# Experimental Investigation of Compressible Boundary Layers under the Influence of Pressure Gradients 

Raymond C. Wier

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OF COMPRESSIBLE BOUNDARY LAYERS
UNDER THE INFLUENCE
OF PRESSURE GRADIENTS

THESIS
Raymond C. Weer
Captain, USAF
AFIT/GAE/ENY/96D-04

#  <br> Approved so r pueise relacan  <br> DEPARTMENT OF THE AIR FORCE <br> AIR UNIVERSITY 

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# EXPERIMENTAL INVESTIGATION <br> OF COMPRESSIBLE BOUNDARY LAYERS <br> UNDER THE INFLUENCE <br> OF PRESSURE GRADIENTS 

THESIS

# Presented to the Faculty of the School of Engineering of the Air Force Institute of Technology Air University <br> In Partial Fulfillment of the Requirements for the Degree of Master of Science in Aeronautical Engineering 

Raymond C. Wier, B.S.
Captain, USAF

December 1996

Approved for public release; distribution unlimited

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Raymond C. Wier

## Table of Contents

Acknowledgments ..... ii
Table of Contents ..... iii
List of Figures ..... vi
List of Tables ..... ix
List of Symbols ..... xii
Abstract ..... xvi

1. Introduction ..... 1-1
1.1 Motivation ..... 1-1
1.2 Research Objectives ..... 1-4
1.3 Overview of Current Study ..... 1-5
2. Background ..... 2-1
2.1 Governing Equations ..... 2-1
2.1.1 The Compressible Navier-Stokes Equations ..... 2-1
2.1.2 Reynolds-Averaged Navier-Stokes Equations (RANS) ..... 2-2
2.1.3 Favré-Averaged Navier-Stokes Equations (FANS) ..... 2-3
2.1.4 Comparison Between Favré-Averaged and Reynolds-Averaged N-S Equations ..... 2-4
2.2 Literature Review ..... 2-5
3. Facilities and Equipment. ..... 3-1
3.1 Mach 3.0 Wind Tunnel. ..... 3-1
3.1.1 High-Pressure Air Supply ..... 3-1
3.1.2 Vacuum System ..... 3-2
3.1.3 Plenum Chamber ..... 3-2
3.1.4 Mach 3.0 Nozzle. ..... 3-2
3.1.5 Test Sections ..... 3-3
3.2 Mach 1.7 Wind Tunnel. ..... 3-4
3.2.1 Plenum Chamber ..... 3-4
3.2.2 Mach 1.7 Nozzle. ..... 3-5
3.3 Mach 5.0 Wind Tunnel ..... 3-5
3.3.1 Air Supply ..... 3-6
3.3.2 Pebble Bed Heaters ..... 3-6
3.3.3 Plenum Chamber ..... 3-6
3.3.4 Mach 5.0 Nozzle ..... 3-7
3.3.5 Test Sections ..... 3-7
3.4 Data Acquisition Equipment ..... 3-8
3.5 Probes ..... 3-8
3.5.1 Pressure Probes ..... 3-8
3.5.2 Hot-Wire and Hot-Film Probes ..... 3-9
3.6 Particle Image Velocimetry ..... 3-10
3.7 Computer Equipment ..... 3-11
4. Data Reduction Techniques ..... 4-1
4.1 Mean Flow Data Reduction ..... 4-1
4.1.1 Calculation of Mach Number, Velocity and Fluid Properties ..... 4-1
4.1.1.1 Pressure and Temperature Methodology - ZPG Flow ..... 4-2
4.1.1.2 Pressure and Temperature Methodology - FPG and CPG Flow ..... 4-2
4.1.2 Calculation of Boundary Layer, Displacement and Momentum Thickness. ..... 4-3
4.1.3 Calculation of Van Driest Velocity Profile and Wall Shear Stress ..... 4-4
4.2 Hot-Wire Data Reduction ..... 4-5
4.2.1 Hot-Wire Theory. ..... 4-6
4.2.2 Separation of Turbulence Variables ..... 4-8
4.2.3 Single Overheat Theory ..... 4-9
4.2.4 Mach 2.8 Hot-Wire Probe Data ..... 4-10
4.2.5 Mach 1.7 Hot-Wire Probe Data ..... 4-10
4.2.6 Power Spectra Data Reduction ..... 4-11
4.3 PIV Data Reduction ..... 4-11
5. Mach 2.8 Results and Discussion ..... 5-1
5.1 Flow Visualization ..... 5-1
5.2 Mean Flow Results ..... 5-3
5.2.1 Conventional Pressure Probe ..... 5-3
5.2.2 Hot-Wire Mean Profiles ..... 5-6
5.2.3 PIV Mean Profiles ..... 5-9
5.2.4 Comparison Between Data Collection Methods ..... 5-14
5.2.4.1 Mean Flow Comparison ..... 5-14
5.2.4.2 Van Driest Correlation. ..... 5-16
5.3 Turbulence Results ..... 5-19
5.3.1 Hot-Film Turbulence Results ..... 5-20
5.3.1.1 Mass Flux Turbulence Intensity Profiles ..... 5-20
5.3.1.2 Energy Spectra ..... 5-23
5.3.2 PIV Turbulence Results ..... 5-26
5.3.3 Comparison Between Data Collection Methods - Turbulence Quantities ..... 5-28
5.3.4 Density Fluctuations ..... 5-30
6. Mach 1.7 Results and Discussion ..... 6-1
6.1 Mach 1.7 Conventional Pressure Data ..... 6-1
6.2 Hot-Wire Results ..... 6-3
6.2.1 Mean Flow Profiles ..... 6-3
6.3 Turbulence Results ..... 6-4
6.3.1 Mass Flux Turbulence Intensity Profiles ..... 6-4
6.3.2 Energy Spectra ..... 6-5
6.3.3 Separated Turbulence Results ..... 6-6
7. Mach 5.0 Results and Discussion ..... 7-1
8. Conclusions and Recommendations ..... 8-1
8.1 Mach 2.8 Conclusions ..... 8-1
8.2 Mach 1.7 Conclusions ..... 8-2
8.3 Mach 5.0 Conclusions ..... 8-3
8.4 Recommendations ..... 8-3
Appendix A : Hot-Wire Calibration ..... A-1
A. 1 Mach 2.8 Calibration ..... A-1
A. 2 Mach 1.7 Calibration ..... A-3
Appendix B : Uncertainty Analysis ..... B-1
B. 1 Elementary Uncertainties ..... B-1
B. 2 Derived Uncertainties ..... B-3
B.2.1 Mean Flow Derived Uncertainties ..... B-3
B.2.2 Hot-Wire Derived Uncertainties ..... B-4
B. 3 PIV Signal-to-Noise Uncertainty ..... B-4
Appendix C: Data Files ..... C-1
Bibliography ..... BIB-1
Vita. ..... VITA

## List of Figures

Figure 3-1: AFIT Mach 2.8 wind tunnel schematic ..... 3-1
Figure 3-2: Pressure gradient contours \& data collection locations ..... 3-4
Figure 3-3: Optical glass floor schematic. ..... 3-4
Figure 3-4: Mach 5.0 wind tunnel schematic ..... 3-6
Figure 3-5: PIV schematic. ..... 3-10
Figure 5-1: ZPG and CPG flow visualization images ..... 5-1
Figure 5-2: ZPG and FPG flow visualization images ..... 5-2
Figure 5-3: Pitot pressure profile ..... 5-4
Figure 5-4: Cone-static pressure profile. ..... 5-4
Figure 5-5: Mach number profile. ..... 5-5
Figure 5-6: ZPG hot-film data ..... 5-7
Figure 5-7: FPG hot-film data ..... 5-7
Figure 5-8: $\mathrm{CPG}(\mathrm{x}=68 \mathrm{~cm})$ hot-film data. ..... 5-9
Figure 5-9: CPG ( $x=70 \mathrm{~cm}$ ) hot-film data. ..... 5-9
Figure 5-10: ZPG velocity contour plot. ..... 5-12
Figure 5-11: FPG velocity contour plot ..... 5-12
Figure 5-12: CPG velocity contour plot ..... 5-12
Figure 5-13: ZPG velocity vectors ..... 5-13
Figure 5-14: FPG velocity vectors. ..... 5-13
Figure 5-15: CPG velocity vectors ..... 5-13
Figure 5-16: ZPG velocity comparison ..... 5-15
Figure 5-17: FPG velocity comparison. ..... 5-15
Figure 5-18: CPG velocity comparison ..... $5-16$
Figure 5-19: ZPG Van Driest velocity profile ..... 5-17
Figure 5-20: FPG Van Driest velocity profile ..... 5-18
Figure 5-21: CPG $(x=68 \mathrm{~cm})$ Van Driest velocity profile ..... 5-18
Figure 5-22: CPG $(x=70 \mathrm{~cm})$ Van Driest velocity profile ..... 5-19
Figure 5-23: ZPG mass flux turbulence intensity ..... 5-20
Figure 5-24: FPG mass flux turbulence intensity ..... 5-21
Figure 5-25: CPG $(x=68 \mathrm{~cm})$ mass flux turbulence intensity ..... 5-22
Figure 5-26: $\mathrm{CPG}(x=70 \mathrm{~cm})$ mass flux turbulence intensity ..... 5-23
Figure 5-27: ZPG power spectra data ..... 5-24
Figure 5-28: FPG power spectra data ..... 5-24
Figure 5-29: CPG power spectra data ..... 5-25
Figure 5-30: ZPG u turbulence intensity contours (\%) ..... 5-27
Figure 5-31: FPG u turbulence intensity contours (\%) ..... 5-27
Figure 5-32: CPG u turbulence intensity contours (\%) ..... 5-27
Figure 5-33: ZPG u turbulence intensity profiles ..... 5-29
Figure 5-34: FPG u turbulence intensity profiles ..... 5-29
Figure 5-35: CPG $(x=68 \mathrm{~cm})$ u turbulence intensity profile ..... 5-30
Figure 5-36: CPG $(x=70 \mathrm{~cm})$ u turbulence intensity profile ..... 5-31
Figure 5-37: Density fluctuations ..... 5-31
Figure 6-1: Mach 1.7 Pitot pressure profiles ..... 6-1
Figure 6-2: Mach 1.7 Mach number profile ..... 6-2
Figure 6-3: Mach 1.7 velocity profile. ..... 6-2
Figure 6-4: Mach 1.7 ZPG hot-film data ..... 6-4
Figure 6-5: Mach 1.7 ZPG mass flux turbulence intensity ..... 6-5
Figure 6-6: Mach 1.7 power spectra data ..... 6-6
Figure 6-7: Mach 1.7 u turbulence intensity profiles ..... 6-7
Figure 6-8: Mach 1.7 density fluctuation profiles. ..... 6-7
Figure 6-9: Mach 1.7 total temperature fluctuation profiles ..... 6-8
Figure 7-1: Mach 5.0 Pitot pressure profiles ..... 7-1
Figure 7-2: Mach 5.0 Mach number profile ..... 7-2
Figure 7-3: Mach 5.0 velocity profile. ..... 7-2
Figure 7-4: Mach 5.0 CPG and FPG static pressure profiles ..... 7-3
Figure A-1: Mach 2.8 hot-wire calibration curve $(\mathrm{OHR}=2.03)$ ..... A-2
Figure A-2: Mach 2.8 hot-wire calibration curve $(\mathrm{OHR}=1.66)$ ..... A-2
Figure A-3: Mach 1.7 hot-wire calibration curves. ..... A-4

## List of Tables

Table 1-1: Kolmogoroff scale calculations ..... 1-1
Table 3-1: Pressure gradient test section curve coefficients ..... 3-3
Table 4-1: Mach 2.8 probe and overheat ratio data ..... 4-10
Table 4-2: Mach 1.7 probe and overheat ratio data ..... 4-10
Table 5-1: Boundary layer thickness parameters. ..... 5-5
Table 5-2: Number of PIV images used to create the combined image ..... 5-10
Table 6-1: Mach 1.7 boundary layer thickness ..... 6-3
Table 7-1: Mach 5.0 boundary layer thickness. ..... 7-1
Table 8-1: Verification of Settles and Dodson criteria. ..... 8-2
Table A-1: Mach 2.8 sample raw calibration file $(\mathrm{OHR}=2.03)$ ..... A-1
Table A-2: Mach 2.8 sample raw calibration file $(\mathrm{OHR}=1.66)$ ..... A-1
Table A-3: Mach 2.8 calibration constants ..... A-3
Table A-4: Mach 1.7 sample raw calibration file $(\mathrm{OHR}=1.95)$ ..... A-3
Table A-5: Mach 1.7 sample raw calibration file $(\mathrm{OHR}=1.66)$ ..... A-4
Table A-6: Mach 1.7 calibration constants ..... A-5
Table B-1: Mach 2.8 ZPG reference conditions ..... B-2
Table B-2: Elementary uncertainties. ..... B-3
Table B-3: Mean flow derived uncertainties ..... B-4
Table B-4: Boundary layer height uncertainty ..... B-4
Table B-5: Van Driest uncertainties ..... B-4
Table B-6: Hot-wire derived uncertainties ..... B-5
Table C-1: Mach 2.8 ZPG pressure data. ..... C-1
Table C-2: Mach 2.8 FPG pressure data. ..... C-2
Table C-3: Mach 2.8 ZPG hot-film data - traverse up ..... C-4
Table C-4: Mach 2.8 ZPG separated turbulence variables single overheat- traverse up ..... C-4
Table C-5: Mach 2.8 ZPG hot-film data - traverse down ..... C-5
Table C-6: Mach 2.8 ZPG separated turbulence variables single overheat - traverse down ..... C-5
Table C-7: Mach 2.8 FPG hot-film data-traverse up. ..... C-6
Table C-8: Mach 2.8 FPG separated turbulence variables single overheat - traverse up. ..... C-7
Table C-9: Mach 2.8 CPG ( $\mathrm{x}=68 \mathrm{~cm}$ ) hot-film data - traverse up. ..... C-8
Table C-10: Mach 2.8 CPG ( $\mathrm{x}=68 \mathrm{~cm}$ ) separated turbulence variables single overheat- traverse up ..... C-8
Table C-11: Mach $2.8 \mathrm{CPG}(\mathrm{x}=68 \mathrm{~cm})$ hot-film data - traverse down ..... C-9
Table C-12: Mach $2.8 \mathrm{CPG}(\mathrm{x}=68 \mathrm{~cm}$ ) separated turbulence variables single overheat- traverse down. ..... C-9
Table C-13: Mach 2.8 CPG $(x=70 \mathrm{~cm})$ hot-film data - traverse up. ..... C-10
Table C-14: Mach 2.8 CPG ( $\mathrm{x}=70 \mathrm{~cm}$ ) separated turbulence variables single overheat- traverse up ..... C-10
Table C-15: Mach 2.8 CPG $(x=70 \mathrm{~cm})$ hot-film data - traverse down ..... C-11
Table C-16: Mach 2.8 CPG ( $\mathrm{x}=70 \mathrm{~cm}$ ) separated turbulence variables single overheat- traverse down C -11
Table C-17: Mach 2.8 discrete data turbulence intensity points ..... C-12
Table C-18: Mach 2.8 hot-film van Driest data ..... C-12
Table C-19: Mach 2.8 power spectra data. ..... C-13
Table C-20: Mach 2.8 ZPG PIV data. ..... C-17
Table C-21: Mach 2.8 FPG PIV data - 25 images. ..... C-17
Table C-22: Mach 2.8 FPG PIV data - 93 images ..... C-18
Table C-23: Mach 2.8 CPG $(x=68 \mathrm{~cm})$ PIV data ..... C-18
Table C-24: Mach 2.8 CPG ( $\mathrm{x}=68 \mathrm{~cm}$ ) PIV data ..... C-19
Table C-25: PIV van Driest data ..... C-19
Table C-26: Mach 1.7 ZPG pressure data ..... C-20
Table C-27: Mach 1.7 FPG pressure data ..... C-22
Table C-28: Mach 1.7 ZPG hot-film data - traverse up ..... C-24
Table C-29: Mach 1.7 ZPG hot-film data - traverse down ..... C-25
Table C-30: Mach 1.7 ZPG separated turbulence variables - traverse down ..... C-26
Table C-31: Mach 1.7 ZPG separated turbulence variables - traverse down ..... C-26
Table C-32: Mach 1.7 discrete turbulence intensity data points ..... C-27
Table C-33: Mach 1.7 power spectra data. ..... C-27
Table C-34: Mach 5.0 ZPG pressure data ..... C-30
Table C-35: Mach 5.0 FPG pressure data ..... C-31
Table C-36: Mach 5.0 CPG $\left(\mathrm{X}_{\mathrm{ts}}=5.1 \mathrm{~cm}\right)$ pressure data ..... C-33
Table C-37: Mach 5.0 CPG $\left(\mathrm{X}_{\mathrm{ts}}=6.35 \mathrm{~cm}\right)$ pressure data ..... C-34
Table C-38: Mach 5.0 wall static pressure data. ..... C-34

## List of Symbols

| Symbol | Description |
| :---: | :---: |
| $\mathrm{a}, \mathrm{b}$ | hot-wire calibration constants |
| $\mathrm{A}^{\prime}, \mathrm{B}^{\prime}$ | Constants in Crocco relationship and van Driest correlation |
| $\mathrm{C}_{\mathrm{f}}$ | Coefficient of friction |
| CFD | Computational fluid dynamics |
| $\mathrm{C}_{0}$ | Constant used in Chapter 4 |
| CPG | Combined pressure gradient |
| f | frequency, hot-wire sensitivity |
| FANS | Favré averaged Navier-Stokes equations |
| FPG | Favorable pressure gradient |
| g | hot-wire sensitivity |
| GLS | General least squares |
| h | Enthalpy |
| k | Constant in van Driest correlation, thermal conductivity |
| $\overline{\text { L }}$ | Average vector length |
| LDV | Laser doppler velocimetry |
| M | Mach number |
| MOH | Multiple overheat ratios |
| N | Number of overheat ratios |
| Nu | Nusselt number |
| P | Stagnation or total pressure |
| p | Static pressure |
| PIV | Particle image velocity |
| Pr | Prandtl number |


| q | Heat flux vector |
| :---: | :---: |
| Q | Heat flux |
| r | Recovery factor |
| R | Resistance |
| RANS | Reynolds averages Navier-Stokes equations |
| Re | Reynolds number |
| RHS | Right hand side (of an equation) |
| SOH | Single overheat ratio |
| T | Characteristic time, temperature |
| t | time |
| TI | Turbulence intensity |
| u | axial velocity |
| ${ }^{*}$ | Van Driest friction velocity |
| $\mathrm{u}_{\mathrm{eff}}$ | Van Driest effective velocity |
| $\mathrm{u}_{\text {eff }}{ }^{+}$ | Van Driest effective scaled velocity |
| V | Velocity vector |
| $\mathrm{V}_{\mathrm{w}}$ | Wire voltage |
| x | coordinate tangent to the flow direction |
| y | coordinate normal to the wind tunnel ceiling |
| $\mathrm{y}^{+}$ | Van Driest length scale |
| ZPG | Zero pressure gradient |
| $\alpha$ | Constant in single overheat ratio data reduction |
| $\beta$ | Clauser's equilibrium factor, constant in single overheat ratio data reduction |
| $\delta$ | boundary layer thickness |
| $\delta^{*}$ | Compressible displacement thickness |
| $\delta_{1}{ }^{*}$ | Incompressible displacement thickness |


| $\delta_{i j}$ | Kronecker delta |
| :---: | :---: |
| $\Delta \mathrm{L}_{\mathrm{G}-\mathrm{R}}$ | Individual vector length |
| $\varepsilon$ | Percent error |
| $\gamma$ | Ratio of specific heats ( $=1.4$ ) |
| k | Constant in van Driest correlation |
| $\lambda$ | Wavelength |
| $\mu$ | absolute viscosity |
| $\Pi$ | Coles wake function |
| $\theta$ | Compressible momentum thickness |
| $\theta_{1}$ | Incompressible momentum thickness |
| $\rho$ | density |
| $\tau$ | shear stress |
| TI | Turbulence intensity |
| $\omega$ | Constant in van Driest correlation |
| Subscripts |  |
| 1,2 | Wind tunnel station |
| bar | Average |
| c | Cone |
| e | Edge |
| 1 | Leads |
| M | Mach number |
| o | Reference or total |
| rms | Root mean square |
| S | Series |
| t | Total |

Boundary layer height

## Superscripts

| ()$^{\prime \prime}$ | Favré fluctuating component |
| :--- | :--- |
| ()$^{\prime}$ | Reynolds fluctuating component |
| $\overline{()^{\prime}}$ | Time averaged variable |
| $\sim$ | Favré averaged variable |
| ()$^{\sim}$ | Degree |
| TF | Turbulent, Favré averaged |
| TR | Turbulent, Reynolds averaged |

## Abstract

This study examined the effect of mild pressure gradients on the mean and turbulent flow of high-speed boundary layers. Three Mach numbers (1.7, 3.0 and 5.0) were investigated. Three pressure gradients were examined; a zero pressure gradient (ZPG), a favorable pressure gradient (FPG), and a combined pressure gradient (CPG). The CPG consisted of an adverse pressure gradient followed by a favorable pressure gradient. Conventional pressure probes, hot-wire and particle image velocimetry (PIV) were used to examine the flow. Measurement included mean velocity, velocity turbulence intensity, mass flux turbulence intensity and energy spectra. Instantaneous ( 10 nsec ) Mie scattering flow visualizations were acquired. Qualitatively, the flow visualizations indicated that the turbulent flow structures were strongly affected by the pressure gradients. For the Mach 2.8 case, the PIV contours and the hot-wire profiles both indicated that the boundary layer thickness increased by $40 \%$ and decreased by $100 \%$ relative to the ZPG for the favorable and adverse pressure gradients, respectively. Further, the PIV and hot-wire data indicated that the axial turbulence intensity levels increased by $22 \%$ for the CPG and decreased by $25 \%$ for the FPG. The energy spectra data indicated that once a pressure gradient was applied (favorable or adverse) the low frequency energy increased followed by a rapid decay. Lastly, it was found that nominally 20 to 30 PIV images were sufficient for mean flow boundary layer velocities, but 93 images (the maximum recorded in this study) were insufficient to adequately resolve Reynolds shear stresses.

# EXPERIMENTAL INVESTIGATION OF COMPRESSIBLE BOUNDARY LAYERS UNDER THE INFLUENCE OF PRESSURE GRADIENTS 

## 1. Introduction

### 1.1 Motivation

Current U.S. high-speed aerospace vehicles are designed to operate at such high velocities and Reynolds numbers that conventional wind tunnel testing rapidly becomes too expensive or time consuming. To reduce the cost of testing, computational fluid dynamics (CFD) play an ever increasing role in the design philosophy. CFD has demonstrated its reliability and accuracy for simple geometries such as flow over a flat plate, but results for more complex shapes are as yet unproven. For the exterior of the vehicle to the internal shape of the engine's inlet or nozzle, turbulent flow over curved surfaces dominates the design. The underlying physics of these types of flows are not understood (Bradshaw, 1974 \& Spina et al., 1994); thus, accurate predictions are not possible.

Turbulence has proven to be the controlling factor in the understanding of high-speed highReynolds number flows. Turbulent flow is characterized by a wide range of temporal and size scales, with the smallest defined by Kolmogoroff scales for length, speed and time (Wilcox, 1994). To accurately resolve the smallest eddies in a turbulent flow, the programmer must space the CFD grid nodes at a resolution sufficient to capture these eddies. Detailed estimates of the Kolmogoroff scales for Mach 2.9 and Mach 5.0 flow, with corresponding CFD solution times, are listed in Table 1-1. The total number of grid points is based on a $6.35 \mathrm{~cm} \times 6.35 \mathrm{~cm} \times 75 \mathrm{~cm}$ volume and a $7.62 \mathrm{~cm} \times 7.62 \mathrm{~cm} \times 25.4 \mathrm{~cm}$ volume for the Mach 3.0 and Mach 5.0 tunnels, respectively. These volumes represent the total test section volume for each tunnel. The total number of calculations is based on the crude estimate of $10^{7}$
calculations per grid point. The time to complete the solution is based on IBM's latest computational goal of 3 trillion calculations per second (Pressman, 1996).

Table 1-1: Kolmogoroff scale calculations

|  | Mach 3.0 tunnel | Mach 5.0 tunnel |
| :--- | :--- | :--- |
| Freestream Reynolds number $\left(\mathrm{Re}_{\delta}\right.$, estimated $)$ | $1.9 \times 10^{6}$ | $1.21 \times 10^{6}$ |
| Boundary layer thickness $(\delta)(\mathrm{m}$, estimated $)$ | 0.13 | $2.24 \times 10^{-3}$ |
| Kolmogoroff length scale $(\mathrm{m})$ | $2.48 \times 10^{-6}$ | $1.09 \times 10^{-8}$ |
| Total number of grid points | $2.00 \times 10^{14}$ | $1.15 \times 10^{21}$ |
| Total number of calculations | $2.00 \times 10^{21}$ | $1.15 \times 10^{28}$ |
| Time to complete the solution (years) | 21,000 | 12,000 |

The time to complete the solution in Table 1-1 is only for the test section in question. Solution times over curved surfaces or entire vehicles, even with clustering, would naturally take longer. Hence, computational speed and memory are important limitations of modern CFD.

One successful method used to reduce the CFD solution time is by using approximate, timeaveraged forms of the governing Navier-Stokes equations. The two most prevalent averaging techniques are Reynolds averaging and Favré averaging. Since the full Navier-Stokes equations are non-linear, the averaging techniques result in extra cross correlation fluctuation terms (see Sections 2.1.2 and 2.1.3). This leads to the closure problem wherein more unknowns than equations exist. The science and art of turbulence modeling is to define new variables to help solve this under-defined system.

The turbulence model evolution followed the classical approach of solution development. Simple flows were examined first and then extrapolated to more complex cases. For CFD this means the incompressible zero pressure gradient (ZPG) case was examined first. Once this case was satisfactorily understood, the existing models were extrapolated to flow with pressure gradients on an ad-hoc basis. Due to the quantity and quality of subsonic data with and without pressure gradients, this turned out to be a natural progression of model generality. However, for high-speed flow the situation is more complicated and not well understood. This lack of understanding can be traced to the clear lack of accurate turbulence data (Spina et al., 1994).

Morkovin's hypothesis provided the rationale for extrapolating incompressible models to compressible flows. Morkovin's hypothesis states that "the turbulence structure is unaffected by compressibility as long as the fluctuation Mach number is much less than unity" (Morkovin, 1964). While this hypothesis holds for most cases, Liou and Shih note that Morkovin's hypothesis does not apply for high-speed shear layers with shock or expansion wave interaction (Liou and Shih, 1991). Further, Spina et al. (Spina et al., 1994) indicate that the limitations of this hypothesis are more restrictive than originally believed. In any event, this hypothesis has led to the "ad-hoc" extrapolation of low-speed turbulence models, at all levels, to high-speed flow where the effects of compressibility are treated with correction factors (Wilcox, 1994).

Patel, Rodi and Scheuerer found that while some of these "extrapolation" models work well for flat plates, they fail for curved surfaces. They suggest refining the governing equations in the k -e model to adjusting equation constants based on empirical data (Patel, Rodi and Scheuerer, 1984). Degani and Smits found that a one-equation model provides better results than an algebraic model for compressible flow with regions of concave curvature (Degani and Smits, 1990). Again, they recommend empirical modifications to the equations to better match experimental data. These non-physical empirical corrections are prevalent in the literature (Bradshaw, 1973).

Shock waves will naturally occur in any high performance flight vehicle. As a result, they will interact with the boundary layer on the surfaces of the vehicle. The interaction between the shock wave and boundary layer is complex and difficult to model. However, its importance cannot be underemphasized. If the shock interacts with a control surface for a missile, then the resulting adverse pressure gradient could cause the missile to miss its target due to lack of control authority. If a larger control surface were used to counteract the shock effects, a negative impact on missile speed and range would have to be accepted.

### 1.2 Research Objectives

As discussed in Section 1.1, the need for quality data is clear. Thus, addressing this need is the primary goal of this thesis. Data obtained from this thesis will assist in validation of current and future CFD codes as well as to aid the understanding of the physics of these flow conditions. An additional goal is to compare turbulence data obtained via hot-wire anemometry to turbulence data obtained by Laser Doppler Velocimetry (LDV). The final goal is to verify the applicability of PIV for compressible flow analysis.

Settles and Dodson established a set of criteria that data must adhere to in order to maximize the usefulness of the data (Settles and Dodson, 1994). Their criteria are:

1. Baseline applicability: All candidate studies for use must be experiments involving turbulent flow in either the supersonic or hypersonic Mach number range.
2. Simplicity: Experimental geometries must be sufficiently simple that they may be modeled by CFD methods.
3. Specific applicability: All candidate studies passing these criteria must be capable of providing some useful test of turbulence modeling.
4. Well-defined experimental boundary conditions: All incoming conditions must be carefully documented. This includes that state of the incoming boundary layer. For studies claiming "two-dimensional" flow, data indicating the extent of the spanwise flow variations should be provided.
5. Well-defined error bounds: The experimenter must provide an analysis of the accuracy and repeatability of the data, or error bounds on the data themselves. Further, error bounds on the data must be substantiated in a quantifiable manner.
6. Adequate documentation of the data: Data must be documented and tabulated in a machine-readable format.
7. Adequate spatial resolution of the data: Experiments must present data of sufficiently high resolution, compared with the scaled flow in question, such that the key features of the flow are clearly resolved.

The data in this research were acquired with the goal of meeting these criteria.

### 1.3 Overview of Current Study

PIV was used to make mean and fluctuating velocity measurement in the AFIT Mach 3.0 wind tunnel for three pressure gradients. The results of the PIV analysis were compared to LDV measurements. The pressure gradient models that were studied include the flat plate ( ZPG ) a favorable pressure gradient
(FPG) and a combined pressure gradient (CPG). These models have been the source of extensive study at AFIT (Dotter, 1994; Miller, 1994; Luker, 1995 \& Hale, 1995).

Multiple overheat single hot-wire anemometry was used to investigate the mean and fluctuating mass flux in the boundary layer with the same three pressure gradients in Mach 2.8 flow. In addition, the ZPG was investigated under Mach 1.7 flow. Also, the power spectra of the boundary layer were investigated for three pressure gradients in Mach 2.8 flow. Power spectra data was acquired for the ZPG and FPG in Mach 1.7 flow. Finally, conventional pressure measurements (Pitot pressure and cone-static pressure) were made for the pressure gradient models in Mach 1.7, Mach 2.8 and Mach 5.0 flow.

### 1.4 Thesis Synopsis

An overview of the governing equations and a brief literature review are given in Chapter 2. The equipment used in this research effort is presented in Chapter 3. Chapter 4 documents the data analysis techniques. Chapters 5, 6 and 7 contain the results for Mach 2.8, Mach 1.7 and Mach 5.0 flow, respectively. Finally, chapter 8 contains the conclusion and recommendations drawn from this research.

## 2. Background

This section presents a discussion on the governing equations and a literature review.

### 2.1 Governing Equations

The Navier-Stokes equations are a set of coupled non-linear partial differential equations that govern all types of fluid flow, including turbulent flow. The nature of the Navier-Stokes equations makes them difficult to solve in closed form. As shown in Section 1.1, the solution time with current or near future computers for a simple grid would take far beyond the life span of any researcher. To combat these difficulties, scientists and engineers use an approximate form of the equations. The approximate forms create the problem of turbulence closure wherein the unknowns outnumber the equations.

This section assists the reader in development of the Reynolds and Favre time-averaged NavierStokes equations. The development highlights the closure problem as well as defines the crosscorrelation terms that time averaging creates. In addition, the difference between compressible flows and incompressible flows becomes apparent.

### 2.1.1 The Compressible Navier-Stokes Equations

The Navier-Stokes equations in Cartesian form can be written as (Anderson, 1984):

$$
\begin{gather*}
\frac{\mathrm{D} \rho}{\mathrm{Dt}}+\rho \nabla \cdot \mathbf{V}=0 \\
\rho \frac{\mathrm{DV}}{\mathrm{Dt}}=\nabla \cdot \frac{\bar{\pi}}{\pi}+\rho \mathbf{f}_{\mathrm{b}}  \tag{2.1}\\
\rho \frac{\mathrm{Dh}}{\mathrm{Dt}}=\frac{\mathrm{Dp}}{\mathrm{Dt}}+\frac{\mathrm{DQ}}{\mathrm{Dt}}-\nabla \cdot \mathbf{q}+\Phi
\end{gather*}
$$

where $f_{b}$ is the body force vector (normally neglected) acting on the fluid and

$$
\begin{gather*}
\frac{\mathrm{D}}{\mathrm{Dt}} \equiv \frac{\partial}{\partial \mathrm{t}}+\mathbf{V} \cdot \nabla \\
\pi_{\mathrm{ij}}=-\mathrm{p} \delta_{\mathrm{ij}}+\tau_{\mathrm{ij}} \\
\tau_{\mathrm{ij}}=\mu\left\{\left(\frac{\partial \mathrm{u}_{\mathrm{i}}}{\partial \mathrm{x}_{\mathrm{i}}}+\frac{\partial \mathrm{u}_{\mathrm{j}}}{\partial \mathrm{x}_{\mathrm{i}}}\right)-\frac{2}{3} \delta_{\mathrm{ij}} \frac{\partial \mathrm{u}_{\mathrm{k}}}{\partial \mathrm{x}_{\mathrm{k}}}\right\} \tag{2.2}
\end{gather*}
$$

$$
\Phi=\tau_{\mathrm{ij}} \frac{\partial \mathrm{u}_{\mathrm{i}}}{\partial \mathrm{x}_{\mathrm{j}}}
$$

### 2.1.2 Reynolds-Averaged Navier-Stokes Equations (RANS)

One method of simplifying the Navier-Stokes equations is known as Reynolds averaging. The Reynolds-averaged Navier-Stokes equations are derived by separating a variable into time invariant and time variant parts, as demonstrated for velocity below.

$$
\begin{equation*}
u(x, y, z, t)=\bar{u}(x, y, z, t)+u^{\prime}(x, y, z, t) \tag{2.3}
\end{equation*}
$$

$\bar{u}$ is the time-averaged velocity. Note that the time-averaged velocity itself might also be a function of time. The time-averaged velocity is defined as (Wilcox, 1994)

$$
\begin{equation*}
\bar{u}(x, y, z, t) \equiv \frac{1}{T} \int_{\mathrm{t}}^{\mathrm{t}+\mathrm{T}} \mathrm{u}(\mathrm{x}, \mathrm{y}, \mathrm{z}, \mathrm{t}) \mathrm{dt} \tag{2.4}
\end{equation*}
$$

$T$ is a characteristic time that must be larger than the period of the velocity fluctuations, but smaller than the period of the mean flow.

When the separated variables defined in Equation (2.3) are substituted into the Navier-Stokes equations and the rules of time averaging are applied, the RANS are created. The key rules of time averaging is that an average of a single fluctuating component is zero but an average of two or more fluctuating components are not equal to zero. Assuming zero body forces, the RANS equations are found to be (Wilcox, 1994)

$$
\begin{gather*}
\frac{\partial \bar{\rho}}{\partial \mathrm{t}}+\frac{\partial\left(\overline{\rho \mathrm{u}_{\mathrm{j}}}+\mathrm{m}_{\mathrm{j}}^{\mathrm{TR}}\right)}{\partial \mathrm{x}_{\mathrm{j}}}=0 \\
\frac{\partial\left(\overline{\rho \mathrm{u}_{\mathrm{i}}}+\overline{\rho^{\prime} \mathrm{u}_{\mathrm{i}}^{\prime}}\right)}{\partial \mathrm{t}}+\frac{\partial\left(\overline{\rho \mathrm{u}_{\mathrm{i}}} \overline{\mathrm{u}_{\mathrm{j}}}\right)}{\partial \mathrm{x}_{\mathrm{j}}}=-\frac{\partial \overline{\mathrm{p}}}{\partial \mathrm{x}_{\mathrm{i}}}+\frac{\partial\left(\overline{\tau_{\mathrm{ij}}}+\tau_{\mathrm{ij}}^{\mathrm{TR}}\right)}{\partial \mathrm{x}_{\mathrm{j}}}  \tag{2.5}\\
\frac{\partial\left(\overline{\rho \mathrm{e}_{\mathrm{o}}}+\overline{\rho^{\prime} \mathrm{h}_{\mathrm{o}}{ }^{\prime}}\right)}{\partial \mathrm{t}}+\frac{\partial\left(\overline{\rho \mathrm{h}_{\mathrm{o}}} \overline{\mathrm{u}_{\mathrm{j}}}\right)}{\partial \mathrm{x}_{\mathrm{j}}}=\frac{\partial\left(\overline{\mathrm{u}_{\mathrm{i}} \tau_{\mathrm{ij}}}+\overline{\mathrm{u}_{\mathrm{i}}{ }^{\prime} \tau_{\mathrm{ij}}{ }^{\prime}}-\overline{\mathrm{q}_{\mathrm{j}}}-\mathrm{q}_{\mathrm{j}}^{\mathrm{TR}}\right)}{\partial \mathrm{x}_{\mathrm{j}}}
\end{gather*}
$$

where

$$
\begin{align*}
& \mathrm{m}_{\mathrm{j}}^{\mathrm{TR}}=\overline{\rho^{\prime} \mathbf{u}^{\prime}} \\
& \overline{\tau_{i j}}=\mu\left\{\left(\frac{\partial \overline{u_{i}}}{\partial \mathbf{x}_{j}}+\frac{\partial \overline{u_{j}}}{\partial \mathbf{x}_{\mathrm{i}}}\right)-\frac{2}{3} \delta_{\mathrm{ij}} \frac{\partial \overline{\mathrm{u}_{\mathrm{k}}}}{\partial \mathrm{x}_{\mathrm{k}}}\right\} \\
& \tau_{\mathrm{ij}}{ }^{\prime}=\mu\left\{\left(\frac{\partial \mathrm{u}_{\mathrm{i}}{ }^{\prime}}{\partial \mathrm{x}_{\mathrm{j}}}+\frac{\partial \mathrm{u}_{\mathrm{j}}{ }^{\prime}}{\partial \mathrm{x}_{\mathrm{i}}}\right)-\frac{2}{3} \delta_{\mathrm{ij}} \frac{\partial \mathrm{u}_{\mathrm{k}}{ }^{\prime}}{\partial \mathrm{x}_{\mathrm{k}}}\right\} \\
& \tau_{i j}^{T R}=-\bar{\rho} \overline{u_{i}{ }^{\prime} u_{j}{ }^{\prime}}-\overline{u_{i} \rho^{\prime} u_{j}{ }^{\prime}}-\overline{u_{j} \rho^{\prime} u_{i}{ }^{\prime}}-\overline{\rho^{\prime} u_{i}{ }^{\prime} u_{j}{ }^{\prime}}  \tag{2.6}\\
& \overline{q_{i}}=-k \frac{\partial \bar{T}}{\partial \mathrm{x}_{\mathrm{i}}} \\
& q_{i}{ }^{T R}=-\bar{\rho} \overline{u_{i}{ }^{\prime} h_{o}{ }^{\prime}}-\overline{h_{o} \rho^{\prime} u_{i}{ }^{\prime}}-\overline{u_{j} \rho^{\prime} h_{o}{ }^{\prime}}-\overline{\rho^{\prime} u_{i}{ }^{\prime} h_{o}{ }^{\prime}}
\end{align*}
$$

The subscript $o$ in Equations (2.5) and (2.6) indicates stagnation quantities. The turbulence terms (noted by a superscript $T R$ ), $\mathrm{m}_{\mathrm{j}}^{\mathrm{TR}}, \tau_{\mathrm{ij}}^{\mathrm{TR}}$ and $\mathrm{q}_{\mathrm{i}}^{\mathrm{TR}}$ are known as the apparent mass flux, turbulent shear stress (Reynolds) and the turbulent heat flux, respectively. These are the "extra" terms discussed in the closure problem. For incompressible flow, the apparent mass flux and the final three terms of the turbulent Reynolds shear stress and turbulent heat flux are zero. The triple correlation in the RHS of the turbulent Reynolds shear stress and turbulent heat flux is usually assumed to be much smaller than the other terms. PIV and LDV directly measure the velocity and are more appropriate for Reynolds stress calculations.

### 2.1.3 Favré-Averaged Navier-Stokes Equations (FANS)

An alternate method of simplifying the Navier-Stokes equations is known as Favré averaging. The Favré-averaged Navier-Stokes equations are similar to the Reynolds averaged equations except the mass dependent terms (e.g., terms with $\rho$ ) are replaced by a mass averaged term, $\tilde{\phi}$, plus a fluctuation term, $\phi^{\prime \prime}$ :

$$
\begin{equation*}
\phi=\tilde{\phi}+\phi^{\prime \prime} \quad \phi^{\prime \prime} \equiv \frac{\overline{\rho \phi}}{\bar{\rho}} \tag{2.7}
\end{equation*}
$$

Mass independent terms (such as $\rho$ and p ) are replaced by time variant and time invariant terms following the same approach as in Reynolds averaging. Note that the incompressible Favré -averaged equations reduce to the Reynolds-averaged equations. A key difference between Reynolds and Favré averaging becomes clear after inspection of Equations (2.3) and (2.7). This inspection shows that $\overline{\phi^{\prime}}=\mathbf{0}$ but $\overline{\phi^{\prime \prime}} \neq 0$. With Equation (2.7) and the time averaging rules, the FANS equations are found to be (Wilcox, 1994)

$$
\begin{align*}
& \frac{\partial \bar{\rho}}{\partial \mathrm{t}}+\frac{\partial\left(\bar{\rho} \tilde{\mathrm{u}}_{\mathrm{j}}+\mathrm{m}_{\mathrm{j}}^{\mathrm{TF}}\right)}{\partial \mathrm{x}_{\mathrm{j}}}=0 \\
& \frac{\partial\left(\bar{\rho} \tilde{\mathbf{u}}_{\mathrm{i}}\right)}{\partial \mathrm{t}}+\frac{\partial\left(\bar{\rho} \tilde{\mathrm{u}}_{\mathrm{i}} \tilde{\mathrm{u}}_{\mathrm{j}}\right)}{\partial \mathbf{x}_{\mathrm{j}}}=-\frac{\partial \overline{\mathrm{p}}}{\partial \mathbf{x}_{\mathrm{i}}}+\frac{\partial\left(\overline{\tau_{\mathrm{ij}}}+\tau_{\mathrm{ij}}{ }^{\mathrm{TF}}\right)}{\partial \mathbf{x}_{\mathrm{j}}}  \tag{2.8}\\
& \frac{\partial\left(\bar{\rho} \widetilde{\mathrm{e}}_{\mathrm{o}}\right)}{\partial \mathrm{t}}+\frac{\partial\left(\overline{\mathrm{\rho} \mathrm{~h}_{\mathrm{o}}} \tilde{\mathrm{u}}_{\mathrm{o}}\right)}{\partial \mathbf{x}_{\mathrm{j}}}=\frac{\partial\left(\widetilde{\mathrm{u}}_{\mathrm{i}} \bar{\tau}_{\mathrm{ij}}+\overline{\mathrm{u}_{\mathrm{i}}^{\prime \tau_{\mathrm{ij}}}}-\overline{\mathrm{q}_{\mathrm{j}}}-\mathrm{q}_{\mathrm{j}}^{\mathrm{TF}}\right)}{\partial \mathbf{x}_{\mathrm{j}}}
\end{align*}
$$

where

$$
\begin{align*}
& \mathrm{m}_{\mathrm{j}}^{\mathrm{TF}}=0 \\
& \overline{\tau_{i j}}=\mu\left\{\left(\frac{\partial \widetilde{\mathbf{u}}_{i}}{\partial \mathrm{x}_{\mathrm{j}}}+\frac{\partial{\widetilde{\mathbf{u}_{j}}}_{\partial \mathrm{x}_{\mathrm{i}}}}{}\right)-\frac{2}{3} \delta_{\mathrm{ij}} \frac{\partial \widetilde{u}_{\mathrm{k}}}{\partial \mathrm{x}_{\mathrm{k}}}\right\}+\mu\left\{\left(\frac{\partial \overline{u_{i}^{\prime \prime}}}{\partial \mathrm{x}_{\mathrm{j}}}+\frac{\partial \overline{u_{j}^{\prime \prime}}}{\partial \mathrm{x}_{\mathrm{i}}}\right)-\frac{2}{3} \delta_{\mathrm{ij}} \frac{\partial \overline{u_{k}^{\prime \prime}}}{\partial \mathrm{x}_{\mathrm{k}}}\right\}  \tag{2.9}\\
& \tau_{i j}{ }^{\mathrm{TF}}=-\bar{\rho} \overline{u_{i}^{\prime \prime} u_{j}^{\prime \prime}} \\
& \mathbf{q}_{\mathrm{i}}{ }^{\mathrm{TF}}=-\bar{\rho} \overline{\mathrm{u}_{\mathrm{j}}^{\prime \prime} \mathrm{h}_{\mathrm{o}}^{\prime \prime}}
\end{align*}
$$

Note that the apparent mass in the Favré averaged equation is identically zero.
Researchers have found that the FANS work well for compressible shear layers but fail when used to model shear layers in the vicinity of shocks or expansion waves (Wilcox, 1994).

### 2.1.4 Comparison Between Favré-Averaged and Reynolds-Averaged N-S Equations

An interesting relationship can be developed by equating the Reynolds averaged terms to the Favré averaged terms. The following example uses the $u$ component of velocity. By the definition of Reynolds averaged and Favré averaged variables

$$
\begin{equation*}
\mathbf{u}=\overline{\mathbf{u}}+\mathbf{u}^{\prime}=\widetilde{\mathbf{u}}+\mathbf{u}^{\prime \prime} \tag{2.10}
\end{equation*}
$$

Taking the time average and rearranging Equation (2.10) results in

$$
\begin{equation*}
\overline{\mathbf{u}}-\tilde{\mathbf{u}}=\mathbf{u}^{\prime \prime}=\frac{\overline{\rho \mathbf{u}}}{\bar{\rho}}=\frac{\overline{\left(\rho+\rho^{\prime}\right)\left(u+u^{\prime}\right)}}{\bar{\rho}} \tag{2.11}
\end{equation*}
$$

Expanding the RHS of Equation (2.11) and by use of the properties of time averaging

$$
\begin{equation*}
u^{\prime \prime}=\frac{\overline{\rho^{\prime} u^{\prime}}}{\bar{\rho}} \tag{2.12}
\end{equation*}
$$

Thus, Equation (2.12) defines the relationship between Reynolds averaged variables and Favré averaged variables.

Through a similar set of identities, the relationship between the Reynolds shear stress and Favré shear stress can be determined.

$$
\begin{equation*}
\bar{\rho} \overline{u_{i}^{\prime} u_{j}^{\prime}}=\bar{\rho} \overline{u_{i}^{\prime \prime} u_{j}^{\prime \prime}}-\bar{\rho} \overline{u_{i}^{\prime \prime}} \overline{u_{i}^{\prime \prime}} \tag{2.13}
\end{equation*}
$$

The last term on the RHS of Equation (2.13) is of fourth order. Usually, this term is much smaller than the other terms and is neglected. Thus,

$$
\begin{equation*}
\tau_{i j}^{\mathrm{TR}} \approx \tau_{\mathrm{ij}}^{\mathrm{TF}} \tag{2.14}
\end{equation*}
$$

Thus, hot-wire anemometry can be used to measure Reynolds stress as well as Favré stress.

### 2.2 Literature Review

Undoubtedly, many experiments have been conducted to measure mean flow quantities for supersonic and hypersonic flows; however, these experiments mainly concern ZPG flow (Robinson et al., 1983; Johnson and Rose, 1975; Fernando and Smits, 1990; Kistler, 1959 \& Parrott et al., 1989) and rarely
measure turbulence quantities. Smith and Smits give as the reason for this lack of data as "... is that they are difficult to obtain. Reliable and accurate measurements of turbulence properties are difficult to make in any supersonic flow, and the difficulties are usually more extreme in the presence of flow distortions" (Smith and Smits, 1994). For many studies where turbulence quantities are measured, the quality of the data was found to be flawed and unusable (Bradshaw and Ferris, 1971). Settles and Dodson (Settles and Dodson, 1994) echo this concern. They surveyed over one hundred shock/boundary layer interaction studies and found only 19 that met their turbulence criteria. A brief summary of relevant research on this topic is presented below. The review presented here is not all inclusive. Instead papers relevant to the present thesis objectives (Section 1.2) are presented. A complete review can be found in Spina et al., 1994. In addition, a review of flow visualization research is given.

Bradshaw and Ferriss studied the turbulent kinetic energy equation for boundary layers in compressible flow with arbitrary pressure gradients (Bradshaw and Ferriss, 1971). They obtained correct trends for the ZPG case but were unable to find quality pressure gradient data to compare with their results.

Bradshaw conducted a series of calculations designed to show that bulk compression or dilation affects turbulent shear layers (Bradshaw, 1974). Bradshaw found bulk compression or dilation had a larger effect than was expected and should be accounted for in future studies. Fluctuating turbulence measurement were not made.

Adverse pressure gradients have been found to destabilize the boundary layer. In other words, the turbulent fluctuating properties increase (Dotter, $1994 \&$ Hale, 1995). Adverse pressure gradients can be created by a shock interacting with the boundary layer or by concave curvature. Lewis and Gran (Lewis and Gran, 1972) created an adverse pressure gradient by shock wave boundary layer interaction. They found that turbulent results compare favorably with low-speed results. Fernando and Smits (Fernando and Smits, 1990) suggest that the turbulence quantities undergo a mild amplification while the mean flow structure remains unchanged. Both of these experiments are suspect because the do not include the effect of streamline curvature on the turbulence quantities.

Recall that adverse pressure gradients can be created by regions of concave curvature. Jayaram, Taylor and Smits found that the turbulence levels increased due to the combined effect of the pressure gradient, streamline curvature and bulk compression (Jayaram, Taylor and Smits, 1986). Interestingly, they found, for two of their three cases, that the weak shock wave had no effect on the turbulent quantities.

Few FPG studies have been made. The few that do exist suggest that the FPG acts to damp out the turbulent fluctuations. Bulk dilation and supersonic stream tube expansion provide the most likely explanations of this phenomena (Luker et al., 1997). Studies by Luker (Luker, 1995) and Miller (Miller, 1994) confirm the dampening effect of the pressure gradient.

Flow visualization provides an important diagnostic tool to understand the effect of the pressure gradient on the structure of the flow. Spina (Spina, 1988) discovered a wide range of temporal and spatial scales in a ZPG flow. Application of a FPG was found to increase the boundary layer thickness and assist the development of the elongated structures (Arnette et al., 1995).

Development of particle induced velocimetry is the latest technique available for researchers and scientists. This technique has been proven for a wide range of applications ranging from wake flow (Yao and Pachal, 1994) to supersonic flat plate flow (Glawe et al., 1996). However, its applicability to compressible flow with pressure gradients has yet to be established.

## 3. Facilities and Equipment

This chapter discusses the facilities and equipment used during the course of this research. Data were collected in both the AFIT Mach 3.0 blow-down, pressure-vacuum wind tunnel, the AFIT Mach 1.7 indraft wind tunnel and the AFIT Mach 5.0 blow-down wind tunnel. In addition, this chapter describes wind tunnel configurations, the data collection system, probes, particle image velocimetry (PIV) and computer equipment.

### 3.1 Mach 3.0 Wind Tunnel

The AFIT Mach 3.0 wind tunnel is a blow-down, vacuum system. One end of the tunnel is connected to vacuum tanks, and the other to a continuous high-pressure air supply. This combination of high-pressure air ( 0.69 MPa ) and vacuum provided the pressure differential for the tunnel to operate at its design conditions. A schematic of the AFIT Mach 3.0 wind tunnel is shown in Figure 3-1.


Figure 3-1: AFIT Mach 3.0 wind tunnel schematic

### 3.1.1 High-Pressure Air Supply

One Atlas Copco GAU 807 air compressor provided high pressure air nominally at 0.69 MPa with a flow rate of $0.5 \mathrm{~kg} / \mathrm{sec}$. Two Pioneer Air Systems, Inc. model R500A Refrigerant Air Dryers dried the air prior to the air entering the high-pressure air supply system. The air was filtered a second time by
a centrifugal moisture and particle separator in conjunction with several layers of Filtrite ${ }^{8}$ particle air filters. This system supplied high pressure air to other laboratories in AFIT causing minor fluctuations in air pressure which had a negligible effect on tunnel operation.

### 3.1.2 Vacuum System

Three Stokes Penwalt model 212-11 MicroVac pumps were used to evacuate the 16.0 cubic meter vacuum tank system. With all three vacuum pumps in operation, the tanks could be evacuated to less than 5.0 mm of Hg in approximately six minutes. Wind tunnel operation began only when the pressure in the vacuum tanks reached 5.0 mm of Hg . This level of pressure allowed approximately 25 seconds of uninterrupted operation before the tunnel unstarted.

### 3.1.3 Plenum Chamber

The Plenum chamber's pressure ( $\mathrm{P}_{\mathrm{t}}$ ) and temperature $\left(\mathrm{T}_{\mathrm{t}}\right)$ were maintained at $2.19 \times 10^{5} \pm 11.8$ $\times 10^{3} \mathrm{~Pa}$ and $295 \pm 2 \mathrm{~K}$, respectively during the testing period. An Endevco 0.69 MPa pressure transducer was used to measure the pressure and an Omega Enginecring type K thermocouple was used to measure the temperature. The pressure and temperature were measured upstream of the nozzle but downstream of the flow straighteners.

### 3.1.4 Mach 3.0 Nozzle

A converging-diverging nozzle created the supersonic flow with a nominal Mach number of 2.9 with a measurement uncertainty of $1.8 \%$ (Huffman, Mach 2.9 lab manual). The distance from the throat to the exit was 27.0 cm and the nozzle exit cross section was $6.35 \times 6.35 \mathrm{~cm}$. The freestream turbulence was determined to be $0.8 \%$ with a standard deviation of $0.2 \%$ (Huffiman, Mach 2.9 lab manual). At the measurement locations for the present experiments ( $\mathrm{x}_{\mathrm{ts}}>68 \mathrm{~cm}$ ), the Mach number was 2.8. This lower value in Mach number was the result of boundary layer growth.

### 3.1.5 Test Sections

The test sections of this tunnel were designed to be modular and completely interchangeable. Each test section was 32.8 cm long with inner dimensions of 6.35 cm square. The floor and ceiling sections were constructed out of aluminum alloy 1.91 cm thick. All the seams were fitted with rubber orings to tightly seal the sections together. The sections were also fitted with oversize holes to allow for individual section adjustment. This adjustment minimized the effects of expansions or shocks caused by non-aligned seams. The first section consisted of a flat plate ceiling and floor with either aluminum or Plexiglas walls. No measurements were made in this section.

The second test section consisted of a flat plate floor and three different ceilings; a flat plate, a favorable pressure gradient section and an adverse pressure gradient section. The pressure gradient sections matched a third-order polynomial defined by $y(x)=a_{0}+a_{1} x+a_{2} x^{2}+a_{3} x^{3} . y(x)$ represents the distance from a flat plate ceiling, and x is a scaled distance from the end of the test section. The coefficients for the pressure gradient polynomial are given in Table 3-1.

Table 3-1: Pressure gradient test section curve coefficients

| Model | $\mathrm{a}_{0}$ | $\mathrm{a}_{1}$ | $\mathrm{a}_{2}$ | $\mathrm{a}_{3}$ |
| :--- | :--- | :--- | :--- | :--- |
| FPG | -0.2078 | 0.0897 | -0.0095 | -0.0360 |
| CPG | 1.1858 | -0.5410 | 0.0748 | -0.0028 |

These pressure gradient surfaces have been studied extensively in the Mach 3.0 wind tunnel (Hale, 1995; Luker 1995; Dotter, 1994; Miller, 1994). Figure 3-2 shows the pressure gradient contours and the data collection locations in the Mach 2.8 tunnel.

Data were collected at 71 cm from the throat for the flat plate surface, 71.5 cm from the throat for the FPG and $68 \& 71 \mathrm{~cm}$ from the throat for the CPG. The walls for this section were either Plexiglas or optical glass as discussed by Hale and Luker (Hale, 1995; Luker, 1995). The optical glass was used when flow visualization data was collected. In addition, an optical glass floor was inserted for Particle Image Velocimetry (PIV) runs. A schematic of the optical glass floor is shown as Figure 3-3.


Figure 3-2: Pressure gradient contours \& data collection locations


Figure 3-3: Optical glass floor schematic

### 3.2 Mach 1.7 Wind Tunnel

The AFIT Mach 1.7 wind tunnel is an indraft, vacuum system. One end of the tunnel is connected to vacuum tanks, and the other end is open to the atmosphere. This combination of atmospheric air and vacuum provided the pressure differential for the tunnel to operate at its design conditions. The AFIT Mach 1.7 tunnel was identical to the AFIT Mach 2.8 tunnel with some exceptions. The exceptions are a different nozzle and a different source of high-pressure air. The same models and test sections were used in the Mach 1.7 tunnel as were used in the Mach 2.8 tunnel.

