The uppermost mantle seismic velocity structure of West Antarctica from Rayleigh wave tomography: insights into tectonic structure and geothermal heat flow

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¹ Abstract

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We present a shear wave model of the West Antarctic upper mantle to $\sim 200 \,\mathrm{km}$ depth 3 with enhanced regional resolution from the 2016-2018 UK Antarctic Seismic Network. 4 The model is constructed from the combination of fundamental mode Rayleigh wave 5 phase velocities extracted from ambient noise (periods 8-25s) and earthquake data 6 by two-plane wave analysis (periods 20-143s). We seek to (i) image and interpret 7 structures against the tectonic evolution of West Antarctica, and (ii) extract infor-8 mation from the seismic model that can serve as boundary conditions in ice sheet 9 and glacial isostatic adjustment modelling efforts. The distribution of low veloc-10 ity anomalies in the uppermost mantle suggests that recent tectonism in the West 11 Antarctic Rift System (WARS) is mainly concentrated beneath the rift margins and 12 largely confined to the uppermost mantle $(<180 \,\mathrm{km})$. On the northern margin of 13 the WARS, a pronounced low velocity anomaly extends eastward from beneath the 14 Marie Byrd Land dome toward Pine Island Bay, underlying Thwaites Glacier, but not 15 Pine Island Glacier. If of plume-related thermal origin, the velocity contrast of $\sim 5\%$ 16 between this anomaly and the inner WARS translates to a temperature difference 17 of $\sim 125-200$ °C. However, the strike of the anomaly parallels the paleo-Pacific con-18 vergent margin of Gondwana, so it may reflect subduction-related melt and volatiles 19 rather than anomalously elevated temperatures, or a combination thereof. Motivated 20 by xenolith analyses, we speculate that high velocity zones imaged south of the Marie 21 Byrd Land dome and in the eastern Ross Sea Embayment might reflect the composi-22 tional signature of ancient continental fragments. A pronounced low velocity anomaly 23 underlying the southern Transantarctic Mountains (TAM) is consistent with a pub-24 lished lithospheric foundering hypothesis. Taken together with a magnetotelluric 25 study advocating flexural support of the central TAM by thick, stable lithosphere, 26

this points to along-strike variation in the tectonic history of the TAM. A high veloc-27 ity anomaly located in the southern Weddell Sea Rift System might reflect depleted 28 mantle lithosphere following the extraction of voluminous melt related to Gondwana 29 fragmentation. Lithospheric thickness estimates extracted from 1D shear wave veloc-30 ity profiles representative of tectonic domains in West Antarctica indicate an average 31 lithospheric thickness of $\sim 85 \,\mathrm{km}$ for the WARS, Marie Byrd Land, and Thurston Is-32 land block. This increases to ~ 96 km in the Ellsworth Mountains. A surface heat flow 33 of ${\sim}60\,\mathrm{mW/m^2}$ and attend ant geotherm best explains lithospheric mantle shear wave 34 velocities in the central WARS and in the Thurston Island block adjacent to Pine 35 Island Glacier; a $\sim 50 \,\mathrm{mW/m^2}$ geotherm best explains the velocities in the Ellsworth 36 Mountains, and a $\sim 60 \,\mathrm{mW/m^2}$ geotherm best explains a less well-constrained velocity 37 profile on the southern Antarctic Peninsula. We emphasise that these are regional 38 average (many hundreds of km) heat flow estimates constrained by seismic data with 39 limited sensitivity to upper crustal composition. 40

41 **1** Introduction

West Antarctica owes much of its tectonic heritage to the Jurassic breakup of Gond-42 wana and ensuing dispersal of microplate fragments (e.g., Dalziel & Elliot, 1982; 43 Dalziel, 1992). The development of the West Antarctic Rift System (WARS), the 44 uplift of the Transantarctic Mountains (TAM) and the impact of a putative mantle 45 plume beneath Marie Byrd Land (MBL) have dominated the late Cretaceous to Pa-46 leogene evolution of West Antarctica (Figure 1) (e.g., LeMasurier & Landis, 1996; 47 Fitzgerald, 2002). With geological exposures limited by the West Antarctic Ice Sheet 48 (WAIS), delineation of tectonic domains and recent tectonism is reliant on geophysi-49 cal probing. Owing to the deployment of broadband seismometer arrays, the seismic 50 structure of much of the Antarctic crust and upper mantle is now reasonably well 51 mapped (e.g., An et al., 2015b; Heeszel et al., 2016; Shen et al., 2018). 52

We construct a shear wave model based on fundamental mode Rayleigh wave phase 53 velocities focussing on the uppermost mantle structure (<200 km) of West Antarctica. 54 The model offers enhanced regional resolution through the inclusion of stations from 55 the 2016-2018 UK Antarctic Seismic Network (UKANET, Figure 1). In the first half 56 of this paper we describe the seismic data, processing and inversion, and interpret the 57 structures imaged against the tectonic evolution of West Antarctica. In the second 58 half we extract information that can be used to improve the accuracy of ice sheet and 59 glacial isostatic adjustment (GIA) modelling efforts. Geothermal heat flow moderates 60 ice sheet behaviour: it affects the viscosity of basal ice and, if sufficiently high, can 61 generate lubricating meltwater that reduces friction with the bed (e.g., Martos et al., 62 2017). Pine Island Glacier and Thwaites Glacier in West Antarctica (Figure 1) are of 63 particular concern because they are thought susceptible to marine ice sheet instability 64 (e.g., Barletta et al., 2018). GIA is sensitive to lithospheric thickness and its lateral 65 variation (e.g., Nield et al., 2018). From our shear wave model, we extract lithospheric 66

⁶⁷ thicknesses and model the regional average geotherms and heat flows best describing
⁶⁸ 1D velocity profiles at representative tectonic locations in West Antarctica.

⁶⁹ 2 Tectonic Setting

East Antarctica was amalgamated from Archean nuclei in the Mesoproterozoic, even-70 tually forming the core of Gondwana (e.g., Dalziel, 1992). The Mesozoic fragmenta-71 tion of Gondwana was preceded by the emplacement of the Karoo-Ferrar large igneous 72 province in East Antarctica and southern Africa at \sim 185-177 Ma (e.g., Storey & Kyle, 73 1997; Fitzgerald, 2002, and references therein) and the development of the Weddell 74 Sea Rift System (WSRS), a broad extensional/transtensional province within a dis-75 tributed plate boundary between East and West Antarctica (e.g., Jordan et al., 2017). 76 Karoo-Ferrar magmatism has been linked with a putative mantle plume in the proto-77 Weddell Sea region, potentially a driver for Gondwana breakup (e.g., Storey & Kyle, 78 1997). 79

