

## ENVIRONMENT AND GALAXY EVOLUTION AT INTERMEDIATE REDSHIFT IN THE CNOC2 SURVEY

R. G. CARLBERG,<sup>1,2</sup> H. K. C. YEE,<sup>1,2</sup> S. L. MORRIS,<sup>1,3,4</sup> H. LIN,<sup>1,2,5,6</sup> P. B. HALL,<sup>1,2,7</sup> D. R. PATTON,<sup>1,2,8</sup>  
M. SAWICKI,<sup>1,2,9</sup> AND C. W. SHEPHERD<sup>1,2</sup>

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### ABSTRACT

The systematic variation of galaxy colors and types with clustering environment could either be the result of local conditions at formation or subsequent environmental effects as larger scale structures draw together galaxies whose stellar mass is largely in place. Below redshift 0.7 galaxy luminosities ( $k$ -corrected and evolution compensated) are relatively invariant, whereas galaxy star formation rates, as reflected in their colors, are a “transient” property that have a wide range for a given luminosity. The relations between these galaxy properties and the clustering properties are key statistics for understanding the forces driving late-time galaxy evolution. At  $z \sim 0.4$  the comoving galaxy correlation length,  $r_0$ , measured in the CNOC2 sample is strongly color dependent, rising from  $2 h^{-1}$  Mpc to nearly  $10 h^{-1}$  Mpc as the volume-limited subsamples range from blue to red. The luminosity dependence of  $r_0$  at  $z \sim 0.4$  is weak below  $L_*$  in the  $R$  band, although there is an upturn at high luminosity, where its interpretation depends on separating it from the  $r_0$ -color relation. In the  $B$  band there is a slow, smooth increase of  $r_0$  with luminosity, at least partially related to the color dependence. Study of the evolution of galaxies within groups, which create much of the strongly nonlinear correlation signal, allows a physical investigation of the source of these relations. The dominant effect of the group environment on star formation is seen in the radial gradient of the mean galaxy colors, which on the average become redder than the field toward the group centers. The color differentiation begins around the dynamical radius of virialization of the groups. The redder-than-field trend applies to groups with a line-of-sight velocity dispersion,  $\sigma_1 > 150 \text{ km s}^{-1}$ . There is an indication, somewhat statistically insecure, that the high-luminosity galaxies in groups with  $\sigma_1 < 125 \text{ km s}^{-1}$  become bluer toward the group center. Monte Carlo orbit integrations initiated at the measured positions and velocities show that the rate of galaxy merging in the  $\sigma_1 > 150 \text{ km s}^{-1}$  groups is very low, whereas for  $\sigma_1 < 150 \text{ km s}^{-1}$  about 25% of the galaxies will merge in 0.5 Gyr. We conclude that the higher velocity dispersion groups largely act to suppress star formation relative to the less clustered field, leading to “embalmed” galaxies. On the other hand, the low velocity dispersion groups are prime sites of both strong merging and enhanced star formation that leads to the formation of some new massive galaxies at intermediate redshifts. The tidal fields within the groups appear to be a strong candidate for the physical source of the reduction of star formation in group galaxies relative to field. Tides operate effectively at all velocity dispersions to remove gas-rich companions and low-density gas in galactic halos. We find a close resemblance of the color-dependent galaxy luminosity function evolution in the field and groups, suggesting that the clustering-dependent star formation reduction mechanism is important for the evolution of field galaxies as a whole.

*Subject headings:* galaxies: evolution — large-scale structure of universe

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### 1. INTRODUCTION

Environment plays a major role in the formation and evolution of galaxies, as is dramatically manifest in the differences between galaxies in rich clusters and the field

<sup>1</sup> Visiting Astronomer, Canada-France-Hawaii Telescope, which is operated by the National Research Council of Canada, le Centre National de Recherche Scientifique, and the University of Hawaii.

<sup>2</sup> Department of Astronomy, University of Toronto, Toronto, ON M5S 3H8, Canada.

<sup>3</sup> Dominion Astrophysical Observatory, Herzberg Institute of Astrophysics, National Research Council of Canada, 5071 West Saanich Road, Victoria, BC V8X 4M6, Canada.

<sup>4</sup> Department of Physics, University of Durham, South Road, Durham DH1 3LE, UK.

<sup>5</sup> Steward Observatory, University of Arizona, Tucson, AZ 85721.

<sup>6</sup> Hubble Fellow.

<sup>7</sup> Princeton University Observatory, Princeton, NJ 08544-1001; and Pontificia Universidad Católica de Chile, Departamento de Astronomía y Astrofísica, Facultad de Física, Casilla 306, Santiago 22, Chile.

<sup>8</sup> Department of Physics and Astronomy, University of Victoria, Victoria, BC V8W 3P6, Canada.

<sup>9</sup> Mail Code 320-47, California Institute of Technology, Pasadena, CA 91125.

(Dressler 1980). The environmental effect is mainly one to alter the distribution over galaxy type, with even the somewhat exotic E+A galaxies being found in the field (Zabludoff et al. 1996). A notable exception is the cD galaxies, but even those are believed to originate from an early epoch field population (Dubinski 1998). It is ambiguous at any single redshift whether the systematically redder and earlier morphological types of cluster galaxies, as compared to the field galaxy population, are the result of different local conditions at the time of formation, or, whether all galaxies begin as a fairly uniform distribution of types that subsequent evolutionary forces differentiate as a function of environment. Examining the redshift dependence of galaxy populations is a powerful tool to separate low-redshift environmentally driven changes from high-redshift formation mechanisms as the primary cause of the differences.

An example of the study of the environmental dependence of populations with redshift is the star formation rate of cluster galaxies at intermediate redshift. On the average, cluster galaxy star formation rates are smaller than those in field galaxies, with a progressive decline to the cluster center

(Dressler 1980; Whitmore Gilmore, & Jones 1993; Balogh et al. 1999; Poggianti et al. 1999; Ellingson et al. 2001). This suggests that a field “mix” of galaxies is capable of being transformed into an earlier type cluster population in response to the cluster environment. Cluster galaxies provide important clues of environmental forcing, but they are a moderately rare population subject to exceptionally strong environmental forces.

Galaxy groups show many of the differences from the field galaxy population that cluster galaxies exhibit (Zabludoff & Mulchaey 1998; Hashimoto & Oemler 2000) (Tran, Simard, & Zabludoff 2001). The evolution of group galaxies is necessarily a significant component of the evolution of the entire galaxy population. Approximately 20% of  $L \gtrsim L_*/3$  galaxies reside in small virialized galaxy groups (Nolthenius & White 1987; Carlberg et al. 2001), and virtually all galaxies are in groups if the luminosity limit for inclusion is lowered sufficiently. For instance, the Milky Way, with the Large and Small Magellanic Clouds, constitutes a small virialized group (which falls just slightly short of meeting the luminosity requirements of the groups we define below). The Local Group also usefully illustrates that the lowest velocity dispersion groups overlap the velocity dispersion range of individual and pairs of galaxies. The mean pairwise one-dimensional peculiar velocity dispersion of galaxies is 300–500 km s<sup>-1</sup> (Davis & Peebles 1983; Marzke et al. 1995; Peacock et al. 2001; Zehavi et al. 2001), indicating that groups with one-dimensional velocity dispersions,  $\sigma_1$ , in the range of 200–300 km s<sup>-1</sup> dominate the clustering of field galaxies.

Some combination of internal clocks and external agents drives galaxy evolution. Internal timescales are regulated by dynamical and thermal processes, with generally fairly fast clocks, whereas the external timescales are controlled by the structures in which galaxies find themselves and the rate of development of large-scale clustering, clocks running at approximately the Hubble rate. At a given galaxy type, both field and group galaxies have essentially identical internal processes. Consequently, clustering-dependent population differences must be due to either conditions at formation or subsequent environmental interactions. Galaxy formation is not a process that occurs at a sharply defined time. However, for the intermediate-redshift range we examine here normal luminous galaxies of a wide range of type can be taken to essentially “formed” with varying degrees of ongoing star formation and merging that do not grossly alter the stellar mass of the galaxies.

Gravitational tides, dynamical friction, and intergalactic gas pressure are two physical forces present in groups that can affect their galaxy contents. Small galaxy groups have been investigated extensively at low redshift (Burbidge & Burbidge 1960; Huchra & Geller 1982; Ramella, Geller, & Huchra 1989; Mulchaey & Zabludoff 1998, Mahdavi & Geller 2001) and at intermediate redshift (Carlberg et al. 2001) to show that the total mass of a virialized group is dominated by dark matter. One action of the dark matter is to cause dynamical friction, whose in-spiral timescale varies in proportion to  $\sigma^3/\rho$ . Detailed dynamical analysis (Merritt 1984, 1985) and  $n$ -body simulations (Mamon 1987; Dubinski 1998) provide quantitative evidence that high velocity dispersion clusters do not suffer much internal dynamical evolution of their galaxy populations after their primary formation phase. However, low velocity dispersion groups are found to produce mergers on a cosmologically short

timescale, even at low redshift, where the virialization density, approximately 200 times the critical density, is at a minimum (Barnes 1985, 1990; Bode, Cohn, & Lugger 1993) (Bode et al. 1994; Mamon 1987; Diaferio, Geller, & Ramella 1994) (Moore, Lake, & Katz 1998). Galaxy groups also contain X-ray-emitting hot gas that work on the gas content of galaxies (Mulchaey 2000). Ram pressure stripping declines in proportion to the velocity dispersion (squared), so plays a relatively less important role in groups than in large clusters (Abadi, Moore, & Bower 1999; Balogh, Navarro, & Morris 2000).

The evidence that groups play a significant role in galaxy evolution is strong, but the observationally confirmed effects of groups vary with the group definition and are largely confined to low-redshift samples. For instance, the Hickson (1982) compact groups are deliberately designed to be quite high-density groups. They are found to include a number of strongly interacting and likely-to-be-merging systems with enhanced star formation (Hickson 1997; Hashimoto & Oemler 2000). The ultimate fate predicted for the galaxies in isolated low velocity dispersion groups is to join to form a single massive elliptical galaxy (Barnes 1990; Mulchaey & Zabludoff 1999). Groups defined with a lower density threshold are observed to have suppressed star formation relative to the field (Allam et al. 1999; Tran et al. 2001) reminiscent of the effect that is seen in rich clusters.