### 3.2.1 Plenum Chamber

The Plenum chamber's pressure $\left(\mathrm{P}_{\mathbf{t}}\right)$ and temperature $\left(\mathrm{T}_{\mathrm{t} 1}\right)$ were maintained at $9.47 \times 10^{4} \pm 0.51 \times 10^{3} \mathrm{~Pa}$ and $295 \pm 2 \mathrm{~K}$, respectively during the testing period. An Endevco 0.104 MPa
pressure transducer was used to measure the pressure and an Omega Engineering type K thermocouple was used to measure the temperature. The pressure and temperature were measured upstream of the nozzle but downstream of the flow straighteners.

### 3.2.2 Mach 1.7 Nozzle

A converging-diverging nozzle created the supersonic flow with a nominal Mach number of 1.7. The distance from the throat to the exit was 27.0 cm and the nozzle exit cross section was $6.35 \times 6.35 \mathrm{~cm}$. The freestream turbulence was determined to be on the order of $1.0 \%$.

### 3.3 Mach 5.0 Wind Tunnel

The AFIT Mach 5.0 wind tunnel is a blow-down wind tunnel system. One end of the tunnel is connected to ambient air, and the other end to the 17.9 MPa air tank. This combination of high-pressure air and ambient air provides the pressure differential for the tunnel to operate at its design conditions. A schematic of the AFIT Mach 5.0 wind tunnel is shown in Figure 3-4.


Figure 3-4: Mach 5.0 wind tunnel schematic

### 3.3.1 Air Supply

A 1.25 cubic meter tank supplied the high pressure air, at a mass flow rate of $1.25 \mathrm{~kg} / \mathrm{sec}$. A four cycle Eagle compressor Model HW40HH3 compressed the ambient source air and a Gas Driving Inc., dryer dried the air prior to it being inserted into the tank. The compressor and dryer system could charge the external tank at a rate of $0.69 \mathrm{MPa} /$ hour. Two 20 second runs reduced the tank pressure to 1,800 psig. An Eagle brand regulator Model 03174506 throttled the air to 7.59 MPa before the nozzle throat.

### 3.3.2 Pebble Bed Heaters

Three pebble bed heaters were used to heat the high pressure air before it entered the nozzle throat. This raised the total temperature of the air to 375 K thereby reducing the possibility of oxygen liquefaction in the test section. The pebble bed heaters consisted of three cylinders 81 cm long with a diameter of 15 cm filled with ceramic pebbles. These ceramic pebbles were heated by 0.69 MPa lab air passing through two Reheat electric heaters, model HDA-2-12. The electric heaters could heat two of the three pebble beds at a rate of $19 \mathrm{~K} /$ hour when the pebble beds temperature was 315 K . Note this rate decreased as the temperature of the pebble beds increased. Each heater allowed 30 seconds of continuous tunnel operation before the heaters cooled to ambient conditions.

### 3.3.3 Plenum Chamber

The Plenum chamber's pressure and temperature were monitored during each test run. An Endevco 3.45 MPa pressure transducer measured the pressure and an Omega Engineering type K thermocouple measured the temperature upstream of the nozzle. The temperature was recorded by the data acquisition system.

### 3.3.4 Mach 5.0 Nozzle

A converging-diverging nozzle created the supersonic flow with a nominal Mach number of 5.0. The distance from the throat to the exit was 30.5 cm and the nozzle exit cross section was $7.62 \times 7.62 \mathrm{~cm}$.

### 3.3.5 Test Sections

The test sections of this tunnel were also designed to be modular and completely interchangeable. Each test section was 25.4 cm long with inner dimensions of 7.62 cm square. The ceiling sections were constructed out of aluminum alloy 2.54 cm thick. The walls and floor were constructed out of steel 1.25 cm thick. The walls had two circular optical quality glass inserts used for flow visualization. All the seams were fitted with rubber 0 -rings to tightly seal the sections together. The sections were also fitted with oversize holes to allow for individual section adjustment similar to the Mach 2.8 tunnel. This adjustment minimized the effects of expansions or shocks caused by non-aligned seams.

The test section ceilings consisted of a flat plate floor and three different ceilings; a flat plate, a favorable pressure gradient section and an adverse pressure gradient section. The pressure gradient sections matched the profiles used in the Mach 2.8 tunnel and are given in Section 3.1.5.

It was discovered that the ceramic pebbles in the pebble bed heater shed small particles during tunnel operation. These particles were carried by the high speed air through the tunnel. Due to the highspeed of these particles, they abraded the inner surface of the wind tunnel. Due to the risk of hot-wires breaking, only conventional pressure data were taken in the Mach 5.0 tunnel. Data were taken at $\mathrm{x}_{\mathrm{ts}}=$ 6.35 cm for the FPG and ZPG configuration and at $\mathrm{x}_{\mathrm{ts}}=6.35 \mathrm{~cm}$ and 5.1 cm for the CPG configuration. In addition, wall static pressure data were taken along the surface of the FPG and CPG models. The CPG and FPG contours are the same shape as what was used in the Mach 2.8 tunnel except the contours begin at $\mathrm{x}_{\mathrm{ts}}=1.90 \mathrm{~cm}$ for the FPG and $\mathrm{x}_{\mathrm{ts}}=3.81 \mathrm{~cm}$ for the CPG.

### 3.4 Data Acquisition Equipment

A Nicolet Multipro was used to collect all the data used in this experiment. The Nicolet had four A/D converter boards each with four input / output BNC connectors. Inputs to the Nicolet consisted of voltage from the Endevco pressure conditioners or the bridge voltage from the IFA-100. The maximum sampling rate for each board was 1 MHz and the maximum number of discrete data points able to be collected per A/D board was of 262,144 . The sampling rate decreased linearly with the number of active input / output channels. Thus, if four data inputs were collected, the maximum sampling rate per channel was 250 KHz (Nicolet Systems Operation, 1991).

### 3.5 Probes

Pressure probes and hot-film probes were used to measure the flow conditions for this experiment. Two types of pressure probes were used, a Pitot probe and two cone-static probes. A $10^{\circ}$ axisymmetric cone-static probe was used in the Mach 2.8 and Mach 1.7 tunnel and a $20^{\circ}$ axisymmetric cone was used in the Mach 5.0 tunnel. Two types of hot-film probes were used to measure flow velocity, fluctuations and power spectra, single hot-film probe and a single hot-wire probe.

### 3.5.1 Pressure Probes

Measurement of $P_{t 1}$ was obtained using the Pitot tube in the wind tunnel plenum. A $10^{\circ}$ axisymmetric cone-static pressure probe collected $p_{c}$ and a total pressure probe collected $P_{t 2}$ data at the downstream locations noted in Sections 3.1.5 and 3.2.5. The probes were aligned so they were normal to the wind tunnel ceiling. The wind tunnel ceilings were fitted with static pressure ports. The probes were connected to two $103.4 \times 10^{3} \pm 5.6 \times 10^{3} \mathrm{~Pa}$ Endevco pressure transducers which were connected in turn to Endevco model 4225 signal conditioners. The signal conditioners filtered the high frequency noise and provided an output voltage for the data acquisition system. In addition, the signal conditioners were
calibrated for their transducer and recalibrated periodically. The local atmospheric pressure was recorded on a DPI 141 digital barometer. Pressure probes were sampled at 500 Hz by the Nicolet.

The probes were constructed out of stainless steel ( 0.3175 cm diameter stem) with the head of the probe extending approximately 1.9 cm from the stem. The axisymmetric cones each had four pressure taps $90^{\circ}$ apart. The taps opened to a common settling chamber which corrected for any pressure differences between pressure taps.

It was discovered that the probes deflected approximately $2^{\circ}$ during tunnel operation. Accordingly, the probes were pre-positioned so this deflection moved the probe to the desired measuring location during tunnel operation.

### 3.5.2 Hot-Wire and Hot-Film Probes

Two types of hot-film probes were used in this experiment, a single hot-film (TSI model 121820) and a single hot-wire (TSI model 1218-T1.5). Each probe had a thickness of $51.0 \mu \mathrm{~m}$, sensing length of 1.0 mm , nominal resistance of 5.5 Ohms and a temperature coefficient of resistance of $0.24 \%$. The hotwire had a higher frequency response that the hot-film and was exclusively used for power spectra data.

A TSI 100 Intelligent Flow Analyzer constant temperature anemometer (CTA) was used for this study (IFA 100 instruction manual, 1987). The IFA 100 is composed of the Model 150 anemometer and the Model 157 Signal Conditioner. The IFA 100 was operated in the $1: 1$ bridge mode. The external resistance and overheat ratio were set by a series of external resistors. The bridge voltage output as measured by the TSI 100 was recorded by the Nicolet Multipro at a sampling frequency of 25 KHz for velocity and total temperature data and a frequency of 500 KHz for power spectra data.

All hot-films were calibrated before data acquisition occurred. The hot-film probes were calibrated by increasing $\mathrm{P}_{\mathrm{t} 2}$ while the probe was in the free-stream of the respective wind tunnel. The sampling frequency during calibration was 200 Hz . As with the pressure probes, the hot-film probes deflected $2^{\circ}$ during tunnel operation and were pre-positioned to correct for this deflection.

### 3.6 Particle Image Velocimetry

A two laser PIV system was used for PIV measurement as well as flow visualization. A schematic of the WL/POPT PIV system is presented in Figure 3-5.


Figure 3-5: PIV schematic

Two Continuum power supplies Model SLI-10 were used to power the two Surlite Nd:YAG lasers. The Surlite lasers each produced two coherent laser beams with wavelengths ( $\lambda$ ) of 532 and $1,064 \mathrm{~nm}$. The beam splitter was used to eliminate the higher wavelength beam. One of the beams passed through a Princeton Optics Model RC1000 Raman cell that contained $\mathrm{N}_{2}$ and $\mathrm{H}_{2}$ at a pressure of 65.6 MPa . By Raman scattering off the $\mathrm{N}_{2}$, the $\lambda=532 \mathrm{~nm}$ (green) beam was converted to $\lambda=607 \mathrm{~nm}$ (red). The beams were combined and a series of optics formed a laser sheet with a thickness $<1 \mathrm{~mm}$. The laser sheet was manually aligned and focused on the ceiling of the test section.

Each laser was pulsed at 10 Hz with a 10 ns pulse width (Glawe et al, 1996). The time delay between the green and red pulse was adjusted by the Taitech Inc. PIV control box and monitored by the Gould Model 4074 oscilloscope. The nominal time delay was set to 300 ns and was recorded after each data run.

Images were recorded by a Kodak Model 460C CCD camera with a Nikon N90 lens. The lens settings were F 4.2 with an exposure time of 0.125 seconds. A maximum of three images were taken per data run

### 3.7 Computer Equipment

Personal computers both at AFIT and at the authors residence were used for data reduction, traverse operation and IFA 100 operation. A Digital Equipment Corporation alpha station was also used for data reduction.

## 4. Data Reduction Techniques

This chapter presents the data reduction techniques used in this experiment. First, the mean flow data reduction techniques are discussed. They are followed by the hot-wire data reduction techniques and the PIV data reduction techniques.

### 4.1 Mean Flow Data Reduction

This section provides the background and methodology necessary to compute the mean flow properties.

### 4.1.1 Calculation of Mach Number, Velocity and Fluid Properties

Mean flow data were recorded for each pressure gradient model studied in this experiment. The mean flow data recorded included upstream Pitot pressure ( $\mathrm{P}_{\mathrm{t}}$ ), upstream total pressure ( $\mathrm{T}_{11}$ ), downstream Pitot pressure $\left(\mathrm{P}_{\mathrm{t} 2}\right)$ and downstream cone-static pressure $\left(\mathrm{p}_{\mathrm{c}}\right)$. $\mathrm{P}_{\mathrm{t} 2}$ and $\mathrm{p}_{\mathrm{c}}$ were used to calculate the Mach number for the ZPG model in the Mach 2.8 wind tunnel via the following formula (Bowersox, 1992)

$$
\begin{equation*}
M=\left(-0.052976+4.684 x-18.678 x^{2}+50.7006 x^{3}-54.1577 x^{4}\right)^{-1} \tag{4.1}
\end{equation*}
$$

In equation (4.1), $\mathrm{x}=\mathrm{p}_{\mathrm{c}} / \mathrm{P}_{\mathrm{t} 2}$. Equation (4.1) is valid for Mach numbers between 1.5 to 4.4 and has a standard deviation of $0.06 \%$. For the CPG and FPG models, iteration on the Rayleigh supersonic formula (Liepmann and Rosko, 1957) was used to calculate the Mach number.

With the Mach number and $\mathrm{T}_{\mathrm{tI}}$ known, the local temperature was calculated using the following isentropic relationship. (Note, it was assumed that $\mathrm{T}_{\mathrm{t}}$ was constant throughout the boundary layer.)

$$
\begin{equation*}
\frac{\mathrm{T}_{\mathrm{t} 1}}{\mathrm{~T}}=1+\frac{\gamma-1}{2} \mathrm{M}^{2} \tag{4.2}
\end{equation*}
$$

The local velocity was calculated using the definitions of the Mach number and the speed of sound. The derived velocity is

$$
\begin{equation*}
\frac{\mathrm{u}}{\mathrm{u}_{\mathrm{e}}}=\frac{\mathrm{M}}{\mathrm{M}_{\mathrm{e}}} \sqrt{\frac{1+\frac{\gamma-1}{2} \mathrm{M}^{2}}{1+\frac{\gamma-1}{2} \mathrm{M}_{\mathrm{e}}^{2}}} \tag{4.3}
\end{equation*}
$$

where $\mathrm{M}_{\mathrm{e}}$ was determined by selecting the largest Mach number from analysis of the pressure data and the output of Equation (4.1)

### 4.1.1.1 Pressure and Temperature Methodology - ZPG Flow

For the ZPG case it was assumed that the pressure gradient normal to the surface was zero. A modified Crocco-Busemann relation yields (Van Driest, 1951)

$$
\begin{equation*}
\frac{\rho_{w}}{\rho}=\frac{T_{w}}{T}=1+B^{\prime}\left(\frac{u}{u_{e}}\right)-A^{\prime 2}\left(\frac{u}{u_{e}}\right)^{2} \tag{4.4}
\end{equation*}
$$

$\rho_{w}$ was found from the perfect gas relation and measured $p_{w}$ values. $A^{\prime}$ and $B^{\prime}$ are given by

$$
\begin{equation*}
\mathrm{B}^{\prime}=\frac{1+\left(\frac{\gamma-1}{2}\right) \mathrm{rM}_{\mathrm{e}}^{2}}{\mathrm{~T}_{\mathrm{w}} / \mathrm{T}_{\mathrm{e}}}-1 \quad \mathrm{~A}^{\prime 2}=\frac{\left(\frac{\gamma-1}{2}\right) \mathrm{rM}_{\mathrm{e}}^{2}}{\mathrm{~T}_{\mathrm{w}} / \mathrm{T}_{\mathrm{e}}} \tag{4.5}
\end{equation*}
$$

r is the recovery factor and can be expressed as $\sqrt[3]{\mathrm{Pr}}$. For air $\mathrm{Pr}=0.707$ and $\mathrm{r}=0.892$.

### 4.1.1.2 Pressure and Temperature Methodology - FPG and CPG Flow

For pressure gradient surfaces, the pressure gradient normal to the surface is not zero due to streamline curvature. This was verified by Luker and Hale (Luker, 1995 \& Hale, 1995). Thus, a new method must be found to determine the local temperature and pressure. One method is to assume a linear variation in pressure between the boundary layer edge and the wall. Numerical studies of this flow proved this method (Fick, 1995). The wall pressure was measured by static ports and the following isentropic equation was used to compute the edge pressure.

$$
\begin{equation*}
\frac{p_{\mathrm{e}}}{\mathrm{p}_{\mathrm{o}}}=\left(1+\frac{\gamma-1}{2} \mathrm{M}^{2}\right)^{\frac{-\gamma}{\gamma-1}} \tag{4.6}
\end{equation*}
$$

Assuming a linear variation in pressure and Equation (4.2), the perfect gas law was used to calculate the density.

### 4.1.2 Calculation of Boundary Layer, Displacement and Momentum Thickness

Different boundary layer thickness parameters were calculated based on the result of the conventional pressure data. One thickness was based on the Mach number ( $\delta_{\mathrm{M}}$ ) and the other on the velocity $\left(\delta_{u}\right)$. The Mach number boundary layer thickness provides a measure of the temperature boundary layer thickness. The velocity boundary layer thickness was used to provide a more direct comparison to the LDV data (Hale, 1995 \& Luker, 1995). The local Mach number or velocity was nondimensionalized by the edge value and the location where the non-dimensional ratio was $>99 \%$ was selected as the appropriate thickness. Note, Luker and Hale (Luker, 1995 \& Hale, 1995) selected 99.5\% as their boundary layer criteria.

Numerical integration on the resulting velocity profiles and the known density profile were used to compute the momentum $(\theta)$ and displacement $\left(\delta^{*}\right)$ thickness. The compressible momentum and displacement thickness are defined as (White, 1991)

$$
\begin{align*}
& \theta \equiv \int_{0}^{\infty} \frac{\overline{\rho u}}{\rho_{\mathrm{c}} u_{c}}\left(1-\frac{\overline{\mathrm{u}}}{\mathrm{u}_{\mathrm{e}}}\right) \mathrm{dy}  \tag{4.7}\\
& \delta^{*} \equiv \int_{0}^{\infty}\left(1-\frac{\overline{\rho \mathrm{u}}}{\rho_{\mathrm{e}} \mathrm{u}_{\mathrm{e}}}\right) \mathrm{dy} \tag{4.8}
\end{align*}
$$

The incompressible momentum $\left(\theta_{i}\right)$ and displacement $\left(\delta_{i}^{*}\right)$ thickness were computed by setting the density ratio in Equations (4.7) and (4.8) to one. Due to the asymptotic nature of Equations (4.7) and (4.8), the differences in boundary layer thickness definitions between this study and Luker and Hale will be minimized in momentum and displacement thickness.

### 4.1.3 Calculation of Van Driest Velocity Profile and Wall Shear Stress

Van Driest developed the following skin friction coefficient $\left(\mathrm{C}_{\mathrm{f}}\right)$ formula for flow over a flat plate (Van Driest, 1951)

$$
\begin{equation*}
\frac{0.242}{A^{\prime} \sqrt{C_{f}\left(T_{w} / T_{e}\right)}}\left\{\sin ^{-1}\left(\frac{2 \mathrm{~A}^{\prime 2}-\mathrm{B}^{\prime}}{\sqrt{4 \mathrm{~A}^{\prime 2}+\mathrm{B}^{\prime 2}}}\right)+\sin ^{-1}\left(\frac{\mathrm{~B}^{\prime}}{\sqrt{4 \mathrm{~A}^{\prime 2}+\mathrm{B}^{\prime 2}}}\right)\right\}=\kappa+\log \left(\operatorname{Re}_{\mathrm{x}} \mathrm{C}_{\mathrm{f}}\right)-\omega \log \left(\mathrm{T}_{\mathrm{w}} / \mathrm{T}_{\mathrm{e}}\right) \tag{4.9}
\end{equation*}
$$

where $\kappa=0.41$ and $\omega=0.68$. Equation (4.9) does not have a closed form and must be solved by iteration to find $\mathrm{C}_{\mathrm{f}}$. With $\mathrm{C}_{\mathrm{f}}$ and the known edge values, the wall shear stress can be calculated by the following

$$
\begin{equation*}
\tau_{w}=\frac{\rho_{e} u_{e}^{2} C_{f}}{2} \tag{4.10}
\end{equation*}
$$

Van Driest also proposed a velocity transformation in order to more easily compare compressible flows. The effective velocity ( $u_{\text {eff }}$ ) is defined as

$$
\begin{equation*}
u_{\text {eff }}=\frac{u_{e}}{\mathrm{~A}^{\prime}}\left\{\sin ^{-1}\left(\frac{2 \mathrm{~A}^{\prime 2}\left(\overline{\mathrm{u}} / \mathrm{u}_{\mathrm{e}}\right)-\mathrm{B}^{\prime}}{\sqrt{4 \mathrm{~A}^{\prime 2}+\mathrm{B}^{\prime 2}}}\right)+\sin ^{-1}\left(\frac{\mathrm{~B}^{\prime}}{\sqrt{4 \mathrm{~A}^{\prime 2}+\mathrm{B}^{\prime 2}}}\right)\right\} \tag{4.11}
\end{equation*}
$$

The effective velocity is non-dimensionalized by the wall-friction velocity ( $u^{*}$ ) given by

$$
\begin{equation*}
\mathrm{u}^{*} \equiv \sqrt{\frac{\tau_{\mathrm{w}}}{\rho_{\mathrm{w}}}} \tag{4.12}
\end{equation*}
$$

to obtain

$$
\begin{equation*}
\mathbf{u}_{\mathrm{eff}}^{+}=\frac{\mathbf{u}_{\mathrm{eff}}}{\mathbf{u}^{*}} \tag{4.13}
\end{equation*}
$$

similarly, the y-coordinate is non-dimensionalized as

$$
\begin{equation*}
\mathrm{y}^{+}=\frac{\mathrm{yu}^{*} \rho_{\mathrm{w}}}{\mu_{\mathrm{w}}} \tag{4.14}
\end{equation*}
$$

A plot of $\mathrm{y}^{+}$versus $\mathrm{u}^{+}$is known as the law of the wall plot (White, 1991) The empirical relationship between $\mathrm{y}^{+}$and $\mathrm{u}^{+}$was first developed for incompressible flow and extended to compressible flow. White gives the empirical relationship as

$$
\begin{equation*}
\mathrm{u}_{\mathrm{eff}}^{+}=\frac{\ln \left(\mathrm{y}^{+}\right)}{\kappa}+\mathrm{C}+\frac{2 \Pi}{\kappa} \sin ^{2}\left(\frac{\pi \mathrm{y}}{2 \delta}\right) \tag{4.15}
\end{equation*}
$$

$\kappa$ and $C$ are constants given by, $\kappa=0.41$ and $C=5.0 . \quad \Pi$ is known as Cole's wake parameter and is defined by (Coles, 1956)

$$
\begin{equation*}
\Pi \approx 0.8(\beta+0.5)^{0.75} \tag{4.16}
\end{equation*}
$$

$\beta$ in Equation (4.16) is called Clauser's equilibrium parameter and is defined as (Clauser, 1954)

$$
\begin{equation*}
\beta=\frac{\delta^{*}}{\tau_{\mathrm{w}}} \frac{\mathrm{dp}}{\mathrm{e}} \mathrm{dx} \tag{4.17}
\end{equation*}
$$

Evans (Evans, 1985) extended Equation (4.15) for flows with zero or adverse pressure gradient by adding the following correction to Cole's wake parameter

$$
\begin{equation*}
\mathrm{u}_{\mathrm{eff}}^{+}=\frac{\ln \left(\mathrm{y}^{+}\right)}{\kappa}+\mathrm{C}+(1+5 \mathrm{TI}) \frac{2 \Pi}{\kappa} \sin ^{2}\left(\frac{\pi \mathrm{y}}{2 \delta}\right) \tag{4.18}
\end{equation*}
$$

were Tl is the freestream turbulence intensity. For the Mach 2.8 tunnel, TI was approximately 0.015 .
Equation (4.18) is valid for all regions of the boundary layer except the laminar sublayer. The laminar sublayer is the region in the boundary layer where the flow becomes laminar. This region exists below $\mathrm{y}^{+}$ $=10$ and in this region $\mathbf{u}^{+}=\mathrm{y}^{+}$(White, 1991).

### 4.2 Hot-Wire Data Reduction

Hot-wire data were collected at each measurement station. Hot-wires were used to collect power spectra data and hot-films were used to collect turbulence data. A series of eight overheats ranging from 1.5 to 2.0 were used for the MOH data and the highest overheat was used for single overheat ( SOH ) data. A nominal overheat of 1.8 was selected for the power spectra analysis. All of the MOH and SOH data reduction equations (except for the power spectra equations) have been incorporated into the FORTRAN code MSHEaR (Bowersox, 1994)

### 4.2.1 Hot-Wire Theory

For supersonic turbulent flow, it can be shown that the Nusselt number of a cylinder (the shape of a hot-wire) has the following non-dimensional functional relationship:

$$
\begin{equation*}
\mathrm{Nu}=\mathrm{Nu}\left(\mathrm{~L} / \mathrm{d}, \mathrm{M}, \mathrm{Pr}, \mathrm{Re}_{\mathrm{e}}, \tau\right) \tag{4.19}
\end{equation*}
$$

$\mathrm{L} / \mathrm{d}$ is the wire aspect ratio, M is the Mach number, Pr is the Prandtl number, $\mathrm{Re}_{\mathrm{e}}$ is the effective Reynolds number based on wire diameter, and $\tau$ in the temperature loading factor. Since single wires were used in this study, $\mathrm{Re}_{\mathrm{e}}$ can be replaced by $\mathrm{Re}_{\mathrm{x}}$. The temperature loading factor can be written as $\tau=\left(T_{w}-T_{e}\right) / T_{t} . T_{w}$ is the wire temperature, and $T_{e}$ is the equilibrium temperature. The equilibrium temperature is defined as the temperature an unheated wire would reach if placed in the flow (Lomas, 1986). Equation (4.19) can be simplified by an application of a set of constraints. These constraints include a large wire aspect ratio ( $\mathrm{L} / \mathrm{d}$ much greater than one), a Mach number greater than 1.2, a constant Pr and a $\mathrm{Re}_{x}$ greater than 20. With these constraints, Equation (4.19) can be reduced to:

$$
\begin{equation*}
\mathrm{Nu}=\mathrm{Nu}\left(\mathrm{Re}_{\mathrm{x}}, \tau\right) \tag{4.20}
\end{equation*}
$$

Bowersox and Schetz (Bowersox and Schetz, 1994) further reduced Equation (4.20) to the following

$$
\begin{equation*}
N u=a \sqrt{\operatorname{Re}_{\mathrm{x}}}+\mathrm{b} \tag{4.21}
\end{equation*}
$$

The variables $a$ and $b$ in Equation (4.21) are calibration factors that depend on the overheat ratio. The Nusselt number can be related to the power supplied to the wire by the following (assuming $T_{e}=T_{t}$ )

$$
\begin{equation*}
\mathrm{Nu}=\frac{\mathrm{V}_{\mathrm{w}}^{2} \mathrm{R}_{\mathrm{w}}}{\pi \mathrm{k}_{\mathrm{t}} \mathrm{~L}\left(\mathrm{~T}_{\mathrm{w}}-\mathrm{T}_{\mathrm{t}}\right)\left(\mathrm{R}_{\mathrm{w}}+\mathrm{R}_{\mathrm{s}}+\mathrm{R}_{\mathrm{L}}\right)^{2}} \tag{4.22}
\end{equation*}
$$

The power laws of viscosity and thermal conductivity were used

$$
\begin{equation*}
k_{1}=k_{o}\left(\frac{T_{t}}{T_{o}}\right)^{n_{k}} \quad \text { and } \quad \mu_{t}=\mu_{o}\left(\frac{T_{t}}{T_{o}}\right)^{n_{\mu}} \tag{4.23}
\end{equation*}
$$

to simplify Equation (4.22). In Equation (4.23), $\mathrm{n}_{\mu}=0.77, \mathrm{n}_{\mathrm{k}}=0.89$ (Bowersox, 1992), $\mathrm{k}_{\mathrm{o}}, \mathrm{T}_{\mathrm{o}}$ and $\mu_{\mathrm{o}}$ are constant reference values (White, 1991). Combining Equations (4.22) and (4.23), the hot-wire response can be found to equal

$$
\begin{equation*}
\frac{V_{w}^{2}}{C_{o}}=\left(\frac{T_{t}}{T_{o}}\right)^{n_{k}}\left(a \sqrt{\operatorname{Reo}_{x}}\left(\frac{T_{o}}{T_{t}}\right)^{\frac{n_{u}}{2}}+b\right)\left(T_{w}-T_{t}\right) \tag{4.24}
\end{equation*}
$$

Reo $_{x}$ is the Reynolds number based on $\mu_{0}$ and $C_{o}$ is a function of the resistance and properties of the wire. It can be expressed as

$$
\begin{equation*}
\mathrm{C}_{\mathrm{o}}=\frac{\left(\mathrm{R}_{\mathrm{w}}+\mathrm{R}_{\mathrm{s}}+\mathrm{R}_{\mathrm{L}}\right)^{2}}{\mathrm{R}_{\mathrm{w}}} \pi \mathrm{Lk}_{\mathrm{o}} \tag{4.25}
\end{equation*}
$$

In order to extract the mean and fluctuation components of Equation (4.24), the voltages, Reynolds number and total temperature are replaced by their mean and fluctuating components. If the binomial theorem is used (keeping only first order terms) and noting that

$$
\begin{equation*}
\frac{{\overline{V_{w}}}^{2}}{C_{o}}=\left(\frac{\bar{T}_{t}}{T_{o}}\right)^{n_{k}}\left(a \sqrt{\overline{R e o}_{x}}\left(\frac{T_{o}}{T_{t}}\right)^{\frac{n_{u}}{2}}+b\right)\left(T_{w}-\overline{T_{t}}\right) \tag{4.26}
\end{equation*}
$$

then solving for $\mathrm{v}_{\mathrm{m}}^{v} / \mathrm{v}_{\mathrm{w}}$, yields the following hot-wire fluctuation equation

$$
\begin{equation*}
\frac{V_{w}^{\prime}}{\overline{V_{w}}}=f\left(\frac{\mathrm{ReO}_{x}^{\prime}}{\overline{R_{0} o_{x}}}\right)+\mathrm{g}\left(\frac{T_{t}^{\prime}}{\bar{T}_{t}}\right) \tag{4.27}
\end{equation*}
$$

where $f$ and $g$ are hot-wire sensitivities and can be written as

$$
\begin{equation*}
\mathrm{f}=0.25\left(1+\frac{\mathrm{b}}{\mathrm{a} \sqrt{\mathrm{Re}_{\mathrm{x}}}}\right)^{-1} \text { and } \mathrm{g}=\frac{\overline{\mathrm{T}_{\mathrm{t}}}}{2\left(\mathrm{~T}_{\mathrm{w}}-\overline{\mathrm{T}_{\mathrm{t}}}\right)}+\frac{\mathrm{n}_{\mathrm{k}}}{2}-\mathrm{fn}_{\mu} \tag{4.28}
\end{equation*}
$$

Equation (4.26) can be rewritten as

$$
\begin{equation*}
\sqrt{\sqrt{\mathrm{Re}_{\mathrm{x}}}}+x_{i} \overline{T_{t}} \sqrt{\overline{\mathrm{Re}_{\mathrm{x}}}}+y_{i} \overline{\mathrm{~T}}_{\mathrm{t}}=z_{i} \tag{4.29}
\end{equation*}
$$

In Equation (4.29), $x_{i}=-1 / T_{w i}, y_{i}=-b_{i} /\left(a_{i} T_{w i}\right)$ and $z_{i}=\bar{V}_{w i}^{2} /\left(C_{i} a_{i} T_{w i}\right)-b_{i} / a_{i}$. The subscript $i$ in Equation (4.29) is an index representing the number of overheat ratios. Since Equation (4.29) has two unknowns ( $\sqrt{\overline{\mathrm{Re}_{\mathrm{x}}}}$ and $\overline{\mathrm{T}_{\mathrm{t}}}$ ), a minimum of two overheat ratios are necessary to solve this equation. If more overheats are available, then a least squares analysis can be used to solve Equation (4.29). The least squares analysis solves the following set of equations by the secant iteration method (Bowersox, 1992)

$$
\begin{gather*}
N \sqrt{\overline{\operatorname{Re}_{x}}}+\overline{T_{t}}\left(\sum_{N} y_{i}-\sum_{N} x_{i} z_{i}\right)+2 \overline{T_{t}} \sqrt{\overline{R e x_{x}}} \sum_{N} x_{i}+  \tag{4.30}\\
\bar{T}_{t}^{2} \sum_{N} x_{i} y_{i}+\bar{T}_{t}^{2} \sqrt{\overline{R e}_{x}} \sum_{N} x_{i}^{2}=\sum_{N} z_{i} \\
\sqrt{\overline{R e}_{x}}\left(\sum_{N} y_{i}-\sum_{N} x_{i} z_{i}\right)+\bar{T}_{t} \sum_{N} y_{i}^{2}+2 \overline{T_{t}} \sqrt{\operatorname{Re}_{x}} \sum_{N} x_{i} y_{i}+  \tag{4.31}\\
\frac{T_{t} \operatorname{Re}_{x}}{} \sum_{N} x_{i}^{2}=\sum_{N} y_{i} z_{i}
\end{gather*}
$$

N is the number of overheat ratios. To obtain the hot-wire fluctuation equation, Equation (4.27) is squared and averaged. The result of this process is

Unlike Equation (4.29), Equation (4.15) required three overheats to solve. If more overheat ratios are available, then a General Least Squares (GLS) method can be used to solve Equation (4.32). Application of the GLS method yields the following set of equations.

Solution of Equation (4.33) completes the MOH analysis

### 4.2.2 Separation of Turbulence Variables

Multiple overheat single hot-wire anemometry in supersonic flow can be used to determine the following conservative turbulence quantities.

$$
\begin{equation*}
\overline{(\rho u)^{\prime 2}} \quad \overline{\mathrm{~T}_{\mathrm{i}}^{\prime 2}} \quad \overline{(\rho u)^{\prime} \mathrm{T}^{\prime}} \tag{4.34}
\end{equation*}
$$

Two assumptions must be made in order to separate the variables in Equation (4.34) into non-conservative
variables. The first assumption is that the fluid obeys the thermally perfect gas assumption. The second assumption is controversial, it is that the pressure fluctuations are small compared to the density and temperature fluctuations. This hypothesis has been demonstrated by Kistler and Bowersox (Kistler, 1959; Bowersox \& Schetz 1994) Kistler postulated that $\mathrm{p}^{\prime}$ is proportional to $\mathrm{u}^{\prime 2}, \mathrm{u}^{12}$ is second order and is usually negligible compared to first order terms. Bowersox experimentally verified the $\mathrm{p}^{\prime}=0$ assumption in a Mach 4.0 free mixing layer. With these two assumptions, the separated results can be calculated from the following set of equations.

$$
\begin{gather*}
\frac{\mathbf{u}^{\prime}}{\overline{\mathrm{u}}}=\frac{(\rho \mathrm{u})^{\prime}}{\overline{\rho \mathrm{u}}}-\frac{\rho^{\prime}}{\rho}  \tag{4.35}\\
\frac{\rho^{\prime}}{\rho}=\frac{1}{\alpha+\beta}\left(\beta \frac{(\rho \mathrm{u})^{\prime}}{\overline{\rho u}}-\frac{\mathrm{T}_{\mathrm{t}}^{\prime}}{\overline{\mathrm{T}_{\mathrm{t}}}}\right) \tag{4.36}
\end{gather*}
$$

where $\alpha=\left(1+0.5(\gamma-1) \mathrm{M}^{2}\right)-1$ and $\beta=(\gamma-1) \alpha \mathrm{M}^{2}$.

### 4.2.3 Single Overheat Theory

As the hot-wire's overheat is increased, the hot-wire becomes less sensitive to total temperature fluctuations. Thus, the SOH method can be used to obtain accurate results for flow where $T_{t}^{\prime}$ is small. If $\mathrm{T}_{\mathrm{t}}^{\prime}$ is assumed to be zero, Morkovin demonstrated that for $\operatorname{Pr}=1$ (Morkovin, 1961)

$$
\begin{equation*}
\frac{\mathrm{T}^{\prime}}{\overline{\mathrm{T}}}=-(\gamma-1) \mathrm{M}^{2}\left(\frac{\mathbf{u}^{\prime}}{\overline{\mathbf{u}}}\right) \tag{4.37}
\end{equation*}
$$

It can be shown that

$$
\begin{equation*}
\frac{\mathrm{u}^{\prime}}{\overline{\mathrm{u}}}=\frac{1}{1-(\gamma-1) \mathrm{M}^{2}}\left(\frac{(\mathrm{\rho u})^{\prime}}{\overline{\rho \mathrm{u}}}\right) \tag{4.38}
\end{equation*}
$$

and

$$
\begin{equation*}
\frac{T_{t}^{\prime}}{\bar{T}_{t}}=0 \tag{4.39}
\end{equation*}
$$

### 4.2.4 Mach 2.8 Hot-Wire Probe Data

Eight overheat ratios were used to examine the flow for mean flow and turbulence results. They ranged from 1.52 to 2.03 . In addition, the highest overheat was extracted for a SOH analysis (see Section 5.3). Power spectra data was obtained for a single overheat ratio only. The data in Table 4-1 contain the wire information and overheat ratios used in this experiment. Note that only one hot-wire and one hotfilm was used to collect data for this experiment.

Table 4-1: Mach 2.8 probe and overheat ratio data

| Probe | Purpose | OHR |
| :---: | :--- | :--- |
| $1218-20 \mathrm{~S} / \mathrm{N} 75014$ | ZPG, FPG, CPG mean flow and turbulence | $2.03,1.98,1.87,1.81,1.74,1.66,1.58,1.52$ |
| $1218-\mathrm{T} 1.5 \mathrm{~S} / \mathrm{N} 43456$ | FPG, CPG power spectra | 1.85 |
| $1218-\mathrm{T1.5} \mathrm{~S} / \mathrm{N} 43456$ | ZPG power spectra | 1.88 |

### 4.2.5 Mach 1.7 Hot-Wire Probe Data

Table 4-2 presents the probe information for the probes used in this portion of the study. Note the same probe was used for the Mach 1.7 mean flow as the Mach 2.9 mean flow. A hot-film was used to evaluate the power spectra. The film had a frequency response of 140 KHz that is somewhat less than what was used for the Mach 2.9 power spectra data.

Table 4-2: Mach 1.7 probe and overheat ratio data

| Probe | Purpose | OHR |
| :--- | :--- | :--- |
| 1218-20 S/N 75014 | ZPG mean flow and turbulence | $1.98,1.89,1.82,1.73,1.66$ |
| 1214-10 S/N K746 | ZPG power spectra | 1.77 |
| 1214-20 S/N 9446 | FPG power spectra | 2.01 |

### 4.2.6 Power Spectra Data Reduction

SOH hot-wire anemometry was used to evaluate the energy spectra in the boundary layer. The hot-wire response was sampled at 500 KHz for 0.52 seconds. The hot-wire output was divided into 511 blocks of 1,024 data points for ease of data reduction. A discrete Fourier transform (DFT) (Reid, 1983) was performed on each block to convert the hot-wire results from the time domain to the frequency domain. The classic DFT is modeled by the following equation

$$
\begin{equation*}
\mathrm{V}(f)=\frac{1}{\mathrm{~N}} \sum_{\mathrm{N}} \mathrm{~V}(\mathrm{t}) \mathrm{e}^{-\frac{2 \text { nift }}{\mathrm{N}}} \tag{4.40}
\end{equation*}
$$

where N is the sample size of 1,024 points. The output of the $\operatorname{DFT}$ was $((\rho \mathrm{pu}(f)) / \overline{\rho u})^{2}$ versus frequency. To obtain the final result, the average of the 511 blocks was calculated.

### 4.3 PIV Data Reduction

A CCD camera was used to acquire and digitize the particles in the flow. A raw digitized image looked like red and green dots against a black background. The red and green dots were seed particles illuminated by the red or green laser. Due to the quality of the digitized image, software image enhancement was not required. With this image, the velocity field was calculated using a crosscorrelation technique. The cross-correlation technique performs a Fourier transform on the intensity distribution of the red and green images over a small patch of the digitized image. In this experiment, a patch size of $128 \times 128$ pixels (approximately $2 \mathrm{~mm} \times 2 \mathrm{~mm}$ ) was used. The patches were overlapped by 64 pixels to more accurately map the flowfield. An intensity-weighted peak-searching routine developed by Innovative Scientific Solutions Inc., determined the locations of the red and green peaks to sub-pixel accuracy (Glawe et al, 1996). With the red and green peaks located, knowledge of the time delay between the red and green laser and knowledge of which laser fired first, the velocity vector could be drawn for that patch connecting the red and green peaks.

Filtering was performed on the calculated vector field to eliminate unnecessary and incorrect vectors by one of two methods. The software developed by Innovative Scientific Solutions Inc., was used to compare each vector's length and magnitude to its neighbors. If the difference in length or magnitude of any particular vector compared to its neighbor exceeded a specified input, the vector was deleted. Typically, this filtering process was repeated three times. In the event that this filter sequence did not remove a bad vector, a manual deletion of the vectors was performed. Manual deletion was required approximately every three images.

Calibration of the images was necessary due to large amounts of scatter in the laser time base. The average vector length in the freestream was calculated for each image by averaging all the vectors in the lower left hand corner of the image. This location represents undisturbed freestream flow for each pressure gradient model. The freestream velocity was divided by the average vector length to obtain each image's calibration factor. Each image's vector field was multiplied by the calibration factor. Thus, the velocity of each vector in the vector field can be found by the following

$$
\begin{equation*}
\mathrm{u}=\Delta \mathrm{L}_{\mathrm{G}-\mathrm{R}} * \frac{\mathrm{U}_{\mathrm{e}}}{\overline{\mathrm{~L}}} \tag{4.41}
\end{equation*}
$$

$\Delta \mathrm{L}_{\mathrm{G}-\mathrm{R}}$ is one vector length and $\overline{\mathrm{L}}$ is the average vector length.
This process effectively removed the scatter and adjusted the vector field such that the freestream vector's length matched the freestream velocity. After filtering and calibration, the calculated vector fields over the entire image were averaged together to obtain a composite vector field.

## 5. Mach 2.8 Results and Discussion

While measurements were obtained at three Mach numbers during the course of this study, the main test condition was Mach 2.8. This chapter presents the results from the Mach 2.8 phase of this experiment. Flow visualization is discussed first, followed by a discussion of the mean flow and turbulence results. Recall that the objectives of this study are to, quantify the effects of pressure gradients on the flow validate PIV as a flow diagnostic tool and compare turbulence data obtained via hot-wire anemometry to turbulence data obtained by LDV.

### 5.1 Flow Visualization

In order to better understand the flow, it is useful to have some type of flow visualization. Traditional visualization techniques include shadowgraphs and schlieren photographs; however, the Mie scattering images of PIV can also be used for flow visualization. Flow visualization images comparing the CPG to the ZPG and the FPG to the ZPG are presented as Figures 5-1 and 5-2, respectively. Two images were taken for each pressure gradient to illustrate data repeatability.


Figure 5-1: ZPG and CPG flow visualization images

Note in both figures, the ZPG boundary layer (along the tunnel floor) is uniform and the structures have an inclination of between $45^{\circ}-60^{\circ}$. This effect has been noted by Spina (Spina, 1988). The ZPG flow structures span $\sim 1 / 2$ the boundary layer. Application of a CPG causes a decreases in the boundary layer thickness. Further, the CPG structures are inclined approximately $50^{\circ}-70^{\circ}$. Finally, the structures are smaller and span the entire boundary layer.


Figure 5-2: ZPG and FPG flow visualization images
Figure 5-2 shows the comparison between the FPG and ZPG surfaces. As compared to the ZPG, the FPG boundary layer grows in the streamwise direction, as expected. More small scale structures are evident on the boundary layer edge; however, no large scale structures are seen. This is consistent with the theory that a FPG promotes the disassociation of large scale structures into small scale structures (Luker et al., 1997). Thus, the FPG causes the boundary layer to grow while contributing to the conversion large scale structures to smaller scale structures.

Estimates of the boundary layer thickness can be ascertained from Figures 5-1 and 5-2. The estimated boundary layer thickness is 1.0 cm for the ZPG, 0.8 cm for the CPG and 1.1 cm for the FPG.

### 5.2 Mean Flow Results

Mean flow properties were required to map out the flow and make sure that the flow matched previous studies (Miller, 1994; Dotter, 1994; Luker, 1995 \& Hale, 1995). Mean flow properties were evaluated by three different methods. First of all, conventional pressure probes (Pitot and cone-static) were used to compute the Mach number, static pressure (p) and velocity (u). Secondly, the hot-wire was used to calculate the mass flux and total temperature ratio. Finally, PIV was used to provide velocity information. Data were normalized by their edge values (or $\mathrm{P}_{\mathrm{t} 1}$ for pressure data) and plotted versus $\mathrm{y} / \delta$. Unless otherwise stated, $\delta$ is based on $\mathfrak{u}=0.99 \mathrm{U}_{\mathrm{e}}$. Boundary layer thickness parameters are summarized in Table 5.1.

### 5.2.1 Conventional Pressure Probe

Pitot and cone-static measurements were taken along the tunnel centerline at positions 71.0 cm and 71.5 cm downstream of the nozzle for the ZPG and FPG, respectively. Pitot and cone-static data were not taken for the CPG. This data is available in Hale and Dotter (Hale, 1995 \& Dotter, 1994). Figures 53 and 5-4 present the normalized Pitot and cone-static data.

The FPG flow is not as fully developed as the ZPG flow. This is demonstrated by the increasing trends in the Pitot and Cone-static pressure profiles. The data in Figures 5-3 and 5-4 in conjunction with Equation (4.1) and the supersonic Rayleigh Pitot formula was used to calculate the Mach number. The freestream Mach number was found to be 2.79 and 3.00 for the ZPG and FPG, respectively. The Mach number profile is shown in Figure 5-5.


Figure 5-3: Pitot pressure profile


Figure 5-4: Cone-static pressure profile


Figure 5-5: Mach number profile
These data in conjunction with the definitions of the boundary layer thickness parameters
(Section 4.1.2) were used to compute $\delta_{\mathrm{u}}$ and $\delta_{\mathrm{M}}$. Table 5.1 presents the results of these computations.
Data from Luker and Hale (Luker, 1995 and Hale, 1995) are also presented to form a basis of comparison.
The differences between the boundary layer parameters are within the $25 \%$ uncertainty expected by Luker (Luker, 1995).

Table 5-1: Boundary layer parameters

|  | ZPG |  | FPG |  | CPG |  |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: |
|  | Luker | Present | Luker | Present | Hale $(\mathrm{x}=68 \mathrm{~cm})$ | Hale $(\mathrm{x}=70 \mathrm{~cm})$ |
| $\delta_{\mathrm{u}}(\mathrm{mm})$ | 9.85 | 10.1 | 11.9 | 10.5 | 9.09 | 7.48 |
| $\delta_{\mathrm{M}}(\mathrm{mm})$ | 10.8 | 10.6 | 13.1 | 11.6 | 10.04 | 8.44 |
| $\mathrm{M}_{\mathrm{e}}$ | 2.79 | 2.82 | 2.91 | 3.00 | 2.72 | 2.51 |
| $\mathrm{U}_{\mathrm{e}}(\mathrm{m} / \mathrm{sec})$ | 602 | 604 | 606 | 616 | 596 | 575 |

### 5.2.2 Hot-Film Mean Profiles

Single wire data was acquired at the same locations as the mean flow data for the ZPG and FPG. For the CPG, data was collected the tunnel centerline at 68.0 cm and 70.0 cm downstream of the throat. For mean flow data, the hot-film was traversed towards the wall until the probe struck the wall. Then the probe was traversed back to its starting point. This allowed for accurate near wall data collection and for the wall's position to be accurately located. Due to the thickness of the boundary layer and limitations on run time, data collection was only possible during the initial probe motion (traverse up) for the FPG. For the ZPG and CPG, data was collected on the initial and final probe motion (traverse up \& traverse down). A single hot-film was used for all mean flow measurements.

The mean flow data included the local Reynolds number and total temperature. The Reynolds number was converted to mass flux and non-dimensionalized by the freestream value. The total temperature was non-dimensionalized by $T_{t 1}$. Figures 5-6 to 5-9 present the mass flux information for all pressure gradients. In addition, these figures present the total temperature profiles for each pressure gradient model during the traverse up motion.

The ZPG data is presented in Figure 5-6. There is excellent agreement between data collected during the traverse up and the traverse down movement. The single OHR data matches well with the multiple OHR data implying that the total temperature variations were small. This is confirmed by the total temperature profile which was reasonably constant across the boundary layer. The freestream mass flux was found to be approximately $400 \mathrm{~kg} / \mathrm{m}^{2}-\mathrm{sec}$.

The FPG data is presented in Figure 5-7. Recall that data was only obtained during the traverse up movement. Again, the single OHR data matches the multiple OHR data well implying that the FPG has a stabilizing effect on total temperature ratio. This is confirmed by the accompanying total temperature profile. The freestream mass flux was found to be approximately $320 \mathrm{~kg} / \mathrm{m}^{2}-\mathrm{sec}$.


Figure 5-6: ZPG hot-film data


Figure 5-7: FPG hot-film data

Figures 5-8 and 5-9 present the CPG data at two different distances from the nozzle. Figure 5-8 shows data 68 cm downstream of the nozzle and Figure $5-9$ shows data 70 cm downstream of the nozzle. The data in Figure 5-8 demonstrates the destabilizing effect of the CPG. The data scatter is slightly larger between the different data collection methods implying that the total temperature fluctuations were higher for the CPG than for the ZPG or FPG. Again, this is confirmed by the relatively large total temperature variations ( $\sim 7 \%$ ) noted in the total temperature profiles. The largest value of mass flux does not occur in the freestream as in the ZPG or FPG, but the largest value occurs near $\mathrm{y} / \delta=1.0$. This is the result of the compression wave formed by the CPG model. The freestream value of the mass flux is approximately 560 $\mathrm{kg} / \mathrm{m}^{2}$-sec.

The data in Figure 5-9 has similarities to the FPG as well as the upstream CPG data. The shape of the curve does not have the elbow present in Figure 5-8. As with the ZPG and FPG, the highest value of mass flux $\left(420 \mathrm{~kg} / \mathrm{m}^{2}-\mathrm{sec}\right)$ occurs in the freestream. This indicates that this measurement location is downstream of the compression wave. Like in Figure 5-8, the scatter in the data is slightly larger throughout the boundary layer.


Figure 5-8: CPG ( $\mathrm{x}=68 \mathrm{~cm}$ ) hot-film data


Figure 5-9: $\mathrm{CPG}(\mathrm{x}=70 \mathrm{~cm})$ hot-film data

### 5.2.3 PIV Mean Profiles

A main advantage of PIV is that it can resolve data along an entire plane. Recall that one of the objectives of this study was to further evaluate PIV as a diagnostic tool for supersonic flow. The PIV portion of this analysis required a great deal of effort and time. PIV has only been used once at AFIT (Glawe et al., 1996) and rarely for supersonic flow. Thus, the procedures (seeding, camera use) were not well established. Several weeks were expended in trying various seeding techniques and camera types before acceptable results were obtained.

Recall that PIV takes images of the flow. The images were averaged together to obtain mean and turbulent statistics. The number of images used to create the combined image are listed in Table 5-2. For the present experimental set-up 1 to 3 images per run were attainable and each image required about 6 Mbytes of disk space (stored in a compressed TIFF file format).

Table 5-2: Number of PIV images used to create the combined image

| Pressure gradient | Number of images |
| :--- | :--- |
| ZPG | 25 |
| CPG | 30 |
| FPG | 25 or 93 |

The two different data set sizes for the FPG case allowed study of the effect of image sample size on flow field convergence. The results from the averaging are presented as Figures 5-10 to 5-12 for the ZPG, FPG and CPG, respectively. Locations where there was insufficient data to properly compute velocity or turbulence intensity are indicated by diamond shapes with locally high velocity gradients. The boundary layer thickness in Figures 5-10 to $5-12$ is based on the location labeled 'LDV and hot-film measurement location(s)'. The horizontal x axis does not corrleate to any exact wind tunnel station. All profiles, unless otherwise stated, are extracted from the location labeled 'LDV and hot-film measurement location(s)'.

The ZPG mean flow results are plotted in Figure 5-10. Note that PIV provides an entire plane of information. The velocity gradient shown in Figure 5-10 is uniform and relatively smooth throughout the boundary layer even with the relatively small number of images taken of this flowfield. The $u=605$ $\mathrm{m} / \sec \left(\mathrm{u}=0.98 \mathrm{U}_{\mathrm{e}}\right)$ contour does not grow across the image. Thus, the boundary layer is steady and does not grow significantly in the 2 cm PIV region. A $128 \times 128$ correlation block was used to compute the averages; however, the first data point was roughly 69 pixels from the wall (i.e., $\mathrm{y} / \delta=0.03$ ). The freestream velocity was calibrated to $600 \mathrm{~m} / \mathrm{sec}$.

The FPG mean flow results are plotted in Figure 5-11. This figure shows the effect of the favorable pressure gradient on the velocity profile. Both the boundary layer height and freestream velocity increased from their ZPG values. The boundary layer slightly grows from the $\mathrm{x}=1.0 \mathrm{~cm}$ location to the x $=3.0 \mathrm{~cm}$ location, illustrating the effect of the FPG on the boundary layer structure. The acceleration due to the expansion is evident by a large area of velocity greater than $600 \mathrm{~m} / \mathrm{sec}$ downstream of the LDV measurement location. Finally, the velocity gradient is evident in this figure.

The CPG results are plotted in Figure $5-12$. This figure shows the effect of the combined pressure gradient on the velocity profile. The boundary layer height decreased from its ZPG and FPG value; however, the freestream velocity remained similar to the ZPG freestream velocity. Velocity gradients are clearly indicated in Figure 5-12. The boundary layer growth indicated in the aft hot-wire measurement location should be treated with caution as $\delta$ is based off the first measurement location. In addition, the expected compression wave is not apparent in this figure.

PIV derived velocity vectors are presented as Figures $5-13$ to $5-15$ for the ZPG, FPG and CPG, respectively. The vectors illustrate the grid size discussed in Chapter 4. Note the turning of the velocity vectors due to the pressure gradient in Figures $5-14$ and $5-15$. The wall boundary was sketched in after the vectors were drawn and do not exactly match the wall contour.


Figure 5-10: ZPG velocity contour plot


Figure 5-11: FPG velocity contour plot


Figure 5-12: CPG velocity contour plot


Figure 5-13: ZPG velocity vectors


Figure 5-14: FPG velocity vectors


Figure 5-15: CPG velocity vectors

### 5.2.4 Comparison Between Data Collection Methods

This section compares the mean velocity for all of the data collection methods. In addition, the data are compared to the Van Driest profiles (Section 4.1.3).

### 5.2.4.1 Mean Flow Comparison

Since hot-wire anemometry and LDV only acquire data at one point and PIV provides data along an entire plane, a line (or profile) of data was extracted from the PIV images. This allows a better method to compare PIV results with the LDV results of Luker et al. (Luker et al., 1997). The PIV profile was extracted at the locations indicated in Figures 5-10 to 5-12, normal to the wall. The velocity comparisons between LDV, PIV and conventional pressure probes are presented as Figures 5-16 to 5-18 for the ZPG, FPG and CPG, respectively.

Figure 5-16 presents the ZPG velocity comparisons. There is excellent agreement between the LDV, PIV and conventional pressure velocity data above $\mathrm{y} / \delta=0.2$. Below this location in the boundary layer, PIV over-predicts the velocity ratio. This indicates that 25 images provides satisfactory mean flow data for the ZPG configuration.

Figure 5-17 presents the FPG data. Recall that the PIV FPG data was reduced with two different number of images to gauge the effect of number of images on mean flow quantities. The data for both image sample sizes predicts a fuller velocity profile than the LDV and pressure data; however the overall agreement is excellent. This indicates that velocity convergence has been reached with 25 images.

Figure 5-18 presents the CPG data. At both locations there is excellent agreement between the LDV and PIV data. This indicates that 30 images provides satisfactory mean flow data for the CPG configuration. Note that PIV resolved mean flow data closer to the wall.

From Figures 5-16 to 5-18, it was demonstrated that PIV can provide accurate mean flow data for as few as 25 images. Increasing the number of images does not assist in velocity convergence. This is an important finding since testing was time consuming and costly.


Figure 5-16: ZPG velocity comparison


Figure 5-17: FPG velocity comparison


Figure 5-18: CPG velocity comparison

### 5.2.4.2 Van Driest Correlation

Another method used to validate PIV was to use the scaling laws proposed by Van Driest (Van Driest, 1951) to present the data. Evan's correction factor (Evans, 1985) was set equal to $1.5 \%$ for all pressure gradient cases. In addition, the hot-film data from this study was also scaled and plotted along with the LDV data by Luker et al. (Luker et al., 1997). This allows a comparison between the PIV data, LDV data and hot-film data. The profile for the ZPG is plotted in Figure 5-19 along with the theoretical flat plate values. For the flat plate, Clauser's parameter, $\beta$, equals 0.0 . This value leads to a Coles wake parameter, $\Pi$, equal to 0.476 . There is excellent agreement among the three data types and with the theory. The PIV lies closer to the theoretical values than either the LDV or hot-film data demonstrating the accuracy of the PIV method. In addition, PIV allows for data to be collected nearer to the wall ( $\mathrm{y}^{+} \approx 31$ ) than LDV or the hot-film data. This is another illustration of the power of PIV. In addition to the ability to visualize a plane of data, near wall data can be collected even with the relatively large grid size ( $128 \times 128$ pixels) used in this study.


Figure 5-19: ZPG Van Driest velocity profile
For the FPG, $\beta$ was found by Luker (Luker, 1995) to be -1.34 , leading to an undefined Coles wake parameter. Thus, the FPG data was compared to the theoretical flat plate values. Figure 5-20 presents the FPG comparisons.

As with the ZPG, there is excellent agreement with the PIV and hot-film data to the theory. The PIV agreement is better than the LDV data found by Luker. However, in this case, LDV was able to resolve data closer to the wall. This is due to the effects of the wall slope and the $128 \times 128$ pixel patch size in the PIV data reduction.