West Antarctica is regarded as an assemblage of discrete crustal blocks separated 80 by subglacial depressions. Three of the main four blocks - the Antarctic Peninsula, 81 Thurston Island and Marie Byrd Land - are fore-arc and magmatic-arc terranes de-82 veloped along the paleo-Pacific margin of Gondwana (e.g., Dalziel, 1992). The fourth 83 block, the Haag-Ellsworth Whitmore (HEW) block, is regarded as an allochthonous 84 continental fragment translated and rotated to its present location from an original 85 pre-Gondwana-breakup position close to the East Antarctic plate and/or southern 86 Africa. Exposed lithologies in the HEW block include a ~ 13 km thick Paleozoic sedi-87 mentary sequence in the Ellsworth Whitmore Mountains, and Precambrian basement 88 dated to ~ 1 Ga in the Haag Nunataks (e.g., Storey & Kyle, 1997; Jordan et al., 2017, 89 and references therein). 90

⁹¹ The tectonic regime in West Antarctica switched from compressional to extensional

following subduction of the Pacific-Phoenix spreading center at $\sim 110-105$ Ma. The 92 West Antarctic Rift System formed as MBL and Thurston Island moved away from the 93 East Antarctica craton, with the major WARS extensional phase occurring between 94 $\sim 105-85$ Ma (e.g., Fitzgerald, 2002, and references therein). Paleogene extension was 95 limited to the western Ross Sea and accompanied by rapid exhumation and uplift of 96 the Transantarctic Mountains. In MBL an estimated maximum $\sim 3 \,\mathrm{km}$ of tectonic 97 uplift associated with alkaline volcanism beginning at ca. 28-30 Ma is cited as evidence 98 of a mantle plume (e.g., LeMasurier & Landis, 1996). Others favour a model of 99 subduction-related alkaline magma genesis in MBL (e.g., Finn et al., 2005). Inferred 100 Neogene reactivation of subglacial troughs in central West Antarctica has been linked 101 with Neogene extensional pulses in the western Ross Sea (e.g., Lloyd et al., 2015, and 102 references therein). 103

¹⁰⁴ 3 Seismic Arrays

The International Polar Year 2007-2008 motivated the first deployment of yearround broadband seismometer arrays in the interior of Antarctica. As part of the POLENET-ANET project, a backbone array was deployed across Antarctica (Figure 1). The extant array comprises a mixture of cold-rated Güralp CMG-3T 120 s and Nanometrics Trillium 240 s seismometers sampling at 1 and 40 samples per second (sps).

Denser temporary arrays have intermittently supplemented the POLENET-ANET backbone array in West Antarctica, the most recent of which was the 2016-2018 UKANET array. This consisted of 10 cold-rated Güralp CMG-3T 120 s seismometers sampling at 1 and 100 sps (Figure 1 and Table S1). The 2015-2017 POLENET-ANET mini-array was complementary in design and location to the UKANET array.

¹¹⁶ Additional coverage is provided by the Antarctic Seismographic Argentinean Ital-

ian Network (ASAIN), station PMSA of the Global Seismographic Network (GSN)
and the 1997-1999 Seismic Experiment in Patagonia and Antarctica network (SEPA)
shown in Figure 1.

¹²⁰ 4 Two-Plane-Wave Tomography

Surface wave amplitudes and phases observed across seismic arrays often exhibit ef-121 fects reminiscent of interference. This motivated Forysth & Li (2005) to model the 122 wavefield as the superposition of two interfering plane waves. We applied this two-123 plane-wave method to fundamental mode Rayleigh waves recorded on the UKANET, 124 POLENET-ANET, ASAIN and SEPA arrays and PMSA station over the periods 125 1997-1999 and 2010-2018. To garner good quality waveforms, we examined earth-126 quakes with magnitudes >5.5 located within a distance of 120° of the composite 127 seismic array. Earthquakes located within $\sim 30^{\circ}$ of the array were excluded because 128 the wave fronts cannot be considered planar at incidence. 129

An initial cull of earthquakes giving poor signal-to-noise ratio seismograms was carried 130 out by visual inspection. Instrument responses were deconvolved from the remaining 131 seismograms and these filtered into $12 \times 10 \,\mathrm{mHz}$ wide frequency bands with centre 132 periods ranging from 20 to 143 s using a zero-phase-shift, four-pole Butterworth filter 133 centred at the period of interest (Figure 2). Next, for each earthquake a window was 134 manually defined at each period to isolate the fundamental mode Rayleigh waves from 135 other seismic phases and/or interfering lateral refractions. At each period, only those 136 earthquakes vielding high signal-to-noise ratio Rayleigh waves at at least five stations 137 were considered for two-plane-wave tomography (2PWT). Out of a total of ~ 2700 138 earthquakes screened, 457 were deemed suitable for analysis (Figure 2). Following 139 Forysth & Li (2005), we assigned a prior data uncertainty of 10% to the phase and 140 amplitude of each Rayleigh wave. 141

In the 2PWT inversion, at each period the Rayleigh wave phase velocity map best 142 explaining phase and amplitude variations between stations was inferred on a grid 143 with a node spacing of 100 km spanning West Antarctica. Being predicated on the 144 assumption of planar wave fronts, the validity of 2PWT varies inversely with the areal 145 extent of the seismic array. In response, we subdivided the expansive composite array 146 into three sub-arrays approximately coincident with the Antarctic Peninsula, eastern 147 West Antarctica and central West Antarctica. In this scheme, a given earthquake 148 is effectively treated as three separate earthquakes, each incident on one of the sub-149 arrays. Following Yang & Forsyth (2006), finite frequency sensitivity kernels were 150 used to represent the sensitivities of Rayleigh wave phases and amplitudes to struc-151 ture. A smoothing length scale of 140 km gave the best compromise between unduly 152 rough models arising from over-fitting data at the shortest length scales and under-fit 153 models at the longest length scales (Figure S1). Using the 1D average phase velocity 154 curve inferred by Heeszel et al. (2016) as a starting model, we initially inverted for 155 a 1D average phase velocity curve representing our study area to serve as a starting 156 model for the 2D tomographic inversions (Figure 3). 157

¹⁵⁸ 5 Rayleigh Wave Phase Velocities

Figure 4 shows the inferred 2D Rayleigh wave phase velocity uncertainty, calculated from the posterior model covariance matrix, at periods 25, 80 and 125 s. As expected, the uncertainty is least where the concentration of seismic stations is greatest and increases toward the grid periphery. Superimposed on the lateral variations is a trend of increasing uncertainty with increasing period, a reflection of the progressively increasing wavelength of the Rayleigh waves and hence decreasing resolution.

Figure 4 also shows the resolving capability of the inversion. The resolution matrix
indicates that the morphology of velocity anomalies of length scale 400 km is recovered

with high fidelity within the polygon on Figure 4 at all periods. At periods 125 and 143 s there is some diminution in amplitudes at this length scale, but at all shorter periods amplitude recovery is generally better than 90%. Amplitude resolution at periods 125 and 143 s reaches this level for a length scale of 500 km.

The resolution matrix gives an overly optimistic picture of resolution at peripheral grid regions beyond the footprint of the seismic array. In subsequent plots we confine our discussion to the region enclosed by the polygon. Within this region (i) phase velocity uncertainty is generally less than $\sim 0.02-0.03$ km/s at periods below 80 s and less than ~ 0.05 km/s at periods 100-143 s, (ii) the resolution matrix indicates that velocity structure of length scale 400-500 km is imaged with high fidelity and (iii) imaged velocity structure transitions credibly between periods.

