum

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ABSTRACT

This work considers which higher-order effects in modelling the cosmic shear angular power spectra must be taken into account for *Euclid*. We identify which terms are of concern, and quantify their individual and cumulative impact on cosmological parameter inference from *Euclid*. We compute the values of these higher-order effects using analytic expressions, and calculate the impact on cosmological parameter estimation using the Fisher matrix formalism. We review 24 effects and find the following potentially need to be accounted for: the reduced shear approximation, magnification bias, source-lens clustering, source obscuration, local Universe effects, and the flat Universe assumption. Upon computing these explicitly, and calculating their cosmological parameter biases, using a maximum multipole of $\ell = 5000$, we find that the magnification bias, source-lens clustering, source obscuration, and local Universe terms individually produce significant ($> 0.25\sigma$) cosmological biases in one or more parameters, and accordingly must be accounted for. In total, over all effects, we find biases in Ω_m , Ω_b , h , and σ_8 of 0.73σ , 0.28σ , 0.25σ , and -0.79σ , respectively, for flat Λ CDM. For the w_0w_a CDM case, we find biases in Ω_m , Ω_b , h , n_s , σ_8 , and w_a of 1.49σ , 0.35σ , -1.36σ , 1.31σ , -0.84σ , and -0.35σ , respectively; which are increased relative to the Λ CDM due to additional degeneracies as a function of redshift and scale.

Key words. Cosmology: observations – Gravitational lensing: weak – Methods: analytical

1. Introduction

Our current best-in-class framework for parameterising the Universe, the Lambda cold dark matter (Λ CDM) model, leaves several open questions. A key component, yet to be fully explained, is the acceleration of the expansion of the Universe and its proposed driver: dark energy. A powerful tool for studying this is cosmic shear – the distortion of the ellipticities we observe for distant galaxies, by weak gravitational lensing from the large-scale structure of the Universe (LSS; see for example Albrecht et al. 2006).

To-date, the most recent generation of cosmic shear surveys (Hikage et al. 2019; Asgari et al. 2021; Abbott et al. 2022) has been able to achieve precision cosmology competitive with cosmic microwave background experiments, for a combination of σ_8 and Ω_m (Planck Collaboration et al. 2018). Now, upcoming Stage IV surveys (Albrecht et al. 2006) will probe a greater area and depth than previously possible. For example, telescopes such as *Euclid*¹ (Laureijs et al. 2011), *Nancy Grace Roman*² (Akeson et al. 2019), and the Vera C. Rubin Observatory³ (LSST Science Collaboration et al. 2009) will achieve more than an order-of-magnitude increase in precision over existing surveys (Euclid Collaboration: Blanchard et al. 2020, hereafter EC20). We must therefore ensure that any sources of bias in our theoretical formalism are properly accounted for.

In this work, we consider the common approximations made, and effects typically neglected when modelling the cosmic shear angular power spectrum. This is of importance when deriving cosmological parameters from a shear-only analysis, but also in a 3×2 pt analyses where modelling the weak lensing power spectrum sufficiently is also essential. Throughout the literature, these effects have been studied independently and using varying specifications. Here, we evaluate them in a consistent framework, and quantify their cumulative impact on cosmology inferred from *Euclid*'s weak lensing probe. As a first step, we review the literature and pinpoint which terms are potentially significant, as well as those for which the impact on the shear power spectrum has not been evaluated. A comparison of the typical magnitudes of the studied corrections is given in Table 1. To create Table 1 we manually read the published numbers from graphs in the referenced papers for the auto-correlation cosmic shear power spectrum for a redshift bin closest to $z = 1$ for a in each paper. These values are *only approximate* due to the inherent inaccuracy of reading value from graphs and varying assumptions the papers⁴.

From our survey of the literature we identify the following as potentially significant systematic effects requiring full analysis:

- Reduced shear approximation: the effect of assuming that the measured two-point statistics of reduced shear are equal to those of the shear field (Shapiro 2009; Krause & Hirata 2010). It has been previously shown that this approximation will require relaxation for *Euclid* (Deshpande et al. 2020a),
- Magnification bias: the change in the observed number density of sources, due to galaxies at the flux limit of the survey having their flux increased or decreased due to magni-

fication by lensing (Turner et al. 1984). This effect has also shown to significantly bias cosmological information from *Euclid* (Deshpande et al. 2020a; Duncan et al. 2022) if not accounted for. Additionally, magnification bias must also be accounted for in probes of galaxy clustering. Its impact on the *Euclid* galaxy clustering probe is discussed in Euclid Collaboration: Lepori et al. (2022).

- Source-lens clustering: the intrinsic clustering of source galaxies correlated with the density field (Bernardeau 1998; Hamana et al. 2002; Yu et al. 2015). Typically, it is assumed source galaxies are distributed homogeneously across the sky.
- Source obscuration: A reduction in the observed galaxy distribution due to closely-spaced and blended or overlapping source galaxies (Hartlap et al. 2011).
- Local Universe effects: a possible bias in our measurements of summary statistics of the LSS due to residing in a region with a higher-than-average density (Reischke et al. 2019; Hall 2020).
- Flat Universe assumption: the impact of assuming that non-flat geometries are sufficiently well represented by modifying the expression for comoving distance, and neglecting the additional change in the lensing kernels used to calculate the shear power spectrum (Taylor et al. 2018b).

We made this determination by first excluding the terms which are fourth-order in the lens potential or higher, as these have consistently been shown to be sub-dominant (Cooray & Hu 2002; Shapiro & Cooray 2006). Among these are time delay-lens coupling, and deflection-deflection coupling (Bernardeau et al. 2010), which result from foregoing the small-angle and thin lens approximations and solving the Sachs equation explicitly. Similarly, fourth-order correction terms resulting from relaxing the Born approximation and accounting for line-of-sight coupling of two foreground lenses are negligible (Shapiro & Cooray 2006). Additionally, fourth-order reduced shear corrections (Krause & Hirata 2010) were also neglected. Re-enforcing the sub-dominance of these terms is the fact that the standard reduced shear correction matches forward models and N-body simulations sufficiently well (Dodelson et al. 2006; Deshpande et al. 2020a). We further excluded fourth-order and higher terms resulting from the contribution of dark energy pressure to the lensing potential (Simpson et al. 2010).

Furthermore, we neglected finite beam corrections (Fleury et al. 2017; Fleury et al. 2019), that manifest as a fractional correction to the lensing power spectra on very small scales which is approximately $-(1/3)(\ell\theta)^2$, where ℓ is the angular multipole of the power spectrum and θ the mean angular size of galaxies (this was verified by Breton & Fleury 2021, using ray-tracing simulations); for *Euclid* this results in a fractional correction of $-(\ell/1.2 \times 10^6)^2$ to the power spectrum. We also neglect the effects of spatially-varying survey depth (Heydenreich et al. 2020), which can be accounted for directly in the covariance matrix through forward-modelling (Loureiro et al. 2021).

Of the remaining effects, we then neglected those significantly smaller than the reduced shear correction. This proves a good comparison because this correction has been consistently demonstrated to produce biases close to the significance threshold (Shapiro 2009; Deshpande et al. 2020a). Accordingly, we neglected the impact of the Doppler-shift of galaxies on their two-point statistics. Due to the inhomogeneity of the Universe, galaxies have peculiar velocities which affect the measurement of their redshifts. If this is taken into account, it results in an additional contribution to the reduced shear (Bernardeau et al.

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¹ <https://www.euclid-ec.org/>

² <https://roman.gsfc.nasa.gov/>

³ <https://www.lsst.org/>

⁴ Every paper made various slightly differing assumptions regarding survey area, depth and number density, in all cases details can be found in the references. In the case that only correction function analyses were available, not power spectrum, we performed a Hankel transform over the correlation functions of the graphs over the angular range available.

2010). As can be seen in Table 1, this effect can be safely neglected as it is two or three orders of magnitude below the reduced shear correction. Additionally, the cosmological parameter biases resulting from it are two orders of magnitude below the level of concern (Deshpande & Kitching 2021).

Likewise, the effect of unequal-time correlators was neglected, as the resulting correction to the angular power spectrum is more than four orders of magnitude smaller than that for the reduced shear (Kitching & Heavens 2017); as illustrated in Table 1. This correction is a consequence of relaxing the equal-time approximation, which approximates the cross-correlation matter power spectrum evaluated at different times as either the power spectrum at a fixed time, or by a geometric mean.

We also did not need to explicitly evaluate the effects of relaxing the Limber, and flat-sky approximations, as the combined corrections for these are an order-of-magnitude smaller than the reduced shear correction over the majority of the range of scales observed by *Euclid* (Kitching et al. 2017). This is again demonstrated in Table 1. The Limber approximation considers only correlations in the plane of the sky as contributing to the lensing signal, and projects others onto the plane of the sky by replacing spherical Bessel functions with Delta functions (Limber 1953; Kaiser 1998; LoVerde & Afshordi 2008).

Additionally, the Limber approximation is also employed when computing higher-order corrections to the angular power spectrum, such as the reduced shear correction. In this work, we also did not relax the use of the Limber approximation here, as the cosmological parameter biases from this are safely negligible (Deshpande & Kitching 2020).

Another series of corrections that we deemed safely negligible stem from corrections to the theoretical expressions describing light propagation (Cuesta-Lazaro et al. 2018). Among these is the effect of second-order corrections to the effective speed of light. This relaxes the assumption that, as lensing potentials are small, the lensing effect can be studied in an effective Minkowskian spacetime and, accordingly, the effective speed of light need only be computed to the first-order. Similarly, a second correction to the effective speed of light presents itself from the energy-momentum tensor. Typically, this quantity is calculated under the assumption that lenses are moving slowly, so that the kinetic contribution to gravity can be ignored. Addressing this creates another correction to the angular power spectrum.

A further effect is that the observed ellipticity is non-linearly related to the shear, but that a linear approximation is often made. For a discussion of the impact of the non-linear relation on the observed shear distribution see Viola et al. (2014). The impact of this effect on the power spectrum (if one assumes a linear relation instead of the non-linear relation) is investigated in Krause & Hirata (2010) (Section 3.3) who find that the impact is three orders of magnitude smaller than the shear power spectrum⁵.

The temporal-Born approximation is another correction to the description of light propagation. While the correction for the standard Born approximation accounts for the spatial discrepancy between the true perturbed path of a photon from source to observer compared to the mathematically convenient straight one, this discrepancy also produces a temporal one. The photon on the perturbed path will at times be ahead of the photon on the

⁵ In fact there are two definitions of ellipticity: third eccentricity and third flattening (see e.g. Viola et al. 2014) that relate the observed ellipticity to the shear in different ways. Krause & Hirata (2010) find that for third flattening the correction to the power spectrum is zero (since the moments of the third flattening are exactly the moments of the reduced shear as shown by Seitz & Schneider 1997), but that for third eccentricity the effect is non-zero.

idealised path, and at other times lag behind. Accordingly, the two would encounter different evolutionary stages of the LSS at different times, necessitating a correction in the two-point statistics. The remaining two corrections that fall under this umbrella are the corrections of the Sachs-Wolfe and integrated Sachs-Wolfe effects. The former describes the redshift of an emitted photon due to the source galaxy’s gravitational potential, while the latter encodes the effect on the photon of interaction with the evolving gravitational potential along its path. All of these light propagation corrections are many orders of magnitude below the reduced shear correction, as can be seen in a representative example in Table 1.

Finally, we deemed the flexion corrections to be negligible without requiring explicit calculation. This additional correction term arises from the fact that, for larger sources, the image distortion consists of both shear and a higher-order component labelled flexion (Schneider & Er 2008). This term should be negligible because its effect on the cosmic shear signal will be dependent on third-order or higher-order brightness moments.

Here, we did not consider biases arising from general modelling of unknown shape measurement systematic effects (i.e. multiplicative and additive biases). Instead, we focused only on these well defined theoretical assumptions. For more details on shape measurement effects, see Kitching et al. (2019, 2020, 2021); Kitching & Deshpande (2022). Additionally, we did not evaluate the additional selection effects of flux cuts and size cuts, as the former of these can be calibrated from deep fields, and size cuts are primarily a concern for ground-based telescopes, rather than space-based ones.

In this work, we also did not examine the impact of neglecting effects that are already well-established as requiring evaluation, i.e. photometric redshift uncertainties, intrinsic alignments (IA) modelling, baryonic feedback, and modelling of the non-linear component of the matter power spectrum. Determining the exact specification for these is outside of the scope of this work, and each of those effects requires its own through investigation.

This work is structured as follows: In Sect. 2, we detail the theoretical formalism used. We review the basic, first-order cosmic shear angular power spectrum calculation. Then, the expressions for the six correction terms of interest are detailed. We also describe the Fisher matrix formalism used to predict cosmological parameter constraints and biases. In Sect. 3, we discuss the modelling and computational specifics used in this work. Finally, we discuss our results in Sect. 4. We show the cosmological parameter biases that result from neglecting the studied corrections, and discuss their implications for *Euclid*.

2. Theoretical formalism

Here, we begin by reviewing the standard first-order calculation of the cosmic shear angular power spectrum. Additional contributions to the lensing signal resulting from IAs and shot-noise are then described. We then detail the analytical forms of the six corrections requiring full evaluation: reduced shear, magnification bias, source-lens clustering, source obscuration, local Universe effects, and the flat Universe assumption. Finally, we review the Fisher matrix formalism used to predict cosmological parameter constraints and biases.

