# Drift-dependent changes in iceberg size-frequency distributions

James D. Kirkham<sup>1,2\*</sup> Nick J. Rosser<sup>1</sup> John Wainwright<sup>1</sup> Emma C. Vann Jones (née Norman)<sup>1</sup> Stuart A. Dunning<sup>3</sup> Victoria S. Lane<sup>4</sup> David E. Hawthorn<sup>5</sup> Mateusz C. Strzelecki<sup>6</sup> Witold Szczuciński<sup>7</sup>

<sup>1</sup>Geography Department and Institute of Hazard Risk and Resilience, Durham University, Durham, DH1 3LE, UK. <sup>2</sup>Scott Polar Research Institute, University of Cambridge, Cambridge, CB2 1ER, UK. <sup>3</sup>School of Geography, Politics and Sociology, Newcastle University, Newcastle, NE1 7RU. <sup>4</sup>SEIS-UK, Department of Geology, University of Leicester, Leicester, LE1 7RH, UK. <sup>5</sup>British Geological Society, The Lyell Centre, Edinburgh, EH14 4AP, UK. <sup>6</sup>Institute of Geography and Regional Development, University of Wrocław, 50-137 Wrocław, Poland. <sup>7</sup>Institute of Geology, Adam Mickiewicz University in Poznań, 61-680 Poznań, Poland.

<sup>\*</sup>email: jk675@cam.ac.uk

# 1 Abstract

2	Although the size-frequency distributions of icebergs can provide insight into how they
3	disintegrate, our understanding of this process is incomplete. Fundamentally, there is a discrepancy
4	between iceberg power-law size-frequency distributions observed at glacial calving fronts and
5	lognormal size-frequency distributions observed globally within open waters that remains
6	unexplained. Here we use passive seismic monitoring to examine mechanisms of iceberg
7	disintegration as a function of drift. Our results indicate that the shift in the size-frequency
8	distribution of iceberg sizes observed is a product of fracture-driven iceberg disintegration and
9	dimensional reductions through melting. We suggest that changes in the characteristic
10	size-frequency scaling of icebergs can be explained by the emergence of a dominant set of driving
11	processes of iceberg degradation towards the open ocean. Consequently, the size-frequency
12	distribution required to model iceberg distributions accurately must vary according to distance from
13	the calving front.
14	
15	
15	
16	
17	
18	
19	
20	
21	
22	
23	

#### 24 Introduction

25 The rate at which icebergs drift and disintegrate influences the risk of collisions with high-latitude hydrocarbon infrastructure and shipping<sup>1</sup>, the extent of zones of nutrient-enhanced 26 27 carbon sequestration<sup>2,3</sup>, and the interpretation of palaeoclimate indicators such as ice-rafted debris<sup>4</sup>. Although iceberg drift-decay models exist<sup>5</sup>, our mechanical understanding of iceberg disintegration 28 remains unable to explain the size-frequency distributions of icebergs commonly observed; most 29 30 notably the discrepancy between the power-law distributed icebergs sizes observed at glacial calving fronts<sup>6</sup> and the lognormal iceberg-size distributions observed globally within open waters<sup>7,8</sup>. 31 Although it has been speculated that the lognormal distribution of iceberg sizes observed away from 32 33 glacial calving fronts is the product of the mechanisms by which icebergs fracture and disintegrate<sup>7</sup>, the absence of appropriate methods with which to study free-floating iceberg disintegrations has 34 35 limited efforts to study the mechanics of this phenomenon.

Over the last four decades<sup>9</sup>, passive seismic investigations of glaciological phenomena have 36 revealed that different glaciological processes are characterised by unique and highly distinctive 37 38 signal properties including dominant spectral frequency, event duration and the shape of the signal onset and coda<sup>10</sup> (Table 1). The application of passive seismic techniques has significantly 39 40 increased our understanding of inaccessible glaciological processes including crevasse propagation<sup>9,11,12</sup>, basal sliding<sup>13,14</sup> and iceberg calving from tidewater glaciers<sup>15,16</sup>. Seismic 41 42 methods have also been used to describe flexure and breakage of free-floating tabular icebergs<sup>14,17</sup>, 43 demonstrating their potential to provide insight into the mechanisms responsible for iceberg 44 disintegration.

The Greenland Ice Sheet has experienced persistent and increasing mass loss since the 1990s<sup>[18]</sup> in a spatially complex pattern driven by rising surface air temperatures<sup>19</sup> and accelerations in outlet glacier velocities<sup>20,21</sup>. During this time, freshwater fluxes into the North Atlantic Ocean sourced from surface and submarine melting of the Greenland Ice Sheet, as well as the melting of

icebergs and ice mélange, have been observed to increase<sup>22,23</sup>. In addition to their implications for 49 circulation dynamics within the global ocean<sup>24</sup> and mass-loss feedbacks within the fjords of 50 marine-terminating outlet glaciers<sup>25</sup>, elevated meltwater fluxes are likely to increase the input of 51 bioavailable particulate iron into the North Atlantic Ocean<sup>3</sup>, potentially affecting marine biological 52 productivity, ecosystem dynamics and the oceanic uptake of CO<sub>2</sub><sup>[2]</sup>. Meltwater fluxes sourced from 53 54 the melting of icebergs and ice mélange within Greenlandic glacial fjords including that of Greenland's large outlet glacier, Jakobshavn Isbræ<sup>26</sup>, may potentially exceed the flux associated 55 with glacier surface and submarine melting<sup>22</sup>. The drift and decay of icebergs during transit from 56 57 the calving terminus therefore represents an important mechanism by which nutrients and 58 freshwater are transported into the North Atlantic Ocean; however, due to a poor understanding of 59 iceberg-disintegration mechanics, these processes are relatively poorly quantified around Greenland at present<sup>3</sup>. 60

61 If an accurate understanding of the mechanisms of iceberg breakup could be obtained, numerical models could be used to predict the expected distribution of iceberg sizes resulting from 62 63 the disintegration process, their trajectories, and their longevity; information which informs risk to 64 shipping and delineates the areas influenced by the delivery of ice-rafted debris and nutrients. One 65 such modelling approach is the use of probabilistic magnitude-frequency scaling laws, which 66 provide a means to quantify the likelihood that an event of a given magnitude will occur over time 67 or, in the context of icebergs, that an iceberg of known dimensions will be produced as a result of 68 the disintegration process. This approach has been widely applied in attempts to forecast the occurrence of natural hazards such as rockfalls<sup>27</sup> and landslides<sup>28</sup>. We therefore apply a probabilistic 69 scaling approach to capture, characterise and model the manner in which icebergs calved from 70 71 Jakobshavn Isbræ disintegrate as they drift through the Vaigat Strait towards Baffin Bay, West 72 Greenland, determined through passive seismic monitoring (Fig. 1). Based on a lognormal 73 distribution of energy released by iceberg cracking and calving, we conclude that the lognormality 74 associated with free-floating iceberg size-frequency distributions is a product of the process of

75 iceberg disintegration and dimensional reductions through melting after their initial calving. We

76 propose that the emergence of a dominant set of iceberg-degradation processes over space

transforms the characteristic distribution of iceberg dimensions from a power-law at glacial calving

fronts to a lognormal distribution as icebergs drift towards the open ocean.