Figure 5 shows Rayleigh wave phase velocity maps at selected periods. At periods 178 \sim 20-30 s Rayleigh wave propagation is most sensitive to variations in crustal thick-179 ness: if the crust is thick, Rayleigh waves at these periods largely sample lower crustal 180 rock, whereas if the crust is thin they largely sample seismically-faster mantle rock. 181 At 25s for example, relatively slower phase velocities coincident with the TAM, the 182 HEW block, MBL, the southern Antarctic Peninsula and northern WSRS are consis-183 tent with thicker crust (e.g., Chaput et al., 2014; O'Donnell & Nyblade, 2014; Shen 184 et al., 2018). In contrast, relatively faster phase velocities underlying the Ross and 185 Amundsen Sea Embayments in the WARS and in the southern WSRS are likely the 186 signature of mantle rock, and hence thinner crust. 187

At periods 40 s and above the Rayleigh wave phase velocities predominantly reflect uppermost mantle structure. The geological dichotomy of Antarctica is here apparent: slower phase velocities characterising the West Antarctic uppermost mantle contrast with faster velocities underlying East Antarctica. Prominent slow phase velocity anomalies at these periods occur beneath MBL and a portion of the southern TAM. Notably, the slow velocity anomaly underlying MBL extends eastward beyond the MBL topographic dome toward Pine Island Bay. Offshore MBL a slow velocity
anomaly coincides with the location of the Marie Byrd Seamounts and is conceivably
the source thereof.

¹⁹⁷ 6 Shear Wave Velocities

At each grid node, a phase velocity dispersion curve (periods 20-143 s) was extracted 198 by sampling the 2PWT phase velocity maps. These curves were merged with counter-199 parts extracted from ambient noise tomography (ANT) Rayleigh wave phase velocity 200 maps developed by the authors (O'Donnell et al., 2018). The shorter period ANT 201 data (periods 8-25s) have a greater sensitivity to crustal structure than the 2PWT 202 data. Figure 6 compares ANT- and 2PWT-inferred phase velocity maps at 25 s and 203 shows an example of a composite 8-143 s phase velocity dispersion curve obtained by 204 weighted least squares polynomial regression of the ANT- and 2PWT-curves. Dif-205 ferences in processing, inversion and regularisation schemes result in minor disparity 206 between ANT- and 2PWT-inferred velocities, but they generally agree within uncer-207 tainty bounds at overlapping periods. The areal extent of the ANT model domain, 208 however, is less extensive than the 2PWT domain, so merged ANT-2PWT dispersion 209 curves are restricted to the ANT domain. The phase velocity dispersion curves were 210 subsequently inverted for 1D shear wave velocity structure. Because Rayleigh waves 211 are most sensitive to vertically-polarised shear wave velocity, V_{SV} , we inferred V_{SV} 212 rather than isotropic V_S . 213

The V_{SV} models were parameterised by ice and/or water layers overlying crustal and uppermost mantle layers. Ice thicknesses and water depths were taken from BEDMAP2 and allowed to vary within their uncertainty limits (Fretwell et al., 2013). The ice shear wave velocity was permitted to range between 1.82-2.02 km/s with a density fixed at 910 kg/m³. We opted to not invert for a sedimentary layer because (1)

Rayleigh waves have limited sensitivity to shallow crustal structure in the period range 219 considered and (2) sediment thickness estimates to guide the inversion are extremely 220 limited. The 1D V_{SV} structure of the underlying crustal layer was parameterised 221 using 4 cubic B-splines and a crustal thickness permitted to vary $\pm 5 \,\mathrm{km}$ from initial 222 estimates extracted from the An et al. (2015b) Antarctic crustal model. The 1D 223 uppermost mantle V_{SV} structure was parameterised using 5 cubic B-splines to a depth 224 of 250 km, below which PREM V_{SV} values were adopted. In a Bayesian framework, 225 we permitted crustal and uppermost mantle V_{SV} velocities to explore a broad $\pm 20\%$ 226 range around initial PREM V_{SV} velocities, a range which encompasses published 227 Antarctic velocity models (e.g., An et al., 2015b). This suite of constraints informed 228 the prior model probability density function (PDF). 229

The likelihood function for dispersion curve prediction used the *Mineos* package 230 (https://geodynamics.org/cig/software/mineos/). Crustal compressional wave 231 velocities and densities were scaled from inferred shear wave velocities using regres-232 sions reported in Brocher (2005), while upper mantle counterparts were scaled using 233 a Vp/Vs ratio of 1.74 and Birch's law (Birch, 1961). PREM Q values were used 234 to correct for anelastic attenuation. A Markov chain Monte Carlo sampling scheme 235 based on the Delayed Rejection Adaptive Metropolis algorithm built the posterior 236 model PDF from the final 2,500 accepted models of 100,000 simulations (Guo et al., 237 2016, and references therein). 238

239 6.1 Tectonic Interpretation

Figure 7 shows a selection of 1D V_{SV} profiles representative of their parent tectonic domains in West Antarctica: station PIG3 lies adjacent to Pine Island Glacier in the Thurston Island block; station FOWL is close to the Haag Nunataks of the HEW block; node 1624 is in the Ellsworth Mountains of the HEW block; station BREN is at Brenneke Nunatak on the southern Antarctic Peninsula; station SILY is at Mount Sidley in MBL; station BYRD is in the central WARS; station DUFK is at the Dufek Intrusion at the margin of the WSRS; and station SURP is at the southern TAM front (see Figure 1 and Table S1). The average standard deviation of inferred mantle V_{SV} velocities is generally less than ~0.075 km/s, increasing to ~0.1 km/s for locations (e.g., BREN) at the periphery of the modelled domain. The average standard deviation of inferred crustal velocities is generally less than ~0.1 km/s.