For the CPG, $\beta$ was found by Hale (Hale, 1995) to be 1.12 at $\mathrm{x}=68 \mathrm{~cm}$, leading to a Coles wake parameter equal of 1.15 . Figure $5-21$ presents the CPG comparisons for $\mathrm{x}=68 \mathrm{~cm}$. Again, note the excellent agreement between the PIV data and the theory. As with the FPG, PIV predicts the scaled velocity with more accuracy than the LDV predictions. The hot-film data also agrees well with the theory.

At $x=71 \mathrm{~cm}, \beta$ was found to be -0.94 , leading to an undefined Coles wake parameter. Thus, the PIV data at $x=70 \mathrm{~cm}$ was compared to the theoretical flat plate values. The agreement between the


Figure 5-20: FPG Van Driest velocity profile


Figure 5-21: $\mathrm{CPG}(\mathrm{x}=68 \mathrm{~cm})$ Van Driest velocity profile
data types is good; however, the agreement between the data and the theory is not as good for this flow field condition. This is most likely due to inaccuracies in predicting the theory or due to the thinness of the boundary layer. Recall from Section 5.1 that the CPG reduced the boundary layer thickness by approximately $20 \%$. Figure $5-22$ presents the CPG comparisons for $\mathrm{x}=70 \mathrm{~cm}$.


Figure 5-22: CPG ( $x=70 \mathrm{~cm}$ ) Van Driest velocity profile
As demonstrated in the previous four figures, use of Van Driest's scaling factors collapses the PIV data with the theory extremely well. This is a further indication of the accuracy and applicability of PIV for mean flow data reduction.

### 5.3 Turbulence Results

Hot-films and PIV were used to evaluate the fluctuating components of fluid properties. Hotfilms were used to evaluate the mass flux, total temperature fluctuations and power spectra. Power spectra data was obtained by sampling the boundary layer at discrete locations in the boundary layer. A single hot-wire was used for all power spectra data. Finally, PIV was used to evaluate velocity fluctuations. This will assist in the validation of PIV as a flow diagnostic tool.

### 5.3.1 Hot-Film Turbulence Results

This section contains the hot-film turbulence results. Included in this section are the mass flux turbulence intensity profiles and the energy spectra results.

### 5.3.1.1 Turbulence Intensity Profiles

The fluctuating mass flux profiles are given in Figures 5-23 to 5-26 for the ZPG, FPG, and CPG (both locations), respectively.

Figure 5-23 presents the ZPG mass flux turbulence intensity (TI). There is excellent agreement between the multiple and single overheat methods as well as the two different traversing methods. The discrete points from the spectra data also show excellent agreement with the traverse data. The data indicates that the probe just barely reached the freestream. The maximum turbulence intensity occurs between $\mathrm{y} / \delta=0.3$ and $\mathrm{y} / \delta=0.5$ with a value of between $15 \%$ to $16 \%$. Although this value is $5 \%$ higher than that found by Dotter (Dotter, 1994), it falls within the expected uncertainty bounds (see Appendix B).


Figure 5-23: ZPG mass flux turbulence intensity

Figure 5-24 presents the FPG mass flux turbulence intensity. There is excellent agreement between the single overheat data, multiple overheat data and the discrete point data. The maximum turbulence intensity is $12 \%$ at a $\mathrm{y} / \delta$ of 0.5 . Thus, the FPG reduced the maximum turbulence intensity by $25 \%$ as compared to the ZPG. Note that traversing the probe into the wall provides near wall data.


Figure 5-24: FPG mass flux turbulence intensity
Figure 5-25 presents the CPG mass flux turbulence intensity at $\mathrm{x}=68 \mathrm{~cm}$. Again, there is excellent agreement between the discrete data points, the traversed multiple overheat and single overheat data. The maximum turbulence intensity is $20 \%$, a $25 \%$ increase over the ZPG value, at $\mathrm{y} / \delta=0.3$.

Figure 5-26 presents the CPG mass flux turbulence intensity at $x=70 \mathrm{~cm}$. As with the ZPG and FPG, there is excellent agreement between the single and multiple overheat data. Discrete data points were not taken at this location. The maximum turbulence intensity is $22 \%$. This level of turbulence intensity is higher than observed at the $\mathrm{x}=68 \mathrm{~cm}$ location. This data contradicts the data by Hale (Hale, 1995) who found that the turbulence intensity was reduced at this location. This discrepancy could be explained by differing data collection and reduction methods. In addition, uncertainty in x location may play a part. Hale acquired data at $\mathrm{x}=71 \mathrm{~cm}$. Dotter (Dotter, 1994) collected data with a cross-wire at $\mathrm{x}=$

71 cm . The magnitude of the peak turbulence intensity was roughly $20 \%$ which agrees well with the present study. As with the $x=68 \mathrm{~cm}$ data, the turbulence intensity increased as compared to the ZPG data in Figure 5-26.

The results in Figures 5-23 and 5-26 lead to the following conclusions. For the ZPG, the maximum mass flux turbulence intensity is between $15 \%$ to $16 \%$. Application of a FPG damps out the turbulence as evidenced by the $25 \%$ reduction in the mass flux axial turbulence intensity. Application of the CPG has the opposite effect. The mass flux turbulence intensity increases by $22 \%$ as compared to the ZPG.


Figure 5-25: CPG $(x=68 \mathrm{~cm})$ mass flux turbulence intensity


Figure 5-26: $\operatorname{CPG}(x=70 \mathrm{~cm})$ mass flux turbulence intensity

### 5.3.1.2 Energy Spectra

Hot-wires were used to interrogate the boundary layer to obtain energy spectra data. This will determine the effect of the pressure gradient on the power spectra. Figure 5-27 presents the power spectra data for the ZPG. The data in Figure $5-27$ indicates that, for a ZPG, there is little variation in energy across the frequency spectrum. The roll-off frequency is approximately $40,000 \mathrm{~Hz}$.

Figure $5-28$ presents the power spectra for the FPG. The FPG data contains several key differences. First, the energy level is not constant across the frequency spectrum. The energy is more heavily weighted towards the lower frequency structures. This indicates that the higher frequency structures are feeding their energy in a FPG flow field to the low frequency structures their energy in a FPG flow field. This confirms the theory that application of a FPG causes disassociation of large scale structures into small scale structures. In addition, the roll-off frequency is lower, approximately 10,000 Hz. Finally, a cross-over between the inner region data and outer ration data occurs. This indicates that the large frequency structures contain less energy in the inner region, than they do in the outer region.


Figure 5-27: ZPG power spectra data


Figure 5-28: FPG power spectra data

Figure 5-29 presents the power spectra for the CPG. Data was taken at $x=68 \mathrm{~cm}$ only. The CPG power spectra data contains several key differences as compared to the ZPG flow. As with the FPG, the energy level is not constant across the frequency spectrum. The energy is more heavily weighted towards the lower frequency structures and the energy level is higher for the CPG than the FPG or ZPG . The roll-off frequency is approximately unchanged from the FPG conditions. Also the cross-over noted in the FPG data is repeated in Figure 5-29. These data confirm that the large scale structures contain most of the energy. This finding is consistent with the flow visualization results.


Figure 5-29: CPG power spectra data
The following statements sum up the effect of the pressure gradient on the power spectra. First, for the ZPG , the energy is evenly distributed among the frequency spectrum. Application of a pressure gradient shifts energy from the high frequency to the low frequency structures. The CPG results in a higher low frequency energy than the FPG. A cross-over in energy occurs between the inner region and outer regions of the boundary layer.

### 5.3.2 PIV Turbulence Results

PIV was used to compute the axial turbulence intensity (TI). Like LDV, PIV directly measures the velocity fluctuations, whereas hot-films measure mass flux TI. Figures 5-30 to $5-32$ present PIV velocity TI contours for the ZPG, FPG and CPG, respectively. Note that the power of PIV allows an entire plane of data to be resolved.

In all figures, the TI in the freestream is between $3 \%$ and $4 \%$, which is higher than that reported by Luker and Hale (Luker, 1995 \& Hale, 1995). This suggests that PIV either is not sensitive to TI (low signal to noise ratio) or there were an insufficient number of samples taken. However, the boundary layer data is more reasonable as will be shown.

The ZPG contour plot of TI is presented in Figure 5-30. The maximum TI is near the wall with a magnitude of $15 \%$. This shows close agreement with the mass flux TI obtained via hot-film techniques (Figure 5-26). Regions of TI in excess of $20 \%$ in the freestream or near the wall are due to either an insufficient number of samples or lack of seed. The TI gradient is resolved well throughout the boundary layer. As previously mentioned, the freestream TI is approximately $4 \%$, twice as high as compared to Figure 5-26.

Figure 5-31 presents the FPG TI contour plot. The maximum TI in the boundary layer is approximately $13 \%$, a $15 \%$ reduction from the ZPG results. Thus, the stabilizing effect of the FPG is observed in the PIV measurements. The PIV results presented in Figure 5-31 also indicate that the boundary layer is growing in the flow direction, again as previously demonstrated via hot-film techniques. The freestream TI is higher for the FPG, which is unexpected. Recall that in Figure 5-27 the freestream TI is approximately $1.5 \%$.

Figure 5-32 shows the effect of the CPG on the TI. The highest TI is $26 \%$, a $50 \%$ increase over the ZPG results. Again, PIV confirms the TI enhancement caused by the CPG. The freestream TI is approximately $5 \%$, which is a $66 \%$ increase over the ZPG and confirms the destabilizing effect of the CPG on the flow.


Figure 5-30: ZPG u turbulence intensity contours (\%)


Figure 5-31: FPG u turbulence intensity contours (\%)


Figure 5-32: CPG u turbulence intensity contours (\%)

### 5.3.3 Comparison Between Data Collection Methods - Turbulence Quantities

Data were extracted from the PIV contour plots and compared to the hot-film data and relevant data from the open literature (Luker, 1995, Hale, 1995 \& Klebanoff, 1954). The data was extracted from the PIV in the same manner as described in Section 5.2.4. The line plot comparisons are presented as Figures 5-33 to 5-35 for the ZPG, FPG and CPG, respectively. In addition, the single overheat data was reduced as in Section 4.2.3 as another basis of comparison.

Figure 5-33 presents the ZPG turbulence intensity contours. In the boundary layer ( $\mathrm{y} / \delta<1.0$ ) there is close agreement between the LDV and PIV data. Above $\mathrm{y} / \delta=1.0$, the LDV data drops to $1 \%$ while the PIV stabilizes at $3 \%$. Klebanoff 's data was for incompressible flow and appears to set a lower boundary for the TI. The hot-film data from the present experiment matches extremely well with the LDV data as well as the Klebanoff data.

Figure 5-34 contains the FPG turbulence intensity comparison. Note that the stabilizing effect of the FPG is apparent in all measurement techniques. In the boundary layer, the PIV compares favorably with the LDV and hot-film data. The divergence between the PIV and LDV data occurs above $\mathrm{y} / \delta=1.0$ with the LDV predicting a TI of $1 \%$ while the PIV predicts $3 \%$. The hot-film data agrees with the LDV for near wall data and in the freestream; however, in the upper region of the boundary layer $(\mathrm{y} / \delta>0.2)$ the PIV under predicts the turbulence intensity.

Figure $5-35$ contains the CPG u turbulence intensity profile for $\mathrm{x}=68 \mathrm{~cm}$. The hot-film and LDV data match well; however, the PIV data consistently over predicts the TI by approximately $33 \%$ suggesting that an insufficient number of images was used to evaluate the flow.

Figure 5-36 contains the CPG $u$ turbulence intensity profile for $x=70 \mathrm{~cm}$. As with the $\mathrm{x}=68$ cm data, agreement between the hot-film and LDV data is excellent. Again, the PIV data consistently over predicts the TI by approximately $33 \%$. Thus, PIV does not capture the TI accurately for this configuration.


Figure 5-33: ZPG u turbulence intensity profiles


Figure 5-34: FPG u turbulence intensity profiles


Figure 5-35: $\mathrm{CPG}(\mathrm{x}=68 \mathrm{~cm}) \mathrm{u}$ turbulence intensity profile


Figure 5-36: $\mathrm{CPG}(\mathrm{x}=70 \mathrm{~cm}) \mathrm{u}$ turbulence intensity profile

The data in Figures 5-33 to 5-36 leads the conclusion that while PIV is accurate for mean flow quantities, it is inaccurate for turbulence quantities. Accuracy can be improved by increasing the number of images.

### 5.3.4 Density Fluctuations

Finally, the density fluctuations in the boundary layer were extracted from the hot-film data. Since PIV is only sensitive to the velocity, it cannot measure this fluctuation. The density fluctuation data is presented as Figure 5-37. The density fluctuation data exhibits the same FPG stabilizing trends and CPG destabilizing trends as indicated in Figures 5-26 to 5-29. The maximum density fluctuation is $10 \%$ for the ZPG, $7 \%$ for the FPG and $14 \%$ for the CPG. Finally, there is good agreement between the different traversing methods.


Figure 5-37: Density fluctuations

## 6. Mach 1.7 Results and Discussion

This chapter contains the results of the Mach 1.7 analysis. Since the main test condition for this study was Mach 2.8, only selected measurements were taken at Mach 1.7. Pressure and hot-wire data were acquired with the FPG model. Data were taken along the tunnel centerline at locations 2.5 and 11.5 cm downstream of the nozzle. The 2.5 cm location corresponded to ZPG flow and the 11.5 cm location was used as the FPG. The FPG location was the same as used in the Mach 2.8 configuration. The CPG model was not used because it "unstarted" the tunnel. Tunnel unstart is a condition where supersonic flow cannot be achieved due to flow blockage from either a model or a probe.

### 6.1 Mach 1.7 Conventional Pressure Data

Pitot pressure was taken at each measuring location. From the Pitot pressure, the methods of Sections 4.1.1.1 and 4.1.1.2 were used to compute the Mach number profile and velocity profile. Figures 6-1 to 6-3 present these results.


Figure 6-1: Mach 1.7 Pitot pressure profiles


Figure 6-2: Mach 1.7 Mach number profile


Figure 6-3: Mach 1.7 velocity profile

The data in these figures present the same trends as noted in the Mach 2.8 data. The data in these figures were used to calculate the boundary layer thickness. Table 6-1 presents the results of these calculations. Note that the values in Table 6-1 are lower than the boundary layer thickness in Mach 2.8 flow.

Table 6-1: Mach 1.7 boundary layer parameters

|  | ZPG | FPG |
| :--- | :--- | :--- |
| $\delta_{u}(\mathrm{~mm})$ | 3.91 | 4.70 |
| $\delta_{\mathrm{M}}(\mathrm{mm})$ | 4.06 | 4.90 |
| $\mathrm{M}_{\mathrm{e}}$ | 1.66 | 1.77 |
| $\mathrm{U}_{\mathrm{e}}(\mathrm{m} / \mathrm{sec})$ | 468 | 478 |

### 6.2 Hot-Wire Results

This section presents the Mach 1.7 hot-wire results. Since the boundary layer was not as thick as for the Mach 2.9 configuration, the averaging packet was reduced from 4,096 points to 2,048 points. This allowed a sufficient number of data points to be evaluated in the boundary layer.

### 6.2.1 Mean Flow Profiles

Figure 6-4 presents the Mach 1.7 ZPG mass flux ratio and total temperature ratio. Note the excellent agreement between the traverse up and traverse down data. The single overheat methodology aligns quite well with the multiple overheat data indicating that the total temperature is essentially constant through the boundary layer. This is confirmed by the total temperature profiles. Note that while the total temperature ratio is essentially constant for either the traverse up or traverse down profiles, there is a very slight difference between the traverse up and traverse down profiles. This difference is most likely due to a slight cooling of the flow straightener during the test run.


Figure 6-4: Mach 1.7 ZPG hot-film data

### 6.3 Turbulence Results

Hot-films were used to evaluate the fluctuation components of fluid properties. Hot-films were used to evaluate the mass flux, total temperature, and velocity fluctuations. Power spectra data were taken for both configurations, but turbulence information is only presented for the ZPG configuration.

### 6.3.1 Mass Flux Turbulence Intensity Profiles

The mass flux turbulence results are presented in Figure 6-5. There is excellent agreement between the traverse up and traverse down as well as the single versus multiple overheat data. The maximum TI decreased from $17 \%$ to $7 \%$ from Mach 2.9 to Mach 1.7 , a $60 \%$ reduction. The freestream TI was found to be $1.5 \%$.


Figure 6-5: Mach 1.7 ZPG mass flux turbulence intensity

### 6.3.2 Energy spectra

Figure 6-6 presents the power spectra results. As with the Mach 2.8 results, the ZPG data indicates that the energy evenly spread throughout the frequency spectrum. In the outer region, the energy is more heavily weighted towards the larger scale structures. However, the Mach 1.7 FPG data exhibits opposite trends than the Mach 2.8 FPG data. In Figure 6-5, the FPG consistently maintains a higher level of energy than the ZPG. The data in Figures 5-11 and 5-12 indicate that the FPG has a lower level of energy than the ZPG. The fact that the Mach number is subsonic over a substantial portion of the boundary layer may be responsible for this opposite trend (Spina et al., 1994). Another possible explaination is compressiblity. Since the Mach number is closer to one, the effects of compressilibity may be lessened.


Figure 6-6: Mach 1.7 power spectra data

### 6.3.3 Separated Turbulence Results

Figures 6-7 to 6-9 present the separated turbulence intensity results for velocity, density and total temperature, respectively.

Figure 6-7 presents the Mach 1.7 u turbulence intensity profiles. There is excellent agreement between separated turbulence intensity profiles during the differnent traversing methods. The $\mathbf{u}$ TI is approximately $10 \%$ at $\mathrm{y} / \delta=0.2$ and $1 \%$ in the freestream. This is an increase of $20 \%$ from the Mach 2.8 ZPG condition.

Figure 6-8 presents the Mach 1.7 density fluctuation profiles. The density turbulence intensity is the lowest fluctuation as compared to the velocity and total temperature fluctuations with a value of $5 \%$ at $\mathrm{y} / \delta=0.2$. The freestream values is approximately $0.75 \%$. This represents a $50 \%$ decrease from the Mach 2.8 ZPG value of density fluctuation. The agreement between the traverse up and traverse down data is not as good as it was for the $u$ turbulence intensity. This could be due to the variation in total temperature between the traverse up and traverse down data noted in Figure 6-4.


Figure 6-7: Mach 1.7 u turbulence intensity profiles


Figure 6-8: Mach 1.7 density fluctuation profiles

Figure 6-9 presents the Mach 1.7 total temperature fluctuation profiles. The temperature fluctuations lie between the density and velocity with a maximum value of $8 \%$. The freestream value of total temperature fluctuation is approximately $2 \%$. The scatter in the total temperature fluctuations is greater than the scatter in Figures 6-7 and 6-8. Again this is attributed to the change in total temperature between the traverse up and traverse down data.


Figure 6-9: Mach 1.7 total temperature fluctuation profiles
Since FPG or CPG data was not taken, no conclusions on the effect of pressure gradients on the turbulent quantities can be made.

From the data in this chapter it can be noted that reducing the Mach number reduces the mass flux turbulence intensity by $20 \%$. This was attributed to the $50 \%$ reduction in density turbulence intensity. Interestingly, the $u$ turbulence intensity increases when the Mach number is decreased.

## 7. Mach 5.0 Results and Discussion

Pitot pressure was taken at each measuring location. From the Pitot pressure, Rayleigh supersonic Pitot formula was used to compute the Mach number profile and velocity profile. Figures 7-1 to 7-3 present these results.


Figure 7-1: Mach 5.0 Pitot pressure profiles
The data in these figures present the same trends as noted in the Mach 2.8 data. Note the effect of the shock in the CPG Pitot pressure and Mach number profiles. The data in these figures were used to calculate the boundary layer thickness parameters. Table 7-1 presents the results of these calculations.

Table 7-1: Mach 5.0 boundary layer parameters

|  | ZPG | FPG | CPG $\left(\mathrm{x}_{\mathrm{ts}}=5.1 \mathrm{~cm}\right)$ | CPG $\left(\mathrm{x}_{\mathrm{ts}}=6.35 \mathrm{~cm}\right)$ |
| :--- | :---: | :---: | :---: | :---: |
| $\delta_{\mathrm{u}}(\mathrm{mm})$ | 8.46 | 10.8 | 6.17 | 5.51 |
| $\delta_{\mathrm{M}}(\mathrm{mm})$ | 12.2 | 13.1 | 7.09 | 6.60 |
| $\mathrm{M}_{\mathrm{e}}$ | 5.10 | 5.32 | 4.30 | 5.20 |
| $\mathrm{U}_{e}(\mathrm{~m} / \mathrm{sec})$ | 777 | 784 | 756 | 780 |



Figure 7-2: Mach 5.0 Mach number profile


Figure 7-3: Mach 5.0 velocity profile

The surface static pressure profile is presented as Figure 7-4. The FPG data clearly shows the effect of the shock near the end of the FPG surface. The CPG data shows the effect of the compression followed by the expansion on the wall static pressure. Wall static pressure ports are marked with an x in Figure 7-4.


Figure 7-4: Mach 5.0 CPG and FPG static pressure profiles

## 8. Conclusions and Recommendations

Bradshaw (Bradshaw, 1977) has noted that pressure gradients distort turbulent flow properties of a supersonic boundary layer by an order of magnitude more than indicated by the additional production. Spina (Spina et al., 1994) restated this concept almost 20 years later. During this time, little quality research has been done on these types of flows (Settles and Dodson, 1994). The intentions of this study were to help fill this void of information. Recall that the main objective of this study was to obtain a better understanding of supersonic boundary layers under the influence of pressure gradients. The objective was met by using a variety of measurement techniques (PIV, hot-wire, conventional probes and Mie-scattering flow visualization) to provide empirical information. A second goal was the evaluation of PIV as a flow diagnostic tool. The measurements taken in this study included conventional pressure probe pressure profiles, hot-wire mass flux turbulence intensity and energy spectra data and PIV for mean and fluctuating velocities.

### 8.1 Mach 2.8 Conclusions

PIV was validated for mean flow quantities as demonstrated by the comparison between LDV and PIV data. Application of Van Driest's scaling laws (Van Driest, 1951), satisfactorily reduced the PIV derived velocity data to theoretical values. PIV was not validated for fluctuation quantities with the number of samples used in this study. PIV consistently overpredicted the freestream turbulence intensity and underpredicted the turbulence intensity in the boundary layer. A likely explanations for this effect is PIV's low signal-to-noise ratio. For a further discussion of this effect see Appendix B. An increase in the number of images could aid in solving this problem.

Multiple overheat hot-wire anemometry was used to investigate the effects of pressure gradients on mean flow and fluctuation turbulence quantities. Traversing the hot-film probe into the wall allowed for accurate, near wall measurements to be taken. The ZPG mass flux turbulence intensity showed excellent agreement between single overheat and multiple overheat data reduction methodologies. The

FPG was found to reduce the maximum mass flux turbulence intensity by $25 \%$, from $17 \%$ to $12.5 \%$. The CPG was found to increase the maximum mass flux turbulence intensity $20 \%(x=68 \mathrm{~cm})$ and $22 \%(x=$ 70 cm ) from the $17 \% \mathrm{ZPG}$ value.

Separating the mass flux components into velocity and density fluctuations was accomplished for a single overheat ratio. The separated turbulence intensities accurately matched the LDV data by Luker (Luker et al., 1997). These data also confirmed the effects of the pressure gradients on the turbulent quantities.

Power spectra measurements were taken for the three pressure gradient models with a hot-wire tuned to give a high frequency response. The ZPG spectra indicated that the energy was evenly distributed among the different frequencies. Once any pressure gradient is applied the energy is shifted to the lower frequencies. This confirms the hypothesis that a FPG causes the large scale structures to disassociate into smaller scale structures. For the CPG, this effect can be explained by the reduction of the boundary layer thickness and the corresponding reduction in structure sizes. At a sufficiently high frequency, the ZPG maintained a higher level of energy than the FPG or CPG.

Recall that Settles and Dodson (Settles and Dodson, 1994) established a set of criteria used to judge data. It was a goal of this study to satisfy these criteria. Table 8-1 lists the criteria and how they were satisfied for the Mach 2.8 data.

Table 8-1: Verification of Settles and Dodson criteria

| Settles and Dodson criteria | How satisfied |
| :--- | :--- |
| Baseline applicability | Mach 2.8 flow |
| Simplicity | Simple model design |
| Specific applicability | Simple models and obtaining of turbulence data |
| Well defined experimental boundary conditions | See Luker (Luker, 1995) |
| Well defined error bounds | See appendix B |
| Adequate documentation of data | See appendix C |
| Adequate spatial resolution of the data | Near wall data acquired, PIV data |

### 8.2 Mach 1.7 Conclusions

Multiple overheat hot-wire measurements were taken for the ZPG configuration. The u velocity turbulence intensity was found to be $10 \%$ and the maximum mass flux turbulence intensity was found to be $7 \%$. Recall that for the Mach 2.8 tunnel, the ZPG u velocity turbulence intensity was $8 \%$ and the mass flux turbulence intensity was found to be $10 \%$. Note that in the Mach 1.7 tunnel the velocity turbulence intensity is higher while the mass flux turbulence intensity is higher in the Mach 2.8 tunnel. This is due to the density fluctuations. The maximum density fluctuation is approximately $9 \%$ in the Mach 2.9 tunnel while it peaks at $5 \%$ in the Mach 1.7 tunnel.

Power spectra measurements were taken for the ZPG and FPG in the Mach 1.7 tunnel. The cross-over noted in the Mach 2.9 tunnel was found in the Mach 1.7 tunnel. However, the FPG always maintained a higher energy level across the frequencies measured. This is opposite of what was noted in the Mach 2.9 tunnel. This was a result of the reduced level of compressibility and the fact that the Mach number was subsonic for a substantial portion of the boundary layer for this flow.

### 8.3 Mach 5.0 Conclusions

Conventional pressure measurements were taken in Mach 5.0 flow for the three pressure gradients. The effects of the pressure gradient on boundary layer thickness noted in Mach 2.8 and Mach 1.7 flow were also noted in Mach 5.0 flow. Using the ZPG as a reference, the FPG caused the boundary layer thickness to increase while the CPG caused the boundary layer thickness to decrease.

### 8.4 Recommendations

Several recommendations can be made to improve on the study. First, more images should be taken for PIV analysis. More images will allow a better estimate of turbulence intensity to be made and allow a better understanding of data closure. Currently, it is not known how many images are required to achieve data closure.

Second, multiple overheat cross-wire anemometry should be accomplished for all pressure gradient models. This will allow measurements of $v_{\text {ms }}^{\prime}$ and the cross correlation terms found in the Reynolds averaged Navier-Stokes equations. In addition, more accurate measurements of the wall shear stress should be made. This will allow a more accurate estimate of the Van Driest velocity and length correlation.

Third, the boundary layer power spectrum should be investigated in Mach 5.0 flow with a hotwire. The Mach 1.7 power spectra data should be repeated with a hot-wire, as opposed to the hot-film used in this study, to confirm the pressure gradient effect on the spectra.

Finally, hot-wire measurements should be taken in the Mach 5.0 tunnel. This will allow a better understanding of the effect of Mach number on the mean flow and turbulence quantities. In addition, the ceramic pebbles in the pebble bed heater should be replaced with stainless steel pebbles. This will eliminate the abrasion problem currently noted in the Mach 5.0 tunnel.

## Appendix A : Hot-Wire Calibration

This appendix presents additional hot-wire calibration data, including calibration curves.

## A.1 Mach 2.8 Calibration

Hot-wires were calibrated by placing them in the freestream and varying the total pressure in the plenum by adjusting the Kinney valve. With the Total temperature and Mach number as inputs, the MSHEaR code converted the plenum pressure and output bridge voltage to a Reynolds number and a Nusselt number. Recall that Equation (4.21) indicates that a the Nusselt number is a linear function of the square root of the Reynolds number. Tables A-1 and A-2 present raw calibration data for two different overheat ratios. Figures A-1 and A-2 present the linear curvefit and the raw data.

Table A-1: Mach 2.8 sample raw calibration file $(\mathrm{OHR}=2.03)$

| HW voltage | $\mathrm{P}_{\mathrm{t}}$ | $\mathrm{P}_{\mathrm{t} 2} / \mathrm{P}_{\mathrm{t} 1}$ | Mach | $\mathrm{T}_{\mathrm{t}}$ | $\mathrm{\rho u}$ | Re | $\mathrm{Re}^{1 / 2}$ | curve | Nu |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| (voit) | $(\mathrm{psi})$ |  |  | $(\mathrm{K})$ | $\left(\mathrm{kg} / \mathrm{m}^{2} \mathrm{~s}\right)$ |  |  | fit | data |
| 4.161 | 13.6 | 0.456 | 2.61 | 294.26 | 75.85 | 212.69 | 14.584 | 1.516 | 1.54 |
| 4.434 | 17.4 | 0.456 | 2.61 | 294.26 | 96.84 | 271.54 | 16.478 | 1.744 | 1.75 |
| 4.656 | 21.1 | 0.456 | 2.61 | 294.26 | 117.17 | 328.57 | 18.126 | 1.942 | 1.93 |
| 4.891 | 25.2 | 0.456 | 2.61 | 294.26 | 140.24 | 393.26 | 19.831 | 2.147 | 2.13 |
| 5.103 | 29.3 | 0.456 | 2.61 | 294.26 | 162.57 | 455.88 | 21.351 | 2.330 | 2.31 |
| 5.28 | 32.7 | 0.456 | 2.61 | 294.26 | 181.45 | 508.82 | 22.557 | 2.475 | 2.48 |
| 5.418 | 35.9 | 0.456 | 2.61 | 294.26 | 199.35 | 559.02 | 23.644 | 2.606 | 2.61 |
| 5.555 | 39.3 | 0.456 | 2.61 | 294.26 | 218.29 | 612.13 | 24.741 | 2.738 | 2.74 |
| 5.666 | 42.2 | 0.456 | 2.61 | 294.26 | 234.59 | 657.83 | 25.648 | 2.847 | 2.85 |
| 5.774 | 45.3 | 0.456 | 2.61 | 294.26 | 251.83 | 706.17 | 26.574 | 2.958 | 2.96 |

Table A-2: Mach 2.8 sample raw calibration file $(\mathrm{OHR}=1.66)$

| HW voltage | $\mathrm{P}_{\mathrm{t} 1}$ | $\mathrm{P}_{\mathrm{t} 2} / \mathrm{P}_{\mathrm{t} 1}$ | Mach | $\mathrm{T}_{\mathrm{t} 1}$ | $\rho \mathrm{u}$ | Re | $\mathrm{Re}^{1 / 2}$ | curve | Nu |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| (volt) | $(\mathrm{psi})$ |  |  | $(\mathrm{K})$ | $\left(\mathrm{kg} / \mathrm{m}^{2} \mathrm{~s}\right.$ |  |  | fit | data |
| 3.324 | 11.3 | 0.456 | 2.61 | 294.26 | 62.52 | 175.33 | 13.241 | 1.297 | 1.35 |
| 3.457 | 14.2 | 0.456 | 2.61 | 294.26 | 78.87 | 221.18 | 14.872 | 1.469 | 1.46 |
| 3.615 | 17 | 0.456 | 2.61 | 294.26 | 94.62 | 265.33 | 16.289 | 1.618 | 1.6 |
| 3.777 | 20.2 | 0.456 | 2.61 | 294.26 | 112.2 | 314.62 | 17.738 | 1.771 | 1.75 |
| 3.915 | 22.7 | 0.456 | 2.61 | 294.26 | 126.01 | 353.34 | 18.797 | 1.882 | 1.88 |
| 4.044 | 25.8 | 0.456 | 2.61 | 294.26 | 143.64 | 402.8 | 20.070 | 2.016 | 2 |
| 4.198 | 29.6 | 0.456 | 2.61 | 294.26 | 164.74 | 461.96 | 21.493 | 2.166 | 2.16 |
| 4.322 | 33.1 | 0.456 | 2.61 | 294.26 | 184.12 | 516.31 | 22.722 | 2.296 | 2.29 |
| 4.465 | 37.4 | 0.456 | 2.61 | 294.26 | 208.04 | 583.39 | 24.153 | 2.447 | 2.44 |
| 4.574 | 40.6 | 0.456 | 2.61 | 294.26 | 225.76 | 633.07 | 25.161 | 2.553 | 2.56 |
| 4.662 | 43.3 | 0.456 | 2.61 | 294.26 | 240.52 | 674.46 | 25.970 | 2.638 | 2.66 |
| 4.721 | 45.6 | 0.456 | 2.61 | 294.26 | 253.38 | 710.53 | 26.656 | 2.710 | 2.73 |



Figure A-1: Mach 2.8 hot-wire calibration curve $(\mathrm{OHR}=2.03)$


Figure A-2: Mach 2.8 hot-wire calibration curve $(\mathrm{OHR}=1.66)$
Note the excellent agreement between the data and the least-squares curvefit. The minimum correlation coefficient (squared) was always greater than 0.99 . This indicates that the curve fit is extremely accurate in modeling the data.

With this curve, a given bridge voltage could be easily and accurately converted to a mass flux. Table A-3 presents the constants used in Equation (4.21) for all the overheat ratios used in the Mach 2.8 portion of the experiment.

Table A-3: Mach 2.8 calibration constants

| OHR | a | b |
| :--- | :--- | :--- |
| 2.03 | 0.120 | -0.238 |
| 1.98 | 0.117 | -0.192 |
| 1.87 | 0.112 | -0.138 |
| 1.81 | 0.110 | -0.113 |
| 1.74 | 0.107 | -0.090 |
| 1.66 | 0.105 | -0.098 |
| 1.58 | 0.102 | -0.071 |
| 1.52 | 0.101 | -0.087 |

The constant a is the slope and b is the y -intercept. Note that a is only weakly dependent on OHR. Further study is required to fully examine this relationship. It might be possible to specify the slope and keep it constant throughout the overheat ratios.

## A.2 Mach 1.7 Calibration

Although the calibration theory was identical between the Mach 1.7 and Mach 2.9 calibration, application was different. The Mach 1.7 tunnel did not have an easy method to vary the plenum total pressure. Calibration cards were devised to solve this problem. A card with a series of holes punched in it was placed over the entrance to the tunnel. This card reduced the amount of flow that could enter the tunnel and thereby the plenum total pressure was reduced. A series of cards were created that established a different plenum pressure. Tables A-4 and A-5 present sample Mach 1.7 calibration data. Figure A-3 presents the least squares curvefit and the raw calibration data for all overheat ratios used in the Mach 1.7 phase of this study.

Table A-4: Mach 1.7 sample raw calibration file ( $\mathrm{OHR}=1.95$ )

| HW voltage | $\mathrm{P}_{\mathrm{tl}}$ | $\mathrm{P}_{\mathrm{t}} / \mathrm{P}_{\mathrm{t1}}$ | Mach | $\mathrm{T}_{\mathrm{t} 1}$ | $\rho \mathrm{pu}$ | Re | $\mathrm{Re}^{1 / 2}$ | curve | Nu |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| (volt) | $(\mathrm{psi})$ |  |  | $(\mathrm{K})$ | $\left(\mathrm{kg} / \mathrm{m}^{2} \mathrm{~s}\right.$ |  |  | fit | data |
| 5.138 | 13.4 | 0.817 | 1.79 | 299.26 | 151.59 | 419.59 | 20.484 | 2.474 | 2.49 |
| 4.861 | 11 | 0.817 | 1.79 | 299.26 | 123.7 | 342.4 | 18.504 | 2.252 | 2.23 |
| 4.355 | 6.5 | 0.817 | 1.79 | 299.26 | 73.52 | 203.51 | 14.266 | 1.777 | 1.79 |
| 3.655 | 3 | 0.817 | 1.79 | 299.26 | 33.85 | 93.69 | 9.679 | 1.262 | 1.26 |

Table A-5: : Mach 1.7 sample raw calibration file $(\mathrm{OHR}=1.66)$

| HW voltage | $\mathrm{P}_{\mathrm{t} 1}$ | $\mathrm{P}_{\mathrm{t} 2} / \mathrm{P}_{\mathrm{t} 1}$ | Mach | $\mathrm{T}_{\mathrm{t}}$ | $\rho \mathrm{l}$ | Re | $\mathrm{Re}^{1 / 2}$ | curve | Nu |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| (volt) | $\mathrm{psi})$ |  |  | $(\mathrm{K})$ | $\left(\mathrm{kg} / \mathrm{m}^{2} \mathrm{~s}\right.$ |  |  | fit | data |
| 4.352 | 13.5 | 0.817 | 1.79 | 299.26 | 151.71 | 419.94 | 20.492 | 2.320 | 2.33 |
| 4.106 | 10.8 | 0.817 | 1.79 | 299.26 | 122.15 | 338.11 | 18.388 | 2.110 | 2.07 |
| 3.782 | 6.6 | 0.817 | 1.79 | 299.26 | 73.95 | 204.7 | 14.307 | 1.702 | 1.76 |
| 3.138 | 3 | 0.817 | 1.79 | 299.26 | 33.35 | 92.31 | 9.608 | 1.233 | 1.21 |



Figure A-3: Mach 1.7 hot-wire calibration curves

The different overheats were plotted on the same axis in Figure A-3 to illustrate the insensitivity of the overheat ratio on the slope. Table A-6 presents the calibration constants for all five overheat ratios.

Table A-6: Mach 1.7 calibration constants

| OHR | a | b |
| :--- | :--- | :--- |
| 1.95 | 0.112 | 0.177 |
| 1.89 | 0.105 | 0.272 |
| 1.82 | 0.098 | 0.362 |
| 1.73 | 0.104 | 0.220 |
| 1.66 | 0.100 | 0.274 |

## Appendix B : Uncertainty Analysis

This section provides uncertainty estimates for the data presented in this thesis. This analysis augments the well-documented analysis by Luker (Luker, 1995) and Huffman (Dotter, 1995). Elementary uncertainties are presented first followed by the derived uncertainties.

Uncertainty analysis is necessary to determine the validity of the data. If the uncertainty is large or not presented it can call into question the appropriateness of the data and reflect on the reputation of the researcher. The absolute uncertainty, $\Delta f$ is found by taking the partial derivative of $f$ with respect to the dependent variable $\mathrm{x}_{\mathrm{i}}$ :

$$
\begin{equation*}
\Delta \mathrm{f}=\left.\sum_{\mathrm{i}=1}^{\mathrm{N}} \frac{\partial \mathrm{f}}{\partial \mathrm{x}_{\mathrm{i}}}\right|_{\mathrm{f}_{\mathrm{o}}} \Delta \mathrm{x}_{\mathrm{i}} \tag{B.1}
\end{equation*}
$$

where $f_{o}$ is the mean value of $f$. Taking a partial derivative is necessary to isolate the individual dependency of x on f . The uncertainty defined in Equation (B.1) can be converted to a percentage value by the following

$$
\begin{equation*}
\varepsilon_{\mathrm{f}}=100 \% * \frac{\Delta \mathrm{f}}{\mathrm{f}_{\mathrm{o}}} \tag{B.2}
\end{equation*}
$$

Summing each of these uncertainties over each dependent variable gives the maximum possible uncertainty. This assumes that each uncertainty reaches its maximum value simultaneously. A more reasonable approach uses the Euclidean norm defined as

$$
\begin{equation*}
\varepsilon_{\mathrm{f}}=\left\|\varepsilon_{\mathrm{i}}\right\|=\sqrt{\sum_{\mathrm{i}} \varepsilon_{\mathrm{i}}{ }^{2}} \tag{B.3}
\end{equation*}
$$

The Euclidean norm will be used to compute the uncertainties in this study.

## B. 1 Elementary Uncertainties

This section documents the elementary uncertainties. Uncertainties are based off the Mach 2.8 ZPG values given in Table B-1.

Table B-1: Mach 2.8 ZPG reference conditions

| Variable | Value |
| :--- | :--- |
| $\mathrm{T}_{\mathrm{t} 1}$ | 296 K |
| $\mathrm{P}_{\mathrm{t} 1}$ | $3.15 \times 10^{5} \mathrm{~Pa}$ |
| $\mathrm{P}_{\mathrm{t} 2}$ | $1.83 \times 10^{5} \mathrm{~Pa}$ |
| pc | $1.83 \times 10^{4} \mathrm{~Pa}$ |
| Mach | 2.78 |
| $\mathrm{P}_{\mathrm{e}}$ | $1.21 \times 10^{4} \mathrm{~Pa}$ |
| $\mathrm{P}_{\mathrm{w}}$ | $8.02 \times 10^{3} \mathrm{~Pa}$ |
| $\mathrm{~T}_{\mathrm{w}}$ | 296 K |
| $\mathrm{~T}_{\mathrm{e}}$ | 114 K |
| $\mathrm{P}_{\mathrm{e}}$ | $0.245 \mathrm{~kg} / \mathrm{m}^{3}$ |
| $\rho_{\mathrm{w}}$ | $0.095 \mathrm{~kg} / \mathrm{m}^{3}$ |
| $\mathrm{U}_{\mathrm{e}}$ | $604 \mathrm{~m} / \mathrm{sec}$ |
| $\mathrm{V}_{\mathrm{w}}$ | 5.0 Volts |
| a | 0.11 |
| b | -0.15 |
| $\mathrm{R}_{\mathrm{t}}$ | $0.2 \Omega$ |
| $\mathrm{R}_{\mathrm{w}}$ | $5.5 \Omega$ |
| $\mathrm{R}_{\mathrm{s}}$ | $50.0 \Omega$ |
| L | 1.0 mm |
| d | $5.1 \times 10^{-5} \mathrm{~m}$ |

Cone-static, Pitot and plenum pressure were measured by pressure transducers. The pressure transducers had two sources of uncertainty. They were a gain uncertainty of $0.5 \%$ and a gain stability uncertainty of $0.2 \%$. The ambient pressure was measured by a RDS corporation DP141 ambient pressure meter. The least significant digit was 0.0001 in Hg leading to an uncertainty in ambient pressure of 3.43 $\times 10^{-4} \%$. Since this uncertainty is two orders of magnitude less than the gain uncertainty and gain stability uncertainty, it will be neglected. Huffman (Huffman, 1995) found that the digital conversion uncertainty was $\pm 0.0041 \mathrm{~atm}$ for plenum pressure and $\pm 0.00046 \mathrm{~atm}$ for Pitot pressure. Combining the pressure source of uncertainty leads to a Euclidean pressure uncertainty of $0.54 \%$.

Two sources of uncertainty were inherent in the temperature measurement. Luker (Luker, 1995) found that the assumption of constant total temperature variation leads to an uncertainty of up to $4 \%$. Recall from Figure 5-10 the maximum scatter in total temperature is $5 \%$. Temperature measurement
uncertainty was found to equal $0.75 \%$ (Temperature handbook, 1992). Combining these leads to a Euclidean uncertainty of $5.1 \%$.

The position uncertainty of the probe was mainly due to probe deflection. It was found that the deflection was equal to $1.5^{\circ}$. Assuming that the tangent of this angle equals the uncertainty in $x$ location leads to an $x$ uncertainty of $2.6 \%$

Huffman found that the y position uncertainty was $4.1 \%$ and the hot-wire voltage uncertainty was 1.0\% (Huffman, 1995).

Table B-2 summarizes the elementary uncertainty data.

Table B-2: Elementary uncertainties

| Elementary | Value |
| :--- | :--- |
| $\varepsilon_{\mathrm{Pt} 1}$ | $0.54 \%$ |
| $\varepsilon_{\mathrm{Pt} 2}$ | $0.54 \%$ |
| $\varepsilon_{\mathrm{Pcs}}$ | $0.54 \%$ |
| $\varepsilon_{\mathrm{Tt} 1}$ | $5.1 \%$ |
| $\varepsilon_{\mathrm{x}}$ | $2.6 \%$ |
| $\varepsilon_{\mathrm{y}}$ | $4.1 \%$ |
| $\varepsilon_{\mathrm{VW}}$ | $1.0 \%$ |

## B. 2 Derived Uncertainties

The uncertainties described in Section B-1 propagate into derived uncertainties. Derived uncertainties are uncertainties in quantities derived from the variables listed in Table B-2.

## B.2.1 Mean Flow Derived Uncertainties

Recall from chapter 4, that the Mach number was calculated based on the cone-static pressure and Pitot pressure. Equation 4.1 was linearized and the uncertainties listed in Table B-2 were applied. With the Mach number uncertainty, the isentropic equations were linearized and used to compute the uncertainty in pressure and temperature. With the pressure and temperature uncertainty known, the density uncertainty could be found from the linearized equation of state. The uncertainty in the viscosity was calculated by linearizing Sutherland's law. Table B-3 presents the mean flow derived uncertainties.

Table B-3: Mean flow derived uncertainties

| Uncertainty | Equation | Value |
| :--- | :--- | :--- |
| $\varepsilon_{\mathrm{Me}}$ | See equation 4.1 | $0.62 \%$ |
| $\varepsilon_{\mathrm{Pw}}$ | $\left\\|\varepsilon_{\mathrm{Pw}}\right\\|$ | $0.58 \%$ |
| $\varepsilon_{\mathrm{Pe}}$ | $\left\\|\varepsilon_{\mathrm{Pt} 2}, 4.25 \varepsilon_{\mathrm{Me}}\right\\|$ | $2.69 \%$ |
| $\varepsilon_{\mathrm{Te}}$ | $\left\\|\varepsilon_{\mathrm{Pt} 2}, 1.214 \varepsilon_{\mathrm{Me}}\right\\|$ | $0.93 \%$ |
| $\varepsilon_{\mathrm{Tw}}$ | $\left\\|\varepsilon_{\mathrm{T}}\right\\|$ | $5.1 \%$ |
| $\varepsilon_{\mathrm{\rho w}}$ | $\left\\|\varepsilon_{\mathrm{Tw}}, \varepsilon_{\mathrm{Pw}}\right\\|$ | $5.1 \%$ |
| $\varepsilon_{\mathrm{Pe}}$ | $\left\\|\varepsilon_{\mathrm{Te}}, \varepsilon_{\mathrm{Pe}}\right\\|$ | $2.9 \%$ |
| $\varepsilon_{\mathrm{Ue}}$ | $\left\\|\varepsilon_{\mathrm{Me}}, 0.5 \varepsilon_{\mathrm{Te}}\right\\|$ | $0.78 \%$ |
| $\varepsilon_{\mu}$ | $\\|$ Sutherland's Law $\\|$ | $2.0 \%$ |

Uncertainty in the boundary layer was calculated graphically as in Luker (Luker, 1995). The boundary height uncertainties are listed in Table B-4.

Table B-4: Boundary layer height uncertainty

| Uncertainty | Value |
| :--- | :--- |
| $\varepsilon_{\delta \mathrm{M}}$ | $21.7 \%$ |
| $\varepsilon_{\delta \mathrm{u}}$ | $5.1 \%$ |

The uncertainty in the van Driest correlation was calculated as in Luker (Luker, 1995). The van Driest uncertainties are listed in Table B-5.

Table B-5: Van Driest uncertainties

| Uncertainty | Equation | Value |
| :--- | :--- | :--- |
| $\varepsilon_{\mathrm{ueff}}$ | $\left\\|0.5 \varepsilon_{\rho e}, 0.5 \varepsilon_{\rho \mathrm{pw}}, 2.0 \varepsilon_{\mathrm{u}}\right\\|$ | $3.3 \%$ |
| $\varepsilon_{\mathrm{qw}}$ | See Luker $($ Luker, 1995$)$ | $7.0 \%$ |
| $\varepsilon_{\mathrm{u}^{*}}$ | $\left\\|0.5 \varepsilon_{\mathrm{w}}, 0.5 \varepsilon_{\rho \mathrm{w}}\right\\|$ | $4.3 \%$ |
| $\varepsilon_{\mathrm{y}}^{+}$ | $\left\\|\varepsilon_{\mathrm{y}}, \varepsilon_{\mathrm{u}}, \varepsilon_{\mathrm{E}}, \varepsilon_{\rho \mathrm{pw}}\right\\|$ | $8.8 \%$ |
| $\varepsilon_{\mathrm{ueff}}{ }^{+}$ | $\left\\|0.5 \varepsilon_{\mathrm{rw}}, 0.5 \varepsilon_{\mathrm{ueff}}\right\\|$ | $7.7 \%$ |

## B.2.2 Hot-Wire Derived Uncertainties

The uncertainty in the Nusselt number was determined from Equation 4.22. The Reynolds number uncertainty was found by linearizing King's law (Equation 4.21) The separation of variables
uncertainty was determined by the outline given in Huffman (Huffman, 1995). The results of the hot-wire derived uncertainty analyses are given in Table B-6.

Table B-6: Hot-wire derived uncertainties

| Hot-wire | derived | Equation | Value |
| :---: | :---: | :---: | :---: |
| $\varepsilon_{\mathrm{Nu}}$ |  | $\\| 2.0 \varepsilon_{\mathrm{Vw}, \varepsilon_{\text {Te }} \\|}$ | 5.4\% |
| $\varepsilon_{\mathrm{Re}}$ |  | $\left\\|\varepsilon_{\mathrm{Vw}}, 2.0 \varepsilon_{\mathrm{Nu}}\right\\|$ | 6.5\% |
| $\varepsilon_{(\rho \mathrm{pl}}$ |  | $\varepsilon_{\text {Re }}$ | 6.5\% |
| $\varepsilon_{(\text {(pu) }}$ |  | $\varepsilon_{\text {Re }}$ | 6.5\% |
| $\varepsilon_{\rho^{\prime}}$ |  | $\left\\|0.46 \varepsilon_{\text {Me }}, 0.77 \varepsilon_{(\text {(pu) }}, \varepsilon_{\text {Ti }}\right\\|$ | 7.2\% |
| $\varepsilon_{u}$ |  | $\left\\|\varepsilon_{\rho^{\prime}}, \varepsilon_{(\text {(u) }}\right\\|^{\prime \prime}$ | 9.7\% |

## B. 3 PIV Signal-to-Noise Ratio Uncertainty

The signal-to-noise ratio for the PIV system is a function of the temporal separation consistency of the seed particles, the camera alignment, the camera focus, the pixel resolution of the camera and the soffware data reduction algorithm. The uncertainty of these effects cannot be quantified at this time; however, by inspection of the PIV results, it can be shown that the combination of these effects accounts for a $3 \%$ increase in the axial turbulence intensity.

## Appendix C : Data Files

This section contains the data necessary to reproduce all the figures in this thesis. The Mach 2.8
data is presented first (pressure followed by hot-wire and PIV) followed by the Mach 1.7 data (pressure and hot-wire) and the Mach 5.0 data (pressure only). All rms values have been non-dimensionalized by their respective freestream value.