2.1. The first-order cosmic shear calculation

As a consequence of weak gravitational lensing by the LSS, the observed ellipticity of distant galaxies is distorted. This change

Table 1. List of higher-order correction terms to the shear angular power spectrum resulting from relaxing approximations. To illustrate their typical sizes and facilitate comparison, the values of these correction terms at redshift $z \sim 1$ are also stated here. These are provided for ℓ -modes 10, 100, and 1000, and as a percentage of sample variance i.e. $100\delta C_\ell/(\text{sample variance})$. Sample variance here is calculated using the definition of Kaiser (1992), see equation (61). These values are taken from the available literature for the case of Stage IV cosmic shear experiments. Unavailable values are represented by ‘N/A’. Corrections which have a functional form that is fourth-order in lensing potential, and therefore sub-dominant, are denoted by $O(\phi^4)$. The references provided refer to the values stated where available, or to the work describing the correction where explicit values are not available. For the case where the absolute value of an effect is more than four orders of magnitude smaller than the sample variance, it is denoted by < 0.01 . The sign denotes whether the effect is to decrease the power (a negative sign) or to increase (no sign). The rows are ordered in decreasing amplitude for $\ell = 1000$. These numbers were read from graphs published in the referenced papers, due to the inherent inaccuracy of this approach we quote only one decimal place; in the case that only correlation functions were provide these were converted to power spectrum results using a Hankel transform over the quoted angular range.

Correction	Reference	Percentage of Sample Variance		
		$\ell = 10$	$\ell = 100$	$\ell = 1000$
Source-lens clustering	Yu et al. (2015)	N/A	6.0	57.2
Reduced shear + magnification bias	Deshpande et al. (2020a)	0.4	1.8	15.3
Post-Limber reduced shear	Deshpande & Kitching (2020)	0.2	0.6	1.9
Non-linear ellipticity-shear relation	Krause & Hirata (2010)	0.2	0.6	1.9
Limber + flat-sky	Kitching et al. (2017)	9.8	3.0	1.0
Local Universe effects	Hall (2020)	7.8	24.2	N/A
Higher-order reduced shear	Krause & Hirata (2010)	$O(\phi^4)$	$O(\phi^4)$	$O(\phi^4)$
Time delay-lens coupling	Bernardeau et al. (2010)	$O(\phi^4)$	$O(\phi^4)$	$O(\phi^4)$
Deflection-deflection coupling	Bernardeau et al. (2010)	$O(\phi^4)$	$O(\phi^4)$	$O(\phi^4)$
Born approximation	Cooray & Hu (2002)	$O(\phi^4)$	$O(\phi^4)$	$O(\phi^4)$
Lensing by dark energy pressure	Simpson et al. (2010)	$O(\phi^4)$	$O(\phi^4)$	$O(\phi^4)$
Second-order speed-of-light	Cuesta-Lazaro et al. (2018)	< 0.01	< 0.01	< 0.01
Temporal-Born approximation	Cuesta-Lazaro et al. (2018)	< 0.01	< 0.01	< 0.01
Finite-beam corrections	Fleury et al. (2019)	< 0.01	< 0.01	< 0.01
Doppler-shift	Deshpande & Kitching (2021)	< 0.01	< 0.01	< 0.01
Unequal-time correlators	Kitching & Heavens (2017)	< 0.01	< 0.01	< 0.01
Sachs-Wolfe effect	Cuesta-Lazaro et al. (2018)	< 0.01	< 0.01	< 0.01
Integrated Sachs-Wolfe effect	Cuesta-Lazaro et al. (2018)	< 0.01	< 0.01	< 0.01
Flexion correction	Schneider & Er (2008)	N/A	N/A	N/A
Flat-geometry assumption	Taylor et al. (2018b)	-2.0	-6.0	-19.1
Source obscuration	Hartlap et al. (2011)	-2.0	-6.0	-19.1
Spatially-varying survey depth	Heydenreich et al. (2020)	-5.9	-18.1	-57.2

is dependent on the reduced shear, g , according to

$$g^\alpha(\boldsymbol{\theta}) = \frac{\gamma^\alpha(\boldsymbol{\theta})}{1 - \kappa(\boldsymbol{\theta})}, \quad (1)$$

where $\boldsymbol{\theta}$ is the position of the galaxy on the sky, γ is the spin-2 shear with index $\alpha \in \{1, 2\}$ which describes the anisotropic stretching that turns circular distributions of light elliptical, and κ is the convergence – responsible for the isotropic change in the size of the image. Since in the weak lensing regime $|\kappa| \ll 1$, it is standard practice to make the reduced shear approximation, whereby

$$g^\alpha(\boldsymbol{\theta}) \approx \gamma^\alpha(\boldsymbol{\theta}). \quad (2)$$

Additionally, the convergence is a projection of the density contrast of the Universe, δ , along the line-of-sight over comoving distance, χ , to the comoving distance to the horizon, χ_h . For a particular tomographic redshift bin i , it is mathematically described by

$$\kappa_i(\boldsymbol{\theta}) = \int_0^{\chi_h} d\chi \delta[S_K(\chi)\boldsymbol{\theta}, \chi] W_i(\chi), \quad (3)$$

where S_K is a function that encodes the effect of the curvature of the Universe, K , on comoving distances according to

$$S_K(\chi) = \begin{cases} |K|^{-1/2} \sin(|K|^{-1/2}\chi) & K > 0 \text{ (closed Universe)} \\ \chi & K = 0 \text{ (flat Universe)} \\ |K|^{-1/2} \sinh(|K|^{-1/2}\chi) & K < 0 \text{ (open Universe)}. \end{cases} \quad (4)$$

We remind the reader that for the quantity $\delta[S_K(\chi)\boldsymbol{\theta}, \chi]$ in equation (3) the second χ means not only that there is an evaluation at a comoving radius χ , but also at a conformal time $\eta = \eta_0 - \chi$, meaning that all the integration over χ in this the paper are performed down the background light cone.

The $W_i(\chi)$ in Eq. (3) is the lensing projection kernel for tomographic bin i . It takes the form

$$W_i(\chi) = \frac{3}{2} \Omega_m \frac{H_0^2}{c^2} \frac{S_K(\chi)}{a(\chi)} \int_{\chi'}^{\chi_h} d\chi' n_i(\chi') \frac{S_K(\chi' - \chi)}{S_K(\chi')}, \quad (5)$$

which is dependent on the dimensionless present-day matter density of the Universe Ω_m , the speed of light in a vacuum c , the Hubble constant H_0 , the scale factor of the Universe $a(\chi)$, and the probability distribution of galaxies within redshift bin i $n_i(\chi)$.

The spin-2 shear is directly related to the convergence in spherical-harmonic space. For a specified lensing mass distribution, assuming the flat-sky and prefactor-unity approximations (Kitching et al. 2017), and under the small-angle limit, this relationship takes the form

$$\bar{\gamma}_i^\alpha(\boldsymbol{\ell}) = T^\alpha(\boldsymbol{\ell}) \bar{\kappa}_i(\boldsymbol{\ell}), \quad (6)$$

where $\boldsymbol{\ell}$ is the spherical-harmonic conjugate of $\boldsymbol{\theta}$, with magnitude ℓ and angular component ϕ_ℓ . The functions T^α are two trigonometric weighting functions corresponding to each of the

shear components. These take the form:

$$T^1(\boldsymbol{\ell}) = \cos(2\phi_\ell), \quad (7)$$

$$T^2(\boldsymbol{\ell}) = \sin(2\phi_\ell). \quad (8)$$

In the case of an arbitrary shear field, for example a field constructed from data, two linear combinations of the individual shear components are pertinent. Specifically, these are a divergence-free B -mode, and a curl-free E -mode:

$$\widetilde{E}_i(\boldsymbol{\ell}) = \sum_\alpha T^\alpha(\boldsymbol{\ell}) \widetilde{\gamma}_i^\alpha(\boldsymbol{\ell}), \quad (9)$$

$$\widetilde{B}_i(\boldsymbol{\ell}) = \sum_\alpha \sum_\beta \varepsilon^{\alpha\beta} T^\alpha(\boldsymbol{\ell}) \widetilde{\gamma}_i^\beta(\boldsymbol{\ell}). \quad (10)$$

Here, the summations are over the shear components, and $\varepsilon^{\alpha\beta}$ is the Levi-Civita symbol in the two-dimensional case; such that: $\varepsilon^{11} = \varepsilon^{22} = 0$ and $\varepsilon^{12} = -\varepsilon^{21} = 1$.

Assuming that higher-order systematic effects in the data have been accounted for, the B -mode of Eq. (10) vanishes. For the remaining E -mode, observables of interest are defined in the form of angular auto and cross-correlation power spectra, $C_{\ell;ij}^{\gamma\gamma}$, such that

$$\langle \widetilde{E}_i(\boldsymbol{\ell}) \widetilde{E}_j(\boldsymbol{\ell}') \rangle = (2\pi)^2 \delta_D^{(2)}(\boldsymbol{\ell} + \boldsymbol{\ell}') C_{\ell;ij}^{\gamma\gamma}, \quad (11)$$

where the angular brackets on the left-hand-side denote the ensemble average, which under the assumption of ergodicity becomes a spatial average, and $\delta_D^{(2)}$ is the Dirac delta for two dimensions. Under the extended Limber approximation (LoVerde & Afshordi 2008), where $k = (\ell + 1/2)/S_K(\chi)$, the power spectra themselves are further defined as

$$C_{\ell;ij}^{\gamma\gamma} = \int_0^{\chi_h} d\chi \frac{W_i(\chi) W_j(\chi)}{S_K^2(\chi)} P_{\delta\delta}(k, \chi), \quad (12)$$

where $P_{\delta\delta}$ is the three-dimensional matter power spectrum, and k is the magnitude of the spatial momentum vector \mathbf{k} which also shares the angular component ϕ_ℓ . Detailed reviews of this standard calculation can be found in Kilbinger (2015); Munshi et al. (2008); Bartelmann & Schneider (2001).

2.2. Intrinsic alignments and shot noise

When the angular power spectra are actually measured from surveys of galaxies, they contain non-lensing signals together with the pertinent cosmic shear power spectra. It is necessary to model each of these components to ensure accurate cosmological inference. A key non-lensing contribution arises from the fact that galaxies forming close to each other are forming in a similar tidal environment. Consequently, they have intrinsically correlated alignments (Joachimi et al. 2015; Kirk et al. 2015; Kiessling et al. 2015).

The observed ellipticity of an individual source, ϵ , can then, to first-order, be written as a combination of its underlying ellipticity in the absence of any cosmic shear or IA, ϵ^s , the cosmic shear, $\gamma = \gamma^1 + i\gamma^2$, and the effect of IA, ϵ^I according to

$$\epsilon = \epsilon^s + \gamma + \epsilon^I. \quad (13)$$

The angular power spectra corresponding to this observed ellipticity, $C_{\ell;ij}^{\epsilon\epsilon}$, are then the sum of contributions resulting from its components,

$$C_{\ell;ij}^{\epsilon\epsilon} = C_{\ell;ij}^{\gamma\gamma} + C_{\ell;ij}^{\gamma^1} + C_{\ell;ij}^{\gamma^2} + C_{\ell;ij}^{\text{II}} + N_{\ell;ij}^\epsilon, \quad (14)$$

in which $C_{\ell;ij}^{\gamma\gamma}$ are the cosmic shear angular power spectra defined in Eq. (12); in all cases the notation denotes $z_i \leq z_j$. The $C_{\ell;ij}^{\gamma^1}$ are the angular power spectra of correlations between foreground shear and background IA, which are only non-zero if photometric redshift estimates result in the scattering of observed redshifts between bins. On the other hand, the $C_{\ell;ij}^{\gamma^2}$ arise from the correlation between background shear and foreground IA, and the $C_{\ell;ij}^{\text{II}}$ represent the auto-correlation of the IA; both must be accounted for. To accomplish this, the non-linear alignment (NLA) model (Bridle & King 2007) can be employed. Under this model, these IA spectra take the form

$$C_{\ell;ij}^{\gamma^1} = \int_0^{\chi_h} \frac{d\chi}{S_K^2(\chi)} [W_i(\chi)n_j(\chi) + n_i(\chi)W_j(\chi)] P_{\delta\text{I}}(k, \chi), \quad (15)$$

$$C_{\ell;ij}^{\text{II}} = \int_0^{\chi_h} \frac{d\chi}{S_K^2(\chi)} n_i(\chi)n_j(\chi) P_{\text{II}}(k, \chi), \quad (16)$$

which, in a similar manner to the shear power spectra, are projections of three-dimensional IA power spectra, $P_{\delta\text{I}}$ and P_{II} . Both of these are related to the matter power spectrum as follows:

$$P_{\delta\text{I}}(k, \chi) = \left[-\frac{\mathcal{A}_{\text{IA}} C_{\text{IA}} \Omega_m}{D(\chi)} \right] P_{\delta\delta}(k, \chi), \quad (17)$$

$$P_{\text{II}}(k, \chi) = \left[-\frac{\mathcal{A}_{\text{IA}} C_{\text{IA}} \Omega_m}{D(\chi)} \right]^2 P_{\delta\delta}(k, \chi), \quad (18)$$

where the product of \mathcal{A}_{IA} and C_{IA} is a free parameter typically set by fitting to simulations or data, and $D(\chi)$ is the density perturbation growth factor. We note that the NLA model is a limited description of IAs, and accordingly has its own associated modelling uncertainties. Extensions of this model have been proposed (see e.g. Fortuna et al. 2021, EC20). However, investigating the modelling of IAs for *Euclid* in detail is out of the scope of this work, and necessitates a separate future investigation of itself.