79 **Results** 

#### 80 Description and interpretation of seismic events

81 Seismic signals generated by the processes of iceberg decay were recorded over a 49-day 82 period using a network of six seismometers installed in coastal locations along a 50 km stretch of 83 the Vaigat Strait. Based on their distinctive characteristic spectral frequencies, event durations and 84 signal onset and coda geometries, the observed iceberg-related seismic events may be classified into 85 three groupings, implying that three predominant processes are responsible for generating seismicity in the Vaigat Strait. The waveform geometry, typical duration and characteristic 86 87 frequencies of the three classes of signals compare favourably to previously examined glaciological 88 processes (Table 1), suggesting that the seismic signatures of iceberg decay observed within the 89 Vaigat Strait relate to cracking, microfracturing and iceberg-calving processes (Table 2) 90 (Supplementary Figs. 1 and 2).

91 The first class of icequake, Type 1, exhibits an impulsive, short duration (~1 s) waveform 92 with a characteristic frequency of 30-40 Hz. The brittle nature of ice means that mechanisms of ice 93 deformation are dominated by fracturing, resulting in micro-cracking, coalescence of fractures and 94 fragmentation as the ice exceeds a critical threshold of viscoelastic strain<sup>30</sup>. Ice crevassing and surface fracture is typically associated with short duration (~0.1–2.5 s), 10–30 Hz seismic tremors 95 with highly impulsive onsets<sup>9,29</sup>. These properties are consistent with the characteristics of Type 2 96 97 icequakes, implying that this signal type likely corresponds to tensile fracturing and the enlargement 98 of pre-existing cracks and crevasses. Depending on the mode of failure, fractures may open through tension-based or shearing-dominated mechanisms<sup>31</sup>. These two mechanisms may be differentiated 99

by inspecting the polarity of the first motion of the seismic signal — with consistent first motion polarity across all sensors indicative of tensile failure and mixed polarity signifying that the source has some shearing component<sup>32,33</sup>. Although it is often difficult to distinguish the onset of a signal from the pre-event noise, Type 2 signals generally exhibit a consistent polarity of first tremor motion, supporting the interpretation of this signal as originating from the tensile failure of ice.

105 The similar waveform geometry, duration and first motion polarity of Type 1 and Type 2 106 icequakes suggests that these signals share a similar genesis. The characteristic frequency of seismic 107 waves resulting from brittle material failure scales in accordance with the size of the fracture and 108 the shear modulus of the medium<sup>10,12</sup>. The basic response frequency for fractures in ice, f [Hz], has 109 been shown to respond to changes in crack length<sup>34,35</sup>, L [m]:

$$f = \frac{V}{2L} \tag{1}$$

111 Where *V* is the typical crack propagation velocity for ice  $[m s^{-1}]$ . Laboratory and large-scale 112 geophysical experiments have demonstrated that the mean velocity of a simple crack in ice is 113 approximately 50 m s<sup>-1</sup> and thus, assuming a constant shear modulus, smaller length cracks will 114 result in a higher characteristic frequency relative to larger crack lengths<sup>35,36</sup>. On the basis of this 115 relationship, Type 2 signals correspond to 0.8–2.5 m crack lengths whereas Type 1 signals relate to 116 smaller microfractures with lengths less than 0.8 m.

Type 1 signals are frequently detected prior to and after the onset of Type 2 signals, suggesting that these types of events may be mechanically linked. Cracking and micro-fracturing are progressive processes in which fractures radiate outwards from the tips of cracks following the exceedance of interatomic bonding forces by local tensile stresses<sup>35</sup>. Consequently, the co-occurrence of these two signals likely reflects micro-fracture nucleation at the tips of an enlarging crack, instigated by the volumetric enlargement of the pre-existing rupture. Hence, signal types 1 and 2 appear to be part of a genetically related continuum of tensile fracture processes.

124 The emergent onset, gradually declining coda and predominantly monochromatic 1–5 Hz 125 spectral frequency of the Type 3 signals is consistent with the characteristics of iceberg calving from glacial termini<sup>7,37</sup>. The low frequency (1-5 Hz) spectral peak associated with this type of event 126 127 has been attributed in previous studies to iceberg-water interactions through both displacement of water following iceberg collisions with the water surface<sup>38</sup> and the tilting and rolling of unstable 128 icebergs following detachment from the calving terminus<sup>39</sup>. Whilst the dominant spectral frequency 129 130 of Type 3 events corresponds to the 1–5 Hz frequency band, this type of signal commonly contains 131 a number of short-lived 20–40 Hz peaks, similar to the Type 1 and 2 icequakes, prior to and within 132 the main body of the signal. Type 3 events may therefore consist of the tensile expansion of cracks 133 and microfractures up to a critical threshold where failure of the iceberg occurs through calving. The incidence of Type 1 and 2 events within the dominant 1–5 Hz frequency envelope possibly 134 135 reflects continued cracking and damage accumulation produced by tensile stresses as the iceberg 136 rolls to reach a new buoyant equilibrium following the loss of an ice block through calving 137 (Supplementary Fig. 2). Thus, crack nucleation and expansion appear to progressively weaken 138 icebergs in transit through the Vaigat Strait until mechanical stresses exceed the strength of the ice, 139 culminating in a calving event.