Table C-1: Mach 2.8 ZPG pressure data

| $y$ (in) | $\mathrm{P}_{\mathrm{t} 1}(\mathrm{psia})$ | $\mathrm{P}_{12}$ (psia) | PT2/PT1 | Ps/PT2 | Mach | M/Me | u/Ue | $\rho \mathrm{w} / \rho$ | $\rho / \mathrm{p}$ ¢ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.0625 | 31.8164 | 3.7234 | 0.1170 | 0.3015 | 1.4750 | 0.5231 | 0.7027 | 0.6882 | 0.5609 |
| 0.0699 | 31.7446 | 4.1989 | 0.1323 | 0.2674 | 1.5838 | 0.5616 | 0.7376 | 0.6589 | 0.5859 |
| 0.0801 | 31.7430 | 4.9182 | 0.1549 | 0.2283 | 1.7352 | 0.6153 | 0.7824 | 0.6194 | 0.6232 |
| 0.0910 | 31.7384 | 5.2463 | 0.1653 | 0.2140 | 1.8002 | 0.6384 | 0.8003 | 0.6030 | 0.6402 |
| 0.1021 | 31.7430 | 5.6387 | 0.1776 | 0.1991 | 1.8741 | 0.6646 | 0.8198 | 0.5848 | 0.6601 |
| 0.1128 | 31.8429 | 5.9274 | 0.1861 | 0.1894 | 1.9271 | 0.6834 | 0.8331 | 0.5721 | 0.6748 |
| 0.1234 | 31.6556 | 6.0913 | 0.1924 | 0.1843 | 1.9550 | 0.6932 | 0.8400 | 0.5655 | 0.6827 |
| 0.1337 | 31.7946 | 6.6252 | 0.2084 | 0.1694 | 2.0491 | 0.7266 | 0.8622 | 0.5437 | 0.7100 |
| 0.1445 | 31.7352 | 6.4666 | 0.2038 | 0.1736 | 2.0232 | 0.7174 | 0.8562 | 0.5496 | 0.7024 |
| 0.1554 | 31.8008 | 6.5669 | 0.2065 | 0.1709 | 2.0404 | 0.7235 | 0.8602 | 0.5457 | 0.7074 |
| 0.1666 | 31.7836 | 7.1332 | 0.2244 | 0.1574 | 2.1344 | 0.7569 | 0.8812 | 0.5247 | 0.7357 |
| 0.1768 | 31.7477 | 6.7134 | 0.2115 | 0.1672 | 2.0668 | 0.7329 | 0.8663 | 0.5397 | 0.7153 |
| 0.1873 | 31.7477 | 7.2332 | 0.2278 | 0.1552 | 2.1502 | 0.7625 | 0.8846 | 0.5213 | 0.7406 |
| 0.1980 | 31.6837 | 7.2684 | 0.2294 | 0.1544 | 2.1563 | 0.7647 | 0.8859 | 0.5199 | 0.7425 |
| 0.2090 | 31.8508 | 8.0338 | 0.2522 | 0.1397 | 2.2749 | 0.8067 | 0.9102 | 0.4949 | 0.7800 |
| 0.2195 | 31.6338 | 8.1745 | 0.2584 | 0.1373 | 2.2949 | 0.8138 | 0.9141 | 0.4908 | 0.7865 |
| 0.2300 | 31.7836 | 8.3197 | 0.2618 | 0.1349 | 2.3162 | 0.8214 | 0.9182 | 0.4865 | 0.7934 |
| 0.2407 | 31.7056 | 8.4487 | 0.2665 | 0.1329 | 2.3355 | 0.8282 | 0.9218 | 0.4827 | 0.7998 |
| 0.2517 | 31.6900 | 8.5679 | 0.2704 | 0.1310 | 2.3532 | 0.8345 | 0.9252 | 0.4792 | 0.8056 |
| 0.2626 | 31.7212 | 8.7184 | 0.2748 | 0.1288 | 2.3743 | 0.8419 | 0.9291 | 0.4750 | 0.8126 |
| 0.2732 | 31.7493 | 9.1628 | 0.2886 | 0.1225 | 2.4398 | 0.8652 | 0.9409 | 0.4624 | 0.8348 |
| 0.2837 | 31.7212 | 9.5033 | 0.2996 | 0.1181 | 2.4880 | 0.8823 | 0.9492 | 0.4534 | 0.8514 |
| 0.2947 | 31.7352 | 9.8194 | 0.3094 | 0.1143 | 2.5314 | 0.8977 | 0.9565 | 0.4455 | 0.8666 |
| 0.3058 | 31.7540 | 9.7267 | 0.3063 | 0.1154 | 2.5238 | 0.8950 | 0.9552 | 0.4468 | 0.8639 |
| 0.3165 | 31.7618 | 10.1273 | 0.3189 | 0.1108 | 2.5747 | 0.9130 | 0.9636 | 0.4377 | 0.8819 |
| 0.3271 | 31.7680 | 10.3382 | 0.3254 | 0.1086 | 2.6029 | 0.9230 | 0.9681 | 0.4327 | 0.8920 |
| 0.3376 | 31.6603 | 10.3509 | 0.3269 | 0.1085 | 2.6057 | 0.9240 | 0.9685 | 0.4323 | 0.8931 |
| 0.3487 | 31.7009 | 10.7962 | 0.3406 | 0.1040 | 2.6602 | 0.9433 | 0.9769 | 0.4229 | 0.9128 |
| 0.3598 | 31.7743 | 10.8290 | 0.3408 | 0.1037 | 2.6681 | 0.9461 | 0.9781 | 0.4216 | 0.9157 |
| 0.3704 | 31.8258 | 11.2041 | 0.3520 | 0.1002 | 2.7122 | 0.9618 | 0.9847 | 0.4142 | 0.9320 |
| 0.3809 | 31.7087 | 11.2488 | 0.3548 | 0.0998 | 2.7216 | 0.9651 | 0.9861 | 0.4127 | 0.9355 |
| 0.3916 | 31.6697 | 11.3907 | 0.3597 | 0.0986 | 2.7386 | 0.9711 | 0.9886 | 0.4099 | 0.9418 |
| 0.4026 | 31.7961 | 11.6797 | 0.3673 | 0.0961 | 2.7722 | 0.9830 | 0.9933 | 0.4045 | 0.9544 |
| 0.4134 | 31.6509 | 11.6012 | 0.3665 | 0.0968 | 2.7678 | 0.9815 | 0.9927 | 0.4052 | 0.9528 |
| 0.4240 | 31.6275 | 11.8733 | 0.3754 | 0.0945 | 2.8005 | 0.9931 | 0.9973 | 0.4000 | 0.9652 |
| 0.4344 | 31.8039 | 12.0154 | 0.3778 | 0.0934 | 2.8181 | 0.9993 | 0.9997 | 0.3972 | 0.9718 |
| 0.4455 | 31.7134 | 12.0840 | 0.3810 | 0.0929 | 2.8267 | 1.0024 | 1.0009 | 0.3959 | 0.9751 |
| 0.4562 | 31.7056 | 12.1360 | 0.3828 | 0.0925 | 2.8331 | 1.0047 | 1.0018 | 0.3949 | 0.9776 |
| 0.4670 | 31.7540 | 12.1487 | 0.3826 | 0.0924 | 2.8348 | 1.0053 | 1.0020 | 0.3946 | 0.9782 |
| 0.4775 | 31.7868 | 12.2189 | 0.3844 | 0.0919 | 2.8431 | 1.0082 | 1.0031 | 0.3933 | 0.9814 |
| 0.4882 | 31.7868 | 12.2594 | 0.3857 | 0.0916 | 2.8482 | 1.0100 | 1.0038 | 0.3926 | 0.9834 |
| 0.4991 | 31.6572 | 12.3343 | 0.3896 | 0.0910 | 2.8571 | 1.0131 | 1.0050 | 0.3912 | 0.9868 |
| 0.5101 | 31.7774 | 12.3456 | 0.3885 | 0.0909 | 2.8588 | 1.0137 | 1.0052 | 0.3909 | 0.9875 |


| 0.5209 | 31.6884 | 12.4275 | 0.3922 | 0.0903 | 2.8683 | 1.0171 | 1.0065 | 0.3895 | 0.9911 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.5313 | 31.7430 | 12.5343 | 0.3949 | 0.0896 | 2.8809 | 1.0216 | 1.0082 | 0.3876 | 0.9960 |
| 0.5422 | 31.6619 | 12.6078 | 0.3982 | 0.0890 | 2.8899 | 1.0248 | 1.0094 | 0.3862 | 0.9995 |
| 0.5531 | 31.7150 | 12.6781 | 0.3998 | 0.0885 | 2.8983 | 1.0278 | 1.0104 | 0.3850 | 1.0028 |
| 0.5641 | 31.8632 | 12.6832 | 0.3981 | 0.0885 | 2.8994 | 1.0282 | 1.0106 | 0.3848 | 1.0032 |
| 0.5747 | 31.7227 | 12.6436 | 0.3986 | 0.0888 | 2.8950 | 1.0266 | 1.0100 | 0.3854 | 1.0015 |
| 0.5851 | 31.7071 | 12.6118 | 0.3978 | 0.0890 | 2.8909 | 1.0251 | 1.0095 | 0.3861 | 0.9999 |
| 0.5960 | 31.6744 | 12.5701 | 0.3969 | 0.0893 | 2.8859 | 1.0234 | 1.0088 | 0.3868 | 0.9980 |
| 0.6071 | 31.7415 | 12.5273 | 0.3947 | 0.0896 | 2.8807 | 1.0215 | 1.0082 | 0.3876 | 0.9960 |
| 0.6181 | 31.6494 | 12.4968 | 0.3949 | 0.0898 | 2.8771 | 1.0202 | 1.0077 | 0.3881 | 0.9945 |
| 0.6286 | 31.7649 | 12.4535 | 0.3921 | 0.0901 | 2.8718 | 1.0184 | 1.0070 | 0.3889 | 0.9925 |
| 0.6392 | 31.8383 | 12.4083 | 0.3897 | 0.0905 | 2.8664 | 1.0164 | 1.0063 | 0.3898 | 0.9904 |
| 0.6500 | 31.9475 | 12.4074 | 0.3884 | 0.0905 | 2.8663 | 1.0164 | 1.0062 | 0.3898 | 0.9904 |
| 0.6608 | 31.6322 | 12.2533 | 0.3874 | 0.0916 | 2.8472 | 1.0096 | 1.0037 | 0.3927 | 0.9830 |
| 0.6717 | 31.6962 | 12.2557 | 0.3867 | 0.0916 | 2.8479 | 1.0099 | 1.0038 | 0.3926 | 0.9833 |
| 0.6819 | 31.6728 | 12.1899 | 0.3849 | 0.0921 | 2.8399 | 1.0070 | 1.0027 | 0.3938 | 0.9802 |
| 0.6927 | 31.8242 | 12.1768 | 0.3826 | 0.0922 | 2.8383 | 1.0065 | 1.0025 | 0.3941 | 0.9796 |
| 0.7037 | 31.6962 | 12.1156 | 0.3822 | 0.0927 | 2.8308 | 1.0038 | 1.0015 | 0.3952 | 0.9767 |
| 0.7147 | 31.7446 | 12.1409 | 0.3825 | 0.0925 | 2.8340 | 1.0049 | 1.0019 | 0.3947 | 0.9779 |
| 0.7251 | 31.7477 | 12.1180 | 0.3817 | 0.0926 | 2.8311 | 1.0039 | 1.0015 | 0.3952 | 0.9768 |
| 0.7357 | 31.6681 | 12.0417 | 0.3802 | 0.0932 | 2.8217 | 1.0006 | 1.0002 | 0.3966 | 0.9732 |
| 0.7466 | 31.6494 | 12.0302 | 0.3801 | 0.0933 | 2.8204 | 1.0001 | 1.0000 | 0.3969 | 0.9727 |
| 0.7576 | 31.6307 | 11.9588 | 0.3781 | 0.0939 | 2.8115 | 0.9970 | 0.9988 | 0.3982 | 0.9693 |
| 0.7683 | 31.6978 | 11.9728 | 0.3777 | 0.0938 | 2.8133 | 0.9976 | 0.9991 | 0.3980 | 0.9700 |
| 0.7787 | 31.6681 | 11.9801 | 0.3783 | 0.0937 | 2.8142 | 0.9979 | 0.9992 | 0.3978 | 0.9704 |
| 0.7894 | 31.6728 | 11.9625 | 0.3777 | 0.0938 | 2.8120 | 0.9972 | 0.9989 | 0.3982 | 0.9695 |
| 0.8002 | 31.7243 | 11.9639 | 0.3771 | 0.0938 | 2.8122 | 0.9972 | 0.9989 | 0.3981 | 0.9696 |
| 0.8113 | 31.5916 | 11.9185 | 0.3773 | 0.0942 | 2.8066 | 0.9952 | 0.9982 | 0.3990 | 0.9675 |
| 0.8220 | 31.6385 | 11.9393 | 0.3774 | 0.0940 | 2.8092 | 0.9962 | 0.9985 | 0.3986 | 0.9684 |
| 0.8287 | 31.6338 | 11.9347 | 0.3773 | 0.0941 | 2.8086 | 0.9960 | 0.9984 | 0.3987 | 0.9682 |
| 0.8349 | 31.7633 | 11.9855 | 0.3773 | 0.0937 | 2.8148 | 0.9982 | 0.9993 | 0.3977 | 0.9706 |
| 0.8457 | 31.7711 | 11.9857 | 0.3773 | 0.0937 | 2.8149 | 0.9982 | 0.9993 | 0.3977 | 0.9706 |
| 0.8565 | 31.6681 | 11.9546 | 0.3775 | 0.0939 | 2.8111 | 0.9968 | 0.9988 | 0.3983 | 0.9692 |
| 0.8674 | 31.6431 | 11.9445 | 0.3775 | 0.0940 | 2.8098 | 0.9964 | 0.9986 | 0.3985 | 0.9687 |

Table C-2: Mach 2.8 FPG pressure data

| $y$ (in) | $\mathrm{P}_{11}$ (psia) | $\mathrm{P}_{12}$ (psia) | $\mathrm{P}_{12} / \mathrm{P}_{11}$ | Pc (psia) | $\mathrm{Pc} / \mathrm{P}_{11}$ | Mach | P (Pa) | $\rho\left(\mathrm{kg} / \mathrm{m}^{3}\right)$ | $\rho / \mathrm{pe}$ | $\mathrm{M} / \mathrm{M}_{\mathrm{e}}$ | $\mathrm{u} / \mathrm{U}_{\mathrm{e}}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.0625 | 31.5642 | 2.7763 | 0.0880 | 1.0600 | 3.36E-02 | 1.46674 | 4984 | 0.0844 | 0.4304 | 0.4889 | 0.6841 |
| 0.0695 | 31.5736 | 2.7742 | 0.0879 | 1.0572 | 3.35E-02 | 1.46736 | 5002 | 0.0847 | 0.4321 | 0.4891 | 0.6843 |
| 0.0799 | 31.5720 | 3.0985 | 0.0981 | 1.0420 | 3.30E-02 | 1.55645 | 5030 | 0.0884 | 0.4509 | 0.5188 | 0.7125 |
| 0.0906 | 31.5767 | 3.8072 | 0.1206 | 1.0267 | 3.25E-02 | 1.79615 | 5058 | 0.0985 | 0.5025 | 0.5987 | 0.7811 |
| 0.1015 | 31.5642 | 4.1069 | 0.1301 | 1.0118 | 3.21E-02 | 1.91195 | 5087 | 0.1043 | 0.5317 | 0.6373 | 0.8105 |
| 0.1124 | 31.5501 | 4.3158 | 0.1368 | 0.9989 | 3.17E-02 | 1.99732 | 5115 | 0.1089 | 0.5553 | 0.6658 | 0.8309 |
| 0.1232 | 31.5501 | 4.4317 | 0.1405 | 0.9937 | 3.15E-02 | 2.04183 | 5144 | 0.1117 | 0.5696 | 0.6806 | 0.8410 |
| 0.1338 | 31.5548 | 4.5862 | 0.1453 | 0.9902 | 3.14E-02 | 2.0939 | 5172 | 0.1149 | 0.5861 | 0.6980 | 0.8525 |
| 0.1448 | 31.5330 | 4.7565 | 0.1508 | 0.9912 | 3.14E-02 | 2.14761 | 5201 | 0.1184 | 0.6037 | 0.7159 | 0.8639 |
| 0.1558 | 31.5642 | 4.9763 | 0.1577 | 0.9919 | $3.14 \mathrm{E}-02$ | 2.2123 | 5230 | 0.1225 | 0.6250 | 0.7374 | 0.8772 |
| 0.1667 | 31.5298 | 5.0779 | 0.1611 | 0.9937 | 3.15E-02 | 2.24252 | 5259 | 0.1249 | 0.6369 | 0.7475 | 0.8832 |
| 0.1773 | 31.5517 | 5.2868 | 0.1676 | 0.9935 | 3.15E-02 | 2.30479 | 5287 | 0.1291 | 0.6584 | 0.7683 | 0.8952 |
| 0.1880 | 31.5642 | 5.4781 | 0.1736 | 0.9991 | 3.17E-02 | 2.35252 | 5315 | 0.1326 | 0.6762 | 0.7842 | 0.9040 |
| 0.1990 | 31.5533 | 5.6155 | 0.1780 | 1.0045 | 3.18E-02 | 2.38425 | 5344 | 0.1352 | 0.6896 | 0.7947 | 0.9097 |
| 0.2100 | 31.5470 | 5.7387 | 0.1819 | 1.0111 | 3.20E-02 | 2.41033 | 5373 | 0.1375 | 0.7014 | 0.8034 | 0.9144 |
| 0.2207 | 31.5611 | 5.8452 | 0.1852 | 1.0181 | 3.23E-02 | 2.42947 | 5401 | 0.1395 | 0.7112 | 0.8098 | 0.9177 |
| 0.2312 | 31.5392 | 6.1150 | 0.1939 | 1.0267 | 3.26E-02 | 2.49267 | 5429 | 0.1442 | 0.7352 | 0.8309 | 0.9284 |
| 0.2421 | 31.5408 | 6.2417 | 0.1979 | 1.0354 | 3.28E-02 | 2.51458 | 5458 | 0.1464 | 0.7463 | 0.8382 | 0.9320 |
| 0.2531 | 31.5564 | 6.4709 | 0.2051 | 1.0438 | 3.31E-02 | 2.56395 | 5487 | 0.1504 | 0.7670 | 0.8546 | 0.9400 |
| 0.2641 | 31.5470 | 6.6711 | 0.2115 | 1.0530 | 3.34E-02 | 2.60334 | 5516 | 0.1538 | 0.7846 | 0.8678 | 0.9461 |
| 0.2749 | 31.5330 | 6.6851 | 0.2120 | 1.0607 | 3.36E-02 | 2.59559 | 5545 | 0.1541 | 0.7859 | 0.8652 | 0.9449 |


| 0.2854 | 31.5423 | 6.8195 | 0.2162 | 1.0703 | $3.39 \mathrm{E}-02$ | 2.61472 | 5572 | 0.1562 | 0.7966 | 0.8716 | 0.9479 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.2963 | 31.5377 | 6.9329 | 0.2198 | 1.0808 | 3.43E-02 | 2.62802 | 5914 | 0.1668 | 0.8505 | 0.8760 | 0.9499 |
| 0.3074 | 31.5252 | 7.1153 | 0.2257 | 1.0886 | 3.45E-02 | 2.66311 | 5914 | 0.1694 | 0.8637 | 0.8877 | 0.9552 |
| 0.3183 | 31.5345 | 7.4182 | 0.2352 | 1.0984 | 3.48E-02 | 2.7259 | 5914 | 0.1741 | 0.8879 | 0.9086 | 0.9643 |
| 0.3289 | 31.5392 | 7.6224 | 0.2417 | 1.1056 | 3.51E-02 | 2.76532 | 5914 | 0.1771 | 0.9034 | 0.9218 | 0.9698 |
| 0.3396 | 31.5298 | 7.7102 | 0.2445 | 1.1155 | 3.54E-02 | 2.77207 | 5914 | 0.1777 | 0.9060 | 0.9240 | 0.9708 |
| 0.3506 | 31.5267 | 7.9233 | 0.2513 | 1.1230 | 3.56E-02 | 2.81412 | 5914 | 0.1810 | 0.9228 | 0.9380 | 0.9765 |
| 0.3723 | 31.5252 | 7.9605 | 0.2525 | 1.1377 | 3.61E-02 | 2.79831 | 5914 | 0.1797 | 0.9165 | 0.9328 | 0.9744 |
| 0.3829 | 31.537 | 8.0095 | 0.2540 | 1.1433 | 3.63E-02 | 2.79964 | 5914 | 0.1798 | 0.9170 | 0.9332 | 0.9745 |
| 0.3938 | 31.5267 | 8.4824 | 0.2691 | 1.1515 | 3.65E-02 | 2.90437 | 5914 | 0.1882 | 0.9597 | 0.9681 | 0.9883 |
| 0.4047 | 31.5142 | 8.4452 | 0.2680 | 1.1574 | 3.67E-02 | 2.88447 | 5914 | 0.1866 | 0.9514 | 0.9615 | 0.9857 |
| 0.4157 | 31.5205 | 8.6681 | 0.2750 | 1.1658 | $3.70 \mathrm{E}-02$ | 2.92461 | 5914 | 0.1898 | 0.9681 | 0.9749 | 0.9908 |
| 0.4264 | 31.5236 | 8.6150 | 0.2733 | 1.1717 | $3.72 \mathrm{E}-02$ | 2.90193 | 5914 | 0.1880 | 0.9587 | 0.9673 | 0.9879 |
| 0.4368 | 31.5236 | 8.7864 | 0.2787 | 1.1759 | $3.73 \mathrm{E}-02$ | 2.93638 | 5914 | 0.1908 | 0.9730 | 0.9788 | 0.9923 |
| 0.4479 | 31.5018 | 8.8622 | 0.2813 | 1.1829 | $3.76 \mathrm{E}-02$ | 2.94384 | 5914 | 0.1914 | 0.9762 | 0.9813 | 0.9932 |
| 0.4590 | 31.5127 | 9.0559 | 0.2874 | 1.1895 | 3.77E-02 | 2.9784 | 591 | 0.1943 | 0.9908 | 0.9928 | 0.9974 |
| 0.4698 | 31.4986 | 9.2628 | 0.2941 | 1.2014 | 3.81E-02 | 3.00698 | 5914 | 0.1967 | 1.0030 | 1.0023 | 1.0008 |
| 0.4804 | 31.5064 | 9.2139 | 0.2924 | 1.2070 | 3.83E-02 | 2.98276 | 5914 | 0.1946 | 0.9926 | 0.9943 | 0.9979 |
| 0.4911 | 31.5142 | 9.4104 | 0.2986 | 1.2157 | 3.86E-02 | 3.01616 | 5914 | 0.1975 | 1.0069 | 1.0054 | 1.0019 |
| 0.5022 | 31.5111 | 9.4396 | 0.2996 | 1.2253 | 3.89E-02 | 3.00424 | 5914 | 0.1964 | 1.0018 | 1.0014 | 1.0005 |
| 0.5131 | 31.5002 | 9.5108 | 0.3019 | 1.2347 | 3.92E-02 | 3.00413 | 5914 | 0.1964 | 1.0018 | 1.0014 | 1.0005 |
| 0.5238 | 31.4924 | 9.5902 | 0.3045 | 1.2410 | 3.94E-02 | 3.01189 | 5914 | 0.1971 | 1.0051 | 1.0040 | 1.0014 |
| 0.5343 | 31.4924 | 9.6703 | 0.3071 | 1.2518 | 3.97E-02 | 3.01159 | 5914 | 0.1971 | 1.0050 | 1.0039 | 1.0014 |
| 0.5451 | 31.4877 | 9.7810 | 0.3106 | 1.2607 | 4.00E-02 | 3.02307 | 5914 | 0.1980 | 1.0099 | 1.0077 | 1.0027 |
| 0.5562 | 31.5142 | 9.7735 | 0.3101 | 1.2717 | 4.04E-02 | 2.99956 | 5914 | 0.1961 | 0.9998 | 0.9999 | 0.9999 |
| 0.5672 | 31.4986 | 9.8583 | 0.3130 | 1.2792 | $4.06 \mathrm{E}-02$ | 3.00609 | 5914 | 0.1966 | 1.0026 | 1.0020 | 1.0007 |
| 0.5778 | 31.5002 | 9.8892 | 0.3139 | 1.2857 | $4.08 \mathrm{E}-02$ | 3.00187 | 5914 | 0.1962 | 1.0008 | 1.0006 | 1.0002 |
| 0.5885 | 31.4908 | 9.9681 | 0.3165 | 1.2934 | 4.11E-02 | 3.00714 | 5914 | 0.1967 | 1.0031 | 1.0024 | 1.0008 |
| 0.5993 | 31.4924 | 9.9765 | 0.3168 | 1.3030 | 4.14E-02 | 2.99205 | 5914 | 0.1954 | 0.9966 | 0.9973 | 0.9991 |
| 0.6103 | 31.5033 | 10.0662 | 0.3195 | 1.3089 | 4.15E-02 | 3.00114 | 5914 | 0.1962 | 1.0005 | 1.0004 | 1.0001 |
| 0.6211 | 31.5018 | 10.1421 | 0.3220 | 1.3197 | $4.19 \mathrm{E}-02$ | 3.00059 | 5914 | 0.1961 | 1.0003 | 1.0002 | 1.0001 |
| 0.6316 | 31.4627 | 10.1887 | 0.3238 | 1.3272 | $4.22 \mathrm{E}-02$ | 3.00084 | 5914 | 0.1962 | 1.0004 | 1.0003 | 1.0001 |
| 0.6423 | 31.4939 | 10.2676 | 0.3260 | 1.3375 | $4.25 \mathrm{E}-02$ | 2.99924 | 5914 | 0.1960 | 0.9997 | 0.9997 | 0.9999 |
| 0.6531 | 31.4877 | 10.3294 | 0.3280 | 1.3459 | 4.27E-02 | 2.99796 | 5914 | 0.1959 | 0.9991 | 0.9993 | 0.9998 |
| 0.6642 | 31.4971 | 10.3989 | 0.3302 | 1.3553 | $4.30 \mathrm{E}-02$ | 2.99865 | 5914 | 0.1960 | 0.9994 | 0.9996 | 0.9998 |
| 0.6749 | 31.5018 | 10.4776 | 0.3326 | 1.3670 | 4.34E-02 | 2.99539 | 5914 | 0.1957 | 0.9980 | 0.9985 | 0.9995 |
| 0.6853 | 31.4846 | 10.6162 | 0.3372 | 1.3794 | $4.38 \mathrm{E}-02$ | 3.00636 | 5914 | 0.1966 | 1.0027 | 1.0021 | 1.0008 |
| 0.6961 | 31.4846 | 10.7541 | 0.3416 | 1.3906 | 4.42E-02 | 3.01747 | 5914 | 0.1976 | 1.0075 | 1.0058 | 1.0021 |
| 0.7071 | 31.4752 | 10.8871 | 0.3459 | 1.4107 | $4.48 \mathrm{E}-02$ | 3.013 | 5914 | 0.1972 | 1.0056 | 1.0043 | 1.0015 |
| 0.7178 | 31.4596 | 10.9288 | 0.3474 | 1.4311 | $4.55 \mathrm{E}-02$ | 2.99045 | 5914 | 0.1953 | 0.9959 | 0.9968 | 0.9989 |
| 0.7284 | 31.4861 | 10.9445 | 0.3476 | 1.4463 | 4.59E-02 | 2.96745 | 5914 | 0.1934 | 0.9861 | 0.9892 | 0.9961 |
| 0.7390 | 31.4752 | 10.9382 | 0.3475 | 1.4569 | 4.63E-02 | 2.95144 | 5914 | 0.1920 | 0.9794 | 0.9838 | 0.9941 |
| 0.7498 | 31.4659 | 10.9316 | 0.3474 | 1.4707 | 4.67E-02 | 2.9309 | 5914 | 0.1904 | 0.9707 | 0.9770 | 0.9916 |
| 0.7608 | 31.4799 | 10.9309 | 0.3472 | 1.4836 | 4.71E-02 | 2.91071 | 5914 | 0.1887 | 0.9623 | 0.9702 | 0.9891 |
| 0.7719 | 31.4674 | 10.8789 | 0.3457 | 1.4854 | 4.72E-02 | 2.89899 | 5914 | 0.1877 | 0.9574 | 0.9663 | 0.9876 |
| 0.7822 | 31.4643 | 10.8653 | 0.3453 | 1.4882 | 4.73E-02 | 2.89284 | 5914 | 0.1872 | 0.9549 | 0.9643 | 0.9868 |
| 0.7930 | 31.4721 | 10.8569 | 0.3450 | 1.4911 | $4.74 \mathrm{E}-02$ | 2.88481 | 5914 | 0.1866 | 0.9516 | 0.9616 | 0.9858 |
| 0.8039 | 31.4580 | 10.8199 | 0.3439 | 1.4948 | 4.75E-02 | 2.87429 | 5914 | 0.1857 | 0.9473 | 0.9581 | 0.9844 |
| 0.8149 | 31.4721 | 10.8017 | 0.3432 | 1.4990 | 4.76E-02 | 2.86502 | 5914 | 0.1850 | 0.9435 | 0.9550 | 0.9832 |
| 0.8255 | 31.4534 | 10.7633 | 0.3422 | 1.5058 | 4.79E-02 | 2.8494 | 5914 | 0.1838 | 0.9371 | 0.9498 | 0.9812 |
| 0.8360 | 31.4705 | 10.7670 | 0.3421 | 1.5049 | 4.78E-02 | 2.84947 | 5914 | 0.1838 | 0.9371 | 0.9498 | 0.9812 |
| 0.8468 | 31.4799 | 10.7825 | 0.3425 | 1.5114 | $4.80 \mathrm{E}-02$ | 2.842 | 5914 | 0.1832 | 0.9341 | 0.9473 | 0.9802 |
| 0.8578 | 31.4612 | 10.7789 | 0.3426 | 1.5159 | 4.82E-02 | 2.8377 | 5914 | 0.1828 | 0.9323 | 0.9459 | 0.9796 |
| 0.8687 | 31.4674 | 10.8021 | 0.3433 | 1.5206 | $4.83 \mathrm{E}-02$ | 2.83525 | 5914 | 0.1826 | 0.9313 | 0.9451 | 0.9793 |
| 0.8790 | 31.4612 | 10.8164 | 0.3438 | 1.5231 | $4.84 \mathrm{E}-02$ | 2.83509 | 5914 | 0.1826 | 0.9313 | 0.9450 | 0.9793 |
| 0.8897 | 31.4549 | 10.8419 | 0.3447 | 1.5220 | $4.84 \mathrm{E}-02$ | 2.84156 | 5914 | 0.1831 | 0.9339 | 0.9472 | 0.9801 |
| 0.9007 | 31.4268 | 10.8576 | 0.3455 | 1.5271 | $4.86 \mathrm{E}-02$ | 2.84001 | 5914 | 0.1830 | 0.9333 | 0.9467 | 0.9799 |
| 0.9116 | 31.4456 | 10.8965 | 0.3465 | 1.5330 | 4.87E-02 | 2.83755 | 5914 | 0.1828 | 0.9323 | 0.9459 | 0.9796 |
| 0.9224 | 31.4393 | 10.9241 | 0.3475 | 1.5365 | 4.89E-02 | 2.83978 | 5914 | 0.1830 | 0.9332 | 0.9466 | 0.9799 |

Table C-3: Mach 2.8 ZPG hot-film data - traverse up

| Multiple overheat data |  |  |  |  |  | Single overheat data |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $y$ (in) | $\mathrm{T}_{1} / \mathrm{T}_{11}$ | Re ${ }^{1 / 2}$ | $(\mathrm{pu})_{\text {ms }}$ | $\left(\mathrm{T}_{1}\right)_{\text {mus }}$ | ¢u/( $\mathrm{\rho u})_{\text {ces }}$ | Re ${ }^{1 / 2}$ | (pu) ${ }_{\text {ms }}$ | ¢u $/(\mathrm{\rho u})_{\text {c }}$ | $\mathrm{y} / \delta_{y}$ |
| 0.0231 | 0.9760 | 13.1060 | 0.1594 | 0.0341 | 0.3977 | 12.9210 | 0.1600 | 0.3929 | 0.0582 |
| 0.0446 | 0.9180 | 14.1360 | 0.1680 | 0.1280 | 0.4442 | 13.8270 | 0.1519 | 0.4499 | 0.1123 |
| 0.0666 | 0.9760 | 14.7220 | 0.1363 | 0.0452 | 0.5018 | 14.6330 | 0.1458 | 0.5039 | 0.1678 |
| 0.0880 | 0.9570 | 15.2220 | 0.1671 | 0.1084 | 0.5294 | 15.0680 | 0.1527 | 0.5343 | 0.2217 |
| 0.1101 | 0.9670 | 15.8030 | 0.1537 | 0.0682 | 0.5747 | 15.7840 | 0.1473 | 0.5863 | 0.2773 |
| 0.1312 | 0.9700 | 16.2830 | 0.1622 | 0.0761 | 0.6115 | 15.9440 | 0.1467 | 0.5982 | 0.3305 |
| 0.1532 | 1.0180 | 16.5600 | 0.1714 | 0.0756 | 0.6531 | 16.5610 | 0.1659 | 0.6454 | 0.3859 |
| 0.1751 | 1.0200 | 17.0550 | 0.1528 | 0.0553 | 0.6938 | 17.1090 | 0.1419 | 0.6888 | 0.4411 |
| 0.1970 | 1.0080 | 17.4580 | 0.1444 | 0.0273 | 0.7211 | 17.3870 | 0.1458 | 0.7114 | 0.4962 |
| 0.2192 | 0.9960 | 18.0940 | 0.1470 | 0.0671 | 0.7680 | 18.1360 | 0.1351 | 0.7740 | 0.5521 |
| 0.2409 | 1.0020 | 18.3590 | 0.1236 | 0.0718 | 0.7941 | 18.3900 | 0.1437 | 0.7959 | 0.6068 |
| 0.2632 | 1.0120 | 18.7220 | 0.1506 | 0.0518 | 0.8312 | 18.5280 | 0.1397 | 0.8079 | 0.6630 |
| 0.2846 | 1.0370 | 19.0870 | 0.1203 | 0.0031 | 0.8782 | 18.9200 | 0.1379 | 0.8424 | 0.7169 |
| 0.3067 | 1.0330 | 19.4170 | 0.1182 | 0.0303 | 0.9066 | 19.1160 | 0.1321 | 0.8599 | 0.7725 |
| 0.3286 | 0.9860 | 19.7430 | 0.0987 | 0.0060 | 0.9085 | 19.6440 | 0.1042 | 0.9081 | 0.8277 |
| 0.3508 | 0.9770 | 19.9060 | 0.1087 | 0.0594 | 0.9183 | 19.8200 | 0.1004 | 0.9244 | 0.8836 |
| 0.3729 | 0.9790 | 19.9870 | 0.0959 | 0.0703 | 0.9266 | 20.1770 | 0.0715 | 0.9580 | 0.9393 |
| 0.3947 | 0.9710 | 20.1270 | 0.1010 | 0.0884 | 0.9344 | 19.8410 | 0.0872 | 0.9264 | 0.9942 |
| 0.4169 | 0.9720 | 20.4300 | 0.0629 | 0.0210 | 0.9639 | 20.3860 | 0.0596 | 0.9780 | 1.0501 |
| 0.4387 | 0.9960 | 20.6650 | 0.0618 | 0.0680 | 1.0022 | 20.5250 | 0.0410 | 0.9914 | 1.1050 |
| 0.4609 | 0.9700 | 20.8020 | 0.0719 | 0.0681 | 0.9982 | 20.4900 | 0.0592 | 0.9880 | 1.1610 |
| 0.4735 | 0.9830 | 20.7700 | 0.0052 | 0.0438 | 1.0032 | 20.6290 | 0.0417 | 1.0015 | 1.1927 |
| 0.4740 | 0.9750 | 20.7920 | 0.0682 | 0.0685 | 1.0000 | 20.4990 | 0.0578 | 0.9889 | 1.1940 |

Table C-4: Mach 2.8 ZPG separated turbulence variables single overheat- traverse up

| y (inch) | Mach | ( $\rho$ ) mas | $\mathrm{U}_{\text {mms }}$ | $\rho^{\prime} U^{\prime}$ | u"/u | pbarubar | pubar | $\mathrm{y} / \delta_{u}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.045 | 1.21 | 5.62E-02 | 9.57E-02 | 5.38E-03 | -0.4150 | 39.9 | 68.2 | 0.112 |
| 0.067 | 1.54 | 7.08E-02 | 7.50E-02 | $5.31 \mathrm{E}-03$ | -0.2410 | 58.0 | 76.4 | 0.168 |
| 0.088 | 1.78 | $8.55 \mathrm{E}-02$ | 6.72E-02 | 5.75E-03 | -0.1230 | 71.0 | 81.0 | 0.222 |
| 0.110 | 1.91 | 8.75E-02 | $5.98 \mathrm{E}-02$ | 5.23E-03 | -0.1140 | 78.7 | 88.9 | 0.277 |
| 0.131 | 2.03 | 9.12E-02 | 5.55E-02 | 5.06E-03 | -0.0611 | 85.1 | 90.7 | 0.330 |
| 0.153 | 2.04 | 1.04E-01 | $6.24 \mathrm{E}-02$ | $6.46 \mathrm{E}-03$ | -0.1210 | 86.0 | 97.8 | 0.385 |
| 0.175 | 2.08 | 8.99E-02 | $5.20 \mathrm{E}-02$ | $4.68 \mathrm{E}-03$ | -0.1560 | 88.1 | 104.0 | 0.441 |
| 0.197 | 2.16 | $9.48 \mathrm{E}-02$ | $5.10 \mathrm{E}-02$ | 4.83E-03 | -0.1300 | 93.8 | 108.0 | 0.496 |
| 0.219 | 2.29 | 9.16E-02 | 4.35E-02 | 3.98E-03 | -0.1240 | 103.0 | 117.0 | 0.552 |
| 0.241 | 2.34 | 9.85E-02 | 4.52E-02 | 4.45E-03 | -0.1220 | 106.0 | 121.0 | 0.607 |
| 0.263 | 2.38 | 9.69E-02 | $4.28 \mathrm{E}-02$ | 4.15E-03 | -0.1150 | 108.0 | 122.0 | 0.662 |
| 0.285 | 2.49 | 9.83E-02 | $3.96 \mathrm{E}-02$ | $3.89 \mathrm{E}-03$ | -0.0868 | 117.0 | 128.0 | 0.718 |
| 0.307 | 2.53 | $9.50 \mathrm{E}-02$ | $3.71 \mathrm{E}-02$ | 3.53E-03 | -0.0876 | 119.0 | 130.0 | 0.773 |
| 0.329 | 2.60 | 7.61E-02 | 2.81E-02 | 2.14E-03 | -0.0981 | 124.0 | 138.0 | 0.829 |
| 0.351 | 2.66 | 7.42E-02 | 2.62E-02 | 1.94E-03 | -0.0737 | 130.0 | 140.0 | 0.884 |
| 0.373 | 2.71 | 5.34E-02 | 1.81E-02 | 9.68E-04 | -0.0858 | 133.0 | 145.0 | 0.940 |
| 0.395 | 2.75 | 6.55E-02 | 2.17E-02 | 1.42E-03 | -0.0287 | 136.0 | 140.0 | 0.995 |
| 0.417 | 2.78 | 4.51E-02 | $1.46 \mathrm{E}-02$ | $6.57 \mathrm{E}-04$ | -0.0676 | 138.0 | 148.0 | 1.050 |
| 0.439 | 2.82 | 3.12E-02 | $9.80 \mathrm{E}-03$ | 3.06E-04 | -0.0575 | 142.0 | 150.0 | 1.106 |
| 0.461 | 2.83 | 4.52E-02 | $1.41 \mathrm{E}-02$ | $6.35 \mathrm{E}-04$ | -0.0494 | 142.0 | 150.0 | 1.161 |
| 0.474 | 2.84 | $3.18 \mathrm{E}-02$ | $9.86 \mathrm{E}-03$ | 3.14E-04 | -0.0583 | 143.0 | 152.0 | 1.194 |

Table C-5: Mach 2.8 ZPG hot-film data - traverse down

| Multiple overheat data |  |  |  |  |  | Single overheat data |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $y$ (in) | $\mathrm{T}_{\mathrm{t}} / \mathrm{T}_{\text {t }}$ | $\mathrm{Re}^{1 / 2}$ | ( pu$)_{\text {rms }}$ | $\left(\mathrm{T}_{1}\right)_{\text {ms }}$ | pu/(pu) ${ }_{\text {c }}$ | $\mathrm{Re}^{1 / 2}$ | (pu) ${ }_{\text {rms }}$ | pu /(pu) ${ }_{\text {c }}$ | $y / \delta_{11}$ |
| 0.0407 | 0.9980 | 13.8700 | 0.1569 | 0.0648 | 0.4527 | 13.807 | 0.1506 | 0.4493 | 0.1025 |
| 0.0629 | 1.0020 | 14.6160 | 0.1394 | 0.0048 | 0.5040 | 14.548 | 0.1430 | 0.4989 | 0.1584 |
| 0.0844 | 1.0030 | 15.2100 | 0.1565 | 0.0686 | 0.5463 | 15.166 | 0.1484 | 0.5421 | 0.2126 |
| 0.1063 | 0.9660 | 15.5460 | 0.1510 | 0.0278 | 0.5568 | 15.445 | 0.1554 | 0.5623 | 0.2678 |
| 0.1281 | 0.9700 | 15.9580 | 0.1480 | 0.0389 | 0.5880 | 15.980 | 0.1481 | 0.6019 | 0.3227 |
| 0.1501 | 1.0030 | 16.5310 | 0.1808 | 0.1130 | 0.6454 | 16.398 | 0.1693 | 0.6338 | 0.3781 |
| 0.1722 | 0.9790 | 16.9310 | 0.1539 | 0.0745 | 0.6659 | 16.426 | 0.1500 | 0.6360 | 0.4338 |
| 0.1939 | 1.0090 | 17.4400 | 0.1423 | 0.0539 | 0.7213 | 17.604 | 0.1441 | 0.7305 | 0.4884 |
| 0.2163 | 1.0000 | 17.8130 | 0.1569 | 0.0639 | 0.7477 | 17.979 | 0.1451 | 0.7619 | 0.5448 |
| 0.2379 | 1.0080 | 18.3170 | 0.1190 | 0.0373 | 0.7951 | 18.276 | 0.1294 | 0.7873 | 0.5992 |
| 0.2598 | 1.0110 | 18.5820 | 0.0788 | 0.1201 | 0.8199 | 18.863 | 0.1334 | 0.8387 | 0.6544 |
| 0.2813 | 1.0320 | 19.0190 | 0.1210 | 0.0380 | 0.8708 | 18.955 | 0.1383 | 0.8469 | 0.7086 |
| 0.3036 | 1.0070 | 19.3840 | 0.1296 | 0.0797 | 0.8895 | 19.395 | 0.1099 | 0.8867 | 0.7647 |
| 0.3256 | 0.9940 | 19.4780 | 0.1410 | 0.0967 | 0.8906 | 19.267 | 0.1202 | 0.8750 | 0.8202 |
| 0.3478 | 1.0240 | 19.8270 | 0.0653 | 0.0442 | 0.9414 | 20.024 | 0.0866 | 0.9451 | 0.8761 |
| 0.3700 | 0.9880 | 19.9670 | 0.0755 | 0.0407 | 0.9323 | 19.998 | 0.0796 | 0.9426 | 0.9320 |
| 0.3917 | 0.9660 | 20.1350 | 0.0897 | 0.0513 | 0.9336 | 19.946 | 0.0823 | 0.9377 | 0.9866 |
| 0.4141 | 0.9730 | 20.2960 | 0.0427 | 0.0668 | 0.9533 | 20.283 | 0.0549 | 0.9697 | 1.0431 |
| 0.4358 | 0.9890 | 20.4620 | 0.0127 | 0.0578 | 0.9795 | 20.342 | 0.0605 | 0.9754 | 1.0977 |
| 0.4582 | 0.9970 | 20.6940 | 0.0695 | 0.0722 | 1.0076 | 20.589 | 0.0499 | 0.9992 | 1.1542 |
| 0.4733 | 0.9760 | 20.7680 | 0.0429 | 0.0281 | 1.0000 | 20.663 | 0.0397 | 1.0064 | 1.1922 |

Table C-6: Mach 2.8 ZPG separated turbulence variables single overheat - traverse down

| y (inch) | Mach | ( $\rho$ ) ${ }_{\text {mos }}$ | $\mathrm{U}_{\mathrm{rms}}$ | م'U | u"/u | pbarubar | pubar | $\mathrm{y} / \delta_{u}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.063 | 1.48 | $6.68 \mathrm{E}-02$ | 7.62E-02 | 5.09E-03 | -0.2660 | 55.4 | 75.5 | 0.158 |
| 0.084 | 1.76 | 8.22E-02 | 6.62E-02 | 5.44E-03 | -0.1480 | 69.9 | 82.0 | 0.213 |
| 0.106 | 1.89 | 9.16E-02 | 6.38E-02 | 5.84E-03 | -0.0913 | 77.3 | 85.1 | 0.267 |
| 0.128 | 2.00 | 9.11E-02 | 5.70E-02 | 5.19E-03 | -0.0770 | 84.1 | 91.1 | 0.322 |
| 0.150 | 2.03 | 1.05E-01 | $6.38 \mathrm{E}-02$ | 6.73E-03 | -0.1100 | 85.3 | 95.9 | 0.378 |
| 0.172 | 2.10 | 9.56E-02 | $5.44 \mathrm{E}-02$ | $5.20 \mathrm{E}-03$ | -0.0667 | 89.8 | 96.2 | 0.433 |
| 0.194 | 2.15 | 9.36E-02 | 5.05E-02 | 4.73E-03 | -0.1580 | 93.1 | 111.0 | 0.489 |
| 0.216 | 2.29 | 9.82E-02 | 4.69E-02 | $4.60 \mathrm{E}-03$ | -0.1090 | 103.0 | 115.0 | 0.544 |
| 0.238 | 2.33 | 8.86E-02 | $4.08 \mathrm{E}-02$ | 3.61E-03 | -0.1180 | 105.0 | 119.0 | 0.599 |
| 0.260 | 2.37 | 9.23E-02 | 4.11E-02 | 3.79E-03 | -0.1470 | 108.0 | 127.0 | 0.655 |
| 0.281 | 2.48 | 9.83E-02 | 4.00E-02 | 3.93E-03 | -0.1010 | 115.0 | 128.0 | 0.708 |
| 0.304 | 2.53 | 7.89E-02 | 3.10E-02 | $2.44 \mathrm{E}-03$ | -0.1160 | 119.0 | 134.0 | 0.766 |
| 0.326 | 2.60 | 8.77E-02 | 3.25E-02 | $2.85 \mathrm{E}-03$ | -0.0602 | 124.0 | 132.0 | 0.821 |
| 0.348 | 2.66 | $6.39 \mathrm{E}-02$ | 2.27E-02 | 1.45E-03 | -0.1010 | 129.0 | 143.0 | 0.877 |
| 0.370 | 2.71 | 5.94E-02 | 2.02E-02 | 1.20E-03 | -0.0662 | 133.0 | 143.0 | 0.932 |
| 0.392 | 2.74 | 6.17E-02 | 2.06E-02 | 1.27E-03 | -0.0483 | 135.0 | 142.0 | 0.987 |
| 0.414 | 2.77 | 4.14E-02 | 1.35E-02 | 5.58E-04 | -0.0593 | 138.0 | 147.0 | 1.043 |
| 0.436 | 2.82 | 4.60E-02 | 1.45E-02 | $6.67 \mathrm{E}-04$ | -0.0454 | 141.0 | 148.0 | 1.098 |
| 0.458 | 2.83 | 3.81E-02 | 1.18E-02 | 4.51E-04 | -0.0553 | 143.0 | 151.0 | 1.154 |
| 0.473 | 2,84 | 3.03E-02 | 9.40E-03 | 2.85E-04 | -0.0633 | 143.0 | 152.0 | 1.191 |

Table C-7: Mach 2.8 FPG hot-film data - traverse up

| Multiple overheat data |  |  |  |  |  | Single overheat data |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $y$ (in) | $\mathrm{T}_{1} / \mathrm{T}_{\mathrm{tl}}$ | $\mathrm{Re}^{1 / 2}$ | (pu) mim | $\left(\mathrm{T}_{\mathrm{t}}\right)_{\text {mas }}$ | $\mathrm{pu} /(\mathrm{pu})_{x}$ | $\mathrm{Re}^{1 / 2}$ | (pu) ${ }_{\text {rms }}$ | $\mathrm{pu} /(\mathrm{pu})_{x}$ | $\mathrm{y} / \delta_{\mathrm{u}}$ |
| 0.0072 | 1.074 | 8.653 | 8.58E-02 | 5.18E-02 | 0.1974 | 8.835 | 6.55E-02 | 0.1963 | 0.0173 |
| 0.0287 | 1.055 | 10.392 | $4.76 \mathrm{E}-02$ | $9.78 \mathrm{E}-02$ | 0.2814 | 10.582 | 1.30E-01 | 0.2816 | 0.0692 |
| 0.0506 | 1.009 | 11.988 | $7.63 \mathrm{E}-02$ | 6.93E-02 | 0.3635 | 12.066 | 1.01E-01 | 0.3661 | 0.1219 |
| 0.0726 | 0.979 | 12.682 | $9.46 \mathrm{E}-02$ | 1.94E-02 | 0.3986 | 12.848 | 9.43E-02 | 0.4151 | 0.1749 |
| 0.0946 | 0.981 | 13.26 | 9.93E-02 | 1.45E-02 | 0.4365 | 13.461 | 1.01E-01 | 0.4556 | 0.2280 |
| 0.1163 | 0.975 | 13.696 | $8.60 \mathrm{E}-02$ | 6.03E-02 | 0.4636 | 13.739 | 9.89E-02 | 0.4746 | 0.2802 |
| 0.1384 | 0.999 | 14.036 | 1.23E-01 | 6.97E-02 | 0.4950 | 14.151 | 1.07E-01 | 0.5035 | 0.3335 |
| 0.1602 | 0.999 | 14.402 | $9.32 \mathrm{E}-02$ | 5.08E-02 | 0.5213 | 14.634 | 1.03E-01 | 0.5385 | 0.3860 |
| 0.1823 | 0.973 | 14.693 | 1.04E-01 | 3.26E-02 | 0.5331 | 14.647 | 1.09E-01 | 0.5394 | 0.4393 |
| 0.2041 | 0.995 | 14.954 | 1.09E-01 | 4.78E-02 | 0.5605 | 15.108 | 1.09E-01 | 0.5739 | 0.4918 |
| 0.2262 | 1.007 | 15.192 | 1.19E-01 | 4.56E-02 | 0.5828 | 15.378 | 1.12E-01 | 0.5946 | 0.5451 |
| 0.2483 | 1.008 | 15.562 | 7.33E-02 | 7.57E-02 | 0.6122 | 15.75 | 1.02E-01 | 0.6237 | 0.5983 |
| 0.2702 | 1.007 | 15.783 | $1.04 \mathrm{E}-01$ | $5.01 \mathrm{E}-02$ | 0.6292 | 15.969 | 1.17E-01 | 0.6412 | 0.6511 |
| 0.2924 | 1.025 | 16.226 | 1.04E-01 | 3.87E-02 | 0.6731 | 16.469 | 1.10E-01 | 0.6820 | 0.7046 |
| 0.3143 | 1.042 | 16.621 | 9.54E-02 | 5.31E-02 | 0.7140 | 16.66 | 1.09E-01 | 0.6979 | 0.7573 |
| 0.3366 | 1.023 | 16.603 | 1.08E-01 | 2.57E-02 | 0.7036 | 16.797 | 1.08E-01 | 0.7094 | 0.8111 |
| 0.3586 | 0.968 | 16.86 | $9.62 \mathrm{E}-02$ | 7.35E-02 | 0.6995 | 16.882 | 1.07E-01 | 0.7166 | 0.8641 |
| 0.381 | 1.011 | 17.207 | 1.20E-01 | 5.74E-02 | 0.7497 | 17.27 | 1.15E-01 | 0.7499 | 0.9181 |
| 0.4029 | 0.966 | 17.467 | $9.20 \mathrm{E}-02$ | $6.09 \mathrm{E}-02$ | 0.7499 | 17.567 | 1.08E-01 | 0.7760 | 0.9708 |
| 0.4249 | 0.999 | 17.738 | 7.84E-02 | $6.48 \mathrm{E}-02$ | 0.7904 | 17.75 | 1.02E-01 | 0.7922 | 1.0239 |
| 0.4463 | 0.997 | 17.966 | 7.01E-02 | 5.08E-02 | 0.8101 | 18.114 | 7.96E-02 | 0.8250 | 1.0754 |
| 0.4682 | 1.015 | 18.099 | 6.81E-02 | 7.54E-03 | 0.8319 | 18.24 | 7.58E-02 | 0.8366 | 1.1282 |
| 0.4904 | 0.97 | 18.093 | 7.73E-02 | 3.40E-02 | 0.8067 | 18.004 | 8.52E-02 | 0.8150 | 1.1817 |
| 0.5122 | 0.983 | 18.311 | 8.54E-02 | 6.53E-02 | 0.8333 | 18.309 | 7.55E-02 | 0.8429 | 1.2342 |
| 0.5345 | 0.979 | 18.407 | $6.79 \mathrm{E}-02$ | $3.34 \mathrm{E}-02$ | 0.8401 | 18.345 | 6.71E-02 | 0.8462 | 1.2880 |
| 0.5562 | 1.001 | 18.443 | 2.74E-02 | 4.67E-02 | 0.8561 | 18.661 | $5.08 \mathrm{E}-02$ | 0.8756 | 1.3402 |
| 0.5784 | 0.97 | 18.556 | $5.90 \mathrm{E}-02$ | 8.83E-03 | 0.8481 | 18.582 | 5.33E-02 | 0.8682 | 1.3937 |
| 0.6003 | 0.979 | 18.641 | $5.38 \mathrm{E}-02$ | $1.03 \mathrm{E}-02$ | 0.8612 | 18.586 | 6.36E-02 | 0.8686 | 1.4465 |
| 0.6224 | 0.993 | 18.816 | 1.14E-02 | 3.53E-02 | 0.8860 | 19.013 | 3.57E-02 | 0.9090 | 1.4998 |
| 0.6444 | 0.989 | 18.888 | 3.54E-02 | 1.89E-02 | 0.8905 | 19.044 | 2.76E-02 | 0.9119 | 1.5528 |
| 0.6662 | 1.008 | 19.031 | $2.64 \mathrm{E}-02$ | 2.06E-02 | 0.9152 | 19.312 | $3.24 \mathrm{E}-02$ | 0.9378 | 1.6053 |
| 0.6884 | 0.997 | 19.24 | 1.99E-02 | 1.70E-02 | 0.9291 | 19.39 | 3.00E-02 | 0.9454 | 1.6588 |
| 0.7101 | 0.996 | 19.41 | 2.57E-02 | 2.41E-02 | 0.9450 | 19.532 | 2.51E-02 | 0.9593 | 1.7111 |
| 0.7323 | 0.997 | 19.509 | 2.43E-02 | 3.22E-02 | 0.9551 | 19.707 | $1.78 \mathrm{E}-02$ | 0.9765 | 1.7646 |
| 0.7539 | 0.994 | 19.518 | 1.55E-02 | 7.36E-03 | 0.9542 | 19.646 | $1.76 \mathrm{E}-02$ | 0.9705 | 1.8166 |
| 0.776 | 0.998 | 19.548 | $9.34 \mathrm{E}-03$ | 1.06E-02 | 0.9593 | 19.765 | 1.55E-02 | 0.9823 | 1.8699 |
| 0.7979 | 0.995 | 19.589 | $9.64 \mathrm{E}-03$ | 2.10E-02 | 0.9618 | 19.772 | 1.50E-02 | 0.9830 | 1.9227 |
| 0.8198 | 0.998 | 19.63 | 1.06E-02 | $1.72 \mathrm{E}-02$ | 0.9676 | 19.826 | 1.37E-02 | 0.9884 | 1.9754 |
| 0.8418 | 0.995 | 19.682 | 6.71E-03 | 1.20E-02 | 0.9707 | 19.861 | 1.42E-02 | 0.9919 | 2.0284 |
| 0.8635 | 0.994 | 19.742 | 9.18E-03 | 1.03E-02 | 0.9762 | 19.915 | 1.45E-02 | 0.9973 | 2.0807 |
| 0.8848 | 0.995 | 19.797 | 8.25E-03 | 1.31E-02 | 0.9821 | 19.942 | 1.47E-02 | 1.0000 | 2.1320 |
| 0.9059 | 0.994 | 19.832 | 7.37E-03 | 1.19E-02 | 0.9847 | 20.004 | 1.41E-02 | 1.0062 | 2.1829 |
| 0.928 | 0.996 | 19.872 | 9.22E-03 | 1.07E-02 | 0.9900 | 20.058 | 1.42E-02 | 1.0116 | 2.2361 |
| 0.9497 | 0.995 | 19.893 | 8.57E-03 | 1.05E-02 | 0.9919 | 20.09 | 1.38E-02 | 1.0149 | 2.2884 |
| 0.9717 | 0.993 | 19.921 | $7.29 \mathrm{E}-03$ | 1.22E-02 | 0.9929 | 20.075 | 1.38E-02 | 1.0133 | 2.3414 |
| 0.9908 | 0.994 | 19.943 | $9.57 \mathrm{E}-03$ | 1.02E-02 | 0.9962 | 20.109 | 1.46E-02 | 1.0168 | 2.3875 |
| 0.9939 | 0.998 | 19.959 | 1.25E-02 | 1.88E-03 | 1.0001 | 20.149 | 1.56E-02 | 1.0208 | 2.3949 |
| 0.9939 | 0.997 | 19.965 | 1.05E-02 | 8.53E-03 | 1.0000 | 20.178 | 1.42E-02 | 1.0238 | 2.3949 |

Table C-8: Mach 2.8 FPG separated turbulence variables single overheat - traverse up

| y (inch) | Mach | (0) $\mathrm{mms}^{\text {cmin }}$ | $\mathrm{U}_{\text {mms }}$ | $\rho^{\prime} \mathrm{U}^{\prime}$ | u"/u | pbarubar | pubar | $y / \delta_{u}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.029 | 1.46 | 5.98E-02 | $6.97 \mathrm{E}-02$ | 4.17E-03 | 0.0479 | 41.8 | 39.9 | 0.069 |
| 0.051 | 1.47 | 4.69E-02 | $5.45 \mathrm{E}-02$ | 2.56E-03 | -0.1970 | 41.7 | 51.9 | 0.122 |
| 0.073 | 1.49 | 4.45E-02 | 4.98E-02 | 2.22E-03 | -0.2740 | 42.8 | 58.9 | 0.175 |
| 0.095 | 1.84 | $5.80 \mathrm{E}-02$ | 4.29E-02 | 2.49E-03 | -0.1650 | 54.0 | 64.6 | 0.228 |
| 0.116 | 2.01 | $6.12 \mathrm{E}-02$ | 3.77E-02 | $2.31 \mathrm{E}-03$ | -0.1380 | 58.1 | 67.3 | 0.280 |
| 0.138 | 2.12 | $6.87 \mathrm{E}-02$ | $3.84 \mathrm{E}-02$ | 2.64E-03 | -0.1470 | 60.9 | 71.4 | 0.333 |
| 0.160 | 2.22 | $6.84 \mathrm{E}-02$ | 3.45E-02 | $2.36 \mathrm{E}-03$ | -0.1560 | 64.5 | 76.4 | 0.386 |
| 0.182 | 2.33 | $7.44 \mathrm{E}-02$ | 3.43E-02 | 2.55E-03 | -0.1120 | 68.0 | 76.5 | 0.439 |
| 0.204 | 2.40 | 7.59E-02 | 3.30E-02 | 2.51E-03 | -0.1270 | 71.1 | 81.4 | 0.492 |
| 0.226 | 2.46 | 7.91E-02 | 3.26E-02 | $2.58 \mathrm{E}-03$ | -0.1200 | 74.2 | 84.3 | 0.545 |
| 0.248 | 2.54 | 7.36E-02 | 2.84E-02 | 2.09E-03 | -0.1160 | 78.2 | 88.5 | 0.598 |
| 0.270 | 2.60 | 8.57E-02 | 3.17E-02 | 2.72E-03 | -0.1040 | 81.5 | 90.9 | 0.651 |
| 0.292 | 2.62 | 8.10E-02 | 2.94E-02 | $2.38 \mathrm{E}-03$ | -0.1340 | 83.8 | 96.7 | 0.704 |
| 0.314 | 2.70 | 8.11E-02 | 2.77E-02 | 2.25E-03 | -0.1100 | 88.1 | 99.0 | 0.757 |
| 0.337 | 2.77 | 8.17E-02 | $2.66 \mathrm{E}-02$ | $2.17 \mathrm{E}-03$ | -0.0862 | 92.0 | 101.0 | 0.812 |
| 0.359 | 2.81 | 8.14E-02 | 2.58E-02 | $2.10 \mathrm{E}-03$ | -0.0707 | 94.5 | 102.0 | 0.865 |
| 0.381 | 2.80 | $8.70 \mathrm{E}-02$ | 2.78E-02 | $2.42 \mathrm{E}-03$ | -0.1030 | 95.4 | 106.0 | 0.918 |
| 0.403 | 2.89 | $8.27 \mathrm{E}-02$ | $2.48 \mathrm{E}-02$ | $2.05 \mathrm{E}-03$ | -0.0918 | 100.0 | 110.0 | 0.971 |
| 0.425 | 2.91 | $7.88 \mathrm{E}-02$ | $2.33 \mathrm{E}-02$ | $1.84 \mathrm{E}-03$ | -0.0945 | 102.0 | 112.0 | 1.024 |
| 0.446 | 2.94 | 6.17E-02 | 1.78E-02 | 1.10E-03 | -0.1110 | 104.0 | 117.0 | 1.075 |
| 0.468 | 3.00 | 5.93E-02 | 1.65E-02 | 9.77E-04 | -0.0899 | 108.0 | 119.0 | 1.128 |
| 0.490 | 3.01 | $6.68 \mathrm{E}-02$ | 1.84E-02 | 1.23E-03 | -0.0507 | 110.0 | 116.0 | 1.181 |
| 0.512 | 3.00 | 5.91E-02 | 1.64E-02 | $9.68 \mathrm{E}-04$ | -0.0707 | 111.0 | 120.0 | 1.234 |
| 0.535 | 3.01 | $5.26 \mathrm{E}-02$ | 1.45E-02 | 7.62E-04 | -0.0583 | 113.0 | 120.0 | 1.289 |
| 0.556 | 3.00 | 3.98E-02 | 1.11E-02 | 4.40E-04 | -0.0799 | 114.0 | 124.0 | 1.340 |
| 0.578 | 3.00 | 4.17E-02 | 1.16E-02 | 4.83E-04 | -0.0607 | 116.0 | 123.0 | 1.393 |
| 0.600 | 2.99 | 4.98E-02 | 1.39E-02 | 6.91E-04 | -0.0514 | 117.0 | 123.0 | 1.446 |
| 0.622 | 3.00 | 2.79E-02 | 7.76E-03 | 2.17E-04 | -0.0793 | 119.0 | 129.0 | 1.499 |
| 0.644 | 3.00 | 2.16E-02 | 6.00E-03 | 1.29E-04 | -0.0703 | 120.0 | 129.0 | 1.552 |
| 0.666 | 3.00 | $2.54 \mathrm{E}-02$ | $7.06 \mathrm{E}-03$ | $1.79 \mathrm{E}-04$ | -0.0838 | 122.0 | 133.0 | 1.605 |
| 0.688 | 3.01 | $2.35 \mathrm{E}-02$ | $6.48 \mathrm{E}-03$ | 1.52E-04 | -0.0709 | 125.0 | 134.0 | 1.658 |
| 0.710 | 3.01 | $1.96 \mathrm{E}-02$ | 5.43E-03 | 1.07E-04 | -0.0635 | 127.0 | 136.0 | 1.711 |
| 0.732 | 2.96 | 1.38E-02 | 3.95E-03 | $5.46 \mathrm{E}-05$ | -0.0726 | 128.0 | 139.0 | 1.764 |
| 0.754 | 2.92 | $1.36 \mathrm{E}-02$ | 3.98E-03 | 5.41E-05 | -0.0640 | 129.0 | 138.0 | 1.817 |
| 0.776 | 2.90 | 1.19E-02 | 3.55E-03 | 4.23E-05 | -0.0776 | 129.0 | 139.0 | 1.870 |
| 0.798 | 2.88 | 1.15E-02 | 3.48E-03 | 4.01E-05 | -0.0801 | 128.0 | 139.0 | 1.923 |
| 0.820 | 2.86 | $1.05 \mathrm{E}-02$ | 3.22E-03 | $3.39 \mathrm{E}-05$ | -0.0874 | 128.0 | 140.0 | 1.976 |