The crust thickens from $\sim 25 \,\mathrm{km}$ in the Thurston Island block (PIG3), to $\sim 29 \,\mathrm{km}$ at 251 the Haag Nunataks (FOWL), to $\sim 37 \,\mathrm{km}$ in the Ellsworth Mountains (node 1624). In 252 the southern Antarctic Peninsula (BREN) the crust is $\sim 39 \,\mathrm{km}$; however, this profile 253 is the least well constrained of those displayed due to the peripheral location (see 254 Figures 1 and 5). The crust is $\sim 27 \,\mathrm{km}$ thick in MBL (SILY), $\sim 26 \,\mathrm{km}$ in the central 255 WARS (BYRD), and $\sim 36 \,\mathrm{km}$ thick at the Dufek Intrusion (DUFK). The signature 256 of a sharp crust-mantle transition is absent at the southern TAM front (SURP), so 257 the estimated crustal thickness of $\sim 26 \,\mathrm{km}$ is less well constrained than the other 258 locations. These estimates of crustal thickness are consistent with preceding studies 259 (e.g., Chaput et al., 2014; Ramirez et al., 2017; O'Donnell et al., 2017). 260

All V_{SV} depth profiles show a high-velocity seismic mantle "lid". Defining the seismic 261 lithosphere-asthenosphere boundary (LAB) at the strongest negative velocity gradient 262 at the base of the high-velocity lid (e.g., Eaton et al., 2009), the seismic LAB is at 263 $\sim 85 \,\mathrm{km}$ depth beneath the Thurston Island block (PIG3), MBL (SILY), the central 264 WARS (BYRD) and southern TAM front (SURP). The seismic LAB depth increases 265 to $\sim 92 \text{ km}$ at the Dufek Intrusion (DUFK) and $\sim 96 \text{ km}$ at the Ellsworth Mountains 266 (node 1624) (Figure 7). Alternative definitions of the seismic LAB exist (e.g., Eaton 267 et al., 2009); for example, adopting the onset of the negative velocity gradient at the 268 lid base would reduce our seismic LAB depth estimates by $\sim 10-20$ km. The lid at 269 the southern TAM front (SURP), and at MBL (SILY) to a lesser extent, is underlain 270

by a pronounced low velocity zone: at $\sim 130 \text{ km}$ depth, V_{SV} is ~ 4.05 -4.15 km/s at SIL? SURP and ~ 4.15 -4.20 km/s at SILY. In contrast to SURP and SILY, at BYRD in the central WARS V_{SV} is ~ 4.20 -4.30 km/s at 130 km depth.

2D V_{SV} maps were constructed by gridding the suite of 1D V_{SV} profiles (Figures 274 8 and 9). At 25 km depth, velocities strongly characteristic of crustal lithologies 275 $(V_{SV} < \sim 4.0 \text{ km/s})$ are evident beneath the southern TAM, the WSRS, the HEW 276 block and the Antarctic Peninsula. The slowest velocities at this depth are located 277 beneath the southern TAM and Ellsworth Mountains. However, the ANT resolu-278 tion degrades on the Peninsula (O'Donnell et al., 2018), so the inferred crustal V_{SV} 279 velocities there are likely overestimated; gravity data suggest that crustal thickness 280 on the southern Peninsula is comparable to that beneath the Ellsworth Mountains 281 (e.g., O'Donnell & Nyblade, 2014). Faster velocities - indicative of thinner crust 282 - characterise the WARS at this depth, with velocities indicative of mantle rock 283 $(V_{SV} > \sim 4.3 \text{ km/s})$ apparent in the Ross and Amundsen Sea Embayments. Crust 284 thinner than 25 km at these locations is consistent with preceding studies (e.g., Cha-285 put et al., 2014; Shen et al., 2018). Our model suggests that thicker crust in the WARS 286 is found in a region extending south from the MBL topographic dome, consistent with 287 Chaput et al. (2014). 288

The outstanding feature at 60 km depth is the high velocity anomaly located between 289 the Ellsworth Mountains and the Dufek Intrusion/Pensacola Mountains, also seen 290 in cross-section AA' in Figure 9. Storey & Kyle (1997) posit that plume-generated 291 Ferrar magmas could have ponded in large magma chambers, like that from which 292 the Dufek Intrusion crystallized (see Figure 1 for location), and from these spread 293 along the length of the TAM, explaining the chemical uniformity of Ferrar exposures 294 over large distances. Shear velocities of the magnitude we infer ($\sim 4.6-4.8 \text{ km/s}$) in the 295 lithospheric mantle beneath the southern WSRS are characteristic of depleted, cra-296 tonic lithosphere. We speculate that the high velocity anomaly might reflect depleted 297

mantle lithosphere following the extraction of voluminous melt related to Gondwanabreakup.

The absence of a sharp velocity contrast at the eastern margin of the WSRS is consistent with the WSRS being a broad extensional/transtensional province within a distributed plate boundary between East and West Antarctica (Jordan et al., 2017). The conventional interpretation of the TAM as the margin of East Antarctica in the Weddell Sea Embayment may need to be re-visited.

The seismic signature of the cratonic margin of East Antarctic is clear along the south-305 ern and northern TAM front at depth slices 120 and 150 km. However, the boundary 306 is located behind the southern TAM front. Depth slices at 90, 120 and 150 km reveal 307 a pronounced low velocity anomaly underlying the southern TAM front (minimum 308 V_{SV} is ~4.05 km/s). Shen et al. (2017, 2018) also image this low velocity anomaly 309 and attribute it to lithospheric foundering, a mechanism they invoke to explain the 310 uplift of the TAM. The southern portions of our cross-sections CC' and DD' in Fig-311 ure 9 does not contradict their interpretation. Taken together with a magnetotelluric 312 study advocating flexural support of the central TAM by thick, high electrical resis-313 tivity lithosphere (Wannamaker et al., 2017), and seismic studies advocating flexural 314 support of the northern TAM by warm, buoyant upper mantle impinging from the 315 adjacent WARS (e.g., Lawrence et al., 2006), this points to along-strike variation in 316 the tectonic history of the TAM. 317

We do not interpret structure below 200 km depth, but seismic velocities characteristic of cratonic lithosphere are inferred to persist to depths of \sim 220-250 km beneath East Antarctica (e.g., Ritzwoller et al., 2001; Shen et al., 2018). The thickness of the seismic lid beneath the Ellsworth Mountains (\sim 95-100 km) is substantially less than that underlying the East Antarctic craton (see cross-section AA in Figure 9). This points to modification of the Precambrian lithosphere beneath the Ellsworth Whitmore Mountains, which Lloyd et al. (2015) suggest reflects lithospheric foundering related to Gondwana breakup, magmatic intrusion, and subsequent development ofthe WARS.

At 90 km depth, high velocity zones (V_{SV} \sim 4.5-4.55 km/s) are apparent south of the 327 MBL dome and in the eastern Ross Sea Embayment. White-Gaynor et al. (2019) 328 propose that relatively faster upper mantle V_P velocities imaged beneath the eastern 320 Ross Sea Embayment by body-wave tomography reflect lithosphere that may not 330 have been reheated by the Cenozoic rifting that affected other parts of the WARS. 331 Xenolith analyses suggest that lithospheric mantle beneath MBL and circum-Pacific 332 Phanerozoic continental crustal terranes in south east Australia and other locations in 333 Zealandia preserves ancient Archean-Proterozoic peridotite components (e.g., Handler 334 et al., 2003; Liu et al., 2015, and references therein). Handler et al. (2003) suggest 335 that the Proterozoic mantle beneath MBL might have a provenance in the East 336 Antarctic craton, while Liu et al. (2015) invoke a model whereby ancient depleted 337 mantle domains are dispersed in the convecting mantle and reappear beneath young 338 continents. As a possible alternative to the White-Gaynor et al. (2019) model, we 339 suggest that the high velocity zones imaged south of the MBL dome and in the eastern 340 Ross Sea Embayment might reflect the compositional signature of ancient continental 341 fragments. 342