## 140 Magnitude-frequency scaling

141 Testing of different frequency distribution functions (outlined in Methods and 142 Supplementary Methods 1) demonstrates that lognormal distributions provide the most robust 143 analogue for the various iceberg disintegration processes. The spectrum of energies released by cracking and microfracturing are lognormally distributed over six orders of magnitude (Fig. 2a–b), 144 145 with alternative power-law fits only providing a robust approximation of the data between signal energies of 5 x  $10^{-10}$ J to 1 x  $10^{-8}$  J and 3 x  $10^{-10}$  J to 2 x  $10^{-9}$  J, respectively. Power-laws thus 146 147 overestimate the likelihood of occurrence for the smallest and highest magnitude cracking and microfracturing events — a pattern that is also observed for all events combined (Fig. 2d). A 148 149 power-law approximation of the energy released by iceberg calving and rolling provides a robust fit

to events with energies >8 x  $10^{-11}$  J (n = 440) (Fig. 2c), but fails to predict the rollover of lower magnitude energies where the majority of the data (n = 961) falls. Despite overpredicting the likelihood of energies between 4 x  $10^{-11}$  J and 1 x  $10^{-9}$  J and underpredicting the incidence probability of the nine largest iceberg calving events, a lognormal distribution provides a better approximation of the data than the fitted power-law.

#### 155 Event timing

156 Event timing gives insight into the drivers of iceberg disintegration. Correlations between 157 the number of icequakes detected and the height of the semi-diurnal tidal range are moderate to 158 weak (r < 0.5) and vary considerably between the different processes detected (Supplementary 159 Fig. 3a–c). However, greater numbers of events are observed across all six seismometers during 160 periods coincident with the daily tidal range maxima, implying spatially consistent forcing driven by tides. The timing of iceberg calving, rolling and microfracturing exhibits comparable phasing to 161 162 the lunar fortnightly (M<sub>f</sub>, 13.70 day) constituent tide, causing greater numbers of icebergs to calve, 163 roll and fracture when transported into shallower coastal waters during periods of higher tidal 164 amplitude, enabling keel grounding upon the seabed. These processes also exhibit a significant 165 2-4 day periodicity (Supplementary Fig. 3d) that may reflect progressive cycles of damage 166 accumulation due to tidal grounding, culminating in iceberg disintegration after 2 to 4 days of repeated tensile loading, or amplified wave-notching as a result of increased ocean turbulence 167 during the passage of transient storms<sup>15</sup>. Wave-driven turbulence disturbs the build-up of a static 168 169 cold-water layer around icebergs that would diminish melt rates<sup>40</sup>, driving higher rates of heat 170 transfer into the ice, increasing notch cutting and iceberg instability<sup>1</sup>. Similar periodic behaviour is 171 not present for cracking-induced signals, which increase in prevalence throughout the study period 172 (Supplementary Fig. 3a). This pattern likely relates to the cumulative expansion of cracks as a result of progressive microfracture nucleation and growth in response to storms through the summer 173 174 season and iceberg grounding during periods of high tidal amplitude, permitting fractures to 175 coalesce to produce the lower frequency signals associated with large-scale cracking.

#### 176 **Discussion**

Analysis of satellite imagery demonstrates that the distribution of planform iceberg areas in 177 178 Vaigat is well fitted by a lognormal distribution except for the likelihood of the very largest 179 icebergs, which are slightly overpredicted (Fig. 3a). The planform areas of icebergs situated in the 180 zone proximal to the outlet of Jakobshavn Isbræ are power-law distributed over two orders of 181 magnitude, with minor deviations from the fitted distribution occurring for icebergs with a 182 planimetric area of less than 10,000 m<sup>2</sup> (Fig. 3b). This result concurs with satellite-based analysis of the size-frequency distributions of icebergs located in other Greenlandic fjords situated within 183 200 km north of Jakobshavn Isbræ<sup>41</sup>, and within the Ilulissat Isfjord proximal to the calving 184 terminus of Jakobshavn Isbræ<sup>22</sup>, suggesting that the size-frequency distributions of icebergs calved 185 186 from Greenlandic outlet glaciers likely conform to power-law scaling. A power-law distribution of 187 iceberg areas is consistent with observed and theoretical fragment-size distributions calved from tidewater glaciers and ice sheets<sup>6</sup>. The size-frequency distributions of iceberg sizes calved from 188 189 glacial termini measured by seismic and imagery-based methods of monitoring are consistent 190 regardless of the method used to conduct the monitoring (see supplementary methods in ref. 6), 191 indicating that: (i) seismic and satellite-derived measures of iceberg size-frequency distributions are 192 compatible measures of iceberg size, and therefore that (ii) the process of iceberg calving operating 193 at these glacial termini may be different from the lognormal distribution of energies associated with 194 iceberg decay in the open waters of the Vaigat Strait.

195 The incidence of power-law scaling is indicative of a scale-invariant self-organised system 196 fluctuating between regimes of sub-critical damage accumulation and super-critical instability 197 collapse<sup>6</sup>. Systems exhibiting self-organised criticality evolve towards a critical state through the 198 interaction of multiple simultaneous processes<sup>42</sup>. When the critical state is attained, accumulated 199 instabilities may dynamically relax through scale-invariant avalanching, which in the context of a 200 calving terminus may range from minor ice falls to the collapse of the entire calving front<sup>6</sup>. The 201 progression to calving front instability is achieved through numerous mechanisms including surface

ablation, longitudinal stretching, crevasse formation and submarine melt<sup>43</sup>. The connection of the
structural damage generated by the various processes of instability propagation culminates in the
mechanical failure of portions of the glacier terminus, producing icebergs at the calving margin<sup>44</sup>.
Under these conditions, as no single mechanism of damage accretion dominates the iceberg calving
process, the size-frequency distribution of the icebergs produced will not reflect a single formative
process, resulting in the production of a scale-invariant power-law distribution of calved iceberg
sizes<sup>6</sup>.

209 In contrast to the distributions generated at calving termini, power-law approximations of 210 iceberg sizes observed across the North Atlantic Ocean over-predict both the smallest and largest iceberg dimensions<sup>45</sup>. The lognormal distribution of iceberg areas observed within Vaigat concurs 211 with observations from the Arctic<sup>7</sup> and Antarctic<sup>8</sup>, wherein distal icebergs obey lognormal scaling 212 with minimal year-to-year variability<sup>5</sup>. The production of lognormal distributions has been 213 theoretically<sup>46</sup> and experimentally<sup>47</sup> associated with multiplicative breakage and repeated fracturing. 214 The theory of breakage represents an inverse application of the law of proportionate effect<sup>48</sup> in 215 216 which the value of a variable undergoing change corresponds to a random proportion of its previous value<sup>49</sup>. For a system governed by this law, assuming each transformation induced by breakage is 217 218 small, application of the central limit theorem demonstrates that the logarithm of the variable undergoing change will be asymptotically normally distributed<sup>50</sup>, with the breadth of the lognormal 219 220 distribution reflecting the number of independent transformations that are responsible for its formation<sup>48,51</sup>. The incidence of lognormal scaling in the distribution of energies released by iceberg 221 fracturing suggests that the dominant mechanisms by which icebergs decay can be approximated as 222 223 a process operating under the law of proportionate effect.