Of the terms in Eq. (14), $N_{\ell;ij}^\epsilon$ remains; this shot noise term arises from the zero-lag autocorrelation of the unlensed, uncorrelated source ellipticity ϵ^s in Eq. (13) (see e.g. Hu 1999, equation 4). For a survey with equi-populated tomographic redshift bins, such as *Euclid* (EC20), this is expressed by

$$N_{\ell;ij}^\epsilon = \frac{\sigma_\epsilon^2}{\bar{n}_g / N_{\text{bin}}} \delta_{ij}^K, \quad (19)$$

within which σ_ϵ^2 is the variance of the observed ellipticities in the survey, \bar{n}_g is the surface density of galaxies in the survey, N_{bin} is the survey's number of tomographic redshift bins, and δ_{ij}^K is the Kronecker delta – which here indicates that the shot noise vanishes for cross-correlation spectra, as the ellipticities of galaxies at differing redshifts should not be correlated.

2.3. The reduced shear approximation

While relaxing the reduced shear approximation completely and explicitly is intractable, this can be sufficiently well modelled by applying a second-order Taylor expansion (Dodelson et al. 2006; Shapiro 2009; Krause & Hirata 2010; Deshpande et al. 2020a) to Eq. (1), resulting in

$$g^\alpha(\boldsymbol{\theta}) = \gamma^\alpha(\boldsymbol{\theta}) + (\gamma^\alpha \kappa)(\boldsymbol{\theta}) + \mathcal{O}(\kappa^3). \quad (20)$$

Computing the angular E -mode power spectra using this expanded expression results in the standard two-point expression of Eq. (11), plus three-point terms. These additional terms, $\delta \langle \tilde{E}_i(\boldsymbol{\ell}) \tilde{E}_j(\boldsymbol{\ell}') \rangle$, are given by

$$\begin{aligned} \delta \langle \tilde{E}_i(\boldsymbol{\ell}) \tilde{E}_j(\boldsymbol{\ell}') \rangle &= \sum_{\alpha} \sum_{\beta} T^{\alpha}(\boldsymbol{\ell}) T^{\beta}(\boldsymbol{\ell}') \langle (\widehat{\gamma^{\alpha\kappa}})_i(\boldsymbol{\ell}) \widehat{\gamma}_j^{\beta}(\boldsymbol{\ell}') \rangle \\ &+ T^{\alpha}(\boldsymbol{\ell}') T^{\beta}(\boldsymbol{\ell}) \langle (\widehat{\gamma^{\alpha\kappa}})_j(\boldsymbol{\ell}') \widehat{\gamma}_i^{\beta}(\boldsymbol{\ell}) \rangle \\ &= (2\pi)^2 \delta_{\text{D}}^{(2)}(\boldsymbol{\ell} + \boldsymbol{\ell}') \delta C_{\ell;ij}^{\text{RS}}, \end{aligned} \quad (21)$$

where $\delta C_{\ell;ij}^{\text{RS}}$ is the corresponding correction to $C_{\ell;ij}^{\gamma\gamma}$, and is given by

$$\delta C_{\ell;ij}^{\text{RS}} = \int_0^{\infty} \frac{d^2 \ell'}{(2\pi)^2} \cos(2\phi_{\ell'}) B_{ij}^{\text{KKK}}(\boldsymbol{\ell}, \boldsymbol{\ell}', -\boldsymbol{\ell} - \boldsymbol{\ell}'), \quad (22)$$

where we are always free to choose a coordinate system such that $\phi_{\ell} = 0$, and accordingly the correction only depends on the magnitude, ℓ . It depends on the two-redshift convergence bispectrum, B_{ij}^{KKK} , which is the three-point counterpart of the convergence power spectrum. Higher-order terms in the Taylor expansion of Eq. (20) would here result in corrections dependent on the matter trispectrum, as well as Wick contraction terms of $\mathcal{O}(P_{\delta\delta}^2)$. Both types of terms have been shown to be sub-dominant (Cooray & Hu 2002; Shapiro & Cooray 2006; Dodelson et al. 2006; Krause & Hirata 2010; Deshpande et al. 2020a). The latter type of term, although of the same perturbative order in the power spectrum as the bispectrum, is still of $\mathcal{O}(W(\chi)^4)$, and given that typically $\chi W(\chi) \ll 1$, it will still be significantly smaller than the correction of Eq. (22).

Additionally, just as the convergence power spectrum is the projection of the matter power spectrum, the convergence bispectrum is analogously the projection of the matter bispectrum, $B_{\delta\delta\delta}$. Under the Limber approximation, this takes the form

$$\begin{aligned} B_{ij}^{\text{KKK}}(\boldsymbol{\ell}_1, \boldsymbol{\ell}_2, \boldsymbol{\ell}_3) &= B_{ii}^{\text{KKK}}(\boldsymbol{\ell}_1, \boldsymbol{\ell}_2, \boldsymbol{\ell}_3) + B_{ijj}^{\text{KKK}}(\boldsymbol{\ell}_1, \boldsymbol{\ell}_2, \boldsymbol{\ell}_3) \\ &= \int_0^{\chi_{\text{h}}} \frac{d\chi}{S_{\text{K}}^4(\chi)} W_i(\chi) W_j(\chi) [W_i(\chi) + W_j(\chi)] \\ &\times B_{\delta\delta\delta}(\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3, \chi). \end{aligned} \quad (23)$$

For a relaxation of the Limber approximation see Deshpande & Kitching (2020).

It should also be noted that the use of the reduced shear approximation can produce a B -mode signal contribution. However, it has been demonstrated that this is negligible (Schneider et al. 2002).

2.4. Source-lens clustering

Since, in practice, cosmic shear is only measured where galaxies are present, care must be taken to account for biases from any correlations between background source galaxies and the foreground lensing field. Given that, in reality, tomographic bins must be wide enough to include a sufficient number of galaxies so that shape-measurement noise is minimised, there will be overlap between the source and lensing distributions. The situation is further aggravated by broadening of bins due to photometric redshift uncertainties.

As a consequence of this effect, the observed number density of galaxies used in a given estimator which determines the shear angular power spectra from data is correlated with the intrinsic source galaxy overdensity, δ_i^g , such that (Bernardeau 1998;

Hamana et al. 2002; Schmidt et al. 2009)

$$n_i^{\text{obs}}(\boldsymbol{\theta}, \chi) = n_i(\chi) [1 + \delta_i^g(\boldsymbol{\theta})]. \quad (24)$$

Accordingly, the shear used in the theoretical formalism for inference, is similarly replaced with an ‘observed’ shear,

$$\gamma_{\text{obs};i}^{\alpha}(\boldsymbol{\theta}) = \gamma_i^{\alpha}(\boldsymbol{\theta}) + \gamma_i^{\alpha}(\boldsymbol{\theta}) \delta_i^g(\boldsymbol{\theta}). \quad (25)$$

This is similar in form to the Taylor expansion of the reduced shear expressed in Eq. (20), and results in an analogous correction term, $\delta C_{\ell;ij}^{\text{SLC}}$, to the angular power spectra,

$$\delta C_{\ell;ij}^{\text{SLC}} = \int_0^{\infty} \frac{d^2 \ell'}{(2\pi)^2} \cos(2\phi_{\ell'}) B_{ij}^{\text{K}\delta\text{K}}(\boldsymbol{\ell}, \boldsymbol{\ell}', -\boldsymbol{\ell} - \boldsymbol{\ell}'), \quad (26)$$

where $B_{ij}^{\text{K}\delta\text{K}}$ is now the two-redshift convergence-galaxy bispectrum. By adopting a linear galaxy bias model (so that $\delta_g = b\delta$) as used in EC20, and noting that δ^g is the 2D projection of δ_g , the convergence-galaxy bispectrum can also be expressed as a projection of the matter bispectrum

$$\begin{aligned} B_{ij}^{\text{K}\delta\text{K}}(\boldsymbol{\ell}_1, \boldsymbol{\ell}_2, \boldsymbol{\ell}_3) &= B_{ii}^{\text{K}\delta\text{K}}(\boldsymbol{\ell}_1, \boldsymbol{\ell}_2, \boldsymbol{\ell}_3) + B_{ijj}^{\text{K}\delta\text{K}}(\boldsymbol{\ell}_1, \boldsymbol{\ell}_2, \boldsymbol{\ell}_3) \\ &= \int_0^{\chi_{\text{h}}} \frac{d\chi}{S_{\text{K}}^4(\chi)} [b_i n_i(\chi) + b_j n_j(\chi)] W_i(\chi) W_j(\chi) \\ &\times B_{\delta\delta\delta}(\mathbf{k}_1, \mathbf{k}_2, \mathbf{k}_3, \chi), \end{aligned} \quad (27)$$

where b_i and b_j are the galaxy biases for tomographic bins i and j , respectively. While more complex models of the galaxy bias exist, we proceed with the linear bias in this work, in order to mitigate the already significant computational load of these three-point terms. We note that modelling the galaxy bias requires more complexity at smaller scales, where the SLC effect is most relevant. When ultimately computing this term in the *Euclid* cosmological analysis, the final *Euclid* galaxy bias model should be used. The linear galaxy bias for each tomographic bin is given by

$$b_i = \sqrt{1 + \bar{z}_i}, \quad (28)$$

where \bar{z}_i is tomographic bin i ’s central redshift. For a review of galaxy bias models, see Desjacques et al. (2018).

In addition to this contribution to the E -mode angular power spectra, source-lens clustering produces a B -mode signal as well. This term is comparable to the E -mode correction in magnitude, and accordingly, its detection in the absence of other B -mode contributions could allow for direct correction of the E -mode signal, rather than requiring the computation of Eq. (26). However, typical B -mode signals are dominated by other contributions (Schneider et al. 2002; Yu et al. 2015).

2.5. Magnification bias

An additional consequence of gravitational lensing is that the density of galaxies observed by a particular survey is no longer representative of the true underlying galaxy density (Turner et al. 1984). In particular, magnification resulting from the convergence modifies the density in two contrasting ways.

One manifestation of the effect is that individual sources are magnified, and as a consequence of this, their flux increases. Accordingly, any sources lying just beyond the flux limit of the survey may have their fluxes increased to the point of then being within the flux limit; increasing the observed density. As sources are magnified, the patch of sky around them too is magnified.

This causes the second, competing manifestation. Within the magnified patch of sky, the galaxy density is reduced. The total effect, known as magnification bias, is dependent on the slope of the unlensed galaxy luminosity function. This assumes that the magnification $\mu > 1$.

Assuming that, on our scales of interest, fluctuations in the intrinsic galaxy overdensity are small, and taking into account that, for weak lensing, $|\kappa| \ll 1$, the observed galaxy overdensity for a given tomographic bin, $\delta_{\text{obs};i}^g$, is given by (Hui et al. 2007; Schmidt et al. 2009)

$$\delta_{\text{obs};i}^g(\boldsymbol{\theta}) = \delta_i^g(\boldsymbol{\theta}) + (5s_i - 2)\kappa_i(\boldsymbol{\theta}), \quad (29)$$

where δ_i^g is the intrinsic galaxy overdensity in the absence of magnification or any other systematic effects, and s_i is the slope of the luminosity function for redshift bin i . This is given by the derivative of the cumulative galaxy number counts with respect to magnitude, m , evaluated at the survey's limiting magnitude, m_{lim} such that

$$s_i = \left. \frac{\partial \log_{10} n(\bar{z}_i, m)}{\partial m} \right|_{m_{\text{lim}}}, \quad (30)$$

in which $n(\bar{z}_i, m)$ is the true, underlying distribution of galaxies, evaluated at the tomographic bin's central redshift, \bar{z}_i . Here, we have suppressed an additional dependence on the wavelength band in which the galaxy is observed. This should be considered when determining the slope from observational data.

Accordingly, Eq. (25) gains an extra term

$$\gamma_{\text{obs};i}^\alpha = \gamma_i^\alpha(\boldsymbol{\theta}) + \gamma_i^\alpha(\boldsymbol{\theta}) \delta_i^g(\boldsymbol{\theta}) + (5s_i - 2)\gamma_i^\alpha(\boldsymbol{\theta})\kappa_i(\boldsymbol{\theta}). \quad (31)$$

This additional term is near-identical to the second term in Eq. (20), but for the prefactor of $(5s_i - 2)$. Accordingly, it too spawns a correction to the angular power spectra. This correction for magnification bias, $\delta C_{\ell;ij}^{\text{MB}}$, takes a similar form to the reduced shear correction of Eq. (22),

$$\delta C_{\ell;ij}^{\text{MB}} = \int_0^\infty \frac{d^2 \ell'}{(2\pi)^2} \cos(2\phi_{\ell'}) [(5s_i - 2)B_{ij}^{\text{KKK}}(\boldsymbol{\ell}, \boldsymbol{\ell}', -\boldsymbol{\ell} - \boldsymbol{\ell}') + (5s_j - 2)B_{ij}^{\text{KKK}}(\boldsymbol{\ell}, \boldsymbol{\ell}', -\boldsymbol{\ell} - \boldsymbol{\ell}')]. \quad (32)$$

Given the similarity of the magnification bias correction to the source-lens clustering and reduced shear corrections, it too would produce a contribution to the B -mode signal. While this has not been explicitly evaluated, we would expect this term, as with its reduced shear and source-lens clustering counterparts, to be sub-dominant.