An empirical study of groups uniformly selected from the field at  $z \sim 0.5$  opens up the possibility of detecting differential evolutionary effects between groups and the field. Galaxy evolution over the range from  $z = 0$  to  $z \sim 0.7$  is substantial (Broadhurst, Ellis, & Shanks 1988), particularly in relatively blue galaxies less luminous than  $L_*$  (Lilly et al. 1995; Lin et al. 1999). The physical origin of this evolution remains uncertain, although it is clear that in individual galaxies it is superimposed on a slowly aging underlying stellar population and that it primarily requires regulation of the star formation rate (Lilly et al. 1996). Galaxy groups in this redshift range have been defined with photometric data (de Mello, Infante, & Menanteau 1997) and with radio galaxies as centers (Allington-Smith et al. 1993), but large, kinematically accurate redshift catalogs are needed to put intermediate-redshift groups on the same standing as those at low redshift.

In the next section we review the Canadian Network for Observational Cosmology field galaxy redshift survey (CNOC2) sample, which contains about 6000 galaxies with velocities measured to a precision of about 70–100 km s<sup>-1</sup> over the 0.1–0.55 redshift range. The sample was selected to minimize clustering-dependent differential selection effects so that the relative evolution of group and field galaxies can be readily examined. We first use the field sample as a whole to show the correlation function’s strong dependence on color and relatively weak luminosity dependence. § 3 discusses the color and luminosity distribution of field and group galaxies and their change with redshift. The variation of the mean galaxy colors with distance from the group center is discussed for high and low velocity dispersion groups in § 4. The orbits of the group galaxies are Monte Carlo integrated in § 5 to determine their likelihood of merging. Section 6 discusses the implications of our findings for galaxy evolution and states our conclusions. Most quantities in this paper are derived with  $H_0 = 100 h$  km s<sup>-1</sup> Mpc<sup>-1</sup> and an  $\Omega_M = 0.2$ ,  $\Omega_\Lambda = 0$  cosmology, with exceptions noted.

## 2. THE GALAXY CLUSTERING-COLOR RELATION AT $z \sim 0.4$

The CNOC2 field galaxy redshift survey covers approximately 1.5 square degrees of sky spread over four widely separated sky patches (Yee et al. 2000). The unbiased redshift sample covers the 0.1–0.55 redshift range. The survey has an accurately measured selection function that varies from nearly 100% at  $R = 19$  mag down to about 20% at the limit of  $R = 21.5$  mag. The selection function is used to correct the spectroscopic sample numbers back to what a complete sample would, on the average, contain. For each galaxy  $k$ -corrections are evaluated in all observed colors using the closest approximating Coleman, Wu, & Weedman (1980) spectral energy distribution (SED) using a finely interpolated set of SEDs. The evolution of the  $k$ -corrected galaxy luminosity function as a function of non-evolving SED type has been discussed elsewhere (Lin et al. 1999). We use these results to define a luminosity evolution compensation term, intended to identify statistically the same population over our entire redshift range. This is valid for the population as a whole or two color subclasses that are our samples of interest. An accurate galaxy-by-galaxy correction requires more information, ideally using a mass indicator such as IR luminosity or dynamical masses. The  $k$ -corrected and evolution-compensated  $R$ -band luminosities,  $M_R^{ke}$ , are calculated as  $M_R - k_R(z) + E_z z$ , where the mean evolution rate,  $E_z$ , is approximately compensated with a fading of 1 mag per unit redshift for all SED types. This approach is statistically acceptable for the relatively high-luminosity galaxies we use in this paper, but cannot be extended to lower luminosity blue galaxies, which require either an allowance for more luminosity evolution or number evolution.

The central observational results presented in this paper address the dependence of colors on clustering environment. Within our photometric data set the color that best serves as our star formation rate indicator is the  $k$ -corrected color  $(B-R)_0 = (B-R) - k_B + k_R$ . We are motivated to use this color because all our galaxies have  $B$  and  $R$  observations, with less deep coverage in  $U$  and a limited redshift range for the  $[O\ II] \lambda 3727$  line. Of course  $(B-R)_0$  has a dependence on factors other than star formation alone, such as metallicity, dust, and the details of the star formation history. Our study depends on a comparison of various subsamples where  $(B-R)_0$  serves to indicate the relative star formation rates.

To explicitly relate the colors to star formation rates we use the PEGASE.2 models (Fioc & Rocca-Volmerange 1997). We also examine the stellar mass-to-light ratio,  $M/L$ , for its population dependence. Figure 1 plots the star formation rate and the mass-to-light ratio,  $M/L$ , in the  $R$  and  $B$  bands, against the rest-frame  $B-R$  color. Points are plotted every  $3 \times 10^8$  model years of age. All the photometric models are assumed to be dust free and have exponentially declining star formation rates,  $e^{-t/\tau}$ , with  $\tau = 1, 2, 4$  and 6 Gyr at solar metallicity and a  $\tau = 1$  Gyr at twice solar. The colors are plotted for model ages of 3 to 12 Gyr. The figure shows that  $(B-R)_0$  color (assuming no dust) indicates the current star formation rate, largely independent of the details of the history of the star formation, although we see the well-known effect that the metallicity is important in the color of galaxies with low star formation rates. We note that the  $M/L$  value of the stellar population

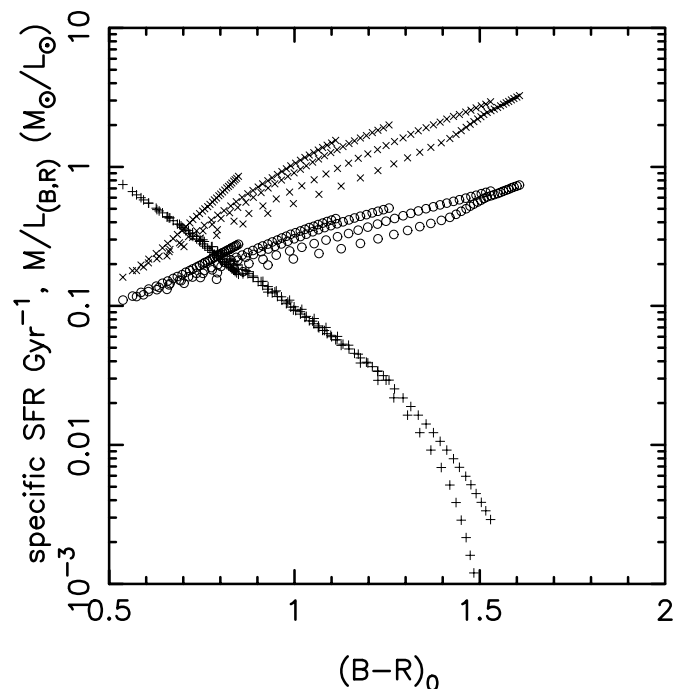


FIG. 1.—Relation between rest frame  $(B-R)_0$  and the instantaneous star formation rate, normalized to the mass, which falls as  $(B-R)_0$  increases, is shown as the plus signs. The mass-to-light ratio is shown, using the same axes, for blue (crosses) and red (circles) rest-frame luminosities. A variety of exponential models are shown for solar and 2 times solar (the reddest set of model results) abundance. Points are plotted every  $3 \times 10^8$  yr. [See the electronic edition of the Journal for a color version of this figure.]

varies about a factor of 2 less in the  $R$  band than in the  $B$  band.

Our analysis of the redshift dependence of the two-point correlation function of the CNOC2 sample finds that the correlation length,  $r_0(z)$ , is nearly constant in comoving coordinates (Carlberg et al. 2000a). The sample is volume-limited on the basis of  $k$ -corrected and evolution-compensated luminosities,  $M_R^{ke} \leq -20$  mag. Our correlation measurements are consistent with the theoretical predictions of a slow increase of biasing of galaxies relative to the dark matter increasing with redshift.

The intermediate-redshift clustering of the CNOC2 sample varies strongly with the mean SED type of the galaxies (Shepherd et al. 2001), with late type galaxies being far less clustered than early type galaxies. The clustering analysis divides galaxies based on nonevolving SED types; hence, the dependence of clustering with type evolves with redshift as galaxies move to later types with increasing redshift (Shepherd et al. 2001). Here we complement our earlier analysis with a larger sample, in a narrower redshift range. We are particularly interested in the relative sensitivity of  $r_0$  to star formation and luminosity. We use the redshift range 0.25–0.45, in which the bulk of the CNOC2 sample falls. Figure 2 shows the dependence of the comoving correlation length on  $(B-R)_0$ , analyzed in a flat  $\Omega_M = 0.2$ ,  $\Omega_\Lambda = 0.8$  cosmology ( $\Lambda = 0$  gives  $r_0$  that are about 10% smaller at the average redshift of 0.39). The sample is volume-limited, with  $M_R^{ke} \leq -18.5$  mag. The correlation measurements are made with exactly the same procedures as discussed in Carlberg et al. (2000a, 2000b). We fit the measured projected correlation function to the projection of a power-law correlation function,  $(r_0/r)^\gamma$ , where all lengths are measured in

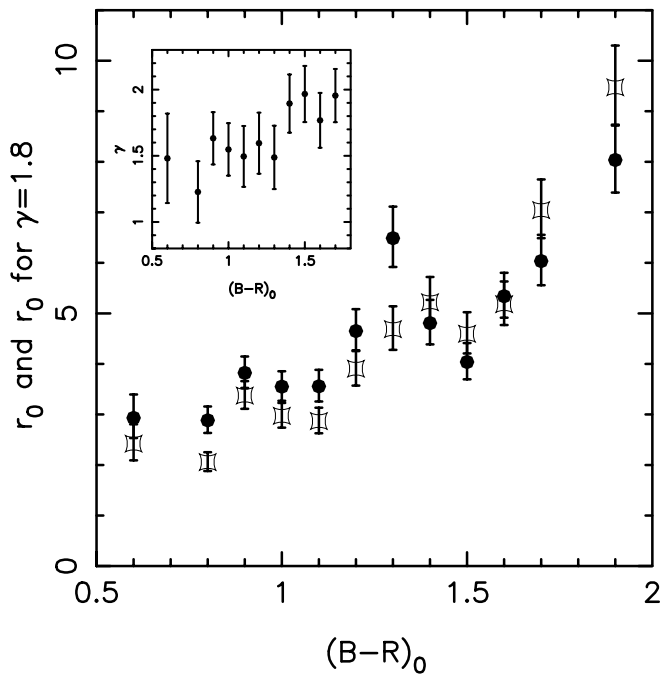


FIG. 2.—The fitted  $r_0$  (filled circles) as a function of the  $k$ -corrected color  $(B-R)_0$ . The fitted slopes,  $\gamma$ , are shown in the inset. The correlation lengths normalized to  $\gamma = 1.8$  are shown as the pincushion symbols. The color bins are 50% overlapped in width.

comoving units. The derived  $r_0$  and  $\gamma$  are shown in Figure 2 along with the errors estimated on the basis of the variance of the four observational patches in the sample. The slope of the correlation function,  $\gamma$ , becomes steeper with increasingly red color. Both of these effects are known at lower redshift as morphological type or emission/absorption line

galaxy correlations, although the large range of correlation amplitudes we find is not so readily apparent at low redshift (Davis & Geller 1976; Saunders, Rowan-Robinson, & Lawrence 1992; Fisher et al. 1994; Guzzo et al. 1997; Willmer, da Costa, & Pellegrini 1998; Loveday, Tresse, & Maddox 1999). The most relevant measure of correlation amplitude in a situation where the  $\gamma$  vary systematically is to show the correlation length normalized to a fixed  $\gamma$ , here chosen to be  $\gamma = 1.8$ . As shown in Figure 2, these values range from about  $2 h^{-1}$  Mpc for the bluest galaxies to nearly  $10 h^{-1}$  Mpc for the reddest galaxies. The errors in these figures are estimated using the patch to patch variance over our four widely separated patches, which is a robust empirical estimator.