A lognormal distribution of iceberg sizes produced by fracturing processes is likely reinforced by the dimensional reduction of iceberg dimensions through melting. Although this process in itself has no detectable seismic signature, smaller icebergs generated through the

fracturing process will exhibit greater surface area to volume ratios, making them more susceptible to mass loss through melting. The preferential removal of smaller icebergs from the total population through melting is conducive to the production of a rollover tail in the observed iceberg size-frequency distribution. The characteristic size-frequency distribution of iceberg fragments observed within the open ocean is therefore a function of the preferential loss of smaller icebergs through melting, facilitated by the tensile fracturing of larger icebergs.

233 Although initially unintuitive, the production of two different size-frequency distributions for icebergs, despite both being driven by fracture-dominated decay processes, may reflect the 234 intrinsically connected nature of lognormal and power-law distributions as demonstrated by the fact 235 that both may be produced using similar basic generative models<sup>51</sup>. The distributional breadth of 236 237 lognormal distributions increases as the active processes responsible for their generation become more intricate and numerous<sup>52</sup>. As the complexity, and thus breadth, of the distribution increases, 238 239 lognormal distributions begin to exhibit properties that are more commonly associated with power-law behaviour, providing a greater extent of overlap in which these two distributions are 240 indistinguishable<sup>53</sup> (Fig. 1c). 241

242 Reversing this logic, power-law distributed phenomena exposed to a breakage process tend 243 towards lognormality as the complexity of the degradation mechanisms reduces owing to the 244 emergence of a prevailing process subset. The dominance of three iceberg-fragmentation mechanisms within Vaigat, compared to the multitude of damage accretion mechanisms operating 245 246 at glacial termini, suggests that the number of processes driving iceberg disintegration reduces as icebergs drift away from the calving front. It is this simplifying phenomenon that drives the 247 248 transition from a power-law distribution at and proximal to a calving front, to the lognormal 249 distribution of iceberg sizes observed beyond. As the breadth of a lognormal distribution decreases as the number of processes responsible for its formation reduces<sup>51</sup>, the size-frequency 250 characteristics of iceberg populations will become increasingly lognormal as the mechanisms of 251 252 iceberg decay continue simplify with further distance away from the calving front (Fig. 1c). We

therefore anticipate that upon successful transport through Vaigat to deeper open waters, the absence of tidal grounding will further promote the dominance of a smaller number of wave and melt-based processes, consequently reinforcing the lognormality of the observed iceberg distribution.

257 The largest icebergs calved from Jakobshavn Isbræ are commonly over 1000 m in length, several hundred metres wide and exhibit keel depths of up to 900 m<sup>[37]</sup>. However, the seaward 258 259 transportation of the largest icebergs into Disko Bay is impeded by the relatively shallow water of the Isfjeldsbanken bank, stranding those with draughts >200 m until a sufficient reduction in size 260 occurs through fragmentation and/or melt<sup>4</sup>. Consequently, icebergs in the heavy tail of the 261 power-law distribution proximal to Jakobshavn Isbræ are left stranded, which may explain why the 262 263 lognormal approximation of iceberg areas within Vaigat marginally overestimates the likelihood of 264 the largest icebergs (Fig. 3a). The transition from a power-law to a lognormal distribution of iceberg 265 sizes can therefore begin in coastal waters, here close to the calving margin at Isfjeldsbanken, 266 although localised differences in coastal bathymetry will constrain the distance away from the ice 267 margin that this transition will initiate for other calving fronts.

268 The exclusion of the largest icebergs gives an example of a shift in the mechanisms of 269 iceberg decay operating within and beyond Disko Bay. Flexure of icebergs by waves can cause larger icebergs to fatigue, fracturing along pre-existing flaws<sup>40</sup>. However, for icebergs <1,000 m in 270 length, the impact of this process becomes negligible, leaving mass loss to be dominated by wave-271 272 related mechanisms such as the collapse of wavecut overhangs, buoyant failure of protruding 273 underwater rams and forced thermodynamic convection due to differential iceberg-water 274 velocities<sup>1</sup>. As the shallow water at Isfjeldsbanken prevents the very largest icebergs being 275 transported beyond Disko Bay, the complexity of iceberg-fragmentation processes operating within these waters is significantly reduced in comparison to those present at the calving front, permitting a 276 277 small number of decay mechanisms to dominate. Thus, whilst calving may generate power-law

frequency distributions of iceberg size, those leaving coastal waters may more likely adhere tolognormal size distributions.

280 Passive seismic monitoring therefore suggests that, owing to the fracture-driven iceberg 281 disintegration processes and dimensional reductions through melting, iceberg size-frequency 282 distributions will exhibit an increasingly definitive lognormal shift with drift away from the calving 283 front. This shift can be explained by the emergence of a dominant set of driving processes of 284 iceberg degradation as icebergs transit towards the open ocean. Although lognormal and power-law distributions both provide credible models for the mid-range values of many empirical data, 285 286 adequately representing the tail of a distribution has significant consequences for predicting the future behaviour of a phenomenon<sup>51</sup>. Whilst the heavy-tailed nature of power laws is required to 287 model icebergs in regions close to calving fronts, the use of a power-law distribution to estimate the 288 289 occurrence probability of seaward icebergs overpredicts the numbers of the largest and smallest 290 iceberg dimensions. A lognormal alternative is therefore needed to model iceberg distributions 291 accurately, and from this to derive risk and rates of iceberg disintegration.