2.6. Source obscuration

There is another systematic effect which can change the observed galaxy number density. Blending of close galaxy pairs can lead to multiple galaxies being discarded, or counted as a lower number than they are (Hartlap et al. 2011). The resulting change in the observed number density of galaxies, $\Delta n^{\text{SO}}(z, \boldsymbol{\theta})$, can be modelled by

$$\frac{\Delta n^{\text{SO}}(z, \boldsymbol{\theta})}{n(z, \boldsymbol{\theta})} = -\pi \left[(2\vartheta)^2 n_{\text{tot}} + \frac{A n(z)}{2 - \zeta} (2\vartheta)^{2-\zeta} \right], \quad (33)$$

where n_{tot} is the total number density of galaxies at all redshifts, $n(z)$ is the observed density of galaxies at redshift z ignoring source-lens clustering, we assume a redshift-independent radius ϑ for all galaxies as in Hartlap et al. (2011), and A and ζ are the

amplitude and power-law index of a power-law model for the two-point galaxy angular correlation function. This expression is obtained by considering the probability that the centroid of another source lies within 2ϑ of a given one by integrating over the probability that another source centroid lies in an annulus of $d\theta$ around the centroid of another one.

Instead, we assume that the blending strategy for *Euclid* will account for blended pairs sufficiently well, such that the only obscuration of concern is substantial overlap; when the centroid of a source is behind another source (i.e. within ϑ rather than 2ϑ). In this case, we are only concerned with the probability of this overlap. We note that in reality this blending has a complex interaction with shape measurement, but evaluating this is out of the scope of this work. Then, assuming that sources are approximately circular, the probability, dP , of a galaxy at redshift z overlapping with one at redshift z' is

$$dP(z, z', \boldsymbol{\theta}) = \pi \vartheta^2 n(z', \boldsymbol{\theta}) dz'. \quad (34)$$

Here it is also assumed that the expected number of galaxies overlapping with a given galaxy is $\ll 1$, such that the probability of at least one galaxy overlapping a given source (resulting in the removal of that source from the sample) is equal to the probability of just one overlap, which by Poisson statistics is then the expected number of overlaps. Accordingly, the total change in the number of sources at z is then given by

$$\begin{aligned} \Delta n^{\text{SO}}(z, \boldsymbol{\theta}) &= -\pi \vartheta^2 n(z, \boldsymbol{\theta}) \int_0^\infty dz' n(z', \boldsymbol{\theta}) \\ &= -\pi \vartheta^2 n(z, \boldsymbol{\theta}) [1 + \delta^g(\boldsymbol{\theta})] \int_0^\infty dz' n(z') \\ &= -\pi \vartheta^2 n(z, \boldsymbol{\theta}) [1 + \delta^g(\boldsymbol{\theta})] n_{\text{tot}}. \end{aligned} \quad (35)$$

We neglect the second term on the right-hand side of Eq. (33) as it specifically accounts for the correlated overlap of galaxies at the same redshift within a fixed disk around the source, in addition to the random one already included. Given that the fractional change is calculated by integrating over redshift slices, that this term would only appear for the slice where $z' = z$, and that the source obscuration term itself is small (see Sect. 4), we expect it to be safely negligible.

Adopting the tomographic redshift binning approach, for a given source can be obscured by another in the same bin, or by sources in lower redshift bins than the one the source belongs in. The fractional change in the number density of galaxies in redshift bin i then becomes

$$\begin{aligned} \frac{\Delta n_i^{\text{SO}}(z, \boldsymbol{\theta})}{n_i(z, \boldsymbol{\theta})} &= -\pi \vartheta^2 \sum_{q=1}^i [1 + \delta_q^g(\boldsymbol{\theta})] n_{\text{tot};q} \\ &= -\pi \vartheta^2 \sum_{q=1}^i [1 + b_q \delta^g(\boldsymbol{\theta})] n_{\text{tot};q} \\ &= -\pi \vartheta^2 n_{\text{cumul};i} - \pi \vartheta^2 \delta^g f_{\text{SO};i}, \end{aligned} \quad (36)$$

where the $n_{\text{tot};q}$ is the total surface density of galaxies for redshift bin q (that is the integral over $n_q(z)$ for bin q), $n_{\text{cumul};i}$ is the cumulative total surface density of galaxies for all redshift bins up to and including bin i , and

$$f_{\text{SO};i} = \sum_{\rho=1}^i b_\rho n_{\text{tot};\rho}. \quad (37)$$

Including the effect of source obscuration in addition to source-lens clustering and magnification bias, the observed number density for a given tomographic redshift bin i , becomes

$$n_i^{\text{obs}}(\boldsymbol{\theta}, \chi) = n_i(\chi) [1 + (5s_i - 2)\kappa_i(\boldsymbol{\theta}) + \delta_i^g(\boldsymbol{\theta})] \times [1 - \pi \vartheta^2 n_{\text{cumul.};i} - \pi \vartheta^2 \delta^g(\boldsymbol{\theta}) f_{\text{SO};i}]. \quad (38)$$

From here, only terms to first-order in the lensing potential are retained in order to suppress fourth-order or higher terms appearing in the two-point statistic. Then, it can be seen that source obscuration adds prefactors to the base angular shear power spectra, the source-lens clustering correction from Eq. (26), and the magnification bias correction of Eq. (32). Accordingly, source obscuration produces three new correction terms:

$$\delta C_{\ell;ij}^{\text{SO}} = (\pi^2 \vartheta^4 n_{\text{cumul.};i} n_{\text{cumul.};j} - \pi \vartheta^2 n_{\text{cumul.};i} - \pi \vartheta^2 n_{\text{cumul.};j}) C_{\ell;ij}^{\gamma\gamma}, \quad (39)$$

$$\delta C_{\ell;ij}^{\text{SO-SLC}} = \int_0^\infty \frac{d^2 \boldsymbol{\ell}'}{(2\pi)^2} \cos(2\phi_{\ell'}) \times \left[(\pi^2 \vartheta^4 n_{\text{cumul.};i} n_{\text{cumul.};j} - \pi \vartheta^2 n_{\text{cumul.};i} b_i - \pi \vartheta^2 n_{\text{cumul.};j} b_j - \pi \vartheta^2 (1 - n_{\text{cumul.};j}) \times f_{\text{SO};i} B_{ij}^{\kappa\delta\kappa}(\boldsymbol{\ell}_1, \boldsymbol{\ell}_2, \boldsymbol{\ell}_3) + i \leftrightarrow j \right], \quad (40)$$

$$\delta C_{\ell;ij}^{\text{SO-MB}} = (\pi^2 \vartheta^4 n_{\text{cumul.};i} n_{\text{cumul.};j} - \pi \vartheta^2 n_{\text{cumul.};i} - \pi \vartheta^2 n_{\text{cumul.};j}) \delta C_{\ell;ij}^{\text{MB}}, \quad (41)$$

where $i \leftrightarrow j$ indicates a repetition of the preceding bispectrum term and its pre-factor, with all instances of the i and j bin indices exchanged.

2.7. Local Universe effects

A further effect to consider is that the observed two-point statistic at our location may be biased due to local over or underdensities. Accordingly, the angular power spectra must be calculated conditioned on the local density (Hall 2020). The local density contrast, δ_0 , can be defined as the matter density contrast smoothed by a top-hat kernel of comoving radius R according to

$$\delta_0(R, \chi) \equiv \frac{3}{4\pi R^3} \int d^3 \mathbf{r} \Theta(R - |\mathbf{r}|) \delta(\mathbf{r}, \chi), \quad (42)$$

where the matter density contrast is now expressed in terms of spatial distance, \mathbf{r} , rather than angle on the sky, and Θ is the Heaviside step-function.

Then, the conditional angular power spectra can be obtained using the Edgeworth expansion for conditional distributions. This calculation is mathematically intensive, and accordingly is not reproduced here. The full derivation can be found in Hall (2020). Under the Limber approximation, and assuming $\ell \gg 1$ (and note that cosmic shear is only defined for $\ell \geq 2$), this expression consists of two terms, the standard power spectra of

Eq. (12) and a correction term, $\delta C_{\ell;ij}^{\text{LU}}$ which is defined as

$$\delta C_{\ell;ij}^{\text{LU}} = 2 \frac{\delta_0(R, \chi)}{\sigma^2(R, \chi)} \int_0^{\chi_h} d\chi \frac{W_i(\chi) W_j(\chi)}{\chi^2} \times \left\{ \left[\frac{34}{21} \xi_R(\chi) - \frac{4}{21} \psi_R(\chi) \right] P_{\delta\delta}(k, \chi) + \left[\frac{\chi}{\ell} \xi'_R(\chi) - \frac{\ell}{\chi} \Omega_R(\chi) \right] \frac{\partial P_{\delta\delta}(k, \chi)}{\partial k} \frac{1}{\chi} - \frac{4}{7} \psi_R(\chi) \frac{\partial^2 P_{\delta\delta}(k, \chi)}{\partial k^2} \frac{1}{\chi^2} \right\}, \quad (43)$$

where σ^2 is the variance of the local density contrast, it is assumed the ratio of the local density contrast to its variance is constant with comoving distance, ξ_R , ψ_R , ξ'_R , and Ω_R are correlation functions defined in Hall (2020), and the expression assumes a flat-geometry, which is valid under current constraints on Ω_K , as lenses are much less than curvature distance away. Additionally, we note that this expression is derived using only the tree-level Eulerian perturbation theory expression for the matter bispectrum.

2.8. The flat Universe assumption

Typically when computing cosmic shear angular power spectra, spatially non-flat universes are accounted for through modifying comoving distances based on curvature, as described by Eq. (4). In practice, however, curvature also modifies the projection kernel (Taylor et al. 2018b).

Under the assumption of a spatially flat Universe, the Poisson equation gives the relationship between the comoving Newtonian gravitational potential, ϕ , and the matter density contrast,

$$\nabla_\chi^2 \phi(\mathbf{r}, \chi) = \frac{3\Omega_m H_0^2}{2c^2 a(t)} \delta(\mathbf{r}, \chi), \quad (44)$$

where ∇_χ^2 is the Laplacian for a spatially flat Universe. This allows the shear angular power spectra to be expressed in terms of the matter power spectrum, as in Eq. (12). However, the matter density contrast has rectilinear coordinates, whereas the lensing potential is defined in terms of angular coordinates (r, θ, φ) , from the observer's frame of reference. Relating the two as above requires expressing the potential in spherical Bessel space as

$$\phi_{\ell m}(k) = \sqrt{\frac{2}{\pi}} \int d^3 r \phi(r) j_\ell(kr) Y_{\ell m}(\theta, \varphi), \quad (45)$$

where j_ℓ are spherical Bessel functions, and $Y_{\ell m}$ are spherical harmonics. Owing to the fact that these spherical Bessel functions and spherical harmonics are eigenfunctions of the Laplacian, the following relationship is obtained:

$$(\nabla_r^2 + k^2) j_\ell(kr) Y_{\ell m}(\theta, \varphi) = 0. \quad (46)$$

This allows the relation of the lensing potential to the matter density contrast in spherical harmonic space and, under the Limber approximation, the eventual calculation of Eq. (12). See Kitching et al. (2017) for a full derivation.

However, in the case of a spatially non-flat Universe, the Laplacian in Eqs. (44 – 46) must be replaced by one corresponding to a curved geometry, $\nabla_{S_K}^2$. Accordingly, the projection kernel must also be modified; by replacing the spherical Bessel

functions in Eq. (46) with hyper-spherical Bessel functions, Φ_ℓ^β , so that

$$\left(\nabla_{S_K}^2 + k^2\right) \Phi_\ell^\beta(r) Y_{\ell m}(\theta, \varphi) = 0, \quad (47)$$

where $\beta = \sqrt{(k^2 + K)/|K|}$ (Lesgourgues & Tram 2014). Consequently, the shear angular power spectra for a spatially non-flat Universe, $C_{\ell;ij}^{\gamma\gamma;\text{NF}}$, under the Limber approximation, under the Limber approximation, is given by modifying equation (12) to be

$$C_{\ell;ij}^{\gamma\gamma} = \int_0^{\chi_h} d\chi W_\ell^{\text{NF}}(\chi; K) \frac{W_i(\chi) W_j(\chi)}{S_K^2(\chi)} P_{\delta\delta}(k, \chi), \quad (48)$$

where

$$W_\ell^{\text{NF}}(\chi; K) = \left[1 - \text{sgn}(K) \frac{\ell^2}{(\ell + 1/2)^2 / S_K^2(\chi) + K}\right]^{-1/2} \quad (49)$$

where $\text{sgn}(K)$ is the sign of the curvature K . Alternatively, for consistency with the previously discussed corrections, this can be expressed as a correction term, $\delta C_{\ell;ij}^{\text{NF}}$, to the spatially flat angular power spectra such that

$$\delta C_{\ell;ij}^{\text{NF}} = C_{\ell;ij}^{\gamma\gamma;\text{NF}} - C_{\ell;ij}^{\gamma\gamma}. \quad (50)$$

2.9. The Fisher matrix and bias formalism

The constraining power of cosmological surveys, in terms of uncertainties on inferred cosmological parameters, is often predicted by using the Fisher matrix formalism. It also allows the quantification of biases in this inference resulting from neglecting systematic effects within the signal itself. Here, we use this technique to predict how biased cosmological parameters inferred from *Euclid* would be when the previously discussed systematic effects are neglected.