Shepherd et al. (2001) found that there is a weak luminosity increase (or none) of clustering within each SED class over a luminosity range of 1.5 mag. Here we examine the luminosity dependence of clustering over a 3.0 mag range, displaying the result in Figure 3. Recall that in the  $R$  band  $M_*^{ke}(z = 0.3) \simeq -20.3$  mag and in the  $B$  band  $M_*^{ke}(z = 0.3) \simeq -19.1$  mag. At lower luminosities,  $0.2-2 L_*$ , the figure shows that  $r_0$  could rise about 10% for the factor of 10 luminosity increase (for  $\gamma = 1.8$ ), but the difference is statistically consistent with an even weaker relation. Galaxies more luminous than  $2 L_*$  show a highly significant increase in clustering, rising from  $4 h^{-1}$  Mpc for  $L \sim L_*$  to  $8 h^{-1}$  Mpc. However, the mean color in this bin is  $(B-R)_0 = 1.55$  mag, which is about 0.25 mag redder than the lower luminosity members of the sample. Referring to the color dependence of Figure 2 indicates that at a mean color of  $(B-R)_0 = 1.7$  we expect  $r_0 \simeq 7.2 \pm 0.6 h^{-1}$  Mpc. The highest luminosity galaxies here have  $r_0 \simeq 8.0 \pm 0.7 h^{-1}$  Mpc. Consequently the increase in clustering in this high-luminosity sample could be largely the result of the correlation-color relation. This is difficult to disentangle,

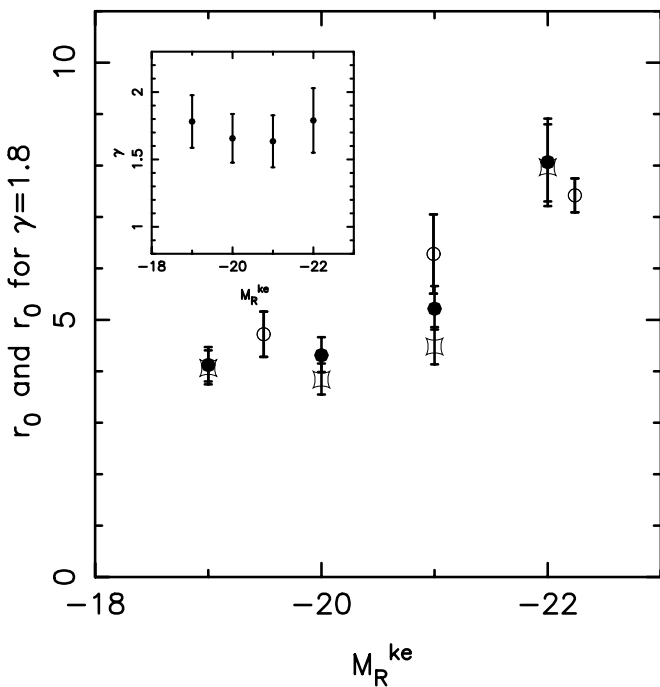


FIG. 3a

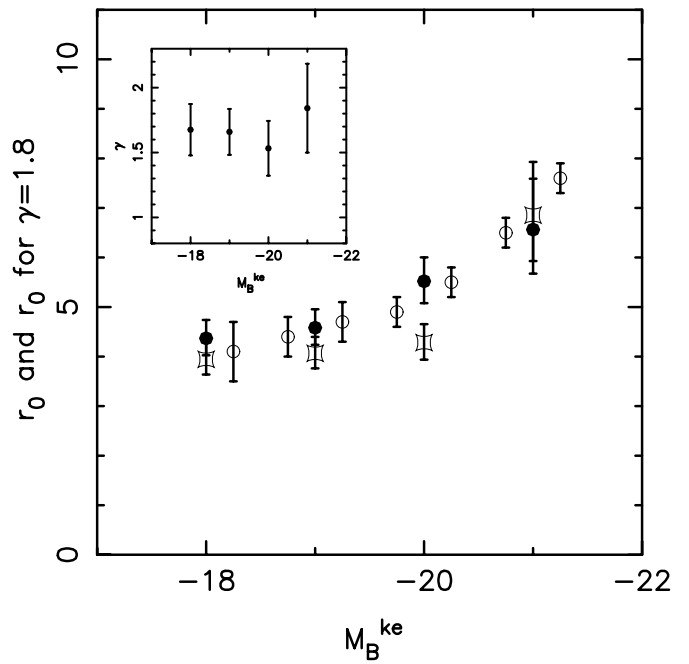


FIG. 3b

FIG. 3.—The fitted  $r_0$  (filled circles) as a function of the  $k$ -corrected and evolution-compensated absolute magnitude ( $R$  band in 3a and  $B$  in 3b) in 1 mag bins around the central color. The fitted slopes,  $\gamma$ , are shown in the inset. The correlation lengths normalized to  $\gamma = 1.8$  are shown as the pincushion symbols. The open circles are the SDSS (3a) and 2dF (3b) measurements at low redshift.

since virtually all high-luminosity galaxies are quite red, leaving too few luminous blue galaxies in the present sample for an accurate clustering measurement.

We can compare our intermediate-redshift results with similar low-redshift results. We show in Figure 3a the clustering-luminosity relation found in the low-redshift SDSS (Zehavi et al. 2001) using  $R$ -band luminosity, where we convert SDSS  $r_*$  to Cousins  $R$  using  $R = r_* - 0.24$  (Fukugita, Shimasaku, & Ichikawa 1995). A similar low-redshift  $R$ -band luminosity dependence of clustering was seen in the LCRS (Lin et al. 1996). In Figure 3b we show our measurements of the dependence of clustering on  $B$ -band luminosity along with the relation found in the 2dF survey (Norberg et al. 2001). The 2dF paper interprets the luminosity dependence as a result of the natural bias mechanism, in which dark matter halo clustering increases with halo mass. However, galaxies of all masses are subject to considerable clustering-dependent evolutionary forces. Here we note that our clustering results at  $z \sim 0.4$  have a luminosity dependence statistically identical with those of the SDSS and 2dF surveys at low redshift, confirming our “no co-moving coordinate clustering evolution” result. However, we also note that  $r_0$  is a function of both galaxy mass, a relative fixed property at late times, and the current star formation rate, which can vary substantially. The advantage of some range of redshift is that we can use it to deepen the study of the relation of clustering, mass, star formation and their evolution.

The theoretical predictions for clustering are clearest for dark halos, which are increasingly clustered with mass (Mo & White 1996; Jing 1998). The theory works quite well to predict the relative clustering of our small galaxy groups relative to the clustering of individual galaxies (Carlberg et al. 2001). However, the luminosity dependence of clustering requires a procedure to associate a luminosity with dark halos, ultimately requiring a fairly detailed parameterization of the galaxy formation history. Within semi-analytic models (Kauffmann et al. 1999; Benson et al. 2000a, 2000b) (Somerville et al. 2001) galaxy clustering is expected to have a dependence on the star formation rate and a weak luminosity dependence, but the details (even the sign of the effect) depend significantly on the adopted cosmology, as well as the merger and gas cooling history. At this stage no published model simultaneously quantitatively predicts the CNOC2 overall correlation evolution and its star formation dependence. Some models find that the bluest galaxies are very strongly clustered at late times, likely as a result of a high merger rate driving star formation, which is in significant contrast to our observational results. Understanding the origin of the strong color- $r_0$  relation is fundamental to understanding the late time evolution of galaxies. Since galaxies in small groups and pairs generate much of the correlation function, this paper empirically investigates this question through the evolution of galaxies in small groups.

The clustering-color (or density-star formation) relation shown in Figure 2 must evolve with redshift. At the current epoch a similar relation holds, but the bluest galaxies are a relatively smaller fraction of the galaxy population as a consequence of ongoing evolution. Clustering is predicted to change slowly with redshift. Clearly, for some fixed total stellar mass, the galaxies with the lowest star formation rates at intermediate redshift must have had higher than average star formation rates at high redshift. If the clus-

tering properties of those galaxies change slowly relative to the rate at which their colors change, then it follows that at some higher redshift the most strongly clustered galaxies are also the ones with the highest star formation rates. The Mo-White formalism predicts that the high-redshift progenitors of the galaxies with  $r_0 \simeq 9 h^{-1}$  Mpc must have correlations at redshifts of two or three that increase over their lower redshift values. Even Lyman break galaxies at  $z \sim 3$  are only clustered with  $r_0 \simeq 4 h^{-1}$  Mpc (Adelberger et al. 1998).

### 3. GLOBAL PROPERTIES OF GROUP AND FIELD GALAXIES

The CNOC2 redshifts have a velocity precision of 70–100  $\text{km s}^{-1}$  in the rest frame of the galaxies. Our group selection procedure (Carlberg et al. 2001) is applied to the  $M_R^{ke} \leq -18.5$  mag sample used in the clustering analysis above. The CNOC2 approach to finding galaxy groups is, in brief, as follows. Groups with three or more members are identified using an algorithm that begins with a friends-of-friends linking in redshift space. These trial groups are refined to a “virialized” set using an iterative procedure designed to ensure that all galaxies are within the estimated radius of virialization for a low-density cosmology,  $1.5r_{200}$ , where  $r_{200}$  is the radius at which the mean interior overdensity is 200 times the critical value. Kinematic group definitions are the most objective approach available, but remain subject to the redshift space interlopers. The resulting groups have centers in position and redshift that do not depend strongly on the group identification details. However, the velocity dispersion of the groups is sensitive to the search distance in the redshift direction and as individuals the groups are subject to large statistical errors, see Carlberg et al. (2001) for a discussion. In spite of these complications the statistical properties of the groups appear to be fairly robust.