292 Methods

# 293 Data collection & processing

Seismic signals generated by the processes of iceberg decay were recorded over a 49-day 294 period between 18<sup>th</sup> July and 4<sup>th</sup> September 2013 using six Güralp ESPCD broadband seismometers 295 296 installed in coastal locations along a 50 km stretch of the Vaigat Strait. The vertical component of 297 ground motion was detrended before being filtered using a 1–50 Hz Butterworth band-pass filter to 298 attenuate noise generated by ocean waves and distal earthquakes. Events were detected using the 299 ratio of the root mean square of short-term moving average (2 s) and long-term moving average 300 (60 s) windows, with events being retained for further analysis when the ratio exceeded a threshold 301 of 10. Cross-correlation of signal arrival times for each seismometer revealed that the detected 302 events are highly localised and generally only exceed the retention threshold at a single station.

Arrival time differences between the P-wave and S-wave component of signals demonstrates that
the source of the detected events is located within ~15 km of the associated seismometer (Fig. 1).
The detected signals are therefore sourced from processes operating within the Vaigat Strait and are
not duplicated across multiple seismometers.

#### **307** Event classification

308 Previous investigations of glaciological phenomena using passive seismic techniques have 309 demonstrated that different glacial processes are characterised by unique and highly distinctive 310 signal properties including dominant spectral frequency, event duration and the shape of the signal 311 onset and coda (Table 1). On detection, the characteristic frequency of individual signals was 312 examined using a combination of spectrograms and power spectral density estimations. Signal 313 duration and the profile of each signal onset and coda were manually described on the basis of 314 visual inspection. By using a threshold of signal power relative to the background noise in each detection envelope, it was possible to describe each detected signal in terms of an envelope 315 316 dominant frequency and duration, whilst the shape of the onset and coda of each signal was 317 classified as either impulsive or emergent. Using characteristic spectral frequency, duration and the 318 shape of the signal onset and coda as distinguishing properties, the 6842 events detected by the 319 seismometer array were grouped into three signal categories (Supplementary Fig. 1; Table 2). 320 Dominant spectral frequency was the clearest descriptor of signal type. However, the detected 321 events appear to be drawn from a continuum of processes; hence distinctions between signal types 322 were often ambiguous, with some events consisting of a sequence of all three types of events 323 combined (Supplementary Fig. 2).

324 Magnitude-frequency analysis

The energy released by each detected signal was calculated using methods introduced by Amitrano *et al.*<sup>54</sup> in a study of cliff collapse in Normandy, France. Following the grouping of all events into process-related classifications, the complementary cumulative size-frequency

328 distribution (CSFD), of signal energies was used to assess the size-frequency characteristics of each 329 iceberg disintegration process. The CSFD denotes the probability  $Pr(E \ge e)$  that the energy of an event, E (J), exceeds a given energy, e (J) <sup>[29]</sup>. Technical assessment<sup>55</sup> of the cumulative 330 331 size-frequency distributions generated by each iceberg disintegration mechanism indicated that both 332 power-law and lognormal distributions could provide potential models for the iceberg disintegration 333 process. In order to establish which of these competing distributions provided the most credible 334 model for the data, best-fit parameters for each distribution, including the minimum boundary for which the model applies  $(x_{min})$ , were derived using maximum likelihood estimation<sup>56</sup>. Directly 335 336 competing power-law and lognormal models for the data were then compared using Vuong's test<sup>57</sup> 337 — a likelihood-ratio test using the Kullback-Leiber criterion<sup>58</sup>. The sign of the likelihood ratio, R, indicates which distributional model provides the best fit to the data<sup>56</sup>. Here, R is positive if the 338 339 power-law model provides the better fit, negative if a lognormal model provides the best fit, and 340 zero if the fit provided by a distribution is indistinguishable from its alternative. The statistical 341 significance of the sign of R is given by a p-value. If p is small (p < 0.1), it is unlikely that the 342 observed sign of R may vary due to statistical fluctuations and thus may be used to comment on which distribution provides the most robust fit<sup>56</sup>. This analysis was conducted using the poweRlaw 343 package in the statistical software R<sup>58,59</sup>. This analysis is presented in further detail in 344 345 Supplementary Methods 1.

# 346 Comparison with satellite imagery

Icebergs in transit through the Vaigat Strait predominantly originate from Jakobshavn Isbræ.
A comparison between the dimensions of icebergs present within the Vaigat Strait compared to
those located proximal to Jakobshavn Isbræ was derived from a 1500 km<sup>2</sup> contemporaneous
Landsat 8 image of Vaigat (image ID: LC80110112013259LGN00, 09.16.2013) and a 1800 km<sup>2</sup>
Landsat 7 image of Jakobshavn Isbræ (image ID: LE70100112013196EDC00, 06.15.2013). Both
images share the same 30 m resolution, permitting a direct comparison to be made between the
areal properties of the icebergs present in each area. Iceberg areas were delineated using an

354 automated algorithm based on the contrast between the icebergs and the surrounding seawater. In 355 order to ensure that the iceberg dimensions were accurately delineated, the contrast between the 356 icebergs and the surrounding seawater was first increased using a global image threshold based on 357 Otsu's method. The validity of the mapping algorithm was then manually checked to ensure that 358 closely grouped patches of icebergs were not interpreted as one large ice mass; any areas in which 359 this issue was present were excluded from the analysis. The magnitude-frequency characteristics of 360 the extracted iceberg area populations were then analysed in the same manner as the seismic signal 361 energies, detailed in Methods: Magnitude-frequency analysis.

#### 362 **Correlation with tides**

Modelled hourly tide data for Ilulissat (~120 km from Vaigat) between the 18<sup>th</sup> of July and 363 30<sup>th</sup> of August was provided by the Danish Meteorological Institute, Copenhagen. A time-lag 364 correction of 2 hours relative to Ilulissat was applied in order to make the data applicable to Vaigat, 365 based upon analysis of the tidal signal in the seismic data<sup>60</sup>. The timing of detected icequakes was 366 367 compared against periodic components of the modelled tidal cycle in order to examine any potential 368 relationship between tidal forcing and seismic signal incidence. As the uneven time interval 369 between observations inhibits the application of typical fast Fourier transform techniques to assess 370 the periodicity of icequake signals, the Lomb-Scargle periodogram, which is designed to examine unevenly spaced time series<sup>61</sup>, was used to estimate the power spectrum of the icequake time series, 371 372 binned into 6-hour intervals. Bin width had a negligible effect on the calculated spectral power.

373 Funding & Acknowledgements

The study was supported by Polish National Science Centre grant no. 2011/01/B/ST10/01553.