Explicitly, the Fisher matrix is defined as the expected value of the Hessian of the log likelihood (defined for a Gaussian likelihood, and applied to CMB data in Tegmark et al. 2015). For Stage IV weak lensing cosmology, it has been demonstrated that the likelihood can safely be assumed to be Gaussian (Lin et al. 2019; Taylor et al. 2019a; Upham et al. 2021; Hall & Taylor 2022). Accordingly, the Fisher matrix for cosmic shear is defined as

$$F_{\mu\nu} = \sum_{\ell'=\ell_{\min}}^{\ell_{\max}} \sum_{\ell=\ell_{\min}}^{\ell_{\max}} \sum_{ij, mn} \frac{\partial C_{\ell;ij}^{\text{NF}}}{\partial \theta_\mu} \text{Cov}^{-1} [C_{\ell;ij}^{\text{NF}}, C_{\ell';mn}^{\text{NF}}] \frac{\partial C_{\ell';mn}^{\text{NF}}}{\partial \theta_\nu}, \quad (51)$$

where the μ and ν indices denote element in the Fisher matrix associated with cosmological parameters θ_μ and θ_ν , respectively, ℓ_{\min} is the minimum angular wavenumber of the survey, ℓ_{\max} is the maximum angular wavenumber used, the sums are over the ℓ -blocks of power spectrum bands, and $\text{Cov}^{-1} [C_{\ell;ij}^{\text{NF}}, C_{\ell';mn}^{\text{NF}}]$ is the inverse of the covariance of the angular power spectra signal.

In practice, this covariance term is non-Gaussian (Barreira et al. 2018a; Takada & Hu 2013; Upham et al. 2022), with an additional contribution arising from the super-sample covariance (SSC; Hu & Kravtsov 2003). This SSC term encapsulates the effects on the covariance of density fluctuations with wavelengths larger than the extent of the galaxy survey. Such fluctuations result in the background density of the survey ceasing to be representative of the underlying density of the Universe. The

total covariance is then the sum of the Gaussian, Cov_G and SSC, Cov_{SSC} , terms:

$$\text{Cov} [C_{\ell;ij}^{\text{NF}}, C_{\ell';mn}^{\text{NF}}] = \text{Cov}_G [C_{\ell;ij}^{\text{NF}}, C_{\ell';mn}^{\text{NF}}] + \text{Cov}_{\text{SSC}} [C_{\ell;ij}^{\text{NF}}, C_{\ell';mn}^{\text{NF}}], \quad (52)$$

where the Gaussian component is given by

$$\text{Cov}_G [C_{\ell;ij}^{\text{NF}}, C_{\ell';mn}^{\text{NF}}] = \frac{C_{\ell;im}^{\text{NF}} C_{\ell';jn}^{\text{NF}} + C_{\ell;in}^{\text{NF}} C_{\ell';jm}^{\text{NF}}}{(2\ell + 1) f_{\text{sky}} \Delta\ell} \delta_{\ell\ell'}^K, \quad (53)$$

where f_{sky} is the fraction of the sky observed by the galaxy survey, $\Delta\ell$ is the bandwidth of the ℓ -modes sampled, and $\delta_{\ell\ell'}^K$ is the Kronecker delta. Other non-Gaussian terms in the covariance can be neglected (see e.g. Barreira et al. 2018b). The SSC component is well-approximated by (Lacasa & Grain 2019)

$$\text{Cov}_{\text{SSC}} [C_{\ell;ij}^{\text{NF}}, C_{\ell';mn}^{\text{NF}}] \approx R_\ell C_{\ell;ij}^{\text{NF}} R_{\ell'} C_{\ell';mn}^{\text{NF}} S_{ijmn}, \quad (54)$$

where S_{ijmn} is the dimensionless volume-averaged covariance of the background matter density contrast, and R_ℓ is the effective relative response of the observed power spectrum. We assume that there is no interrelation between local Universe effects and the SSC, but this is a caveat that should be verified in future.

The diagonal of the inverse of the Fisher matrix is used to predict the 1σ uncertainties on each of the parameters. Explicitly, the uncertainty, σ_μ , on parameter θ_μ is given by

$$\sigma_\mu = \sqrt{F_{\mu\mu}^{-1}}. \quad (55)$$

By extending this formalism, the biases on inferred parameters resulting from neglecting systematic effects can also be predicted (Taylor et al. 2007). For a given systematic, $\delta C_{\ell;ij}$, the bias, b_μ , on parameter θ_μ , is given by

$$b_\mu = \sum_\nu (F^{-1})_{\mu\nu} \mathcal{B}_\nu, \quad (56)$$

where

$$\mathcal{B}_\nu = \sum_{\ell'=\ell_{\min}}^{\ell_{\max}} \sum_{\ell=\ell_{\min}}^{\ell_{\max}} \sum_{ij, mn} \delta C_{\ell;ij} \text{Cov}^{-1} [C_{\ell;ij}^{\text{NF}}, C_{\ell';mn}^{\text{NF}}] \frac{\partial C_{\ell';mn}^{\text{NF}}}{\partial \theta_\nu}. \quad (57)$$

We note that we assume a Gaussian likelihood function but with a correlated covariance matrix (that includes non-Gaussian contributions). The extent to which this assumption is robust to relaxing the Gaussian likelihood assumption was explored in Martinelli et al. (2021) who found good agreement between Fisher matrix (Gaussian likelihood) predictions and full MCMC predictions, and in Taylor et al. (2019b) who found a similar result but allowing for the possibility of a fully non-Gaussian likelihood function.

3. Methodology

In this section, we review the computational and modelling specifics used within this investigation. We begin by describing the survey specifications adopted. Then, details are given about our choice of fiducial cosmology, modelling of background quantities, and Fisher matrices. Lastly, we describe modelling choices made in the computation of the magnification bias, source obscuration, and local Universe effects corrections.

3.1. Survey specifications

For *Euclid*, forecasting specifications are specified in EC20; we adopted these here, but we note that our IA model is more simple (see Section 2.2), and we vary Ω_K rather than dark energy critical density. Specifically, we considered the ‘optimistic’ scenario described in that work, as this is the case where the cosmic shear probe is able to meet its precision goals by itself. Under this scenario, the survey is taken to extend up to ℓ -modes of 5000.

Additionally, the intrinsic variance of observed ellipticities is taken to consist of two components; each with a magnitude of 0.21. Correspondingly, the RMS intrinsic ellipticity variance is $\sigma_\epsilon = \sqrt{2} \times 0.21 \approx 0.3$ ⁶. *Euclid* is also expected to have a survey area such that $f_{\text{sky}} = 0.36$. The survey’s galaxy surface density is anticipated to be $\bar{n}_g = 30 \text{ arcmin}^{-2}$.

The cosmic shear probe of *Euclid* is planned to observe sources between redshifts of 0 and 2.5, and utilise 10 equi-populated tomographic redshift bins with the following edges: {0.001, 0.418, 0.560, 0.678, 0.789, 0.900, 1.019, 1.155, 1.324, 1.576, 2.50}.

Given that *Euclid* will use photometric redshifts, the model for the source distributions within these tomographic bins must account for photometric redshift uncertainties. Accordingly, for a particular bin, i , the galaxy redshift distribution, $n_i(z)$, was described by

$$n_i(z) = \frac{\int_{z_i^-}^{z_i^+} dz_p n(z) p_{\text{ph}}(z_p|z)}{\int_{z_{\text{min}}}^{z_{\text{max}}} dz \int_{z_i^-}^{z_i^+} dz_p n(z) p_{\text{ph}}(z_p|z)}, \quad (58)$$

where z_p is measured photometric redshift, z_i^- and z_i^+ are the limits of the i -th redshift bin, and z_{min} and z_{max} are the redshift limits of the survey itself. Additionally, $n(z)$ is the underlying distribution of galaxies which here we modeled according to the formalism established in Laureijs et al. (2011):

$$n(z) \propto \left(\frac{z}{z_0}\right)^2 \exp\left[-\left(\frac{z}{z_0}\right)^{3/2}\right], \quad (59)$$

where $z_0 = z_m/\sqrt{2}$, and $z_m = 0.9$ is the median redshift of the survey. The remaining function in Eq. (58), $p_{\text{ph}}(z_p|z)$, encapsulates the probability that a source measured to have a photometric redshift of z_p actually has a redshift of z . This distribution

⁶ We use the specification in EC20, but note that Euclid Collaboration et al. (2019) uses a value of 0.26 per component. Since we are looking at biases caused by the differences in the signal the shot noise component does not affect the reported biases, however the relative significance (bias divided by error) will be lower for a larger shot noise term.

Table 2. Values of model parameters used in defining the uncertainty of photometric redshift estimates through Eq. (60). Chosen according to EC20.

Parameter	Fiducial Value
c_b	1.0
z_b	0.0
σ_b	0.05
c_o	1.0
z_o	0.1
σ_o	0.05
f_{out}	0.1

Table 3. Λ CDM and w_0w_a CDM cosmological parameter fiducial values used in this investigation. These values correspond to EC20. It should be noted that $\sum m_\nu \neq 0$ is assumed to be fixed, and uncertainties and biases are not calculated for it. Additionally, two possible values are provided for Ω_K , because the non-zero value must be used when evaluating the non-flat Universe correction. This value is selected using the upper-bound of the Planck Collaboration et al. (2018) 1σ uncertainty.

Cosmological Parameter	Fiducial Value
Ω_m	0.32
Ω_b	0.05
h	0.67
n_s	0.96
σ_8	0.816
w_0	-1
w_a	0
$\sum m_\nu$ (eV)	0.06
Ω_K	{0, 0.05}

takes the form (Kitching et al. 2008)

$$p_{\text{ph}}(z_p|z) = \frac{1 - f_{\text{out}}}{\sqrt{2\pi}\sigma_b(1+z)} \exp\left\{-\frac{1}{2}\left[\frac{z - c_b z_p - z_b}{\sigma_b(1+z)}\right]^2\right\} + \frac{f_{\text{out}}}{\sqrt{2\pi}\sigma_o(1+z)} \exp\left\{-\frac{1}{2}\left[\frac{z - c_o z_p - z_o}{\sigma_o(1+z)}\right]^2\right\}. \quad (60)$$

Here, the distribution is expressed as the sum of two terms – the first is the uncertainty resulting from multiplicative and additive bias in redshift determination for the fraction of sources with a well measured redshift, whilst the second represents the same, but for a fraction of catastrophic outliers in the sample, f_{out} . The values used for the individual parameters in this parameterisation match the selection of EC20, and are stated in Table 2, which are fixed throughout our analysis.

3.2. Cosmological modelling and Fisher matrices

Throughout this investigation, we considered the Λ CDM cosmological model and its extension: the w_0w_a CDM model, which also allows for varying dark energy pressure and a separately parameterised dark energy equation of state at early times. The Λ CDM model uses 7 parameters, which are defined thusly: the present-day total matter density parameter Ω_m , the present-day baryonic matter density parameter Ω_b , the dimensionless curvature parameter $\Omega_K = -K(c/H_0)^2$, the Hubble parameter $h = H_0/100 \text{ km s}^{-1} \text{ Mpc}^{-1}$, the spectral index n_s , the RMS value of density fluctuations on $8 h^{-1} \text{ Mpc}$ scales σ_8 , and massive neutrinos with a sum of masses $\sum m_\nu \neq 0$. The w_0w_a CDM model additionally adds in the present-day value of the dark energy equation of state w_0 , and the high-redshift value of the dark energy equation of state w_a . Typically the present-day densities Ω_i , $i \in \{m, b, K\}$, are denoted with an additional subscript 0; we omit this here for brevity. Primarily, we are interested the w_0w_a CDM case when discussing corrections in this investigation, as a key goal of Stage IV surveys is exploring models of dark energy. However, when examining the cosmological parameter biases, we also present the Λ CDM case, for completeness.

The specific values used for each of these parameters are also chosen for consistency with EC20, and are given in Table 3. As in EC20, the value of $\sum m_\nu \neq 0$ was treated as fixed, and we did not calculate uncertainties or biases for it. When computing biases for all corrections except for the non-flat Universe term,

we set Ω_K to 0. Only when testing the significance of the additional non-flat Universe correction term was it set to 0.05. In this case, the value is selected as being the upper-limit of the Planck Collaboration et al. (2018) 1σ constraint on the parameter.

In the cases when the non-flat Universe correction was not being evaluated, our Fisher matrices matched those of the *Euclid* forecasting specification (EC20), and contained: $\Omega_m, \Omega_b, h, n_s, \sigma_8$ and \mathcal{A}_{IA} for the Λ CDM case, and additionally w_0 and w_a , for the w_0w_a CDM case. When the correction for spatially curvature needed to be tested, the Fisher matrices also included Ω_K . It was not necessary to include any further nuisance parameters within the matrix, as EC20 showed that the inclusion of various different nuisance parameters (such as those modelling the non-linear part of the matter power spectrum) typically altered the forecasted uncertainties on cosmological parameters by less than 10%. The S_{ijmn} were calculated using the publicly available PySSC⁷ code (Lacasa & Grain 2019), with an R_ℓ of 3.