The groups used in this paper are identified with a maximum projected distance linking length of  $0.25 h^{-1}$  Mpc and a “standard” redshift distance maximum linking length of  $5 h^{-1}$  Mpc. These parameters, applied within our group identification procedure, yield 195 groups containing 754 galaxies in the 0.1–0.55 redshift range. The sample is complete to about redshift 0.45.

On the average, the groups have 3.8 members with redshifts. Given the small numbers, there are large errors in the derived size and velocity dispersion. Figure 4 shows the number of group members with redshifts and the completeness-corrected numbers versus the estimated velocity dispersion of the groups. The total number of groups with a given number of redshifts is given at the top of the figure. The numbers of group members could be increased at the cost of more complete spectroscopy at our current  $R = 21.5$  mag limit, or, by going deeper. However, the small numbers of members are fundamental to these small groups. Similarly small groups have an unavoidable and significant contamination from redshift space interlopers. We have demonstrated (Carlberg et al. 2001) that as a population these groups have the statistical properties one would expect to find, and they cannot possibly be simply chance groupings of randomly distributed field galaxies.

#### 3.1. Color-Luminosity and Color-Color Relations

Differential evolution could lead to either completely different galaxy populations with types nonexistent in the other group, or, a milder version in which the distribution

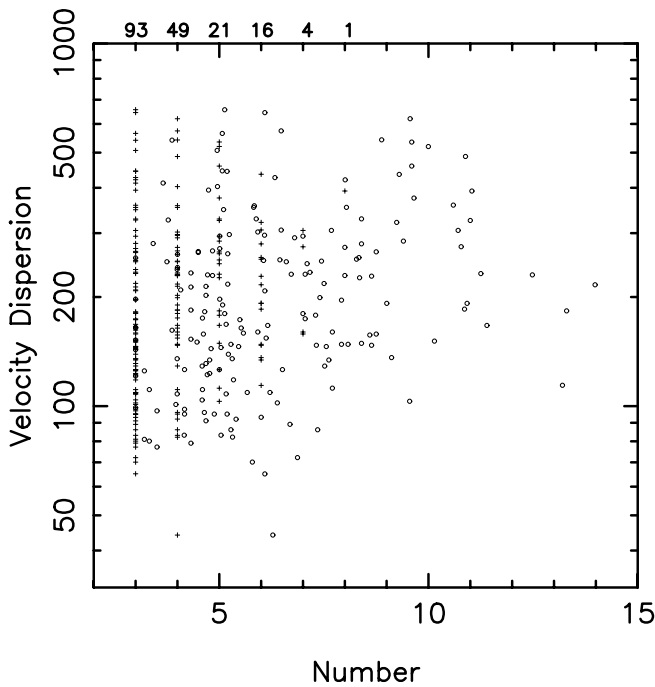


FIG. 4.—Distribution of groups as a function of number. The plus marks indicate the number of spectroscopic redshifts in the identified group. The circles give the completeness-corrected numbers in the group. The groups lie in the redshift range 0.15–0.55.

of galaxies over type is different, but all types are represented in both populations. To test for this the  $k$ -corrected colors,  $(B-R)_0$ , of the galaxies are plotted against the  $k$ -corrected and evolution-compensated absolute  $R$ -band absolute magnitudes for field and group galaxies in Figure 5. Groups have relatively more red galaxies, but, irrespective

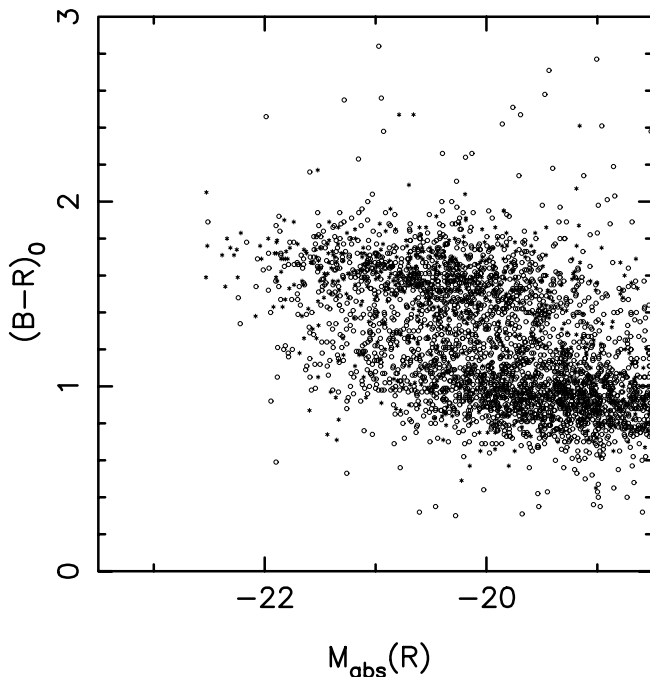


FIG. 5.—Color vs. absolute magnitude for group (asterisks) and field (circles) galaxies. [See the electronic edition of the Journal for a color version of this figure.]

of environment, galaxies in both environments occupy the same total range of colors and luminosities. At low redshift this effect is seen as a morphological type difference, albeit in somewhat higher velocity dispersion groups (Zabludoff & Mulchaey 1998). We classify galaxies into blue and red subsamples using a color of  $(B-R)_0 = 1.25$  mag. Since this color is near the minimum of the number-color relation, it can be shifted 0.05 mag either direction with minimal change in the two subsamples.

The plot of the observed  $U-V$  versus  $V-I$  in the redshift range 0.25–0.45 displayed in Figure 6 further demonstrates that field and group galaxies have essentially the same range of properties present in different proportions. The limited redshift range allows us to compare the colors without  $k$ -corrections. Overplotted are exponential star formation models from the PEGASE.2 spectral synthesis code at solar metallicity and no dust (Fioc & Rocca-Volmerange 1997). Virtually the entire color-color plane is spanned by the exponential models. The bluest points, which are predominantly field galaxies, clearly imply fairly strong star formation over an extended period and require a burst component. The reddest points can be reached from low star formation rate models with a modest amount of dust, or, increasing the metallicity above solar. However, many of these points have fairly large errors in the observed  $U$  flux so may simply be scattered out of the main distribution.

The important conclusion from these color relations is that group and field galaxies cover the same range of properties; however, the relative frequency of blue and red galaxies is quite different in the two environments. One interpretation is that all the galaxies have been drawn from the same parent population and have been subject to similar evolutionary forces, but that these have been applied to different degrees in field and group environments. We

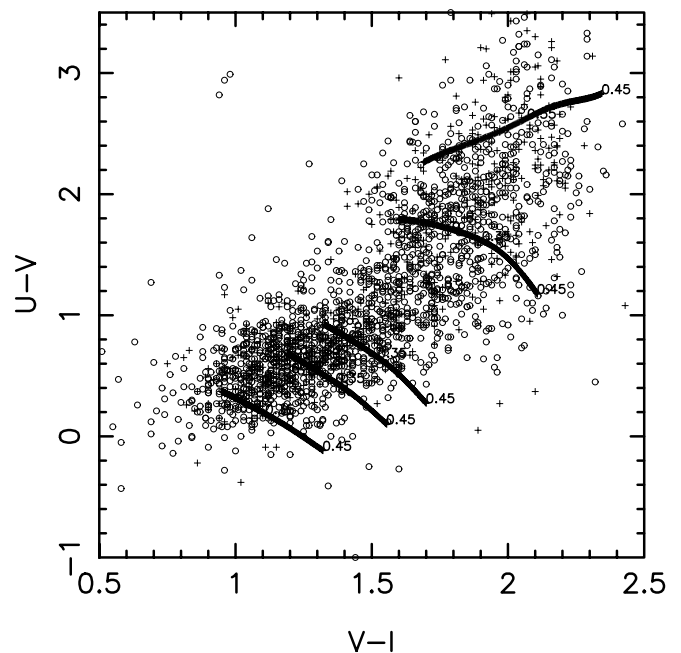


FIG. 6.—Observed  $U-V$  vs.  $V-I$  colors of the group (plus signs) and field (circles) galaxies. The lines are for a series of exponentially declining star formation rate models with decay rates ranging from 1 Gyr (reddest) to 6 Gyr (bluest) with star formation starting at  $t = 0$ . [See the electronic edition of the Journal for a color version of this figure.]

can exclude the possibility that the color-luminosity relations of the two populations require unique evolutionary scenarios, even for a small subset of their population.

### 3.2. Color and Luminosity Distributions

The color distributions, in  $k$ -corrected  $(B-R)_0$ , of field and group galaxies are shown for two approximately equal-volume redshift ranges,  $z = 0.15 - 0.36$  and  $0.36 - 0.47$ , in Figure 7. The sample is restricted to galaxies with  $M_R^k \leq -20$  mag to ensure completeness to redshift 0.55. In this case we have not applied any overall redshift compensation to the luminosities in order that the redshift dependence, if present, be explicitly visible. The figure shows that the group galaxies are on the average redder than the field, for both redshift ranges. The rise in the numbers (selection function weighted) of blue field galaxies with redshift is clearly evident. However, the blue fraction in the groups, those with  $(B-R)_0 < 1.25$  mag, only rises from 28% to 31% across the two redshift bins, which is not significant.

Luminosity functions allow us to sharpen the comparison of group and field galaxy evolution. We first divide the galaxies into red and blue subsamples. A second division of the group galaxies is required since groups with low velocity dispersions,  $\sigma_1 \lesssim 150 \text{ km s}^{-1}$ , are subject to strong dynamical friction which will lead to galaxy merging being important at a level not present in the higher velocity dispersion groups. We therefore split the group galaxies into two roughly equal subsamples at a line-of-sight group velocity dispersion  $\sigma_1 = 150 \text{ km s}^{-1}$ . Even though we do not directly volume normalize the samples, the sample is split into the redshift ranges of  $0.1-0.36$  and  $0.36-0.47$ , which yields two nearly equal volumes of about  $1.5 \times 10^5 h^{-3} \text{ Mpc}^3$ . The average redshifts of the two bins are 0.26 and 0.40. The luminosity functions are calculated as the

luminosity-binned sums of the magnitude selection function weights (Lin et al. 1999), which is adequate for our purposes, although it leads to some uncorrected incompleteness at lower luminosities at higher redshift.

The field and two group luminosity functions are shown as summed weights in Figure 8, where the thick and thin lines are, respectively, for the low- and high-redshift volumes. There is little evolution of the red galaxy luminosity functions, for either the field or group population. We measure a small increase in the mean luminosity with redshift, approximately 0.1 mag, but this difference is not statistically significant. On the other hand the blue galaxies show about  $0.3 \pm 0.1$  mag of increase of the mean luminosity, even across this small redshift range. The low velocity dispersion groups have poor statistics, but there is weak evidence that they have a negative luminosity evolution with increasing redshift. That is, they contain relatively more high-luminosity galaxies at lower redshift.