Cenozoic alkaline volcanism in MBL, which started at $\sim 28-30$ Ma, was preceded by 343 uplift of the peneplained surface of the MBL block. This, and the isotopic signa-344 ture of a high-U/Pb (HIMU) mantle reservoir in the rocks, suggests plume-related 345 volcanism (e.g., LeMasurier & Landis, 1996, and references therein). Anomalously 346 low seismic velocity upper mantle beneath the MBL dome is consistently imaged, 347 but the unambiguous signature of a plume "tail" extending deeper into the mantle 348 has thus far evaded detection (e.g., Lloyd et al., 2015; Shen et al., 2018). At the 349 northern margin of the WARS, we image a pronounced low velocity anomaly stretch-350 ing eastward from beneath the MBL dome to Pine Island Bay, underlying Thwaites 351

Glacier, but not Pine Island Glacier. The velocity contrast between this perturbed 352 upper mantle and that of the inner WARS ($\sim 5\%$) is consistent with estimates from 353 Lloyd et al. (2015) and Shen et al. (2018). Assuming temperature is the dominant 354 control on lateral variations in seismic velocity in the upper mantle, this contrast 355 translates to a thermal anomaly of $\sim 125-200$ °C (e.g., Faul & Jackson, 2005). Finn 356 et al. (2005) favour a model of subduction-related alkaline magma genesis in MBL. 357 They suggest that protracted Paleozoic-Mesozoic subduction along the Paleo-Pacific 358 margin of Gondwana resulted in metasomatic enrichment of the upper mantle; detach-359 ment of subducted slabs in the late Cretaceous along the former Gondwana margin 360 induced Rayleigh-Taylor instabilities, triggering lateral and vertical flow of warm Pa-361 cific mantle. They suggest that this catalysed melting of the metasomatised upper 362 mantle, resulting in Cenozoic alkaline magmatism. Emry et al. (2014) also suggest 363 that subduction-related volatiles might explain negative peaks in receiver functions 364 above the mantle transition zone in West Antarctica. The velocity anomaly we image 365 strikes approximately parallel to the convergent paleo-Pacific margin of Gondwana, 366 so it conceivably encodes the signature of subduction-related melt and volatiles rather 367 than, or in addition to, plume-related anomalously elevated temperatures. Additional 368 data (e.g., compressional wave velocities, resistivity measurements) are needed to dif-369 ferentiate between chemical and thermal contributions to the observed low shear wave 370 velocity anomaly, and hence between subduction and plume hypotheses. A less pro-371 nounced low velocity zone underlying the southern Antarctica Peninsula to $\sim 100 \, \mathrm{km}$ 372 depth may similarly encode the signature of Mesozoic subduction and/or a remnant 373 thermal signature of the mid-Cretaceous Palmer Land orogeny affecting the southern 374 Peninsula (e.g., Vaughan et al., 2002). 375

A low velocity anomaly underlying the Bentley Subglacial Trench in the central WARS is evident at depth slices 90, 120 and 150 km (minimum V_{SV} is ~4.15-4.20 km/s). Lloyd et al. (2015) imaged the same velocity anomaly, arguing that it represents a thermal anomaly associated with focussed Neogene extension. They suggest that
surrounding faster velocities in the WARS may reflect Late Cretaceous/early Cenozoic
extension whose thermal perturbation due to rifting has largely dissipated.

The V_{SV} maps suggests that - the Bentley Subglacial Trench aside - current tectonism in the WARS is concentrated beneath the rift margins. By 180 km depth, lateral variations in velocity across West Antarctica are much reduced, as is the contrast with East Antarctica. The reduced lateral velocity variations within West Antarctica suggest that rift-related tectonism is largely confined to the uppermost mantle (<180 km depth).

388 6.2 Geotherms and Heat Flow

Accurate estimation of geothermal heat flow in West Antarctica is pressing given the 389 considered vulnerability of the WAIS to marine ice sheet instability (e.g., Barletta 390 et al., 2018). We seek the steady-state conductive geotherms, and hence surface 391 heat flows, best explaining inferred V_{SV} profiles at representative tectonic locations 392 in West Antarctica. The selected stations/grid nodes have V_{SV} profiles typical of 393 their parent tectonic domains: the southern Antarctic Peninsula (BREN), the central 394 WARS (BYRD), the Thurston Island block (PIG3, located adjacent to Pine Island 395 Glacier), and the Ellsworth Mountains of the HEW block (grid node 1624) (Figure 396 7). Based on the location of low V_{SV} velocity anomalies in Figure 8, steady-state 397 conduction is probably a reasonable assumption at these locations. Locations for 398 which steady-state conduction is unlikely, for example, in MBL and the southern 399 TAM, are beyond the scope of the present study. A companion study to define 3D 400 variations in mantle viscosity beneath West Antarctica will use the V_{SV} model as a 401 3D gauge of uppermost mantle temperatures. 402

⁴⁰³ We use the Abers & Hacker (2016) MATLAB toolbox to predict the elastic, isotropic

V_S of average spinel peridotite and garnet peridotite compositions of lithospheric mantle for candidate geotherms. The spinel peridotite composition represents average continental lithospheric mantle based on spinel lherzolite xenoliths (McDonough, 1990), and the garnet peridotite composition represents "tecton" (i.e., formed or modified at < 1 Ga) lithospheric mantle based on garnet xenocrysts (Griffin et al., 2009). For fertile peridotites, the transition from spinel peridotite to garnet peridotite occurs at ~1.5 GPa (~45-50 km depth) (e.g., Lee, 2003, and references therein).

For a layer of thickness Δz with constant radiogenic heat production, A, and constant thermal conductivity, k, undergoing 1D steady-state heat conduction, the temperature and heat flow at the bottom of the layer (T_b and q_b, respectively) can be determined from the temperature and heat flow at the top of the layer (T_t and q_t, respectively) using

$$T_b = T_t + \frac{q_t}{k}\Delta z - \frac{A}{2k}\Delta z^2 \tag{1}$$

416 and

$$q_b = q_t - A\Delta z \tag{2}$$

(e.g., Hasterok & Chapman, 2011; Furlong & Chapman, 2013). A 1D steady-state
conductive geotherm is obtained by applying these equations to successive layers
and iterating to account for the temperature and pressure dependence of thermal
conductivity.