Table C-9: Mach 2.8 CPG ( $\mathrm{x}=68 \mathrm{~cm}$ ) hot-film data - traverse up

| Multiple overheat data |  |  |  |  |  | Single overheat data |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| y (in) | T/ / $\mathrm{T}_{11}$ | Re ${ }^{1 / 2}$ | ( P$)_{\text {mus }}$ | $\left(\mathrm{T}_{\mathrm{t}}\right)_{\text {ms }}$ | ¢u/( pu$)_{\text {x }}$ | $\mathrm{Re}^{1 / 2}$ | (pu) ${ }_{\text {ms }}$ | $\mathrm{pu} /(\mathrm{pu})_{\text {c }}$ | $\mathrm{y} / \delta_{1}$ |
| 0.0126 | 1.0090 | 14.9490 | 0.1239 | 0.0115 | 0.3990 | 15.1240 | 0.1494 | 0.4046 | 0.0352 |
| 0.0347 | 0.9680 | 16.6310 | 0.1648 | 0.0654 | 0.4803 | 16.9100 | 0.1796 | 0.5058 | 0.0970 |
| 0.0562 | 0.9930 | 17.6210 | 0.1812 | 0.0173 | 0.5484 | 17.7430 | 0.1858 | 0.5568 | 0.1570 |
| 0.0786 | 1.0350 | 18.6590 | 0.1477 | 0.0992 | 0.6320 | 18.7750 | 0.1886 | 0.6235 | 0.2196 |
| 0.1001 | 0.9810 | 19.3830 | 0.1756 | 0.0663 | 0.6579 | 19.5860 | 0.1903 | 0.6785 | 0.2797 |
| 0.1224 | 1.0250 | 20.2060 | 0.1935 | 0.0313 | 0.7362 | 20.4510 | 0.1964 | 0.7398 | 0.3420 |
| 0.1441 | 1.0150 | 20.6300 | 0.1750 | 0.0434 | 0.7629 | 21.3920 | 0.1956 | 0.8094 | 0.4027 |
| 0.1662 | 1.0530 | 21.9350 | 0.1230 | 0.0962 | 0.8837 | 22.7740 | 0.1711 | 0.9174 | 0.4644 |
| 0.1883 | 1.0240 | 22.4460 | 0.1483 | 0.0732 | 0.9083 | 22.8300 | 0.1661 | 0.9219 | 0.5262 |
| 0.2101 | 1.0270 | 23.3320 | 0.1320 | 0.0754 | 0.9832 | 23.7560 | 0.1664 | 0.9982 | 0.5871 |
| 0.2325 | 0.9550 | 24.0080 | 0.1414 | 0.0889 | 0.9920 | 23.7770 | 0.1502 | 1.0000 | 0.6497 |
| 0.2542 | 0.9630 | 24.6430 | 0.1623 | 0.1075 | 1.0509 | 24.8680 | 0.1384 | 1.0938 | 0.7103 |
| 0.2767 | 0.9940 | 25.1460 | 0.1053 | 0.0592 | 1.1177 | 25.4250 | 0.1162 | 1.1434 | 0.7732 |
| 0.2985 | 0.9930 | 25.4950 | 0.1032 | 0.0414 | 1.1479 | 25.7520 | 0.1036 | 1.1730 | 0.8341 |
| 0.3210 | 1.0020 | 25.7990 | 0.0571 | 0.0632 | 1.1824 | 25.9480 | 0.0819 | 1.1909 | 0.8970 |
| 0.3427 | 0.9970 | 25.8650 | 0.0406 | 0.0579 | 1.1847 | 26.1910 | 0.0523 | 1.2133 | 0.9576 |
| 0.3649 | 0.9880 | 25.8990 | 0.0770 | 0.0575 | 1.1808 | 25.8960 | 0.0585 | 1.1861 | 1.0196 |
| 0.3867 | 0.9990 | 25.7610 | 0.0156 | 0.0116 | 1.1771 | 25.8740 | 0.0299 | 1.1841 | 1.0805 |
| 0.4081 | 0.9890 | 25.3720 | 0.0290 | 0.0265 | 1.1338 | 25.4850 | 0.0373 | 1.1488 | 1.1403 |
| 0.4303 | 0.9880 | 24.9570 | 0.0078 | 0.0395 | 1.0960 | 24.9400 | 0.0318 | 1.1002 | 1.2024 |
| 0.4513 | 0.9950 | 24.4940 | 0.0357 | 0.0030 | 1.0608 | 24.6580 | 0.0389 | 1.0754 | 1.2611 |
| 0.4738 | 0.9840 | 24.1380 | 0.0385 | 0.0381 | 1.0226 | 24.2430 | 0.0273 | 1.0395 | 1.3239 |
| 0.4908 | 0.9790 | 23.9080 | 0.0166 | 0.0080 | 1.0000 | 23.9570 | 0.0258 | 1.0152 | 1.3714 |

Table C-10: Mach 2.8 CPG ( $\mathrm{x}=68 \mathrm{~cm}$ ) separated turbulence variables single overheat- traverse up

| y (inch) | Mach | (p) ${ }_{\text {mis }}$ | $\mathrm{U}_{\text {ms }}$ | $\rho^{\prime} U$ | u"/u | pbarubar | pubar | $\mathrm{y} / \delta_{u}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.035 | 1.54 | 8.76E-02 | $9.20 \mathrm{E}-02$ | 8.06E-03 | -0.0109 | 101.0 | 102.0 | 0.097 |
| 0.056 | 1.77 | 1.03E-01 | $8.27 \mathrm{E}-02$ | 8.53E-03 | 0.0265 | 115.0 | 112.0 | 0.157 |
| 0.079 | 1.87 | 1.10E-01 | 7.84E-02 | $8.64 \mathrm{E}-03$ | -0.0283 | 122.0 | 126.0 | 0.220 |
| 0.100 | 1.97 | 1.16E-01 | 7.45E-02 | 8.63E-03 | -0.0739 | 127.0 | 137.0 | 0.279 |
| 0.122 | 2.09 | 1.25E-01 | 7.15E-02 | 8.93E-03 | -0.1130 | 132.0 | 149.0 | 0.341 |
| 0.144 | 2.20 | 1.29E-01 | $6.68 \mathrm{E}-02$ | 8.60E-03 | -0.1400 | 140.0 | 163.0 | 0.402 |
| 0.166 | 2.19 | 1.13E-01 | $5.85 \mathrm{E}-02$ | 6.59E-03 | -0.2030 | 147.0 | 185.0 | 0.464 |
| 0.188 | 2.35 | 1.14E-01 | 5.17E-02 | 5.91E-03 | -0.1740 | 154.0 | 186.0 | 0.525 |
| 0.210 | 2.38 | 1.16E-01 | $5.08 \mathrm{E}-02$ | 5.87E-03 | -0.2290 | 155.0 | 201.0 | 0.587 |
| 0.233 | 2.44 | 1.06E-01 | 4.43E-02 | 4.69E-03 | -0.2050 | 160.0 | 202.0 | 0.651 |
| 0.254 | 2.49 | 9.87E-02 | $3.97 \mathrm{E}-02$ | 3.92E-03 | -0.2410 | 167.0 | 221.0 | 0.710 |
| 0.277 | 2.50 | $8.30 \mathrm{E}-02$ | $3.32 \mathrm{E}-02$ | 2.76E-03 | -0.2510 | 173.0 | 231.0 | 0.774 |
| 0.299 | 2.53 | $7.46 \mathrm{E}-02$ | $2.90 \mathrm{E}-02$ | 2.17E-03 | -0.2560 | 176.0 | 237.0 | 0.835 |
| 0.321 | 2.54 | 5.90E-02 | 2.29E-02 | 1.35E-03 | -0.2600 | 178.0 | 240.0 | 0.897 |
| 0.343 | 2.53 | 3.76E-02 | 1.47E-02 | 5.53E-04 | -0.2620 | 181.0 | 245.0 | 0.958 |
| 0.365 | 2.57 | 4.24E-02 | 1.61E-02 | 6.82E-04 | -0.2360 | 183.0 | 239.0 | 1.020 |
| 0.387 | 2.60 | $2.18 \mathrm{E}-02$ | $8.06 \mathrm{E}-03$ | 1.76E-04 | -0.2200 | 186.0 | 239.0 | 1.081 |
| 0.408 | 2.61 | 2.73E-02 | $9.99 \mathrm{E}-03$ | 2.73E-04 | -0.1950 | 186.0 | 232.0 | 1.140 |
| 0.430 | 2.60 | 2.32E-02 | 8.61E-03 | $2.00 \mathrm{E}-04$ | -0.1520 | 188.0 | 222.0 | 1.202 |
| 0.451 | 2.58 | 2.83E-02 | $1.06 \mathrm{E}-02$ | $3.00 \mathrm{E}-04$ | -0.1290 | 189.0 | 217.0 | 1.260 |
| 0.474 | 2.59 | 1.99E-02 | $7.40 \mathrm{E}-03$ | 1.47E-04 | -0.0935 | 190.0 | 210.0 | 1.324 |
| 0.491 | 2.59 | $1.88 \mathrm{E}-02$ | 7.01E-03 | 1.32E-04 | -0.0729 | 190.0 | 205.0 | 1.372 |

Table C-11: Mach 2.8 CPG ( $\mathrm{x}=68 \mathrm{~cm}$ ) hot-film data - traverse down

| Multiple overheat data |  |  |  |  |  | Single overheat data |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $y$ y(in) | $\mathrm{T}_{\mathrm{t}} / \mathrm{T}_{11}$ | $\mathrm{Re}^{1 / 2}$ | ( $\rho \mathrm{u})_{\text {mos }}$ | $\left(\mathrm{T}_{2}\right)_{\text {mms }}$ | $\mathrm{pu} /(\mathrm{pu})_{x}$ | $\mathrm{Re}^{1 / 2}$ | ( pu$)_{\text {mos }}$ | $\rho \mathrm{u} /(\mathrm{pu})_{\infty}$ | $\mathrm{y} / \delta_{u}$ |
| 0.0084 | 1.028 | 14.309 | 8.94E-02 | 8.16E-02 | 0.3655 | 14.35 | 1.23E-01 | 0.3610 | 0.0235 |
| 0.0307 | 0.996 | 16.305 | 1.56E-01 | 6.69E-02 | 0.4646 | 16.454 | 1.77E-01 | 0.4746 | 0.0858 |
| 0.0521 | 1.004 | 17.133 | 1.94E-01 | 7.23E-02 | 0.5158 | 17.318 | 1.95E-01 | 0.5258 | 0.1456 |
| 0.0744 | 0.975 | 18.172 | 1.92E-01 | 5.09E-02 | 0.5690 | 18.247 | 1.97E-01 | 0.5837 | 0.2079 |
| 0.0961 | 1.023 | 19.332 | 1.86E-01 | 4.69E-02 | 0.6652 | 19.417 | 2.07E-01 | 0.6610 | 0.2685 |
| 0.1181 | 0.961 | 19.802 | 2.00E-01 | 5.54E-02 | 0.6694 | 19.615 | $2.00 \mathrm{E}-01$ | 0.6745 | 0.3300 |
| 0.1402 | 0.992 | 20.633 | 1.81E-01 | 2.66E-02 | 0.7424 | 20.781 | 1.84E-01 | 0.7571 | 0.3918 |
| 0.162 | 1.011 | 21.534 | 1.27E-01 | 1.10E-01 | 0.8190 | 21.753 | 1.74E-01 | 0.8296 | 0.4527 |
| 0.1844 | 1.04 | 22.052 | 1.91E-01 | $7.28 \mathrm{E}-02$ | 0.8751 | 22.479 | 1.86E-01 | 0.8859 | 0.5153 |
| 0.2059 | 1.012 | 22.931 | 1.69E-01 | $5.49 \mathrm{E}-02$ | 0.9289 | 23.039 | 1.73E-01 | 0.9305 | 0.5753 |
| 0.2284 | 1.025 | 23.688 | 1.43E-01 | $6.44 \mathrm{E}-02$ | 1.0001 | 23.71 | 1.64E-01 | 0.9855 | 0.6382 |
| 0.2501 | 0.974 | 24.373 | 1.37E-01 | 9.16E-02 | 1.0232 | 24.307 | 1.50E-01 | 1.0358 | 0.6988 |
| 0.2725 | 0.989 | 24.746 | 1.33E-01 | 5.61E-02 | 1.0653 | 24.849 | 1.42E-01 | 1.0825 | 0.7614 |
| 0.2945 | 0.981 | 25.068 | 1.15E-01 | 6.93E-02 | 1.0874 | 24.801 | 1.37E-01 | 1.0783 | 0.8229 |
| 0.3167 | 1.006 | 25.663 | 5.50E-02 | 1.07E-01 | 1.1588 | 25.825 | 8.71E-02 | 1.1692 | 0.8849 |
| 0.3389 | 0.992 | 25.634 | 6.53E-02 | 4.97E-02 | 1.1460 | 25.412 | 9.52E-02 | 1.1321 | 0.9470 |
| 0.3606 | 0.995 | 25.925 | 4.40E-02 | 3.52E-02 | 1.1740 | 25.952 | $6.25 \mathrm{E}-02$ | 1.1807 | 1.0076 |
| 0.383 | 0.99 | 25.845 | 1.56E-02 | $4.80 \mathrm{E}-02$ | 1.1630 | 25.912 | 3.42E-02 | 1.1771 | 1.0702 |
| 0.4046 | 0.99 | 25.522 | 2.44E-02 | 1.02E-02 | 1.1344 | 25.502 | $3.20 \mathrm{E}-02$ | 1.1401 | 1.1306 |
| 0.4271 | 0.983 | 25.15 | $4.88 \mathrm{E}-02$ | 4.94E-02 | 1.0960 | 25.097 | $3.04 \mathrm{E}-02$ | 1.1042 | 1.1934 |
| 0.4487 | 0.976 | 24.683 | 1.88E-02 | $3.69 \mathrm{E}-02$ | 1.0511 | 24.537 | 3.41E-02 | 1.0555 | 1.2538 |
| 0.471 | 0.98 | 24.25 | 3.38E-02 | 2.56E-02 | 1.0174 | 24.172 | 2.83E-02 | 1.0243 | 1.3161 |
| 0.4899 | 0.984 | 24.009 | 2.62E-02 | 1.86E-03 | 1.0000 | 24.075 | 2.57E-02 | 1.0161 | 1.3689 |

Table C-12: Mach 2.8 CPG ( $\mathrm{x}=68 \mathrm{~cm}$ ) separated turbulence variables single overheat- traverse down

| $y$ (inch) | Mach | $(\rho)_{\text {ms }}$ | $\mathrm{U}_{\text {rms }}$ | $\rho^{\prime} \mathrm{U}$ | u"/u | pbarubar | pubar | $\mathrm{y} / \delta_{u}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.031 | 1.51 | 8.41E-02 | 9.27E-02 | 7.80E-03 | 0.0236 | 98.8 | 96.6 | 0.086 |
| 0.052 | 1.73 | 1.06E-01 | 8.84E-02 | 9.39E-03 | 0.0516 | 112.0 | 107.0 | 0.146 |
| 0.074 | 1.85 | 1.14E-01 | 8.34E-02 | 9.47E-03 | 0.0175 | 121.0 | 119.0 | 0.208 |
| 0.096 | 1.96 | 1.26E-01 | 8.16E-02 | 1.02E-02 | -0.0630 | 126.0 | 134.0 | 0.269 |
| 0.118 | 2.08 | 1.27E-01 | 7.34E-02 | 9.32E-03 | -0.0421 | 131.0 | 137.0 | 0.330 |
| 0.140 | 2.18 | $1.21 \mathrm{E}-01$ | 6.33E-02 | $7.64 \mathrm{E}-03$ | -0.0944 | 139.0 | 154.0 | 0.391 |
| 0.162 | 2.24 | 1.16E-01 | $5.79 \mathrm{E}-02$ | $6.71 \mathrm{E}-03$ | -0.1420 | 145.0 | 169.0 | 0.453 |
| 0.184 | 2.34 | $1.28 \mathrm{E}-01$ | $5.80 \mathrm{E}-02$ | $7.40 \mathrm{E}-03$ | -0.1510 | 153.0 | 180.0 | 0.514 |
| 0.206 | 2.36 | 1.19E-01 | 5.35E-02 | $6.38 \mathrm{E}-03$ | -0.1890 | 153.0 | 189.0 | 0.576 |
| 0.228 | 2.44 | 1.16E-01 | 4.85E-02 | 5.62E-03 | -0.2020 | 160.0 | 200.0 | 0.637 |
| 0.250 | 2.48 | 1.06E-01 | $4.31 \mathrm{E}-02$ | $4.59 \mathrm{E}-03$ | -0.2080 | 167.0 | 211.0 | 0.699 |
| 0.273 | 2.50 | 1.01E-01 | 4.04E-02 | 4.10E-03 | -0.2180 | 172.0 | 220.0 | 0.763 |
| 0.294 | 2.54 | 9.87E-02 | 3.84E-02 | $3.79 \mathrm{E}-03$ | -0.1960 | 176.0 | 219.0 | 0.822 |
| 0.317 | 2.54 | $6.28 \mathrm{E}-02$ | 2.43E-02 | 1.53E-03 | -0.2560 | 177.0 | 238.0 | 0.886 |
| 0.339 | 2.53 | $6.85 \mathrm{E}-02$ | 2.67E-02 | 1.83E-03 | -0.2190 | 180.0 | 230.0 | 0.947 |
| 0.361 | 2.56 | 4.52E-02 | 1.73E-02 | 7.81E-04 | -0.2430 | 182.0 | 240.0 | 1.009 |
| 0.383 | 2.59 | 2.49E-02 | 9.28E-03 | 2.31E-04 | -0.2260 | 185.0 | 239.0 | 1.070 |
| 0.405 | 2.61 | 2.35E-02 | 8.58E-03 | 2.01E-04 | -0.1960 | 186.0 | 232.0 | 1.132 |
| 0.427 | 2.59 | 2.21E-02 | 8.22E-03 | 1.82E-04 | -0.1630 | 188.0 | 225.0 | 1.193 |
| 0.449 | 2.58 | 2.48E-02 | 9.31E-03 | 2.31E-04 | -0.1220 | 189.0 | 215.0 | 1.255 |
| 0.471 | 2.59 | 2.07E-02 | 7.68E-03 | 1.59E-04 | -0.0885 | 190.0 | 208.0 | 1.316 |
| 0.490 | 2.59 | 1.87E-02 | 6.97E-03 | 1.30E-04 | -0.0819 | 190.0 | 207.0 | 1.369 |

Table C-13: Mach 2.8 CPG ( $\mathrm{x}=70 \mathrm{~cm}$ ) hot-film data - traverse up

| Multiple overheat data |  |  |  |  |  | Single overheat data |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $y$ (in) | $\mathrm{T}_{\mathrm{t}} / \mathrm{T}_{\mathrm{tl}}$ | $\mathrm{Re}^{1 / 2}$ | (pu) ${ }_{\text {rms }}$ | $\left(\mathrm{T}_{4}\right)_{\text {mms }}$ | $\mathrm{pu} /(\mathrm{pu})_{\text {c }}$ | $\mathrm{Re}^{\text {T/2 }}$ | (pu) ${ }_{\text {mas }}$ | pu $/(\mathrm{pu})_{\text {ex }}$ | $\mathrm{y} / \delta_{\mathrm{u}}$ |
| 0.0127 | 1.0220 | 15.2550 | 0.1302 | 0.0377 | 0.3591 | 15.3110 | 0.1427 | 0.3567 | 0.0431 |
| 0.0348 | 1.0370 | 17.1360 | 0.1826 | 0.0472 | 0.4577 | 17.4130 | 0.1937 | 0.4613 | 0.1182 |
| 0.0565 | 1.0430 | 18.0530 | 0.1852 | 0.0634 | 0.5099 | 17.8690 | 0.1963 | 0.4858 | 0.1919 |
| 0.0790 | 1.0490 | 18.9200 | 0.1937 | 0.0622 | 0.5623 | 19.1090 | 0.2052 | 0.5555 | 0.2683 |
| 0.1007 | 0.9720 | 19.3670 | 0.1923 | 0.0759 | 0.5600 | 19.1260 | 0.2026 | 0.5565 | 0.3419 |
| 0.1232 | 0.9740 | 19.8640 | 0.2242 | 0.0644 | 0.5898 | 19.3240 | 0.2193 | 0.5681 | 0.4184 |
| 0.1448 | 0.9750 | 20.6440 | 0.1862 | 0.1485 | 0.6373 | 20.5790 | 0.1992 | 0.6443 | 0.4917 |
| 0.1670 | 1.0000 | 21.5470 | 0.2048 | 0.0563 | 0.7065 | 21.7090 | 0.2168 | 0.7170 | 0.5671 |
| 0.1889 | 1.0590 | 22.2840 | 0.1284 | 0.1227 | 0.7850 | 22.5540 | 0.1879 | 0.7739 | 0.6415 |
| 0.2108 | 0.9820 | 22.7990 | 0.1499 | 0.1441 | 0.7809 | 22.5420 | 0.1921 | 0.7731 | 0.7158 |
| 0.2330 | 0.9430 | 23.5590 | 0.1955 | 0.0973 | 0.8118 | 22.0720 | 0.1921 | 0.7412 | 0.7912 |
| 0.2545 | 1.0280 | 24.2820 | 0.1461 | 0.0367 | 0.9134 | 23.9720 | 0.1622 | 0.8743 | 0.8642 |
| 0.2770 | 1.0380 | 24.8180 | 0.0560 | 0.1064 | 0.9603 | 24.7690 | 0.1382 | 0.9334 | 0.9406 |
| 0.2984 | 0.9930 | 25.1110 | 0.1602 | 0.1132 | 0.9548 | 24.7780 | 0.1306 | 0.9341 | 1.0133 |
| 0.3208 | 1.0170 | 25.5180 | 0.0600 | 0.0721 | 1.0021 | 25.8960 | 0.0710 | 1.0203 | 1.0893 |
| 0.3424 | 1.0030 | 25.9500 | 0.0425 | 0.0611 | 1.0266 | 25.9710 | 0.0578 | 1.0262 | 1.1627 |
| 0.3645 | 1.0020 | 26.0150 | 0.0187 | 0.0640 | 1.0311 | 25.9150 | 0.0570 | 1.0218 | 1.2377 |
| 0.3858 | 1.0100 | 26.0170 | 0.0431 | 0.0593 | 1.0366 | 26.0040 | 0.0334 | 1.0288 | 1.3101 |
| 0.4075 | 1.0020 | 26.0270 | 0.0207 | 0.0360 | 1.0315 | 25.9530 | 0.0173 | 1.0248 | 1.3838 |
| 0.4297 | 1.0000 | 25.9290 | 0.0152 | 0.0240 | 1.0229 | 25.8060 | 0.0153 | 1.0132 | 1.4591 |
| 0.4512 | 1.0030 | 25.8370 | 0.0156 | 0.0133 | 1.0174 | 25.7280 | 0.0152 | 1.0071 | 1.5321 |
| 0.4735 | 1.0020 | 25.7410 | 0.0215 | 0.0140 | 1.0092 | 25.6180 | 0.0170 | 0.9985 | 1.6079 |
| 0.4918 | 0.9980 | 25.6580 | 0.0229 | 0.0194 | 1.0000 | 25.4950 | 0.0199 | 0.9889 | 1.6700 |

Table C-14: Mach $2.8 \mathrm{CPG}(\mathrm{x}=70 \mathrm{~cm})$ separated turbulence variables single overheat- traverse up

| $y$ (inch) | Mach | (p) $)_{\text {ras }}$ | $\mathrm{U}_{\text {mms }}$ | pU | u"/u | pbarubar | pubar | $\mathrm{y} / \delta_{\mathrm{u}}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.035 | 1.54 | $9.46 \mathrm{E}-02$ | 9.91E-02 | 9.37E-03 | -0.0667 | 101.0 | 108.0 | 0.118 |
| 0.057 | 1.77 | $1.09 \mathrm{E}-01$ | 8.72E-02 | 9.51E-03 | 0.0134 | 115.0 | 114.0 | 0.191 |
| 0.079 | 1.88 | $1.20 \mathrm{E}-01$ | 8.52E-02 | 1.02E-02 | -0.0610 | 122.0 | 130.0 | 0.268 |
| 0.101 | 1.97 | 1.23E-01 | 7.92E-02 | $9.78 \mathrm{E}-03$ | -0.0280 | 127.0 | 130.0 | 0.342 |
| 0.123 | 2.09 | 1.40E-01 | 7.97E-02 | 1.11E-02 | -0.0055 | 132.0 | 133.0 | 0.417 |
| 0.145 | 2.20 | 1.31E-01 | $6.79 \mathrm{E}-02$ | 8.92E-03 | -0.0697 | 141.0 | 151.0 | 0.491 |
| 0.167 | 2.19 | 1.42E-01 | $7.45 \mathrm{E}-02$ | 1.06E-02 | -0.1200 | 148.0 | 168.0 | 0.566 |
| 0.189 | 2.35 | 1.29E-01 | 5.84E-02 | 7.57E-03 | -0.1530 | 154.0 | 181.0 | 0.640 |
| 0.211 | 2.39 | 1.34E-01 | 5.85E-02 | 7.82E-03 | -0.1420 | 155.0 | 181.0 | 0.715 |
| 0.233 | 2.45 | 1.35E-01 | $5.66 \mathrm{E}-02$ | 7.67E-03 | -0.0777 | 160.0 | 174.0 | 0.789 |
| 0.255 | 2.49 | 1.16E-01 | 4.66E-02 | 5.38E-03 | -0.1830 | 167.0 | 205.0 | 0.864 |
| 0.277 | 2.50 | 9.87E-02 | 3.95E-02 | 3.90E-03 | -0.2110 | 173.0 | 219.0 | 0.938 |
| 0.298 | 2.53 | $9.40 \mathrm{E}-02$ | $3.66 \mathrm{E}-02$ | 3.44E-03 | -0.1960 | 176.0 | 219.0 | 1.009 |
| 0.321 | 2.54 | 5.12E-02 | 1.98E-02 | 1.02E-03 | -0.2570 | 178.0 | 239.0 | 1.087 |
| 0.342 | 2.53 | 4.16E-02 | 1.62E-02 | 6.75E-04 | -0.2500 | 181.0 | 241.0 | 1.158 |
| 0.364 | 2.57 | $4.13 \mathrm{E}-02$ | 1.57E-02 | $6.47 \mathrm{E}-04$ | -0.2370 | 183.0 | 240.0 | 1.233 |
| 0.386 | 2.60 | $2.44 \mathrm{E}-02$ | 9.04E-03 | 2.20E-04 | -0.2290 | 186.0 | 241.0 | 1.307 |
| 0.407 | 2.61 | $1.26 \mathrm{E}-02$ | 4.63E-03 | 5.85E-05 | -0.2240 | 186.0 | 240.0 | 1.378 |
| 0.430 | 2.60 | 1.11E-02 | 4.14E-03 | 4.61E-05 | -0.2080 | 188.0 | 238.0 | 1.456 |
| 0.451 | 2.58 | 1.11E-02 | 4.15E-03 | $4.59 \mathrm{E}-05$ | -0.2000 | 189.0 | 236.0 | 1.527 |
| 0.474 | 2.59 | $1.24 \mathrm{E}-02$ | 4.61E-03 | 5.70E-05 | -0.1880 | 190.0 | 234.0 | 1.605 |
| 0.492 | 2.59 | 1.45E-02 | $5.40 \mathrm{E}-03$ | 7.83E-05 | -0.1810 | 190.0 | 232.0 | 1.666 |

Table C-15: Mach 2.8 CPG ( $\mathrm{x}=70 \mathrm{~cm}$ ) hot-film data - traverse down

| Multiple overheat data |  |  |  |  |  | Single overheat data |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $y$ (in) | $\mathrm{T}_{1} / \mathrm{T}_{1}$ | $\mathrm{Re}^{\text {4/2}}$ | $(\mathrm{pu})_{\text {rms }}$ | ( $\left.\mathrm{T}_{1}\right)_{\text {mm }}$ | pu /( $\rho \mathrm{u})_{\text {c }}$ | $\mathrm{Re}^{1 / 2}$ | (pu) ${ }_{\text {mus }}$ |  | $\mathrm{y} / \delta_{\mathrm{u}}$ |
| 0.0085 | 1.0450 | 14.7050 | 0.1279 | 0.0459 | 0.3354 | 15.1180 | 0.1335 | 0.3442 | 0.0289 |
| 0.0308 | 1.0260 | 16.6440 | 0.1852 | 0.0239 | 0.4243 | 16.8680 | 0.1924 | 0.4285 | 0.1046 |
| 0.0524 | 0.9970 | 17.7130 | 0.1903 | 0.0657 | 0.4715 | 17.5570 | 0.1999 | 0.4642 | 0.1779 |
| 0.0748 | 1.0320 | 18.5150 | 0.2043 | 0.0197 | 0.5273 | 18.6260 | 0.2083 | 0.5224 | 0.2540 |
| 0.0965 | 1.0090 | 19.0500 | 0.1936 | 0.0741 | 0.5498 | 19.0410 | 0.2049 | 0.5460 | 0.3277 |
| 0.1188 | 1.0350 | 19.9100 | 0.1887 | 0.0677 | 0.6105 | 20.7410 | 0.1925 | 0.6478 | 0.4034 |
| 0.1408 | 1.0280 | 20.7000 | 0.2134 | 0.0606 | 0.6572 | 20.9290 | 0.2132 | 0.6596 | 0.4781 |
| 0.1626 | 1.0150 | 21.2930 | 0.1870 | 0.1022 | 0.6897 | 21.0930 | 0.2098 | 0.6700 | 0.5521 |
| 0.1848 | 0.9550 | 21.8960 | 0.1889 | 0.0606 | 0.7003 | 21.5020 | 0.1931 | 0.6962 | 0.6275 |
| 0.2064 | 0.9800 | 22.7090 | 0.1987 | 0.0707 | 0.7660 | 22.2080 | 0.1995 | 0.7427 | 0.7009 |
| 0.2288 | 1.0020 | 23.1020 | 0.1575 | 0.1007 | 0.8046 | 23.1370 | 0.1852 | 0.8062 | 0.7769 |
| 0.2502 | 1.0140 | 24.1080 | 0.1333 | 0.0873 | 0.8835 | 24.4080 | 0.1631 | 0.8972 | 0.8496 |
| 0.2726 | 0.9820 | 24.4180 | 0.1849 | 0.1180 | 0.8869 | 24.1910 | 0.1627 | 0.8813 | 0.9257 |
| 0.2944 | 1.0370 | 25.1970 | 0.1213 | 0.0571 | 0.9795 | 25.2850 | 0.1176 | 0.9628 | 0.9997 |
| 0.3164 | 1.0010 | 25.3840 | 0.1150 | 0.0324 | 0.9710 | 25.0190 | 0.1302 | 0.9426 | 1.0744 |
| 0.3384 | 1.0020 | 25.6290 | 0.0299 | 0.1124 | 0.9901 | 25.5070 | 0.0945 | 0.9798 | 1.1491 |
| 0.3601 | 0.9970 | 25.8900 | 0.0302 | 0.0778 | 1.0074 | 25.9810 | 0.0647 | 1.0165 | 1.2228 |
| 0.3823 | 1.0030 | 25.9060 | 0.0083 | 0.0651 | 1.0125 | 25.6920 | 0.0752 | 0.9940 | 1.2982 |
| 0.4038 | 0.9890 | 26.0530 | 0.0624 | 0.0403 | 1.0147 | 25.9840 | 0.0436 | 1.0168 | 1.3712 |
| 0.4263 | 0.9940 | 26.0220 | 0.0620 | 0.0554 | 1.0158 | 25.9810 | 0.0279 | 1.0165 | 1.4476 |
| 0.4477 | 1.0040 | 25.9770 | 0.0318 | 0.0403 | 1.0185 | 25.9450 | 0.0179 | 1.0137 | 1.5203 |
| 0.4698 | 1.0000 | 25.8800 | 0.0134 | 0.0105 | 1.0082 | 25.8240 | 0.0202 | 1.0043 | 1.5953 |
| 0.4900 | 0.9960 | 25.8070 | 0.0249 | 0.0185 | 1.0000 | 25.7210 | 0.0229 | 0.9963 | 1.6639 |

Table C-16: Mach 2.8 CPG ( $\mathrm{x}=70 \mathrm{~cm}$ ) separated turbulence variables single overheat- traverse down

| y (inch) | Mach | (p) ${ }_{\text {rms }}$ | $\mathrm{U}_{\text {rms }}$ | $\rho^{\prime} \mathrm{U}$ | 4"/u | pbarubar | pubar | $\mathrm{y} / \delta_{u}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.031 | 1.51 | 9.16E-02 | 1.01E-01 | 9.23E-03 | -0.0255 | 98.9 | 101.0 | 0.104 |
| 0.052 | 1.74 | 1.09E-01 | 9.06E-02 | 9.90E-03 | 0.0265 | 113.0 | 110.0 | 0.177 |
| 0.075 | 1.85 | $1.20 \mathrm{E}-01$ | $8.80 \mathrm{E}-02$ | 1.06E-02 | -0.0225 | 121.0 | 124.0 | 0.253 |
| 0.097 | 1.96 | 1.24E-01 | $8.07 \mathrm{E}-02$ | $1.00 \mathrm{E}-02$ | -0.0251 | 126.0 | 129.0 | 0.327 |
| 0.119 | 2.08 | 1.22E-01 | 7.05E-02 | 8.60E-03 | -0.1420 | 132.0 | 153.0 | 0.403 |
| 0.141 | 2.19 | $1.40 \mathrm{E}-01$ | 7.33E-02 | 1.03E-02 | -0.1060 | 140.0 | 156.0 | 0.478 |
| 0.163 | 2.23 | $1.40 \mathrm{E}-01$ | $7.01 \mathrm{E}-02$ | $9.80 \mathrm{E}-03$ | -0.0849 | 145.0 | 159.0 | 0.552 |
| 0.185 | 2.35 | 1.33E-01 | $6.03 \mathrm{E}-02$ | 8.01E-03 | -0.0714 | 153.0 | 165.0 | 0.627 |
| 0.206 | 2.36 | 1.38E-01 | 6.17E-02 | $8.50 \mathrm{E}-03$ | -0.1270 | 154.0 | 176.0 | 0.698 |
| 0.229 | 2.44 | $1.30 \mathrm{E}-01$ | 5.47E-02 | 7.14E-03 | -0.1620 | 160.0 | 191.0 | 0.776 |
| 0.250 | 2.48 | 1.16E-01 | $4.70 \mathrm{E}-02$ | 5.46E-03 | -0.2140 | 167.0 | 212.0 | 0.847 |
| 0.273 | 2.50 | $1.16 \mathrm{E}-01$ | $4.64 \mathrm{E}-02$ | $5.40 \mathrm{E}-03$ | -0.1750 | 172.0 | 209.0 | 0.925 |
| 0.294 | 2.54 | $8.47 \mathrm{E}-02$ | 3.29E-02 | $2.79 \mathrm{E}-03$ | -0.2270 | 176.0 | 228.0 | 0.996 |
| 0.316 | 2.54 | $9.39 \mathrm{E}-02$ | $3.63 \mathrm{E}-02$ | $3.41 \mathrm{E}-03$ | -0.2080 | 177.0 | 223.0 | 1.070 |
| 0.338 | 2.53 | $6.80 \mathrm{E}-02$ | 2.65E-02 | $1.80 \mathrm{E}-03$ | -0.2250 | 180.0 | 232.0 | 1.145 |
| 0.360 | 2.56 | $4.68 \mathrm{E}-02$ | $1.79 \mathrm{E}-02$ | 8.37E-04 | -0.2450 | 182.0 | 241.0 | 1.219 |
| 0.382 | 2.59 | $5.48 \mathrm{E}-02$ | 2.04E-02 | 1.12E-03 | -0.2130 | 185.0 | 235.0 | 1.294 |
| 0.404 | 2.61 | $3.19 \mathrm{E}-02$ | 1.17E-02 | 3.73E-04 | -0.2260 | 186.0 | 241.0 | 1.368 |
| 0.426 | 2.59 | 2.03E-02 | 7.56E-03 | 1.54E-04 | -0.2190 | 188.0 | 241.0 | 1.443 |
| 0.448 | 2.58 | $1.30 \mathrm{E}-02$ | 4.88E-03 | 6.33E-05 | -0.2150 | 188.0 | 240.0 | 1.517 |
| 0.470 | 2.59 | $1.47 \mathrm{E}-02$ | 5.47E-03 | 8.04E-05 | -0.2020 | 190.0 | 238.0 | 1.592 |
| 0.490 | 2.59 | $1.67 \mathrm{E}-02$ | 6.21E-03 | 1.04E-04 | -0.1950 | 190.0 | 236.0 | 1.659 |

Table C-17: Mach 2.8 discrete data turbulence intensity points

| $y$ (in) | a | b | sqrt(Re) | f | $\mathrm{V}_{\mathrm{ms}}$ | $\mathrm{V}_{\text {ber }}$ | T.I. | $y / \delta_{u}$ FPG |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.278 | 0.027806 | -0.14622 | 18.064 | 0.352665 | 0.0607 | 1.83 | 9.405\% | 0.670 |
| 0.446 | 0.027806 | -0.14622 | 20.152 | 0.338272 | 0.0517 | 1.996 | 7.657\% | 1.075 |
| 0.111 | 0.027806 | -0.14622 | 15.622 | 0.376857 | 0.0506 | 1.64 | 8.187\% | 0.267 |
|  |  |  | fbar | 0.355931 |  |  |  |  |
|  |  |  |  |  |  |  |  |  |
| y (in) | a | b | sqrt(Re) | f | $\mathrm{V}_{\text {rms }}$ | $\mathrm{V}_{\text {bar }}$ | T.I. | $\mathrm{y} / \delta_{\mathrm{u}} \mathrm{ZPG}$ |
| 0.03125 | 0.027897 | -0.10393 | 17.33 | 0.318456 | 0.08133 | 1.69 | 15.112\% | 0.078696 |
| 0.13125 | 0.027897 | -0.10393 | 20.3 | 0.30619 | 0.09089 | 1.857 | 15.985\% | 0.330521 |
| 0.23125 | 0.027897 | -0.10393 | 21.27 | 0.303083 | 0.09676 | 1.991 | 16.035\% | 0.582347 |
| 0.33125 | 0.027897 | -0.10393 | 21.35 | 0.302842 | 0.09012 | 2.119 | 14.043\% | 0.834173 |
|  |  |  | fbar | 0.307643 |  |  |  |  |
|  |  |  |  |  |  |  |  |  |
| $y$ (in) | a | b | sqrt(Re) | $f$ | $\mathrm{V}_{\mathrm{nms}}$ | $\mathrm{V}_{\text {bar }}$ | T.I. | $\mathrm{y} / \delta_{11} \mathrm{CPG}$ |
| 0.075 | 0.02387 | -0.12132 | 15.819 | 0.368353 | 0.0984 | 1.749 | 15.274\% | 0.209556 |
| 0.185 | 0.02387 | -0.12132 | 19.674 | 0.337084 | 0.1155 | 2.02 | 16.963\% | 0.516904 |
| 0.298 | 0.02387 | -0.12132 | 23.043 | 0.320749 | 0.0916 | 2.23 | 12.806\% | 0.832635 |
|  |  |  | fbar | 0.342062 |  |  |  |  |

Table C-18: Mach 2.8 hot-film van Driest data

| ZPG |  |  | FPG |  |  | CPG ( $x=68$ | $68 \mathrm{~cm})$ |  | CPG ( $\mathrm{x}=$ | $70 \mathrm{~cm})$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $y+$ | ueff+ |  | y+ | Ueff+ |  | y+ | ueff+ |  | y+ |  |
| 126.6 | 12.2 |  | 103.5 | 16.9 |  | 140.4 | 17.1 |  | 132.2 | 19.7 |
| 189.1 | 15.8 |  | 182.5 | 16.6 |  | 227.3 | 19.5 |  | 214.7 | 22.0 |
| 249.8 | 17.2 |  | 261.9 | 16.5 |  | 318.0 | 21.4 |  | 300.1 | 23.2 |
| 312.6 | 18.3 |  | 341.2 | 19.3 |  | 404.9 | 20.8 |  | 382.6 | 21.8 |
| 372.5 | 19.2 |  | 419.5 | 20.4 |  | 495.1 | 22.8 |  | 468.1 | 22.8 |
| 434.9 | 20.4 |  | 499.2 | 21.5 |  | 582.9 | 23.3 |  | 550.1 | 23.5 |
| 497.1 | 20.7 |  | 577.9 | 22.1 |  | 672.3 | 24.4 |  | 634.5 | 24.2 |
| 559.3 | 20.9 |  | 657.6 | 22.4 |  | 761.7 | 24.6 |  | 717.7 | 27.3 |
| 622.3 | 21.4 |  | 736.2 | 23.1 |  | 849.9 | 24.9 |  | 800.9 | 25.0 |
| 683.9 | 21.8 |  | 815.9 | 23.6 |  | 940.5 | 23.0 |  | 885.2 | 24.1 |
| 747.2 | 22.3 |  | 895.6 | 24.1 |  | 1028.3 | 23.5 |  | 966.9 | 27.2 |
| 808.0 | 23.7 |  | 974.6 | 24.4 |  | 11193 | 24.5 |  | 1052.4 | 27.6 |
| 870.7 | 23.8 |  | 1054.7 | 24.8 |  | 1207.5 | 24.7 |  | 1133.7 | 26.2 |
| 932.9 | 22.7 |  | 1133.7 | 25.5 |  | 1298.5 | 25.0 |  | 1218.8 | 27.1 |
| 995.9 | 22.8 |  | 1214.1 | 25.5 |  | 1386.3 | 24.8 |  | 1300.8 | 26.6 |
| 1058.7 | 23.0 |  | 1293.5 | 24.8 |  | 1476.1 | 24.7 |  | 1384.8 | 26.8 |
| 1120.5 | 23.0 |  | 1374.3 | 25.5 |  | 1564.3 | 25.3 |  | 1465.7 | 27.2 |
| 1183.6 | 23.1 |  | 1453.3 | 25.1 |  | 1650.9 | 25.0 |  | 1548.1 | 27.0 |
| 1245.5 | 24.1 |  | 1532.6 | 25.8 |  | 1740.7 | 24.9 |  | 1632.5 | 26.9 |
| 1308.5 | 23.3 |  | 1609.8 | 25.9 |  | 1825.6 | 25.0 |  | 1714.2 | 26.9 |
| 1344.3 | 23.7 |  | 1688.8 | 26.5 |  | 1916.6 | 24.7 |  | 1798.9 | 26.9 |
|  |  |  | 1768.9 | 25.7 |  | 1985.4 | 24.6 |  | 1868.4 | 26.7 |
|  |  |  | 1847.5 | 25.9 |  |  |  |  |  |  |
|  |  |  | 1928.0 | 25.9 |  |  |  |  |  |  |
|  |  |  | 2006.3 | 26.2 |  |  |  |  |  |  |
|  |  |  | 2086.3 | 25.7 |  |  |  |  |  |  |
|  |  |  | 2165.3 | 25.8 |  |  |  |  |  |  |
|  |  |  | 2245.0 | 26.1 |  |  |  |  |  |  |
|  |  |  | 2324.4 | 26.0 |  |  |  |  |  |  |
|  |  |  | 2403.0 | 26.3 |  |  |  |  |  |  |
|  |  |  | 2483.1 | 26.2 |  |  |  |  |  |  |
|  |  |  | 2561.4 | 26.2 |  |  |  |  |  |  |
|  |  |  | 2641.5 | 26.0 |  |  |  |  |  |  |
|  |  |  | 2719.4 | 25.8 |  |  |  |  |  |  |
|  |  |  | 2799.1 | 25.7 |  |  |  |  |  |  |

Table C-19: Mach 2.8 power spectra data

| frequency | $\begin{aligned} & \text { ZPG, y/d = } \\ & 0.33 \end{aligned}$ | $\begin{aligned} & \mathrm{ZPG}, \mathrm{y} / \mathrm{d}= \\ & 0.58 \end{aligned}$ | $\left\lvert\, \begin{aligned} & \text { ZPG, } \mathrm{y} / \mathrm{d}= \\ & 0.83 \end{aligned}\right.$ | $\begin{aligned} & \text { FPG, y/d = } \\ & 0.27 \end{aligned}$ | $\begin{aligned} & \text { FPG, } y / d= \\ & 0.67 \end{aligned}$ | $\begin{aligned} & \text { CPG, } \mathrm{y} / \mathrm{d}= \\ & 0.21 \end{aligned}$ | $\begin{aligned} & \text { CPG, y/d = } \\ & 0.52 \end{aligned}$ | $\begin{aligned} & \text { CPG, } y / d= \\ & 0.83 \end{aligned}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 488.3 | 6.03E-06 | 4.97E-06 | 3E-06 | 1.34E-05 | 2.47E-05 | 8.09E-05 | 0.000158 | 5.4E-05 |
| 976.6 | 3.4E-06 | 2.51E-06 | 1.48E-06 | 5.31E-06 | 3.34E-06 | 2.97E-05 | 3.05E-05 | 1.29E-05 |
| 1465 | 4.3E-06 | 2.6E-06 | $1.58 \mathrm{E}-06$ | 3.65E-06 | 2.33E-06 | 2.66E-05 | 1.62E-05 | 5.03E-06 |
| 1953 | 2.48E-06 | 1.92E-06 | 1.32E-06 | 2.7E-06 | 1.48E-06 | 1.94E-05 | 1.14E-05 | 3.48E-06 |
| 2441 | 2.15E-06 | 1.33E-06 | 1.13E-06 | 1.93E-06 | $8.64 \mathrm{E}-07$ | 1.6E-05 | 7.43E-06 | 2.2E-06 |
| 2930 | 2.4E-06 | 1.42E-06 | $9.74 \mathrm{E}-07$ | 1.77E-06 | 7.95E-07 | 1.2E-05 | 6.13E-06 | 1.99E-06 |
| 3418 | 2.35E-06 | $1.45 \mathrm{E}-06$ | 9.83E-07 | 1.6E-06 | 1.13E-06 | 1.03E-05 | 4.47E-06 | 1.7E-06 |
| 3906 | $2.58 \mathrm{E}-06$ | 1.39E-06 | 9.3E-07 | 1.24E-06 | 7.12E-07 | 8.64E-06 | 4.42E-06 | 1.55E-06 |
| 4395 | 2.15E-06 | 1.05E-06 | 8.45E-07 | 1.14E-06 | 8.12E-07 | 9.03E-06 | 2.73E-06 | 1.43E-06 |
| 4883 | 1.64E-06 | 1.32E-06 | 7.75E-07 | 1.04E-06 | 6.68E-07 | 7.24E-06 | 3.43E-06 | 1.48E-06 |
| 5371 | 1.89E-06 | $1.34 \mathrm{E}-06$ | 7.44E-07 | 9.93E-07 | $7.28 \mathrm{E}-07$ | 5.55E-06 | 2.79E-06 | 1.19E-06 |
| 5859 | 1.75E-06 | 1.19E-06 | 8.65E-07 | 9.39E-07 | 6.91E-07 | 4.39E-06 | 2.41E-06 | 1.11E-06 |
| 6348 | 1.48E-06 | 1.3E-06 | 7.08E-07 | 1.04E-06 | $9.13 \mathrm{E}-07$ | 4.5E-06 | 2.53E-06 | 1.21E-06 |
| 6836 | 1.35E-06 | 1.31E-06 | 5.83E-07 | 6.5E-07 | $8.56 \mathrm{E}-07$ | 3.84E-06 | $2.38 \mathrm{E}-06$ | 1.29E-06 |
| 7324 | 1.51E-06 | 1.05E-06 | 4.8E-07 | 7.27E-07 | 7.67E-07 | 2.64E-06 | 2.15E-06 | 1.23E-06 |
| 7813 | 1.7E-06 | 1.25E-06 | 6.97E-07 | 5.56E-07 | 7.22E-07 | 2.61E-06 | 2.2E-06 | 1.09E-06 |
| 8301 | 1.33E-06 | 1.1E-06 | 8.63E-07 | 5.74E-07 | 6.07E-07 | 2.14E-06 | 2.25E-06 | 1.08E-06 |
| 8789 | 1.26E-06 | 1.2E-06 | 7.52E-07 | 5.41E-07 | $6.28 \mathrm{E}-07$ | 1.82E-06 | 1.8E-06 | 1.26E-06 |
| 9277 | 1.12E-06 | 1.02E-06 | 6.71E-07 | 4.66E-07 | 7.76E-07 | 1.68E-06 | 1.66E-06 | 1.32E-06 |
| 9766 | 1.32E-06 | $1.29 \mathrm{E}-06$ | 6.48E-07 | 4.73E-07 | 6.05E-07 | 1.63E-06 | $1.66 \mathrm{E}-06$ | $9.34 \mathrm{E}-07$ |
| 10250 | $1.48 \mathrm{E}-06$ | 1.15E-06 | 7.2E-07 | 4.27E-07 | 6.5E-07 | 1.44E-06 | 1.49E-06 | 1.28E-06 |
| 10740 | $1.44 \mathrm{E}-06$ | 9.29E-07 | 7.77E-07 | 6.43E-07 | 7.52E-07 | 1.21E-06 | 1.76E-06 | 1.1E-06 |
| 11230 | 1.44E-06 | 1.27E-06 | 6.16E-07 | 6.14E-07 | 1.09E-06 | 2.05E-06 | 4.01E-06 | 3.02E-06 |
| 11720 | $1.38 \mathrm{E}-06$ | 1.13E-06 | $6.89 \mathrm{E}-07$ | 3.52E-07 | 5.11E-07 | 9.51E-07 | 1.66E-06 | 9.29E-07 |
| 12210 | 1.38E-06 | 1.02E-06 | 6.26E-07 | 3.53E-07 | 5.94E-07 | 7.77E-07 | 1.09E-06 | 1.11E-06 |
| 12700 | 9.17E-07 | 1.09E-06 | $6.31 \mathrm{E}-07$ | 3.16E-07 | 6.06E-07 | 8.87E-07 | 1.09E-06 | 8.37E-07 |
| 13180 | 9.81E-07 | 1.05E-06 | 7.23E-07 | 3.06E-07 | 5.89E-07 | 5.26E-07 | $1.06 \mathrm{E}-06$ | 8.54E-07 |
| 13670 | 1.33E-06 | 1.06E-06 | 8.24E-07 | 1.91E-07 | 6.49E-07 | 5.35E-07 | 1.3E-06 | 8.51E-07 |
| 14160 | 9.98E-07 | 1.03E-06 | 6.05E-07 | 2.27E-07 | 5.4E-07 | 4.33E-07 | 1.21E-06 | 9.42E-07 |
| 14650 | 1.01E-06 | 1.12E-06 | 5.97E-07 | $2.38 \mathrm{E}-07$ | 4.18E-07 | 4.61E-07 | 1.03E-06 | 6.93E-07 |
| 15140 | 9.69E-07 | 9.86E-07 | 6.3E-07 | 2.68E-07 | 4.62E-07 | 4.34E-07 | 8.97E-07 | 9.19E-07 |
| 15630 | 9.09E-07 | 9.31E-07 | 5.93E-07 | 2.11E-07 | 3.96E-07 | 2.78E-07 | 9.56E-07 | 7.6E-07 |
| 16110 | 1.03E-06 | 1.36E-06 | 6.73E-07 | 1.86E-07 | 4.42E-07 | 3.25E-07 | $8.16 \mathrm{E}-07$ | 7.61E-07 |
| 16600 | 8.21E-07 | 9.59E-07 | 5.95E-07 | 1.94E-07 | 4.58E-07 | 2.94E-07 | 7.12E-07 | 7.8E-07 |
| 17090 | 8.92E-07 | 8.69E-07 | 5.06E-07 | 1.5E-07 | $4.38 \mathrm{E}-07$ | 2.61E-07 | 6.15E-07 | 8.02E-07 |
| 17580 | 6.53E-07 | 8.73E-07 | $4.93 \mathrm{E}-07$ | 1.37E-07 | $3.04 \mathrm{E}-07$ | 2.1E-07 | 7.08E-07 | 7.16E-07 |
| 18070 | 7.35E-07 | 1E-06 | $4.49 \mathrm{E}-07$ | $1.25 \mathrm{E}-07$ | 4.39E-07 | 2.36E-07 | 9E-07 | 6.01E-07 |
| 18550 | 8.43E-07 | 6.61E-07 | $5.26 \mathrm{E}-07$ | 1.05E-07 | 2.52E-07 | 2.08E-07 | 5.7E-07 | 6.04E-07 |
| 19040 | 5.47E-07 | 1.08E-06 | 5.23E-07 | 1.15E-07 | 3.3E-07 | 2.07E-07 | 6.35E-07 | 6.93E-07 |
| 19530 | 7.14E-07 | 9.26E-07 | 3.37E-07 | 1.03E-07 | 2.59E-07 | $1.65 \mathrm{E}-07$ | $5.66 \mathrm{E}-07$ | 5.49E-07 |
| 20020 | 6.63E-07 | 7.4E-07 | $5.11 \mathrm{E}-07$ | 1.04E-07 | $3.14 \mathrm{E}-07$ | $1.56 \mathrm{E}-07$ | 6.02E-07 | $5.78 \mathrm{E}-07$ |
| 20510 | 5.89E-07 | 7.87E-07 | 4.18E-07 | 9.44E-08 | 2.96E-07 | 1.47E-07 | 4.43E-07 | 4.09E-07 |
| 21000 | 5.88E-07 | 7.55E-07 | 5E-07 | 9.67E-08 | 2.25E-07 | $1.68 \mathrm{E}-07$ | 4.09E-07 | 3.69E-07 |
| 21480 | 5.78E-07 | 7.39E-07 | 4.45E-07 | $6.5 \mathrm{E}-08$ | 1.93E-07 | 1.61E-07 | 4.33E-07 | 4.99E-07 |
| 21970 | 3.74E-07 | $6.74 \mathrm{E}-07$ | 4.03E-07 | 6.75E-08 | 2.38E-07 | 9.84E-08 | 3.25E-07 | 5.09E-07 |
| 22460 | 4.53E-07 | 5.74E-07 | 4.73E-07 | 8.21E-08 | 1.77E-07 | 1.18E-07 | 3.84E-07 | 4.74E-07 |
| 22950 | 4.66E-07 | 6.03E-07 | 4.49E-07 | $6.39 \mathrm{E}-08$ | 1.72E-07 | 1.51E-07 | 4.59E-07 | 3.83E-07 |
| 23440 | 4.67E-07 | 5.93E-07 | $3.58 \mathrm{E}-07$ | $6.48 \mathrm{E}-08$ | 1.55E-07 | 1E-07 | 2.85E-07 | 4.03E-07 |
| 23930 | 4.66E-07 | 4.84E-07 | 4.47E-07 | 5.77E-08 | 1.73E-07 | 1.12E-07 | 3.45E-07 | 3.01E-07 |
| 24410 | 4.65E-07 | 5.05E-07 | $4.14 \mathrm{E}-07$ | 4.77E-08 | 1.36E-07 | $7.09 \mathrm{E}-08$ | 4E-07 | 3.85E-07 |
| 24900 | 3.92E-07 | $5.35 \mathrm{E}-07$ | 3.68E-07 | 4.3E-08 | 1.85E-07 | 8.17E-08 | 3.12E-07 | 3.7E-07 |
| 25390 | 3.97E-07 | 5.81E-07 | 3.24E-07 | 4.46E-08 | 1.35E-07 | 6.32E-08 | 2.7E-07 | 2.48E-07 |
| 25880 | 4.14E-07 | 5.5E-07 | 3.83E-07 | $4.93 \mathrm{E}-08$ | 9.23E-08 | 5.56E-08 | 1.9E-07 | 3.53E-07 |
| 26370 | $3.35 \mathrm{E}-07$ | 5.23E-07 | 3.91E-07 | 4E-08 | 1.17E-07 | $5.26 \mathrm{E}-08$ | 2.36E-07 | 3.39E-07 |
| 26860 | 3.05E-07 | 4.23E-07 | 3.66E-07 | 4.19E-08 | 1.1E-07 | $6.63 \mathrm{E}-08$ | 2.6E-07 | 2.63E-07 |
| 27340 | 3E-07 | $5.35 \mathrm{E}-07$ | 3.5E-07 | 2.93E-08 | 1E-07 | 4.72E-08 | 2.18E-07 | 2.55E-07 |