The strong field galaxy evolution of the blue population and weak evolution of the red population of the left-hand panels of Figure 8 is a restatement of the results presented in our more detailed paper (Lin et al. 1999), where quantitative measures of the evolution are given. Here our primary interest is to compare the field galaxy evolution to the group galaxy evolution. In the high velocity dispersion groups evolution is largely equal to the evolution of the field galaxies. The one possible difference is that the groups appear to have relatively more very high luminosity red galaxies than the field. Although this is not statistically significant in this sample, it is reminiscent of the cD galaxies found in clusters.

The important point established in this section is that both the field and  $\sigma_1 > 150 \text{ km s}^{-1}$  group galaxy populations have very similar evolution of their color-dependent

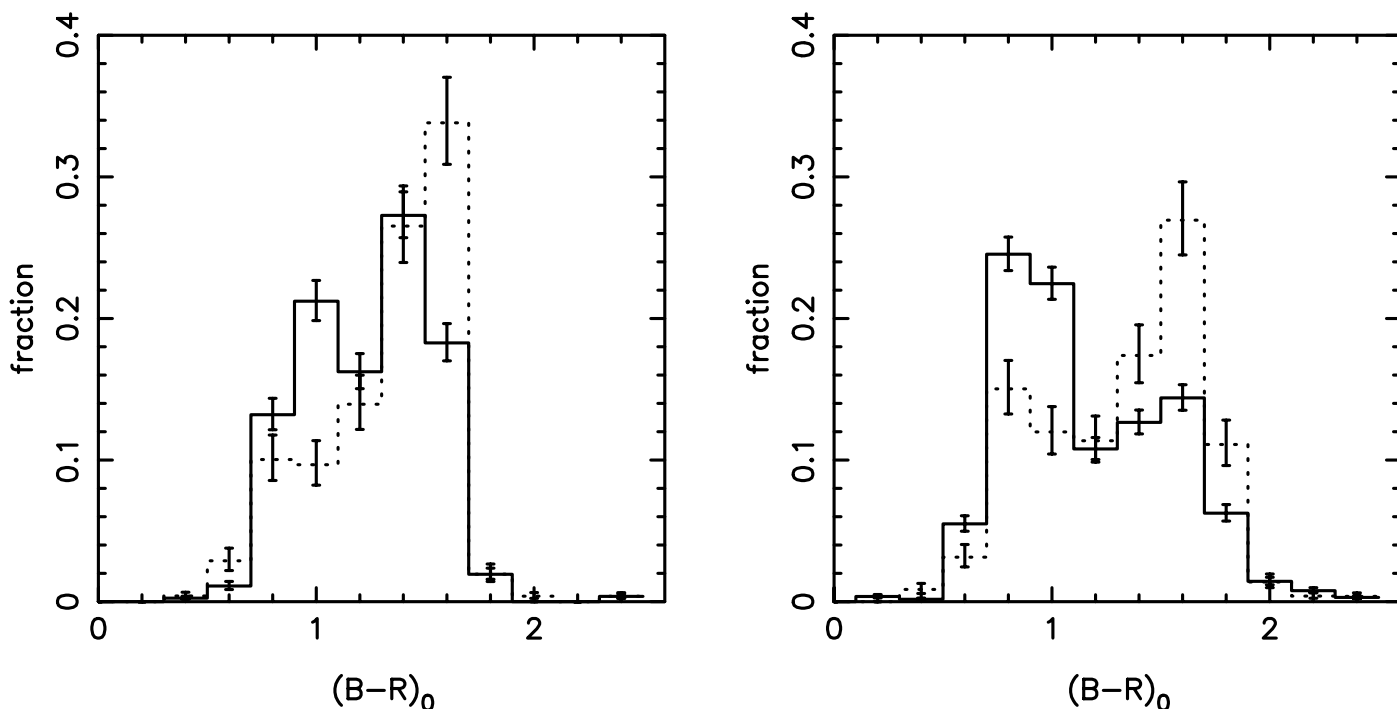


FIG. 7.—Distribution of colors of field galaxies (solid line) and group galaxies (dashed line) that are more luminous than  $M_R^k = -20$  mag. The left-hand panel is for redshifts between 0.15, and 0.36 the right is for redshifts 0.36 to 0.47. The error bars are simply  $\sqrt{N}$  for the bins.

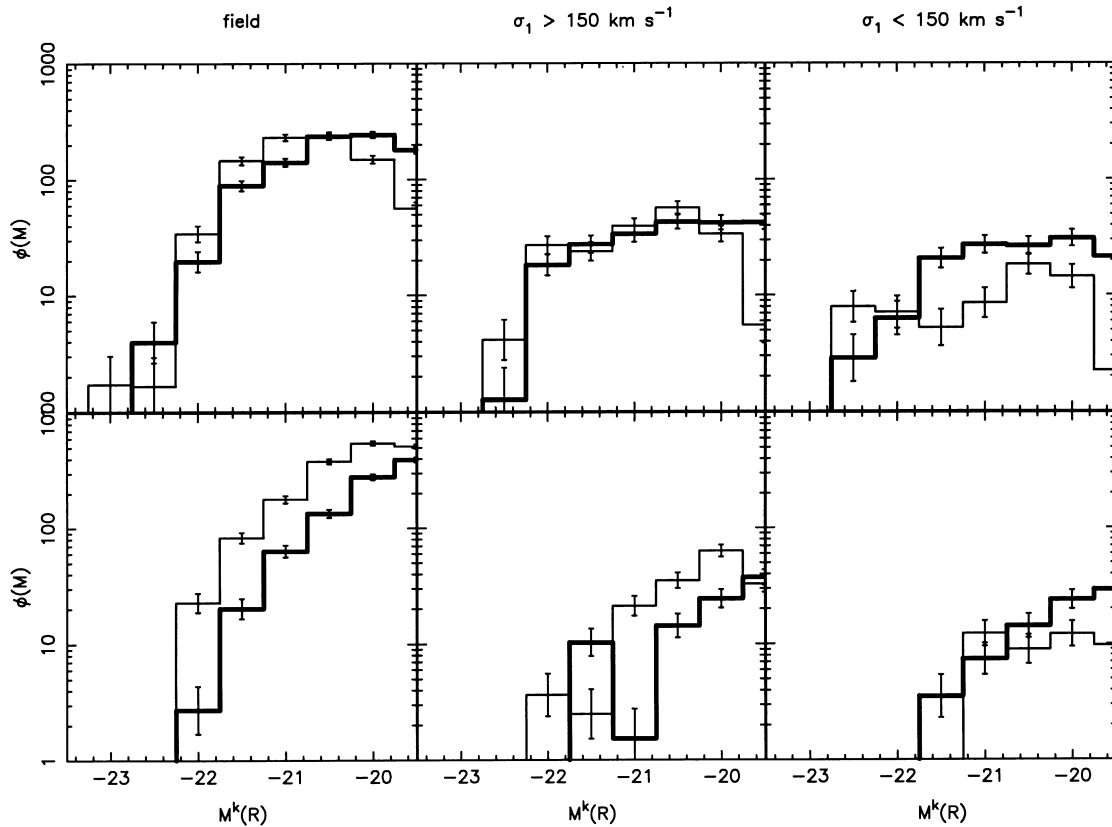


FIG. 8.—The  $k$ -corrected luminosity functions in the  $z = 0.1 - 0.36$  (thick lines) and  $z = 0.36 - 0.47$  (thin lines) redshift ranges. The galaxies are separated into field, high and low velocity dispersion groups. These redshift ranges have nearly equal volumes. Blue galaxy,  $(B - R)_0 < 1.25$ , luminosity functions are shown in the lower panels and red galaxies in the upper panels. The luminosity functions are not normalized to the volume, approximately  $1.1 \times 10^5 h^{-3} \text{ Mpc}^3$ . The errors are simply  $\sqrt{N}$  for the bin. [See the electronic edition of the *Journal* for a color version of this figure.]

luminosity functions with redshift. That is, in these two environments, the red galaxy luminosity function evolves little with redshift, whereas the blue galaxy luminosity function appears to move to higher luminosities with increasing redshift. Groups with velocity dispersion less than  $150 \text{ km s}^{-1}$  have poorer statistics and do not show evidence for any significant luminosity function change with redshift.

#### 4. RADIAL COLOR GRADIENTS WITHIN GROUPS

The radial dependence of the colors allows a test of the various physical forces that may affect evolution. Evolutionary forces that are universal, such as the metagalactic ionizing flux or internal processes of individual galaxies, should have almost no radial dependence across a group. Environmentally dependent evolutionary forces, such as the tidal fields of the group and its member galaxies, or gas stripping due to an intergroup medium, should have a radial dependence within a group. Orbital mixing will dilute the visibility of these effects but statistically they will remain visible, if present.

Figure 9 displays the radial dependence of color in the redshift ranges  $0.10-0.36$  (solid line) and  $0.36-0.47$  (dashed line), the same ranges as used for our luminosity functions. The sample is restricted to  $M_R^{ke} \leq -18.5$  mag. In making these radial plots we make no distinction between group and field galaxies, proceeding as follows. For each group center we search the entire volume-limited sample for galaxies having  $|\Delta v| \leq 3\sigma_1$ . For the galaxies in this sheet we

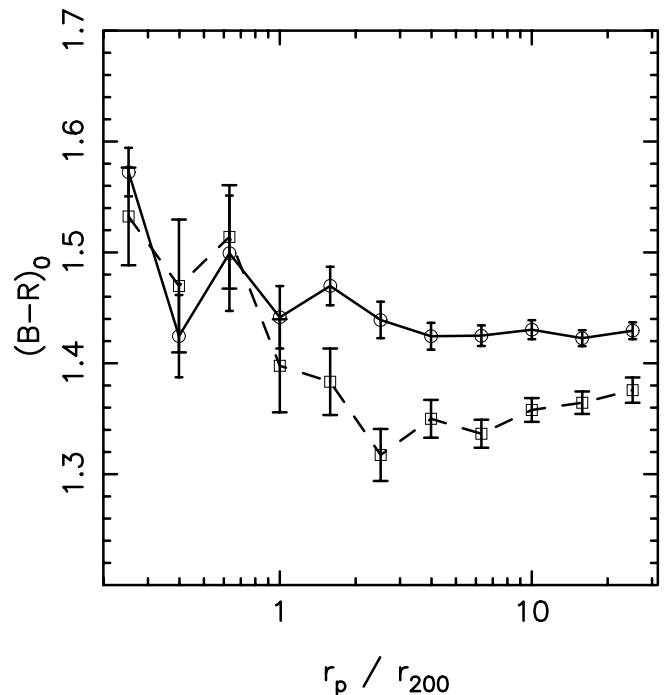


FIG. 9.—Mean colors of galaxies as a function of projected distance from the centers of the groups. The galaxies are selected to have velocities with respect to the group center of  $|\Delta v| \leq 3\sigma_1$ . The solid line is for galaxies with redshifts between 0.1 and 0.36, and the dashed line is for 0.36 to 0.47.