To calculate the cosmological background quantities required for the investigation, including the matter power spectrum and growth factor, we used the CAMB⁸ software package (Lewis et al. 2000). Additionally, we utilised the Halofit (Takahashi et al. 2012) implementation of the non-linear part of the power spectrum, and included additional corrections identified by Bird et al. (2012). Where necessary, we additionally employed Astropy⁹ (Astropy Collaboration et al. 2013, 2018) to compute cosmological distances. The NLA model IA parameters were set to $\mathcal{A}_{IA} = 1.72$ and $C_{IA} = 0.0134$, again in accordance with EC20. The required partial derivatives required were computed numerically, using the procedure described in EC20. Throughout this work, all quantities were evaluated for 200 ℓ -bands. The limits for these were logarithmically spaced, with an ℓ_{\min} of 10, and an ℓ_{\max} of 5000.

3.3. Modelling higher-order corrections

To model the matter bispectrum required by the reduced shear, magnification bias, source-lens clustering, and source obscuration corrections, we used the BiHalofit model and code¹⁰ (Takahashi et al. 2020). This represents the matter bispectrum using one-halo and three-halo terms, which themselves are determined through fitting to N-body simulations.

For the magnification bias correction, we used the slope of the luminosity function as calculated from the fitting formula given in Appendix C of Euclid Collaboration: Lepori et al. (2022). This is determined from the *Euclid* Flagship simulation (Potter et al. 2017) and for the limiting magnitude 24.5 of the VIS instrument (AB in the *Euclid* VIS band (Cropper et al. 2012)). Therefore, it provides the most *Euclid* specific estimate of this quantity to-date. However, once the *Euclid* survey is in-progress, we note that this quantity should be calculated directly from the observed data. We used a single value for the slope for each tomographic redshift bin. This value was calculated at the central redshift of the bin. The slopes for all bins, together with their central redshifts, can be found in Table 4. We note that the magnification bias from Euclid Collaboration: Lepori et al. (2022) was obtained for the $n(z)$ from Euclid Collaboration et al. (2021) but we use the Euclid Collaboration: Blanchard et al. (2020) $n(z)$, however the effect of the small changes in the

Table 4. The values of the slope of the luminosity function used in computing the magnification bias correction. These are calculated at the central redshift of each tomographic bin. The limiting magnitude is taken to be 24.5, and the slopes are calculated with a fitting function (Euclid Collaboration: Lepori et al. 2022) determined from the *Euclid* Flagship simulation.

Bin i	Central Redshift	Slope s_i
1	0.2095	0.108
2	0.489	0.180
3	0.619	0.229
4	0.7335	0.279
5	0.8445	0.335
6	0.9595	0.400
7	1.087	0.480
8	1.2395	0.586
9	1.45	0.753
10	2.038	1.335

assumed $n(z)$ should be small, which is consistent with the small differences in the results between this paper and Deshpande et al. (2020a).

In order to evaluate the source obscuration terms, we set the total number of observed galaxies to 2×10^9 , so that the total number of galaxies per redshift bin was 2×10^8 . We also took the mean galaxy radius to be $\vartheta = 0.32''$ (1.55×10^{-6} rad). This value is the mean half-light radius of galaxies from the *Euclid* Flagship mock (Euclid Collaboration: Bretonnière et al. 2022).

To evaluate the impact of the local Universe correction, we used a smoothing scale of $120 h^{-1}$ Mpc. This is the primary value used in Hall (2020), as it is just large enough for the local overdensity to be linear, while still possessing full-sky spherical coverage within the 2M++ galaxy redshift catalogue (Lavaux & Hudson 2011) used to measure the local overdensity. Accordingly, $\delta_0(R = 120 h^{-1} \text{ Mpc}) = 0.045$ and $\delta_0(R = 120 h^{-1} \text{ Mpc})/\sigma(R = 120 h^{-1} \text{ Mpc}) = 0.85$. Different choices of smoothing scale result in different values of the local overdensity δ_0 , with some choices consistent with zero. As the LU effect scales linearly with the local overdensity, we do not rule out the LU bias as being exactly zero. Furthermore, it might be expected that the LU effect is mostly subsumed within the SSC uncertainty for any choice of smoothing scale, since it is mostly affected by modes that are outside of the survey. Detailed investigation of these points is beyond the scope of the paper, and our intention here is merely to assess how much bias would result from a nominal amplitude for the local overdensity combined with a fiducial implementation of the SSC covariance.

In this work, we compared the magnitude of the studied correction terms to the Gaussian sample variance, $\Delta C_\ell/C_\ell$. This was calculated according to Kaiser (1992), and took the form

$$\Delta C_\ell/C_\ell = \sqrt{2} [f_{\text{sky}}(2\ell + 1)]^{-1/2}. \quad (61)$$

4. Results and discussion

This section presents and discusses the computed values for the studied corrections. First, we show the magnitudes of the correction relative to the magnitude of the cosmic shear angular power spectra, and compare them to the sample variance.

In Fig. 1 the magnitudes of the reduced shear, source-lens clustering, magnification bias, source obscuration, local Universe effect, and non-flat Universe corrections, relative to the angular power spectra are shown. The combined correction is

⁷ <https://github.com/fabienlacasa/PySSC>

⁸ <https://camb.info/>

⁹ <http://www.astropy.org>

¹⁰ http://cosmo.phys.hirosaki-u.ac.jp/takahasi/codes_e.htm

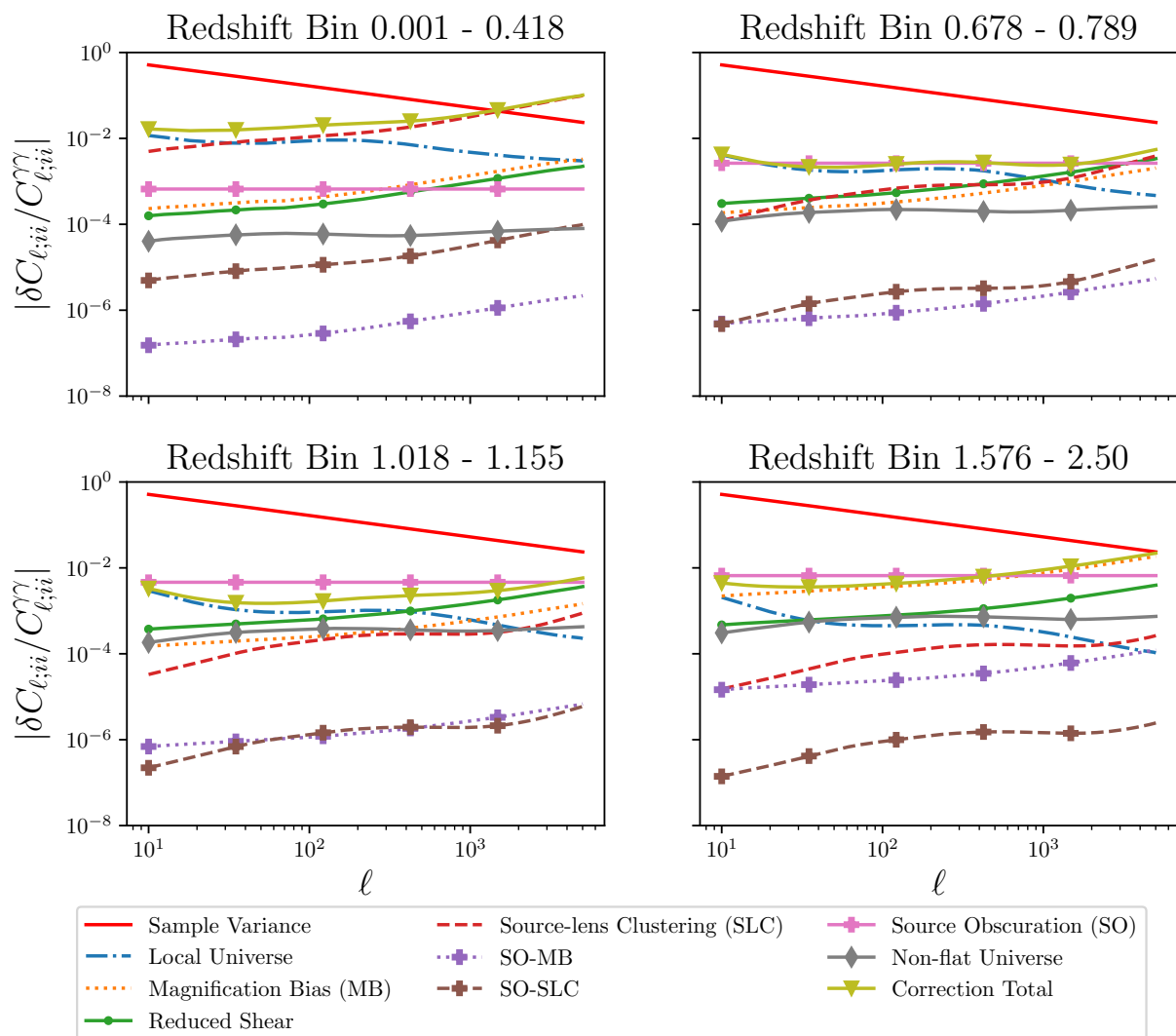


Fig. 1. Absolute magnitudes of the reduced shear, source-lens clustering, magnification bias, local Universe, source obscuration, and non-flat Universe corrections to the shear angular power spectra, relative to those angular power spectra, for *Euclid*. The corrections to the angular power spectra for four redshift bin auto-correlations are shown as representative examples, spanning across the redshift range of *Euclid*. The remaining auto and cross-correlations exhibit the same patterns. The absolute value of the signed sum of the corrections is also shown. These are all compared to the sample variance, calculated according to Eq. (61). Notably, while the magnitudes of individual corrections are either higher at lower redshifts or vice-versa, the magnitude of the sum of the corrections is consistently high. Additionally, the cross-terms between source obscuration, and magnification bias and source-lens clustering are multiple orders of magnitude below other terms and sample variance, suggesting they are negligible. The remainder of the terms are typically of similar magnitudes across redshifts, suggesting they must all be accounted for. We note that these magnitudes are for both the Λ CDM and w_0w_a CDM cases, as the choice of fiducial values for the latter matches the former, and that the non-flat Universe correction here has been computed for a cosmology with $\Omega_K = 0.05$, whilst other corrections are when $\Omega_K = 0$. The markers for the SO-MB, SO-SLC, SO, Non-flat Universe, and total lines are only used to distinguish those from the other terms, and do not have any other significance. The symbols (points) are only included to allow a reader to distinguish the lines (in particular if printing in gray-scale) and do not indicate the ℓ -modes where a computation was made; all quantities were evaluated for 200 ℓ -bands, logarithmically spaced, with an ℓ_{\min} of 10, and an ℓ_{\max} of 5000.

also shown, and we note that this is the sum of the signed values of the corrections, rather than the absolute values which are shown here for comparison. These terms are displayed for the auto-correlations of four redshift bins across the survey's range; specifically, bins 1, 4, 7, and 10. These particular bins are presented for illustrative purposes, and the remaining bins and cross-correlations display consistent trends. It can be seen that all corrections are typically below sample variance both individually and when combined, with the exception of at small

physical scales at the lowest and highest redshifts. A noteworthy detail from this figure is that, while individual corrections are either consistently well below sample variance, or are higher at low redshifts and reduce significantly at high redshifts or vice-versa, the total magnitude of the corrections is consistently high.

Furthermore, another immediately noticeable feature is that the source obscuration cross terms with magnification bias and source-lens clustering are always multiple orders of magnitude

Table 5. The uncertainties on, and biases induced from neglecting the various corrections, in the Λ CDM parameters of Table 3 for *Euclid*. This table considers the case when $\Omega_K = 0$, and accordingly lists the biases resulting from all corrections except the non-flat Universe correction. That can be found in Table 7. All biases are given as a fraction of the 1σ uncertainty on each parameter. A bias is considered significant if it reaches or exceeds 0.25σ , as at this point its uncorrected and corrected confidence contours overlap by less than 90%. ‘RS’ denotes the reduced shear correction, ‘SLC’ is the source-lens clustering term, ‘MB’ is the magnification bias correction, ‘SO’ is the two-point source obscuration correction, ‘SO-MB’ and ‘SO-SLC’ are the source obscuration-magnification bias and source-lens clustering cross terms respectively, and ‘LU’ is the local Universe correction. The total biases from the sum of all corrections, as well as the total biases from only the individually significant corrections are also given.