| 27830 | 3.79E-07 | 4.74E-07 | 3.22E-07 | 3.17E-08 | 9.76E-08 | 4.21E-08 | 1.79E-07 | $3.09 \mathrm{E}-07$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 28320 | 2.94E-07 | 4.93E-07 | 3.07E-07 | 2.81E-08 | $6.98 \mathrm{E}-08$ | 4.52E-08 | 1.51E-07 | 2.19E-07 |
| 28810 | 3.83E-07 | 4.62E-07 | 3.35E-07 | 2.03E-08 | $9.15 \mathrm{E}-08$ | 4.21E-08 | 1.74E-07 | 2.41E-07 |
| 29300 | $2.69 \mathrm{E}-07$ | 4.98E-07 | 2.47E-07 | 2.21E-08 | $8.22 \mathrm{E}-08$ | 3.52E-08 | 1.43E-07 | 2.27E-07 |
| 29790 | 2.93E-07 | 4.47E-07 | 3.36E-07 | 3.22E-08 | 6.25E-08 | 3.29E-08 | 1.77E-07 | 1.67E-07 |
| 30270 | 2.49E-07 | 4.52E-07 | 2.02E-07 | 1.93E-08 | 7.11E-08 | $3.54 \mathrm{E}-08$ | 1.95E-07 | 1.89E-07 |
| 30760 | 2.37E-07 | 3.81E-07 | 2.63E-07 | 1.65E-08 | 7.45E-08 | 2.63E-08 | 1.55E-07 | 2.37E-07 |
| 31250 | 2.66E-07 | 3.11E-07 | 2.54E-07 | 2.17E-08 | $5.99 \mathrm{E}-08$ | $2.73 \mathrm{E}-08$ | $1.35 \mathrm{E}-07$ | 1.89E-07 |
| 31740 | 2.36E-07 | 4.5E-07 | 2.63E-07 | 1.61E-08 | 6.29E-08 | 2.63E-08 | 1.74E-07 | 1.26E-07 |
| 32230 | 2.72E-07 | 3.67E-07 | 2.6E-07 | 1.78E-08 | 5.43E-08 | 2.46E-08 | 1.07E-07 | 1.28E-07 |
| 32710 | 2.79E-07 | 3.08E-07 | 2.16E-07 | 1.49E-08 | 4.63E-08 | 1.87E-08 | $1.08 \mathrm{E}-07$ | 1.25E-07 |
| 33200 | 2.13E-07 | 3.65E-07 | 2.44E-07 | 1.42E-08 | 4.11E-08 | 2.53E-08 | 7.05E-08 | 1.08E-07 |
| 33690 | 2.35E-07 | 3.15E-07 | 2.07E-07 | 1.11E-08 | 4.27E-08 | 1.53E-08 | 1.07E-07 | 1.13E-07 |
| 34180 | 1.71E-07 | 2.91E-07 | 2.27E-07 | 8.26E-09 | 2.89E-08 | $1.58 \mathrm{E}-08$ | 1.11E-07 | 1.14E-07 |
| 34670 | 1.85E-07 | 2.45E-07 | 2.05E-07 | 9.37E-09 | 2.2E-08 | 1.81E-08 | 8.83E-08 | 1.12E-07 |
| 35160 | 2.06E-07 | 2.72E-07 | 1.77E-07 | 8.97E-09 | 3.92E-08 | 1.88E-08 | 1.13E-07 | 1.22E-07 |
| 35640 | $1.74 \mathrm{E}-07$ | 2.9E-07 | 1.63E-07 | 1.09E-08 | $3.19 \mathrm{E}-08$ | $2.65 \mathrm{E}-08$ | 1.28E-07 | 1.12E-07 |
| 36130 | 2.06E-07 | 2.32E-07 | 1.74E-07 | 7.76E-09 | $3.46 \mathrm{E}-08$ | 1.8E-08 | 8.13E-08 | 9.41E-08 |
| 36620 | 1.76E-07 | $2.48 \mathrm{E}-07$ | 1.9E-07 | 7.87E-09 | 2.96E-08 | 1.54E-08 | $6.24 \mathrm{E}-08$ | 8.15E-08 |
| 37110 | 1.64E-07 | 2.35E-07 | 2.02E-07 | 1.21E-08 | 2.37E-08 | 1.72E-08 | 7.25E-08 | 9.45E-08 |
| 37600 | 1.5E-07 | 2.32E-07 | 1.81E-07 | 6.41E-09 | 2.48E-08 | 1.54E-08 | 7.55E-08 | $6.77 \mathrm{E}-08$ |
| 38090 | 1.23E-07 | 2.37E-07 | 1.92E-07 | 7.73E-09 | 2.71E-08 | $1.51 \mathrm{E}-08$ | 8.39E-08 | 7.91E-08 |
| 38570 | 1.13E-07 | 2.36E-07 | 1.26E-07 | 5.96E-09 | $2.4 \mathrm{E}-08$ | $1.39 \mathrm{E}-08$ | 5.19E-08 | $7.96 \mathrm{E}-08$ |
| 39060 | $1.24 \mathrm{E}-07$ | 2.09E-07 | $1.48 \mathrm{E}-07$ | 6.53E-09 | 1.75E-08 | 1.07E-08 | 5.51E-08 | $5.87 \mathrm{E}-08$ |
| 39550 | 1.12E-07 | 1.77E-07 | 1.48E-07 | 6.3E-09 | 2.41E-08 | $1.24 \mathrm{E}-08$ | 5.54E-08 | 6.57E-08 |
| 40040 | 1.26E-07 | 1.8E-07 | 1.29E-07 | 4.47E-09 | 1.86E-08 | $1.43 \mathrm{E}-08$ | $6.27 \mathrm{E}-08$ | $4.84 \mathrm{E}-08$ |
| 40530 | 1.24E-07 | 1.93E-07 | 1.37E-07 | 5.38E-09 | 1.62E-08 | 1.18E-08 | 5.47E-08 | 6.58E-08 |
| 41020 | 1.32E-07 | 2.04E-07 | 1.75E-07 | $4.68 \mathrm{E}-09$ | 1.55E-08 |  |  |  |
| 41500 | 1.36E-07 | 1.81E-07 | 1.2E-07 | 4.46E-09 | 1.77E-08 |  |  |  |
| 41990 | 9.86E-08 | 2.01E-07 | 1.04E-07 | 4.18E-09 | 1.22E-08 |  |  |  |
| 42480 | 1.03E-07 | 1.5E-07 | 1.22E-07 | $4.18 \mathrm{E}-09$ | 1.18E-08 |  |  |  |
| 42970 | 8.14E-08 | 1.48E-07 | 1.34E-07 | $4.38 \mathrm{E}-09$ | 1.34E-08 |  |  |  |
| 43460 | 8.6E-08 | 1.89E-07 | 8.89E-08 | 3.29E-09 | 1.32E-08 |  |  |  |
| 43950 | 7.97E-08 | 1.49E-07 | 1.19E-07 | $3.29 \mathrm{E}-09$ | 1E-08 |  |  |  |
| 44430 | 1.04E-07 | 1.4E-07 | 1.12E-07 | 3.98E-09 | 1.53E-08 |  |  |  |
| 44920 | 8.94E-08 | 1.24E-07 | 1.14E-07 | $3.08 \mathrm{E}-09$ | 1.17E-08 |  |  |  |
| 45410 | 8.3E-08 | 1.27E-07 | 1.23E-07 | 2.86E-09 | 1.18E-08 |  |  |  |
| 45900 | $6.21 \mathrm{E}-08$ | 1.21E-07 | 9.79E-08 | 2.92E-09 | 8.88E-09 |  |  |  |
| 46390 | 7.39E-08 | 1.07E-07 | 1.01E-07 | 2.73E-09 | 9.3E-09 |  |  |  |
| 46880 | 7.35E-08 | 1.51E-07 | $9.94 \mathrm{E}-08$ | 3.31E-09 | 7.83E-09 |  |  |  |
| 47360 | 6.43E-08 | 1.22E-07 | 8.81E-08 | 3.11E-09 | 8.27E-09 |  |  |  |
| 47850 | $6.11 \mathrm{E}-08$ | 1.08E-07 | 9.47E-08 | 3E-09 | 7.97E-09 |  |  |  |
| 48340 | 9.19E-08 | 1.17E-07 | 8.56E-08 | 3.41E-09 | 7.19E-09 |  |  |  |
| 48830 | $5.64 \mathrm{E}-08$ | 1.06E-07 | $8.24 \mathrm{E}-08$ | 4.15E-09 | 6.94E-09 |  |  |  |
| 49320 | 5.82E-08 | 1.13E-07 | 9.32E-08 | 3.57E-09 | 8.44E-09 |  |  |  |
| 49800 | 4.51E-08 | 8.52E-08 | 6.72E-08 | 4.44E-09 | 6.9E-09 |  |  |  |
| 50290 | 5.23E-08 | 8.58E-08 | 8.57E-08 | 4.47E-09 | 6.09E-09 |  |  |  |
| 50780 | $5.35 \mathrm{E}-08$ | 8.15E-08 | 7.14E-08 | 5.1E-09 | 7.27E-09 |  |  |  |
| 51270 | 4.9E-08 | 8.92E-08 | 6.62E-08 | 5.22E-09 | 5.52E-09 |  |  |  |
| 51760 | $3.94 \mathrm{E}-08$ | 8.62E-08 | $6.36 \mathrm{E}-08$ | 5.18E-09 | 6.56E-09 |  |  |  |
| 52250 | $3.84 \mathrm{E}-08$ | 7.21E-08 | $6.62 \mathrm{E}-08$ | 3.92E-09 | 5.61E-09 |  |  |  |
| 52730 | 4.25E-08 | 8.16E-08 | $7.33 \mathrm{E}-08$ | 5.18E-09 | 3.69E-09 |  |  |  |
| 53220 | 3.85E-08 | 7.42E-08 | 4.91E-08 | 3.07E-09 | 6.66E-09 |  |  |  |
| 53710 | 3.5E-08 | $7.73 \mathrm{E}-08$ | 5.62E-08 | 2.91E-09 | 5.58E-09 |  |  |  |
| 54200 | 2.5E-08 | $6.66 \mathrm{E}-08$ | $5.24 \mathrm{E}-08$ | 1.8E-09 | $6.98 \mathrm{E}-09$ |  |  |  |
| 54690 | 3.72E-08 | 4.4E-08 | 3.84E-08 | 1.74E-09 | 6.68E-09 |  |  |  |
| 55180 | 2.74E-08 | 6.57E-08 | 5.63E-08 |  | 6.05E-09 |  |  |  |
| 55660 | 3.16E-08 | 5.81E-08 | 4.61E-08 |  | 7.65E-09 |  |  |  |
| 56150 | 2.41E-08 | 4.72E-08 | 4.7E-08 |  | 6.87E-09 |  |  |  |
| 56640 | $3.14 \mathrm{E}-08$ | 5.57E-08 | 4.74E-08 |  | 8.8E-09 |  |  |  |


| 57130 | 2.31E-08 | 4.39E-08 | $4.02 \mathrm{E}-08$ | 9.42E-09 |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 57620 | 2.4E-08 | 5.11E-08 | 3.24E-08 | $6.73 \mathrm{E}-09$ |  |  |  |
| 58110 | 2.13E-08 | 3.93E-08 | 3.02E-08 | $6.35 \mathrm{E}-09$ |  |  |  |
| 58590 | 1.95E-08 | 4.8E-08 | $3.93 \mathrm{E}-08$ | 5.46E-09 |  |  |  |
| 59080 | 2.22E-08 | 5.45E-08 | $3.41 \mathrm{E}-08$ | 5.99E-09 |  |  |  |
| 59570 | 2.17E-08 | 3.12E-08 | 3E-08 | 4.7E-09 |  |  |  |
| 60060 | 2.07E-08 | 4.01E-08 | $3.11 \mathrm{E}-08$ | 5.19E-09 |  |  |  |
| 60550 | 1.72E-08 | $3.39 \mathrm{E}-08$ | 3.08E-08 | 3.77E-09 |  |  |  |
| 61040 | 2.04E-08 | 4.07E-08 | 3.21E-08 | $5.11 \mathrm{E}-09$ |  |  |  |
| 61520 | 1.33E-08 | 2.53E-08 | 2.88E-08 | 3.53E-09 |  |  |  |
| 62010 | 1.34E-08 | 2.68E-08 | 2.97E-08 | $3.58 \mathrm{E}-09$ |  |  |  |
| 62500 | 1.48E-08 | 2.46E-08 | 1.84E-08 | 3.51E-09 |  |  |  |
| 62990 | 1.21E-08 | 2.1E-08 | 1.81E-08 | 3.7E-09 |  |  |  |
| 63480 | 1.91E-08 | 2.67E-08 | $1.98 \mathrm{E}-08$ |  |  |  |  |
| 63960 | 1.23E-08 | 2.09E-08 | $1.69 \mathrm{E}-08$ |  |  |  |  |
| 64450 | $1.38 \mathrm{E}-08$ | 2.03E-08 | 1.59E-08 |  |  |  |  |
| 64940 | 1.27E-08 | 2.33E-08 | 1.85E-08 |  |  |  |  |
| 65430 | 1.06E-08 | 2.03E-08 | $1.89 \mathrm{E}-08$ |  |  |  |  |
| 65920 | 8.31E-09 | 2.01E-08 | 1.82E-08 |  |  |  |  |
| 66410 | 9.8E-09 | 1.87E-08 | 1.9E-08 |  |  |  |  |
| 66890 | 7.91E-09 | $1.99 \mathrm{E}-08$ | 1.49E-08 |  |  |  |  |
| 67380 | 8.61E-09 | $1.89 \mathrm{E}-08$ | 1.72E-08 |  |  |  |  |
| 67870 | 6.23E-09 | 1.47E-08 | $2.28 \mathrm{E}-08$ |  |  |  |  |
| 68360 | 7.97E-09 | 1.37E-08 | 1.28E-08 |  |  |  |  |
| 68850 | 5.48E-09 | 1.63E-08 | 1.16E-08 |  |  |  |  |
| 69340 | 7E-09 | 1.33E-08 | 1.18E-08 |  |  |  |  |
| 69820 | 7.16E-09 | 1.67E-08 | 1.18E-08 |  |  |  |  |
| 70310 | $3.68 \mathrm{E}-09$ | 1.12E-08 | $1.44 \mathrm{E}-08$ |  |  |  |  |
| 70800 | 4.46E-09 | $1.29 \mathrm{E}-08$ | $1.05 \mathrm{E}-08$ |  |  |  |  |
| 71290 | 4.65E-09 | 1.19E-08 | 9.02E-09 |  |  |  |  |
| 71780 | 4.12E-09 | 8.68E-09 | 8.18E-09 |  |  |  |  |
| 72270 | 4.53E-09 | 9.94E-09 | 9.49E-09 |  |  |  |  |
| 72750 | 3.91E-09 | 1.03E-08 | 8.36E-09 |  |  |  |  |
| 73240 | 3.67E-09 | 1.04E-08 | 6.92E-09 |  |  |  |  |
| 73730 | 3.38E-09 | 7.88E-09 | $6.58 \mathrm{E}-09$ |  |  |  |  |
| 74220 | 2.95E-09 | 6.46E-09 | 6.92E-09 |  |  |  |  |
| 74710 | 3.1E-09 | 8.36E-09 | 6.46E-09 |  |  |  |  |
| 75200 | 3.67E-09 | 7.53E-09 | $6.26 \mathrm{E}-09$ |  |  |  |  |
| 75680 | 3.33E-09 | 6.5E-09 | $6.53 \mathrm{E}-09$ |  |  |  |  |
| 76170 | 3.37E-09 | 4.58E-09 | 6.24E-09 |  |  |  |  |
| 76660 | $2.69 \mathrm{E}-09$ | 4.7E-09 | $6.05 \mathrm{E}-09$ |  |  |  |  |
| 77150 | 4.5E-09 | 5.19E-09 | $6.55 \mathrm{E}-09$ |  |  |  |  |
| 77640 | 3.25E-09 | 3.79E-09 | 5.09E-09 |  |  |  |  |
| 78130 | 2.69E-09 | 4.33E-09 | 4.91E-09 |  |  |  |  |
| 78610 | 3.7E-09 | 4.16E-09 | 4.84E-09 |  |  |  |  |
| 79100 | 3.93E-09 | 3.69E-09 | 5.35E-09 |  |  |  |  |
| 79590 | $3.47 \mathrm{E}-09$ | 3.52E-09 | 4.02E-09 |  |  |  |  |
| 80080 | 4.93E-09 | $3.96 \mathrm{E}-09$ | 3.72E-09 |  |  |  |  |
| 80570 | 4.38E-09 | 4.04E-09 | 2.63E-09 |  |  |  |  |
| 81050 | 3.41E-09 | 4E-09 | 2.16E-09 |  |  |  |  |
| 81540 | 3.59E-09 | 3.51E-09 | 2.61E-09 |  |  |  |  |
| 82030 | 3.09E-09 | 3.79E-09 | 1.8E-09 |  |  |  |  |
| 82520 | 2.89E-09 | 3.57E-09 | $2.54 \mathrm{E}-09$ |  |  |  |  |
| 83010 | 2.16E-09 | 3.76E-09 | 1.63E-09 |  |  |  |  |
| 83500 | 2.65E-09 | 3.7E-09 | 1.89E-09 |  |  |  |  |
| 83980 | 2.45E-09 | $3.05 \mathrm{E}-09$ | 2E-09 |  |  |  |  |
| 84470 | 1.92E-09 | 4.63E-09 | 1.83E-09 |  |  |  |  |
| 84960 | 2.09E-09 | 4.59E-09 | 1.92E-09 |  |  |  |  |
| 85450 | 1.73E-09 | 5.15E-09 | 1.76E-09 |  |  |  |  |
| 85940 | 1.59E-09 | 4.07E-09 | 1.76E-09 |  |  |  |  |
| 86430 | 1.27E-09 | 5.71E-09 | 1.92E-09 |  |  |  |  |


| 86910 | 1.78E-09 | 4.71E-09 | 1.22E-09 |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 87400 | 1.55E-09 | 4.72E-09 | 1.82E-09 |  |  |  |  |  |
| 87890 | 1.59E-09 | 3.59E-09 | 1.6E-09 |  |  |  |  |  |
| 88380 | 1.06E-09 | 3.9E-09 | 1.5E-09 |  |  |  |  |  |
| 88870 | 1.05E-09 | 3.97E-09 | 1.65E-09 |  |  |  |  |  |
| 89360 | $8.73 \mathrm{E}-10$ | 4.65E-09 | 1.47E-09 |  |  |  |  |  |
| 89840 | $9.35 \mathrm{E}-10$ | 3.15E-09 | 1.4E-09 |  |  |  |  |  |
| 90330 | 9.68E-10 | 3.26E-09 | 1.45E-09 |  |  |  |  |  |
| 90820 | $6.27 \mathrm{E}-10$ | $3.25 \mathrm{E}-09$ | 1.31E-09 |  |  |  |  |  |
| 91310 | 7.46E-10 | 2.35E-09 | 1.17E-09 |  |  |  |  |  |
| 91800 | $8.73 \mathrm{E}-10$ | 2.41E-09 | 1.45E-09 |  |  |  |  |  |
| 92290 | $6.94 \mathrm{E}-10$ | 2.43E-09 | 1.35E-09 |  |  |  |  |  |
| 92770 | $6.72 \mathrm{E}-10$ | 2.16E-09 | 1.46E-09 |  |  |  |  |  |
| 93260 | $5.37 \mathrm{E}-10$ | 1.62E-09 | $1.61 \mathrm{E}-09$ |  |  |  |  |  |
| 93750 | $6.19 \mathrm{E}-10$ | 1.75E-09 | 1.11E-09 |  |  |  |  |  |
| 94240 | 4.02E-10 | 1.74E-09 | 1.18E-09 |  |  |  |  |  |
| 94730 | $5.52 \mathrm{E}-10$ | 1.43E-09 | 1.33E-09 |  |  |  |  |  |
| 95210 | 4.94E-10 | 1.36E-09 | 1.85E-09 |  |  |  |  |  |
| 95700 | 4.13E-10 | 1.28E-09 | 1.44E-09 |  |  |  |  |  |
| 96190 | $4.69 \mathrm{E}-10$ | 1.2E-09 | 1.85E-09 |  |  |  |  |  |
| 96680 | $3.42 \mathrm{E}-10$ | $9.6 \mathrm{E}-10$ | 1.82E-09 |  |  |  |  |  |
| 97170 | $3.54 \mathrm{E}-10$ | 9.88E-10 | 1.65E-09 |  |  |  |  |  |
| 97660 | $2.96 \mathrm{E}-10$ | $1.06 \mathrm{E}-09$ | 1.56E-09 |  |  |  |  |  |
| 98140 | 2.76E-10 | $9.01 \mathrm{E}-10$ | 1.52E-09 |  |  |  |  |  |
| 98630 | $2.75 \mathrm{E}-10$ | $6.95 \mathrm{E}-10$ | 1.29E-09 |  |  |  |  |  |
| 99120 | 2.93E-10 | 6.61E-10 | 1.1E-09 |  |  |  |  |  |
| 99610 | 2.87E-10 | $6.16 \mathrm{E}-10$ | 1.2E-09 |  |  |  |  |  |

Table C-20: Mach 2.8 ZPG PIV data

| y (in) | u (m/s) | $\mathrm{v}(\mathrm{m} / \mathrm{s})$ | uTl (\%) | vTI (\%) | shear | $\mathrm{y} / \delta_{\mathrm{u}}$ | $\mathrm{y}^{+}$ | $\mathrm{u}_{\text {eff }}$ | $\mathbf{u}^{+}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.014 | 433.514 | 14.609 | 53.389 | 6.192 | 9.88E-05 | 0.035 | 38.971 | 465.194 | 16.210 |
| 0.046 | 443.560 | 20.154 | 13.568 | 4.647 | -7.8E-05 | 0.116 | 130.522 | 477.645 | 16.644 |
| 0.078 | 481.606 | 15.508 | 8.489 | 3.475 | -8.9E-05 | 0.197 | 222.074 | 526.193 | 18.336 |
| 0.111 | 500.204 | 12.904 | 8.129 | 4.180 | -5.8E-05 | 0.279 | 313.625 | 550.838 | 19.194 |
| 0.143 | 520.826 | 11.498 | 7.187 | 4.334 | -7.6E-06 | 0.360 | 405.177 | 578.987 | 20.175 |
| 0.175 | 536.771 | 13.568 | 6.738 | 3.931 | 0.000115 | 0.441 | 496.729 | 601.422 | 20.957 |
| 0.208 | 551.751 | 13.380 | 6.206 | 3.452 | 0.00013 | 0.523 | 588.280 | 623.094 | 21.712 |
| 0.240 | 567.558 | 11.181 | 5.609 | 2.797 | 0.000105 | 0.604 | 679.832 | 646.670 | 22.534 |
| 0.272 | 579.507 | 10.237 | 4.296 | 2.695 | 1.43E-05 | 0.685 | 771.383 | 665.027 | 23.173 |
| 0.304 | 589.037 | 10.259 | 3.582 | 2.410 | 4.6E-06 | 0.767 | 862.935 | 680.037 | 23.696 |
| 0.337 | 595.952 | 11.776 | 3.621 | 2.090 | -1E-06 | 0.848 | 954.486 | 691.149 | 24.084 |
| 0.369 | 595.304 | 12.785 | 2.735 | 2.620 | 1.19E-06 | 0.929 | 1046.038 | 690.100 | 24.047 |
| 0.401 | 579.558 | 11.051 | 17.165 | 2.993 | 3.09E-06 | 1.011 | 1137.590 | 665.107 | 23.176 |
| 0.434 | 601.411 | 9.860 | 2.602 | 2.742 | -1.3E-05 | 1.092 | 1229.141 | 700.061 | 24.394 |
| 0.466 | 599.442 | 11.992 | 3.121 | 2.641 | -1.2E-05 | 1.173 | 1320.693 | 696.831 | 24.282 |
| 0.498 | 598.978 | 14.631 | 2.603 | 2.543 | 1.03E-06 | 1.255 | 1412.244 | 696.072 | 24.255 |
| 0.531 | 596.450 | 17.444 | 2.569 | 2.537 | 1.33E-05 | 1.336 | 1503.796 | 691.956 | 24.112 |
| 0.563 | 593.541 | 14.807 | 2.836 | 2.583 | -1.1E-05 | 1.417 | 1595.347 | 687.252 | 23.948 |
| 0.595 | 597.547 | 14.877 | 2.805 | 3.034 | -1.5E-05 | 1.499 | 1686.899 | 693.740 | 24.174 |
| 0.627 | 602.018 | 16.747 | 2.892 | 2.479 | -1.2E-05 | 1.580 | 1778.450 | 701.061 | 24.429 |
| 0.660 | 604.290 | 18.742 | 2.960 | 2.890 | -1.2E-05 | 1.661 | 1870.002 | 704.814 | 24.560 |
| 0.692 | 603.014 | 20.725 | 2.801 | 3.206 | -3E-06 | 1.743 | 1961.553 | 702.703 | 24.486 |
| 0.724 | 600.511 | 21.776 | 2.624 | 3.080 | 4.45E-06 | 1.824 | 2053.105 | 698.584 | 24.343 |
| 0.757 | 602.101 | 21.328 | 3.075 | 2.600 | 2.13E-06 | 1.906 | 2144.656 | 701.197 | 24.434 |
| 0.789 | 607.687 | 17.228 | 3.251 | 2.397 | 9.06E-06 | 1.987 | 2236.208 | 710.470 | 24.757 |

Table C-21: Mach 2.8 FPG PIV data - 25 images

| y (in) | $\mathrm{u}(\mathrm{m} / \mathrm{s})$ | $\mathrm{v}(\mathrm{m} / \mathrm{s})$ | uTI (\%) | vTI (\%) | shear | $\mathrm{y} / \delta_{0}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 2.042 | 411.144 | 31.607 | -167.470 | -159.012 | -172.582 | 0.061 |
| 1.961 | 517.126 | 45.202 | 6.557 | 18.738 | -0.00024 | 0.129 |
| 1.880 | 529.859 | 44.705 | 6.535 | 18.528 | -4.7E-05 | 0.197 |
| 1.799 | 542.286 | 40.064 | 6.659 | 16.404 | -3.7E-05 | 0.265 |
| 1.718 | 559.302 | 35.523 | 6.630 | 14.085 | -6.2E-06 | 0.333 |
| 1.638 | 569.197 | 35.365 | 6.173 | 13.815 | -5.8E-05 | 0.401 |
| 1.557 | 576.951 | 35.338 | 6.679 | 13.615 | 3.81E-05 | 0.469 |
| 1.476 | 582.564 | 29.688 | 6.371 | 11.413 | 0.000228 | 0.537 |
| 1.395 | 583.898 | 28.055 | 6.214 | 10.937 | 0.000188 | 0.605 |
| 1.314 | 588.501 | 37.541 | 5.938 | 13.820 | 0.000227 | 0.673 |
| 1.233 | 599.054 | 32.099 | 5.780 | 11.854 | -2E-05 | 0.741 |
| 1.152 | 608.734 | 26.967 | 4.245 | 10.058 | -0.00025 | 0.809 |
| 1.071 | 607.658 | 28.944 | 3.411 | 11.060 | -0.00011 | 0.877 |
| 0.990 | 610.400 | 25.392 | 3.110 | 9.832 | 6.51E-05 | 0.945 |
| 0.909 | 613.976 | 20.376 | 4.210 | 8.805 | 0.000111 | 1.013 |
| 0.828 | 615.290 | 18.943 | 4.222 | 8.472 | -6.2E-06 | 1.081 |
| 0.747 | 611.317 | 17.520 | 3.897 | 7.674 | -5.4E-05 | 1.149 |
| 0.666 | 608.042 | 16.183 | 3.903 | 6.914 | -3.2E-05 | 1.217 |
| 0.585 | 605.090 | 17.893 | 3.944 | 7.197 | -1.4E-05 | 1.285 |
| 0.505 | 608.711 | 19.596 | 4.122 | 7.712 | -2.8E-05 | 1.353 |
| 0.424 | 605.350 | 16.599 | 7.251 | 7.234 | -0.00011 | 1.421 |
| 0.343 | 591.270 | 10.975 | 16.195 | 6.139 | -8E-07 | 1.489 |
| 0.262 | 598.123 | 17.087 | 5.095 | 8.014 | -5.9E-05 | 1.557 |
| 0.181 | 600.596 | 21.672 | 3.347 | 8.635 | -7.8E-05 | 1.625 |
| 0.100 | 605.107 | 7.898 | 3.229 | 4.510 | -0.00013 | 1.693 |

Table C-22: Mach 2.8 FPG PIV data - 93 images

| $y$ (in) | $\mathrm{u}(\mathrm{m} / \mathrm{s})$ | $\mathrm{v}(\mathrm{m} / \mathrm{s})$ | uTI (\%) | VTI (\%) | shear | $\mathrm{y} / \delta_{\mathrm{u}}$ | $\mathrm{y}^{+}$ | $\mathrm{u}_{\text {eff }}$ | $\mathrm{u}^{+}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.025 | 401.798 | 30.281 |  |  |  | 0.061 | 91.59 | 425.05 | 15.62 |
| 0.054 | 508.438 | 40.714 | 7.547 | 16.927 | 6.44E-05 | 0.129 | 193.40 | 557.86 | 20.50 |
| 0.082 | 529.027 | 40.966 | 6.778 | 16.444 | $1.29 \mathrm{E}-04$ | 0.197 | 295.20 | 585.60 | 21.52 |
| 0.110 | 539.129 | 37.625 | 7.992 | 15.021 | 8.79E-05 | 0.265 | 397.00 | 599.53 | 22.03 |
| 0.138 | 554.481 | 34.276 | 7.535 | 13.395 | 8.34E-05 | 0.333 | 498.80 | 621.14 | 22.82 |
| 0.167 | 564.971 | 33.749 | 6.868 | 13.062 | 1.39E-04 | 0.401 | 600.61 | 636.25 | 23.38 |
| 0.195 | 573.944 | 34.437 | 6.584 | 13.129 | 1.45E-04 | 0.469 | 702.41 | 649.40 | 23.86 |
| 0.223 | 576.676 | 32.273 | 6.466 | 12.142 | $1.68 \mathrm{E}-04$ | 0.537 | 804.21 | 653.45 | 24.01 |
| 0.251 | 582.925 | 31.884 | 6.210 | 12.163 | 2.59E-04 | 0.605 | 906.01 | 662.80 | 24.35 |
| 0.279 | 590.046 | 33.639 | 5.922 | 12.505 | -7.36E-05 | 0.673 | 1007.82 | 673.60 | 24.75 |
| 0.308 | 597.133 | 31.207 | 5.467 | 11.446 | -3.22E-05 | 0.741 | 1109.62 | 684.51 | 25.15 |
| 0.336 | 597.509 | 29.104 | 8.321 | 10.793 | -7.24E-05 | 0.809 | 1211.42 | 685.10 | 25.17 |
| 0.364 | 605.792 | 26.175 | 4.483 | 9.832 | -5.13E-06 | 0.877 | 1313.22 | 698.08 | 25.65 |
| 0.392 | 610.852 | 23.638 | 4.294 | 8.949 | -6.67E-07 | 0.945 | 1415.03 | 706.14 | 25.95 |
| 0.421 | 608.306 | 19.529 | 4.388 | 7.851 | 2.10E-06 | 1.013 | 1516.83 | 702.07 | 25.80 |
| 0.449 | 609.043 | 17.798 | 4.486 | 7.401 | -1.75E-05 | 1.081 | 1618.63 | 703.25 | 25.84 |
| 0.477 | 611.238 | 17.073 | 4.364 | 7.075 | -1.58E-05 | 1.149 | 1720.43 | 706.76 | 25.97 |
| 0.505 | 609.992 | 16.746 | 4.309 | 6.774 | -3.54E-08 | 1.217 | 1822.24 | 704.76 | 25.90 |
| 0.533 | 606.995 | 17.616 | 4.475 | 6.981 | -6.35E-06 | 1.285 | 1924.04 | 699.99 | 25.72 |
| 0.562 | 606.128 | 17.559 | 4.164 | 7.126 | -1.34E-05 | 1.353 | 2025.84 | 698.61 | 25.67 |
| 0.590 | 605.532 | 15.035 | 5.499 | 6.606 | -5.57E-06 | 1.421 | 2127.64 | 697.67 | 25.64 |
| 0.618 | 602.118 | 13.934 | 9.953 | 6.121 | -4.52E-06 | 1.489 | 2229.45 | 692.29 | 25.44 |
| 0.646 | 606.671 | 15.681 | 4.563 | 6.729 | 3.53E-05 | 1.557 | 2331.25 | 699.48 | 25.70 |
| 0.675 | 602.161 | 13.846 | 4.990 | 6.138 | 7.94E-06 | 1.625 | 2433.05 | 692.36 | 25.44 |
| 0.703 | 605.395 | 10.005 | 3.831 | 4.653 | 6.01E-06 | 1.693 | 2534.85 | 697.45 | 25.63 |

Table C-23: Mach 2.8 CPG ( $x=68 \mathrm{~cm}$ ) PIV data

| $y(\mathrm{in})$ | $\mathrm{u}(\mathrm{m} / \mathrm{s})$ | $\mathrm{v}(\mathrm{m} / \mathrm{s})$ | $\mathrm{uTI}(\%)$ | $\mathrm{vTI}(\%)$ | shear | $\mathrm{y} / \delta_{\mathrm{u}}$, |
| ---: | ---: | ---: | ---: | ---: | ---: | ---: |
| 3.176 | 386.885 | -24.305 |  |  |  |  |
| 3.104 | 459.956 | -31.319 | 13.657 | 15.259 | $-6.39 \mathrm{E}-05$ | 0.076 |
| 3.031 | 495.582 | -34.490 | 11.405 | 15.840 | $4.37 \mathrm{E}-04$ | 0.236 |
| 2.959 | 519.896 | -29.595 | 10.108 | 12.867 | $5.44 \mathrm{E}-04$ | 0.316 |
| 2.886 | 537.433 | -28.303 | 8.684 | 11.748 | $8.13 \mathrm{E}-04$ | 0.395 |
| 2.814 | 548.670 | -29.813 | 7.616 | 12.239 | $1.33 \mathrm{E}-03$ | 0.475 |
| 2.741 | 556.021 | -26.331 | 11.352 | 11.281 | $6.76 \mathrm{E}-04$ | 0.555 |
| 2.669 | 568.590 | -22.500 | 10.833 | 9.932 | $5.16 \mathrm{E}-04$ | 0.635 |
| 2.596 | 580.834 | -20.355 | 7.214 | 8.250 | $5.28 \mathrm{E}-04$ | 0.715 |
| 2.523 | 585.771 | -17.353 | 6.083 | 7.056 | $3.74 \mathrm{E}-04$ | 0.795 |
| 2.451 | 593.433 | -17.943 | 5.021 | 7.250 | $4.56 \mathrm{E}-04$ | 0.874 |
| 2.378 | 596.967 | -14.052 | 4.714 | 5.861 | $-8.94 \mathrm{E}-06$ | 0.954 |
| 2.306 | 599.615 | -12.494 | 4.444 | 6.224 | $-3.86 \mathrm{E}-04$ | 1.034 |
| 2.233 | 600.659 | -10.091 | 4.183 | 5.746 | $-3.12 \mathrm{E}-04$ | 1.114 |
| 2.161 | 600.977 | -8.592 | 4.124 | 5.480 | $-2.17 \mathrm{E}-04$ | 1.194 |
| 2.088 | 595.692 | -5.136 | 4.262 | 5.039 | $-1.58 \mathrm{E}-04$ | 1.274 |
| 2.015 | 597.548 | -3.509 | 4.090 | 4.432 | $-9.28 \mathrm{E}-05$ | 1.353 |
| 1.943 | 602.682 | -3.222 | 3.964 | 3.980 | $-1.55 \mathrm{E}-04$ | 1.433 |
| 1.870 | 605.255 | -5.606 | 4.097 | 4.340 | $-2.13 \mathrm{E}-04$ | 1.513 |
| 1.798 | 603.632 | -9.661 | 3.922 | 4.839 | $1.89 \mathrm{E}-04$ | 1.593 |
| 1.725 | 600.125 | -12.391 | 3.984 | 5.299 | $3.12 \mathrm{E}-04$ | 1.673 |
| 1.653 | 600.807 | -9.436 | 4.261 | 4.636 | $3.31 \mathrm{E}-04$ | 1.753 |
| 1.580 | 604.552 | -3.561 | 4.105 | 3.177 | $9.61 \mathrm{E}-05$ | 1.832 |
| 1.508 | 602.623 | -0.735 | 3.730 | 3.602 | $3.70 \mathrm{E}-05$ | 1.912 |

Table C-24: Mach 2.8 CPG ( $\mathrm{x}=68 \mathrm{~cm}$ ) PIV data

| $y(\mathrm{in})$ | $\mathrm{u}(\mathrm{m} / \mathrm{s})$ | $\mathrm{v}(\mathrm{m} / \mathrm{s})$ | $\mathrm{uTI}(\%)$ | $\mathrm{vTI}(\%)$ | shear | $\mathrm{y} / \delta_{\mathrm{u}}$ |
| ---: | ---: | ---: | ---: | ---: | ---: | ---: |
| 2.939 | 372.645 | -34.276 |  |  |  | 0.178 |
| 2.818 | 486.216 | -57.836 | 12.977 | 25.031 | $-1.37 \mathrm{E}-03$ | 0.375 |
| 2.696 | 509.971 | -60.384 | 10.802 | 25.028 | $-3.52 \mathrm{E}-03$ | 0.573 |
| 2.575 | 538.253 | -58.376 | 8.350 | 22.637 | $7.75 \mathrm{E}-04$ | 0.770 |
| 2.453 | 562.558 | -56.920 | 6.027 | 21.208 | $5.93 \mathrm{E}-04$ | 0.967 |
| 2.332 | 573.278 | -52.220 | 5.097 | 19.224 | $-4.25 \mathrm{E}-04$ | 1.165 |
| 2.210 | 578.005 | -42.729 | 4.696 | 15.464 | $1.68 \mathrm{E}-04$ | 1.362 |
| 2.089 | 583.401 | -32.907 | 4.432 | 11.949 | $1.47 \mathrm{E}-04$ | 1.559 |
| 1.968 | 591.659 | -25.485 | 3.951 | 9.527 | $1.15 \mathrm{E}-04$ | 1.756 |
| 1.846 | 601.144 | -16.466 | 4.377 | 6.759 | $9.98 \mathrm{E}-05$ | 1.954 |
| 1.725 | 608.967 | -10.690 | 5.712 | 4.687 | $-3.14 \mathrm{E}-05$ | 2.151 |
| 1.603 | 608.813 | -10.988 | 4.269 | 4.988 | $1.98 \mathrm{E}-04$ | 2.348 |
| 1.482 | 608.541 | -6.166 | 4.214 | 4.445 | $1.75 \mathrm{E}-04$ | 2.546 |
| 1.360 | 609.529 | -6.498 | 3.610 | 3.828 | $1.48 \mathrm{E}-04$ | 2.743 |
| 1.239 | 611.187 | -9.528 | 3.760 | 4.695 | $-4.94 \mathrm{E}-05$ | 2.940 |
| 1.118 | 610.677 | -11.504 | 4.475 | 5.098 | $-2.86 \mathrm{E}-05$ | 3.137 |
| 0.996 | 607.883 | -3.981 | 3.850 | 4.503 | $1.80 \mathrm{E}-05$ | 3.335 |
| 0.875 | 599.867 | -2.595 | 3.874 | 4.107 | $2.41 \mathrm{E}-05$ | 3.532 |
| 0.753 | 603.164 | -10.046 | 3.837 | 4.830 | $-3.23 \mathrm{E}-05$ | 3.729 |
| 0.632 | 608.016 | -8.008 | 3.693 | 4.610 | $2.99 \mathrm{E}-05$ | 3.926 |
| 0.510 | 604.416 | -7.928 | 3.611 | 4.503 | $5.31 \mathrm{E}-05$ | 4.124 |
| 0.389 | 606.330 | -5.445 | 4.472 | 3.429 | $5.69 \mathrm{E}-07$ | 4.321 |
| 0.268 | 605.420 | -7.244 | 4.457 | 4.279 | $-6.67 \mathrm{E}-05$ | 4.518 |
| 0.146 | 600.821 | -13.160 | 3.609 | 5.597 | $-2.06 \mathrm{E}-05$ | 4.716 |
| 0.025 | 604.982 | -9.446 | 4.007 | 2.286 | $-1.09 \mathrm{E}-04$ | 4.913 |

Table C-25: PIV van Driest data


Table C-26: Mach 1.7 ZPG pressure data

| $y$ (in) | $\mathrm{P}_{11}$ (psia) | $\mathrm{P}_{12}$ (psia) | $\mathrm{P}_{\mathrm{T} 2} / \mathrm{P}_{\mathrm{T} 1}$ | $\mathrm{pc} / \mathrm{P}_{12}$ | Mach | $\mathrm{M} / \mathrm{M}_{\mathrm{e}}$ | $\mathrm{u} / \mathrm{U}_{\mathrm{e}}$ | pw/p | $p / \rho e$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.072 | 13.735 | 6.615 | 0.482 | 0.425 | 1.182 | 0.712 | 0.784 | 0.775 | 0.832 |
| 0.077 | 13.733 | 6.847 | 0.499 | 0.410 | 1.210 | 0.729 | 0.799 | 0.767 | 0.840 |
| 0.083 | 13.734 | 7.203 | 0.524 | 0.390 | 1.253 | 0.755 | 0.820 | 0.755 | 0.853 |
| 0.088 | 13.735 | 7.508 | 0.547 | 0.374 | 1.287 | 0.775 | 0.837 | 0.746 | 0.864 |
| 0.093 | 13.735 | 7.915 | 0.576 | 0.355 | 1.332 | 0.802 | 0.858 | 0.733 | 0.879 |
| 0.098 | 13.734 | 8.287 | 0.603 | 0.339 | 1.371 | 0.826 | 0.877 | 0.723 | 0.892 |
| 0.103 | 13.733 | 8.713 | 0.634 | 0.322 | 1.414 | 0.852 | 0.897 | 0.711 | 0.907 |
| 0.108 | 13.732 | 9.080 | 0.661 | 0.309 | 1.450 | 0.874 | 0.913 | 0.701 | 0.920 |
| 0.114 | 13.733 | 9.438 | 0.687 | 0.298 | 1.486 | 0.895 | 0.928 | 0.691 | 0.933 |
| 0.119 | 13.735 | 9.788 | 0.713 | 0.287 | 1.519 | 0.915 | 0.943 | 0.682 | 0.945 |
| 0.125 | 13.736 | 10.094 | 0.735 | 0.278 | 1.547 | 0.932 | 0.954 | 0.675 | 0.956 |
| 0.130 | 13.735 | 10.326 | 0.752 | 0.272 | 1.567 | 0.944 | 0.963 | 0.669 | 0.963 |
| 0.135 | 13.734 | 10.548 | 0.768 | 0.266 | 1.586 | 0.956 | 0.971 | 0.664 | 0.971 |
| 0.141 | 13.734 | 10.751 | 0.783 | 0.261 | 1.604 | 0.966 | 0.978 | 0.659 | 0.978 |
| 0.146 | 13.734 | 10.897 | 0.793 | 0.258 | 1.617 | 0.974 | 0.983 | 0.656 | 0.983 |
| 0.151 | 13.734 | 11.030 | 0.803 | 0.255 | 1.629 | 0.981 | 0.988 | 0.653 | 0.988 |
| 0.157 | 13.735 | 11.105 | 0.808 | 0.253 | 1.636 | 0.985 | 0.991 | 0.651 | 0.990 |
| 0.162 | 13.736 | 11.201 | 0.815 | 0.251 | 1.644 | 0.990 | 0.994 | 0.649 | 0.993 |
| 0.168 | 13.736 | 11.247 | 0.819 | 0.250 | 1.647 | 0.992 | 0.995 | 0.648 | 0.995 |
| 0.173 | 13.735 | 11.273 | 0.821 | 0.249 | 1.650 | 0.994 | 0.996 | 0.647 | 0.996 |
| 0.178 | 13.736 | 11.311 | 0.823 | 0.248 | 1.653 | 0.996 | 0.997 | 0.647 | 0.997 |
| 0.184 | 13.734 | 11.318 | 0.824 | 0.248 | 1.654 | 0.996 | 0.998 | 0.646 | 0.998 |
| 0.189 | 13.735 | 11.333 | 0.825 | 0.248 | 1.655 | 0.997 | 0.998 | 0.646 | 0.998 |
| 0.194 | 13.736 | 11.333 | 0.825 | 0.248 | 1.656 | 0.997 | 0.998 | 0.646 | 0.998 |
| 0.199 | 13.735 | 11.344 | 0.826 | 0.248 | 1.656 | 0.998 | 0.999 | 0.646 | 0.999 |
| 0.205 | 13.734 | 11.336 | 0.825 | 0.248 | 1.656 | 0.998 | 0.999 | 0.646 | 0.998 |
| 0.210 | 13.735 | 11.342 | 0.826 | 0.248 | 1.657 | 0.998 | 0.999 | 0.646 | 0.999 |
| 0.216 | 13.735 | 11.335 | 0.825 | 0.248 | 1.656 | 0.998 | 0.998 | 0.646 | 0.998 |
| 0.221 | 13.735 | 11.343 | 0.826 | 0.248 | 1.657 | 0.998 | 0.999 | 0.646 | 0.999 |
| 0.227 | 13.736 | 11.343 | 0.826 | 0.248 | 1.657 | 0.998 | 0.999 | 0.646 | 0.999 |
| 0.232 | 13.738 | 11.354 | 0.826 | 0.247 | 1.658 | 0.999 | 0.999 | 0.645 | 0.999 |
| 0.238 | 13.734 | 11.343 | 0.826 | 0.248 | 1.657 | 0.998 | 0.999 | 0.646 | 0.999 |
| 0.243 | 13.734 | 11.347 | 0.826 | 0.247 | 1.657 | 0.998 | 0.999 | 0.645 | 0.999 |
| 0.248 | 13.732 | 11.344 | 0.826 | 0.248 | 1.657 | 0.998 | 0.999 | 0.646 | 0.999 |
| 0.253 | 13.735 | 11.341 | 0.826 | 0.248 | 1.657 | 0.998 | 0.999 | 0.646 | 0.999 |
| 0.259 | 13.734 | 11.345 | 0.826 | 0.248 | 1.657 | 0.998 | 0.999 | 0.646 | 0.999 |
| 0.265 | 13.733 | 11.333 | 0.825 | 0.248 | 1.656 | 0.997 | 0.998 | 0.646 | 0.998 |
| 0.270 | 13.736 | 11.326 | 0.825 | 0.248 | 1.655 | 0.997 | 0.998 | 0.646 | 0.998 |
| 0.275 | 13.735 | 11.331 | 0.825 | 0.248 | 1.656 | 0.997 | 0.998 | 0.646 | 0.998 |
| 0.281 | 13.735 | 11.336 | 0.825 | 0.248 | 1.656 | 0.998 | 0.998 | 0.646 | 0.998 |
| 0.286 | 13.738 | 11.336 | 0.825 | 0.248 | 1.656 | 0.998 | 0.998 | 0.646 | 0.998 |
| 0.291 | 13.736 | 11.324 | 0.824 | 0.248 | 1.655 | 0.997 | 0.998 | 0.646 | 0.998 |
| 0.296 | 13.736 | 11.321 | 0.824 | 0.248 | 1.655 | 0.997 | 0.998 | 0.646 | 0.998 |
| 0.302 | 13.735 | 11.326 | 0.825 | 0.248 | 1.655 | 0.997 | 0.998 | 0.646 | 0.998 |
| 0.307 | 13.734 | 11.325 | 0.825 | 0.248 | 1.655 | 0.997 | 0.998 | 0.646 | 0.998 |
| 0.313 | 13.733 | 11.329 | 0.825 | 0.248 | 1.655 | 0.997 | 0.998 | 0.646 | 0.998 |
| 0.318 | 13.733 | 11.329 | 0.825 | 0.248 | 1.656 | 0.997 | 0.998 | 0.646 | 0.998 |
| 0.324 | 13.732 | 11.339 | 0.826 | 0.248 | 1.656 | 0.998 | 0.999 | 0.646 | 0.999 |
| 0.329 | 13.733 | 11.342 | 0.826 | 0.248 | 1.657 | 0.998 | 0.999 | 0.646 | 0.999 |
| 0.335 | 13.733 | 11.359 | 0.827 | 0.247 | 1.658 | 0.999 | 0.999 | 0.645 | 0.999 |
| 0.340 | 13.734 | 11.371 | 0.828 | 0.247 | 1.659 | 0.999 | 1.000 | 0.645 | 1.000 |
| 0.345 | 13.734 | 11.361 | 0.827 | 0.247 | 1.658 | 0.999 | 0.999 | 0.645 | 0.999 |
| 0.350 | 13.734 | 11.368 | 0.828 | 0.247 | 1.659 | 0.999 | 1.000 | 0.645 | 1.000 |
| 0.356 | 13.734 | 11.358 | 0.827 | 0.247 | 1.658 | 0.999 | 0.999 | 0.645 | 0.999 |
| 0.361 | 13.736 | 11.354 | 0.827 | 0.247 | 1.658 | 0.999 | 0.999 | 0.645 | 0.999 |
| 0.367 | 13.734 | 11.361 | 0.827 | 0.247 | 1.658 | 0.999 | 0.999 | 0.645 | 0.999 |
| 0.372 | 13.735 | 11.360 | 0.827 | 0.247 | 1.658 | 0.999 | 0.999 | 0.645 | 0.999 |


| 0.377 | 13.733 | 11.366 | 0.828 | 0.247 | 1.659 | 0.999 | 0.999 | 0.645 | 0.999 |
| ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: |
| 0.383 | 13.733 | 11.380 | 0.829 | 0.247 | 1.660 | 1.000 | 1.000 | 0.645 | 1.000 |
| 0.388 | 13.731 | 11.389 | 0.829 | 0.247 | 1.661 | 1.000 | 1.000 | 0.645 | 1.000 |
| 0.393 | 13.732 | 11.402 | 0.830 | 0.246 | 1.662 | 1.001 | 1.001 | 0.644 | 1.001 |
| 0.399 | 13.729 | 11.450 | 0.834 | 0.245 | 1.666 | 1.003 | 1.002 | 0.643 | 1.002 |
| 0.404 | 13.731 | 11.498 | 0.837 | 0.244 | 1.670 | 1.006 | 1.004 | 0.642 | 1.004 |
| 0.410 | 13.733 | 11.528 | 0.839 | 0.244 | 1.672 | 1.007 | 1.005 | 0.641 | 1.005 |
| 0.415 | 13.733 | 11.541 | 0.840 | 0.243 | 1.673 | 1.008 | 1.005 | 0.641 | 1.005 |
| 0.420 | 13.732 | 11.542 | 0.841 | 0.243 | 1.673 | 1.008 | 1.005 | 0.641 | 1.005 |
| 0.426 | 13.731 | 11.558 | 0.842 | 0.243 | 1.675 | 1.009 | 1.006 | 0.641 | 1.006 |
| 0.431 | 13.732 | 11.552 | 0.841 | 0.243 | 1.674 | 1.009 | 1.006 | 0.641 | 1.006 |
| 0.437 | 13.734 | 11.566 | 0.842 | 0.243 | 1.675 | 1.009 | 1.006 | 0.641 | 1.006 |
| 0.442 | 13.732 | 11.563 | 0.842 | 0.243 | 1.675 | 1.009 | 1.006 | 0.641 | 1.006 |
| 0.447 | 13.733 | 11.574 | 0.843 | 0.243 | 1.676 | 1.010 | 1.006 | 0.641 | 1.007 |
| 0.452 | 13.729 | 11.574 | 0.843 | 0.243 | 1.676 | 1.010 | 1.006 | 0.641 | 1.007 |
| 0.458 | 13.732 | 11.585 | 0.844 | 0.242 | 1.677 | 1.010 | 1.007 | 0.640 | 1.007 |
| 0.463 | 13.731 | 11.586 | 0.844 | 0.242 | 1.677 | 1.010 | 1.007 | 0.640 | 1.007 |
| 0.469 | 13.733 | 11.599 | 0.845 | 0.242 | 1.678 | 1.011 | 1.007 | 0.640 | 1.007 |
| 0.474 | 13.733 | 11.606 | 0.845 | 0.242 | 1.679 | 1.011 | 1.007 | 0.640 | 1.008 |
| 0.480 | 13.733 | 11.611 | 0.845 | 0.242 | 1.679 | 1.012 | 1.007 | 0.640 | 1.008 |
| 0.485 | 13.731 | 11.624 | 0.847 | 0.242 | 1.680 | 1.012 | 1.008 | 0.639 | 1.008 |
| 0.490 | 13.734 | 11.620 | 0.846 | 0.242 | 1.680 | 1.012 | 1.008 | 0.640 | 1.008 |
| 0.495 | 13.733 | 11.629 | 0.847 | 0.241 | 1.681 | 1.013 | 1.008 | 0.639 | 1.008 |
| 0.501 | 13.734 | 11.646 | 0.848 | 0.241 | 1.682 | 1.013 | 1.009 | 0.639 | 1.009 |
| 0.506 | 13.731 | 11.646 | 0.848 | 0.241 | 1.682 | 1.013 | 1.009 | 0.639 | 1.009 |
| 0.512 | 13.733 | 11.658 | 0.849 | 0.241 | 1.683 | 1.014 | 1.009 | 0.639 | 1.009 |
| 0.517 | 13.735 | 11.665 | 0.849 | 0.241 | 1.684 | 1.014 | 1.009 | 0.639 | 1.010 |