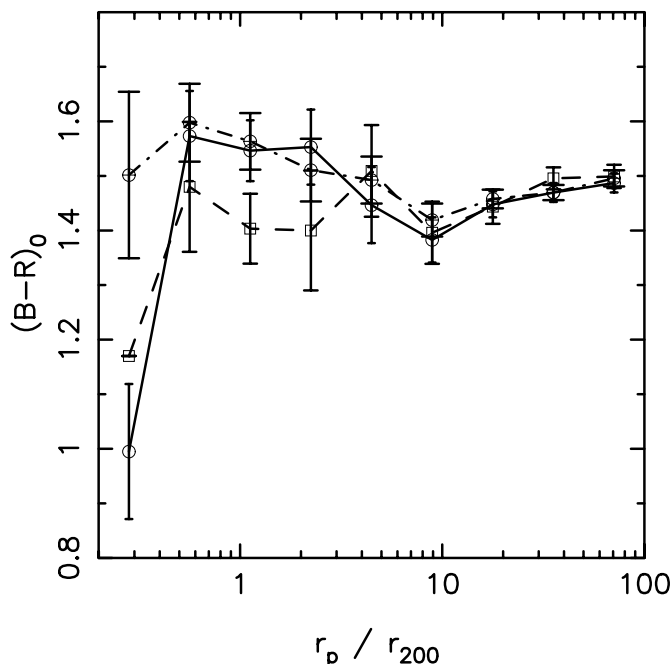


FIG. 10.—Color-radius relation galaxies around groups having  $\sigma_1 < 100, 125,$  and  $150 \text{ km s}^{-1}$  (solid, dashed, and dot-dashed lines, respectively) for high-luminosity galaxies with  $M_R^* \leq -21$ . Lower luminosity galaxies show little effect.

calculate  $r_p$ , the projected separation from the group center. We display  $r_p$  normalized to the nominal virialization radius,  $r_{200} = \sqrt{(3)\sigma_1/[10H(z)]}$ , where  $\sigma_1$  is the line-of-sight velocity dispersion and  $H(z)$  is the Hubble constant at the redshift of the group (Carlberg et al. 2001). Here we restrict the sample to groups with  $\sigma_1 \geq 150 \text{ km s}^{-1}$  to exclude the low velocity dispersion groups that we consider below. The two redshift ranges have different average  $k$ -corrected colors at large radii, a straightforward reflection of the increasing blueness with redshift of the field population. The mean color of the inner  $r \leq 1.3r_{200}$  is  $\langle(B-R)_0\rangle = 1.48 \pm 0.018$  (low- $z$  bin) and  $\langle(B-R)_0\rangle = 1.47 \pm 0.023$  (high- $z$  bin), whereas at larger radii the mean color is  $\langle(B-R)_0\rangle = 1.43 \pm 0.005$  (low- $z$  bin) and  $\langle(B-R)_0\rangle = 1.35 \pm 0.008$  (high- $z$  bin). We conclude the mean color difference between field and group is highly significant. A related effect is seen at low redshifts in the galaxy morphologies, albeit in somewhat richer groups (Tran et al. 2001).

The color gradient found in these groups is not as dramatic as the equivalent measures in rich clusters (Balogh et al. 1999; Ellingson et al. 2001). These groups are not as “tidy” as clusters. At the outer radius of the groups,  $1.5r_{200}$ , the mean local volume density is about  $100\rho_0$ . For a group with a velocity dispersion of  $200 \text{ km s}^{-1}$  the outer radius is about  $0.3 h^{-1} \text{ Mpc}$ . The redshift slice used in the color analysis would extend  $600 \text{ km s}^{-1}$  or about  $5 h^{-1} \text{ Mpc}$  in either direction. Hence the overdensity of the group in that redshift slice is about 6 relative to the unperturbed background. The same is true in clusters, but the total population is so much greater (typically at least 30 members, instead of three) that the overall dilution is much less in clusters. The probability that the galaxy is outside the group but part of the extended group “halo” is about 50% (Carlberg et al. 2001). The halo is a correlated structure that

may well contain a color gradient relative to the field. Hence, the projected color gradient without a background correction is expected to be relatively unaffected by field galaxy interlopers. A second issue is that the velocity dispersions of the groups, and hence the calculated  $r_{200}$  values, have quite large statistical errors. This will lead to some radial scrambling of the summed group that will blur the true color gradient. As an aside we note that the existence of this color gradient is strong support for the physical reality and approximately correct centers (in  $x, y, z$ ) of the majority of these kinematically selected groups.

A striking feature of Figure 9 is that the mean galaxy colors appear to begin to become redder near  $r_{200}$ , the radius of virialization. The current data are insufficient to precisely locate the radius at which the change begins. Within the group the mean colors do not vary much with redshift. In these  $\sigma_1 > 150 \text{ km s}^{-1}$  groups star formation is, on the average, suppressed relative to the field. Figure 9 also shows that the suppression process operates on a timescale short compared to the overall evolutionary timescale, since the group galaxies “quickly” (as measured by group radius) reach the same colors for the two redshift ranges examined. The radius of virialization is where one expects infalling-correlated structure, both associated galaxies and loosely distributed gas, to begin to be stripped away from individual galaxies and to spread through the group. The groups are selected to have mean interior densities that are approximately 200 times the critical density. This means that tidal fields, which scale as mean density alone, will have a similar effect at all velocity dispersions. However, ram pressure stripping, which depends on  $\rho_g \sigma_1^2$ , where  $\rho_g$  is the gas density, will be much less effective in low velocity dispersion systems. The fact that groups with velocity dispersions near  $150 \text{ km s}^{-1}$  have a color gradient that resembles that present in clusters with  $\sigma \geq 800 \text{ km s}^{-1}$ , at least in sign and radial range, suggests that the color change mechanism is not related to velocity dispersion, but to density. Tidal fields are then an acceptable mechanism to differentiate field and group galaxies, whereas ram pressure stripping is not.

The low velocity dispersion groups have an interesting variant on the color gradient. Color gradients of groups with maximum  $\sigma_1 < 100, 125,$  and  $150 \text{ km s}^{-1}$  in the redshift range 0.25–0.45 are shown in Figure 10. Near the centers of low velocity dispersion groups the mean color of the high-luminosity galaxy sample becomes bluer than the field, as shown in Figure 10. Note that the color scale has twice the range of Figure 9. The remarkable result is that the low velocity dispersion groups seem to have centers that are bluer than the surrounding field. The errors are based on the standard deviation on the points in the bins. It should be noted that the inner-most bin for the 27 low velocity dispersion groups bins contain only three galaxies. These measurements suggest that in very low velocity dispersion groups that galaxies can interact (and very likely merge) with an accompanying burst of star formation. The sample is so small that this result has to be viewed with caution.

## 5. GALAXY ORBITS IN GROUPS

Merging is predicted to occur in low velocity dispersion groups on a timescale of a few group crossing times. Theoretical studies have used  $n$ -body simulations to examine quite a large range of idealized groups (Barnes 1985; Mamon 1987; Barnes 1990) (Bode et al. 1993, 1994).

Empirical evidence, such as we have reported here (and reviewed in a low-redshift context in Hickson 1997) confirms the presence of merging in such groups. Here we quantitatively assess, using the parameters of our observed groups, the impact of merging on the overall evolution of the field population in our sample. Our groups are uniformly selected on a kinematic basis from a field galaxy survey, so we know how to relate the group merging to its overall impact on the galaxy population as a whole. We use the measured components of the positions and velocities of the observed galaxies, and then fill in the missing components by randomly sampling three-dimensional distributions in position and velocity which we know from our statistical study of groups (Carlberg et al. 2001).

Rather than full  $n$ -body integrations we follow the orbits of galaxies in a static potential but include the dynamical friction. We undertake the group orbit integrations in a coordinate system where  $r_p$  is the coordinate in the plane of the sky relative to the group center and  $r_z$  is the (physical) coordinate along the line of sight. For each group galaxy we have an observed  $r_p$  and  $v_z$ . To complete the initial conditions we need an  $r_z$  and the two components of velocity that are in the plane of the sky. In Carlberg et al. (2001) we found that group galaxies are distributed with a volume density  $\nu(r)$  that is consistent with the power law  $r^{-2.5}$  and that the velocity distribution is consistent with being drawn from an isotropic Gaussian distribution. Therefore, the two unknown velocity components are set by drawing from a Gaussian velocity distribution having the measured  $\sigma_1$ . The three-dimensional distance to the center,  $r$ , must lie between  $r_p$  and the largest value that keeps it within the virialized region of the group which we take as  $1.5r_{200}$ , as appropriate to a low-density universe. The density distribution  $\nu(r)$  is generated from randomly sampled numbers in  $C(r|r_p) = \int_{r_p}^r \nu(x) dx$ . We therefore find  $r$  from  $C(r|r_p) = \frac{2}{3}(r_p^{-3/2} - r^{-3/2})$ . The line of sight position is calculated from  $r_z = (r^2 - r_p^2)^{1/2}$ , with a randomly assigned sign. Although the

galaxy is initially within the virialized region it is possible that the combination of radius and velocity is such that orbits move outward for the entire duration of the integration.

The gravitational potential of the group is approximated as a static isothermal sphere,  $-2\sigma_1^2/r$ , consistent with the dynamical analysis (Carlberg et al. 2001). This captures the bulk of the overall dynamics of groups but neglects the potential fluctuations of other galaxies orbiting in the group. The mass of each galaxy is taken to be an assigned  $M/L$ , in  $M_\odot/L_\odot R$  band, times the luminosity of the galaxy. We use  $M/L$  values of 10, 30, and 100, where  $M$  is the total mass, including dark halo, of individual galaxies. Here we present the results only for  $M/L = 100$ . Integrations for  $M/L = 30$  have  $\sim 20\%$  reduced merger rates with a factor of about 3 drop in merger rates for  $M/L = 10$ . A vital ingredient in the calculation is to include dynamical friction. We use Chandrasekhar's formula (Binney & Tremaine 1987), adopting a Coulomb logarithm,  $\ln \Lambda = 1$ . The galaxy orbit is then integrated for a set maximum time, usually around 1 Gyr. Two example orbits over a 2 Gyr interval are shown in Figure 11.