375 Seismic equipment was provided by the Natural Environment Research Council Geophysical

376 Equipment Facility, SEIS-UK (loan number 984). The modelled tide data for Ilulissat was kindly

377 provided by Palle Bo Nielsen at the Danish Meteorological Institute. Logistical support was

378 provided by Arctic Station, Qeqertarsuaq. Thanks go to Antony Long and Nick Cox for

- 379 inspiring the ideas that underpin this work. M.C.S. is supported by NCN FUGA Fellowship
- 380 (2013/08/S/ST10/00585).

#### **381** Author contributions

- 382 J.D.K. conceived the study with J.W. and N.J.R. Funding for fieldwork was awarded to W.S. and
- 383 M.C.S. The SEIS-UK equipment loan was obtained by N.J.R and E.V.J., and W.S., M.C.S., N.J.R.,
- 384 S.A.D. and E.V.J. conducted the fieldwork. The data processing was undertaken by J.D.K. and
- 385 E.V.J. with input from V.S.L. and D.E.H. Analysis and interpretation was conducted by all authors.
- 386 J.D.K. wrote the manuscript with input from all co-authors.

#### 387 **Competing financial interests**

388 The authors declare no competing financial interests.

#### 389 Data availability

- 390 The seismic data used in this study are available from the IRIS MDC data repository
- 391 (http://ds.iris.edu/ds/nodes/dmc/data/), initially on request of the corresponding author. Following
- 392 publication, the data will be made fully open-access after 2–3 years, in line with SEIS-UK policy on

393 data availability<sup>62</sup>.

# 394 **References**

Wagner, T.J.W. *et al.* The "footloose" mechanism: Iceberg decay from hydrostatic stresses.
 *Geophys. Res. Lett.* 41, 5522-5529 (2014).

397

400

401 3. Bhatia, M.P. *et al.* Greenland meltwater as a significant and potentially bioavailable source
402 of iron to the ocean. *Nature Geosci.* 6, 274-278 (2013).

403

Duprat, L.P., Bigg, G.R. & Wilton, D.J. Enhanced Southern Ocean marine productivity due
 to fertilization by giant icebergs. *Nature Geosci.* 9, 219-221 (2016).

405	4.	Schumann, K., Völker, D. & Weinrebe, W.R. Acoustic mapping of the Ilulissat Ice Fjord
406		mouth, West Greenland. Quat. Sci. Rev. 40, 78-88 (2012).
407		
408	5.	Bigg, G.R., Wadley, M.R., Stevens, D.P. & Johnson, J.A. Modelling the dynamics and
409		thermodynamics of icebergs. Cold Reg. Sci. Technol. 26, 113-135 (1997).
410		
411	6.	Åström, J.A. et al. Termini of calving glaciers as self-organized critical systems. Nature
412		<i>Geosci.</i> <b>7</b> , 874-878 (2014).
413		
414	7.	Marko, J.R. Small icebergs and iceberg fragments off Newfoundland: Relationships to
415		deterioration mechanisms and the regional iceberg population. Atmos. Ocean 34, 549-579
416		(1996).
417	8.	Tournadre, J., Girard-Ardhuin, F. & Legrésy, B. Antarctic icebergs distributions, 2002-2010.
418		J. Geophys. Res. 117, C05004 (2012).
419		
420	9.	Neave, K.G. & Savage, J.C. Icequakes on the Athabasca Glacier. J. Geophys. Res. 75, 1351-
421		1362 (1970).
422	10.	West, M.E., Larsen, C.F., Truffer, M., O'Neel, S. & LeBlanc, L. Glacier microseismicity.
423		<i>Geology</i> <b>38,</b> 319-322 (2010).
424	11.	Deichmann, N. et al. Evidence for deep icequakes in an Alpine glacier. Ann. Glac. 31, 85-90
425		(2000).
426	12.	Walter, F., Deichmann, N. & Funk, M. Basal icequakes during changing subglacial water
427 428		pressures beneath Gornergletscher, Switzerland. J. Glaciol. 54, 511-521 (2008).
429	13.	Weaver, C.S. & Malone, S.D. Seismic evidence for discrete glacier motion at the rock-ice
430		interface. J. Glac. 23,171-184 (1979).
431		

432	14.	Martin, S., Drucker, R., Aster, R., Davey, F., Okal, E., Scambos, T. & MacAyeal, D.
433		Kinematic and seismic analysis of giant tabular iceberg breakup at Cape Adare, Antarctica.
434		J. Geophys. Res. 115, B06311 (2010).
435		
436	15.	O'Neel, S., Marshall, H.P., McNamara, D.E. & Pfeffer, W.T. Seismic detection and analysis
437		of icequakes at Columbia Glacier, Alaska. J. Geophys. Res. 112, F03S23 (2007).
438		
439	16.	Nettles, M. et al. Step-wise changes in glacier flow speed coincide with calving and glacial
440		earthquakes at Helheim Glacier, Greenland. Geophys. Res. Lett. 35, (2008).
441		
442	17.	Dziak, R.P. et al. Life and death sounds of iceberg A53a. Oceanography 26, 10-13 (2013).
443		
444	18.	van den Broeke, M. et al. Partitioning recent Greenland mass loss. Science 326, 984-986
445		(2009).
446 447	19.	Chen, J.L., Wilson, C.R. & Tapley, B.D. Satellite gravity measurements confirm accelerated
448		melting of Greenland ice sheet. Science 313, 1958-1960 (2006).
449		
450	20.	Rignot, E. & Kanagaratnam, P. Changes in the velocity structure of the Greenland Ice
451		Sheet. Science <b>311</b> , 986-990 (2006).
452		
453	21.	Moon, T., Joughin, I., Smith, B. & Howat, I. 21st-century evolution of Greenland outlet
454		glacier velocities. Science 336, 576-578 (2012).
455		
456	22.	Enderlin, E.M., Hamilton, G.S., Straneo, F. & Sutherland, D.A. Iceberg meltwater fluxes
457		dominate the freshwater budget in Greenland's iceberg-congested glacial fjords. Geophys.
458		<i>Res. Lett.</i> <b>43</b> , (2016).