Table C-27: Mach 1.7 FPG pressure data

| Y (in) | $\mathrm{P}_{11}$ (psia) | $\mathrm{P}_{12}$ (psia) | $\mathrm{P}_{\mathrm{T} 2} / \mathrm{P}_{\mathrm{Tl}}$ | $\mathrm{pc} / \mathrm{P}_{\mathrm{T} 2}$ | MACH | T (K) | $\mathrm{P}(\mathrm{Pa})$ | $\rho\left(\mathrm{kg} / \mathrm{m}^{3}\right)$ | $\rho / \rho e$ | M/Me | $\underline{u} / \mathrm{U}_{\mathrm{e}}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.063 | 13.711 | 5.698 | 0.416 | 0.418 | 1.195 | 230.2 | 10372 | 0.157 | 0.476 | 0.675 | 0.760 |
| 0.068 | 13.712 | 6.052 | 0.441 | 0.394 | 1.244 | 225.9 | 10895 | 0.168 | 0.509 | 0.703 | 0.784 |
| 0.074 | 13.713 | 6.408 | 0.467 | 0.372 | 1.292 | 221.8 | 11397 | 0.179 | 0.543 | 0.730 | 0.806 |
| 0.080 | 13.714 | 6.725 | 0.490 | 0.354 | 1.333 | 218.3 | 11965 | 0.191 | 0.579 | 0.753 | 0.825 |
| 0.086 | 13.712 | 7.020 | 0.512 | 0.339 | 1.370 | 215.2 | 12490 | 0.202 | 0.613 | 0.774 | 0.842 |
| 0.092 | 13.714 | 7.303 | 0.533 | 0.326 | 1.404 | 212.3 | 13027 | 0.214 | 0.648 | 0.793 | 0.857 |
| 0.098 | 13.712 | 7.556 | 0.551 | 0.315 | 1.434 | 209.7 | 13546 | 0.225 | 0.682 | 0.810 | 0.870 |
| 0.104 | 13.712 | 7.874 | 0.574 | 0.303 | 1.470 | 206.6 | 14042 | 0.237 | 0.718 | 0.830 | 0.885 |
| 0.110 | 13.710 | 8.182 | 0.597 | 0.291 | 1.504 | 203.7 | 14570 | 0.249 | 0.755 | 0.850 | 0.899 |
| 0.115 | 13.711 | 8.424 | 0.614 | 0.283 | 1.531 | 201.5 | 15081 | 0.261 | 0.791 | 0.865 | 0.910 |
| 0.122 | 13.710 | 8.688 | 0.634 | 0.274 | 1.559 | 199.2 | 15629 | 0.273 | 0.829 | 0.881 | 0.921 |
| 0.127 | 13.713 | 8.875 | 0.647 | 0.268 | 1.579 | 197.5 | 16134 | 0.285 | 0.863 | 0.892 | 0.929 |
| 0.133 | 13.712 | 9.136 | 0.666 | 0.261 | 1.605 | 195.3 | 16674 | 0.298 | 0.902 | 0.907 | 0.940 |
| 0.140 | 13.713 | 9.422 | 0.687 | 0.253 | 1.634 | 192.9 | 17222 | 0.311 | 0.943 | 0.923 | 0.951 |
| 0.146 | 13.710 | 9.669 | 0.705 | 0.246 | 1.659 | 190.9 | 17758 | 0.324 | 0.983 | 0.937 | 0.960 |
| 0.151 | 13.708 | 9.740 | 0.711 | 0.245 | 1.665 | 190.3 | 18281 | 0.335 | 1.015 | 0.941 | 0.962 |
| 0.157 | 13.710 | 9.948 | 0.726 | 0.239 | 1.687 | 188.6 | 18783 | 0.347 | 1.052 | 0.953 | 0.970 |
| 0.163 | 13.708 | 10.101 | 0.737 | 0.236 | 1.702 | 187.4 | 19308 | 0.359 | 1.088 | 0.962 | 0.976 |
| 0.169 | 13.710 | 10.260 | 0.748 | 0.232 | 1.717 | 186.2 | 19830 | 0.371 | 1.125 | 0.970 | 0.981 |
| 0.175 | 13.710 | 10.414 | 0.760 | 0.229 | 1.732 | 185.0 | 20349 | 0.383 | 1.162 | 0.978 | 0.987 |
| 0.181 | 13.711 | 10.491 | 0.765 | 0.227 | 1.738 | 184.5 | 20883 | 0.394 | 1.196 | 0.982 | 0.989 |
| 0.187 | 13.712 | 10.567 | 0.771 | 0.225 | 1.745 | 183.9 | 21408 | 0.406 | 1.229 | 0.986 | 0.991 |
| 0.193 | 13.710 | 10.631 | 0.775 | 0.224 | 1.751 | 183.4 | 21954 | 0.417 | 1.264 | 0.989 | 0.993 |
| 0.199 | 13.709 | 10.646 | 0.777 | 0.224 | 1.754 | 183.2 | 22496 | 0.428 | 1.297 | 0.991 | 0.994 |
| 0.205 | 13.711 | 10.652 | 0.777 | 0.224 | 1.756 | 183.1 | 23015 | 0.438 | 1.328 | 0.992 | 0.995 |
| 0.210 | 13.709 | 10.670 | 0.778 | 0.223 | 1.757 | 183.0 | 23523 | 0.448 | 1.358 | 0.993 | 0.996 |
| 0.216 | 13.712 | 10.695 | 0.780 | 0.223 | 1.759 | 182.8 | 24045 | 0.458 | 1.389 | 0.994 | 0.996 |
| 0.222 | 13.708 | 10.707 | 0.781 | 0.222 | 1.761 | 182.7 | 24559 | 0.468 | 1.420 | 0.995 | 0.997 |
| 0.228 | 13.708 | 10.731 | 0.783 | 0.222 | 1.763 | 182.5 | 25104 | 0.479 | 1.453 | 0.996 | 0.997 |
| 0.234 | 13.710 | 10.726 | 0.782 | 0.222 | 1.763 | 182.5 | 25615 | 0.489 | 1.483 | 0.996 | 0.998 |
| 0.240 | 13.709 | 10.753 | 0.784 | 0.222 | 1.765 | 182.3 | 26149 | 0.500 | 1.515 | 0.997 | 0.998 |
| 0.246 | 13.709 | 10.753 | 0.784 | 0.222 | 1.765 | 182.3 | 26700 | 0.510 | 1.547 | 0.997 | 0.998 |
| 0.252 | 13.710 | 10.762 | 0.785 | 0.221 | 1.766 | 182.2 | 27236 | 0.521 | 1.579 | 0.998 | 0.999 |
| 0.258 | 13.711 | 10.769 | 0.785 | 0.221 | 1.767 | 182.2 | 27767 | 0.531 | 1.610 | 0.998 | 0.999 |
| 0.264 | 13.710 | 10.774 | 0.786 | 0.221 | 1.767 | 182.1 | 28289 | 0.541 | 1.641 | 0.999 | 0.999 |
| 0.270 | 13.712 | 10.787 | 0.787 | 0.221 | 1.768 | 182.1 | 28806 | 0.551 | 1.671 | 0.999 | 0.999 |
| 0.276 | 13.712 | 10.793 | 0.787 | 0.221 | 1.769 | 182.0 | 29331 | 0.562 | 1.702 | 0.999 | 1.000 |
| 0.282 | 13.711 | 10.791 | 0.787 | 0.221 | 1.769 | 182.0 | 29859 | 0.572 | 1.733 | 1.000 | 1.000 |
| 0.288 | 13.712 | 10.785 | 0.787 | 0.221 | 1.769 | 182.0 | 30381 | 0.582 | 1.763 | 0.999 | 1.000 |
| 0.294 | 13.711 | 10.799 | 0.788 | 0.221 | 1.770 | 181.9 | 30929 | 0.592 | 1.796 | 1.000 | 1.000 |
| 0.300 | 13.713 | 10.804 | 0.788 | 0.221 | 1.770 | 181.9 | 31469 | 0.603 | 1.827 | 1.000 | 1.000 |
| 0.306 | 13.712 | 10.803 | 0.788 | 0.221 | 1.770 | 181.9 | 32005 | 0.613 | 1.859 | 1.000 | 1.000 |
| 0.312 | 13.712 | 10.806 | 0.788 | 0.220 | 1.771 | 181.9 | 32551 | 0.624 | 1.890 | 1.000 | 1.000 |
| 0.318 | 13.708 | 10.805 | 0.788 | 0.220 | 1.771 | 181.9 | 33056 | 0.633 | 1.920 | 1.000 | 1.000 |
| 0.324 | 13.708 | 10.805 | 0.788 | 0.220 | 1.771 | 181.9 | 33569 | 0.643 | 1.949 | 1.000 | 1.000 |
| 0.329 | 13.709 | 10.810 | 0.789 | 0.220 | 1.771 | 181.8 | 34083 | 0.653 | 1.980 | 1.001 | 1.000 |
| 0.335 | 13.708 | 10.810 | 0.789 | 0.220 | 1.771 | 181.8 | 34625 | 0.663 | 2.011 | 1.001 | 1.000 |
| 0.341 | 13.709 | 10.807 | 0.788 | 0.220 | 1.771 | 181.9 | 35156 | 0.674 | 2.042 | 1.000 | 1.000 |
| 0.347 | 13.712 | 10.819 | 0.789 | 0.220 | 1.772 | 181.8 | 35684 | 0.684 | 2.074 | 1.001 | 1.001 |
| 0.353 | 13.712 | 10.822 | 0.789 | 0.220 | 1.772 | 181.8 | 36221 | 0.694 | 2.105 | 1.001 | 1.001 |
| 0.360 | 13.709 | 10.814 | 0.789 | 0.220 | 1.771 | 181.8 | 36766 | 0.705 | 2.136 | 1.001 | 1.001 |
| 0.366 | 13.711 | 10.817 | 0.789 | 0.220 | 1.772 | 181.8 | 37303 | 0.715 | 2.167 | 1.001 | 1.001 |
| 0.371 | 13.711 | 10.823 | 0.789 | 0.220 | 1.772 | 181.8 | 37828 | 0.725 | 2.198 | 1.001 | 1.001 |
| 0.377 | 13.709 | 10.824 | 0.790 | 0.220 | 1.772 | 181.7 | 38318 | 0.735 | 2.227 | 1.001 | 1.001 |
| 0.383 | 13.712 | 10.837 | 0.790 | 0.220 | 1.774 | 181.7 | 38849 | 0.745 | 2.259 | 1.002 | 1.001 |
| 0.389 | 13.712 | 10.832 | 0.790 | 0.220 | 1.773 | 181.7 | 39363 | 0.755 | 2.288 | 1.002 | 1.001 |


| 0.395 | 13.714 | 10.837 | 0.790 | 0.220 | 1.774 | 181.7 | 39902 | 0.765 | 2.320 | 1.002 | 1.001 |
| ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: |
| 0.401 | 13.712 | 10.835 | 0.790 | 0.220 | 1.773 | 181.7 | 40430 | 0.775 | 2.351 | 1.002 | 1.001 |
| 0.407 | 13.713 | 10.849 | 0.791 | 0.220 | 1.775 | 181.6 | 40952 | 0.786 | 2.383 | 1.003 | 1.002 |
| 0.413 | 13.711 | 10.845 | 0.791 | 0.220 | 1.774 | 181.6 | 41509 | 0.796 | 2.415 | 1.002 | 1.002 |
| 0.419 | 13.709 | 10.856 | 0.792 | 0.219 | 1.775 | 181.5 | 42040 | 0.807 | 2.446 | 1.003 | 1.002 |
| 0.425 | 13.712 | 10.848 | 0.791 | 0.220 | 1.775 | 181.6 | 42580 | 0.817 | 2.477 | 1.003 | 1.002 |
| 0.431 | 13.712 | 10.854 | 0.792 | 0.219 | 1.775 | 181.5 | 43084 | 0.827 | 2.507 | 1.003 | 1.002 |
| 0.437 | 13.712 | 10.846 | 0.791 | 0.220 | 1.774 | 181.6 | 43621 | 0.837 | 2.537 | 1.002 | 1.002 |
| 0.442 | 13.712 | 10.850 | 0.791 | 0.220 | 1.775 | 181.6 | 44126 | 0.847 | 2.567 | 1.003 | 1.002 |
| 0.448 | 13.711 | 10.853 | 0.792 | 0.219 | 1.775 | 181.5 | 44671 | 0.857 | 2.599 | 1.003 | 1.002 |
| 0.454 | 13.712 | 10.858 | 0.792 | 0.219 | 1.776 | 181.5 | 45202 | 0.868 | 2.631 | 1.003 | 1.002 |
| 0.460 | 13.711 | 10.863 | 0.792 | 0.219 | 1.776 | 181.5 | 45727 | 0.878 | 2.662 | 1.003 | 1.002 |
| 0.466 | 13.713 | 10.868 | 0.793 | 0.219 | 1.776 | 181.4 | 46273 | 0.889 | 2.694 | 1.004 | 1.002 |
| 0.473 | 13.709 | 10.878 | 0.793 | 0.219 | 1.777 | 181.4 | 46815 | 0.899 | 2.727 | 1.004 | 1.003 |
| 0.478 | 13.709 | 10.882 | 0.794 | 0.219 | 1.778 | 181.3 | 47334 | 0.910 | 2.758 | 1.004 | 1.003 |
| 0.484 | 13.711 | 10.889 | 0.794 | 0.219 | 1.778 | 181.3 | 47859 | 0.920 | 2.789 | 1.005 | 1.003 |
| 0.490 | 13.709 | 10.894 | 0.795 | 0.219 | 1.779 | 181.2 | 48376 | 0.930 | 2.819 | 1.005 | 1.003 |
| 0.496 | 13.710 | 10.894 | 0.795 | 0.219 | 1.779 | 181.2 | 48881 | 0.940 | 2.849 | 1.005 | 1.003 |
| 0.502 | 13.712 | 10.909 | 0.796 | 0.218 | 1.780 | 181.1 | 49426 | 0.951 | 2.882 | 1.006 | 1.004 |
| 0.508 | 13.710 | 10.914 | 0.796 | 0.218 | 1.781 | 181.1 | 49963 | 0.961 | 2.914 | 1.006 | 1.004 |
| 0.514 | 13.710 | 10.916 | 0.796 | 0.218 | 1.781 | 181.1 | 50494 | 0.972 | 2.946 | 1.006 | 1.004 |
| 0.520 | 13.710 | 10.917 | 0.796 | 0.218 | 1.781 | 181.1 | 51042 | 0.982 | 2.978 | 1.006 | 1.004 |
| 0.526 | 13.708 | 10.923 | 0.797 | 0.218 | 1.782 | 181.0 | 51570 | 0.993 | 3.009 | 1.007 | 1.004 |
| 0.532 | 13.712 | 10.943 | 0.798 | 0.218 | 1.783 | 180.9 | 52086 | 1.003 | 3.042 | 1.008 | 1.005 |
| 0.538 | 13.708 | 10.938 | 0.798 | 0.218 | 1.783 | 180.9 | 52608 | 1.013 | 3.072 | 1.007 | 1.004 |
| 0.543 | 13.709 | 10.944 | 0.798 | 0.218 | 1.783 | 180.9 | 53110 | 1.023 | 3.102 | 1.008 | 1.005 |

Table C-28: Mach 1.7 ZPG hot-film data - traverse up

| Multiple overheat data |  |  |  |  |  | Single overheat data |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $y$ (in) | $T_{1} / T_{11}$ | $\mathrm{Re}^{1 / 2}$ | ( pu$)_{\text {mms }}$ | $\left(T_{1}\right)_{\text {mim }}$ | $\mathrm{pu} / \mathrm{pu} \mathrm{u}_{x}$ | $\mathrm{Re}^{1 / 2}$ | ( $\rho$ u) ${ }_{\text {mis }}$ | ¢u/ $\mathrm{u}_{x}$ | $y / \delta_{u}$ |
| 0.0183 | 1.004 | 16.692 | $4.68 \mathrm{E}-02$ | 3.91E-02 | 0.623 | 16.744 | 0.0527 | 0.6304 | 0.1169 |
| 0.0293 | 0.998 | 16.968 | 3.30E-02 | 6.23E-02 | 0.644 | 16.983 | 0.05342 | 0.6486 | 0.1871 |
| 0.0401 | 0.999 | 17.303 | 5.81E-02 | 4.09E-02 | 0.669 | 17.325 | 0.06199 | 0.6750 | 0.2561 |
| 0.0509 | 1.003 | 17.759 | 5.59E-02 | 5.74E-02 | 0.705 | 17.763 | 0.06576 | 0.7095 | 0.3250 |
| 0.0617 | 1.008 | 18.238 | $6.36 \mathrm{E}-02$ | 4.87E-02 | 0.744 | 18.245 | 0.06964 | 0.7485 | 0.3940 |
| 0.0723 | 1.003 | 18.681 | $5.35 \mathrm{E}-02$ | 5.63E-02 | 0.780 | 18.671 | 0.06529 | 0.7839 | 0.4617 |
| 0.0833 | 1.003 | 19.063 | 5.27E-02 | 5.10E-02 | 0.812 | 19.078 | 0.06156 | 0.8185 | 0.5319 |
| 0.0941 | 1.004 | 19.424 | 5.08E-02 | 4.43E-02 | 0.843 | 19.419 | 0.05845 | 0.8480 | 0.6009 |
| 0.105 | 1.007 | 19.743 | $4.40 \mathrm{E}-02$ | 4.62E-02 | 0.871 | 19.759 | 0.05405 | 0.8779 | 0.6705 |
| 0.1163 | 1.009 | 20.016 | $4.40 \mathrm{E}-02$ | 3.62E-02 | 0.896 | 20.037 | 0.05035 | 0.9028 | 0.7427 |
| 0.1273 | 1.006 | 20.251 | 3.57E-02 | 3.58E-02 | 0.917 | 20.272 | 0.04424 | 0.9241 | 0.8129 |
| 0.1381 | 1.007 | 20.446 | 2.86E-02 | 3.67E-02 | 0.935 | 20.455 | 0.03944 | 0.9409 | 0.8819 |
| 0.1489 | 1.003 | 20.598 | 3.01E-02 | 2.54E-02 | 0.948 | 20.6 | 0.03542 | 0.9543 | 0.9508 |
| 0.16 | 1.001 | 20.719 | 3.54E-02 | 1.54E-02 | 0.960 | 20.697 | 0.03305 | 0.9633 | 1.0217 |
| 0.1711 | 1 | 20.803 | 2.56E-02 | $1.60 \mathrm{E}-02$ | 0.967 | 20.778 | 0.02836 | 0.9708 | 1.0926 |
| 0.1821 | 0.999 | 20.867 | 9.37E-03 | 3.23E-02 | 0.973 | 20.848 | 0.02521 | 0.9774 | 1.1628 |
| 0.1929 | 0.999 | 20.91 | 1.89E-02 | 2.29E-02 | 0.977 | 20.897 | 0.02443 | 0.9820 | 1.2318 |
| 0.2038 | 0.999 | 20.937 | 1.86E-02 | 1.85E-02 | 0.980 | 20.914 | 0.02295 | 0.9836 | 1.3014 |
| 0.215 | 0.998 | 20.97 | 1.67E-02 | 1.81E-02 | 0.983 | 20.943 | 0.02137 | 0.9863 | 1.3729 |
| 0.2261 | 0.998 | 20.988 | 1.56E-02 | 1.46E-02 | 0.985 | 20.951 | 0.01952 | 0.9870 | 1.4438 |
| 0.237 | 1 | 21 | 1.92E-02 | 7.79E-03 | 0.986 | 20.97 | 0.01922 | 0.9888 | 1.5134 |
| 0.2477 | 1.002 | 21.008 | 2.22E-02 | 1.17E-02 | 0.987 | 20.992 | 0.01946 | 0.9909 | 1.5817 |
| 0.2588 | 1.007 | 21.022 | 1.86E-02 | 6.26E-03 | 0.988 | 21.015 | 0.01786 | 0.9931 | 1.6526 |
| 0.2699 | 1.01 | 21.038 | $8.24 \mathrm{E}-03$ | 1.89E-02 | 0.989 | 21.041 | 0.01601 | 0.9955 | 1.7235 |
| 0.2808 | 1.012 | 21.048 | 1.20E-02 | 1.60E-02 | 0.990 | 21.05 | 0.01633 | 0.9964 | 1.7931 |
| 0.2916 | 1.014 | 21.064 | 1.43E-02 | 4.03E-03 | 0.992 | 21.073 | 0.01584 | 0.9986 | 1.8621 |
| 0.3024 | 1.015 | 21.063 | 7.52E-03 | 1.62E-02 | 0.992 | 21.077 | 0.01526 | 0.9990 | 1.9310 |
| 0.3136 | 1.017 | 21.067 | $4.90 \mathrm{E}-03$ | 1.71E-02 | 0.992 | 21.098 | 0.01508 | 1.0009 | 2.0026 |
| 0.3246 | 1.017 | 21.071 | $1.88 \mathrm{E}-02$ | 1.31E-02 | 0.993 | 21.111 | 0.01553 | 1.0022 | 2.0728 |
| 0.3355 | 1.014 | 21.082 | 1.69E-02 | 9.18E-03 | 0.994 | 21.121 | 0.01561 | 1.0031 | 2.1424 |
| 0.3462 | 1.011 | 21.097 | 1.39E-02 | 1.51E-03 | 0.995 | 21.125 | 0.01531 | 1.0035 | 2.2107 |
| 0.3573 | 1.008 | 21.122 | 1.45E-02 | 5.32E-03 | 0.997 | 21.142 | 0.01553 | 1.0051 | 2.2816 |
| 0.3684 | 1.006 | 21.143 | 1.68E-02 | $1.28 \mathrm{E}-02$ | 0.999 | 21.154 | 0.01544 | 1.0063 | 2.3525 |
| 0.3793 | 1.005 | 21.157 | 2.07E-02 | 1.51E-02 | 1.001 | 21.162 | 0.01691 | 1.0070 | 2.4221 |
| 0.3901 | 1.004 | 21.166 | 1.33E-02 | 7.34E-03 | 1.002 | 21.163 | 0.0148 | 1.0071 | 2.4911 |
| 0.4009 | 1.005 | 21.173 | 1.27E-02 | $5.21 \mathrm{E}-03$ | 1.002 | 21.174 | 0.01465 | 1.0082 | 2.5600 |
| 0.4121 | 1.005 | 21.175 | 2.36E-02 | 2.28E-02 | 1.002 | 21.174 | 0.01635 | 1.0082 | 2.6315 |
| 0.423 | 1.004 | 21.172 | $1.97 \mathrm{E}-02$ | $1.86 \mathrm{E}-02$ | 1.002 | 21.165 | 0.01508 | 1.0073 | 2.7011 |
| 0.434 | 1.004 | 21.166 | 2.19E-02 | 2.05E-02 | 1.002 | 21.154 | 0.01599 | 1.0063 | 2.7714 |
| 0.4448 | 1.003 | 21.157 | 1.45E-02 | 8.83E-03 | 1.001 | 21.134 | 0.01452 | 1.0044 | 2.8404 |
| 0.4558 | 1.001 | 21.152 | 9.35E-03 | 1.29E-02 | 1.000 | 21.117 | 0.01491 | 1.0028 | 2.9106 |
| 0.4668 | 1 | 21.147 | 2.13E-02 | 2.26E-02 | 1.000 | 21.105 | 0.01507 | 1.0016 | 2.9808 |
| 0.4777 | 0.999 | 21.142 | $1.46 \mathrm{E}-02$ | 4.61E-03 | 0.999 | 21.091 | 0.01468 | 1.0003 | 3.0504 |
| 0.4858 | 0.998 | 21.148 | 9.51E-03 | 8.65E-03 | 1.000 | 21.092 | 0.01363 | 1.0004 | 3.1022 |
| 0.4867 | 0.997 | 21.15 | 1.14E-02 | 1.20E-02 | 1.000 | 21.088 | 0.01492 | 1.0000 | 3.1079 |

Table C-29: Mach 1.7 ZPG hot-film data - traverse down

| $y$ (in) | $\mathrm{T}_{1} / \mathrm{T}_{11}$ | $\mathrm{Re}^{7 / 2}$ | ( L$)_{\text {mis }}$ | $\left(T_{t}\right)_{\text {ms }}$ | pu/ $\mathrm{pu}_{x}$ | $\mathrm{Re}^{1 / 2}$ | (pu) ${ }_{\text {ms }}$ | pu / pux | $y / \delta_{u}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.0194 | 0.984 | 16.687 | 2.63E-02 | 6.54E-02 | 0.6249 | 16.677 | 0.05189 | 0.6263 | 0.1239 |
| 0.0306 | 0.982 | 16.968 | 4.11E-02 | 5.58E-02 | 0.6453 | 16.94 | 0.05576 | 0.6462 | 0.1954 |
| 0.0414 | 0.976 | 17.321 | 3.77E-02 | $7.68 \mathrm{E}-02$ | 0.6696 | 17.249 | 0.06129 | 0.6700 | 0.2644 |
| 0.052 | 0.981 | 17.754 | 5.02E-02 | 7.58E-02 | 0.7059 | 17.703 | 0.06851 | 0.7057 | 0.3321 |
| 0.0632 | 0.984 | 18.214 | 5.33E-02 | 6.97E-02 | 0.7445 | 18.16 | 0.06788 | 0.7426 | 0.4036 |
| 0.0745 | 0.981 | 18.659 | 5.87E-02 | $5.28 \mathrm{E}-02$ | 0.7797 | 18.589 | 0.0668 | 0.7781 | 0.4757 |
| 0.0855 | 0.985 | 19.078 | $5.59 \mathrm{E}-02$ | $5.18 \mathrm{E}-02$ | 0.8174 | 19.027 | 0.06365 | 0.8152 | 0.5460 |
| 0.0962 | 0.98 | 19.434 | 5.35E-02 | 4.96E-02 | 0.8453 | 19.344 | 0.05954 | 0.8426 | 0.6143 |
| 0.1071 | 0.992 | 19.74 | 5.15E-02 | 3.95E-02 | 0.8792 | 19.7 | 0.05594 | 0.8739 | 0.6839 |
| 0.1184 | 0.986 | 20.015 | 4.98E-02 | 2.91E-02 | 0.9002 | 19.962 | 0.05084 | 0.8973 | 0.7561 |
| 0.1295 | 0.986 | 20.217 | $4.30 \mathrm{E}-02$ | 2.99E-02 | 0.9185 | 20.138 | 0.04596 | 0.9132 | 0.8269 |
| 0.1403 | 0.986 | 20.426 | 4.02E-02 | 2.83E-02 | 0.9376 | 20.349 | 0.04231 | 0.9325 | 0.8959 |
| 0.1506 | 0.985 | 20.562 | 3.56E-02 | $2.50 \mathrm{E}-02$ | 0.9495 | 20.464 | 0.03733 | 0.9430 | 0.9617 |
| 0.1611 | 0.983 | 20.679 | 3.67E-02 | 1.43E-02 | 0.9590 | 20.557 | 0.03516 | 0.9516 | 1.0287 |
| 0.1724 | 0.99 | 20.775 | 3.22E-02 | $4.30 \mathrm{E}-03$ | 0.9725 | 20.689 | 0.03019 | 0.9639 | 1.1009 |
| 0.1834 | 0.99 | 20.833 | 2.92E-02 | 9.37E-03 | 0.9780 | 20.757 | 0.02689 | 0.9702 | 1.1711 |
| 0.194 | 0.991 | 20.876 | 2.85E-02 | 1.42E-02 | 0.9827 | 20.796 | 0.02601 | 0.9739 | 1.2388 |
| 0.2048 | 0.989 | 20.926 | 2.68E-02 | 1.45E-02 | 0.9861 | 20.846 | 0.02375 | 0.9786 | 1.3078 |
| 0.2162 | 0.988 | 20.953 | 2.74E-02 | 2.17E-02 | 0.9879 | 20.868 | 0.02244 | 0.9806 | 1.3806 |
| 0.2273 | 0.988 | 20.974 | 2.59E-02 | 1.89E-02 | 0.9899 | 20.897 | 0.02091 | 0.9834 | 1.4515 |
| 0.2382 | 0.993 | 20.983 | $2.20 \mathrm{E}-02$ | $6.81 \mathrm{E}-03$ | 0.9941 | 20.925 | 0.01995 | 0.9860 | 1.5211 |
| 0.2488 | 0.996 | 20.996 | 2.06E-02 | 3.47E-03 | 0.9973 | 20.948 | 0.019 | 0.9882 | 1.5888 |
| 0.2599 | 0.998 | 21.007 | 1.12E-02 | 1.72E-02 | 0.9997 | 20.965 | 0.01705 | 0.9898 | 1.6596 |
| 0.2711 | 1.001 | 21.018 | 1.39E-02 | 1.40E-02 | 1.0028 | 20.983 | 0.01707 | 0.9915 | 1.7312 |
| 0.2821 | 1.004 | 21.027 | 1.57E-02 | 1.09E-02 | 1.0056 | 21.003 | 0.0168 | 0.9934 | 1.8014 |
| 0.2928 | 1.006 | 21.036 | 1.17E-02 | 1.44E-02 | 1.0078 | 21.02 | 0.016 | 0.9950 | 1.8697 |
| 0.3035 | 1.007 | 21.041 | 1.31E-02 | $1.25 \mathrm{E}-02$ | 1.0090 | 21.036 | 0.01586 | 0.9965 | 1.9381 |
| 0.3148 | 1.007 | 21.042 | 1.23E-02 | 1.05E-02 | 1.0091 | 21.041 | 0.01538 | 0.9970 | 2.0102 |
| 0.3259 | 1.008 | 21.048 | 1.11E-02 | 1.08E-02 | 1.0103 | 21.052 | 0.01429 | 0.9980 | 2.0811 |
| 0.3366 | 1.007 | 21.062 | 1.12E-02 | $1.28 \mathrm{E}-02$ | 1.0110 | 21.06 | 0.01531 | 0.9988 | 2.1494 |
| 0.3472 | 1.004 | 21.085 | 1.09E-02 | 9.28E-03 | 1.0112 | 21.073 | 0.01474 | 1.0000 | 2.2171 |
| 0.3583 | 1.001 | 21.109 | 8.94E-03 | 1.29E-02 | 1.0115 | 21.087 | 0.01398 | 1.0013 | 2.2880 |
| 0.3695 | 1.001 | 21.13 | $1.59 \mathrm{E}-02$ | 1.10E-02 | 1.0135 | 21.106 | 0.01515 | 1.0031 | 2.3595 |
| 0.3805 | 1 | 21.141 | 1.17E-02 | 5.32E-03 | 1.0139 | 21.11 | 0.014 | 1.0035 | 2.4298 |
| 0.3912 | 1 | 21.152 | 1.43E-02 | 6.80E-03 | 1.0149 | 21.12 | 0.01469 | 1.0045 | 2.4981 |
| 0.4019 | 1 | 21.176 | 3.56E-02 | 5.07E-02 | 1.0172 | 21.119 | 0.01383 | 1.0044 | 2.5664 |
| 0.4133 | 1.001 | 21.207 | 5.92E-02 | 8.20E-02 | 1.0209 | 21.13 | 0.01489 | 1.0054 | 2.6392 |
| 0.4244 | 1.006 | 21.44 | 1.20E-01 | 1.59E-01 | 1.0469 | 21.13 | 0.01473 | 1.0054 | 2.7101 |
| 0.4352 | 1.01 | 21.637 | $3.47 \mathrm{E}-02$ | $4.66 \mathrm{E}-02$ | 1.0691 | 21.117 | 0.01358 | 1.0042 | 2.7791 |
| 0.4458 | 1.009 | 21.612 | $3.57 \mathrm{E}-02$ | $4.80 \mathrm{E}-02$ | 1.0659 | 21.105 | 0.01352 | 1.0030 | 2.8467 |
| 0.4568 | 1.022 | 21.694 | 3.70E-02 | 7.04E-02 | 1.0832 | 21.101 | 0.01443 | 1.0027 | 2.9170 |
| 0.468 | 1.044 | 21.852 | 5.39E-02 | 1.00E-01 | 1.1147 | 21.09 | 0.01345 | 1.0016 | 2.9885 |

Table C-30: Mach 1.7 ZPG separated turbulence variables - traverse down

| y (inch) | Mach | (P) mss | $\mathrm{U}_{\text {ms }}$ | p'U | u"/u | pbarubar | pubar | $\mathrm{y} / \delta_{1}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.029 | 0.97 | 3.92E-02 | $8.50 \mathrm{E}-02$ | -3.84E-03 | -2.44E-01 | 78.6 | 104.0 | 0.188 |
| 0.040 | 1.03 | 3.14E-02 | 6.43E-02 | -8.74E-04 | -2.07E-01 | 85.7 | 108.0 | 0.257 |
| 0.051 | 1.08 | $3.75 \mathrm{E}-02$ | $8.02 \mathrm{E}-02$ | -2.36E-03 | -1.91E-01 | 92.3 | 114.0 | 0.326 |
| 0.062 | 1.13 | 3.52E-02 | 7.12E-02 | -1.14E-03 | -1.86E-01 | 98.4 | 121.0 | 0.396 |
| 0.072 | 1.18 | $3.32 \mathrm{E}-02$ | 7.61E-02 | -2.02E-03 | -1.72E-01 | 105.0 | 126.0 | 0.463 |
| 0.083 | 1.26 | 3.09E-02 | 6.75E-02 | -1.37E-03 | -1.45E-01 | 113.0 | 132.0 | 0.534 |
| 0.094 | 1.34 | 2.72E-02 | 5.83E-02 | -7.78E-04 | -1.06E-01 | 122.0 | 137.0 | 0.603 |
| 0.105 | 1.43 | $2.31 \mathrm{E}-02$ | 5.72E-02 | -9.37E-04 | -6.81E-02 | 132.0 | 142.0 | 0.673 |
| 0.116 | 1.50 | $2.18 \mathrm{E}-02$ | 4.59E-02 | -3.21E-04 | -3.54E-02 | 141.0 | 146.0 | 0.744 |
| 0.127 | 1.56 | 1.45E-02 | 4.37E-02 | -4.24E-04 | -7.45E-03 | 148.0 | 149.0 | 0.814 |
| 0.138 | 1.60 | 8.52E-03 | $4.28 \mathrm{E}-02$ | -5.43E-04 | $6.78 \mathrm{E}-03$ | 153.0 | 152.0 | 0.885 |
| 0.149 | 1.62 | 1.22E-02 | 3.16E-02 | -1.19E-04 | $2.14 \mathrm{E}-02$ | 157.0 | 154.0 | 0.955 |
| 0.160 | 1.64 | 1.73E-02 | $2.40 \mathrm{E}-02$ | 1.89E-04 | $2.44 \mathrm{E}-02$ | 159.0 | 155.0 | 1.026 |
| 0.171 | 1.65 | 1.15E-02 | $2.14 \mathrm{E}-02$ | 3.11E-05 | $2.54 \mathrm{E}-02$ | 160.0 | 156.0 | 1.096 |
| 0.182 | 1.65 | $8.27 \mathrm{E}-03$ | 3.46E-02 | -5.19E-04 | $2.45 \mathrm{E}-02$ | 161.0 | 157.0 | 1.167 |
| 0.193 | 1.66 | 8.49E-03 | 2.55E-02 | -1.83E-04 | 2.05E-02 | 161.0 | 158.0 | 1.237 |
| 0.204 | 1.66 | 8.15E-03 | 2.14E-02 | -9.00E-05 | $1.82 \mathrm{E}-02$ | 161.0 | 158.0 | 1.308 |
| 0.215 | 1.66 | 7.15E-03 | $2.07 \mathrm{E}-02$ | -9.97E-05 | $1.57 \mathrm{E}-02$ | 161.0 | 159.0 | 1.378 |
| 0.226 | 1.66 | 6.20E-03 | 1.74E-02 | -4.84E-05 | $1.48 \mathrm{E}-02$ | 161.0 | 159.0 | 1.449 |
| 0.237 | 1.66 | 1.07E-02 | 1.16E-02 | 5.93E-05 | 1.21E-02 | 161.0 | 159.0 | 1.519 |
| 0.248 | 1.66 | 9.98E-03 | 1.69E-02 | 5.25E-05 | 8.11E-03 | 161.0 | 160.0 | 1.590 |
| 0.259 | 1.66 | 9.22E-03 | 1.13E-02 | $6.60 \mathrm{E}-05$ | 3.69E-04 | 161.0 | 161.0 | 1.660 |
| 0.270 | 1.66 | 3.52E-03 | 2.04E-02 | -1.68E-04 | -6.38E-03 | 160.0 | 161.0 | 1.731 |
| 0.281 | 1.66 | 4.80E-03 | 1.77E-02 | -9.62E-05 | -8.22E-03 | 160.0 | 161.0 | 1.801 |
| 0.292 | 1.66 | 4.86E-03 | 9.66E-03 | 4.42E-05 | -1.39E-02 | 160.0 | 162.0 | 1.872 |
| 0.302 | 1.66 | $5.31 \mathrm{E}-03$ | 1.80E-02 | -1.20E-04 | -1.53E-02 | 160.0 | 162.0 | 1.936 |

Table C-31: Mach 1.7 ZPG separated turbulence variables - traverse down

| y (inch) | Mach | $(\mathrm{\rho})_{\text {mas }}$ | $\mathrm{U}_{\text {ms }}$ | $\rho \mathrm{U}^{\prime}$ | u ${ }^{1 / и}$ | pbarubar | pubar | $\mathrm{y} / \delta_{\mathrm{o}}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.031 | 0.98 | 3.61E-02 | 7.85E-02 | -2.89E-03 | -2.18E-01 | 80.1 | 103.0 | 0.196 |
| 0.041 | 1.03 | $4.73 \mathrm{E}-02$ | 1.00E-01 | -5.43E-03 | -1.77E-01 | 87.5 | 106.0 | 0.265 |
| 0.052 | 1.08 | $4.71 \mathrm{E}-02$ | 9.83E-02 | -4.68E-03 | -1.62E-01 | 94.0 | 112.0 | 0.333 |
| 0.063 | 1.14 | $4.27 \mathrm{E}-02$ | 9.01E-02 | -3.55E-03 | -1.51E-01 | 100.0 | 118.0 | 0.405 |
| 0.075 | 1.20 | 3.49E-02 | 7.21E-02 | -1.48E-03 | -1.37E-01 | 107.0 | 124.0 | 0.478 |
| 0.086 | 1.27 | $3.34 \mathrm{E}-02$ | 6.77E-02 | -1.29E-03 | -1.13E-01 | 115.0 | 130.0 | 0.548 |
| 0.096 | 1.36 | $3.26 \mathrm{E}-02$ | $6.17 \mathrm{E}-02$ | -1.00E-03 | -6.61E-02 | 125.0 | 134.0 | 0.617 |
| 0.107 | 1.44 | $2.87 \mathrm{E}-02$ | $5.01 \mathrm{E}-02$ | -3.43E-04 | -3.69E-02 | 135.0 | 140.0 | 0.686 |
| 0.118 | 1.51 | $2.77 \mathrm{E}-02$ | 3.85E-02 | 1.16E-04 | 4.19E-03 | 144.0 | 143.0 | 0.756 |
| 0.130 | 1.56 | $2.30 \mathrm{E}-02$ | 3.76E-02 | -4.82E-05 | 2.95E-02 | 150.0 | 146.0 | 0.833 |
| 0.140 | 1.60 | $2.36 \mathrm{E}-02$ | $3.37 \mathrm{E}-02$ | -3.87E-05 | 4.37E-02 | 156.0 | 149.0 | 0.897 |
| 0.151 | 1.63 | 2.04E-02 | 2.99E-02 | -1.95E-05 | 5.13E-02 | 159.0 | 151.0 | 0.968 |
| 0.161 | 1.64 | $2.08 \mathrm{E}-02$ | 2.15E-02 | 2.26E-04 | 5.52E-02 | 161.0 | 152.0 | 1.032 |
| 0.172 | 1.65 | 1.43E-02 | $1.81 \mathrm{E}-02$ | 2.52E-04 | 4.14E-02 | 161.0 | 155.0 | 1.103 |
| 0.183 | 1.65 | $1.38 \mathrm{E}-02$ | 1.80E-02 | 1.69E-04 | 3.90E-02 | 162.0 | 156.0 | 1.173 |
| 0.194 | 1.66 | $1.35 \mathrm{E}-02$ | $2.08 \mathrm{E}-02$ | 1.00E-04 | 3.47E-02 | 162.0 | 156.0 | 1.244 |
| 0.205 | 1.66 | 1.18E-02 | $2.08 \mathrm{E}-02$ | 7.29E-05 | 3.23E-02 | 162.0 | 157.0 | 1.314 |
| 0.216 | 1.66 | 1.26E-02 | $2.65 \mathrm{E}-02$ | -5.63E-05 | 3.18E-02 | 162.0 | 157.0 | 1.385 |
| 0.227 | 1.66 | 1.12E-02 | $2.40 \mathrm{E}-02$ | -1.69E-05 | 2.93E-02 | 162.0 | 157.0 | 1.455 |
| 0.238 | 1.66 | 8.99E-03 | $1.43 \mathrm{E}-02$ | $9.86 \mathrm{E}-05$ | $2.20 \mathrm{E}-02$ | 162.0 | 158.0 | 1.526 |
| 0.249 | 1.66 | 8.48E-03 | 1.22E-02 | 1.02E-04 | 1.69E-02 | 161.0 | 159.0 | 1.596 |
| 0.260 | 1.66 | $2.76 \mathrm{E}-03$ | $1.90 \mathrm{E}-02$ | -1.22E-04 | 1.27E-02 | 161.0 | 159.0 | 1.667 |
| 0.271 | 1.66 | $6.90 \mathrm{E}-03$ | 1.60E-02 | -5.42E-05 | 6.97E-03 | 161.0 | 160.0 | 1.737 |
| 0.282 | 1.66 | 8.59E-03 | $1.31 \mathrm{E}-02$ | -4.17E-07 | $2.94 \mathrm{E}-03$ | 161.0 | 160.0 | 1.808 |
| 0.293 | 1.66 | $4.04 \mathrm{E}-03$ | $1.64 \mathrm{E}-02$ | -7.38E-05 | -9.32E-04 | 160.0 | 161.0 | 1.878 |
| 0.303 | 1.66 | $5.88 \mathrm{E}-03$ | 1.46E-02 | -3.87E-05 | -2.85E-03 | 160.0 | 161.0 | 1.942 |

Table C-32: Mach 1.7 discrete turbulence intensity data points

| $y$ (in) | a | b | sqri(Re) | $f$ | Vrms | Vbar | T.I. | $y / \delta_{u}$ FPG |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.03125 | 0.1861 | 0.2613 | 13.681 | 0.226731 | 0.1254 | 5.475 | 10.102\% | 0.118 |
| 0.08125 | 0.1861 | 0.2613 | 15.499 | 0.229233 | 0.1359 | 5.787 | 10.244\% | 0.307 |
| 0.13125 | 0.1861 | 0.2613 | 16.695 | 0.230606 | 0.1303 | 5.997 | 9.422\% | 0.495 |
| 0.18125 | 0.1861 | 0.2613 | 17.556 | 0.231486 | 0.091 | 6.138 | 6.405\% | 0.684 |
| 0.23125 | 0.1861 | 0.2613 | 17.841 | 0.23176 | 0.04982 | 6.184 | 3.476\% | 0.873 |
| 0.28125 | 0.1861 | 0.2613 | 17.941 | 0.231855 | 0.03855 | 6.2 | 2.682\% | 1.061 |
| 0.33125 | 0.1861 | 0.2613 | 18.000 | 0.23191 | 0.0363 | 6.206 | 2.522\% | 1.250 |
|  |  |  | fbar | 0.230512 |  |  |  |  |
|  |  |  |  |  |  |  |  |  |
| $y$ (in) | a | b | squt(Re) | f | Vrms | Vbar | T.I. | $y / \delta_{u}$ ZPG |
| 0.03125 | 0.0402 | 0.2185 | 17.33 | 0.190311 | 0.06042 | 3.134 | 10.130\% | 0.203 |
| 0.08125 | 0.0402 | 0.2185 | 20.3 | 0.1972 | 0.05398 | 3.332 | 8.215\% | 0.528 |
| 0.13125 | 0.0402 | 0.2185 | 21.27 | 0.199118 | 0.01321 | 3.394 | 1.955\% | 0.852 |
| 0.18125 | 0.0402 | 0.2185 | 21.35 | 0.19927 | 0.008352 | 3.399 | 1.233\% | 1.177 |
|  |  |  | fbar | 0.196475 |  |  |  |  |

Table C-33: Mach 1.7 power spectra data

| frequency | $Z P \mathrm{C}, \mathrm{y} / \mathrm{d}=0.20$ | ZPG, y/d = 0.58 | $Z P G, y / d=0.85$ | FPG, y/d = 0.17 | FPG, $\mathrm{y} / \mathrm{d}=0.44$ | FPG, y/d = 0.71 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 488 | 1.75E-06 | 2.2E-06 | 9.86E-08 | 3.65E-08 | 4.98E-08 | $5.42 \mathrm{E}-08$ |
| 977 | 1.31E-06 | 2.22E-06 | 5.78E-08 | 4.15E-08 | $5.67 \mathrm{E}-08$ | 4.7E-08 |
| 1465 | 1.01E-06 | $1.68 \mathrm{E}-06$ | 3.56E-08 | 4.13E-08 | 4.95E-08 | 6.03E-08 |
| 1953 | 9.05E-07 | 2.03E-06 | 3.78E-08 | 3.83E-08 | 4.14E-08 | $4.64 \mathrm{E}-08$ |
| 2441 | 1.11E-06 | 1.5E-06 | $2.62 \mathrm{E}-08$ | 3.09E-08 | $2.85 \mathrm{E}-08$ | $5.44 \mathrm{E}-08$ |
| 2930 | 7.3E-07 | 1.18E-06 | 2.73E-08 | 3.75E-08 | $3.99 \mathrm{E}-08$ | 5E-08 |
| 3418 | 8.27E-07 | 9.47E-07 | 1.45E-08 | 2.85E-08 | $2.78 \mathrm{E}-08$ | 3.26E-08 |
| 3906 | 7.58E-07 | 6.76E-07 | 2.08E-08 | 3.18E-08 | 2.65E-08 | 2.81E-08 |
| 4395 | 5.22E-07 | 6.39E-07 | $1.48 \mathrm{E}-08$ | $2.39 \mathrm{E}-08$ | 1.9E-08 | $1.43 \mathrm{E}-08$ |
| 4883 | 4.49E-07 | 4.38E-07 | $1.06 \mathrm{E}-08$ | $2.03 \mathrm{E}-08$ | $1.31 \mathrm{E}-08$ | $1.25 \mathrm{E}-08$ |
| 5371 | 4.6E-07 | 4.1E-07 | 1.03E-08 | 1.75E-08 | $1.37 \mathrm{E}-08$ | $1.28 \mathrm{E}-08$ |
| 5859 | 3.14E-07 | 2.77E-07 | 5.85E-09 | 1.58E-08 | 9.22E-09 | 7.15E-09 |
| 6348 | 3.22E-07 | 2.44E-07 | $4.51 \mathrm{E}-09$ | 1.5E-08 | 8.61E-09 | $4.62 \mathrm{E}-09$ |
| 6836 | 2.46E-07 | 1.75E-07 | 3.72E-09 | 1.16E-08 | 5.51E-09 | 3.89E-09 |
| 7324 | 2.23E-07 | 1.22E-07 | 3.14E-09 | 1.08E-08 | 6.31E-09 | 2.86E-09 |
| 7813 | 1.8E-07 | $9.79 \mathrm{E}-08$ | 3.44E-09 | 1.08E-08 | $4.81 \mathrm{E}-09$ | $2.54 \mathrm{E}-09$ |
| 8301 | 1.65E-07 | $9.27 \mathrm{E}-08$ | 2.48E-09 | 8.24E-09 | $4.69 \mathrm{E}-09$ | 3.27E-09 |
| 8789 | 2.12E-07 | 8.12E-08 | 1.49E-09 | $7.78 \mathrm{E}-09$ | $4.25 \mathrm{E}-09$ | 2.14E-09 |
| 9277 | 1.55E-07 | 7.95E-08 | 1.87E-09 | 8.85E-09 | 3.76E-09 | $2.73 \mathrm{E}-09$ |
| 9766 | 1.52E-07 | $5.47 \mathrm{E}-08$ | 1.94E-09 | $6.91 \mathrm{E}-09$ | 4.53E-09 | 3.02E-09 |
| 10250 | $1.38 \mathrm{E}-07$ | 5.98E-08 | 1.68E-09 | 8.03E-09 | 4.82E-09 | 2.32E-09 |
| 10740 | $1.28 \mathrm{E}-07$ | 7.47E-08 | $8.38 \mathrm{E}-10$ | 8.09E-09 | $4.43 \mathrm{E}-09$ | 2.66E-09 |
| 11230 | 1.18E-07 | $5.36 \mathrm{E}-08$ | 2.1E-09 | 7.5E-09 | 3.43E-09 | 2.14E-09 |
| 11720 | 9.77E-08 | 5.17E-08 | 1.7E-09 | 6.19E-09 | 4.33E-09 | 2.56E-09 |
| 12210 | $1.31 \mathrm{E}-07$ | 5.1E-08 | 1.18E-09 | 6.02E-09 | 4.55E-09 | 1.93E-09 |
| 12700 | 9.98E-08 | 3.57E-08 | 1.17E-09 | $4.54 \mathrm{E}-09$ | 4.5E-09 | 2.15E-09 |
| 13180 | $9.89 \mathrm{E}-08$ | 3.86E-08 | 1.96E-09 | 5.01E-09 | 3.39E-09 | 1.88E-09 |
| 13670 | 1.2E-07 | 5.05E-08 | 1.62E-09 | 5.16E-09 | $4.61 \mathrm{E}-09$ | 2.08E-09 |
| 14160 | 8.58E-08 | 3.79E-08 | 1.11E-09 | 5.19E-09 | 4.34E-09 | 1.77E-09 |
| 14650 | 9.78E-08 | 3.11E-08 | 1.19E-09 | $4.61 \mathrm{E}-09$ | 3.68E-09 | 1.76E-09 |
| 15140 | 7.63E-08 | 4.53E-08 | 9.48E-10 | 4.27E-09 | 3.78E-09 | 1.96E-09 |
| 15630 | $7.57 \mathrm{E}-08$ | 3.96E-08 | 1.04E-09 | 4.41E-09 | 3.4E-09 | 1.8E-09 |
| 16110 | $6.92 \mathrm{E}-08$ | 3.8E-08 | 1.17E-09 | $4.24 \mathrm{E}-09$ | 3.32E-09 | 2.02E-09 |
| 16600 | 7.8E-08 | 3.64E-08 | 1.2E-09 | 5.46E-09 | 4.04E-09 | 1.98E-09 |
| 17090 | $8.55 \mathrm{E}-08$ | 4.85E-08 | $4.21 \mathrm{E}-10$ | 3.54E-09 | 3.18E-09 | $1.91 \mathrm{E}-09$ |
| 17580 | 8.28E-08 | $6.08 \mathrm{E}-08$ | 9.99E-10 | 4E-09 | 3.98E-09 | 2.22E-09 |
| 18070 | 1.41E-07 | $5.81 \mathrm{E}-08$ |  | 3.18E-09 | 3.63E-09 | $2.31 \mathrm{E}-09$ |


| 18550 | 1.75E-07 | 1.59E-08 |  | 3.65E-09 | 3.95E-09 | 1.66E-09 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 19040 | 6.79E-08 |  |  | 4.31E-09 | 3.8E-09 | 1.88E-09 |
| 19530 | 5.79E-08 | 1.39E-08 |  | 2.72E-09 | 3.61E-09 | 1.89E-09 |
| 20020 | 5.72E-08 | 1.7E-08 | 6.61E-10 | 3.27E-09 | 3.58E-09 | 1.89E-09 |
| 20510 | 5.22E-08 | 2.17E-08 | 3.27E-10 | 2.73E-09 | 3.19E-09 | 2.11E-09 |
| 21000 | 4.91E-08 | 1.81E-08 | $6.65 \mathrm{E}-10$ | 3.08E-09 | 3.13E-09 | $1.6 \mathrm{E}-09$ |
| 21480 | 5.45E-08 | $2.65 \mathrm{E}-08$ | $3.85 \mathrm{E}-10$ | 3.79E-09 | 3.1E-09 | 1.92E-09 |
| 21970 | 4.41E-08 | 2.33E-08 | $4.78 \mathrm{E}-10$ | 2.52E-09 | 2.92E-09 | 1.81E-09 |
| 22460 | 5.67E-08 | 2.26E-08 | 5.78E-10 | 2.72E-09 | 2.96E-09 | 1.91E-09 |
| 22950 | 4.79E-08 | 2.18E-08 | 7.91E-10 | 2.32E-09 | 2.75E-09 | 1.4E-09 |
| 23440 | $6.36 \mathrm{E}-08$ | 2.41E-08 | $6.78 \mathrm{E}-10$ | 2.75E-09 | 2.63E-09 | $1.76 \mathrm{E}-09$ |
| 23930 | $4.59 \mathrm{E}-08$ | 2.83E-08 | $5.33 \mathrm{E}-10$ | 2.25E-09 | 2.53E-09 | 1.48E-09 |
| 24410 | $5.46 \mathrm{E}-08$ | $2.56 \mathrm{E}-08$ | 6.79E-10 | 2.4E-09 | 2.11E-09 | 1.53E-09 |
| 24900 | 6.16E-08 | 2.5E-08 | $7.49 \mathrm{E}-10$ | 2.53E-09 | 2.23E-09 | 1.12E-09 |
| 25390 | 4.34E-08 | $2.29 \mathrm{E}-08$ | 7.17E-10 | 9.15E-09 | 9.83E-09 | 3.96E-09 |
| 25880 | $6.24 \mathrm{E}-08$ | 2.52E-08 | $3.97 \mathrm{E}-10$ | 4.15E-09 | 5.07E-09 | 2.33E-09 |
| 26370 | $5.09 \mathrm{E}-08$ | 2.49E-08 | 7.22E-10 | 4.46E-09 | 4.02E-09 | 1.72E-09 |
| 26860 | 4.43E-08 | 2.31E-08 | $6.91 \mathrm{E}-10$ | 3.41E-09 | 4.16E-09 | 1.86E-09 |
| 27340 | $4.52 \mathrm{E}-08$ | $2.38 \mathrm{E}-08$ | 8.62E-10 | 2.52E-09 | 3.41E-09 | $1.95 \mathrm{E}-09$ |
| 27830 | $4.86 \mathrm{E}-08$ | 2.31E-08 | 5.16E-10 | 2.4E-09 | 3.27E-09 | 2.05E-09 |
| 28320 | 4.37E-08 | 1.92E-08 | $4.65 \mathrm{E}-10$ | 3.09E-09 | 3.76E-09 | 2.12E-09 |
| 28810 | 3.91E-08 | 2.07E-08 | $6.25 \mathrm{E}-10$ | 2.28E-09 | 4.17E-09 | $1.87 \mathrm{E}-09$ |
| 29300 | $4.24 \mathrm{E}-08$ | 2.28E-08 | $5.41 \mathrm{E}-10$ | 2.45E-09 | 3.12E-09 | 1.96E-09 |
| 29790 | 4.53E-08 | 2.73E-08 | $5.73 \mathrm{E}-10$ | 3.13E-09 | 3.03E-09 | 1.71E-09 |
| 30270 | $4.49 \mathrm{E}-08$ | 2.06E-08 | $3.66 \mathrm{E}-10$ | 2.49E-09 | 3.36E-09 | $2.07 \mathrm{E}-09$ |
| 30760 | 5E-08 | 2.37E-08 | $6.03 \mathrm{E}-10$ | 2.23E-09 | 2.74E-09 | 1.62E-09 |
| 31250 | 4.46E-08 | 1.89E-08 | $5.69 \mathrm{E}-10$ | 2.2E-09 | 2.71E-09 | 1.8E-09 |
| 31740 | 3.95E-08 | $1.96 \mathrm{E}-08$ | 5.54E-10 | 2.95E-09 | 2.69E-09 | $1.68 \mathrm{E}-09$ |
| 32230 | 4.55E-08 | 1.65E-08 | 4.91E-10 | 2.44E-09 | 3.81E-09 | $1.89 \mathrm{E}-09$ |
| 32710 | 3.95E-08 | 2.49E-08 | 5.26E-10 | 2.41E-09 | 3.12E-09 | 2.13E-09 |
| 33200 | 3.35E-08 | 1.69E-08 | 5.62E-10 | 2.57E-09 | 2.99E-09 | 2.01E-09 |
| 33690 | 3.68E-08 | 2.41E-08 | $4.34 \mathrm{E}-10$ | 2.71E-09 | 3.08E-09 | 2.03E-09 |
| 34180 | 4.06E-08 | 2.36E-08 | $3.77 \mathrm{E}-10$ | 2.42E-09 | 2.86E-09 | 1.85E-09 |
| 34670 | 3.92E-08 | 2.61E-08 | $4.56 \mathrm{E}-10$ | 2.62E-09 | 3.35E-09 | 2.03E-09 |
| 35160 | 4E-08 | 2.34E-08 | $5.05 \mathrm{E}-10$ | 2.26E-09 | 3.04E-09 | 1.66E-09 |
| 35640 | 3.91E-08 | 2.16E-08 | 4.05E-10 | 2.42E-09 | 2.87E-09 | $2.15 \mathrm{E}-09$ |
| 36130 | 3.62E-08 | 2.38E-08 | 3.57E-10 | 1.99E-09 | 3.12E-09 | $1.69 \mathrm{E}-09$ |
| 36620 | 4.92E-08 | 2.02E-08 | 4.87E-10 | 2.85E-09 | 3.66E-09 | $2.25 \mathrm{E}-09$ |
| 37110 | 2.83E-08 | $2.69 \mathrm{E}-08$ | $5.94 \mathrm{E}-10$ | 1.87E-09 | 2.88E-09 | 1.74E-09 |
| 37600 | 3.73E-08 | 1.95E-08 | $3.48 \mathrm{E}-10$ | 2.23E-09 | 2.83E-09 | 2.32E-09 |
| 38090 | 3.32E-08 | $1.64 \mathrm{E}-08$ | $4.96 \mathrm{E}-10$ | 1.88E-09 | 3.42E-09 | 2.38E-09 |
| 38570 | 2.83E-08 | 2.22E-08 | 5.01E-10 | 2.55E-09 | 2.62E-09 | $1.48 \mathrm{E}-09$ |
| 39060 | 3.77E-08 | 1.97E-08 | 5.43E-10 | 2.56E-09 | 2.72E-09 | 1.92E-09 |
| 39550 | $3.21 \mathrm{E}-08$ | 2.32E-08 | 4.51E-10 | 2.09E-09 | 2.87E-09 | 2.08E-09 |
| 40040 | 3.21E-08 | 2.51E-08 | 5.46E-10 | $2.39 \mathrm{E}-09$ | $2.75 \mathrm{E}-09$ | 1.66E-09 |
| 40530 | 3.05E-08 | 2.05E-08 | 5.51E-10 | 1.93E-09 | 2.41E-09 | 1.92E-09 |
| 41020 | $4.55 \mathrm{E}-08$ | $2.35 \mathrm{E}-08$ | $4.33 \mathrm{E}-10$ | 2.26E-09 | 3.07E-09 | 2.4E-09 |
| 41500 | 4.06E-08 | 2.37E-08 | 5.02E-10 | 2.18E-09 | 3.43E-09 | 1.52E-09 |
| 41990 | 4.1E-08 | 2.13E-08 | 4.96E-10 | 1.96E-09 | 2.07E-09 | 2.16E-09 |
| 42480 | 3.7E-08 | 2.33E-08 | $6.89 \mathrm{E}-10$ | 2.15E-09 | 2.72E-09 | $1.98 \mathrm{E}-09$ |
| 42970 | 3.35E-08 | 2.07E-08 | $5.05 \mathrm{E}-10$ | 1.72E-09 | 2.2E-09 | 1.71E-09 |
| 43460 | 4E-08 | 2.13E-08 | $5.77 \mathrm{E}-10$ | 2.37E-09 | 2.37E-09 | 1.9E-09 |
| 43950 | 3.72E-08 | 2.22E-08 | $4.52 \mathrm{E}-10$ | 2.43E-09 | 2.61E-09 | 2.04E-09 |
| 44430 | $3.66 \mathrm{E}-08$ | 2.25E-08 | $7.44 \mathrm{E}-10$ | 2.11E-09 | 3.27E-09 | 1.91E-09 |
| 44920 | 4.51E-08 | $2.64 \mathrm{E}-08$ | $5.07 \mathrm{E}-10$ | 1.93E-09 | 2.44E-09 | $1.87 \mathrm{E}-09$ |
| 45410 | 3.29E-08 | 2.47E-08 | $6.93 \mathrm{E}-10$ | 2.15E-09 | 2.8E-09 | 2.02E-09 |
| 45900 | 4.04E-08 | 2.07E-08 | $6.73 \mathrm{E}-10$ | 1.83E-09 | 3.37E-09 | 1.89E-09 |
| 46390 | $4.25 \mathrm{E}-08$ | 2.62E-08 | 7.61E-10 | 2.21E-09 | 2.88E-09 | $2.56 \mathrm{E}-09$ |
| 46880 | $4.84 \mathrm{E}-08$ | 2.55E-08 | 1.07E-09 | 1.98E-09 | 3.24E-09 | 1.79E-09 |
| 47360 | 3.46E-08 | $3.17 \mathrm{E}-08$ | 1.17E-09 | 1.96E-09 | 2.72E-09 | $2.39 \mathrm{E}-09$ |