The statistics of galaxy mergers in groups are shown in Figure 12. We show velocity dispersion–binned merger fractions for integrations of 0.5, 1, and 5 Gyr. Here we only count as mergers those galaxies that simultaneously have  $r < 20 h^{-1}$  kpc and  $|v| \leq 300$  km s $^{-1}$ , which are approximately the criteria for merging of any galaxy with a galaxy sitting at the potential bottom of a group. The figure shows that the only groups in which mergers will occur with any significant rate are those with  $\sigma \leq 150$  km s $^{-1}$ . The expected merger rate declines with time since the reservoir of groups is exhausted in our simulations, which have no ongoing formation of new groups. Most of the groups that will have a merger will do so in the first 0.5 Gyr or so. The rate of merging over a longer time is dependent on the creation of new groups, which is not rapid over this interval

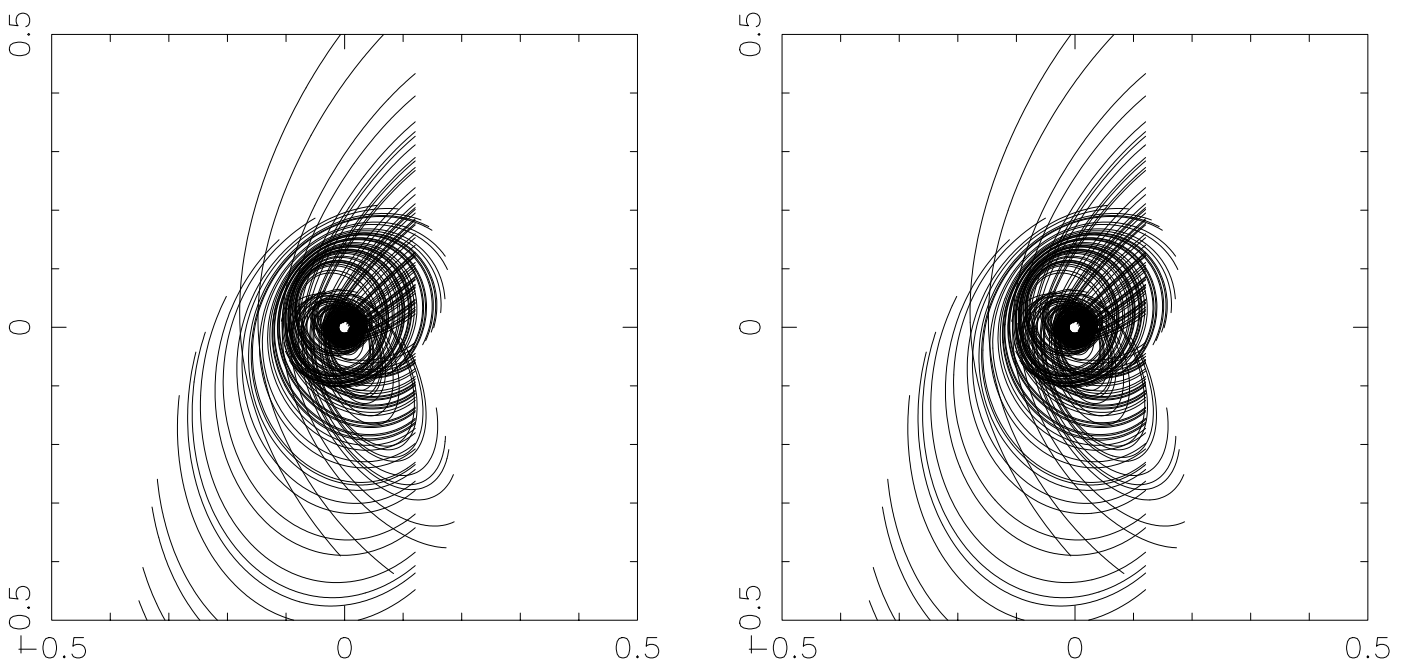


FIG. 11.—Monte Carlo orbits for two of the galaxies in a random group. The box size is given in  $h^{-1}$  Mpc units.

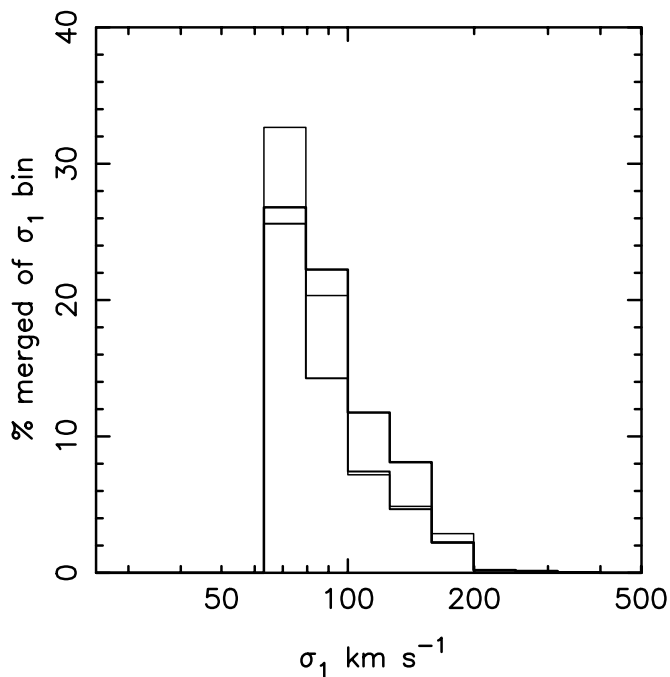


FIG. 12.—Merger rate for the groups as a function of the group velocity dispersion as derived from the Monte Carlo integrations. The lines increase in width as the integration time increases from 0.5, to 1.0, to 5.0 Gyr. The left-hand panel gives the merged fraction in each bin of velocity dispersion, whereas the right-hand panel gives the fraction of the entire sample that merges.

where clustering is barely increasing (Carlberg et al. 2000a). The merger rate in groups, averaged over the entire sample, is about 2.5% per Gyr, with about 25% of all groups below  $150 \text{ km s}^{-1}$  having at least one luminous galaxy merger in 1

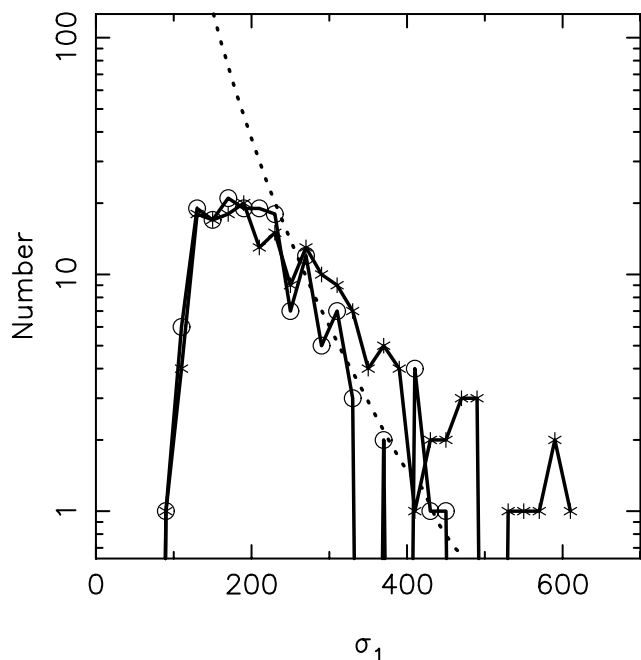


FIG. 13.—Distribution of measured line-of-sight velocity dispersions in  $20 \text{ km s}^{-1}$  bins, plotted at the bin centers. The asterisks are for the standard  $5 h^{-1}$  Mpc groups and the open circles are for the  $3 h^{-1}$  Mpc groups. The dotted line shows the Press-Schechter prediction for the distribution using the cluster normalization. [See the electronic edition of the *Journal* for a color version of this figure.]

Gyr. Although there is no strong astrophysical justification for alternate values of the critical velocity for merging, raising it to  $400 \text{ km s}^{-1}$  increases the merger rate by about 50% and lowering it to  $200 \text{ km s}^{-1}$  reduces the merger rate by about a factor of 4. We conclude that groups below  $150 \text{ km s}^{-1}$  are short lived, although their impact on the population as a whole to redshift one is a 15% or so effect. The higher velocity dispersion groups are not affected by merging at all, so their differences from the field are due to other effects.

Figure 13 shows the distribution of number of groups with a given line-of-sight global velocity dispersions,  $\sigma_1$ , for our standard groups and an alternate group sample defined using a maximum redshift linking length of  $3 h^{-1}$  Mpc. The numbers of high velocity dispersion groups are in good agreement with the rich cluster normalized Press-Schechter (1974) prediction. Below  $200 \text{ km s}^{-1}$  the numbers of groups falls well below the prediction, mainly because they are halos that contain only one or two galaxies above our luminosity limit, rather than the three we require.

## 6. DISCUSSION AND CONCLUSIONS

The theme of this paper is to use the differential evolution of field and group galaxies to examine the dominant processes that drive galaxy evolution at low redshift. We find three empirical relations between clustering environment and the star formation rate, as indicated by  $(B-R)_0$  colors. First, the correlation length-star formation rate relation of Figures 1 and 2 shows that galaxies of low star formation rates usually occur in very strongly clustered regions. Second, the luminosity function evolution of the field population is very similar to that of galaxies in groups having  $\sigma_1 \gtrsim 150 \text{ km s}^{-1}$  (Fig. 8). Third, galaxy groups with  $\sigma_1 \gtrsim 150 \text{ km s}^{-1}$  become redder toward their centers (Fig. 9). There is a statistically weak indication that groups with low velocity dispersions appear to have much bluer than average galaxies at the group center.

It is natural to ask if the changes in group galaxies and the field population as a whole are physically related. Three physical possibilities are outlined below.

1. Field and group galaxies may have a common internal “clock” that causes them to evolve at similar rates. For instance, the rate of cooling of gas and subsequent star formation in the part of the galactic potential well that is occupied by stars may have little environmental dependence. In this approach the population differences between group and field are set in place at some much earlier epoch. This idea does not account for the color gradients in groups and is therefore unattractive.

2. Groups undergo continuing infall of field galaxies, so the evolution of the field is naturally mirrored in groups. In this picture field galaxies are where the bulk of “normal, quiescent” star formation occurs. The origin of the strong evolution of this star formation is not specified, but, once galaxies enter groups their star formation rates tend to decline, as they do in the rich clusters. The combination of infall and subsequent transformation then explains both the common evolution and the population difference between field and group galaxy populations. This idea can account for color gradients, but provides no insight into the mechanism driving the rapid evolution of blue galaxies in the field.