⁴²¹ Under steady-state conditions, surface heat flow represents the sum of heat flow into ⁴²² the base of the lithosphere and the integrated radiogenic heat production within the ⁴²³ lithosphere. Direct measurement of radiogenic heat production indicates generally ⁴²⁴ high values in felsic rocks ($\sim 2-3 \,\mu W/m^3$), low values in mafic rocks ($\sim 0.2 \,\mu W/m^3$), ⁴²⁵ and very low values in ultramafic rocks ($\sim 0.02 \,\mu W/m^3$) (e.g., Furlong & Chapman, ⁴²⁶ 2013). We segregate our 1D V_{SV} crustal profiles into upper (felsic) and lower (mafic) ⁴²⁷ portions based on the observed velocities, with each portion comprising a sequence

of 1 km thick layers (i.e., $\Delta z = 1$ km). A global compilation of seismic velocities 428 suggests that middle continental crust is dominated by $V_P = 6.5-6.8$ km/s and V_P/V_S 429 = 1.65-1.80 (Hacker et al., 2015), implying an upper-middle crust transition at V_S = 430 3.61-3.78 km/s. We adopt V_{SV} < 3.7 km/s as indicative of upper crust and V_{SV} > 431 $3.7 \,\mathrm{km/s}$ as indicative of combined middle and lower crust - hereafter referred to 432 as lower crust. To the lower crust we assign a heat production of $0.4 \,\mu W/m^3$ (e.g., 433 Hasterok & Chapman, 2011). We regard $V_{SV} > 4.3 \text{ km/s}$ as defining the transition to 434 the lithospheric mantle, where we fix heat production at $0.02 \,\mu W/m^3$ (e.g., Hasterok 435 & Chapman, 2011; Furlong & Chapman, 2013). Upper crustal heat production, A_{UC}, 436 is assigned according to 437

$$A_{UC} = (1 - F)q_S/D,\tag{3}$$

where D is the thickness of the upper crust (defined by $V_{SV} < 3.7 \text{ km/s}$), q_S is surface 438 heat flow and F is a partition coefficient defining the ratio of "basal" heat flow (the 439 combination of middle/lower crustal heat production, lithospheric mantle heat pro-440 duction, and sub-lithospheric heat flow) to surface heat flow (e.g., Hasterok & Chap-441 man, 2011; Furlong & Chapman, 2013). With observed seismic velocities controlling 442 the definition of upper crustal, lower crustal and lithospheric mantle layers, the par-443 tition model facilitates the convenient parameterisation of steady-state geotherms in 444 terms of a single variable: surface heat flow. Using a preferred partition coefficient 445 of F = 0.74 (Hasterok & Chapman, 2011), we vary q_S in increments of $5 \,\mathrm{mW/m^2}$ to 446 produce candidate steady-state conductive geotherms at locations representative of 447 the southern Antarctic Peninsula (BREN), the central WARS (BYRD), the Thurston 448 Island block in the vicinity of Pine Island Glacier (PIG3), and the Ellsworth Moun-440 tains in the HEW block (grid node 1624). Crustal thermal conductivity is calculated 450 following Furlong & Chapman (2013) and lattice and radiative contributions to ther-451 mal conductivity in the lithospheric mantle calculated following Hasterok & Chapman 452 (2011).453

Attendant elastic, isotropic V_S velocities for the lithospheric mantle are calculated from the geotherms using Abers & Hacker (2016). To facilitate comparison with the observed anelastic, V_{SV} velocities, the calculated velocities are converted to anelastic, V_{SV} velocities assuming PREM Q values and 4% radial anisotropy in the lithospheric mantle of West Antarctica (Ritzwoller et al., 2001). We do not attempt to model the crustal velocity profiles due to the more complex compositional heterogeneity.

Figure 10 shows geotherms best explaining the observed V_{SV} profiles for the Antarctic Peninsula (BREN), the central WARS (BYRD), the Ellsworth Mountains of the HEW block (node 1624), and the Thurston Island block in the vicinity of Pine Island Glacier (PIG3). We present geotherms corresponding to lower-bound, upper-bound and preferred heat flows.

For a tecton garnet peridotite composition, a surface heat flow of ${\sim}60\,\mathrm{mW/m^2}$ at 465 BYRD and PIG3 and $\sim 50 \,\mathrm{mW/m^2}$ at node 1624 yield geotherms that explain the 466 inferred V_{SV} of the lower lithospheric mantle reasonably well. We define the ther-467 mal LAB as the intersection of the conductive geotherm and a mantle adiabat based 468 on a mantle potential temperature of 1300°C and adiabatic temperature gradient of 469 0.45°C/km (e.g., Katsura et al., 2010). While the seismic and thermal LABs need 470 not coincide (e.g., Eaton et al., 2009), they do covary and occur within \sim 5-15 km 471 of each other at these locations for our preferred heat flows. The V_{SV} profile of the 472 upper lithospheric mantle at these three locations is more problematic. At PIG3 and 473 node 1624 in particular, the predicted upper lithospheric mantle V_{SV} is beyond one 474 standard deviation of the observed mean V_{SV} for the garnet peridotite composition. 475 The spinel peridotite composition reduces the predicted V_{SV} somewhat, but a dis-476 crepancy persists. Potential contributors to the discrepancy include (1) inadequate 477 capture of the true velocity structure at the crust-mantle transition, (2) the adoption 478 of constant radial anisotropy of strength 4% in the lithospheric mantle, (3) the use of 479 PREM Q values to convert from elastic to anelastic velocities, (4) the assumed spinel 480

⁴⁸¹ peridotite and garnet peridotite compositions, and (5) the partition model of heat ⁴⁸² production. Surface waves are less sensitive to sharp impedance contrasts than they ⁴⁸³ are to average velocity structure. The addition of receiver function data would better ⁴⁸⁴ constrain velocity structure at the crust-mantle transition (e.g., Shen et al., 2018) ⁴⁸⁵ and mitigate (1). Within the remit of Antarctic seismology, the development of Love ⁴⁸⁶ wave and attenuation tomography models would eliminate the need for assumptions ⁴⁸⁷ (2) and (3), respectively.

Our preferred surface heat flow of $\sim 60 \,\mathrm{mW/m^2}$ at BYRD is largely consistent with 488 inferences based on satellite magnetic data (\sim 55-65 mW/m²; Fox Maule et al., 2005) 480 and seismic data ($\sim 70 \,\mathrm{mW/m^2}$; An et al., 2015a)) at that location, and an inferred 490 broad scale heat flow of $60-70 \,\mathrm{mW/m^2}$ for east-central West Antarctica based on 491 magnetotelluric data (Wannamaker et al., 2017). Our preferred surface heat flow 492 of $\sim 60 \,\mathrm{mW/m^2}$ is similarly broadly consistent with a heat flow of $\sim 60-65 \,\mathrm{mW/m^2}$ 493 inferred by geodynamic modelling of WARS evolution (van Wijk et al., 2008) and 494 a heat flow of $70 \,\mathrm{mW/m^2}$ invoked as representative of Mesozoic-Cenozoic rifts for 495 Antarctic ice sheet modelling (Pollard et al., 2005). A slightly higher heat flow of 496 $\sim 75 \,\mathrm{mW/m^2}$ at BYRD was estimated from a drill core through the ice sheet to 497 bedrock (Gow et al., 1968). These values contrast with inferred heat flows in the 498 central WARS of ${\sim}{>}120\,{\rm mW/m^2}$ based on airborne magnetic data (Martos et al., 499 2017) and $\sim 110 \,\mathrm{mW/m^2}$ based on the extrapolation of global heat flow measurements 500 to Antarctica via seismic structural similarity (Shapiro & Ritzwoller, 2004). 501

⁵⁰² Our preferred heat flow of $\sim 60 \text{ mW/m}^2$ at PIG3 is broadly consistent with infer-⁵⁰³ ences from satellite magnetic data ($\sim 55-65 \text{ mW/m}^2$; Fox Maule et al., 2005), seismic ⁵⁰⁴ data ($\sim 70 \text{ mW/m}^2$; An et al., 2015a), airborne magnetic data ($\sim 60-75 \text{ mW/m}^2$; Mar-⁵⁰⁵ tos et al., 2017), and in situ measurements in continental shelf sediments in the ⁵⁰⁶ Amundsen Sea Embayment (mean $\sim 65 \text{ mW/m}^2$; Dziadek et al., 2019). Our preferred ⁵⁰⁷ $\sim 60 \text{ mW/m}^2$ heat flow at PIG3 again contrasts with the $\sim 110 \text{ mW/m}^2$ modelled by Shapiro & Ritzwoller (2004); however, their modelled standard deviations are of comparable magnitude to their inferred heat flows.