459 460	23.	Bamber, J., den Broeke, M., Ettema, J., Lenaerts, J. & Rignot, E. Recent large increases in
461		freshwater fluxes from Greenland into the North Atlantic. Geophys. Res. Lett. 39, (2012).
462 463	24.	Luo, H. <i>et al.</i> Oceanic transport of surface meltwater from the southern Greenland ice
404		sneet. <i>Nat. Geosci.</i> 9, 528-532 (2016).
465 466	25.	Straneo, F. et al. Impact of fjord dynamics and glacial runoff on the circulation near Helheim
467		Glacier. Nat. Geosci. 4, 322-327 (2011).
468 469	26.	Echelmeyer, K., Clarke, T.S. & Harrison, W.D. Surficial glaciology of Jakobshavns Isbræ,
470		West Greenland: Part I. Surface morphology. J. Glac. 37, 368-382 (1991).
471		
472	27.	Barlow, J. et al. Modeling cliff erosion using negative power law scaling of
473		rockfalls. <i>Geomorphology</i> . <b>139</b> , 416-424 (2012).
474		
475	28.	Guzzetti, F., Malamud, B.D., Turcotte, D.L. & Reichenbach, P. Power-law correlations of
476		landslide areas in central Italy. Earth Planet. Sci. Lett. 195, 169-183 (2002).
477 478	29.	Canassy, P.D., Faillettaz, J., Walter, F. & Huss, M. Seismic activity and surface motion of a
479		steep temperate glacier: a study on Triftgletscher, Switzerland. J. Glac. 58, 513-528 (2012).
480 481 482	30.	Sinha, N.K. Constant strain-and stress-rate compressive strength of columnar-grained ice. <i>J. Mat. Sci.</i> <b>17</b> , 785-802 (1982).
483 484	31.	Weiss, J. Scaling of fracture and faulting of ice on earth. Surv. Geophys. 24, 185-227 (2003).
485		
486 487	32.	Walter, F., Dalban Canassy, P., Husen, S. & Clinton, J.F. Deep icequakes: What happens at the base of Alpine glaciers?, <i>J. Geophys Res.</i> <b>118</b> , 1720-1728 (2013).
,		

489	33.	Allstadt, K. & Malone, S.D. Swarms of repeating stick-slip icequakes triggered by snow
490		loading at Mount Rainier volcano. J. Geophys. Res. 119, 1180-1203 (2014).
491		
492	34.	Haykin, S.S., Lewis, E.O. & Raney, R.K. (eds). Remote sensing of sea ice and icebergs.
493		(John Wiley and Sons, 1994).
494		
495	35.	Veitch, B., Williams, M., Gardner, A. & Liang, B. Field observations of iceberg
496		deterioration. Technical Report, 20-64, PERD/CHC, (National Research Council Canada,
497		2001).
498		
499	36.	Gagnon, R. E., Williams, F. M. & Sinha, N. K. High speed video observations of fracture
500		from beam bending experiments on sea ice. Proc. Port Ocean Eng. Under Arctic Cond. 25-
501		36 (1999).
502		
503	37.	Amundsen, J.M. et al. Glacier, fjord, and seismic response to recent large calving events,
504		Jakobshavn Isbræ, Greenland. Geophys. Res. Lett. 35, L22501 (2008).
505		
506	38.	Bartholomaus, T.C., Larsen, C.F., O'Neel, S. & West, M.E. Calving seismicity from
507		iceberg-sea surface interactions. J. Geophys. Res. 117, F04029 (2012).
508		
509	39.	Tsai, V.C., Rice, J.R. & Fahnestock, M. Possible mechanisms for glacial earthquakes, J.
510		Geophys. Res. 113, (2008).
511	40.	Robe, R.Q. Iceberg Drift and Deterioration. In Colbeck, S.C. (ed). Dynamics of Snow and
512		Ice Masses 211-257 (Academic Press 1980)
012		
513	41.	Sulak, D.J., Sutherland, D.A., Enderlin, E.M., Stearns, L.A. & Hamilton, G.S. Iceberg
514		properties and distributions in three Greenlandic fjords using satellite imagery, Ann. Glac.
515		1-15 (2017).
516	42.	Bak, P., Tang, C. & Wiesenfeld, K. Self-organized criticality: An explanation of the 1/f
517		noise, Phys. Rev. Lett. 59, 381 (1987).
518		

519	43.	Benn, D.I., Warren, C.R. & Mottram, R.H. Calving processes and the dynamics of calving
520		glaciers. Earth Sci. Rev. 82, 143-179 (2007).
521 522	44.	Walter, F. et al. Iceberg calving during transition from grounded to floating ice: Columbia
523		Glacier, Alaska. Geophys. Res. Lett. 37, (2010).
524		
525	45.	Savage, S.B., Crocker, G.B., Sayed, M. & Carrieres, T. Size distributions of small ice pieces
526		calved from icebergs. Cold Reg. Sci. Technol. 31, 163-172 (2000).
527		
528	46.	Epstein, B. The mathematical description of certain breakage mechanisms leading to the
529		logarithmico-normal distribution. J. Franklin Inst. 244, 471-477 (1947).
530		
531	47.	Kolmogoroff, A.N. Über das logarithmisch normale Verteilungsgesetz der Dimensionen der
532		Teilchen bei Zerstückelung. Izv. Akad. Nauk. S.S.S.R. 31, 99-101 (1941).
533		
534	48.	Gibrat, R. Les inégalités économiques: applications: aux inégalitês des richesses, à la
535		concentration des entreprises, aux populations des villes, aux statistiques des familles, etc:
536		d'une loi nouvelle: la loi de l'effet proportionnel. (Librairie du Recueil Sirey, 1931.).
537		
538	49.	Blackwood, L.G. The lognormal distribution, environmental data, and radiological
539		monitoring. Environ. Monit. Assess. 21, 193-210 (1992).
540		
541	50.	Aitchison, J. & Brown, J.A.C. The Lognormal Distribution with special reference to its uses
542		in economics. (Cambridge Univ. Press, 1957).
543 544	51	Mitzenmacher M. A brief history of generative models for power law and lognormal
545		distributions Internet mathematics 1 226-251 (2004)
JTJ		(200+).