3. In hierarchical cosmologies all galaxies are created from smaller preexisting galaxies. Moreover, many of the

field galaxies in our sample would be found to be in groups if we counted galaxies further down the luminosity function. Hence, most galaxies can be considered to have been group galaxies at some time. In this approach the natural state of weakly clustered galaxies is to be undergoing star formation. However, as galaxies enter a more clustered environment their star formation is reduced. The reduction is controlled by processes that operate on a timescale comparable to the orbital time in the virialized unit, which is initially comparable to the Hubble time, so the rate of reduction is slow. This picture is a “two-stage” evolutionary picture. The first stage occurs in regions of low velocity dispersion where there is ongoing star formation, accompanied by mergers from time to time that can lead to major star bursts. As the velocity dispersion of the group grows, the dynamical friction timescale ascends so that mergers are no longer important. In the second stage the dominant action of the group is to tidally remove small gas-rich companion galaxies and more generally distributed gaseous material at large radii that would normally be available to support ongoing star formation. Given that the pairwise velocity dispersion at low redshift is 300–500 km s<sup>-1</sup> (Davis & Peebles 1983; Marzke et al. 1995; Peacock et al. 2001; Zehavi et al. 2001) this means that relatively few field galaxies are in low velocity dispersion groups where merging is strong. At some yet unidentified higher redshift the pairwise velocity is the range where low velocity dispersion groups are much more prevalent and merging is likely to be the dominant effect of those groups. Below redshift one clustering has grown sufficiently that the groups have increased their velocity dispersion sufficiently that the dominant environmental evolutionary effect of groups is suppression of star formation by tidal effects. Given that a large fraction of all galaxies are found in groups and that clustering increases to lower redshift, the group-dependent star formation suppression will have an increasing affect on the field

population as a whole. Mergers are relatively rare over our redshift range (Patton et al. 1997, 2000; Le Fèvre et al. 2000; Carlberg et al. 2000b); hence clustering-related suppression, rather than merger enhancement, may be the dominant source of low redshift galaxy evolution.

There are two observational tests that should be made of these scenarios. First, a general test is that if clustering is a property that changes slowly with time, then at some redshift greater than one there should be a population of galaxies with  $r_0 \simeq 7 h^{-1}$  Mpc, or more, that are very strongly star-forming. These would be the progenitors of the reddest galaxies at low redshift. Second, beyond redshift one,  $\Omega_M$  begins to rise toward unity, the dark matter clustering declines, the fraction of galaxies in groups with velocity dispersion larger than 150 km s<sup>-1</sup> starts to become small and the pairwise velocity likely declines (for modest clustering biases). This then is the time when merging in groups, possibly with much more gas-rich galaxies than observed below redshift one, leads to a dramatic buildup of the stellar mass of spheroids through merging. Although there are indirect suggestions of this activity, a direct test with redshifts having the kinematic precision to identify groups is required.

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## REFERENCES

- Abadi, M. G., Moore, B., & Bower, R. G. 1999, *MNRAS*, 308, 947  
 Adelberger, K. L., Steidel, C. C., Giavalisco, M., Dickinson, M., Pettini, M., & Kellogg, M. 1998, *ApJ*, 505, 18  
 Allington-Smith, J. R., Ellis, R., Zirbel, E. L., & Oemler, A. J. 1993, *ApJ*, 404, 521  
 Allam, S. S., Tucker, D. L., Lin, H., & Hashimoto, Y. 1999, *ApJ*, 522, L89  
 Balogh, M. L., Morris, S. L., Yee, H. K. C., Carlberg, R. G., & Ellingson, E. 1999, *ApJ*, 527, 54  
 Balogh, M. L., Navarro, J. F., & Morris, S. L. 2000, *ApJ*, 540, 113  
 Barnes, J. 1985, *MNRAS*, 215, 517  
 ———. 1990, *Nature*, 344, 379  
 Benson, A. J., Baugh, C. M., Cole, S., Frenk, C. S., & Lacey, C. G., 2000, *MNRAS*, 316, 107  
 Benson, A. J., Cole, S., Frenk, C. S., Baugh, C. M., & Lacey, C. G., 2000, *MNRAS*, 311, 793  
 Binney, J., & Tremaine, S. 1987, *Galactic Dynamics* (Princeton: Princeton Univ. Press)  
 Bode, P. W., Berrington, R. C., Cohn, H. N., & Lugger, P. M. 1994, *ApJ*, 433, 479  
 Bode, P. W., Cohn, H. N., & Lugger, P. M. 1993, *ApJ*, 416, 17  
 Broadhurst, T. J., Ellis, R. S., & Shanks, T. 1988, *MNRAS*, 235, 827  
 Burbidge, E. M., & Burbidge, G. R. 1961, *AJ*, 66, 541  
 Carlberg, R. G., Yee, H. K. C., Morris, S. L., Lin, H. Hall, P. B., Patton, D., Sawicki, M., & Shepherd, C. W. 2000, *ApJ*, 542, 57  
 ———. 2001, *ApJ*, 552, 427  
 ———. 2000, *ApJ*, 532, L1  
 Coleman, G. D., Wu, C., & Weedman, D. W. 1980, *ApJS*, 43, 393  
 Davis, M., & Geller, M. J. 1976, *ApJ*, 208, 13  
 Davis, M., & Peebles, P. J. E. 1983, *ApJ*, 267, 465  
 de Mello, D. F., Infante, L., & Menanteau, F. 1997, *ApJS*, 108, 99  
 Diaferio, A., Geller, M. J., & Ramella, M. 1994, *AJ*, 107, 868  
 Dressler, A. 1980, *ApJ*, 236, 351  
 Dubinski, J. 1998, *ApJ*, 502, 141  
 Ellingson, E., Lin, H., Yee, H. K. C., & Carlberg, R. G. 2001, *ApJ*, 547, 609  
 Fioc, M., & Rocca-Volmerange, B. 1997, *A&A*, 326, 950  
 Fisher, K. B., Davis, M., Strauss, M. A., Yahil, A., & Huchra, J. 1994, *MNRAS*, 266, 50  
 Fukugita, M., Shimasaku, K., & Ichikawa, T. 1995, *PASP*, 107, 945  
 Guzzo, L., Strauss, M. A., Fisher, K. B., Giovanelli, R., & Haynes, M. P. 1997, *ApJ*, 489, 37  
 Hashimoto, Y., & Oemler, A. J. 2000, *ApJ*, 530, 652  
 Hickson, P. 1982, *ApJ*, 255, 382  
 ———. 1997, *ARA&A*, 35, 357  
 Huchra, J. P., & Geller, M. J. 1982, *ApJ*, 257, 423  
 Jing, Y. P. 1998, *ApJ*, 503, L9  
 Kauffmann, G., Colberg, J. M., Diaferio, A., & White, S. D. M. 1999, *MNRAS*, 307, 529  
 Le Fèvre, O., et al. 2000, *MNRAS*, 311, 565  
 Lilly, S. J., Le Fevre, O., Hammer, F., & Crampton, D. 1996, *ApJ*, 460, L1  
 Lilly, S. J., Tresse, L., Hammer, F., Crampton, D., & Le Fevre, O. 1995, *ApJ*, 455, 108  
 Lin, H., Kirshner, R. P., Shectman, S. A., Landy, S. D., Oemler, A., Tucker, D. L., & Schechter, P. L. 1996, *ApJ*, 471, 617  
 Lin, H., Yee, H. K. C., Carlberg, R. G., Morris, S. L., Sawicki, M., Patton, D., Wirth, G., & Shepherd, C. W. 1999, *ApJ*, 518, 533  
 Loveday, J., Tresse, L., & Maddox, S. 1999, *MNRAS*, 310, 281  
 Mahdavi, A., & Geller, M. J. 2001, *ApJ*, 554, L129  
 Mamon, G. A. 1987, *ApJ*, 321, 622  
 Marzke, R. O., Geller, M. J., da Costa, L. N., & Huchra, J. P. 1995, *AJ*, 110, 477  
 Merritt, D. 1984, *ApJ*, 276, 26  
 ———. 1985, *ApJ*, 289, 18  
 Mo, H. J., & White, S. D. M. 1996, *MNRAS*, 282, 347  
 Moore, B., Lake, G., & Katz, N. 1998, *ApJ*, 495, 139  
 Mulchaey, J. S. 2000, *ARA&A*, 38, 289  
 Mulchaey, J. S., & Zabludoff, A. I. 1998, *ApJ*, 496, 73  
 ———. 1999, *ApJ*, 514, 133  
 Nolthenius, R., & White, S. D. M. 1987, *MNRAS*, 225, 505  
 Norberg, P., et al. (the 2dF Collaboration) 2001, *MNRAS*, submitted (astro-ph/0105500)

- Patton, D. R., Carlberg, R. G., Marzke, R. O., Pritchett, C. J., da Costa, L. N., & Pellegrini, P. S. 2000, *ApJ*, 536, 153
- Patton, D. R., Pritchett, C. J., Yee, H. K. C., Ellingson, E., & Carlberg, R. G. 1997, *ApJ*, 475, 29
- Peacock, J. A., et al. 2001, *Nature*, 410, 169
- Poggianti, B. M., Smail, I., Dressler, A., Couch, W. J., Barger, A. J., Butcher, H., Ellis, R. S., & Oemler, A. 1999, *ApJ*, 518, 576
- Press, W. H., & Schechter, P. 1974, *ApJ*, 187, 425
- Ramella, M., Geller, M. J., & Huchra, J. P. 1989, *ApJ*, 344, 57
- Saunders, W., Rowan-Robinson, M., & Lawrence, A. 1992, *MNRAS*, 258, 134
- Shepherd, C. W., Carlberg, R. G., Yee, H. K. C., Morris, S. L., Lin, Sawicki, M. H., Hall, P. B., & Patton, D. *ApJ*, 560, 72
- Somerville, R. S., Lemson, G., Sigad, Y., Dekel, A., Kauffmann, G., & White, S. D. M. 2001, *MNRAS*, 320, 289
- Tran, K.-V. H., Simard, L., & Zabludoff, A. I. 2001, *ApJ*, 549, 172
- Whitmore, B. C., Gilmore, D. M., & Jones, C. 1993, *ApJ*, 407, 489
- Willmer, C. N. A., da Costa, L. N., & Pellegrini, P. S. 1998, *AJ*, 115, 869
- Yee, H. K. C., et al. 2000, *ApJS*, 129, 475
- Zabludoff, A. I., & Mulchaey, J. S. 1998, *ApJ*, 496, 39
- Zabludoff, A. I., Zaritsky, D., Huan, L., Tucker, D., Hashimoto, Y., Schechtman, S. A., Oemler, A., & Kirshner, R. P. 1996, *ApJ*, 466, 104
- Zehavi, I., et al. (the SDSS Collaboration) 2001, *ApJ*, submitted (astro-ph/0106476)