Our preferred heat flow of $\sim 50 \,\mathrm{mW/m^2}$ at node 1624 in the Ellsworth Mountains 510 is lower than estimates based on satellite magnetic data ($\sim 70 \,\mathrm{mW/m^2}$; Fox Maule 511 et al., 2005) and airborne magnetic data ($\sim 65-70 \,\mathrm{mW/m^2}$; Martos et al., 2017), but 512 reasonably consistent with recent seismic-based inferences ($\sim 55 \,\mathrm{mW/m^2}$; An et al., 513 2015a). High heat producing granites in the upper crust are known to occur in the 514 Ellsworth Mountains (e.g., Leat et al., 2018), a factor which might render the partition 515 model of heat production with F = 0.74 inappropriate for modelling the local thermal 516 regime. 517

A surface heat flow of $\sim 60 \,\mathrm{mW/m^2}$ best explains the observed V_{SV} profiles at BREN. 518 The signature of a clear seismic LAB at BREN is lacking, likely a reflection of the 519 degradation in resolution at the model periphery, but $q_S = 60 \text{ mW/m}^2$ gives a thermal 520 LAB of ~ 85 km. Burton-Johnson et al. (2017) used geological analyses to infer a mean 521 heat flow of $81 \,\mathrm{mW/m^2}$ on the east and south of the Antarctic Peninsula where silicic 522 rocks predominate, and a mean of $67 \,\mathrm{mW/m^2}$ on the west and north where volcanic 523 arc and quartzose sediments dominate. BREN is located approximately on the border 524 between these domains, where the heat flow inferred by Burton-Johnson et al. (2017) 525 is $\sim 60-80 \,\mathrm{mW/m^2}$. Martos et al. (2017) broadly replicate the spatial variation in heat 526 flow on the Peninsula, but their inferred values are consistently higher than those of 527 Burton-Johnson et al. (2017). 528

We emphasise that inferred heat flows are regional average (many hundreds of km) estimates constrained by seismic data with limited sensitivity to the upper crust in conjunction with radiogenic heat productions for felsic, mafic and ultramafic lithologies taken from global compilations (e.g., Hasterok & Chapman, 2011; Furlong & Chapman, 2013). This precludes meaningful comparison with geographically localised high heat flow anomalies (e.g., Fisher et al., 2015), but does not contradict such mea⁵³⁵ surements. Our inferred geotherms and heat flows can serve as regional average
⁵³⁶ benchmarks which can be modified according to local conditions.

537 7 Conclusions

In this work, we combined data from the UKANET, POLENET-ANET, ASAIN, SEPA and GSN seismic arrays to construct from fundamental mode Rayleigh wave phase velocities a 3D shear wave velocity model of the West Antarctic upper mantle to 200 km depth. Our goals were (i) image and interpret structures against the tectonic evolution of West Antarctica, and (ii) extract information from the seismic model that can serve as boundary conditions in ice sheet and GIA modelling efforts.

We speculate that a high velocity anomaly located in the southern WSRS might reflect 544 depleted mantle lithosphere following the extraction of voluminous melt related to 545 Gondwana fragmentation. High velocity anomalies imaged by body-wave tomography 546 in the upper mantle beneath the eastern Ross Sea Embayment have been interpreted 547 as lithosphere that may not have been reheated by the Cenozoic rifting that affected 548 other parts of the WARS (White-Gaynor et al., 2019). Motivated by xenolith analyses, 549 as an alternative model we propose that high velocity zones imaged south of the 550 MBL dome and in the eastern Ross Sea Embayment in this study might reflect the 551 compositional signature of ancient continental fragments. 552

⁵⁵³ While the seismic signature of the cratonic margin of East Antarctic is clear along ⁵⁵⁴ the southern and northern TAM, the absence of a sharp velocity contrast between ⁵⁵⁵ the WSRS and East Antarctica is consistent with the WSRS being a broad exten-⁵⁵⁶ sional/transtensional province within a distributed plate boundary between East and ⁵⁵⁷ West Antarctica (Jordan et al., 2017).

A pronounced low velocity anomaly underlying the southern TAM is consistent with a published lithospheric foundering hypothesis. Taken together with a magnetotelluric study advocating flexural support of the central TAM by thick, stable lithosphere
(Wannamaker et al., 2017), this points to along-strike variation in the tectonic history
of the TAM.

The Bentley Subglacial Trench aside - which may have experienced a pulse of Neogene 563 extension (Lloyd et al., 2015) - the distribution of low velocity anomalies suggests that 564 current tectonism in the WARS is concentrated beneath the rift margins and largely 565 confined to the uppermost mantle ($< 180 \,\mathrm{km}$ depth). On the northern margin of the 566 WARS, a pronounced low velocity anomaly extends eastward from beneath the MBL 567 dome toward Pine Island Bay. If of plume-related thermal origin, the velocity con-568 trast of $\sim 5\%$ between this anomaly and the inner WARS translates to a temperature 569 difference of $\sim 125-200$ °C. However, the strike of the anomaly parallels the paleo-570 Pacific convergent margin of Gondwana, so it conceivably encodes the signature of 571 subduction-related melt and volatiles rather than anomalously elevated temperatures, 572 or a combination thereof. Thermal versus chemical origins will have different impli-573 cations for geothermal heat flow and mantle viscosity modelling efforts to monitor 574 and predict ice sheet evolution. Differentiating between them should be a pressing 575 concern given that the anomaly underlies Thwaites Glacier, a major outlet glacier of 576 the WAIS considered vulnerable to marine ice sheet instability (e.g., Barletta et al., 577 2018). 578

Lithospheric thickness estimates extracted from 1D shear wave velocity profiles rep-579 resentative of tectonic domains in West Antarctica indicate an average lithospheric 580 thickness of $\sim 85 \,\mathrm{km}$ for the WARS, MBL, and Thurston Island block. This in-581 creases to $\sim 96 \text{ km}$ in the Ellsworth Mountains. $\sim 60 \text{ mW/m}^2$ geotherms best explain 582 lithospheric mantle shear wave velocities in the central WARS (BYRD) and adja-583 cent to Pine Island Glacier in the Thurston Island block (PIG3); a $\sim 50 \,\mathrm{mW/m^2}$ 584 geotherm best explains the velocities in the Ellsworth Mountains (node 1624) and 585 a $\sim 60 \,\mathrm{mW/m^2}$ geotherm best explains a less well-constrained velocity profile on the 586

⁵⁸⁷ southern Antarctic Peninsula (1624). We emphasise that inferred heat flows are re-⁵⁸⁸ gional average estimates constrained by seismic data with limited sensitivity to the ⁵⁸⁹ upper crust. They do not preclude geographically-localised elevated heat flows due ⁵⁹⁰ to localised Cenozoic extension or magmatic activity or variations in upper crustal ⁵⁹¹ heat production rooted in compositional variation.