546 547	52.	Montroll, E.W. & Shlesinger, M.F. On 1/f noise and other distributions with long tails. <i>Proc.</i>
548		Nat. Acad. Sci. 9, 3380-3383 (1982).
549		
550	53.	West, B.J. & Schlesinger, M. The Noise in Natural Phenomena. Amer. Sci. 78, 40-45 (1990).
551		
552	54.	Amitrano, D., Grasso, J.R. & Senfaute, G. Seismic precursory patterns before a cliff collapse
553		and critical point phenomena. Geophys. Res. Lett. 32, L08314 (2005).
554		
555	55.	Cirillo, P. Are your data really Pareto distributed?. Physica A 392, 5947-5962 (2013).
556		
557	56.	Clauset, A., Shalizi, C.R. & Newman, M.E. Power-law distributions in empirical data. SIAM
558		<i>review</i> <b>51</b> , 661-703 (2009).
559		
560	57.	Vuong, Q.H. Likelihood ratio tests for model selection and non-nested hypotheses.
561		Econometrica <b>57</b> , 307-333 (1989).
562 563 564 565	58.	Gillespie, C.S. Fitting Heavy Tailed Distributions: The poweRlaw Package. J. Stat. Soft. 64, 1-16 (2015).
566	59.	R Development Core Team. R: A language and environment for statistical computing. R
567		Foundation for Statistical Computing, Austria (2008).
568		
569	60.	Vann Jones (née Norman), E. C., Rosser, N. J., Brain, M. J. & Petley, D. N. Quantifying the
570		environmental controls on erosion of a hard rock cliff. Mar. Geol. 363, 230-242 (2015).
571		
571 572	61.	Press, W.H., Teukolsky, S.A., Vetterling, W.T. & Flannery, B.P. Numerical Recipes: The

574 575	62.	Brisbourne, A. How to store and share geophysical data. <i>Astronomy and Geophysics</i> , <b>53</b> ,
576		4 19-4 20 (2012)
577		
570	(2)	
578	63.	Jakobsson, M. <i>et al.</i> The International Bathymetric Chart of the Arctic Ocean (IBCAO)
579		Version 3.0, Geophys. Res. Lett. 39, L12609 (2012).
580	61	Nattles M. & Electröm, C. Clasici Forthqueles in Crearland and Antoretics, Ann. Dev
582	04.	Farth Planet Sci. <b>38</b> A67 A91 (2010)
583		Lann 1 lanet. Set. 30, 407-471 (2010).
584	65.	MacAyeal., D.R., Okal., E.A., Aster, R.C. & Bassis, J.N. Seismic and hydroacoustic tremor
585		generated by colliding icebergs. J. Geophys. Res. 113, F03011 (2008).
586		
587		
507		
588		
589		
590		
591		
592		
593		
594		
595		
<b>T</b> O 6		
596		
597		
500		
598		



Figure 1 | The lognormal shift in iceberg size magnitude-frequency distributions with distance away
 from the calving front at Jakobshavn Isbræ. a) Location of the study site in West Greenland and b) the
 area surrounding Vaigat displaying regional bathymetry and the locations referred to in the text.

603 Seismometers are displayed as red dots with their 15 km effective detection radiuses shown as blue circles. 604 c) Idealised comparison between the probability distributions, p(x), of an inverse power-law (green line 1.) 605 and two lognormal distributions (curves) of decreasing complexity, labelled 2. and 3. As the hierarchy of 606 processes responsible for the lognormal distribution becomes more complex (curve 3. to 2.), the distribution 607 becomes broader, providing a greater degree of overlap with the inverse power-law distribution<sup>53</sup> (line 1.). 608 The complexity of the breakage process responsible for generating the magnitude-frequency distribution of 609 iceberg sizes decreases with distance away from the calving front of Jakobshavn Isbræ owing to the 610 emergence of a dominant set of decay mechanisms. As a result, the power-law distribution of iceberg sizes 611 initially present proximal to Jakobshavn Isbræ (1.) evolves towards lognormal scaling, and becomes more

612 characteristically lognormal, as icebergs transit from Jakobshavn Isbræ to Vaigat (2.) and towards the open

613 ocean (3.). Bathymetric data is obtained from the IBCAO V. 3.0 dataset<sup>63</sup> and drawn using ArcMap 10.3.





Figure 2 | Cumulative size-frequency distributions,  $Pr(E \ge e)$ , for the energy released by iceberg617fragmentation processes: a) cracking, b) microfracturing, c) calving and rolling, d) all detected events.618Optimal lognormal and power-law approximations of the data are displayed as dashed red and solid blue619lines, respectively.



Figure 3 | Optimal lognormal (red line) and power-law (blue line) approximations of the distribution
of planform iceberg areas located: a) within Vaigat and b) proximal to Jakobshavn Isbræ. The slope of
the fitted power law in b) is 2.4.

**Table 1.** Glaciological processes known to generate seismic tremors and their associated waveform geometry, frequency and duration. Note that iceberg calving, grounding and ice mélange interaction processes are associated with a significantly longer duration (in the order of minutes to hours) than the other phenomena which may be measured on timescales in the order of seconds.

Glaciological process	Waveform geometry	Characteristic frequency (Hz)	Typical duration	Reference
Surface crevassing and ice fracture	Impulsive onset and abruptly declining coda	10–30	0.1 – 2.5 (s)	[7], [9], [12], [29]
Iceberg calving and capsize	Emergent onset, cigar-shaped envelope, long-duration coda, absence of P- or S-waves, peaks often coincide with 'Worthington jets' produced by cavity collapse	1–5	5 - 30 + (s) (up to 1 <u>hour</u> depending on iceberg dimensions)	[7], [37], [38], [64]
Basal sliding	No surface waves	1–25	_	[14]
Iceberg interaction with ice mélange	Multiple harmonic frequencies	0.5–30 with multiple harmonics	30 – 60 ( <u>minutes</u> )	[37]
Hydraulic movement in glacial water channels	Emergent onset, lack of distinct S- waves	6–15	1 – 10 (s)	[10]
Iceberg grounding and ploughing	Long duration, monochromatic frequency	0.5–1.5	~ 2 ( <u>hours</u> )	[14]
Hydrofracturing	Impulsive onset	20–35	1 – 10 (s)	[10]
Iceberg harmonic tremor	Multiple harmonic frequencies with a distinctive 'chevron' pattern	1–10 with multiple harmonics	1500 (s)	[65]

**Table 2.** Characteristic frequencies, duration and waveform descriptions for the three types of events detected by the seismometer array over a 49-day period.

Signal type	Description of signal onset	Description of signal coda	Characteristic frequency (Hz)	Typical duration (s)	Number detected	Process interpretation
1	Impulsive	Abrupt termination	30–40	1	1979	Microfracturing
2	Impulsive	Abrupt termination	10–30	1–5	3271	Tensile fracturing and crack enlargement.
3	Emergent; amplitude increases over time	Gradual decline in amplitude	1–10	3–15	1592	Iceberg calving, capsize and rolling