| 47850 | 5.45E-08 | 3.1E-08 | 1.39E-09 | 2.34E-09 | 2.33E-09 | 2.37E-09 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 48340 | 4.74E-08 | 3.51E-08 | 1.63E-09 | 1.99E-09 | 2.57E-09 | 2E-09 |
| 48830 | 5.97E-08 | 4.74E-08 | 1.23E-09 | 2.03E-09 | $2.48 \mathrm{E}-09$ | 1.82E-09 |
| 49320 | 6.66E-08 | 5.88E-08 | 3.28E-09 | $2.07 \mathrm{E}-09$ | 2.79E-09 | 2.34E-09 |
| 49800 | 8.04E-08 | 7.23E-08 | 3.75E-09 | 2.07E-09 | 2.99E-09 | $1.61 \mathrm{E}-09$ |
| 50290 | 1.08E-07 | 7.39E-08 | 7.27E-09 | 1.96E-09 | 3.18E-09 | 1.76E-09 |
| 50780 | 7.93E-08 | $5.84 \mathrm{E}-08$ | 1.42E-08 | 2.31E-09 | 2.67E-09 | 2E-09 |
| 51270 | $2.96 \mathrm{E}-08$ |  |  | $2.06 \mathrm{E}-09$ | 2.2E-09 | 2.07E-09 |
| 51760 | 8.08E-09 |  | 4.14E-09 | $1.76 \mathrm{E}-09$ | 2.65E-09 | $2.41 \mathrm{E}-09$ |
| 52250 | 1E-08 |  |  | 2.06E-09 | 2.32E-09 | 1.64E-09 |
| 52730 | 8.7E-09 | 4.67E-09 |  | $1.81 \mathrm{E}-09$ | 2.16E-09 | $1.82 \mathrm{E}-09$ |
| 53220 | 9.96E-09 | 6.26E-09 | $5.95 \mathrm{E}-10$ | 1.52E-09 | 2.29E-09 | 1.72E-09 |
| 53710 | 1.13E-08 | 4.61E-09 | 5.15E-10 | 1.95E-09 | 2.01E-09 | $1.59 \mathrm{E}-09$ |
| 54200 | 1.01E-08 | 6.15E-09 | 2.33E-10 | 1.45E-09 | 2.79E-09 | 1.54E-09 |
| 54690 | $1.27 \mathrm{E}-08$ | 6.63E-09 | $2.53 \mathrm{E}-10$ | 2.12E-09 | 2.72E-09 | 1.52E-09 |
| 55180 | 1.26E-08 | 6.34E-09 | $2.68 \mathrm{E}-10$ | 1.94E-09 | 2.4E-09 | 1.45E-09 |
| 55660 | 1.07E-08 | 7.04E-09 | 1.89E-10 | 1.77E-09 | 2.7E-09 | 2.21E-09 |
| 56150 | 8.7E-09 | 5.51E-09 | 2.2E-10 | 1.79E-09 | 2.45E-09 | $1.55 \mathrm{E}-09$ |
| 56640 | 1.27E-08 | 6.54E-09 | 2.57E-10 | 1.89E-09 | 1.81E-09 | 1.9E-09 |
| 57130 | 8.53E-09 | 5.13E-09 | $1.94 \mathrm{E}-10$ | 1.3E-09 | 1.86E-09 | 1.48E-09 |
| 57620 | 1.22E-08 | 6.52E-09 | $2.14 \mathrm{E}-10$ | 1.4E-09 | 2.33E-09 | $1.55 \mathrm{E}-09$ |
| 58110 | 1.22E-08 | 6.08E-09 | $2.36 \mathrm{E}-10$ | 1.33E-09 | 2.31E-09 | 1.81E-09 |
| 58590 | 1.03E-08 | $6.91 \mathrm{E}-09$ | 2.32E-10 | 1.3E-09 | 1.78E-09 | $1.29 \mathrm{E}-09$ |
| 59080 | 7.24E-09 | 5.63E-09 | 1.93E-10 | 1.34E-09 | 1.88E-09 | $1.27 \mathrm{E}-09$ |
| 59570 | 8.54E-09 | 6.7E-09 | 1.86E-10 | 1.48E-09 | 1.59E-09 | 1.16E-09 |
| 60060 | 9.68E-09 | 6.31E-09 | 1.51E-10 | 1.18E-09 | $1.53 \mathrm{E}-09$ | $8.08 \mathrm{E}-10$ |
| 60550 | 8.55E-09 | 5.75E-09 | $1.85 \mathrm{E}-10$ | 1.14E-09 | 1.03E-09 | $7.36 \mathrm{E}-10$ |
| 61040 | 9.86E-09 | 4.8E-09 | 1.74E-10 | 1.7E-09 | 1.35E-09 | 7.6E-10 |
| 61520 | 7.9E-09 | 7.4E-09 | $2.36 \mathrm{E}-10$ | 2.33E-09 | 2.5E-09 | 1.14E-09 |
| 62010 | 8.36E-09 | 7.04E-09 | 2.2E-10 | 2.3E-09 | 3.67E-09 | 2.41E-09 |
| 62500 | 8.26E-09 | $7.08 \mathrm{E}-09$ | 2.06E-10 | 2E-09 | 2.94E-09 | 2.17E-09 |
| 62990 | 8.58E-09 | 6.81E-09 | $1.77 \mathrm{E}-10$ | 1.88E-09 | 2.56E-09 | $2.39 \mathrm{E}-09$ |
| 63480 | 8.36E-09 | $5.74 \mathrm{E}-09$ | 2.31E-10 | 1.79E-09 | 2.26E-09 | $1.79 \mathrm{E}-09$ |
| 63960 | 8E-09 | 7.95E-09 | 1.41E-10 | 1.41E-09 | 2.57E-09 | 2.07E-09 |
| 64450 | 7.35E-09 | 5.94E-09 | $2.25 \mathrm{E}-10$ | 1.42E-09 | 2.35E-09 | 1.94E-09 |
| 64940 | 8.61E-09 | 5.51E-09 | $1.36 \mathrm{E}-10$ | 1.38E-09 | 2.27E-09 | $1.99 \mathrm{E}-09$ |
| 65430 | 7.86E-09 | 7.24E-09 | $1.23 \mathrm{E}-10$ | 1.28E-09 | 1.61E-09 | $1.66 \mathrm{E}-09$ |
| 65920 | 8.02E-09 | 7.68E-09 | 1.8E-10 | 1.3E-09 | 1.68E-09 | 1.7E-09 |
| 66410 | 7.52E-09 | 5.35E-09 | 1.41E-10 | 1.18E-09 | 1.88E-09 | $1.68 \mathrm{E}-09$ |
| 66890 | 5.95E-09 | 5.17E-09 | 2.47E-10 | 9.81E-10 | 1.98E-09 | $2.18 \mathrm{E}-09$ |
| 67380 | 6.93E-09 | 5.43E-09 | $1.58 \mathrm{E}-10$ | $9.42 \mathrm{E}-10$ | 1.68E-09 | 1.63E-09 |
| 67870 | 6.16E-09 | 5.51E-09 | $1.75 \mathrm{E}-10$ | 8.56E-10 | 1.27E-09 | 1.64E-09 |
| 68360 | $6.29 \mathrm{E}-09$ | $5.28 \mathrm{E}-09$ | 1.8E-10 | $8.47 \mathrm{E}-10$ | 1.63E-09 | $1.45 \mathrm{E}-09$ |
| 68850 | 6.46E-09 | 7.08E-09 | $1.61 \mathrm{E}-10$ | 8.81E-10 | 1.76E-09 | 1.59E-09 |
| 69340 | 5.41E-09 | 4.76E-09 | 1.69E-10 | 7.83E-10 | 1.49E-09 | 1.42E-09 |
| 69820 | 5.55E-09 | 5.73E-09 | 2.1E-10 | $7.95 \mathrm{E}-10$ | 1.53E-09 | 1.23E-09 |
| 70310 | 6.53E-09 | 4.98E-09 | $1.73 \mathrm{E}-10$ | 8.33E-10 | 1.72E-09 | 1.39E-09 |
| 70800 | 5.03E-09 | 6.07E-09 | $1.78 \mathrm{E}-10$ | $6.41 \mathrm{E}-10$ | 1.64E-09 | $1.47 \mathrm{E}-09$ |
| 71290 | 6E-09 | 4.85E-09 | $1.82 \mathrm{E}-10$ | $6.27 \mathrm{E}-10$ | 1.18E-09 | 1.15E-09 |
| 71780 | 6.6E-09 | 7E-09 | 1.51E-10 | $5.97 \mathrm{E}-10$ | 1.21E-09 |  |
| 72270 | 5.93E-09 | 5.17E-09 | 1.97E-10 | $5.95 \mathrm{E}-10$ | $1.08 \mathrm{E}-09$ |  |
| 72750 | 7.42E-09 | 4.91E-09 | 1.74E-10 | $5.15 \mathrm{E}-10$ | $9.98 \mathrm{E}-10$ |  |
| 73240 | 6.62E-09 | 5.1E-09 | $2.14 \mathrm{E}-10$ | $3.73 \mathrm{E}-10$ | $1.08 \mathrm{E}-09$ |  |

Table C-34: Mach 5.0 ZPG pressure data

| $y$ (in) | $\mathrm{P}_{\mathrm{t} 1}$ (psia) | $\mathrm{P}_{\mathrm{t} 2}$ (psia) | $\mathrm{P}_{\mathrm{t} 2} / \mathrm{P}_{11}$ | ps / $\mathrm{P}_{12}$ | Mach | $\mathrm{M} / \mathrm{M}_{\mathrm{e}}$ | $\mathrm{u} / \mathrm{U}_{\mathrm{e}}$ | pw/p | plpe |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.0717 | 440 | 15.6941 | 0.0356 | 0.0594 | 3.5782 | 0.7016 | 0.9260 | 0.2873 | 0.5612 |
| 0.0807 | 439 | 16.4591 | 0.0375 | 0.0567 | 3.6515 | 0.7160 | 0.9312 | 0.2798 | 0.5764 |
| 0.0897 | 440 | 17.0137 | 0.0386 | 0.0548 | 3.7121 | 0.7279 | 0.9353 | 0.2737 | 0.5891 |
| 0.0985 | 440 | 17.0168 | 0.0387 | 0.0548 | 3.7157 | 0.7286 | 0.9356 | 0.2734 | 0.5898 |
| 0.1074 | 439 | 17.3033 | 0.0394 | 0.0539 | 3.7458 | 0.7345 | 0.9375 | 0.2704 | 0.5962 |
| 0.1165 | 440 | 17.7787 | 0.0404 | 0.0525 | 3.7958 | 0.7443 | 0.9408 | 0.2657 | 0.6069 |
| 0.1253 | 440 | 17.8884 | 0.0406 | 0.0521 | 3.8109 | 0.7472 | 0.9417 | 0.2643 | 0.6101 |
| 0.1345 | 440 | 18.4004 | 0.0418 | 0.0507 | 3.8618 | 0.7572 | 0.9449 | 0.2596 | 0.6211 |
| 0.1433 | 440 | 19.1654 | 0.0435 | 0.0487 | 3.9377 | 0.7721 | 0.9495 | 0.2529 | 0.6377 |
| 0.1524 | 441 | 19.5067 | 0.0442 | 0.0478 | 3.9783 | 0.7801 | 0.9518 | 0.2494 | 0.6466 |
| 0.1614 | 441 | 19.7657 | 0.0448 | 0.0472 | 4.0070 | 0.7857 | 0.9535 | 0.2469 | 0.6530 |
| 0.1705 | 441 | 20.6800 | 0.0469 | 0.0451 | 4.0880 | 0.8016 | 0.9580 | 0.2402 | 0.6711 |
| 0.1795 | 440 | 20.7227 | 0.0471 | 0.0450 | 4.1063 | 0.8051 | 0.9589 | 0.2388 | 0.6752 |
| 0.1885 | 440 | 21.1250 | 0.0481 | 0.0441 | 4.1422 | 0.8122 | 0.9608 | 0.2359 | 0.6834 |
| 0.1976 | 440 | 21.9204 | 0.0498 | 0.0425 | 4.2109 | 0.8257 | 0.9644 | 0.2306 | 0.6991 |
| 0.2066 | 440 | 22.0667 | 0.0501 | 0.0423 | 4.2369 | 0.8308 | 0.9657 | 0.2287 | 0.7051 |
| 0.2157 | 440 | 22.3654 | 0.0508 | 0.0417 | 4.2652 | 0.8363 | 0.9671 | 0.2266 | 0.7116 |
| 0.2248 | 440 | 23.0755 | 0.0524 | 0.0404 | 4.3231 | 0.8477 | 0.9698 | 0.2224 | 0.7251 |
| 0.2337 | 439 | 23.5174 | 0.0535 | 0.0397 | 4.3685 | 0.8566 | 0.9720 | 0.2192 | 0.7357 |
| 0.2429 | 441 | 24.1452 | 0.0547 | 0.0386 | 4.4231 | 0.8673 | 0.9745 | 0.2154 | 0.7486 |
| 0.2518 | 441 | 24.9436 | 0.0565 | 0.0374 | 4.4902 | 0.8804 | 0.9774 | 0.2109 | 0.7645 |
| 0.2610 | 441 | 25.2575 | 0.0573 | 0.0369 | 4.5299 | 0.8882 | 0.9791 | 0.2083 | 0.7740 |
| 0.2701 | 441 | 25.5410 | 0.0579 | 0.0365 | 4.5596 | 0.8940 | 0.9804 | 0.2064 | 0.7812 |
| 0.2792 | 440 | 26.2968 | 0.0597 | 0.0355 | 4.6149 | 0.9049 | 0.9826 | 0.2029 | 0.7946 |
| 0.2882 | 439 | 26.3303 | 0.0599 | 0.0354 | 4.6346 | 0.9087 | 0.9834 | 0.2017 | 0.7993 |
| 0.2970 | 440 | 26.7600 | 0.0608 | 0.0349 | 4.6666 | 0.9150 | 0.9847 | 0.1998 | 0.8071 |
| 0.3060 | 440 | 27.5036 | 0.0625 | 0.0339 | 4.7201 | 0.9255 | 0.9868 | 0.1966 | 0.8202 |
| 0.3153 | 439 | 27.9730 | 0.0637 | 0.0333 | 4.7647 | 0.9343 | 0.9885 | 0.1940 | 0.8312 |
| 0.3243 | 441 | 28.0918 | 0.0637 | 0.0332 | 4.7871 | 0.9386 | 0.9893 | 0.1927 | 0.8368 |
| 0.3334 | 440 | 28.2869 | 0.0642 | 0.0330 | 4.8058 | 0.9423 | 0.9900 | 0.1916 | 0.8414 |
| 0.3426 | 442 | 28.8141 | 0.0652 | 0.0324 | 4.8411 | 0.9492 | 0.9913 | 0.1897 | 0.8502 |
| 0.3516 | 440 | 29.1859 | 0.0663 | 0.0320 | 4.8738 | 0.9557 | 0.9924 | 0.1878 | 0.8584 |
| 0.3607 | 442 | 29.1463 | 0.0660 | 0.0320 | 4.8839 | 0.9576 | 0.9928 | 0.1873 | 0.8609 |
| 0.3698 | 441 | 29.5455 | 0.0669 | 0.0316 | 4.9087 | 0.9625 | 0.9937 | 0.1860 | 0.8671 |
| 0.3789 | 441 | 29.7772 | 0.0675 | 0.0313 | 4.9301 | 0.9667 | 0.9944 | 0.1848 | 0.8725 |
| 0.3879 | 441 | 30.0819 | 0.0682 | 0.0310 | 4.9539 | 0.9714 | 0.9952 | 0.1835 | 0.8785 |
| 0.3969 | 442 | 30.2648 | 0.0685 | 0.0308 | 4.9724 | 0.9750 | 0.9958 | 0.1826 | 0.8831 |
| 0.4061 | 441 | 30.2678 | 0.0687 | 0.0308 | 4.9798 | 0.9764 | 0.9961 | 0.1822 | 0.8850 |
| 0.4153 | 440 | 30.3654 | 0.0691 | 0.0307 | 4.9876 | 0.9780 | 0.9963 | 0.1818 | 0.8870 |
| 0.4242 | 440 | 30.3654 | 0.0690 | 0.0307 | 4.9907 | 0.9786 | 0.9964 | 0.1816 | 0.8878 |
| 0.4333 | 440 | 30.4354 | 0.0691 | 0.0306 | 4.9955 | 0.9795 | 0.9966 | 0.1814 | 0.8890 |
| 0.4424 | 442 | 30.7951 | 0.0697 | 0.0303 | 5.0153 | 0.9834 | 0.9973 | 0.1804 | 0.8940 |
| 0.4515 | 442 | 30.7067 | 0.0695 | 0.0304 | 5.0188 | 0.9841 | 0.9974 | 0.1802 | 0.8949 |
| 0.4606 | 442 | 30.7372 | 0.0696 | 0.0303 | 5.0217 | 0.9846 | 0.9975 | 0.1800 | 0.8956 |
| 0.4696 | 440 | 30.7067 | 0.0697 | 0.0304 | 5.0213 | 0.9846 | 0.9975 | 0.1801 | 0.8955 |
| 0.4787 | 442 | 30.9048 | 0.0700 | 0.0302 | 5.0374 | 0.9877 | 0.9980 | 0.1792 | 0.8996 |
| 0.4877 | 441 | 31.0876 | 0.0705 | 0.0300 | 5.0525 | 0.9907 | 0.9985 | 0.1785 | 0.9034 |
| 0.4966 | 442 | 31.0998 | 0.0704 | 0.0300 | 5.0536 | 0.9909 | 0.9985 | 0.1784 | 0.9037 |
| 0.5057 | 442 | 31.0236 | 0.0702 | 0.0301 | 5.0474 | 0.9897 | 0.9983 | 0.1787 | 0.9021 |
| 0.5148 | 442 | 31.1791 | 0.0705 | 0.0299 | 5.0601 | 0.9922 | 0.9987 | 0.1781 | 0.9054 |
| 0.5237 | 441 | 31.1943 | 0.0707 | 0.0299 | 5.0614 | 0.9924 | 0.9988 | 0.1780 | 0.9057 |
| 0.5328 | 442 | 31.0815 | 0.0704 | 0.0300 | 5.0521 | 0.9906 | 0.9985 | 0.1785 | 0.9033 |
| 0.5419 | 442 | 31.0511 | 0.0703 | 0.0300 | 5.0496 | 0.9901 | 0.9984 | 0.1786 | 0.9027 |
| 0.5509 | 442 | 31.2918 | 0.0708 | 0.0298 | 5.0693 | 0.9940 | 0.9990 | 0.1776 | 0.9077 |
| 0.5599 | 442 | 31.4320 | 0.0711 | 0.0297 | 5.0808 | 0.9962 | 0.9994 | 0.1771 | 0.9107 |
| 0.5690 | 442 | 31.1974 | 0.0706 | 0.0299 | 5.0617 | 0.9925 | 0.9988 | 0.1780 | 0.9058 |
| 0.5782 | 441 | 31.4930 | 0.0714 | 0.0296 | 5.0859 | 0.9972 | 0.9996 | 0.1768 | 0.9119 |


| 0.5872 | 442 | 31.4991 | 0.0713 | 0.0296 | 5.0864 | 0.9973 | 0.9996 | 0.1768 | 0.9121 |
| ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: |
| 0.5961 | 443 | 31.4808 | 0.0711 | 0.0296 | 5.0849 | 0.9970 | 0.9995 | 0.1769 | 0.9117 |
| 0.6052 | 442 | 31.7246 | 0.0718 | 0.0294 | 5.1047 | 1.0009 | 1.0001 | 0.1759 | 0.9168 |
| 0.6143 | 441 | 31.6606 | 0.0717 | 0.0295 | 5.0996 | 0.9999 | 1.0000 | 0.1761 | 0.9155 |
| 0.6233 | 442 | 31.8831 | 0.0722 | 0.0293 | 5.1177 | 1.0035 | 1.0006 | 0.1752 | 0.9201 |
| 0.6323 | 442 | 31.9654 | 0.0724 | 0.0292 | 5.1244 | 1.0048 | 1.0008 | 0.1749 | 0.9218 |
| 0.6413 | 441 | 31.9745 | 0.0724 | 0.0292 | 5.1252 | 1.0049 | 1.0008 | 0.1749 | 0.9220 |
| 0.6503 | 440 | 31.9836 | 0.0727 | 0.0292 | 5.1259 | 1.0051 | 1.0008 | 0.1748 | 0.9222 |
| 0.6594 | 442 | 32.1086 | 0.0726 | 0.0290 | 5.1360 | 1.0071 | 1.0011 | 0.1744 | 0.9248 |
| 0.6684 | 441 | 32.3463 | 0.0733 | 0.0288 | 5.1550 | 1.0108 | 1.0017 | 0.1734 | 0.9297 |
| 0.6776 | 441 | 32.3006 | 0.0733 | 0.0289 | 5.1517 | 1.0101 | 1.0016 | 0.1736 | 0.9288 |
| 0.6866 | 443 | 32.4591 | 0.0733 | 0.0287 | 5.1644 | 1.0126 | 1.0020 | 0.1730 | 0.9321 |
| 0.6955 | 442 | 32.5993 | 0.0737 | 0.0286 | 5.1756 | 1.0148 | 1.0023 | 0.1725 | 0.9350 |
| 0.7045 | 443 | 32.5871 | 0.0736 | 0.0286 | 5.1747 | 1.0147 | 1.0023 | 0.1725 | 0.9347 |
| 0.7135 | 442 | 32.7151 | 0.0741 | 0.0285 | 5.1850 | 1.0167 | 1.0026 | 0.1720 | 0.9374 |
| 0.7225 | 442 | 32.9132 | 0.0745 | 0.0283 | 5.2008 | 1.0198 | 1.0031 | 0.1713 | 0.9415 |
| 0.7318 | 442 | 33.0534 | 0.0748 | 0.0282 | 5.2120 | 1.0220 | 1.0034 | 0.1707 | 0.9444 |
| 0.7407 | 442 | 33.1570 | 0.0751 | 0.0281 | 5.2203 | 1.0236 | 1.0037 | 0.1704 | 0.9465 |
| 0.7497 | 442 | 33.3246 | 0.0754 | 0.0280 | 5.2335 | 1.0262 | 1.0041 | 0.1697 | 0.9499 |
| 0.7590 | 442 | 33.6050 | 0.0761 | 0.0278 | 5.2555 | 1.0305 | 1.0047 | 0.1687 | 0.9556 |
| 0.7679 | 442 | 33.6172 | 0.0761 | 0.0277 | 5.2569 | 1.0308 | 1.0048 | 0.1687 | 0.9560 |

Table C-35: Mach 5.0 FPG pressure data

| $y$ (in) | $\mathrm{P}_{\mathrm{t1}}$ (psia) | $\mathrm{P}_{12}$ (psia) | $P_{12} / P_{11}$ | ps / $\mathrm{P}_{12}$ | Mach | P (Pa) | T (K) | $\rho$ | $\rho / \rho e$ | $\mathrm{M} / \mathrm{M}_{\mathrm{e}}$ | $\mathrm{u} / \mathrm{U}_{\mathrm{e}}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.063 | 506 | 15.56647 | 0.031 | 0.069 | 3.304 | 5278 | 113.1110 | 0.1626 | 0.5483 | 0.6210 | 0.8983 |
| 0.071 | 505 | 15.70056 | 0.031 | 0.069 | 3.303 | 5260 | 113.1185 | 0.1620 | 0.5464 | 0.6209 | 0.8983 |
| 0.080 | 505 | 15.72799 | 0.031 | 0.069 | 3.305 | 5243 | 113.0546 | 0.1616 | 0.5449 | 0.6212 | 0.8984 |
| 0.089 | 505 | 16.69409 | 0.033 | 0.065 | 3.405 | 5225 | 108.4737 | 0.1678 | 0.5659 | 0.6400 | 0.9067 |
| 0.097 | 506 | 16.96533 | 0.034 | 0.064 | 3.436 | 5206 | 107.0935 | 0.1694 | 0.5712 | 0.6459 | 0.9092 |
| 0.106 | 504 | 17.22742 | 0.034 | 0.063 | 3.463 | 5189 | 105.9113 | 0.1707 | 0.5756 | 0.6510 | 0.9113 |
| 0.115 | 505 | 17.47733 | 0.035 | 0.062 | 3.489 | 5171 | 104.8120 | 0.1719 | 0.5797 | 0.6558 | 0.9133 |
| 0.123 | 506 | 17.9558 | 0.036 | 0.060 | 3.535 | 5153 | 102.8689 | 0.1745 | 0.5886 | 0.6645 | 0.9168 |
| 0.132 | 507 | 18.7299 | 0.037 | 0.058 | 3.608 | 5136 | 99.9156 | 0.1791 | 0.6039 | 0.6781 | 0.9220 |
| 0.141 | 505 | 18.58666 | 0.037 | 0.058 | 3.604 | 5118 | 100.0602 | 0.1782 | 0.6009 | 0.6775 | 0.9217 |
| 0.149 | 506 | 19.16875 | 0.038 | 0.056 | 3.660 | 5100 | 97.8518 | 0.1816 | 0.6124 | 0.6880 | 0.9257 |
| 0.158 | 505 | 19.83009 | 0.039 | 0.055 | 3.722 | 5082 | 95.4583 | 0.1855 | 0.6256 | 0.6997 | 0.9299 |
| 0.167 | 506 | 20.03733 | 0.040 | 0.054 | 3.744 | 5064 | 94.6432 | 0.1864 | 0.6287 | 0.7038 | 0.9313 |
| 0.175 | 505 | 20.34209 | 0.040 | 0.053 | 3.773 | 5046 | 93.5926 | 0.1879 | 0.6335 | 0.7091 | 0.9331 |
| 0.184 | 506 | 21.02171 | 0.042 | 0.051 | 3.833 | 5029 | 91.4118 | 0.1917 | 0.6464 | 0.7205 | 0.9370 |
| 0.193 | 506 | 21.55809 | 0.043 | 0.050 | 3.883 | 5011 | 89.6608 | 0.1947 | 0.6566 | 0.7298 | 0.9400 |
| 0.201 | 506 | 22.12799 | 0.044 | 0.049 | 3.934 | 4993 | 87.9140 | 0.1979 | 0.6673 | 0.7394 | 0.9430 |
| 0.210 | 506 | 22.20418 | 0.044 | 0.049 | 3.946 | 4975 | 87.5005 | 0.1981 | 0.6681 | 0.7417 | 0.9438 |
| 0.219 | 505 | 22.55466 | 0.045 | 0.048 | 3.975 | 4957 | 86.5342 | 0.1996 | 0.6731 | 0.7472 | 0.9454 |
| 0.228 | 506 | 23.39275 | 0.046 | 0.046 | 4.042 | 4939 | 84.3584 | 0.2040 | 0.6879 | 0.7598 | 0.9492 |
| 0.236 | 505 | 23.80418 | 0.047 | 0.045 | 4.083 | 4921 | 83.0748 | 0.2064 | 0.6961 | 0.7674 | 0.9514 |
| 0.245 | 506 | 23.92304 | 0.047 | 0.045 | 4.098 | 4904 | 82.5927 | 0.2069 | 0.6977 | 0.7703 | 0.9522 |
| 0.254 | 507 | 24.64228 | 0.049 | 0.044 | 4.151 | 4886 | 80.9621 | 0.2103 | 0.7091 | 0.7803 | 0.9550 |
| 0.262 | 507 | 25.11161 | 0.050 | 0.043 | 4.193 | 4868 | 79.7012 | 0.2128 | 0.7177 | 0.7882 | 0.9572 |
| 0.271 | 508 | 25.84304 | 0.051 | 0.042 | 4.250 | 4850 | 78.0450 | 0.2165 | 0.7302 | 0.7989 | 0.9600 |
| 0.280 | 506 | 26.58666 | 0.053 | 0.041 | 4.310 | 4833 | 76.3538 | 0.2205 | 0.7437 | 0.8101 | 0.9629 |
| 0.288 | 507 | 27.00114 | 0.053 | 0.040 | 4.350 | 4815 | 75.2542 | 0.2229 | 0.7517 | 0.8176 | 0.9647 |
| 0.297 | 507 | 27.20837 | 0.054 | 0.040 | 4.372 | 4797 | 74.6426 | 0.2239 | 0.7551 | 0.8218 | 0.9658 |
| 0.306 | 506 | 27.93066 | 0.055 | 0.039 | 4.421 | 4779 | 73.3382 | 0.2271 | 0.7656 | 0.8310 | 0.9680 |
| 0.314 | 504 | 28.24456 | 0.056 | 0.038 | 4.452 | 4761 | 72.5187 | 0.2288 | 0.7714 | 0.8369 | 0.9693 |
| 0.323 | 506 | 29.2838 | 0.058 | 0.037 | 4.519 | 4743 | 70.8038 | 0.2334 | 0.7871 | 0.8495 | 0.9722 |
| 0.332 | 506 | 29.75314 | 0.059 | 0.036 | 4.564 | 4725 | 69.6807 | 0.2363 | 0.7968 | 0.8579 | 0.9741 |
| 0.340 | 507 | 30.24076 | 0.060 | 0.036 | 4.604 | 4708 | 68.7135 | 0.2387 | 0.8050 | 0.8654 | 0.9757 |


| 0.349 | 508 | 30.85942 | 0.061 | 0.035 | 4.648 | 4690 | 67.6477 | 0.2416 | 0.8145 | 0.8738 | 0.9775 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.358 | 506 | 31.56952 | 0.062 | 0.034 | 4.699 | 4672 | 66.4726 | 0.2449 | 0.8258 | 0.8832 | 0.9795 |
| 0.367 | 507 | 31.77066 | 0.063 | 0.034 | 4.726 | 4654 | 65.8599 | 0.2462 | 0.8303 | 0.8883 | 0.9805 |
| 0.375 | 507 | 32.66361 | 0.064 | 0.033 | 4.778 | 4636 | 64.6783 | 0.2497 | 0.8422 | 0.8981 | 0.9825 |
| 0.384 | 508 | 33.20609 | 0.065 | 0.033 | 4.822 | 4618 | 63.7106 | 0.2526 | 0.8516 | 0.9064 | 0.9841 |
| 0.393 | 507 | 33.73333 | 0.067 | 0.032 | 4.862 | 4600 | 62.8444 | 0.2551 | 0.8601 | 0.9140 | 0.9855 |
| 0.402 | 507 | 34.05333 | 0.067 | 0.032 | 4.891 | 4600 | 62.2387 | 0.2575 | 0.8684 | 0.9193 | 0.9865 |
| 0.410 | 507 | 34.52266 | 0.068 | 0.031 | 4.922 | 4600 | 61.5899 | 0.2602 | 0.8776 | 0.9252 | 0.9876 |
| 0.419 | 506 | 34.74513 | 0.069 | 0.031 | 4.943 | 4600 | 61.1594 | 0.2621 | 0.8838 | 0.9291 | 0.9883 |
| 0.428 | 508 | 35.30894 | 0.070 | 0.031 | 4.975 | 4600 | 60.5076 | 0.2649 | 0.8933 | 0.9351 | 0.9894 |
| 0.436 | 505 | 35.97332 | 0.071 | 0.030 | 5.015 | 4600 | 59.7030 | 0.2685 | 0.9053 | 0.9427 | 0.9907 |
| 0.445 | 507 | 36.48533 | 0.072 | 0.030 | 5.052 | 4600 | 58.9819 | 0.2718 | 0.9164 | 0.9495 | 0.9919 |
| 0.454 | 508 | 36.92418 | 0.073 | 0.029 | 5.084 | 4600 | 58.3593 | 0.2747 | 0.9262 | 0.9556 | 0.9929 |
| 0.462 | 508 | 37.31733 | 0.074 | 0.029 | 5.112 | 4600 | 57.8169 | 0.2772 | 0.9348 | 0.9609 | 0.9938 |
| 0.471 | 507 | 37.70438 | 0.074 | 0.029 | 5.139 | 4600 | 57.3163 | 0.2797 | 0.9430 | 0.9659 | 0.9946 |
| 0.480 | 506 | 38.12495 | 0.075 | 0.029 | 5.166 | 4600 | 56.8122 | 0.2821 | 0.9514 | 0.9710 | 0.9955 |
| 0.488 | 508 | 38.37485 | 0.076 | 0.028 | 5.186 | 4600 | 56.4342 | 0.2840 | 0.9578 | 0.9748 | 0.9961 |
| 0.497 | 508 | 38.84418 | 0.076 | 0.028 | 5.212 | 4600 | 55.9565 | 0.2864 | 0.9659 | 0.9798 | 0.9969 |
| 0.506 | 509 | 39.11542 | 0.077 | 0.028 | 5.233 | 4600 | 55.5784 | 0.2884 | 0.9725 | 0.9837 | 0.9975 |
| 0.514 | 507 | 39.19771 | 0.077 | 0.028 | 5.245 | 4600 | 55.3727 | 0.2895 | 0.9761 | 0.9858 | 0.9978 |
| 0.523 | 509 | 39.73409 | 0.078 | 0.027 | 5.269 | 4600 | 54.9367 | 0.2918 | 0.9839 | 0.9905 | 0.9986 |
| 0.532 | 508 | 39.89561 | 0.079 | 0.027 | 5.285 | 4600 | 54.6532 | 0.2933 | 0.9890 | 0.9935 | 0.9990 |
| 0.540 | 508 | 40.15162 | 0.079 | 0.027 | 5.301 | 4600 | 54.3753 | 0.2948 | 0.9940 | 0.9965 | 0.9995 |
| 0.549 | 508 | 40.59961 | 0.080 | 0.027 | 5.324 | 4600 | 53.9806 | 0.2969 | 1.0013 | 1.0008 | 1.0001 |
| 0.558 | 508 | 40.42285 | 0.080 | 0.027 | 5.327 | 4600 | 53.9298 | 0.2972 | 1.0022 | 1.0013 | 1.0002 |
| 0.567 | 509 | 40.13332 | 0.079 | 0.027 | 5.317 | 4600 | 54.1056 | 0.2962 | 0.9990 | 0.9994 | 0.9999 |
| 0.575 | 508 | 39.34399 | 0.077 | 0.028 | 5.253 | 4600 | 55.2259 | 0.2902 | 0.9787 | 0.9874 | 0.9981 |
| 0.584 | 509 | 38.4419 | 0.075 | 0.028 | 5.189 | 4600 | 56.3746 | 0.2843 | 0.9588 | 0.9754 | 0.9962 |
| 0.593 | 509 | 37.60685 | 0.074 | 0.029 | 5.130 | 4600 | 57.4809 | 0.2789 | 0.9403 | 0.9642 | 0.9944 |
| 0.601 | 510 | 36.96685 | 0.073 | 0.030 | 5.084 | 4600 | 58.3488 | 0.2747 | 0.9263 | 0.9557 | 0.9930 |
| 0.610 | 510 | 36.49752 | 0.072 | 0.030 | 5.051 | 4600 | 59.0012 | 0.2717 | 0.9161 | 0.9493 | 0.9919 |
| 0.619 | 510 | 36.18666 | 0.071 | 0.030 | 5.028 | 4600 | 59.4423 | 0.2697 | 0.9093 | 0.9451 | 0.9911 |
| 0.627 | 510 | 36.02818 | 0.071 | 0.030 | 5.017 | 4600 | 59.6713 | 0.2686 | 0.9058 | 0.9430 | 0.9908 |
| 0.636 | 509 | 35.83619 | 0.070 | 0.031 | 5.001 | 4600 | 59.9738 | 0.2673 | 0.9012 | 0.9401 | 0.9903 |
| 0.645 | 509 | 35.72038 | 0.070 | 0.031 | 4.992 | 4600 | 60.1531 | 0.2665 | 0.8985 | 0.9384 | 0.9900 |
| 0.653 | 509 | 35.63809 | 0.070 | 0.031 | 4.985 | 4600 | 60.2916 | 0.2659 | 0.8965 | 0.9371 | 0.9897 |
| 0.662 | 510 | 35.50704 | 0.070 | 0.031 | 4.975 | 4600 | 60.5120 | 0.2649 | 0.8932 | 0.9351 | 0.9894 |
| 0.671 | 511 | 35.5558 | 0.070 | 0.031 | 4.978 | 4600 | 60.4402 | 0.2652 | 0.8943 | 0.9357 | 0.9895 |
| 0.680 | 510 | 35.53447 | 0.070 | 0.031 | 4.975 | 4600 | 60.5025 | 0.2649 | 0.8934 | 0.9352 | 0.9894 |
| 0.688 | 508 | 35.46742 | 0.070 | 0.031 | 4.969 | 4600 | 60.6241 | 0.2644 | 0.8916 | 0.9340 | 0.9892 |
| 0.697 | 511 | 35.44608 | 0.069 | 0.031 | 4.966 | 4600 | 60.6795 | 0.2642 | 0.8907 | 0.9335 | 0.9891 |
| 0.706 | 511 | 35.46742 | 0.069 | 0.031 | 4.967 | 4600 | 60.6734 | 0.2642 | 0.8908 | 0.9336 | 0.9891 |
| 0.714 | 510 | 35.36075 | 0.069 | 0.031 | 4.958 | 4600 | 60.8538 | 0.2634 | 0.8882 | 0.9319 | 0.9888 |
| 0.723 | 509 | 35.31199 | 0.069 | 0.031 | 4.953 | 4600 | 60.9496 | 0.2630 | 0.8868 | 0.9310 | 0.9887 |
| 0.732 | 510 | 35.34856 | 0.069 | 0.031 | 4.954 | 4600 | 60.9208 | 0.2631 | 0.8872 | 0.9313 | 0.9887 |
| 0.740 | 510 | 35.40951 | 0.069 | 0.031 | 4.958 | 4600 | 60.8573 | 0.2634 | 0.8881 | 0.9319 | 0.9888 |
| 0.749 | 512 | 35.41256 | 0.069 | 0.031 | 4.957 | 4600 | 60.8779 | 0.2633 | 0.8878 | 0.9317 | 0.9888 |
| 0.758 | 511 | 35.32113 | 0.069 | 0.031 | 4.949 | 4600 | 61.0368 | 0.2626 | 0.8855 | 0.9302 | 0.9885 |

Table C-36: Mach 5.0 CPG $\left(\mathrm{X}_{\mathrm{ts}}=5.1 \mathrm{~cm}\right)$ pressure data

| $y$ (in) | $\mathrm{P}_{\text {t1 }}$ (psia) | $\mathrm{P}_{\mathrm{t} 2}$ (psia) | $\mathrm{P}_{12} / \mathrm{P}_{11}$ | ps / $\mathrm{P}_{12}$ | Mach | P (Pa) | T (K) | p | $p / \mathrm{pe}$ | $\mathrm{M} / \mathrm{M}_{\mathrm{e}}$ | U/ $\mathrm{U}_{\mathrm{e}}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.0717 | 440 | 15.6941 | 0.0356 | 0.0891 | 2.90673 | 12695 | 133.8382 | 0.3305 | 0.5379 | 0.6760 | 0.8934 |
| 0.0807 | 439 | 16.4591 | 0.0375 | 0.0860 | 2.92825 | 12740 | 132.6001 | 0.3348 | 0.5448 | 0.6810 | 0.8958 |
| 0.0897 | 440 | 17.0137 | 0.0386 | 0.0777 | 3.0879 | 12785 | 123.8379 | 0.3597 | 0.5855 | 0.7181 | 0.9129 |
| 0.0985 | 440 | 17.0168 | 0.0387 | 0.0727 | 3.19624 | 12830 | 118.2969 | 0.3779 | 0.6150 | 0.7433 | 0.9236 |
| 0.1074 | 439 | 17.3033 | 0.0394 | 0.0720 | 3.22421 | 12875 | 116.9171 | 0.3837 | 0.6244 | 0.7498 | 0.9262 |
| 0.1165 | 440 | 17.7787 | 0.0404 | 0.0694 | 3.28023 | 12921 | 114.2139 | 0.3942 | 0.6415 | 0.7628 | 0.9313 |
| 0.1253 | 440 | 17.8884 | 0.0406 | 0.0666 | 3.34619 | 12965 | 111.1318 | 0.4065 | 0.6616 | 0.7782 | 0.9371 |
| 0.1345 | 440 | 18.4004 | 0.0418 | 0.0634 | 3.42552 | 13011 | 107.5642 | 0.4215 | 0.6859 | 0.7966 | 0.9438 |
| 0.1433 | 440 | 19.1654 | 0.0435 | 0.0621 | 3.47127 | 13056 | 105.5736 | 0.4309 | 0.7013 | 0.8073 | 0.9476 |
| 0.1524 | 441 | 19.5067 | 0.0442 | 0.0597 | 3.53557 | 13102 | 102.8556 | 0.4438 | 0.7223 | 0.8222 | 0.9526 |
| 0.1614 | 441 | 19.7657 | 0.0448 | 0.0592 | 3.56141 | 13147 | 101.7890 | 0.4500 | 0.7324 | 0.8282 | 0.9546 |
| 0.1705 | 441 | 20.6800 | 0.0469 | 0.0563 | 3.63631 | 13193 | 98.7776 | 0.4654 | 0.7574 | 0.8457 | 0.9601 |
| 0.1795 | 440 | 20.7227 | 0.0471 | 0.0547 | 3.69534 | 13238 | 96.4861 | 0.4781 | 0.7780 | 0.8594 | 0.9643 |
| 0.1885 | 440 | 21.1250 | 0.0481 | 0.0541 | 3.72719 | 13284 | 95.2787 | 0.4858 | 0.7906 | 0.8668 | 0.9665 |
| 0.1976 | 440 | 21.9204 | 0.0498 | 0.0506 | 3.85439 | 13329 | 90.6512 | 0.5123 | 0.8338 | 0.8964 | 0.9749 |
| 0.2066 | 440 | 22.0667 | 0.0501 | 0.0505 | 3.86806 | 13375 | 90.1718 | 0.5168 | 0.8411 | 0.8995 | 0.9758 |
| 0.2157 | 440 | 22.3654 | 0.0508 | 0.0491 | 3.90957 | 13421 | 88.7367 | 0.5270 | 0.8576 | 0.9092 | 0.9784 |
| 0.2248 | 440 | 23.0755 | 0.0524 | 0.0476 | 3.96353 | 13467 | 86.9163 | 0.5399 | 0.8786 | 0.9218 | 0.9817 |
| 0.2337 | 439 | 23.5174 | 0.0535 | 0.0461 | 4.02285 | 13512 | 84.9725 | 0.5540 | 0.9017 | 0.9355 | 0.9852 |
| 0.2429 | 441 | 24.1452 | 0.0547 | 0.0454 | 4.06453 | 13558 | 83.6416 | 0.5648 | 0.9192 | 0.9452 | 0.9875 |
| 0.2518 | 441 | 24.9436 | 0.0565 | 0.0433 | 4.17042 | 13603 | 80.3844 | 0.5896 | 0.9596 | 0.9699 | 0.9934 |
| 0.2610 | 441 | 25.2575 | 0.0573 | 0.0432 | 4.18511 | 13649 | 79.9462 | 0.5949 | 0.9681 | 0.9733 | 0.9941 |
| 0.2701 | 441 | 25.5410 | 0.0579 | 0.0421 | 4.22355 | 13695 | 78.8147 | 0.6054 | 0.9854 | 0.9822 | 0.9961 |
| 0.2792 | 440 | 26.2968 | 0.0597 | 0.0415 | 4.25788 | 13741 | 77.8226 | 0.6152 | 1.0013 | 0.9902 | 0.9979 |
| 0.2882 | 439 | 26.3303 | 0.0599 | 0.0420 | 4.25704 | 13787 | 77.8466 | 0.6171 | 1.0043 | 0.9900 | 0.9978 |
| 0.2970 | 440 | 26.7600 | 0.0608 | 0.0412 | 4.29793 | 13831 | 76.6865 | 0.6284 | 1.0227 | 0.9995 | 0.9999 |
| 0.3060 | 440 | 27.5036 | 0.0625 | 0.0411 | 4.30498 | 13876 | 76.4888 | 0.6321 | 1.0288 | 1.0012 | 1.0002 |
| 0.3153 | 439 | 27.9730 | 0.0637 | 0.0412 | 4.30122 | 13923 | 76.5941 | 0.6334 | 1.0308 | 1.0003 | 1.0001 |
| 0.3243 | 441 | 28.0918 | 0.0637 | 0.0412 | 4.30247 | 13968 | 76.5591 | 0.6357 | 1.0346 | 1.0006 | 1.0001 |
| 0.3334 | 440 | 28.2869 | 0.0642 | 0.0410 | 4.30987 | 14014 | 76.3521 | 0.6395 | 1.0408 | 1.0023 | 1.0005 |
| 0.3426 | 442 | 28.8141 | 0.0652 | 0.0414 | 4.28803 | 14061 | 76.9652 | 0.6366 | 1.0360 | 0.9972 | 0.9994 |
| 0.3516 | 440 | 29.1859 | 0.0663 | 0.0414 | 4.29028 | 14106 | 76.9017 | 0.6391 | 1.0402 | 0.9977 | 0.9995 |
| 0.3607 | 442 | 29.1463 | 0.0660 | 0.0412 | 4.29856 | 14152 | 76.6688 | 0.6432 | 1.0467 | 0.9997 | 0.9999 |
| 0.3698 | 441 | 29.5455 | 0.0669 | 0.0418 | 4.26748 | 14198 | 77.5481 | 0.6379 | 1.0382 | 0.9924 | 0.9984 |
| 0.3789 | 441 | 29.7772 | 0.0675 | 0.0418 | 4.26925 | 14244 | 77.4977 | 0.6404 | 1.0422 | 0.9928 | 0.9985 |
| 0.3879 | 441 | 30.0819 | 0.0682 | 0.0414 | 4.28896 | 14289 | 76.9389 | 0.6471 | 1.0531 | 0.9974 | 0.9995 |
| 0.3969 | 442 | 30.2648 | 0.0685 | 0.0414 | 4.28701 | 14335 | 76.9940 | 0.6487 | 1.0557 | 0.9970 | 0.9994 |
| 0.4061 | 441 | 30.2678 | 0.0687 | 0.0413 | 4.29231 | 14381 | 76.8445 | 0.6521 | 1.0612 | 0.9982 | 0.9996 |
| 0.4153 | 440 | 30.3654 | 0.0691 | 0.0408 | 4.3187 | 14427 | 76.1062 | 0.6605 | 1.0750 | 1.0043 | 1.0009 |
| 0.4242 | 440 | 30.3654 | 0.0690 | 0.0405 | 4.33843 | 13513 | 75.5605 | 0.6231 | 1.0141 | 1.0089 | 1.0019 |

Table C-37: Mach 5.0 CPG $\left(\mathrm{X}_{\mathrm{ts}}=6.35 \mathrm{~cm}\right)$ pressure data

| $y$ (in) | $\mathrm{P}_{11}$ (psia) | $\mathrm{P}_{12}$ (psia) | $\mathrm{P}_{\mathrm{t} 2} / \mathrm{P}_{\mathrm{t} 1}$ | $\mathrm{ps} / \mathrm{P}_{12}$ | Mach | $\mathrm{P}(\mathrm{Pa})$ | T (K) | $\rho$ | $\rho / \mathrm{pe}$ | $\mathrm{M} / \mathrm{M}_{\mathrm{e}}$ | U/ U ${ }_{\text {e }}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.0625 | 507 | 27.8481 | 0.0550 | 0.0738 | 3.210 | 11599.7 | 117.62 | 0.3436 | 1.0582 | 0.6173 | 0.8932 |
| 0.0712 | 506 | 29.1769 | 0.0577 | 0.0695 | 3.287 | 11243.4 | 113.87 | 0.3440 | 1.0595 | 0.6322 | 0.9001 |
| 0.0797 | 507 | 30.2070 | 0.0595 | 0.0662 | 3.366 | 10891.2 | 110.23 | 0.3443 | 1.0603 | 0.6473 | 0.9067 |
| 0.0883 | 507 | 30.6184 | 0.0604 | 0.0644 | 3.415 | 10539.0 | 108.04 | 0.3399 | 1.0467 | 0.6567 | 0.9107 |
| 0.0968 | 505 | 33.1693 | 0.0656 | 0.0586 | 3.563 | 10185.6 | 101.72 | 0.3489 | 1.0745 | 0.6852 | 0.9220 |
| 0.1052 | 506 | 32.9346 | 0.0650 | 0.0583 | 3.595 | 9839.2 | 100.41 | 0.3414 | 1.0515 | 0.6914 | 0.9244 |
| 0.1138 | 504 | 33.0931 | 0.0656 | 0.0572 | 3.630 | 9485.7 | 99.03 | 0.3338 | 1.0279 | 0.6981 | 0.9268 |
| 0.1224 | 505 | 34.5224 | 0.0683 | 0.0540 | 3.722 | 9132.3 | 95.49 | 0.3332 | 1.0262 | 0.7157 | 0.9331 |
| 0.1310 | 505 | 36.5582 | 0.0724 | 0.0502 | 3.876 | 8780.1 | 89.91 | 0.3403 | 1.0479 | 0.7453 | 0.9429 |
| 0.1396 | 505 | 37.7894 | 0.0748 | 0.0479 | 3.953 | 8425.0 | 87.26 | 0.3364 | 1.0361 | 0.7602 | 0.9475 |
| 0.1482 | 506 | 39.4687 | 0.0780 | 0.0452 | 4.058 | 8070.3 | 83.85 | 0.3354 | 1.0328 | 0.7804 | 0.9534 |
| 0.1568 | 505 | 39.6028 | 0.0784 | 0.0443 | 4.117 | 7715.3 | 82.02 | 0.3278 | 1.0094 | 0.7916 | 0.9565 |
| 0.1654 | 506 | 42.0744 | 0.0832 | 0.0411 | 4.284 | 7361.8 | 77.09 | 0.3328 | 1.0248 | 0.8238 | 0.9650 |
| 0.1739 | 506 | 43.1654 | 0.0853 | 0.0394 | 4.356 | 7009.6 | 75.07 | 0.3253 | 1.0019 | 0.8377 | 0.9684 |
| 0.1825 | 506 | 44.1955 | 0.0874 | 0.0379 | 4.437 | 6657.4 | 72.91 | 0.3181 | 0.9798 | 0.8533 | 0.9721 |
| 0.1911 | 505 | 45.9114 | 0.0909 | 0.0359 | 4.584 | 6301.1 | 69.19 | 0.3173 | 0.9773 | 0.8816 | 0.9784 |
| 0.1997 | 505 | 46.2740 | 0.0916 | 0.0350 | 4.635 | 5947.7 | 67.96 | 0.3049 | 0.9391 | 0.8914 | 0.9804 |
| 0.2083 | 506 | 47.5815 | 0.0940 | 0.0335 | 4.716 | 5595.5 | 66.08 | 0.2950 | 0.9086 | 0.9069 | 0.9836 |
| 0.2170 | 504 | 48.9681 | 0.0971 | 0.0319 | 4.855 | 5237.5 | 62.99 | 0.2897 | 0.8922 | 0.9337 | 0.9887 |
| 0.2256 | 506 | 49.7849 | 0.0984 | 0.0309 | 4.920 | 4882.5 | 61.63 | 0.2760 | 0.8501 | 0.9461 | 0.9910 |
| 0.2342 | 507 | 50.7571 | 0.1001 | 0.0297 | 4.998 | 4529.0 | 60.04 | 0.2628 | 0.8094 | 0.9611 | 0.9936 |
| 0.2429 | 507 | 50.4675 | 0.0996 | 0.0294 | 5.053 | 4168.2 | 58.96 | 0.2463 | 0.7586 | 0.9717 | 0.9954 |
| 0.2515 | 506 | 50.4614 | 0.0996 | 0.0288 | 5.103 | 3814.7 | 57.98 | 0.2292 | 0.7060 | 0.9814 | 0.9970 |
| 0.2601 | 505 | 49.8763 | 0.0987 | 0.0286 | 5.138 | 3458.4 | 57.32 | 0.2102 | 0.6475 | 0.9882 | 0.9981 |
| 0.2688 | 507 | 50.0378 | 0.0987 | 0.0280 | 5.185 | 3100.5 | 56.46 | 0.1913 | 0.5893 | 0.9971 | 0.9995 |
| 0.2775 | 506 | 48.5628 | 0.0960 | 0.0283 | 5.196 | 2745.4 | 56.25 | 0.1701 | 0.5237 | 0.9993 | 0.9999 |
| 0.2860 | 506 | 47.5479 | 0.0940 | 0.0283 | 5.201 | 2393.2 | 56.17 | 0.1485 | 0.4572 | 1.0001 | 1.0000 |
| 0.2944 | 504 | 45.4664 | 0.0902 | 0.0290 | 5.171 | 2045.5 | 56.70 | 0.1257 | 0.3871 | 0.9945 | 0.9991 |

Table C-38: Mach 5.0 wall static pressure data

| CPG - $\mathrm{X}_{\text {ts }}$ | $\mathrm{X}_{\text {s }} / \mathrm{L}$ | ps (psia) | $P_{11}$ | $\mathrm{P}_{11} / \mathrm{ps}$ |
| :---: | :---: | :---: | :---: | :---: |
| 0.750 | 0.0750 | 1.4 | 477 | 341 |
| 1.250 | 0.1250 | 1.295 | 477 | 368 |
| 1.840 | 0.1840 | 1.526 | 477 | 313 |
| 2.375 | 0.2375 | 2.048 | 485 | 237 |
| 2.813 | 0.2813 | 2.058 | 485 | 236 |
| 3.313 | 0.3313 | 2.001 | 485 | 242 |
| 3.750 | 0.3750 | 1.451 | 500 | 345 |
| 4.250 | 0.4250 | 1.361 | 500 | 367 |
| 4.750 | 0.4750 | 1.546 | 500 | 323 |
|  |  |  |  |  |
| FPG - $\mathrm{X}_{\text {ts }}$ | $\mathrm{X}_{\text {ts }} / \mathrm{L}$ | ps (psia) | $\mathrm{P}_{11}$ | $\mathrm{P}_{11} / \mathrm{ps}$ |
| 1.000 | 0.1000 | 1.224 | 493 | 402 |
| 1.438 | 0.1438 | 1.061 | 493 | 464 |
| 2.438 | 0.2438 | 0.789 | 486 | 616 |
| 2.938 | 0.2938 | 0.69 | 486 | 704 |
| 3.313 | 0.3313 | 7.295 | 486 | 67 |
| 3.906 | 0.3906 | 0.92 | 479 | 521 |
| 4.375 | 0.4375 | 1.092 | 479 | 439 |
| 4.906 | 0.4906 | 1.471 | 479 | 326 |

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| 4. The fic subtios <br> EXPERIMENTAL INVESTIGATION OF COMPRESSIBLE <br> BOUNDARY LAYERS UNDER THE INFLUENCE OF <br> PRESSURE GRADIENTS |
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| 1. Suenturn 16TES |
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| 1. RESThaE (iverimum 200 words <br> This study examined the effect of mild pressure gradients on the mean and turbulent flow of high-speed boundary layers. Three Mach numbers (1.7, 3.0 and 5.0 ) were investigated. Three pressure gradients were examined; a zero pressure gradient (ZPG), a favorable pressure gradient (FPG), and a combined pressure gradient (CPG). The CPG consisted of an adverse pressure gradient followed by a favorable pressure gradient. Conventional pressure probes, hot-wire and particle image velocimetry (PIV) were used to examine the flow. Measurement included mean velocity, velocity turbulence intensity, mass flux turbulence intensity and energy spectra. Instantaneous ( 10 nsec ) Mie scattering flow visualizations were acquired. Qualitatively, the flow visualizations indicated that the turbulent flow structures were strongly affected by the pressure gradients. For the Mach 2.8 case, the PIV contours and the hot-wire profiles both indicated that the boundary layer thickness increased by $40 \%$ and decreased by $100 \%$ relative to the ZPG for the favorable and adverse pressure gradients, respectively. Further, the PIV and hot-wire data indicated that the axial turbulence intensity levels increased by $22 \%$ for the CPG and decreased by $25 \%$ for the FPG. The energy spectra data indicated that once a pressure gradient was applied (favorable or adverse) the low frequency energy increased followed by a rapid decay. Lastly, it was found that nominally 20 to 30 PIV images were sufficient for mean flow boundary layer velocities, but 93 images (the maximum recorded in this study) were insufficient to adequately resolve Reynolds shear stresses. |
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