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Figures



Figure 1: (a) Map of West Antarctic BEDMAP2 bedrock topography (Fretwell et al., 2013). Following Dalziel & Elliot (1982), yellow lines delineate the major crustal blocks of West Antarctica that pre-date Gondwana fragmentation (AP, Antarctic Peninsula; TI, Thurston Island; MBL, Marie Byrd Land; HN-EM-WM, Haag Nunataks-Ellsworth Whitmore Mountains Block, hereafter HEW). The approximate locations of Pine Island Glacier (PIG) and Thwaites Glacier (TG) in the Amundsen Sea Embayment are outlined in white. Plate boundaries are marked in red and white crosses show the locations of seamounts. Other abbreviated geographic features: BSB, Byrd Subglacial Basin; BST, Bentley Subglacial Trench; DI, Dufek Intrusion; MBS, Marie Byrd Seamounts; PIB, Pine Island Bay; PM, Pensacola Mountains; TAM, Transantarctic Mountains; WARS, West Antarctic Rift System; WSRS, Weddell Sea Rift System. (b) Map showing the location of the UKANET, POLENET-ANET, ASAIN, SEPA and GSN seismic stations used in this study superimposed on grey-scale bedrock topography. At initial deployment in January-February 2016, five UKANET seismic stations were arranged in a quasi-linear array straddling Pine Island Glacier, two stations were located approximately north of the HEW block, and three stations were deployed along the southern Antarctica Peninsula. At the end of the first year of the deployment the UKANET array was re-configured to bolster coverage along the southern Antarctic Peninsula. The UKANET seismic array was demobilised in January-February 2018. Specific stations and grid nodes (blue star) referred to in the text are labelled. For interpretation of the references to colour in this figure, the reader is referred to the web version of the article.



Figure 2: (a) Vertical-component seismograms from a magnitude 6.0 East Pacific Rise earthquake that occurred on August 18^{th} 2016 (green star in (c)) recorded at seven UKANET seismic stations in West Antarctica (see Table S1). Predicted arrival times of compressional (P) and shear (S) body waves according to the Preliminary Reference Earth Model (PREM; Dziewonski & Anderson, 1981) are marked, after which follows the larger amplitude Rayleigh wave. (b) Rayleigh wave dispersion of the same earthquake at UKANET station PIGD. The raw Rayleigh wave seismogram (top) is filtered into 12×10 mHz wide frequency bands with centre periods ranging from 20 to 143 s. (c) Azimuthal and epicentral distance distribution of the 457 earthquakes used in this study. Tomographic resolution is enhanced by a uniform azimuthal distribution of earthquakes. Concentric circles are at 30° intervals from the south pole. (d) Total number of ray paths used at each period in this study.



Figure 3: Average Rayleigh wave phase velocity dispersion curve for West Antarctica compared with PREM. The 1D average dispersion curve served as a starting model for subsequent 2D tomographic phase velocity inversions.



Figure 4: (Top) Rayleigh wave phase velocity model uncertainty at periods 25, 80 and 125 s. Grid node locations are superimposed on the 25 s map. (Bottom) Rayleigh wave phase velocity model resolution at corresponding periods. For ease of visualization, we present the resolution matrix multiplied by a checkerboard pattern of phase velocity anomalies of wavelength 400 km. 100% represents complete amplitude recovery of positive/negative velocity anomalies. We confine our subsequent discussion of imaged structure to the region enclosed by the white polygon.



Figure 5: Rayleigh wave phase velocity model at a range of periods. Unique scale bars are used at each period to emphasise lateral velocity variations. Blue crosses show the locations of seamounts. Regions of higher uncertainty and lower resolution are masked.



Figure 6: Comparison of Rayleigh wave phase velocity maps at period 25 s inferred by (a) ambient noise tomography (ANT; period range 8-25 s) and (b) two-plane-wave tomography (2PWT; period range 20-143 s). (c) Composite 8-143 s Rayleigh wave phase velocity dispersion curve for UKANET station KEAL obtained by weighted least squares polynomial regression (black curve) of ANT- and 2PWT-curves. The yellow diamond in (a) and (b) shows the location of KEAL.



Figure 7: Vertically-polarised shear wave velocity (V_{SV}) profiles inferred from corresponding Rayleigh wave phase velocity dispersion curves. The thick blue line is the mean V_{SV} velocity, the blue dashed lines are one standard deviation bounds. 0 km depth corresponds to the local elevation of the ice sheet surface at each location. The seismic lithosphere-asthenosphere boundary (LAB) is defined here as the depth of the strongest negative velocity gradient at the base of the high velocity seismic lid.



Figure 8: Shear wave velocity (V_{SV}) maps at a selection of depths. We only interpret shallow (<60 km depth) shear wave structure within the footprint of the ANT model. The ANT model domain is more confined than the 2PWT domain, reflected in the varying areal extent of the maps. Shifting scale bars are used to emphasise lateral velocity variations. The locations of the vertical V_{SV} cross-sections shown in Figure 9 are superimposed on the 90 km depth map.



Figure 9: Vertical shear wave velocity (V_{SV}) cross-sections along the four profiles located in Figure 8. The V_{SV} velocities are contoured at 0.1 km/s intervals. Corresponding BEDMAP2 ice and bedrock topography (Topo) profiles are shown in each case. BST, Bentley Subglacial Trench; EM, Ellsworth Mountains; MBL; Marie Byrd Land; PM, Pensacola Mountains; TAM, Transantarctic Mountains; WARS, West Antarctic Rift System; WC, Walgreen Coast; WM, Whitmore Mountains.



Figure 10: Observed and predicted V_{SV} velocities at seismic stations BREN (southern Antarctic Peninsula), BYRD (central WARS), PIG3 (adjacent to Pine Island Glacier in the Thurston Island block) and node 1624 (Ellsworth Mountains in the HEW block) for spinel peridotite (top) and garnet peridotite lithospheric mantle compositions (middle) corresponding to the steady-state conductive geotherms shown on the bottom. The continuous black V_{SV} profiles represent mean velocities, with dashed and dotted black lines representing one- and two-standard deviation bounds, respectively. Predicted velocity profiles and corresponding geotherms are labelled according to the surface heat flow.