Cosmo. Param.	Uncertainty (1σ)	RS Bias/ σ	SLC Bias/ σ	MB Bias/ σ	SO Bias/ σ	SO-SLC Bias/ σ	SO-MB Bias/ σ	LU Bias/ σ	Total All	Total Sig.
Ω_m	0.0051	-0.032	1.14	-0.75	0.36	0.0032	0.0047	0.26	0.76	0.73
Ω_b	0.021	0.0035	0.28	0.080	-0.044	6.5×10^{-4}	-6.6×10^{-4}	-0.016	0.26	0.28
h	0.13	0.0084	0.15	0.19	-0.11	2.8×10^{-4}	-0.0013	-0.042	0.23	0.25
n_s	0.029	0.017	0.0065	-0.23	0.18	2.1×10^{-4}	0.0016	0.038	2.6×10^{-4}	-0.07
σ_8	0.0072	0.076	-1.10	0.78	-0.54	-0.0031	-0.0050	-0.22	-0.78	-0.79

Table 6. The uncertainties on, and biases induced from neglecting the various corrections, in the w_0w_a CDM parameters of Table 3 for *Euclid*. This table considers the case when $\Omega_K = 0$, and accordingly lists the biases resulting from all corrections except the non-flat Universe correction. That can be found in Table 7. All biases are given as a fraction of the 1σ uncertainty on each parameter. A bias is considered significant if it reaches or exceeds 0.25σ , as at this point its uncorrected and corrected confidence contours overlap by less than 90%. ‘RS’ denotes the reduced shear correction, ‘SLC’ is the source-lens clustering term, ‘MB’ is the magnification bias correction, ‘SO’ is the two-point source obscuration correction, ‘SO-MB’ and ‘SO-SLC’ are the source obscuration-magnification bias and source-lens clustering cross terms respectively, and ‘LU’ is the local Universe correction. The total biases from the sum of all corrections, as well as the total biases from only the individually significant corrections are also given.

Cosmo. Param.	Uncertainty (1σ)	RS Bias/ σ	SLC Bias/ σ	MB Bias/ σ	SO Bias/ σ	SO-SLC Bias/ σ	SO-MB Bias/ σ	LU Bias/ σ	Total All	Total Sig.
Ω_m	0.010	-0.079	1.25	0.12	0.12	0.0025	-0.0022	0.14	1.55	1.49
Ω_b	0.021	-0.0013	0.30	0.069	-0.021	7.2×10^{-4}	-5.8×10^{-4}	-0.0032	0.35	0.35
h	0.13	0.0057	0.11	0.11	-0.053	3.2×10^{-4}	-6.3×10^{-4}	-0.015	0.16	0.17
n_s	0.031	0.051	-0.19	-0.19	0.026	-3.0×10^{-4}	0.0015	-0.045	-0.35	-0.35
σ_8	0.012	0.081	-1.19	0.017	-0.19	-0.0024	0.0011	-0.093	-1.38	-1.36
w_0	0.13	-0.076	0.81	0.55	-0.039	0.0011	-0.0051	0.032	1.28	1.31
w_a	0.36	0.022	-0.51	-0.60	0.27	-3.3×10^{-4}	0.0053	0.097	-0.72	-0.84

Table 7. The uncertainties for the fiducial Λ CDM and w_0w_a CDM cosmologies of Table 3 when $\Omega_K = 0.05$, and biases induced from neglecting the non-flat Universe correction for the *Euclid* cosmic shear probe. The biases are given as a fraction of the 1σ uncertainty on each parameter. A bias is considered significant if it reaches or exceeds 0.25σ , as at this point its uncorrected and corrected confidence contours overlap by less than 90%. The predicted biases from this correction are well below significance for all parameters.

Cosmo. Param.	Λ CDM Uncertainty (1σ)	Λ CDM Non-flat Universe Bias/ σ	w_0w_a CDM Uncertainty (1σ)	w_0w_a CDM Non-flat Universe Bias/ σ
Ω_m	0.0050	0.043	0.014	-0.062
Ω_b	0.021	-0.0070	0.021	8.1×10^{-4}
Ω_K	0.034	0.043	0.062	-0.054
h	0.12	-0.012	0.12	0.012
n_s	0.029	-0.010	0.030	0.012
σ_8	0.028	0.026	0.042	-0.044
w_0	N/A	N/A	0.14	-0.090
w_a	N/A	N/A	0.56	0.10

smaller than sample variance, and typically the other terms as well; suggesting that these cross-terms are negligible.

Despite the fact that these terms are generally below sample variance, because they make contributions consistently across ℓ -modes, they can still cause significant biases in inferred cosmological parameters. The bias in an estimated parameter resulting from neglecting a systematic effect is typically considered significant if it exceeds 25% of the 1σ uncertainty on that parameter (Taylor et al. 2007). This is because, at that point, the biased and unbiased 1σ confidence contours overlap by less than 90%.

The predicted cosmological parameter biases resulting from neglecting all corrections except for the non-flat Universe term are stated in Table 5 and Table 6, for the Λ CDM and w_0w_a CDM cases, respectively. These are also represented visually in Fig. 2. The biases from the non-flat Universe correction (which requires a fiducial cosmology with non-zero curvature) are stated in Table 7, for both choices of cosmology. From these tables we see that the source-lens clustering, magnification bias, source obscuration, and local Universe terms are individually significant in the Λ CDM case, while instead only the source-lens clustering, magnification bias, and source obscuration corrections are indi-

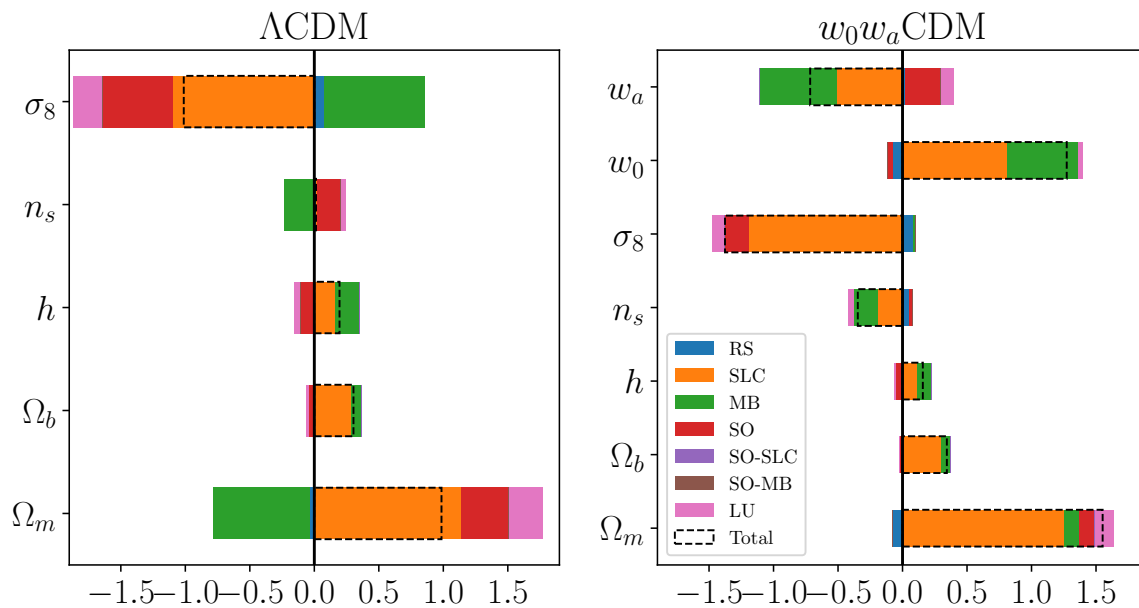


Fig. 2. Stacked bar chart of cosmological parameter biases resulting from the studied higher-order effects, for the flat Λ CDM case (left) of Table 5, and the flat w_0w_a CDM case (right) of Table 6. The non-flat Universe term is not shown here, due to the different cosmology. Biases are presented here as a fraction of the 1σ parameter uncertainty. A bias is non-negligible if its absolute value reaches or exceeds 0.25σ . ‘RS’ denotes the reduced shear correction, ‘SLC’ is the source-lens clustering term, ‘MB’ is the magnification bias correction, ‘SO’ is the two-point source obscuration correction, ‘SO-MB’ and ‘SO-SLC’ are the source obscuration-magnification bias and source-lens clustering cross terms respectively, and ‘LU’ is the local Universe correction. The segments with the dashed outlines show the total parameter biases from these corrections for each parameter.

vidually of concern in the w_0w_a CDM case. This difference is likely due to the presence of the variable dark energy parameters in the w_0w_a CDM scenario reducing sensitivity to scales where the local Universe term is important. Of these, the source-lens clustering term is particularly concerning, as for this term all but two of the parameters have significant biases.

Also shown in these tables are the combined biases when all of the individually significant corrections are taken into account, and when all of these corrections are taken into account. The full total is also shown in Fig. 2. Owing to the fact that some biases are additive while others are subtractive, the total biases in the Λ CDM scenario are, in fact, less severe than some of the individual ones; in particular, source-lens clustering. However, the totals are still significant, and because they do not strongly resemble any one of the biases uniquely, multiple terms must still be computed. In the w_0w_a CDM case, the magnification bias no longer suppresses the source-lens clustering term, instead adding to it and meaning that the total biases in this case are more severe than the individual ones. This change likely occurs due to the dark energy terms increasing sensitivity to scales where the opposite component of the magnification bias (e.g. decrease in galaxy number density due to dilution rather than increase due to increased flux) is dominant.

At inference time, the computational load can be reduced by noting that only the individually significant terms (source-lens clustering, magnification bias, source obscuration, and local Universe) need to be computed in both cases, because this total does not significantly differ from the full total. Although, given that the reduced shear correction is also obtained at no additional cost when computing the magnification bias correction, we recommend including this too. The bias in the two-parameter confidence contours resulting from neglecting the combined effect of the significant biases is shown in Fig. 3 and Fig. 4 for the Λ CDM

and w_0w_a CDM scenarios, respectively. These figures display the contours in the case where no correction has been made, and when the corrections are made. As with Table 5 and Table 6, we see that the cumulative corrections must be accounted for.

We note that many of the modelling specifics used here, for example the slope of the luminosity function or the smoothing scale of local overdensity, will need to be determined directly from the *Euclid* survey itself for self-consistency when these corrections are computed at inference time. Accordingly, it would not be meaningful to place constraints on them here, due to the variable survey specifics.

Additionally, we note that the true value of the local density contrast, and accordingly the radius of the smoothing kernel used to calculate it, remain open questions. Accordingly, there is a large uncertainty in the local Universe correction which cannot be meaningfully constrained and it is possible that it may even be zero. Accurate measurements of the local density contrast are required for this.

Similarly, the source obscuration correction as computed in this work represents a worst-case scenario where *every* galaxy in the foreground of a given redshift slice has an overlap. This represents an upper limit on the bias from this effect, but, in practice, the number of sources with overlap will be a smaller fraction. Accordingly, it is likely that source obscuration will not result in significant biases for true *Euclid* observations, particularly if a robust mitigation strategy is employed.

However, the dominant source of quantifiable modelling uncertainty comes from the modelling of the matter bispectrum and, given that the bispectrum is not currently well constrained by observations, this model is likely to continue to evolve. Accordingly, it is important to constrain the impact of a change in bispectrum model on these terms. To-date, three widely-used matter bispectrum models have been produced (Scoccimarro &

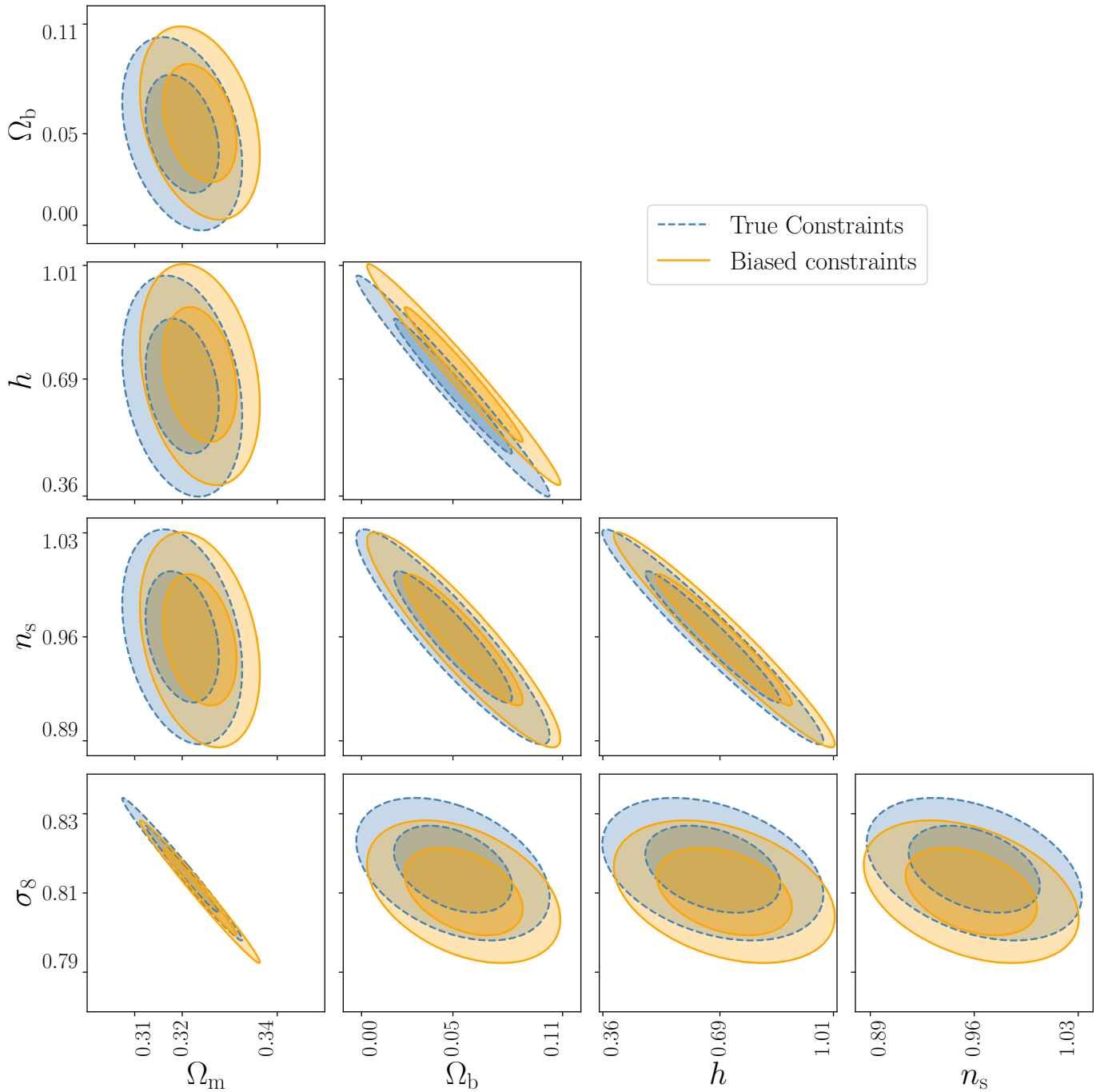


Fig. 3. Projected 1σ and 2σ 2-parameter uncertainty contours for *Euclid* under a Λ CDM cosmology, with and without correcting for the source-lens clustering, magnification bias, source obscuration, and local Universe terms. These are predicted using the Fisher matrix formalism, using the cosmology specified in Table 3, in the case when $\Omega_K = 0$ and is kept fixed. The true location of the constraints is denoted by the blue, dashed contours, while the biased locations if the corrections are not made are given by the solid, gold contours. Significant biases are predicted for Ω_m , Ω_b , h , and σ_8 , and their values can be found in Table 5

Couchman 2001; Gil-Marín et al. 2012; Takahashi et al. 2020). As each subsequent model has become more complex and improved upon the accuracy of its predecessor, comparing correction magnitudes using each of these models would not realistically constrain the uncertainty from the bispectrum model.

Instead, it is useful to set a threshold around the latest of these models (Takahashi et al. 2020), within which any change in the model must be contained, in order to not produce a significant change in the correction terms. We do this by determin-

ing the minimum fractional increase or decrease in the matter bispectrum required across all triangle configurations, for each correction individually, to cause a significant change in any of the cosmological parameter biases, in the w_0w_a CDM case. This corresponds to a change of $\pm 0.25\sigma$ in any one of the biases. Given that we are placing multiplicative limits on the change in the bispectrum, for each correction the smallest limits are found when considering the cosmological parameter that already has the largest bias. The resulting limits are stated in Table 8, along-

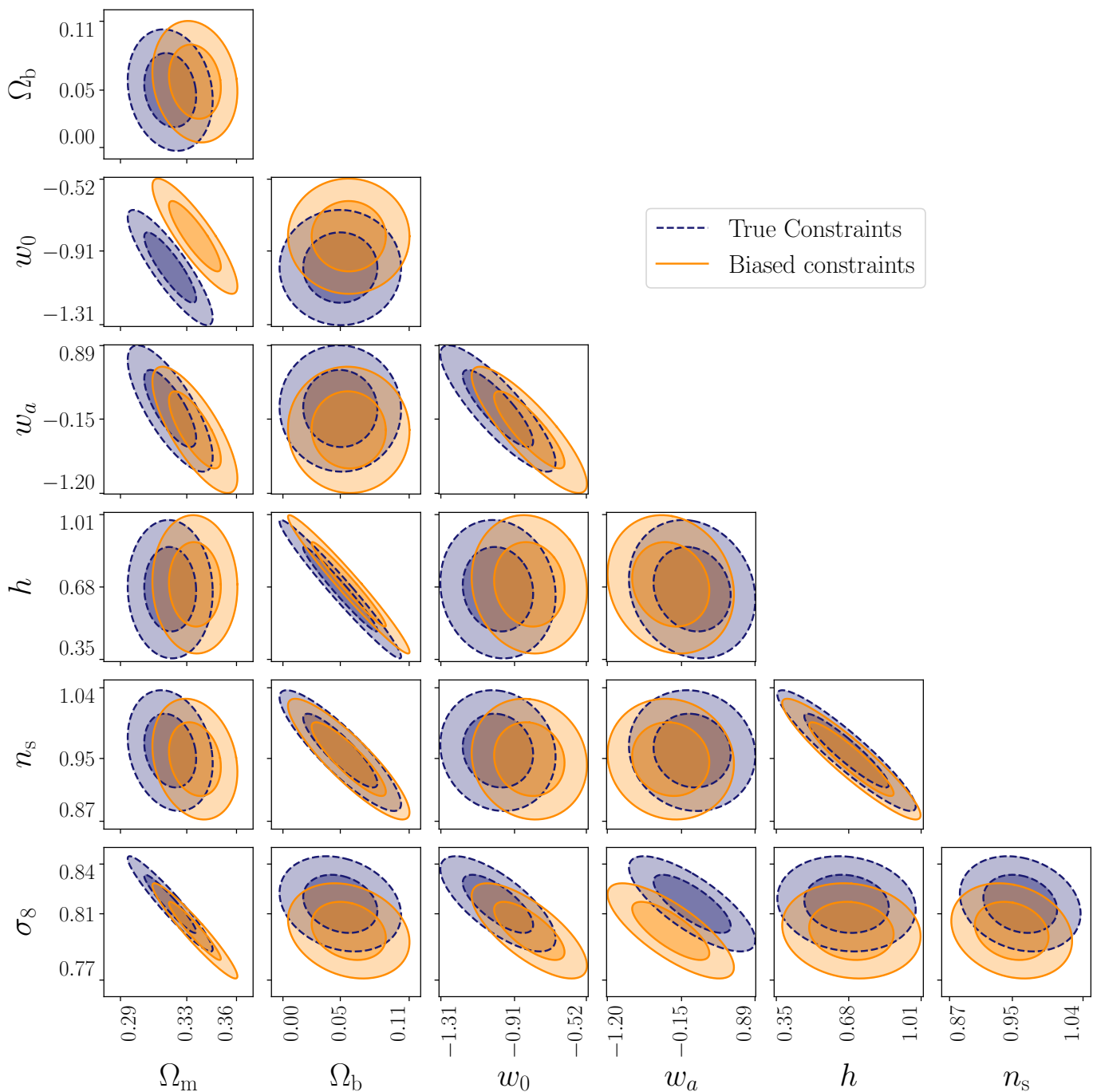


Fig. 4. Projected 1σ and 2σ 2-parameter uncertainty contours for *Euclid*, with and without correcting for source-lens clustering, magnification bias, and source obscuration. These are predicted using the Fisher matrix formalism, for the w_0w_a CDM cosmology specified in Table 3, in the case when $\Omega_K = 0$ and is kept fixed. The true location of the constraints is denoted by the blue, dashed contours, while the biased locations if the corrections are not made are given by the solid, orange contours. Significant biases are predicted for Ω_m , Ω_b , n_s , σ_8 , w_0 , and w_a , and their values can be found in Table 6.

side the parameter which would see the corresponding significant change in its bias. From this, we see that the corrections most susceptible to a change in the bispectrum model are the source-lens clustering and magnification bias terms, as neglecting these already creates the most significant biases individually. Additionally, the two source obscuration cross terms are the least sensitive, requiring a change of an order-of-magnitude. This further reinforces the fact that these terms are safely negligible. We stress that these multiplicative limits are not exhaustive cut-offs

on when a bispectrum model would cause a significant change, because a model with a sufficiently large change for only a select sub-range of scales or configurations could still cause a significant difference in the parameter biases. We recommend explicit reevaluation with any future updated bispectrum models, should they non-trivially exceed these thresholds frequently.

Another consideration is how the inclusion of the studied effects would affect the size of the cosmological parameter uncertainty constraints themselves. In this work, we do not explicitly

Table 8. The increasing and decreasing multiplicative changes required in the matter bispectrum, across all configurations and all ℓ -modes, to significantly affect the biases from each relevant correction term. A significant change is when any one of the biases on the cosmological parameters changes by $\pm 0.25\sigma$. These values encapsulate the fractional change for a fixed cosmology, in this case the w_0w_a CDM cosmology of Table 3, with $\Omega_k = 0$. It should be noted that these values are not exhaustive thresholds, as sufficiently significant changes to the model at a particular sub-range of scales may also be problematic.

Correction	Min. Change Param	Bispec. Mult. Increase	Bispec. Mult. Decrease
RS	Ω_m	-2.16	-4.16
SLC	Ω_m	1.20	0.80
MB	w_a	0.58	1.41
SO-SLC	Ω_m	100.63	-99.67
SO-MB	w_a	47.00	-47.00

calculate this resulting change. However, it has previously been shown that for the bispectrum-dependent terms, even corrections that cause biases of greater than 1σ result in negligible changes to the uncertainty constraints (Shapiro 2009; Deshpande et al. 2020a).

5. Conclusions

In this investigation, we have examined the higher-order corrections to the cosmic shear angular power spectra that must be modelled when performing inference with *Euclid*. By first reviewing the literature, we identified 24 correction terms, and gathered representative values to facilitate comparison. From these, we identified six corrections which were potentially important for *Euclid*, and evaluated them explicitly. These were: the relaxation of the reduced shear approximation, the source-lens clustering correction, the magnification bias correction, the source obscuration correction, the local Universe correction, and the non-flat Universe correction.

After calculating these corrections, we used the Fisher matrix formalism to predict the biases in cosmological parameter biases if each of these terms were to be neglected, in order to identify which ones are necessary to be modelled for *Euclid*. This was done for two scenarios: a Λ CDM cosmology, and a w_0w_a CDM cosmology. For the first of these scenarios, we found that the source-lens clustering, magnification bias, source obscuration, and local Universe terms were significant, while for the second case we found that the source-lens clustering, magnification bias, and source obscuration corrections each produced significant biases in multiple parameters individually. The source-lens clustering term was noted as being of particular concern as multiple biases approached or exceeded 1σ . However, in the Λ CDM case we found that when the biases are combined, they frequently suppressed each other, leading the total bias to be lower than many of the individual biases. Despite this, the total biases were still significant, and did not strongly represent the exact biases from any correction individually. In the w_0w_a CDM case, the total of the three biases was higher than the individual terms. Accordingly, we recommend that the source-lens clustering, magnification bias, source obscuration, and local Universe corrections are all taken into account when modelling the shear angular power spectra for *Euclid*. Additionally, given that the reduced shear correction is obtained at no additional cost when computing the magnification bias correction, we recommend that this too be included.

To provide some constraints on the predictive ability of this work, we quantified how much the matter bispectrum would have to change by in order to illicit a significant change in the biases predicted in this work. We identified that a $\sim 20 - 40\%$ increase or decrease in the amplitude of the matter bispectrum at all scales, and for all triangle configurations would be required for this.

We note that this work did not investigate the impact of higher-order corrections on the IA spectra. Typically, even for effects which, when neglected, produce high biases, i.e. $\sim O(1\sigma)$, the corresponding corrections to the IA spectra cause negligible biases (Deshpande et al. 2020a). Accordingly, an explicit evaluation should not be necessary. Furthermore, we did not consider the impact of baryonic feedback on the bispectrum, and therefore on its dependant corrections. The impact of this remains poorly understood, with inconsistent findings from different simulations (Semboloni et al. 2013; Barreira et al. 2019). Accordingly this is out of the scope of this investigation. However, we note that given the magnitude of change required to the matter bispectrum in order to cause a significant change in the cosmological parameter biases, it is unlikely baryonic feedback would significantly alter the predictions of this investigation.

Given that we find it is necessary to include these higher-order terms in the modelling of the shear power spectra, an open question is the optimal strategy to do so. It has repeatedly been shown that computing these corrections for just one cosmology is relatively time consuming (Deshpande et al. 2020a; Duncan et al. 2022), rendering computation at inference time a serious challenge. While it may be possible to sufficiently optimise the required evaluation code, alternate strategies may also prove useful. Scale cutting techniques such as k -cut cosmic shear (Taylor et al. 2018a) have been shown to mitigate the need to make such corrections without significantly compromising the constraining power of Stage IV surveys (Deshpande et al. 2020b). Alternatively, emulation has recently become a popular tool in cosmology for reducing computation time at inference by replacing analytical models with emulators (see e.g. recent work emulating the matter power spectrum within Spurio Mancini et al. 2022). Emulators could also be developed directly for these correction terms, or intermediate quantities such as the matter or convergence bispectra.

Furthermore, while our analysis here is limited to the angular power spectra, significant corrections for this statistic are also likely to be significant for the two-point correlation function. In fact, due to the mode-mixing that occurs when transforming the power spectra to correlation functions, the effect of the discussed approximations is likely to be more severe. This, combined with the sensitivity of the correlation function to higher ℓ -modes and the additional approximations required (e.g. the flat Hankel transform; Kitching et al. 2017), means that if correlation functions were to be used a similar but separate study would be required to demonstrate modelling of the correlation function to higher-order corrections.

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